

Observation of Odderon Exchange at TeV Energies

based on Real extensions of the Bialas – Bzdak (ReBB) model

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Timeline

Formalism of elastic scattering
Discoveries of Odderon exchange

Model-independent
ReBB model dependent

H(x): Generalized Geometric Scaling

Derivations of H(x) scaling

Model independently

In the ReBB model and
why it works at LHC energies

The logo for MATE (Matters of Acting Theoretical Experimental) is displayed in a stylized, blocky font with a green-to-white gradient.The logo for Wigner Institute for Physics is shown, featuring a stylized particle detector or beam line icon above the word "WIGNER" in a bold, sans-serif font.

Summary

TIMELINE

1986-1988: INTERMITTENCY

Moments of Rapidity Distributions as a Measure of Short Range Fluctuations in High-Energy Collisions #2

A. Bialas (Orsay, LPT and Jagiellonian U.), Robert B. Peschanski (Saclay) (Sep, 1985)

Published in: *Nucl.Phys.B* 273 (1986) 703-718

[DOI](#) [cite](#)



Nuclear Physics B

Volume 273, Issues 3-4, 1 September 1986, Pages 703-718



[reference search](#) [↻ 1,151 citations](#)

Moments of rapidity distributions as a measure of short-range fluctuations in high-energy collisions

Intermittency in Multiparticle Production at High-Energy #1

A. Bialas (Jagiellonian U.), Robert B. Peschanski (Saclay) (Mar, 1988)

Published in: *Nucl.Phys.B* 308 (1988) 857-867

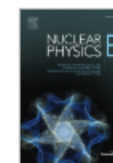
[DOI](#) [cite](#) [claim](#)

Abstract

It is proposed to study the dependence of factorial moments of the rapidity distribution on the size δy of the resolution. It is shown that if the fluctuations are purely statistical no variation of moments δy is expected, and thus observation of such a variation indicates the presence of genuine fluctuations of physical origin. The region in which the change occurs corresponds to the size (in rapidity) of the observed fluctuations. Intermittency, i.e. fluctuations of many different sizes, would show up as a power-law behaviour of moments on δy . A good experimental resolution can be obtained from the rapidity distribution of high multiplicity events.

Nuclear Physics B

Volume 308, Issue 4, 24 October 1988, Pages 857-867



The physics of hadrons

Intermittency in multiparticle production at high energy ☆

A. Bialas, R. Peschanski

[reference search](#) [↻ 739 citations](#)

[Get rights and content](#)

Abstract

We present predictions of random cascading models for multiparticle production at high energy. Standard and correlated factorial moments in rapidity are shown to provide stringent tests of the intermittency patterns characteristic of random cascades. Using the central limit theorem we show how to test directly for the existence of a cascading process. Finally, we discuss how to take into account statistical corrections.

1992 - 2004: INTERMITTENCY & HBT EFFECT

Vol. 23(1992)

ACTA PHYSICA POLONICA B



Bose-Einstein correlations for Levy stable source distributions

#4

T. Csorgo (Budapest, RMKI), S. Hegyi (Budapest, RMKI), W.A. Zajc (Columbia U.) (Oct, 2003)

Published in: *Eur.Phys.J.C* 36 (2004) 67-78 • e-Print: [nucl-th/0310042](#) [nucl-th]

 pdf  DOI  cite  claim

 reference search  157 citations

INTERMITTENCY AND THE HANBURY-BROWN-TWISS EFFECT

A. BIALAS

Institute of Physics, Jagellonian University
Reymonta 4, 30-059 Cracow, Poland

(Received April 15, 1992)

Dedicated to Wiesław Czyż in honour of his 65th birthday.

Relation of the intermittency phenomenon observed in spectra produced in high-energy collisions to the intensity correlation between identical particles is discussed. It is argued that the connection of the two effects requires scale-invariant (power law) fluctuation size of the interaction region.

Eur. Phys. J. C 36, 67–78 (2004)

Digital Object Identifier (DOI) [10.1140/epjc/s2004-01870-9](#)

THE EUROPEAN
PHYSICAL JOURNAL C

Bose–Einstein correlations for Lévy stable source distributions

T. Csörgő^{1,a}, S. Hegyi^{1,b}, W.A. Zajc^{2,c}

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Received: 19 November 2003 / Revised version: 27 April 2004 /




Published online: 23 June 2004 – © Springer-Verlag / Società Italiana di Fisica 2004

Abstract. The peak of the two-particle Bose–Einstein correlation functions has a very interesting structure. It is often believed to have a multivariate Gaussian form. We show here that for the class of stable distributions, characterized by the index of stability $0 < \alpha \leq 2$, the peak has a stretched exponential shape. The Gaussian form corresponds then to the special case of $\alpha = 2$. We give examples for the Bose–Einstein correlation functions for univariate as well as multivariate stable distributions, and we check the model against two-particle correlation data.

Intermittency and the Hanbury-Brown-Twiss effect

A. Bialas (Jagiellonian U.) (Apr, 1992)

Published in: *Acta Phys.Polon.B* 23 (1992) 561-567

 pdf  cite  claim

2006: BIALAS-BZDAK MODEL FOR ELASTIC PP

Constituent quark and diquark properties from small angle proton-proton elastic scattering at high energies

A. Bialas (Jagiellonian U.), A. Bzdak (Jagiellonian U.) (Dec, 2006)

Published in: *Acta Phys.Polon.B* 38 (2007) 159-168 • e-Print: [hep-ph/0612038](https://arxiv.org/abs/hep-ph/0612038) [hep-ph]

pdf links cite claim

reference search 66 citati

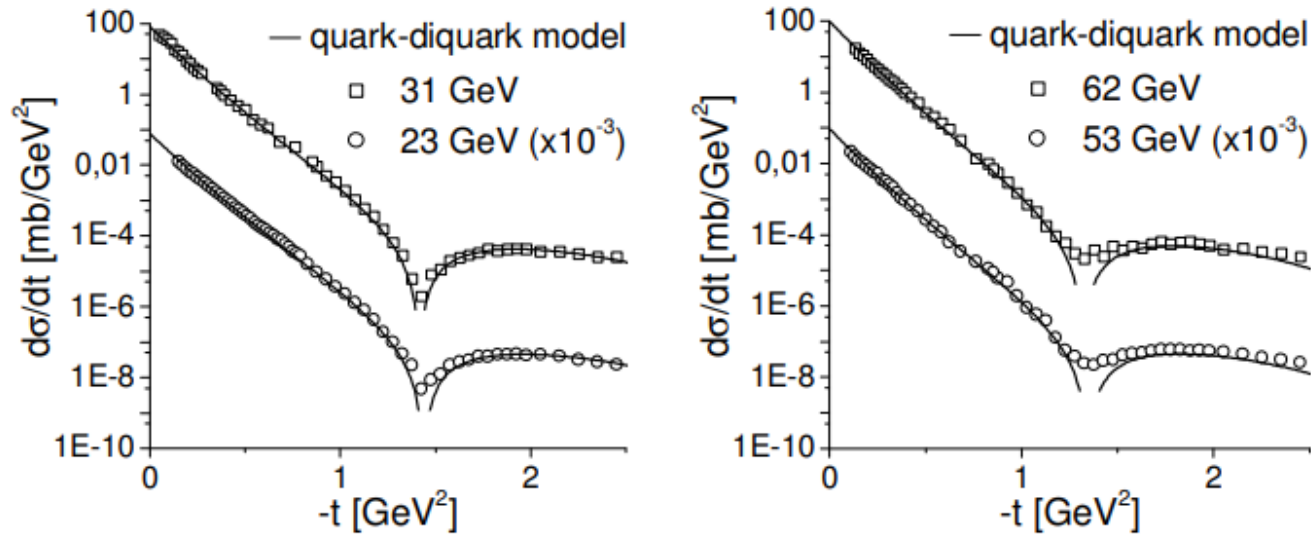


Fig. 3. The quark–diquark model compared to data on differential cross section at four energies. Diquark as a qq system.

**Singularity of BB at the dip
→ Add real part to fix it.**

2. Low momentum transfer elastic scattering in the quark–diquark model

We follow the standard point of view that the imaginary part of the elastic scattering amplitude, dominating at high energy, is generated by the absorption of the incident particle wave, represented by the inelastic (non-diffractive) collisions. The inelastic proton–proton cross-section at a fixed impact parameter b , $\sigma(b)$, is calculated using the rules of the probability calculus. One writes

$$\sigma(b) = \int d^2 s_q d^2 s'_q d^2 s_d d^2 s'_d D(s_q, s_d) D(s'_q, s'_d) \sigma(s_q, s_d; s'_q, s'_d; b), \quad (1)$$

where $D(s_q, s_d)$ denotes the distribution of quark and diquark inside the nucleon, $s_q(s'_q)$, $s_d(s'_d)$ are transverse positions of the quarks and diquarks in the two colliding nucleons, and $\sigma(s_q, s_d; s'_q, s'_d; b)$ is the probability of interaction at fixed impact parameter and the transverse positions of all constituents taking part in the process. This configuration is illustrated in Fig. 1. Since the constituents act independently we have [11,12]:

$$1 - \sigma(s_q, s_d; s'_q, s'_d; b) = [1 - \sigma_{qq}(b + s'_q - s_q)][1 - \sigma_{qd}(b + s'_d - s_q)] \times [1 - \sigma_{dq}(b + s'_q - s_d)][1 - \sigma_{dd}(b + s'_d - s_d)], \quad (2)$$

where $\sigma_{ab}(s) \equiv d^2 \sigma_{ab}(s)/d^2 s$ are inelastic differential cross-sections of the constituents (ab denotes qq , qd or dd).

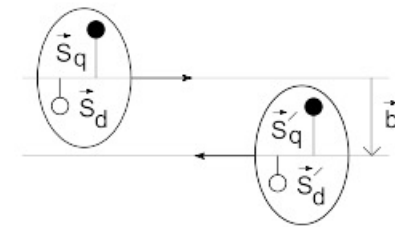


Fig. 1. Proton–proton scattering in the quark–diquark model.

