

'GHOSTS' BEYOND QUANTUM MECHANICS

● WHY DO 'GHOSTS' EXIST?

Equilibration and Locality

Marek Gazdzicki, Mark Gorenstein, Ivan Pidhurskyi, Oleh Savchuk, Leonardo Tinti (Jun 2, 2022)

Published in: *Acta Phys.Polon.B* 53 (2022) 8, 2 • e-Print: 2206.01151 [nucl-th]

● ● HUNTING FOR 'GHOSTS' BEYOND QUANTUM MECHANICS

Strong Locality as a Tetrahedron: A Symmetry-Reduced Geometric Representation of the (3,3,2,2) Bell Scenario

Marek Gazdzicki, Francesco Giacosa, Pawel Piesowicz (May 4, 2026)

e-Print: 2605.03104 [quant-ph]

● WHY DO 'GHOSTS' EXIST ?

LEONARDO

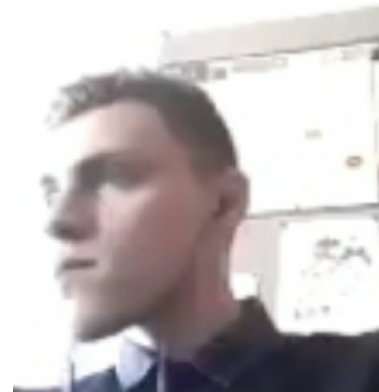


OLEH



MAGDA

MARK



IVAN



MAREK



DEVELOPERS AND PLAYERS
OF THE RANDOM BOARD
GAME,
THE TOY-UNIVERSE,
SUGGESTING THE ANSWER
TO THE QUESTION



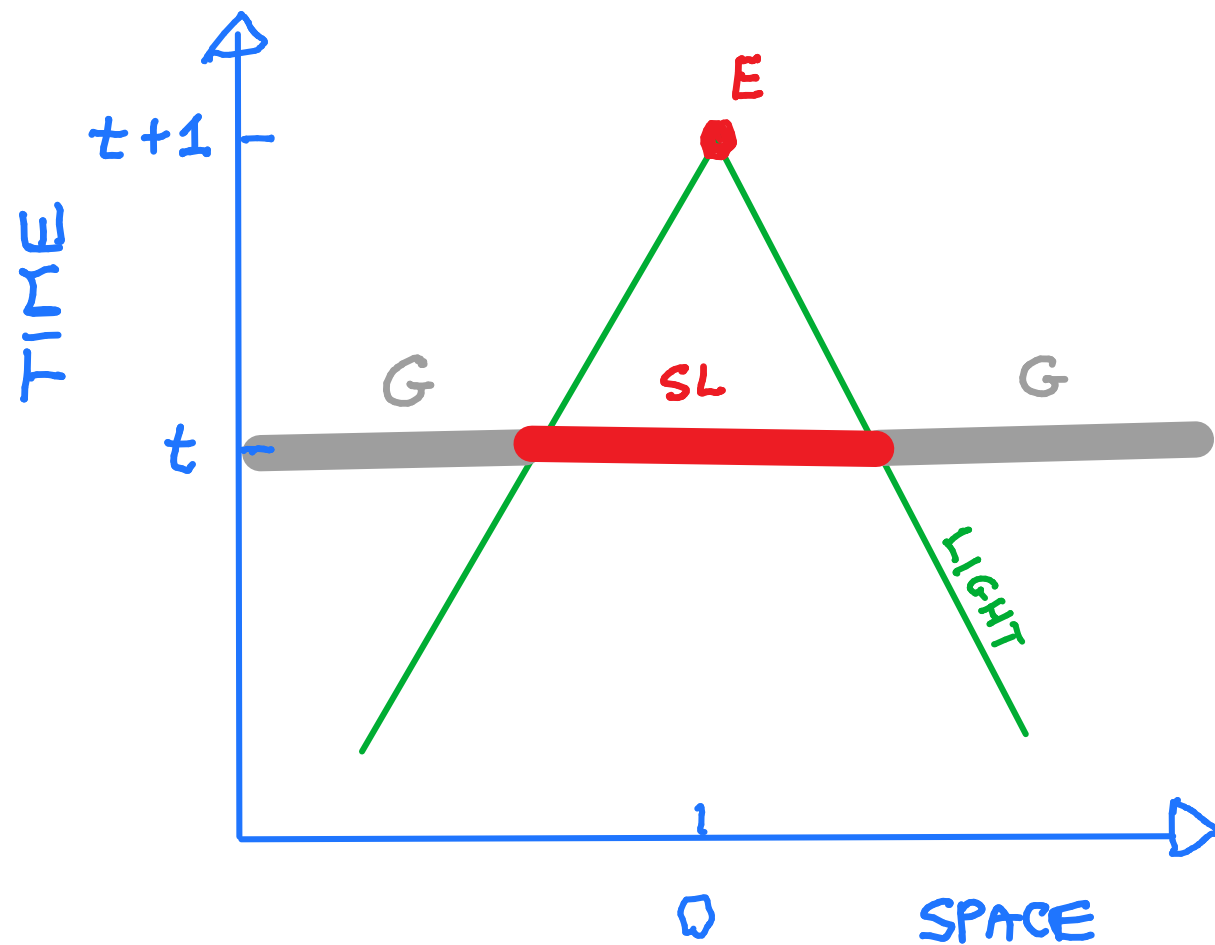
THE COVID TIMES

Equilibration and Locality

Marek Gazdzicki, Mark Gorenstein, Ivan Pidhurskyi, Oleh Savchuk, Leonardo Tinti (Jun 2, 2022)

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'NO-GHOSTS' - DEFINITION:



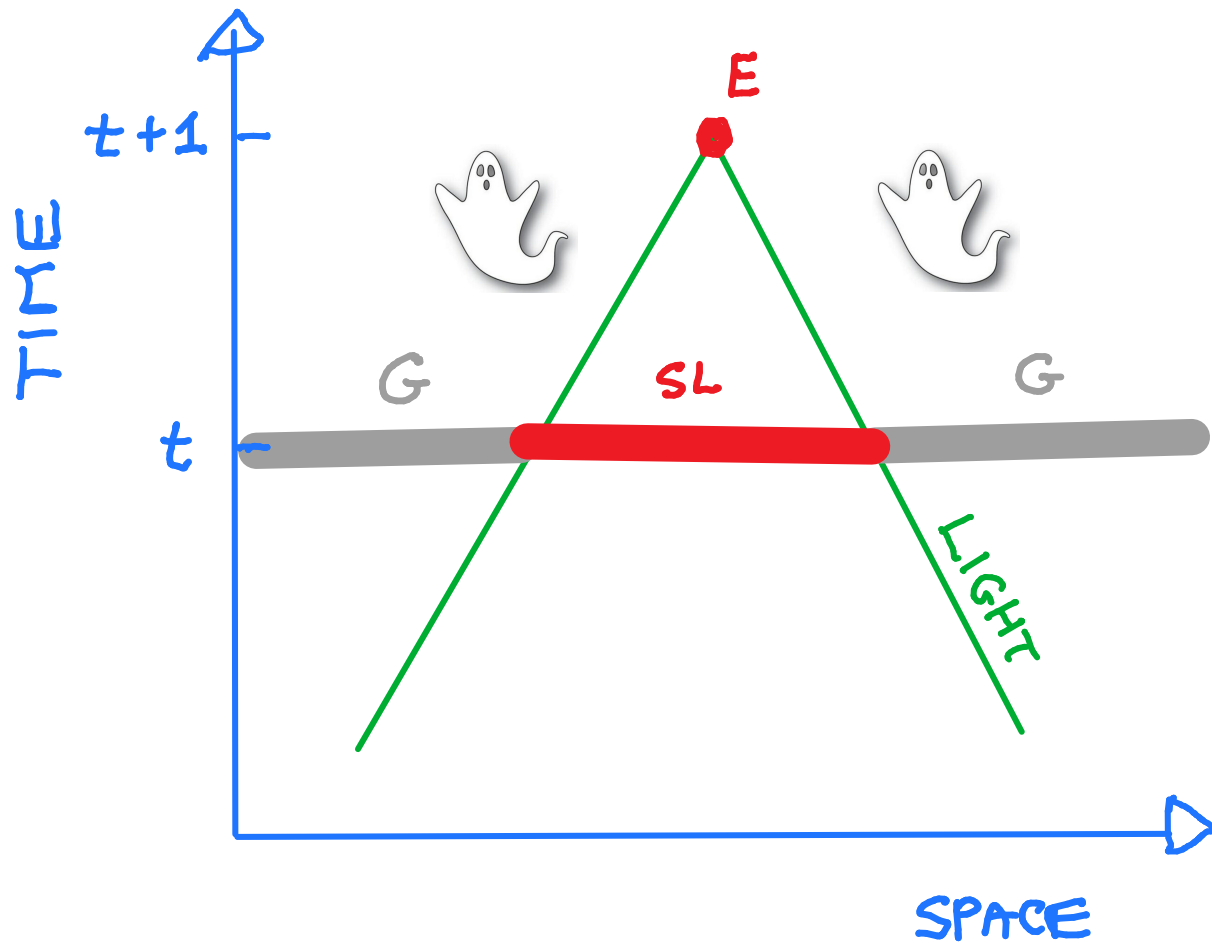
PROBABILITY OF AN EVENT E AT $t+1$,
IN GENERAL, COULD DEPEND ON ALL
EVENTS AT t , SL AND G ,
 $P(E | (SL, G))$.

BUT HAVING SPEED OF LIGHT AS
MAXIMUM TRANSPORT AND SIGNALING
VELOCITY ONE (EINSTEIN ET AL.)
EXPECTS THAT THE DEPENDENCE
ON SL AND G REDUCES TO
THE DEPENDENCE ON SL ,
 $P(E | (SL, G)) = P(E | SL)$.

|||

THE STRONG LOCALITY HYPOTHESIS
'NO GHOSTS' IN NATURE

'GHOSTS' - DEFINITION:



$$\text{IF } P(E | (SL, G)) \neq P(E | SL)$$

- NON-TRIVIAL DEPENDENCE BETWEEN REMOTE EVENTS, 'GHOSTS' IN NATURE (VIOLATION OF STRONG LOCALITY)

HISTORY

≈ 1900: EXPERIMENTS CONTRADICTING CLASSICAL PHYSICS



≈ 1900: QUANTUM MECHANICS



≈ 1950: EINSTEIN-PODOLSKY-ROSEN PARADOX (1935)
BELL'S INEQUALITIES (1964)



≈ 2000: MEASUREMENTS OF 'GHOSTS': ASPECT ET AL. (1982)
CLAUSER ET AL. (1969)
ZEILINGER ET AL. (1998)
...
HANSEN ET AL. (2015)

NOBEL
2022

'GHOSTS' ARE BUILT IN QUANTUM MECHANICS

HERE WE RELATE 'GHOSTS'
TO SYMMETRIES

NEXT SLIDES:

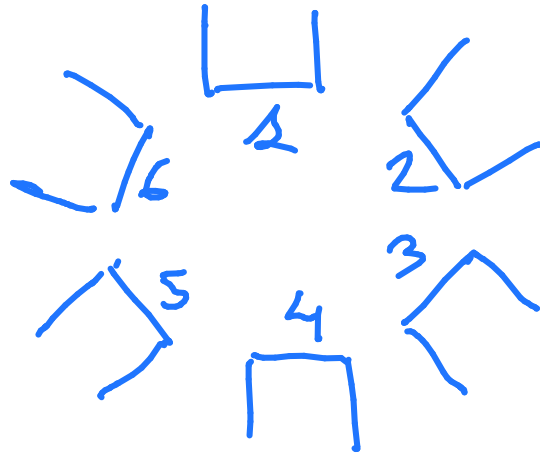
(THE GAME
INSTRUCTION...)