For the distribution of the constituents inside the nucleon we take a Gaussian with radius R :

2015: BIALAS-BZDAK MODEL REAL EXTENSION

Excitation function of elastic pp scattering f

F. Nemes (CERN and Wigner RCP, Budapest), T. Csörgő (2015)

Published in: *Int.J.Mod.Phys.A* 30 (2015) 14, 1550076 • e-

pdf DOI cite claim

The unitarity of the scattering matrix S is expressed

$$SS^\dagger = I,$$

where I is the identity matrix. The decomposition of the transition matrix, leads the unitarity relation Eq.

$$T - T^\dagger = iTT^\dagger,$$

which can be rewritten in the impact parameter b

$$2 \operatorname{Im} t_{el}(s, b) = |t_{el}(s, b)|^2 + \bar{\sigma}_i$$

We have studied several possible choices. (1) the imaginary part of the opacity function so that it is the probability of inelastic scatterings, which is known as the probability of the impact parameter b . A possible interpretation is that the inelastic collisions arising from non-colli- and diquarks follow the same spatial distribution as the same constituents

$$\operatorname{Im} \Omega(s, b) = -\alpha \cdot \bar{\sigma}_{inel}(s, b)$$

p+p → p+p, diquark as a single entity at $\sqrt{s}=23.5$ GeV

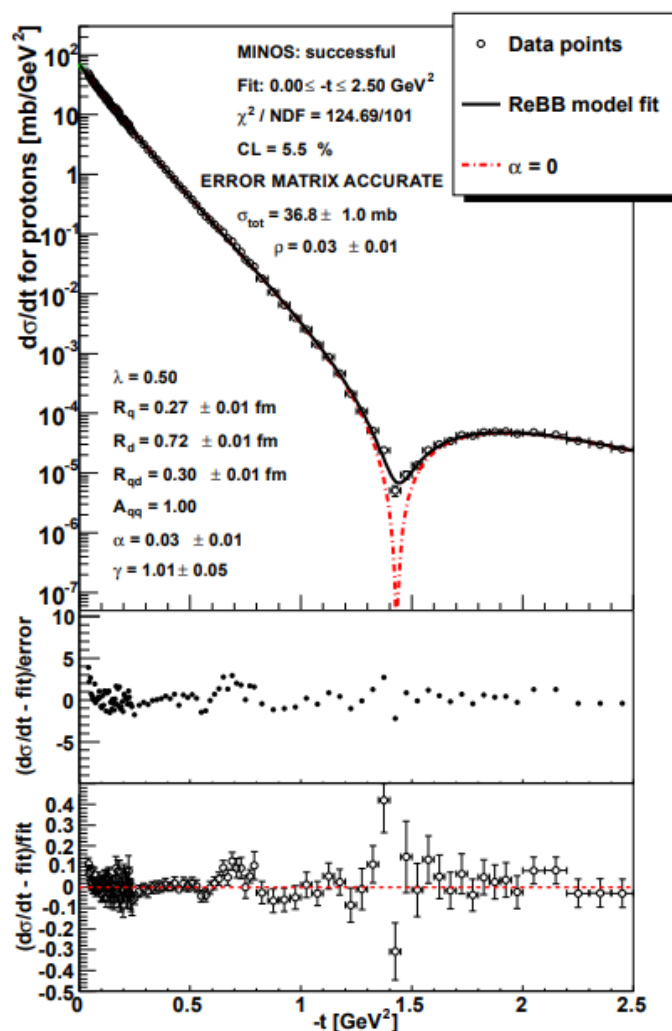


Figure 4: The fit of the ReBB model at $\sqrt{s} = 23.5$ GeV in the $0 < |t| < 2.5$ GeV² squared four-momentum transfer $|t|$ range. The fit uses the statistical errors of the data points and the luminosity error of the systematic uncertainty according to Eq. (32). Parameter values are rounded up to two valuable decimal digits.

Singularity of BB at the dip

Add real part to fix it.

ReBB solves three problems:
Describes dip-bump structure

Provides $\rho_0 \neq 0$

Provides Odderon

second order polynomial equation in terms of $t_{el}(s, b)$. If one introduces the opacity or eikonal

$$t_{el}(s, b) = i \left[1 - e^{-\Omega(s, b)} \right], \quad (4)$$

$$\sigma_{inel}(s, b) = 1 - e^{-2 \operatorname{Re} \Omega(s, b)}. \quad (5)$$

and eikonal form. From Eqs. (5) the real part of the eikonal can be expressed as

$$\operatorname{Re} \Omega(s, b) = -\frac{1}{2} \ln [1 - \bar{\sigma}_{inel}(s, b)]. \quad (6)$$

It is assumed that the real part of t_{el} vanishes. In this case

$$t_{el}(s, b) = i \left[1 - \sqrt{1 - \bar{\sigma}_{inel}(s, b)} \right]. \quad (7)$$

When taken into account in Eq. (4) the result is

$$t_{el}(s, b) = e^{-i \operatorname{Im} \Omega(s, b)} \sqrt{1 - \bar{\sigma}_{inel}(s, b)}, \quad (8)$$

The physical interpretation of $\operatorname{Im} \Omega(s, b)$ is discussed later.

ELASTIC SCATTERING

Formalism: elastic scattering

$$\frac{d\sigma(s)}{dt} = \frac{1}{4\pi} |T_{el}(s, \Delta)|^2, \quad \Delta = \sqrt{|t|}.$$

$$\sigma_{el}(s) = \int_0^\infty d|t| \frac{d\sigma(s)}{dt}$$

$$A(s) = \lim_{t \rightarrow 0} \frac{d\sigma}{dt}(s, t)$$

$$A(s) = \frac{1}{16\pi} (1 + \rho_0^2(s)) \sigma_{tot}^2(s)$$

$$B(s, t) = \frac{d}{dt} \ln \frac{d\sigma(s)}{dt}$$

$$\rho(s, t) \equiv \frac{\text{Re } T_{el}(s, \Delta)}{\text{Im } T_{el}(s, \Delta)}$$

$$B(s) \equiv B_0(s) = \lim_{t \rightarrow 0} B(s, t),$$

$$\rho(s) \equiv \rho_0(s) = \lim_{t \rightarrow 0} \rho(s, t)$$

$$\sigma_{tot}(s) \equiv 2 \text{Im } T_{el}(\Delta = 0, s)$$

Basic problem: $d\sigma/dt$ measures an amplitude, *modulus squared*.
If Odderon exists: signals in elastic scattering at $t = 0$ and at $-t > 0$.

Formalism in b space

$$\frac{d\sigma(s)}{dt} = \frac{1}{4\pi} |T_{el}(s, \Delta)|^2, \quad \Delta = \sqrt{|t|}.$$

$$\begin{aligned} t_{el}(s, b) &= \int \frac{d^2\Delta}{(2\pi)^2} e^{-i\Delta \mathbf{b}} T_{el}(s, \Delta) = \\ &= \frac{1}{2\pi} \int J_0(\Delta b) T_{el}(s, \Delta) \Delta d\Delta, \\ \Delta &\equiv |\mathbf{\Delta}|, \quad b \equiv |\mathbf{b}|. \end{aligned}$$

$$t_{el}(s, b) = i \left[1 - e^{-\Omega(s, b)} \right]$$

$$P(s, b) = 1 - \left| e^{-\Omega(s, b)} \right|^2$$

Impact parameter or b space:
elastic scattering interferes with propagation w/o collisions: Genuine quantum physics.

In general, a complex opacity function $\Omega(\mathbf{s}, \mathbf{b})$ (eikonal, from unitarity)

$0 \leq P(s, b) \leq 1$: *inelastic* scattering has a probabilistic interpretation

Looking for Crossing-Odd(eron) effects

$$\begin{aligned}T_{\text{el}}^{pp}(s,t) &= T_{\text{el}}^+(s,t) - T_{\text{el}}^-(s,t), \\T_{\text{el}}^{p\bar{p}}(s,t) &= T_{\text{el}}^+(s,t) + T_{\text{el}}^-(s,t), \\T_{\text{el}}^+(s,t) &= T_{\text{el}}^P(s,t) + T_{\text{el}}^f(s,t), \\T_{\text{el}}^-(s,t) &= T_{\text{el}}^O(s,t) + T_{\text{el}}^\omega(s,t).\end{aligned}$$

$$\begin{aligned}T_{\text{el}}^P(s,t) &= \frac{1}{2} \left(T_{\text{el}}^{pp}(s,t) + T_{\text{el}}^{p\bar{p}}(s,t) \right) \\T_{\text{el}}^O(s,t) &= \frac{1}{2} \left(T_{\text{el}}^{p\bar{p}}(s,t) - T_{\text{el}}^{pp}(s,t) \right)\end{aligned}$$

for $\sqrt{s} \geq 1 \text{ TeV}$,

Three simple consequences:

$$T_{\text{el}}^O(s,t) = 0 \implies \frac{d\sigma^{pp}}{dt} = \frac{d\sigma^{p\bar{p}}}{dt} \quad \text{for } \sqrt{s} \geq 1 \text{ TeV}$$

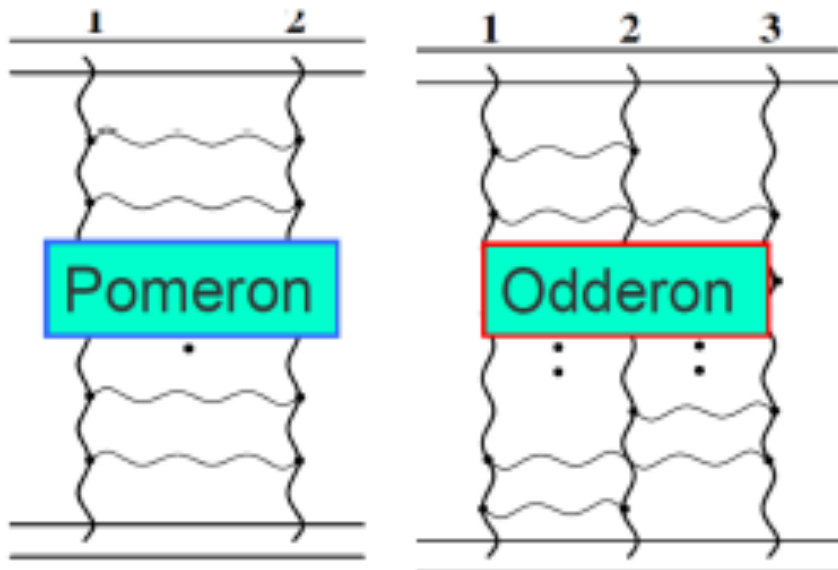
$$\frac{d\sigma^{pp}}{dt} = \frac{d\sigma^{p\bar{p}}}{dt} \quad \text{for } \sqrt{s} \geq 1 \text{ TeV} \not\Rightarrow T_{\text{el}}^O(s,t) = 0.$$

$$\frac{d\sigma^{pp}}{dt} \neq \frac{d\sigma^{p\bar{p}}}{dt} \quad \text{for } \sqrt{s} \geq 1 \text{ TeV} \implies T_{\text{el}}^O(s,t) \neq 0$$

Odderon: extremely elusive, for 48 years

Odderon: L. Lukaszuk, B. Nicolescu,
Lett. Nuovo Cim. 8, 405 (1973)
Received: 31 July 1973

Odderon is an odd component of
elastic scattering:
Changes sign for crossing



СООБЩЕНИЯ
ОБЪЕДИНЕННОГО
ИНСТИТУТА
ЯДЕРНЫХ
ИССЛЕДОВАНИЙ
Дубна



E2-6350

A.V.Efremov, R.Peschanski

EVIDENCE FOR NEW SINGULARITIES
IN REGGE PHENOMENOLOGY

ЛАБОРАТОРИЯ ТЕОРЕТИЧЕСКОЙ ФИЗИКИ

1972

Odderon name coined: D. Joynson, E. Leader, B. Nicolescu, C. Lopez,
Nuovo Cim. 30A, 345 (1975) - Well established in QCD by now !
Honorable mention: A. V. Efremov, R. Peschanski, JINR-E2-6350 (1972)

Odderon: well established in QCD

Odderon proposed in Regge phenomenology:
L. Lukaszuk, B. Nicolescu, Lett. Nuovo Cim. 8, 405 (1973)

Three Gluon Integral Equation and Odd c Singlet Regge Singularities in QCD
BKP evolution equation

J. Bartels, Nucl. Phys. B 175 (1980) 365-401
J. Kwiecinski, M. Praszalowicz, Phys.Lett.B 94 (1980) 413-416

A new Odderon intercept from QCD:
R. A. Janik, J. Wosiek, Phys. Rev. Lett. 82 (1999) 1092

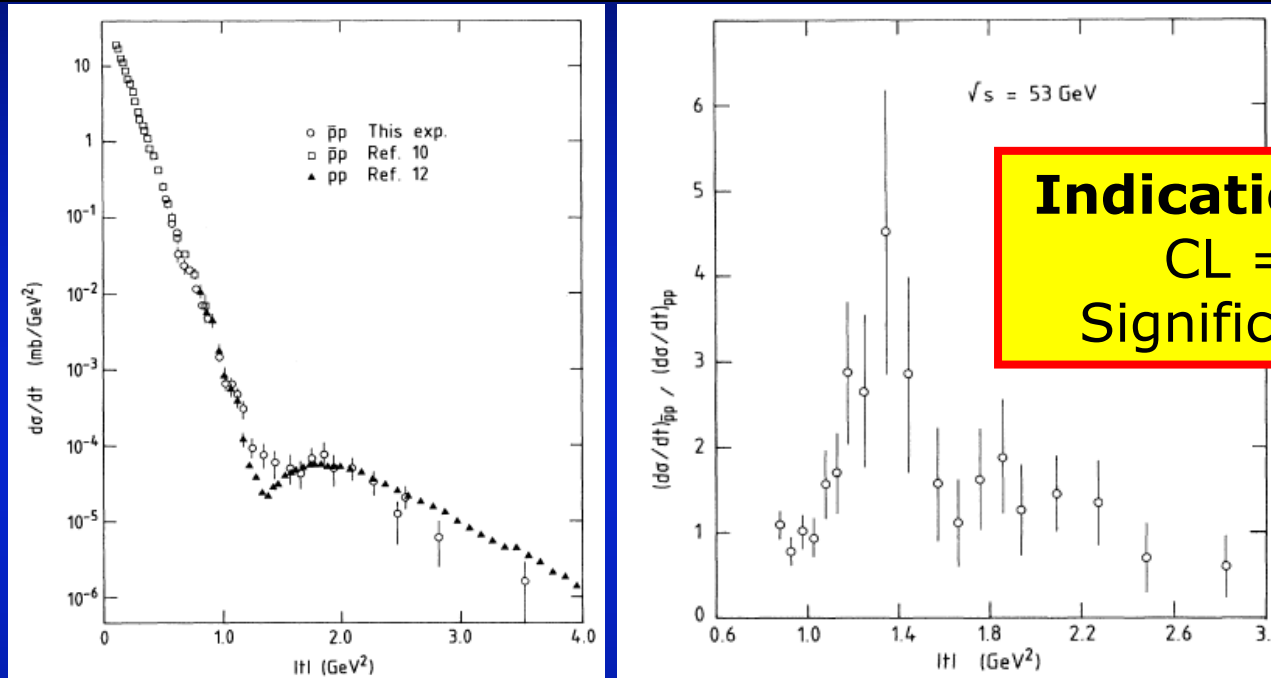
Odderon in QCD:
J. Bartels, L.N. Lipatov, G. P. Vacca: Phys. Lett. B (2000) 178

Odderon in QCD with running coupling:
J. Bartels, C. Contreras, G. P. Vacca, *JHEP* 04 (2020) 183

For an excellent theory intro/review, see Yu. Kovchegov's
CTEQ Webinar, April 28, 2021
<http://youtu.be/yHBO3zcB3V4>

Odderon: elusive experimentally

Odderon search at ISR: indication but no conclusive result
Breakstone et al, Phys. Rev. Lett. 54, 2180 (**1985**): CL = 99.9 %



Terminology for **this** talk:

Agreement if statistical significance is $< 3 \sigma$

Indication of signal if $3 \sigma \leq \text{significance} < 5 \sigma$

Evidence or observation of signal if $5 \sigma \leq \text{significance}$

Discovery of signal if $5 \sigma \leq \text{significance}$, **for the first time**.

Accepted: Discovery if [Clay Mathematical Institute \(CMI\) criteria](#) satisfied.

Miscovery if [CMI criteria for Millenium Prize Problems](#) are **not** satisfied.

A. STER, L. JENKOVSKY, T. Cs. (PRD 2015): PREDICTION OF ODDERON DISCOVERY AT LHC

PHYSICAL REVIEW D 91, 074018 (2015)

Extracting the Odderon from pp and $p\bar{p}$

András Ster^{*}

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H-1525 Budapest 114, P.O. Box

László Jenkovszky

BTFP, National Academy of Sciences of Hungary
and Wigner Research Centre for Physics, Hunga
H-1525 Budapest 114, P.O. Box

Tamás Csörgő[†]

Wigner Research Centre for Physics, Hunga
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and KRF, H-3200 Gyöngyös, Mátrai
(Received 15 January 2015; published 15 February 2015)

Starting from a simple empirical parametrization of the scattering dip-bump structure of elastic pp scattering in t at fixed values between missing energy intervals to extract the Odderon contribution to pp elastic and total cross sections. The model is fitted to data extracted from the Odderon and its ratio to the Pomeron. From our model as predicted by J. Bartels, L. N. Lipatov, and G. P. Vacca in perturbative QCD.

DOI: 10.1103/PhysRevD.91.074018

$d\sigma/dt$ data for pp , $p\bar{p}$ or both

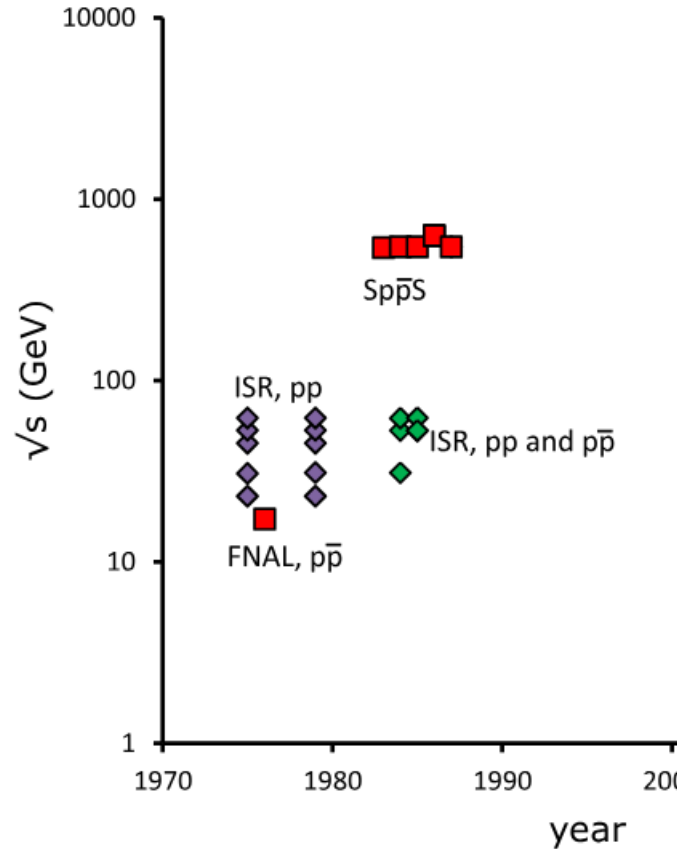
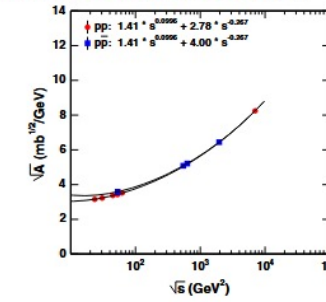


FIG. 1 (color online). Timeline of proton scattering measurements. New accelerators have been built, increasing the maximum available energies; however, at the same time, the maximum energy per nucleon has decreased. At the same time, the pp and the $p\bar{p}$ elastic scattering data were measured at the same $\sqrt{s} = 31, 53, \text{ and } 62 \text{ GeV}$.

EXTRACTING THE ODDERON FROM pp AND $p\bar{p}$...



PHYSICAL REVIEW D 91, 074018 (2015)

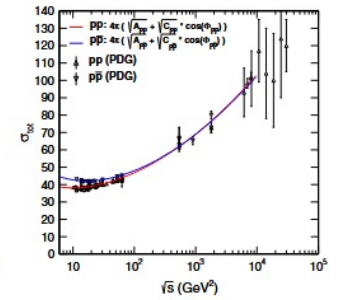
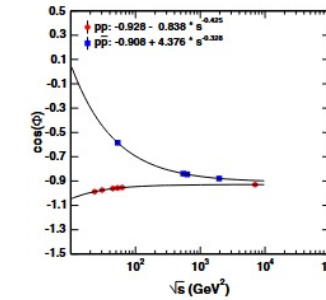
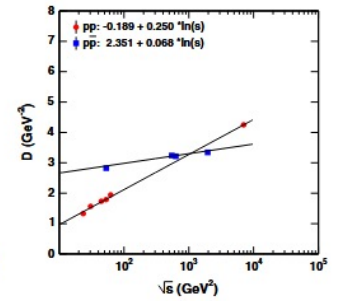
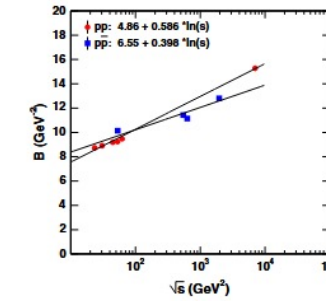
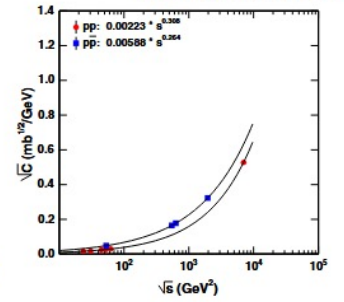


FIG. 4 (color online). Energy-dependent values of the parameters from a fit to pp and $p\bar{p}$ data, constrained by $\sigma_t^{pp} = \sigma_t^{p\bar{p}}$ as $\sqrt{s} \rightarrow \infty$. The lowest right icon shows a cross-check for the (asymptotically converging) total cross sections.