- TOY-UNIVERSE
(THE GAME: BOARD, PAWNS, MOVEMENTS, DICE)
- POPULAR SYMMETRIES
(POPULAR MOVEMENT RULES)
- MICROSTATE-SYMMETRY
(ADVANCED MOVEMENT RULE)
- RESULTS
 - TOY-UNIVERSE + POPULAR SYMMETRIES
 - TOY-UNIVERSE + MICROSTATE SYMMETRY
 - 'GHOSTS' FOLLOW FROM SYMMETRIES

SPACE :

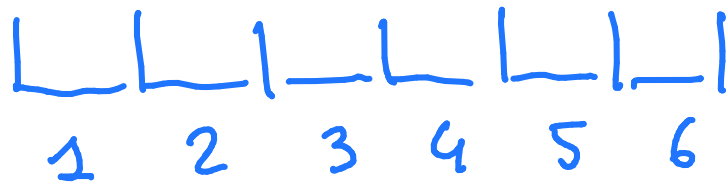
SIX CELLS ARRANGED IN A CIRCLE...

(THE GAME BOARD)



PERIODIC BOUNDARY
CONDITIONS

... WHICH ARE FOR SIMPLICITY DRAWN IN A ROW

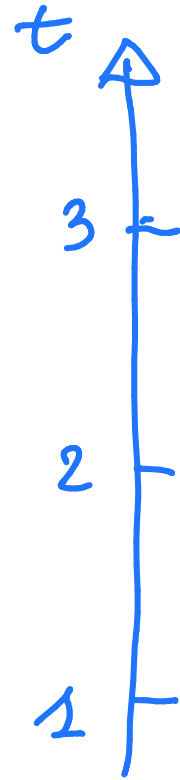


MINIMUM PHASE-SPACE
CELL - $\frac{1}{6}$

TOY-UNIVERSE

TIME:

(GAME MOVES
BY ROLLING
THE DICE)



TIME STEPS.

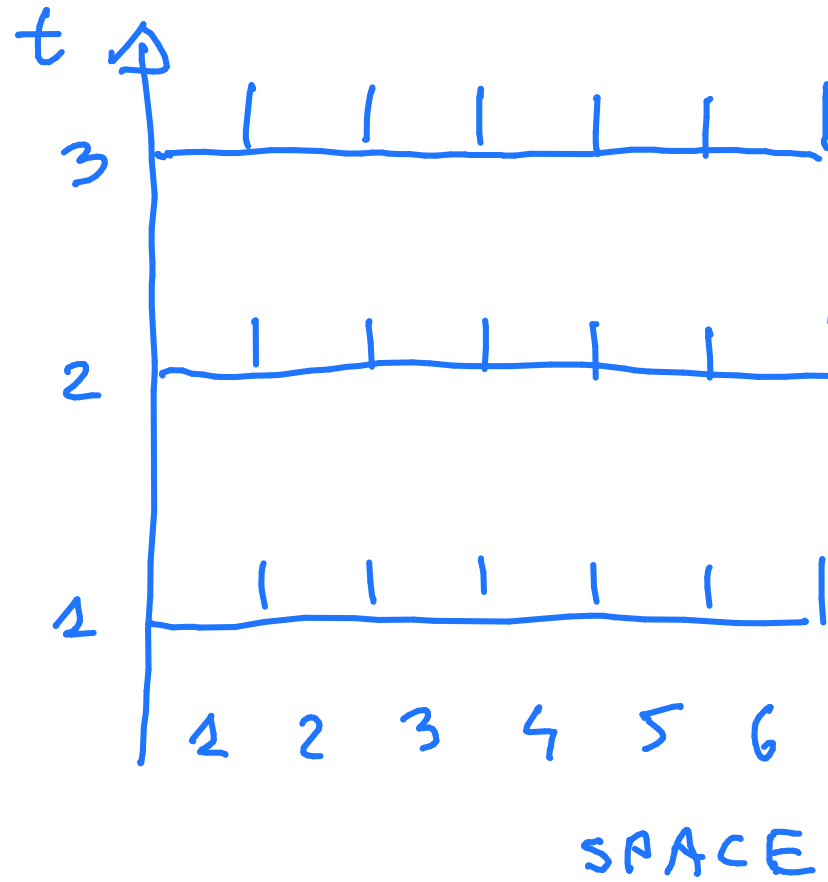
IT IS MEASURED IN STEPS.

THERE IS NO LIMIT TO THE
NUMBER OF STEPS IN TIME

EASY TO INTRODUCE
INDETERMINISM

TOY-UNIVERSE

SPACE - TIME :



TOY - UNIVERSE

PARTICLES:

(PAWNS IN
THE GAME)

● ● ● ● FOUR BALLS AT ALL TIMES

● = ● BALLS ARE IDENTICAL

●
●
●
●
└───┘ NO LIMIT ON NUMBER OF
BALLS IN A CELL

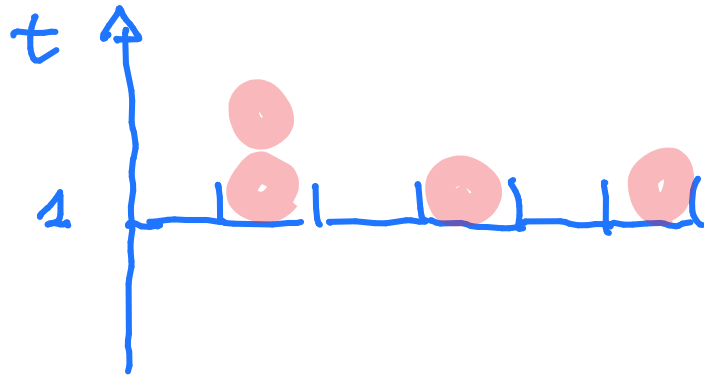
CONSERVATION LAWS

INDISTINGUISHABLE
SPIN-ZERO BOSONS

TOY-UNIVERSE

PARTICLES IN
SPACE-TIME:

MICROSTATE - ANY ARRANGEMENT OF FOUR BALLS
IN SIX CELLS AT A GIVEN TIME
(DENOTED AS x, y, \dots)



TOTAL NUMBER OF MICROSTATES:

$$W_{\text{IND}}(4, 6) = \frac{(4+6-1)!}{4! (6-1)!} = 126$$

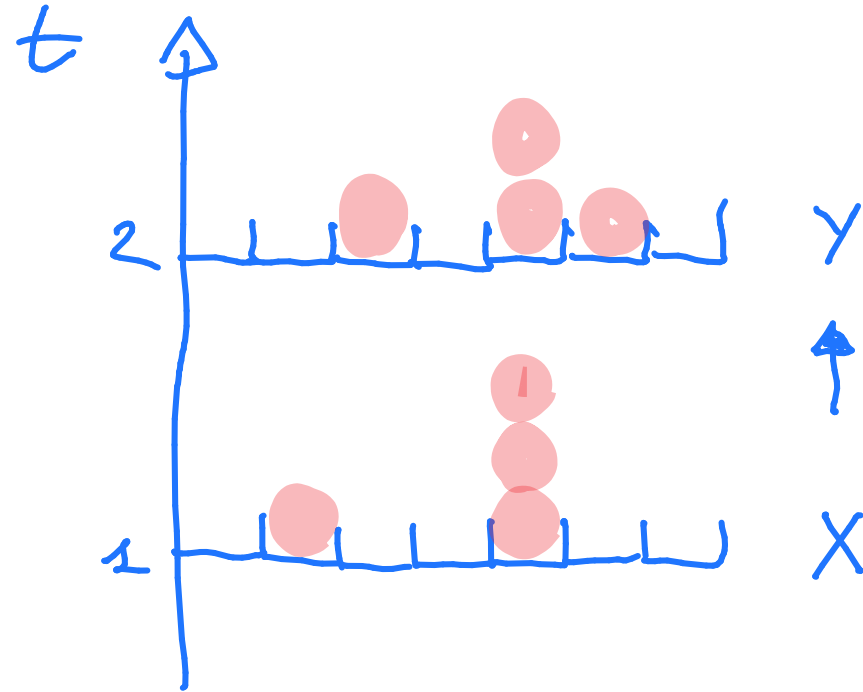
TOY-UNIVERSE

$$\left(W_{\text{DIST}}(4, 6) = 6^4 = 1296 \right)$$

TRANSITIONS:

(MOVES IN THE GAME BY THROWING A DICE)

AT EACH TIME STEP ARRANGEMENT OF BALLS CHANGES — A TRANSITION FROM X TO Y TAKES PLACE WITH TRANSITION PROBABILITY $B(X \rightarrow Y | X)$ DEPENDING ONLY ON X, Y



$$\sum_Y B(X \rightarrow Y | X) = 1$$

TOY-UNIVERSE \rightarrow

MARKOV CHAIN

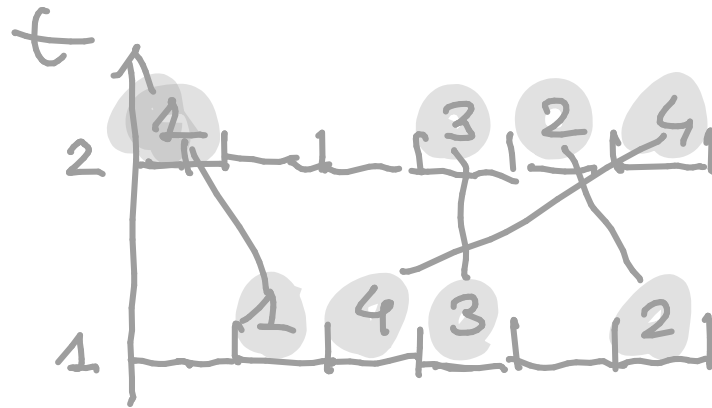
INDETERMINISM OF INTERACTIONS AND MOVEMENTS

TRANSITIONS:

COMMENT ON LABELED AND IDENTICAL BALLS:
(DIST.) (INDIST.)