Odderon search with scaling: a possible strategy

Known trivial s-dependences in
 $\sigma_{\text{tot}}(s), \sigma_{\text{el}}(s), B(s), \rho(s)$

Try to scale this out
Data collapsing (scaling)

Look for scaling violations

In the TeV energy range:
Odderon is equivalent with
a crossing-odd component
Look for violations of C-symmetry

DISCOVERY: ODDERON EXCHANGE

Odderon: first observation with $> 5 \sigma$

EPJ Web of Conf. (2020) **235**: 06005

<https://doi.org/10.1051/epjconf/202023506002>

Proton Holography -- Discovering Odderon from Scaling Properties of Elastic Scattering #4

T. Csorgo (Wigner RCP, Budapest and Eszterhazy Karoly U., Eger), T. Novak (EKU KRC, Gyongyos), R. Pasechnik (Lund U. and Rez, Nucl. Phys. Inst.), A. Ster (Wigner RCP, Budapest), I. Szanyi (Wigner RCP, Budapest and Eotvos U.) (Apr 15, 2020)

Published in: *EPJ Web Conf.* 235 (2020) 06002 • Contribution to: ISMD 2019 • e-Print: 2004.07095 [hep-ph]

First publication of an at least 5.0σ (~~6.26σ~~) Odderon exchange effect:
published on May 11, 2020,

EPJ Web of Conf. 235 (2020) 06002

in an **anonymously refereed / peer reviewed** conference proceedings.

(Proc. ISMD 2019, Santa Fe, USA)

BUT: „Never be the first! It is too early!”

P. Carruthers ~ 1990

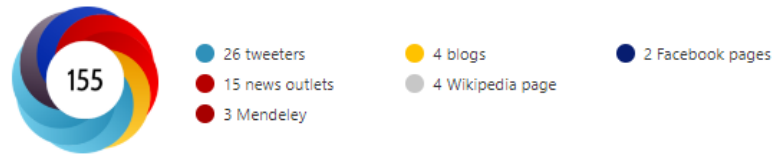
First journal publications, Odderon $> 5 \sigma$

Evidence of Odderon-exchange from scaling properties of elastic scattering at TeV energies #5

T. Csörgő (Wigner RCP, Budapest and CERN), T. Novák (Unlisted, HU), R. Pasechnik (Lund U., Dept. Theor. Phys.), A. Ster (Wigner RCP, Budapest), I. Szanyi (Wigner RCP, Budapest) (Dec 26, 2019)

Published in: *Eur.Phys.J.C* 81 (2021) 2, 180 • e-Print: 1912.11968 [hep-ph]

Online attention



This article is in the 98th percentile (ranked 6,037th) of the 428,075 tracked articles of a similar age in all journals and the 99th percentile (ranked 1st) of the 231 tracked articles of a similar age in *The European Physical Journal C*

Hungarian-Swedish team:

Eur. Phys. J. C (2021) **81**: 180, Published: 23 February 2021
<https://doi.org/10.1140/epjc/s10052-021-08867-6>

Observation of Odderon effects at LHC energies: a real extended Bialas-Bzdak model study #2

T. Csorgo (Wigner RCP, Budapest and EKV KRC, Gyongyos), I. Szanyi (Eotvos U. and Wigner RCP, Budapest) (May 28, 2020)

Published in: *Eur.Phys.J.C* 81 (2021) 7, 611 • e-Print: 2005.14319 [hep-ph]

Hungarian team, Real Extended Bialas-Bzdak model:

Eur. Phys. J. C (2021) **81**:611, Published: 13 July 2021
<https://doi.org/10.1140/epjc/s10052-021-09381-5>

Odderon Exchange from Elastic Scattering Differences between pp and $p\bar{p}$ Data at 1.96 TeV and from pp Forward Scattering Measurements

TOTEM and D0 Collaborations • V.M. Abazov (Dubna, JINR) et al. (Dec 7, 2020)

Published in: *Phys.Rev.Lett.* 127 (2021) 6, 062003 • e-Print: 2012.03981 [hep-ex]



SUMMARY	News	Blogs	Twitter	Wikipedia	Dimensions citations
Title	Odderon Exchange from Elastic Scattering Differences between pp and $p\bar{p}$ Data at 1.96 TeV and from pp Forward Scattering Measurements				
Published in	Physical Review Letters, August 2021				
DOI	10.1103/PhysRevLett.127.062003				
Pubmed ID	34420329				
Authors	V. M. Abazov, B. Abbott, B. S. Acharya, M. Adams, T. Adams, J. P. Agnew, G. D. Alexeev, G. Alkhazov...				

TWITTER DEMOGRAPHICS

MENDELEY READERS

Defense of this D0-TOTEM paper is in progress

D0 and TOTEM Collaborations:

Phys. Rev. Lett. **127** (2021) 6, 062003, Published: 4 August 2021
<https://doi.org/10.1103/PhysRevLett.127.062003>

2022 observations of Odderon with $> 5 \sigma$

Characterisation of the dip-bump structure observed in proton–proton elastic scattering at $\sqrt{s} = 8 \text{ TeV}$ #1

TOTEM Collaboration • G. Antchev (Pilsen U.) et al. (Nov 23, 2021)

Published in: *Eur.Phys.J.C* 82 (2022) 3, 263 • e-Print: 2111.11991 [hep-ex]

Online attention



This article is in the 1st percentile (ranked 279,419th) of the 343,918 tracked articles of a similar age in all journals and the 1st percentile (ranked 73rd) of the 114 tracked articles of a similar age in *The European Physical Journal C*

TOTEM Collaboration:

8 TeV: EPJ C (2022) 82, 263 (2022). Published: March 26, 2022
<https://doi.org/10.1140/epjc/s10052-022-10065-x>
Publishes final data for D0-TOTEM PRL published in 2021

The ReBB model and its H(x) scaling version at 8 TeV: Odderon exchange is a certainty #1

I. Szanyi (Eotvos U. and Wigner RCP, Budapest and Karoly Robert U. Coll.), T. Csörgő (Wigner RCP, Budapest and Karoly Robert U. Coll.) (Apr 21, 2022)

Published in: *Eur.Phys.J.C* 82 (2022) 9, 827, *Eur.Phys.J.C* 82 (2022) 827 • e-Print: 2204.10094 [hep-ph]

Online attention



This article is in the 64th percentile (ranked 57,525th) of the 166,532 tracked articles of a similar age in all journals and the 99th percentile (ranked 1st) of the 1 tracked articles of a similar age in *The European Physical Journal C*

Hungarian team, using the Real extended Bialas-Bzdak (ReBB) model:
New TOTEM 8 TeV data vs ReBB model predictions:
EPJ C 82 (2022) 9, 827. Published: Sept 19, 2022
In the ReBB model, Odderon exchange is a certainty-

What about model independent results?

→ See some of the backup slides

Three Oldest Hungarian Universities

UP Story - 650 years

Home » University » UP Story 650 years

University of Pécs: 1367

University of Pécs:
S: Oldest, C: in Hungary

The history of higher education in Pécs dates back to 1367, when Louis the Great, King of Hungary, established the University of Pécs in the episcopal city of Pécs. As a result of an integration process, which has become one of the most famous, prestigious universities in Hungary. It has ten faculties which cover the full spectrum of higher education.

1367



The University of Debrecen is one of the oldest universities of Hungary.

The institution of higher education in the country operated continuously in the same city, is one of the research universities in Hungary offering the widest spectrum of educational programs in 14 faculties and 24 doctoral schools.

University of Debrecen: 1528

University of Debrecen:

S: Oldest, C: in Hungary,

operating continuously and in the same city

education and culture. The *gerundium*, showing respect

tion
arian
well,
ires,

(S,C) structure evident,

S: statement, valid if

C: condition is satisfied

See talk of [R. Dardashti](#) at ISMD21

Condition changes → Statement changes (!!)

Eötvös University:

S: Oldest, C: in Hungary,

teaching continuously

... in Nagyszombat in 1635 (sixteen thirty-five) by Archbishop of Esztergom, Péter Pázmány, and it is the oldest Hungarian university where the teaching has continued uninterrupted since its inception. More than sixty years

Eötvös Loránd University: 1635

Model independent $H(x|pp)$ scaling yields Odderon $> 6.26 \sigma$

Evidence of Odderon-exchange from scaling properties of elastic scattering at TeV energies #5

T. Csörgő (Wigner RCP, Budapest and CERN), T. Novák (Unlisted, HU), R. Pasechnik (Lund U., Dept. Theor. Phys.), A. Ster (Wigner RCP, Budapest), J. Szanyi (Wigner RCP, Budapest) (Dec 26, 2019)

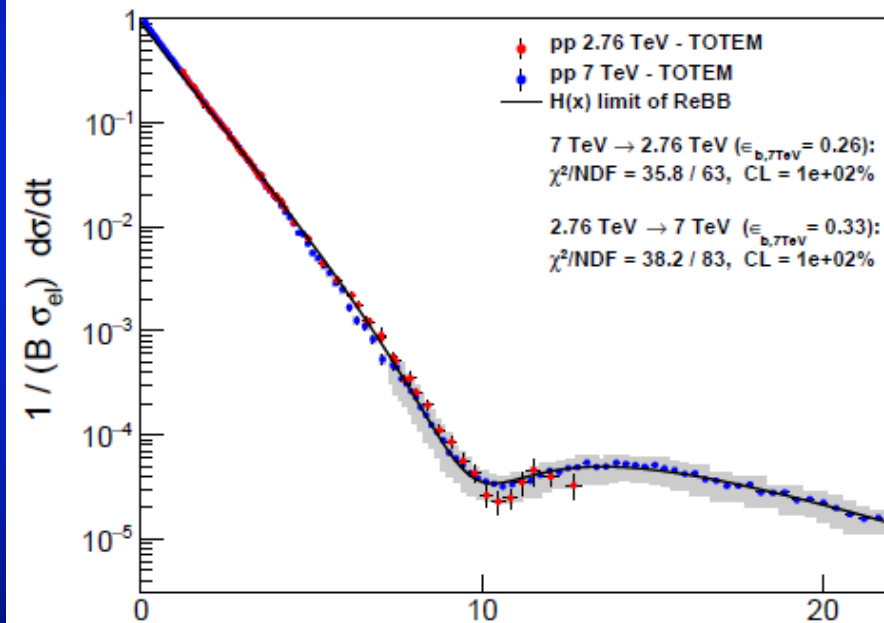
Published in: *Eur.Phys.J.C* 81 (2021) 2, 180 • e-Print: 1912.11968 [hep

Eur. Phys. J. C (2021) 81: 180, published February 2021

<https://doi.org/10.1140/epjc/s10052-021-08867-6>

pdf DOI cite

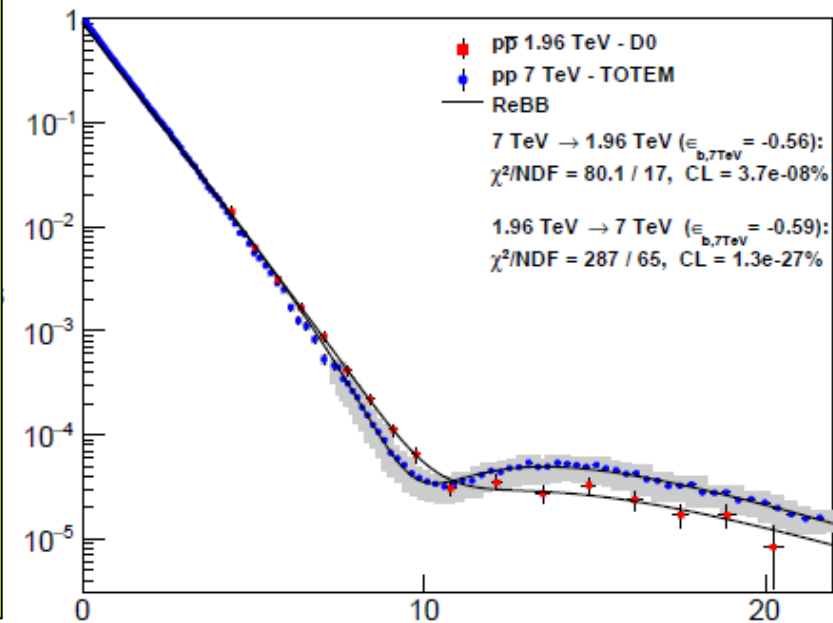
15 citations



$B \equiv B_0(s)$ from now on

$-Bt$

$H(x) = 1/(B \sigma_{el}) d\sigma/dt$ vs $x = -Bt$



$-Bt$

S: Model independent Odderon significance $\geq 6.26 \sigma$

C1: All D0 and TOTEM published data at 1.96, 2.76 and 7.0 TeV

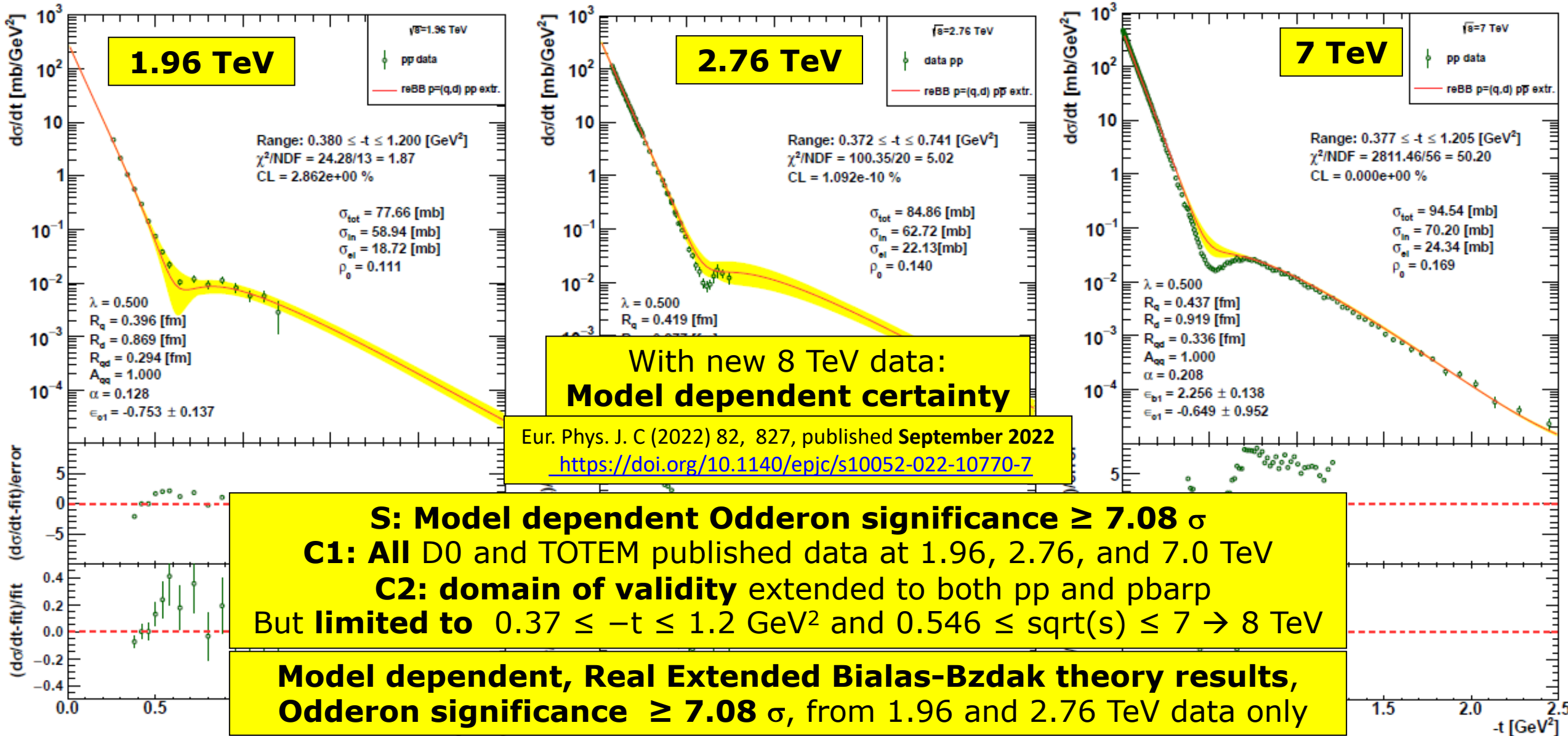
C2: domain of validity is still determined model dependently.

REAL EXTENDED BIALAS-BZDAK AND MINIMAL ODDERON MODELS

Real Extended Bialas-Bzdak: dip fits, Odderon $\geq 7.08 \sigma$

Eur. Phys. J. C (2021) 81:611, published July 2021
<https://doi.org/10.1140/epjc/s10052-021-09381-5>

Observation of Odderon Effects at LHC energies -- A Real Extended Bialas-Bzdak Model Study



Real Extended Bialas-Bzdak Model for Odderon

Observation of Odderon effects at LHC energies: a real extended Bialas-Bzdak model study #4

T. Csorgo (Wigner RCP, Budapest and EKV KRC, Gyongyos), I. Szanyi (Eotvos U. and Wigner RCP, Budapest) (May 28, 2020)

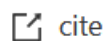
Published in: *Eur.Phys.J.C* 81 (2021) 7, 611 • e-Print: 2005.14319 [hep-ph]



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20 citations

Structure:

Introduction,

Fits with CL > 0.1 % to published pp and pbarp data function

In the dip/bump region (large $-t$ fits)

Linear excitation function

Sanity tests: V

Extrapolations b

Odderon significance f

From combined 1.9

Odderon seen at 7.08 σ

Cross-checks (quadratic trend, ISR data)

The ReBB model and its H(x) scaling version at 8 TeV: Odderon exchange is a certainty #4

I. Szanyi (Eotvos U. and Wigner RCP, Budapest and Karoly Robert U. Coll.), T. Csörgő (Wigner RCP, Budapest and Karoly Robert U. Coll.) (Apr 21, 2022)

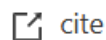
Published in: *Eur.Phys.J.C* 82 (2022) 9, 827, *Eur.Phys.J.C* 82 (2022) 827 • e-Print: 2204.10094 [hep-ph]



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7 citations

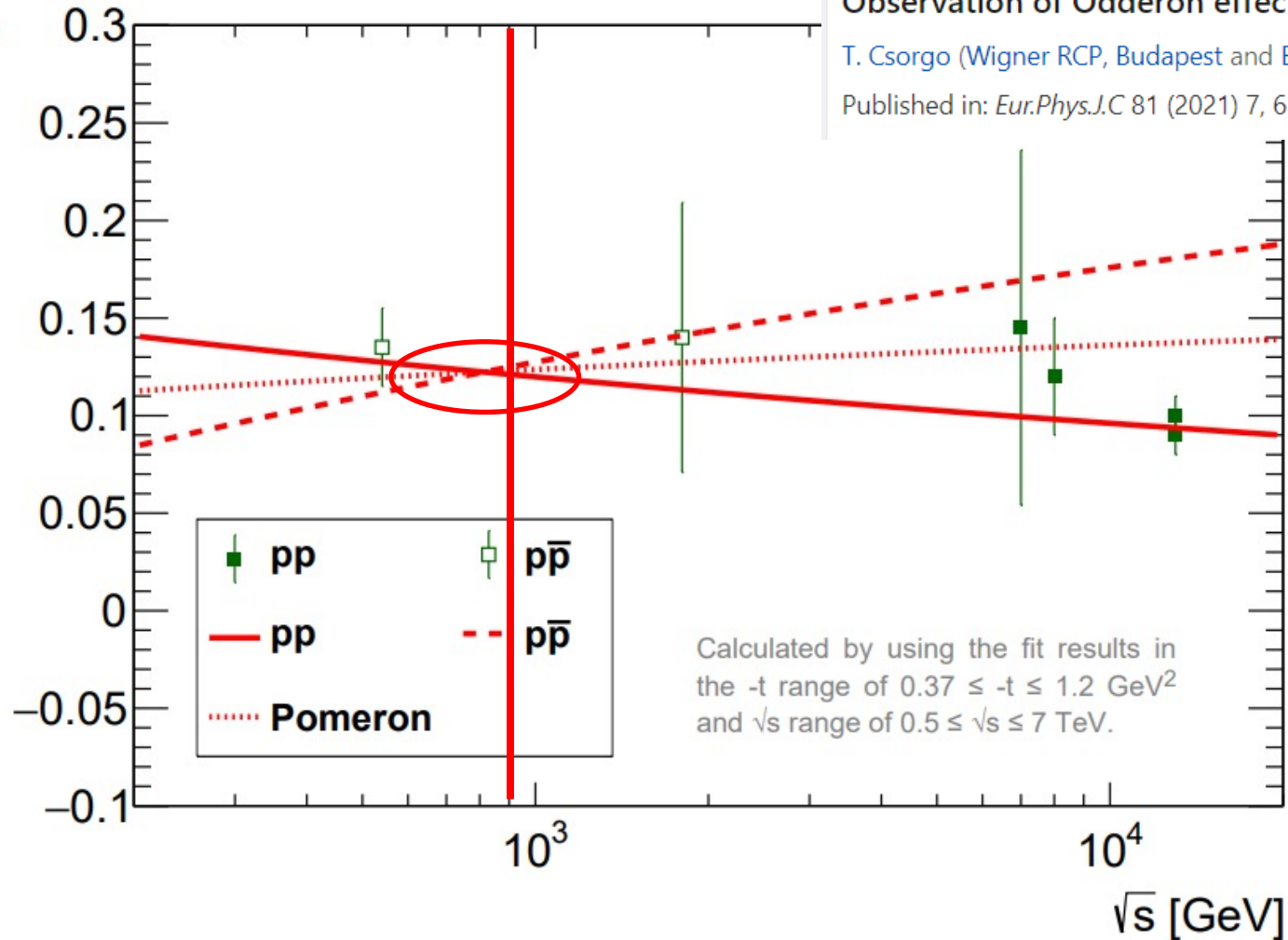
82 pages, 31 figures, model dependent Odderon significance $\geq 7.08 \sigma$,
If 8 TeV TOTEM data are added: significance more than 30σ

ρ_0 from ReBB fits to data – Bonus 2

Observation of Odderon effects at LHC energies: a real extended Bialas–Bzdak model study

T. Csorgo (Wigner RCP, Budapest and EKV KRC, Gyongyos), I. Szanyi (Eotvos U. and Wigner RCP, Budapest) (May 28, 2020)

Published in: *Eur.Phys.J.C* 81 (2021) 7, 611 • e-Print: 2005.14319 [hep-ph]



$$\rho_0^{pp}(s) \neq \rho_0^{p\bar{p}}(s),$$

except at $\sqrt{s} \sim 0.9$ TeV

ReBB fits: R_q, R_d, R_{qd} is the same in pp and pbarp, but $\rho(s)$ is not !
ReBB predicts: at $\sqrt{s} \sim 0.9$ TeV, $\rho(s)$ is the same in pp and pbarb !