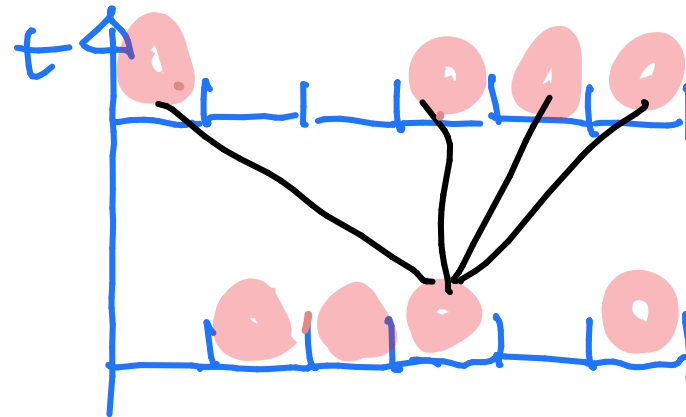
BALL TRAJECTORY \equiv BALL CELL NUMBERS BEFORE AND AFTER A TRANSITION

LABELED BALLS



ALL LABELED BALLS
HAVE TRAJECTORIES

IDENTICAL BALLS



EACH IDENTICAL BALL HAS
FOUR POSSIBLE TRAJECTORIES



NO UNIQUE BALL TRAJECTORIES
FOR IDENTICAL BALLS

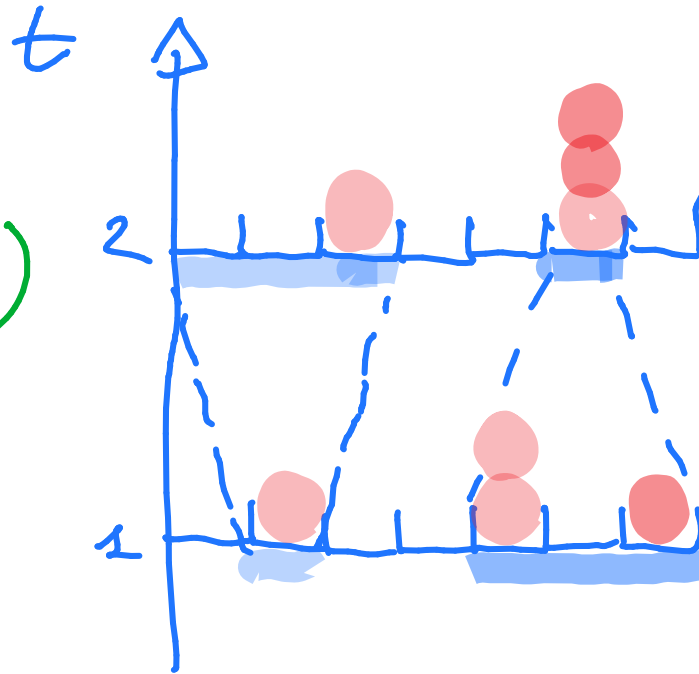
TOY-UNIVERSE

QUANTUM MECHANICS

TRANSPORT LOCALITY:

A TRANSITION CAN TRANSPORT NUMBER OF BALLS IN THE CELL BY NO MORE THAN $\Delta=1$ CELLS.

(SPEED OF LIGHT IN THE GAME)
(FORBIDDEN TELEPORTATION)



LOCALITY CONDITIONS:

$$\sum_{j=i}^{i+k} n_j^X \leq \sum_{l=i-\Delta}^{i+k+\Delta} n_l^Y$$

$$\sum_{l=i}^{i+k} n_l^Y \leq \sum_{j=i-\Delta}^{i+k+\Delta} n_j^X$$

(n - CELL PARTICLE NUMBER)

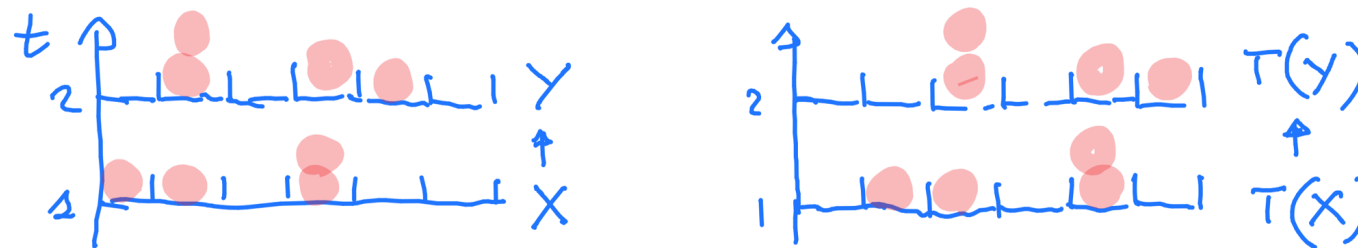
TOY-UNIVERSE

'MAXIMUM VELOCITY' $v \leq c=1$
 FOR IDENTICAL BALLS
 SPECIAL RELATIVITY +
 QUANTUM MECHANICS

POPULAR SYMMETRIES:

- SPACE TRANSLATION
- SPACE REVERSAL
- TIME TRANSLATION
- TIME REVERSAL (KOLMOGOROV CYCLING COND.)
- TRANSITION MATRIX SYMMETRY (+ ERGODICITY \rightarrow
 \rightarrow STEADY STATE = EQUILIBRIUM)

EXAMPLE: SPACE TRANSLATION SYMMETRY:



$$B(x \rightarrow y | x) = B(T(x) \rightarrow T(y) | T(x))$$

RESULT (1)

TOY-UNIVERSE + POPULAR SYMMETRIES



IT IS POSSIBLE TO CONSTRUCT MODELS FOR
INDISTINGUISHABLE PARTICLES (\equiv FIND $B(x \rightarrow y | x)$)
WITHOUT 'GHOSTS'

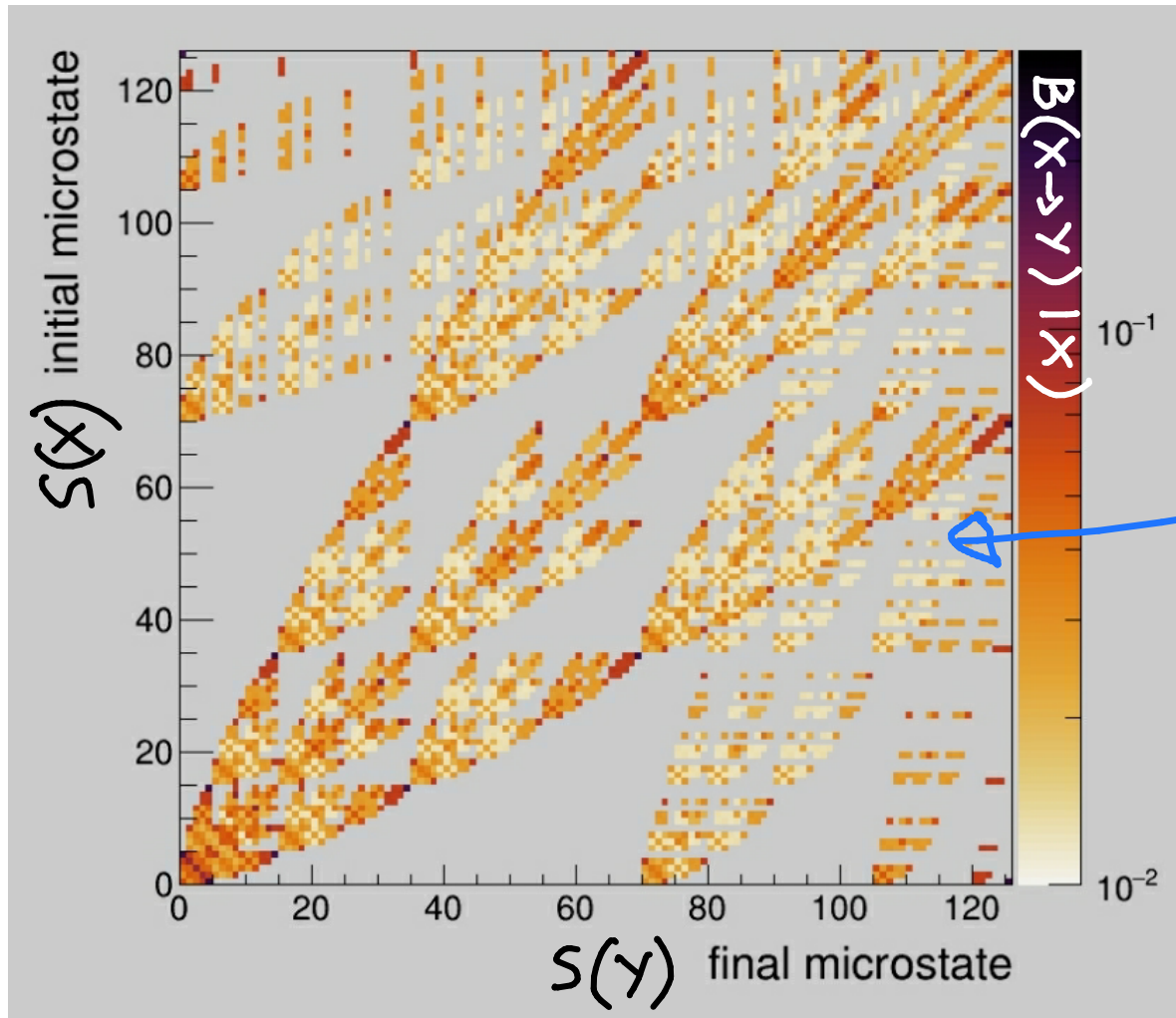
IT IS POSSIBLE, BUT NOT EASY.
IT TOOK US A YEAR OR TWO TO
CONSTRUCT THEM.



RESULTS

EXAMPLE:

TOY-UNIVERSE + POPULAR SYMMETRIES
THE MODELS WITHOUT 'GHOSTS' ARE RATHER COMPLICATED,



MANY
DIFFERENT
COLOURS
(GRAY - FORBIDDEN
TELEPORTATION)



RESULTS

NOTE, MICROSTATE SEQUENTIAL NUMBERS $S(x)$, $S(y)$ UNIQUELY RELATE TO MICROSTATES X AND Y , SEE THE PAPER, BUT THE RELATION IS ARBITRARILY SELECTED

MICROSTATE - SYMMETRY: THE STRONGEST POSSIBLE SYMMETRY

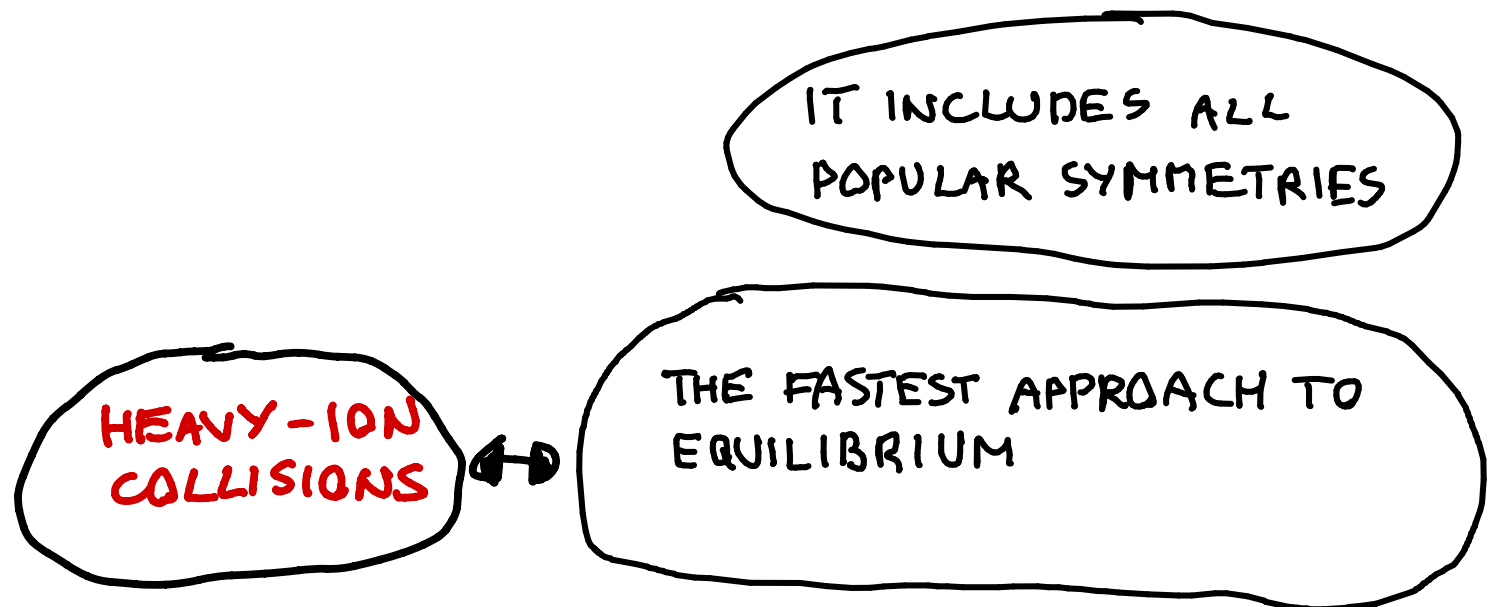
LET M BE ARBITRARY TRANSFORMATION THAT PRESERVES THE SET OF ALLOWED TRANSITIONS,