Minimal Odderon Model (MOM) and H(x) scaling in ReBB

$$t_{el}(s, b) = i \left[1 - e^{-\Omega(s, b)} \right],$$

$$P(s, b) = 1 - |\exp(-\Omega)|^2 = \tilde{\sigma}_{in}(s, b),$$

$$\text{Re } \Omega(s, b) = -\frac{1}{2} \ln [1 - \tilde{\sigma}_{in}(s, b)].$$

$$\text{Im } \Omega(s, b) = -\alpha \cdot \tilde{\sigma}_{in}(s, b),$$

$$t_{el}(s, b) = i \left(1 - e^{i \alpha \tilde{\sigma}_{in}(s, b)} \sqrt{1 - \tilde{\sigma}_{in}(s, b)} \right).$$

Excitation function of elastic pp scattering from a unitarily extended Bialas–Bzdak model

F. Nemes (CERN and Wigner RCP, Budapest), T. Csörgő (Wigner RCP, Budapest and Unlisted, HU), M. Csanád (Eotvos U.) (Apr 16, 2015)

Published in: *Int.J.Mod.Phys.A* 30 (2015) 14, 1550076 • e-Print: 1505.01415 [hep-ph]

Scaling assumption: $\underline{\sigma}_{in}(\mathbf{s}, \mathbf{b}) = \underline{\sigma}_{in}(\mathbf{b}/\mathbf{R}(\mathbf{s}))$

Result: **H(x) scaling follows,**
but $\sigma_{el}(s)/\sigma_{tot}(s)$ is not a constant of s and $\rho_0(s) \neq 0$.

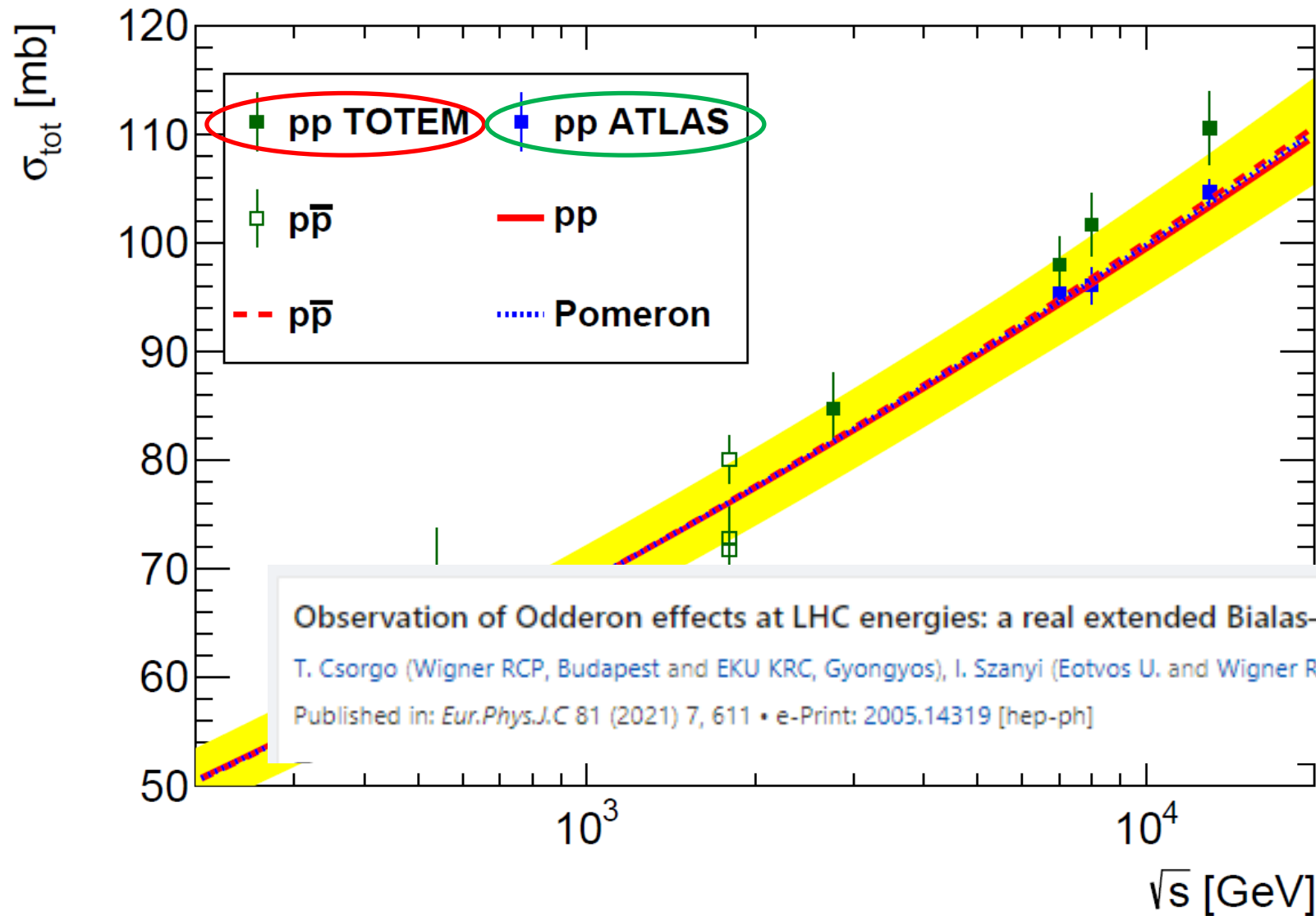
Minimal Odderon Model (MOM):

Odderon exchange signal if $\alpha(s|pp) \neq \alpha(s|p\bar{p})$

Odderon exchange: by only **one s-dependent parameter.**

ReBB is a MOM if $R_q(s) : R_d(s) : R_{qd}(s) = R_q(s_0) : R_d(s_0) : R_{qd}(s_0)$

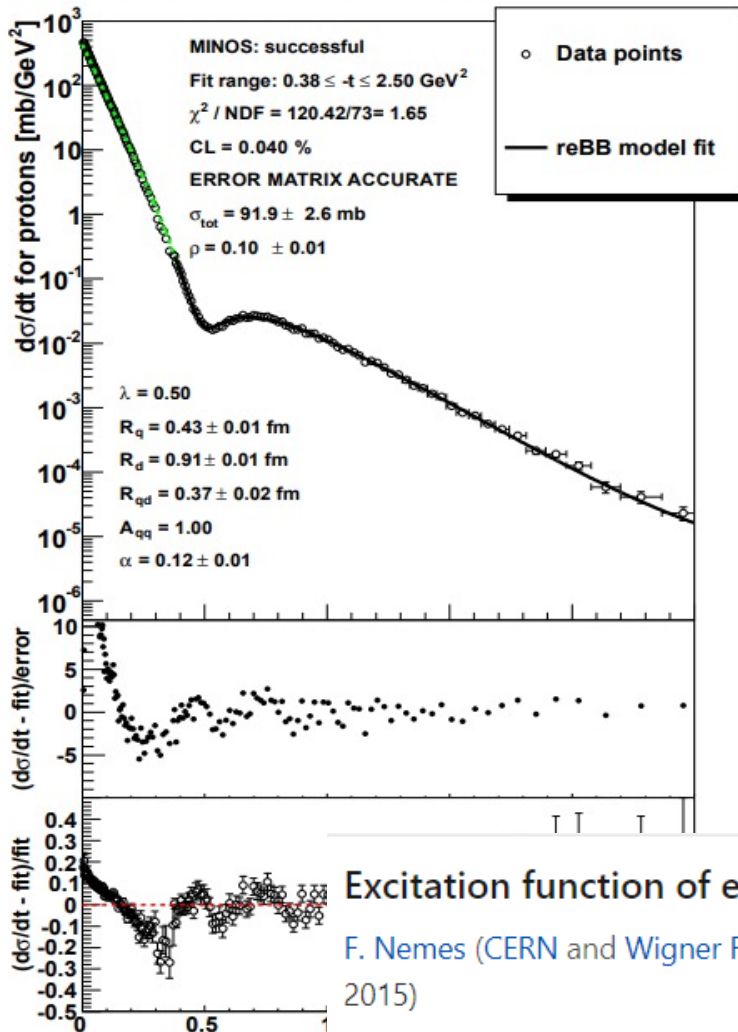
Bonus Extra: ReBB model for σ_{tot} vs TOTEM and ATLAS



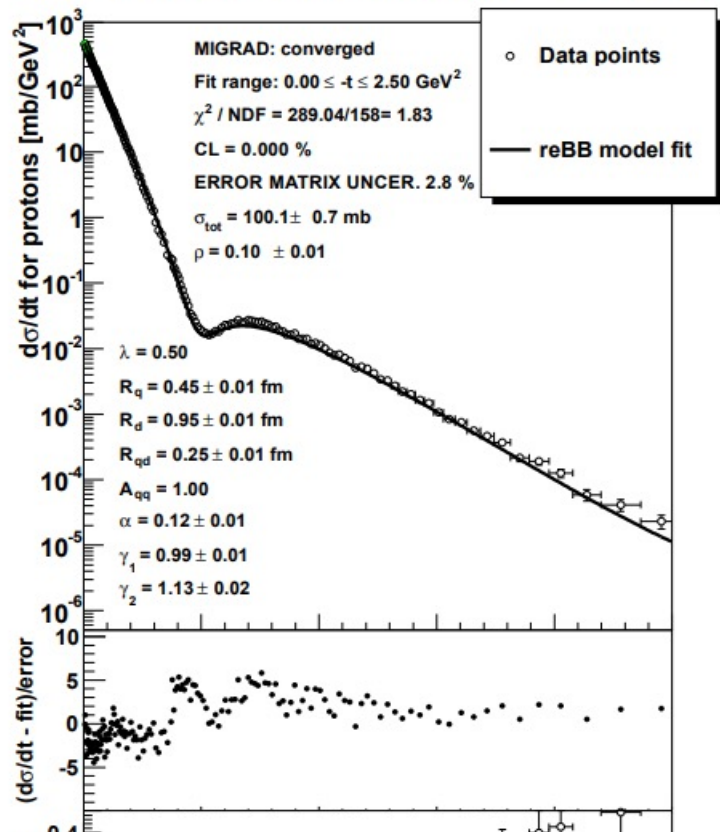
TOTEM data on σ_{tot} are systematically above ReBB result, but ...
ATLAS σ_{tot} data agree with ReBB result, published in EPJ C81 (2021) 7, 611