THEN B IS MICROSTATE - SYMMETRIC IF

$$B(X \rightarrow Y | X) = B(M(X) \rightarrow M(Y) | M(X)) \quad \forall M$$

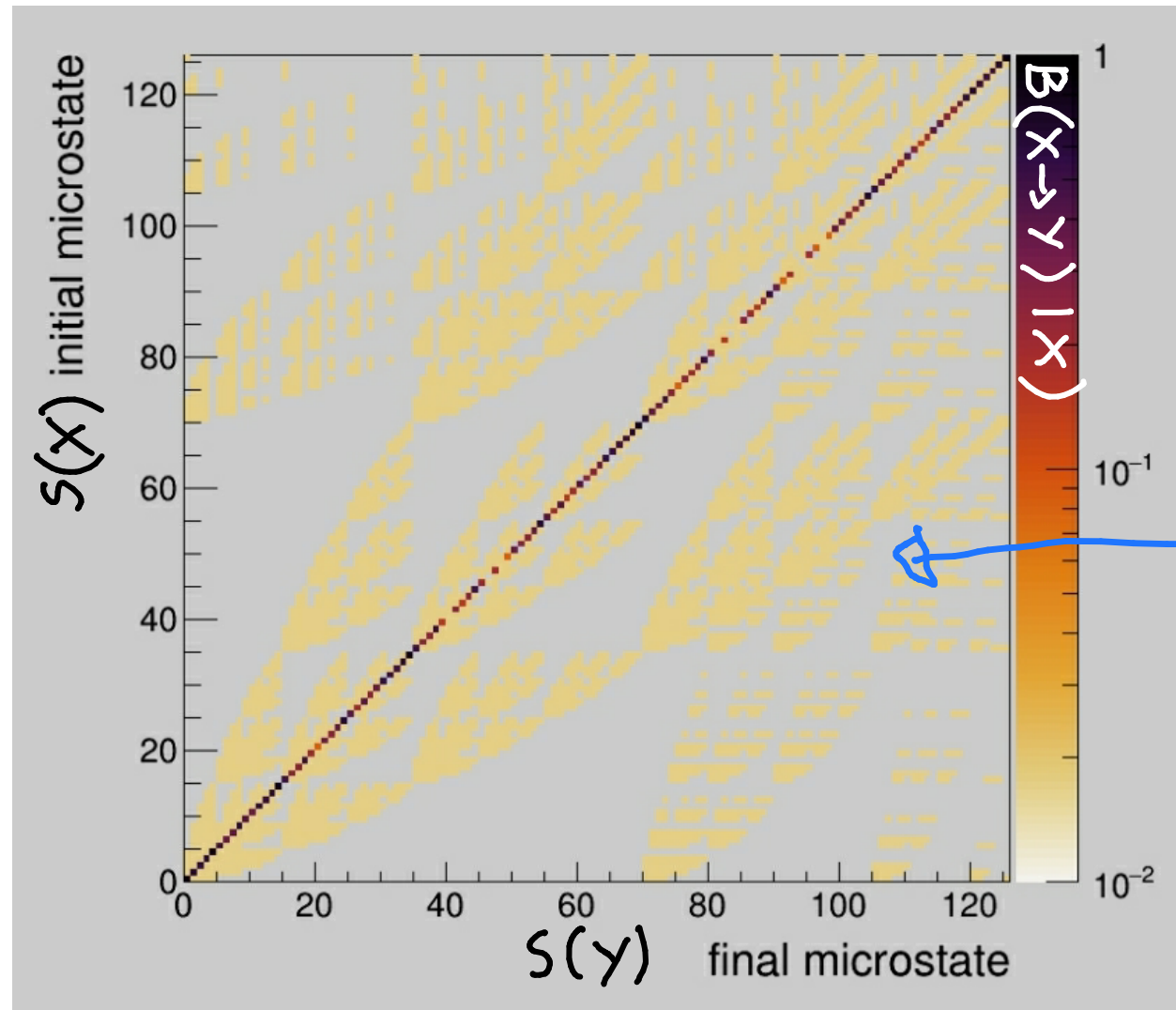
$$\Downarrow$$
$$B(X \rightarrow Y | X) = \text{CONST}(X, Y)$$

MICROSTATE -
SYMMETRY
EQUILIBRIUM



RESULT (II):

TOY UNIVERSE + MICROSTATE-SYMMETRY, BUT DIAGONAL,
ALL MODELS HAVE 'GHOSTS' AND THEY ARE SIMPLE,



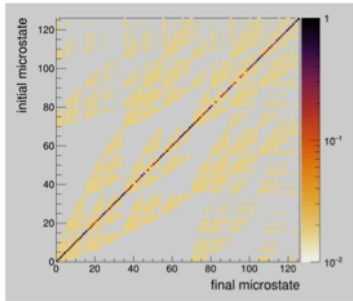
ONLY TWO
NON-DIAGONAL
COLOURS
(GRAY - FORBIDDEN
TELEPORTATION
TRANSITIONS)

HI THERE !



RESULTS

WHY DO 'GHOSTS' EXIST ?



THE SIMPLEST TOY-UNIVERSE
(TELEPORTATION IS FORBIDDEN)
REQUIRES 'GHOSTS'.

HI THERE ♡



HYPOTHESIS:

'GHOSTS' EXIST BECAUSE OF SYMMETRIES.

THEY ARE NOT SPECIFIC TO QUANTUM MECHANICS



HUNTING FOR 'GHOSTS' BEYOND QUANTUM MECHANICS

HUNTING FOR 'GHOSTS' BEYOND QUANTUM MECHANICS

Strong Locality as a Tetrahedron: A Symmetry-Reduced Geometric Representation of the (3,3,2,2) Bell Scenario

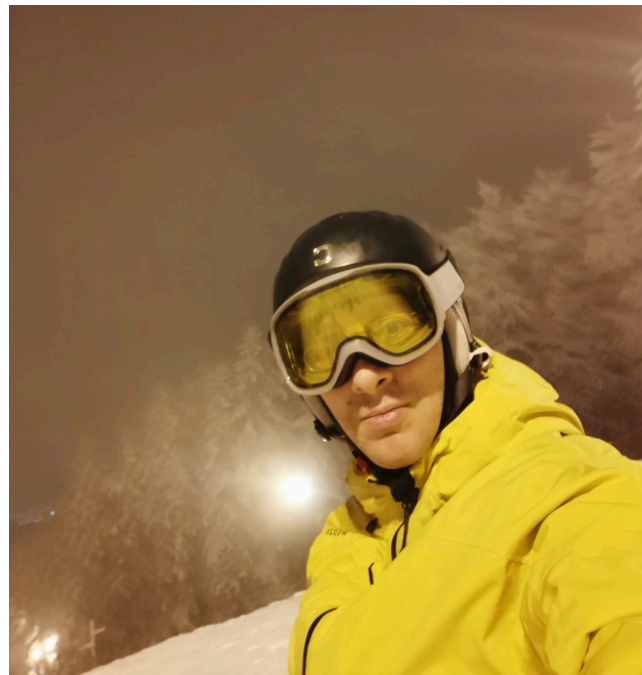
Marek Gazdzicki, Francesco Giacosa, Pawel Piesowicz (May 4, 2026)

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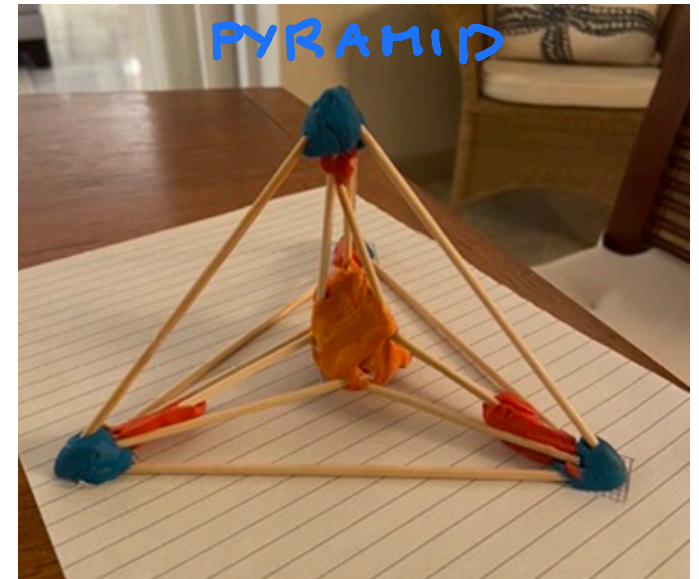
HAUNTERS:



PAWEŁ



FRANCESCO

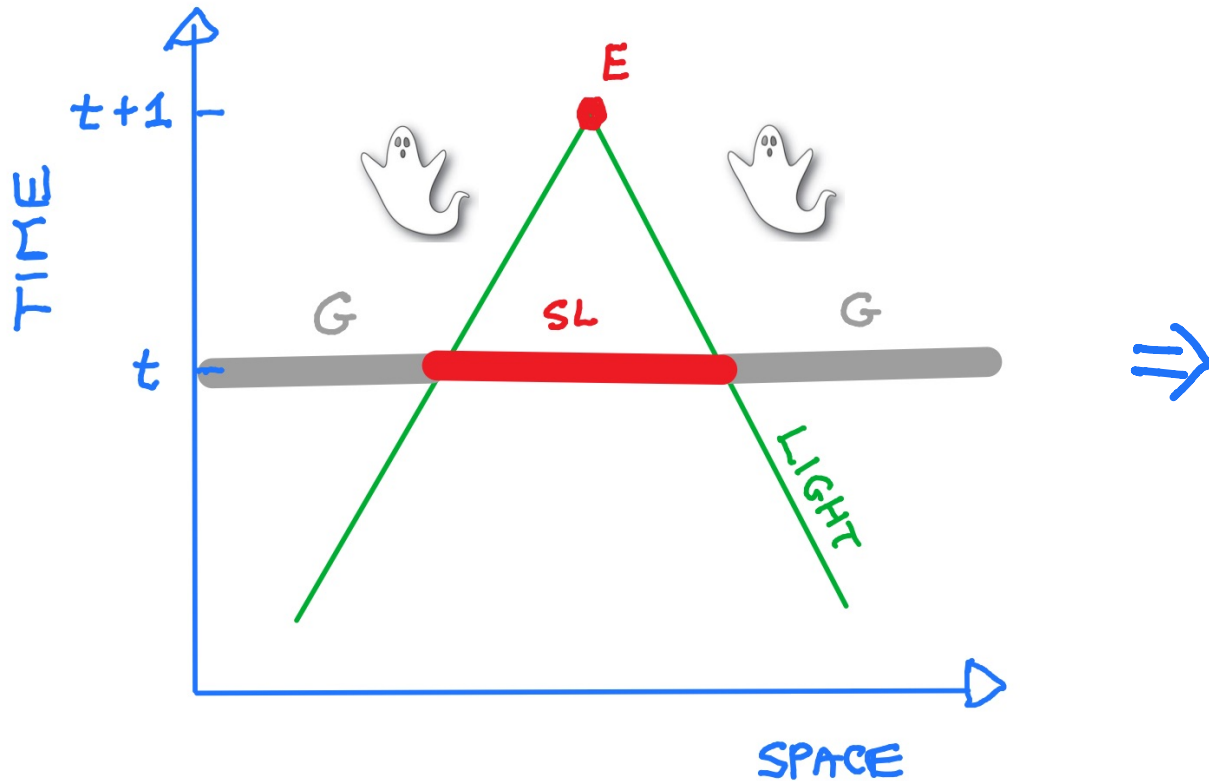


MAREK

THE POST-COVID WINTER VACATIONS

STRONG LOCALITY

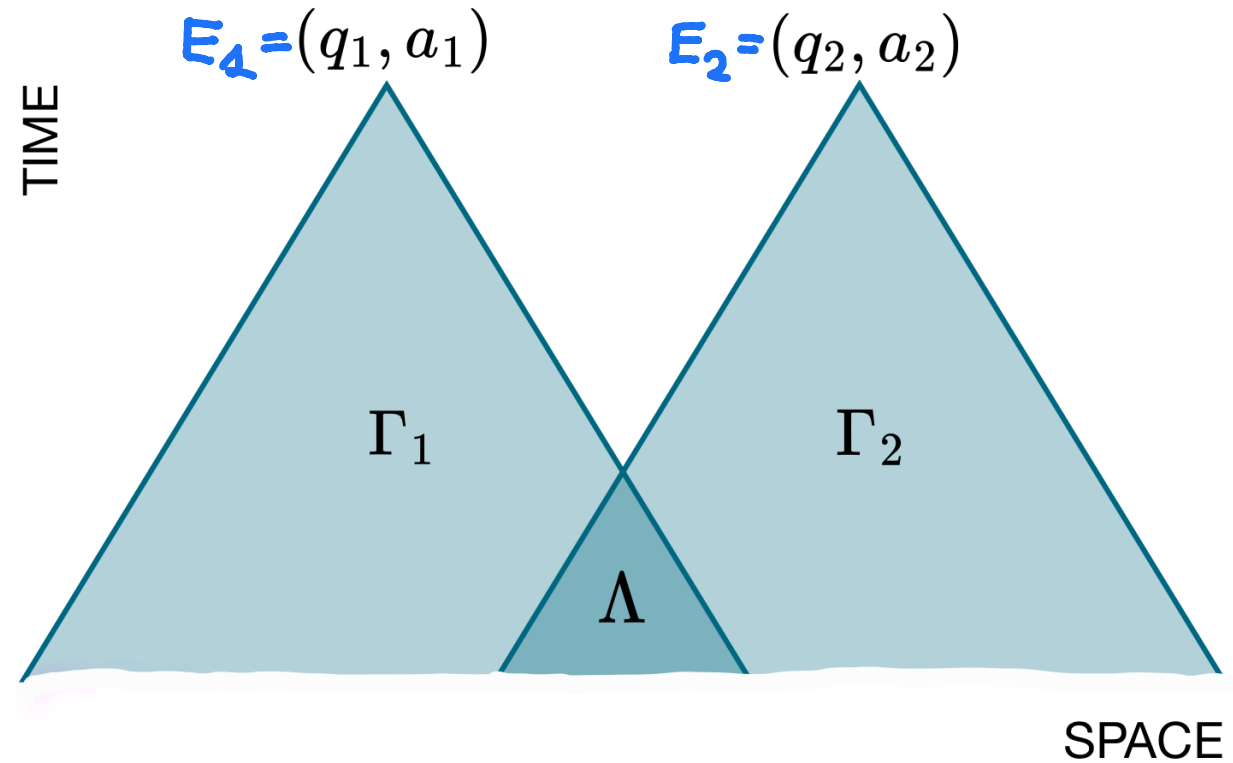
SCREENING-OFF CRITERION



$$P(E | SL, G) = P(E | SL)$$

~ IMPRACTICAL IN
THE REAL UNIVERSE

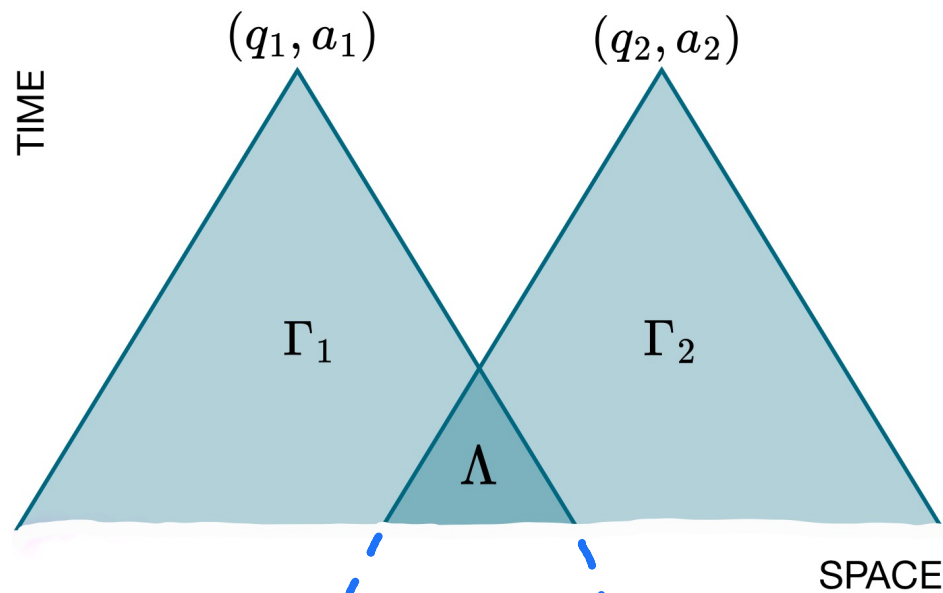
BELL'S PRACTICAL CRITERION



BELL INEQUALITIES WITH E_1, E_2

PRACTICAL, BUT WITH
SUPERDETERMINISTIC LOOPHOLE
(SEE BELOW)

HIDDEN VARIABLE



λ - HIDDEN VARIABLE WHOSE VALUE IS SET IN THE COMMON PAST Λ OF E_1 AND E_2 .

FREE-WILL POSTULATE:



$$h(q_1, q_2 | \lambda) = h(q_1, q_2)$$



$$f(\lambda | q_1, q_2) = f(\lambda)$$

MAY NOT BE TRUE:
SUPERDETERMINISTIC LOOPHOLE

$$P(a_1, a_2 | q_1, q_2)$$

IN SL MODELS

BELL'S DEFINITION OF STRONGLY LOCAL MODELS
IN TERMS OF HIDDEN DISTRIBUTIONS:

$$P(a_1, a_2 | q_1, q_2, \lambda) = P(a_1 | q_1, \lambda) \cdot P(a_2 | q_2, \lambda)$$

(INDISTINGUISHABLE SITES)

↓ AVERAGING OVER HIDDEN VARIABLE λ

THE OBSERVED DISTRIBUTION:

$$P(a_1, a_2 | q_1, q_2) = \sum_{\lambda} P(a_1 | q_1, \lambda) P(a_2 | q_2, \lambda) f(\lambda)$$

CHSH MEASURE S

CLAUSER

HORNE

SHIMONY

HOLT

THE (2,2,2,2) BELL'S CASE:

$$q_1, q_2 \in \{0, 1\}$$

- MEASUREMENT SETTINGS (QUESTIONS)

$$a_1, a_2 \in \{-1, 1\}$$

- MEASUREMENT OUTCOMES (ANSWERS)

OBSERVED DISTRIBUTION

$P(a_1, a_2 | q_1, q_2)$ AND ITS MOMENTS

FOR EXAMPLE $M_{q_1 q_2} \equiv \langle a_1 a_2 \rangle_{q_1 q_2}$

$$S \equiv M_{00} + M_{01} + M_{10} - M_{11}$$

ALGEBRAIC LIMIT $|S| \leq 4$

IT IS 'EASY' TO SHOW THAT $|S| \leq 2$
IN STRONGLY LOCAL MODELS

J. F. Clauser, M. A. Horne, A. Shimony, and R. A. Holt. Proposed experiment to test local hidden-variable theories. *Physical Review Letters*, 23:880–884, 1969.