Statement of the problem, with old χ^2

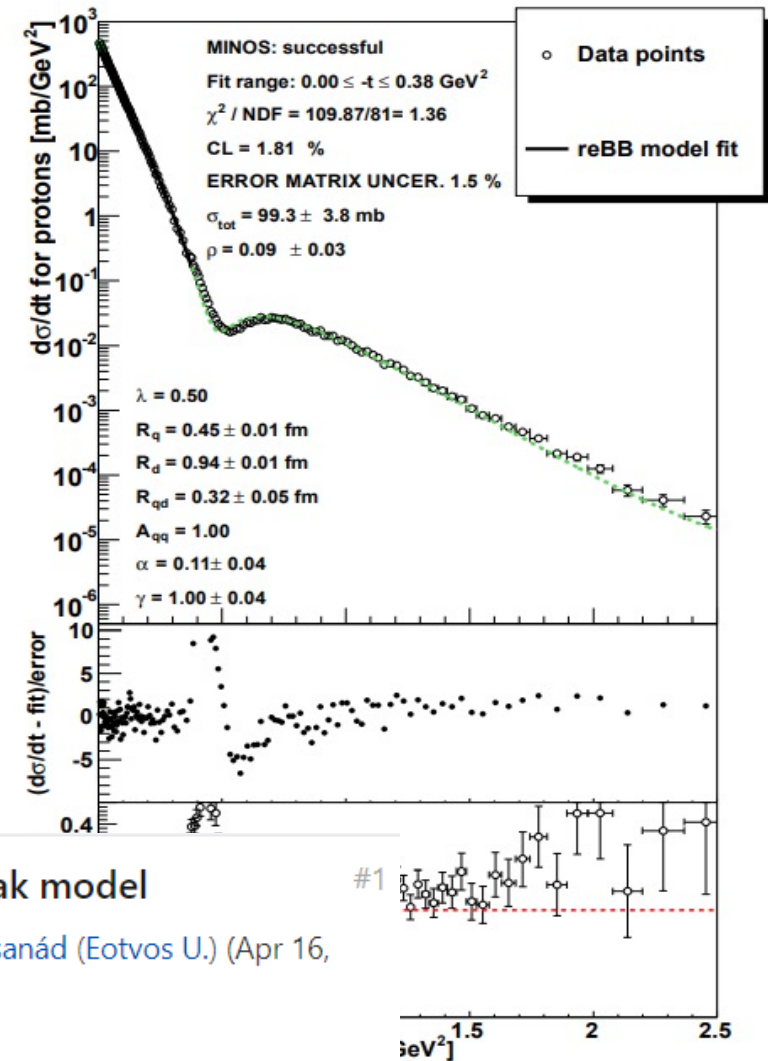
p+p → p+p, diquark as a single entity at $\sqrt{s}=7000.0$ GeV



p+p → p+p, diquark as a single entity at $\sqrt{s}=7000.0$ GeV



p+p → p+p, diquark as a single entity at $\sqrt{s}=7000.0$ GeV



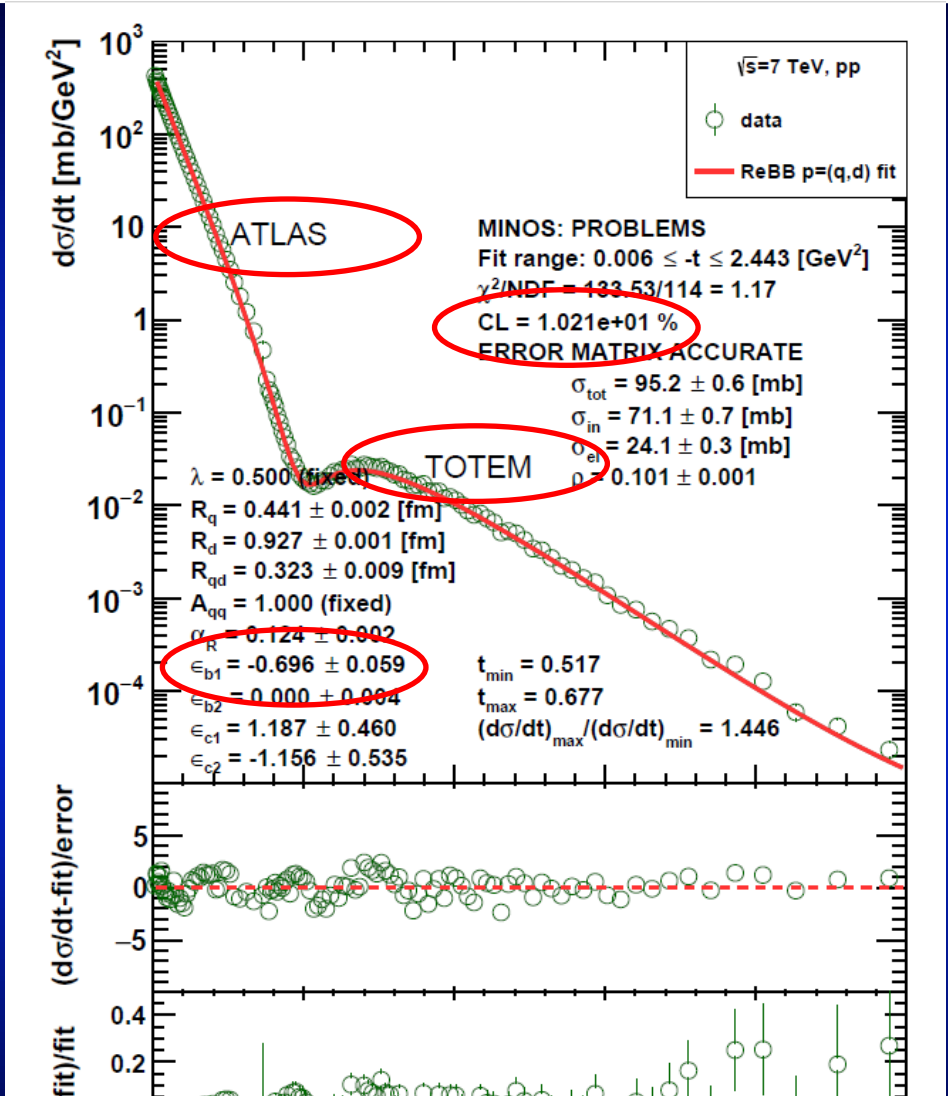
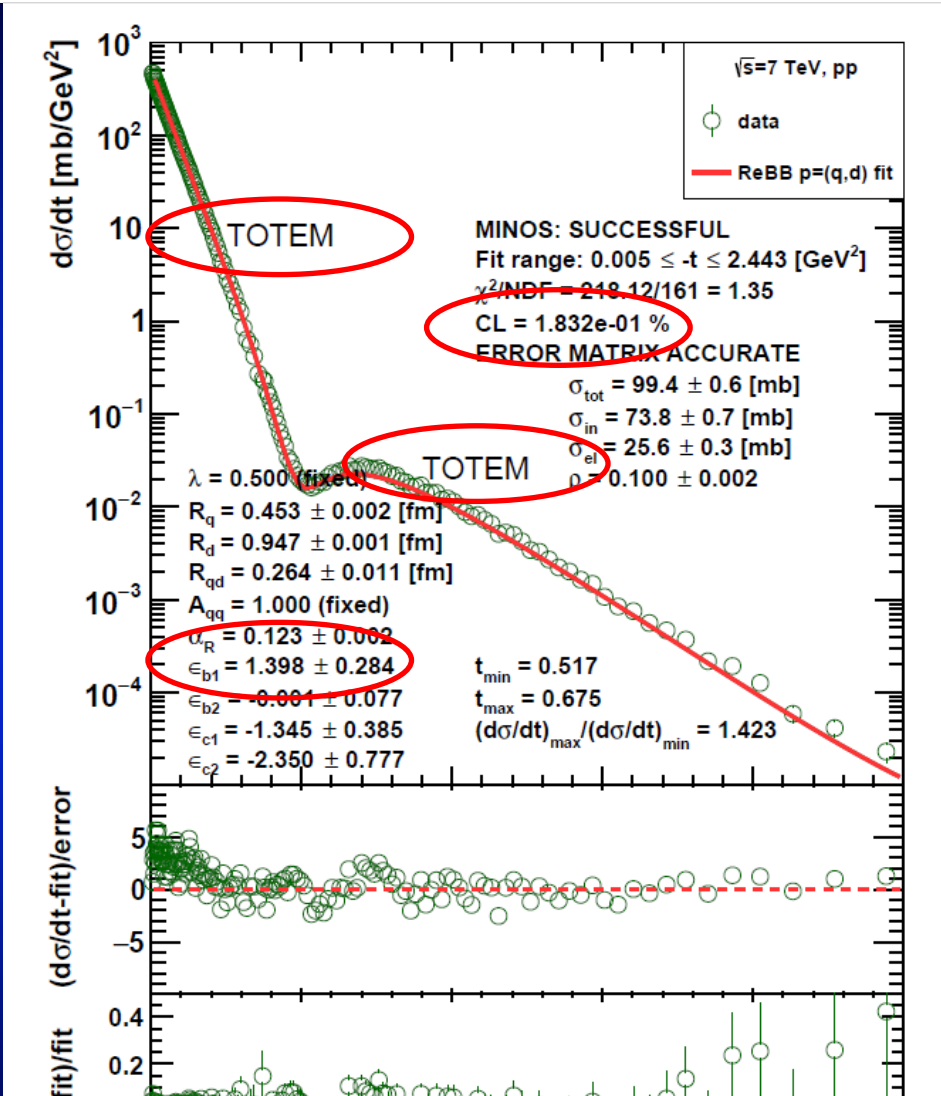
Excitation function of elastic pp scattering from a unitarily extended Bialas–Bzdak model

F. Nemes (CERN and Wigner RCP, Budapest), T. Csörgő (Wigner RCP, Budapest and Unlisted, HU), M. Csanád (Eotvos U.) (Apr 16, 2015)

Published in: *Int.J.Mod.Phys.A* 30 (2015) 14, 1550076 • e-Print: 1505.01415 [hep-ph]