S IN QUANTUM AND NON-SIGNALING MODELS

QUANTUM MODELS:

$$|S| \leq \sqrt{2} \cdot 2 \quad \leftarrow \text{TSIRELSON BOUND}$$

B. S. Tsirelson. Quantum generalizations of Bell's inequality. *Letters in Mathematical Physics*, 4:93–100, 1980.

NON-SIGNALING MODELS:

(MARGINAL DISTRIBUTION AT ONE LOCATION DOES NOT
DEPEND ON SETTING AT OTHER LOCATION)

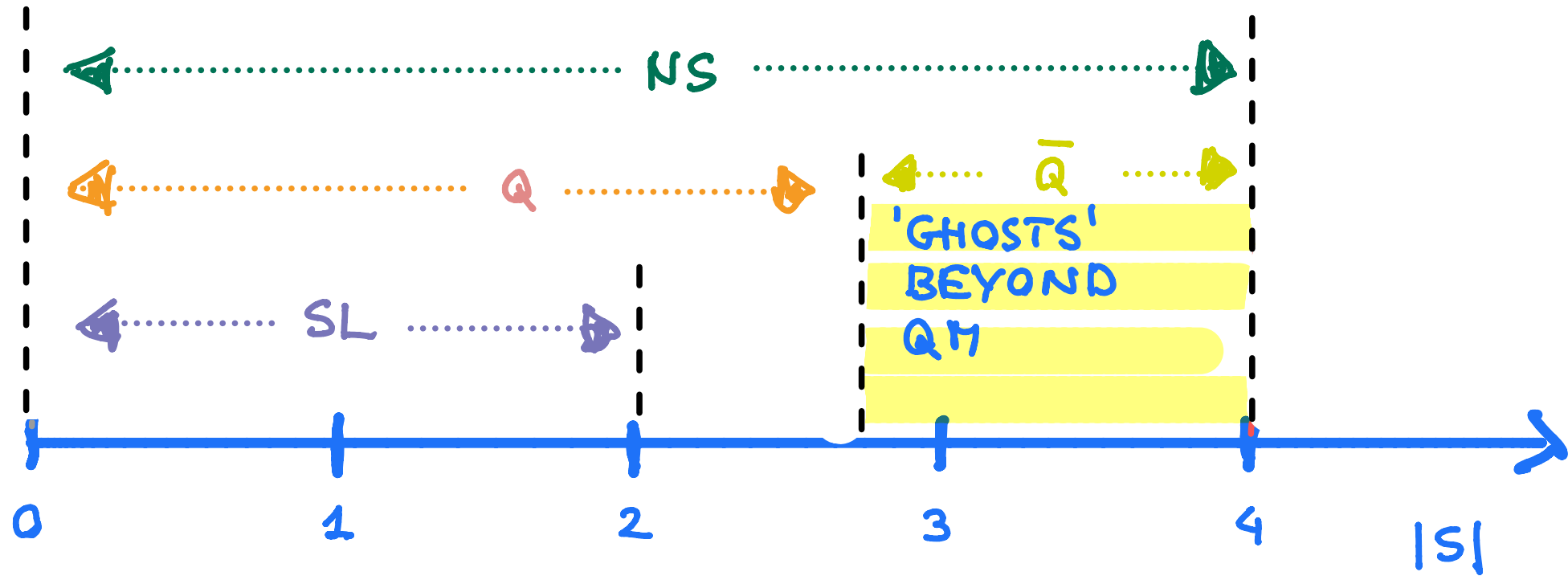
$$|S| \leq 4 \quad \leftarrow \text{ALGEBRAIC MAXIMUM}$$

S. Popescu and D. Rohrlich. Quantum nonlocality as an axiom. *Found. Phys.*, 24(3):379–385, 1994.

THE GENERAL CONDITION FOR INDISTINGUISHABLE SITES:

$$P(a_1, a_2 | q_1, q_2) = P(a_2, a_1 | q_2, q_1)$$

S IN DIFFERENT MODELS



SL - STRONGLY LOCAL

Q - QUANTUM

NS - NON-SIGNALING
($P(q_2|q_1)$ INDEPENDENT OF q_2, \dots)

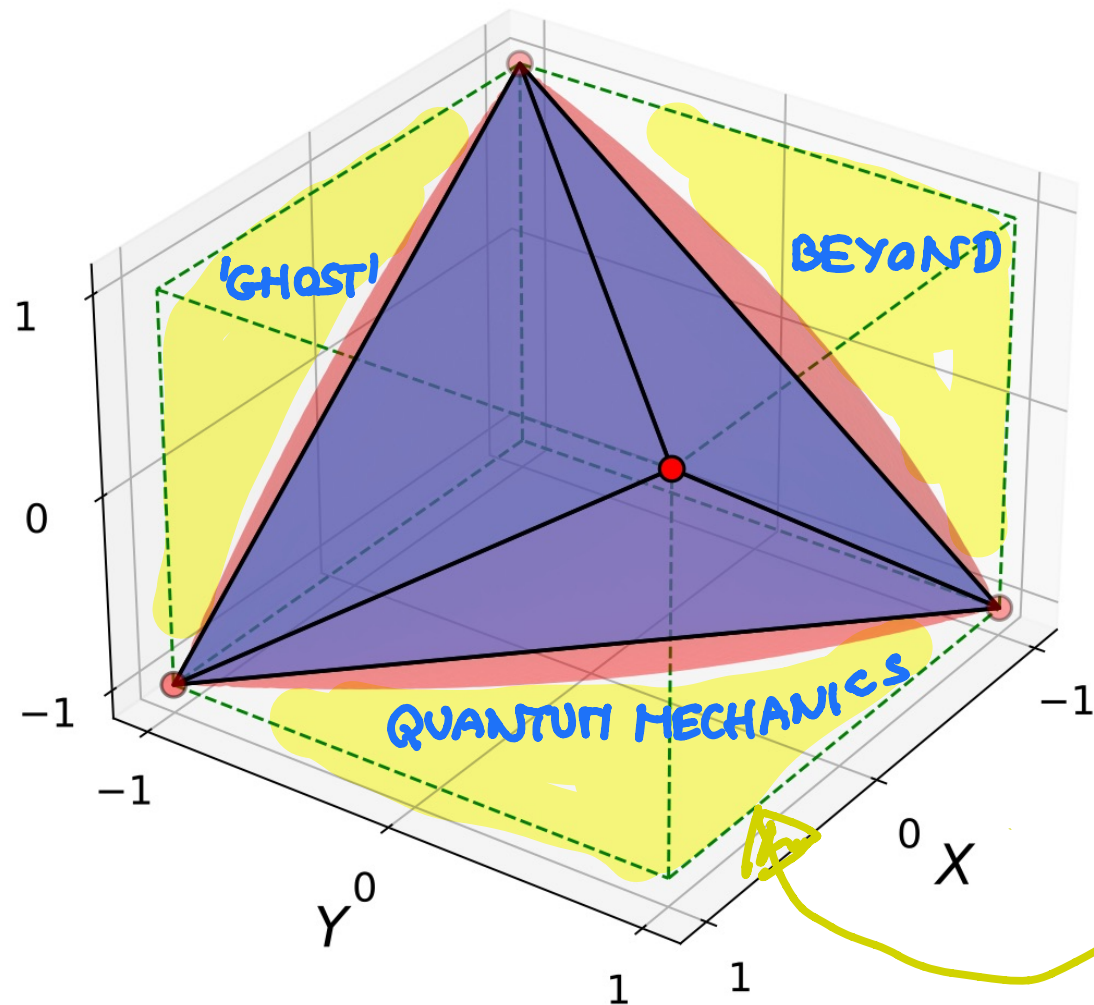
\bar{Q} - POST-QUANTUM

PYRAMID METHOD

THE (3,3,2,2) BELL'S CASE (THREE QUESTIONS, TWO ANSWERS)

ONLY OFF-DIAGONAL MIXED MOMENTS:

$$X = M_{01}, \quad Y = M_{02}, \quad Z = M_{12}$$



IT IS 'EASY' TO SHOW THAT:

SL - TETRAHEDRON

Q - ELLIPTOPE

$$(\det(G) = 1 + 2XYZ - X^2 - Y^2 - Z^2 \succ 0)$$

NS - NON-SIGNALING CUBE

POST-QUANTUM (\bar{Q})

Strong Locality as a Tetrahedron: A Symmetry-Reduced Geometric Representation of the (3,3,2,2) Bell Scenario

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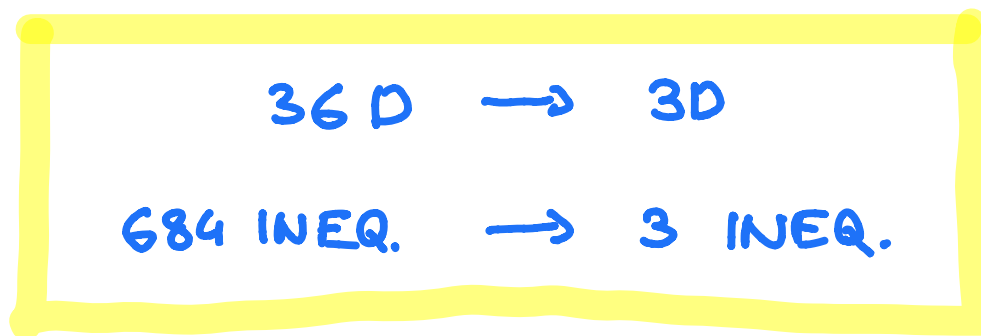
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PYRAMID METHOD

THE (3,3,2,2) BELL'S CASE (THREE QUESTIONS, TWO ANSWERS)

THE FULL POLYTOPE OF STRONGLY LOCAL MODELS LIVES IN 36 D SPACE OF CONDITIONAL PROBABILITIES AND HAS 684 FACET-DEFINING (BELL) INEQUALITIES.

THE PYRAMID METHOD (NORMALISATION, SYMMETRY AND X-Y-Z TRICK) REDUCES THE POLYTOPE TO TETRAHEDRON (3D SPACE) AND NUMBER OF (BELL) INEQUALITIES TO THREE:



CHSH VS PYRAMID

CHSH

PYRAMID

DIMENSIONS

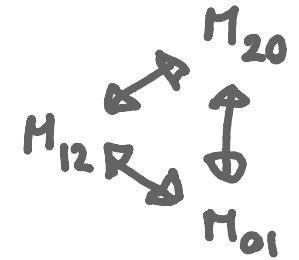
1D

3D

TESTING CORRELATION
TRANSITIVITY

⊖

⊕



DIAGONAL-FREE TESTING

⊖

⊕

SENSITIVITY TO \overline{SL} :

$$1 - \sqrt{V_{SL}/V_{NS}}$$

50% ⊖

66% ⊕

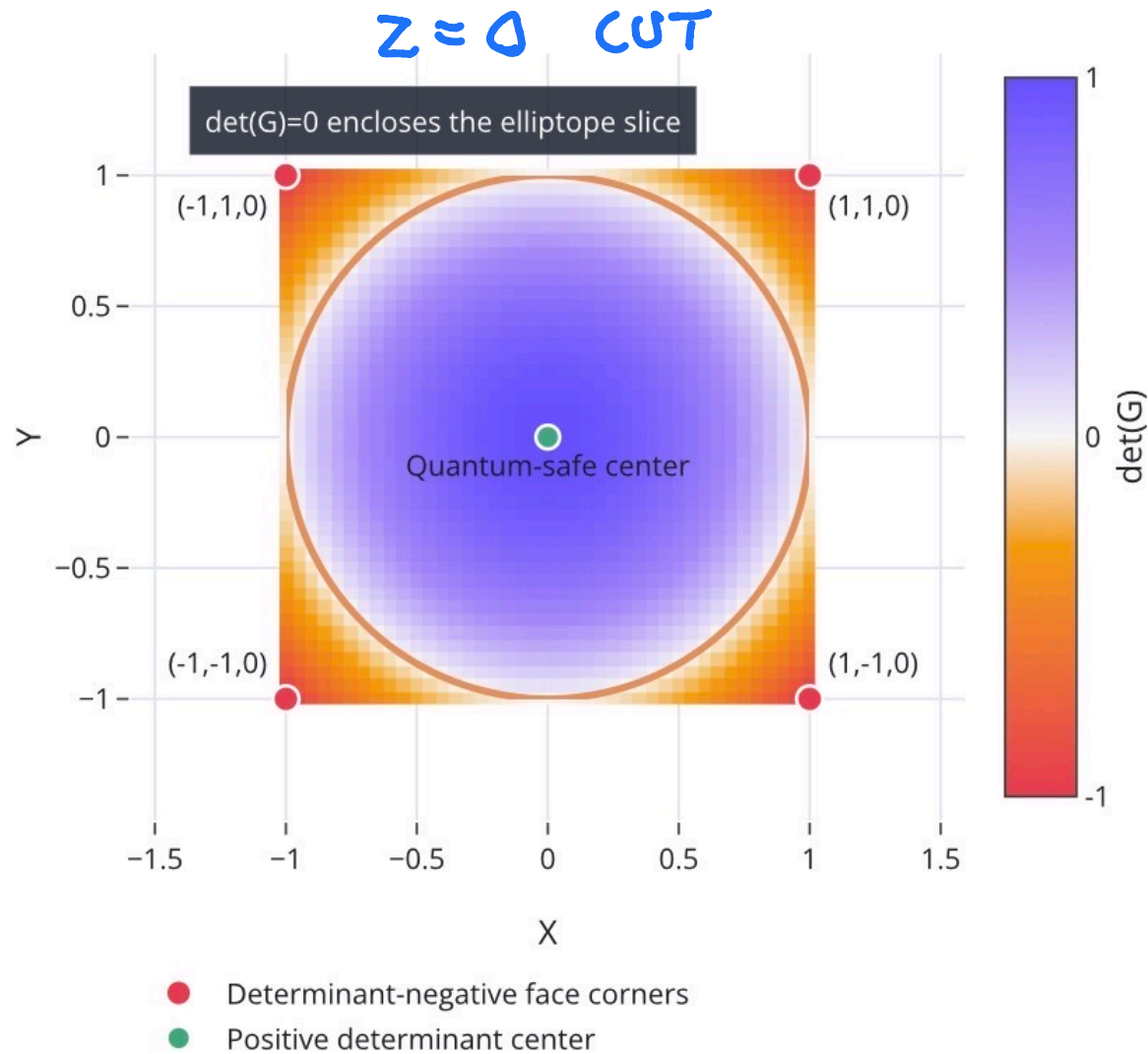
SENSITIVITY TO \overline{Q}

$$1 - \sqrt{V_Q/V_{NS}}$$

30% ⊖

38% ⊕

PYRAMID ADVANTAGES (I) IN POST-QUANTUM SEARCH



THE PYRAMID:

POST-QUANTUM MAXIMALLY VIOLATES
TRANSITIVITY, FOR EXAMPLE:

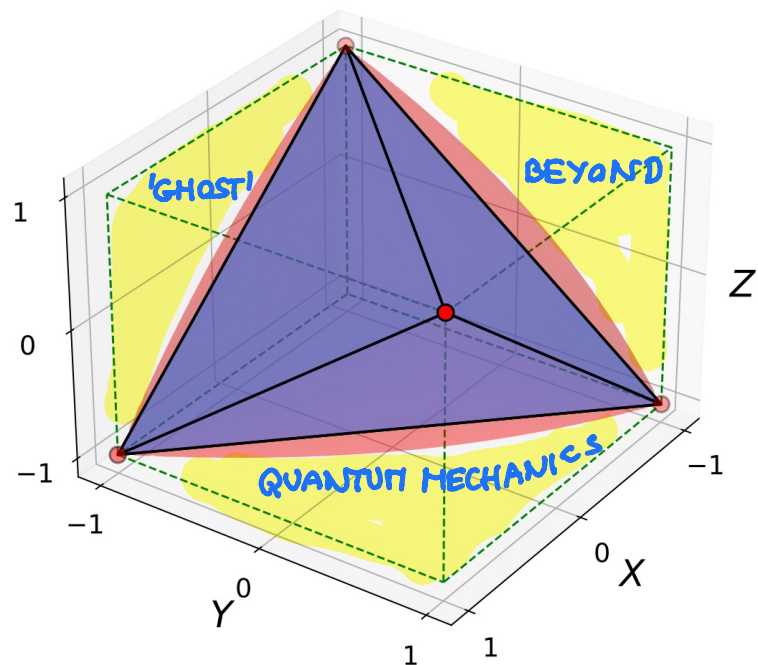
$X = 1$ | FULLY CORRELATED
 $Y = 1$ | OUTCOMES

BUT
 $Z = 0$ | UNCORRELATED
OUTCOMES

NO TRANSITIVITY TESTING
IN THE 1D CHSH CASE

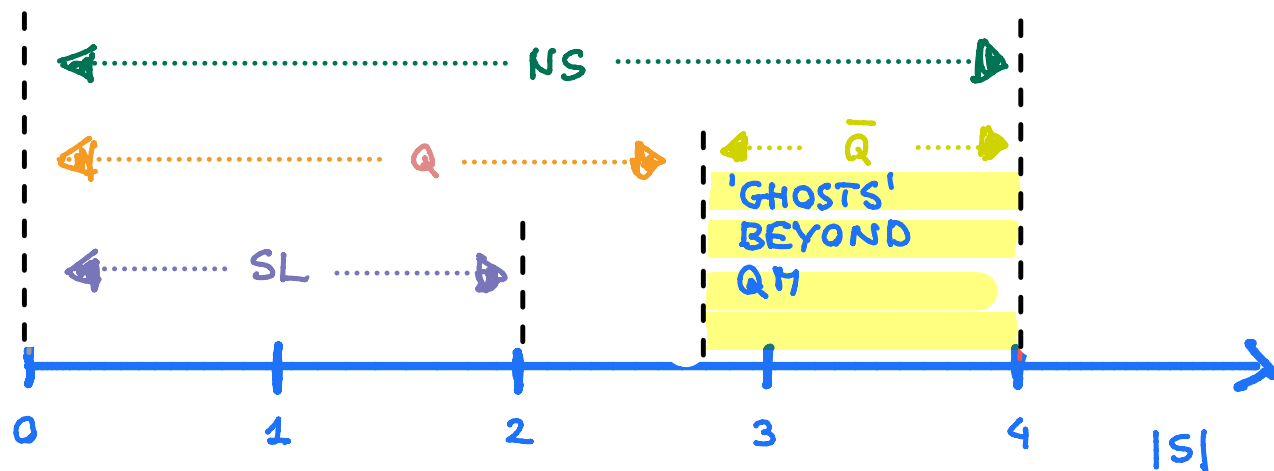
$$(\det(G) = 1 + 2XYZ - X^2 - Y^2 - Z^2 \succ 0)$$

PYRAMID ADVANTAGES (II) IN POST-QUANTUM SEARCH



THE PYRAMID ALLOCATES A LARGE
VOLUME FRACTION TO POST-QUANTUM
 $\approx 38\%$

THIS COMPARES TO $\approx 30\%$ IN
THE CHSH CASE.



SUMMARY

- THE CELL MODEL EXAMPLE SUGGESTS THAT VIOLATION OF STRONG-LOCALITY IN NATURE MAY BE DUE TO SYMMETRY OF THE UNIVERSE → SEARCH FOR 'GHOSTS' BEYOND QUANTUM MECHANICS

Equilibration and Locality

Marek Gazdzicki, Mark Gorenstein, Ivan Pidhurskyi, Oleh Savchuk, Leonardo Tinti (Jun 2, 2022)

Published in: *Acta Phys.Polon.B* 53 (2022) 8, 2 • e-Print: 2206.01151 [nucl-th]

- ● THE PYRAMID METHOD FOR STUDYING VIOLATION OF STRONG-LOCALITY IS WELL SUITED FOR 'GHOSTS' SEARCHING BEYOND QUANTUM MECHANICS.

Strong Locality as a Tetrahedron: A Symmetry-Reduced Geometric Representation of the (3,3,2,2) Bell Scenario

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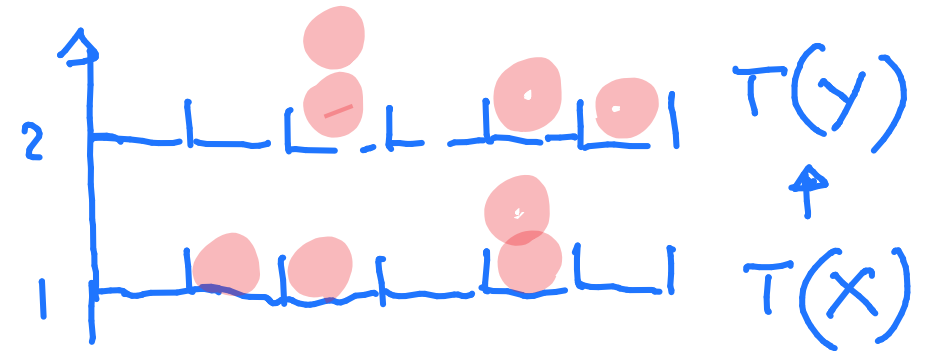
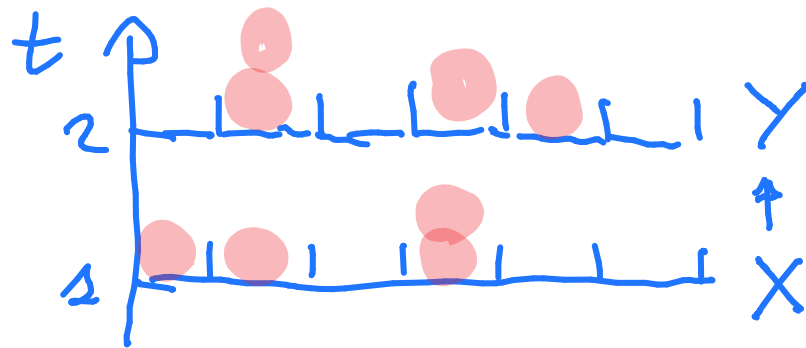
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ADDITIONAL

SLIDES

SPACE-TRANSLATION
SYMMETRY ;

GIVEN A TRANSLATION T OF PARTICLES IN CELLS,
TRANSITION PROBABILITY $X \rightarrow Y$ IS THE SAME
AS PROBABILITY FOR TRANSLATED MICROSTATES
 $T(X) \rightarrow T(Y)$.



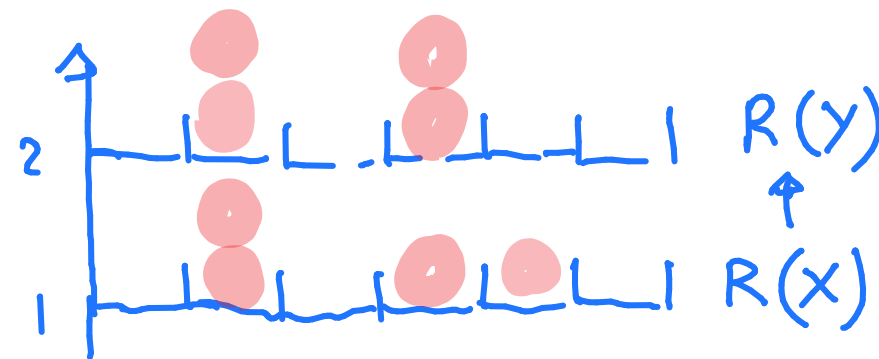
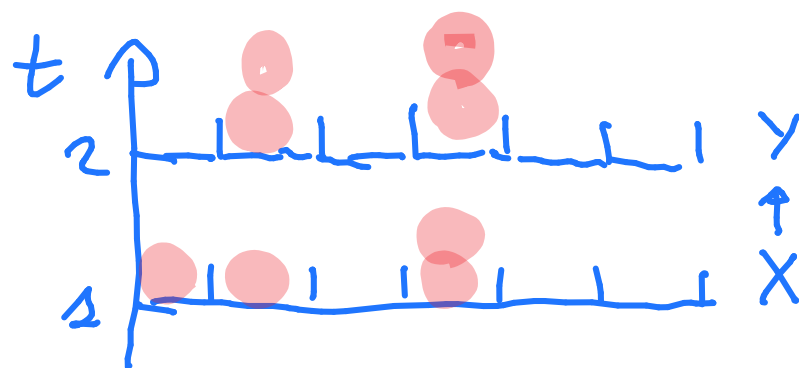
$$B(X \rightarrow Y | X) = B(T(X) \rightarrow T(Y) | T(X))$$

POPULAR
SYMMETRIES
SPACE

SPACE-TRANSLATION
INVARIANCE

SPACE-REVERSAL
SYMMETRY :

GIVEN A REFLECTION R OF PARTICLES WITH RESPECT TO A CELL, TRANSITION PROBABILITY $X \rightarrow Y$ IS THE SAME AS PROBABILITY FOR REFLECTED MICROSTATES $R(X) \rightarrow R(Y)$



$$B(X \rightarrow Y | X) = B(R(X) \rightarrow R(Y) | R(X))$$

POPULAR
SYMMETRIES
SPACE

SPACE-ROTATION
INVARIANCE

TIME - TRANSLATION

SYMMETRY:

THE TRANSITION PROBABILITY $X \rightarrow Y$ IS
INDEPENDENT OF TIME:

$$P(X \rightarrow Y | X)(t) = \text{CONST}(t)$$

THE DICE DOES NOT CHANGE WITH TIME

POPULAR
SYMMETRIES
TIME

TIME - TRANSLATION
INVARIANCE

TIME - REVERSAL

SYMMETRY:

FOR EVERY RECURRING TIME-SEQUENCE OF MICROSTATES $(X_1, X_2 \dots X_m, X_1)$ THE PROBABILITY OF VISITING MICROSTATES IN THE ORIGINAL ORDER OR THE REVERSED ONE IS THE SAME:

$$B(X_1 \rightarrow X_2 | X_1) \cdot B(X_2 \rightarrow X_3 | X_2) \cdots B(X_{m-1} \rightarrow X_m | X_{m-1}) \cdot B(X_m \rightarrow X_1 | X_m) = \\ = B(X_1 \rightarrow X_m | X_1) \cdot B(X_m \rightarrow X_{m-1} | X_m) \cdots B(X_3 \rightarrow X_2 | X_3) \cdot B(X_2 \rightarrow X_1 | X_2)$$

(KOLMOGOROV CYCLING CONDITION \uparrow
FOR ERGODIC SYSTEMS)

POPULAR
SYMMETRIES
TIME

TIME - REVERSAL
INVARIANCE

TRANSITION-MATRIX

SYMMETRY:

$$B(x \rightarrow y | x) = B(y \rightarrow x | y)$$

+ ERGODICITY \rightarrow UNIQUE STEADY STATE



STEADY STATE = EQUILIBRIUM STATE

PROBABILITY TO FIND A MICROSTATE x IN AN ENSEMBLE OF SYSTEMS FOR $t \rightarrow \infty$ IS INDEPENDENT OF TIME AND x :

$$\begin{aligned}\tilde{P}(x, t) &= 1 / W_{IND} \\ &= 1 / W_{DIST}\end{aligned}$$

POPULAR
SYMMETRIES
EQUILIBRIUM

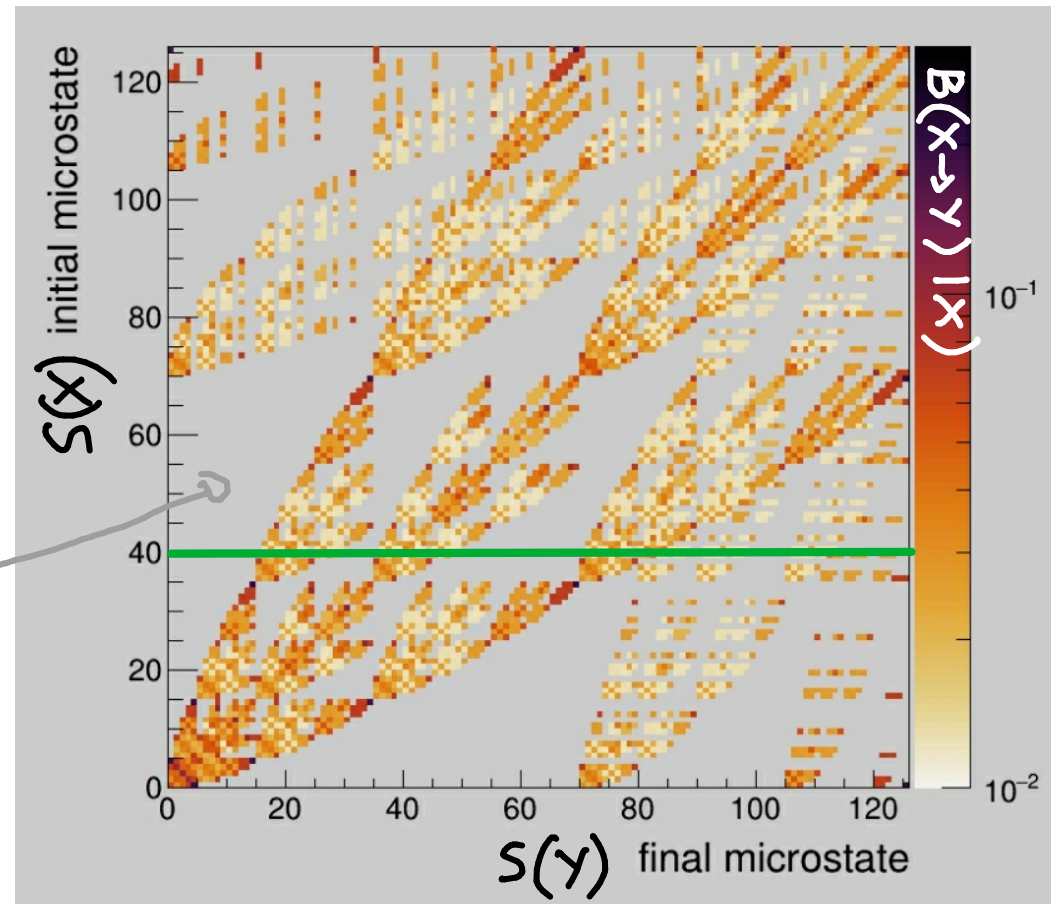
EQUILIBRIUM
STATES
IN NATURE

EXAMPLE OF
TRANSITION MATRIX
WITH POPULAR SYMMETRIES

THE DICE EXAMPLE

FORBIDDEN
TELEPORTATION
TRANSITIONS
($B(x \rightarrow y | x) = 0$)

THE DICE

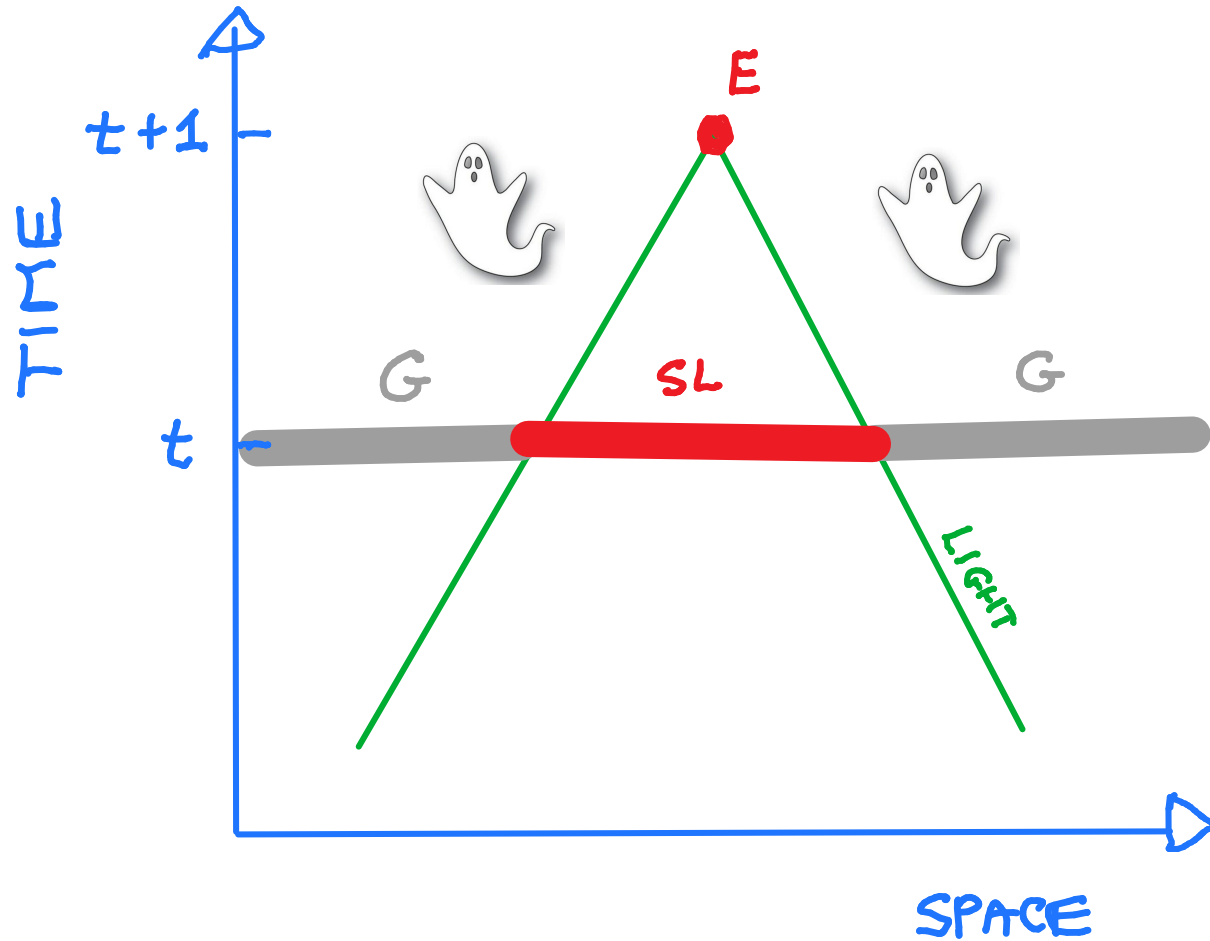


IN SPITE OF ALL SIMPLIFICATIONS,
IT IS RATHER COMPLICATED

NOTE, MICROSTATE SEQUENTIAL NUMBERS $S(x)$, $S(y)$ UNIQUELY RELATE TO MICROSTATES x AND y , SEE THE PAPER, BUT THE RELATION IS ARBITRARILY SELECTED

REMINDER:

GHOSTS - DEFINITION:



$$\text{IF } P(E | (SL, G)) \neq P(E | SL)$$

- NON-TRIVIAL DEPENDENCE BETWEEN REMOTE EVENTS, 'GHOSTS' IN NATURE (VIOLATION OF STRONG LOCALITY)

RESULTS