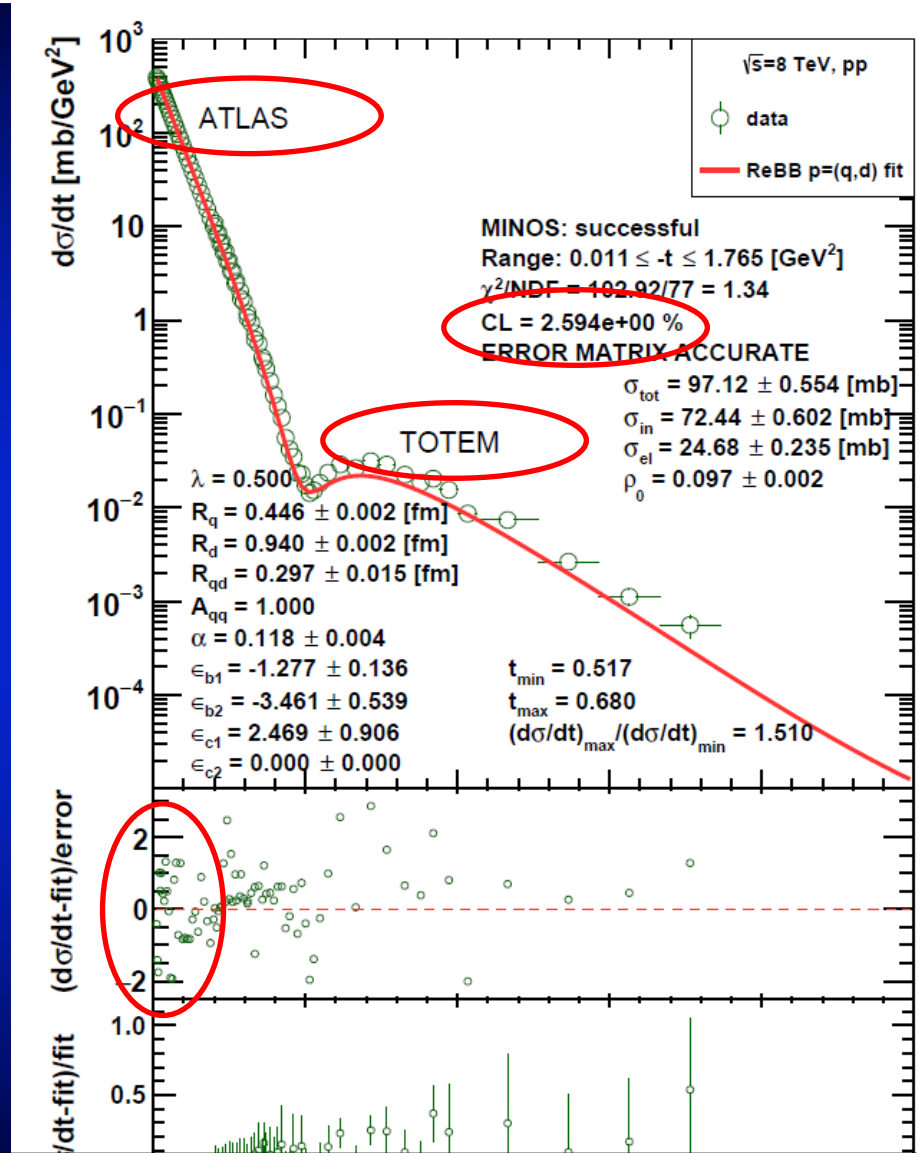
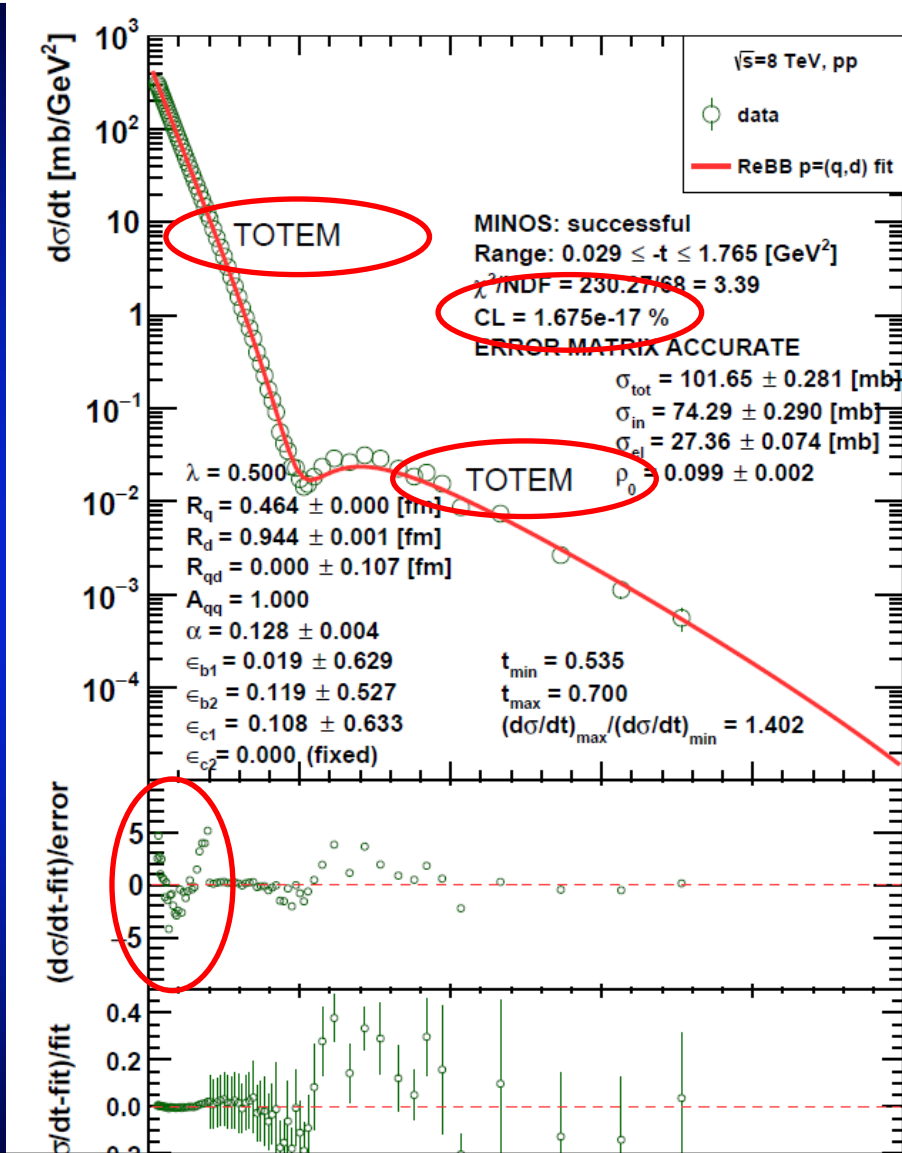
ReBB fits to both TOTEM low-t and TOTEM-large-t fit acceptable at 7 TeV, but The two datasets could not be ReBB fitted simultaneously, without an advanced χ^2 definition !

ATLAS and TOTEM: ReBB model to low $-t$, 7 TeV



TOTEM low- t vs TOTEM-large- t fit acceptable at 7 TeV, obtained with advanced PHENIX χ^2 definition, but with $\epsilon_B > 1$, outside expected range $(-1,1)$...
ATLAS low- t vs TOTEM-large- t fit **successful (CL = 10.2 %) at 7 TeV !**

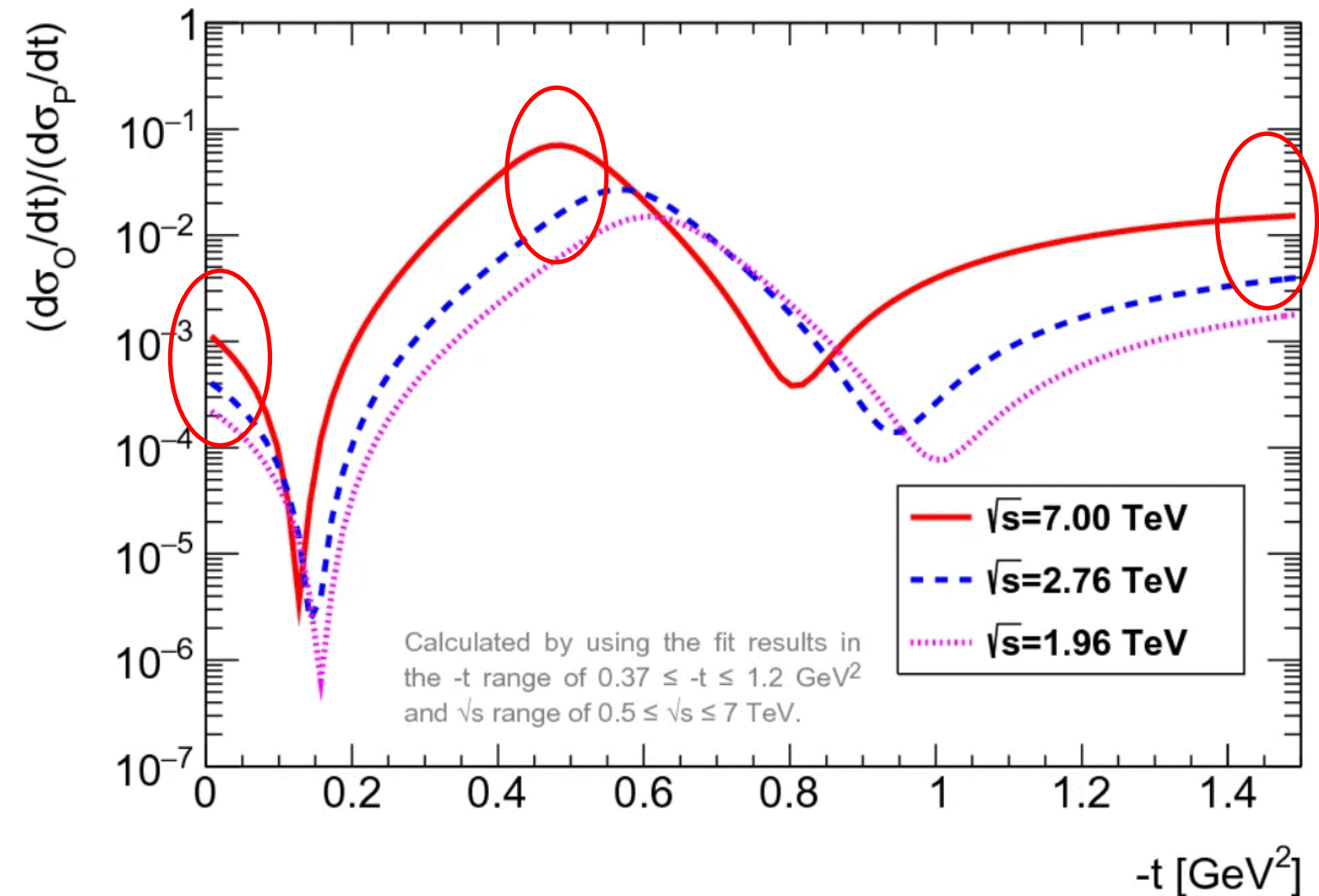
ReBB model extension to low $-t$, 8 TeV



TOTEM low- t vs TOTEM-large- t fit FAILS at 8 TeV, but ...
ATLAS low- t vs TOTEM-large- t fit SUCCESSFUL with CL = 2.59 % at 8 TeV !

ReBB model: where to look for Odderon?

Recent review of Ryskin: asks for Odderon amplitude, intercept etc.
But this has been done (for ReBB) already in 2021!



Current Status of the Odderon

Mikhail G. Ryskin (St. Petersburg, INP) (Aug 4, 2024)

e-Print: 2408.01990 [hep-ph]

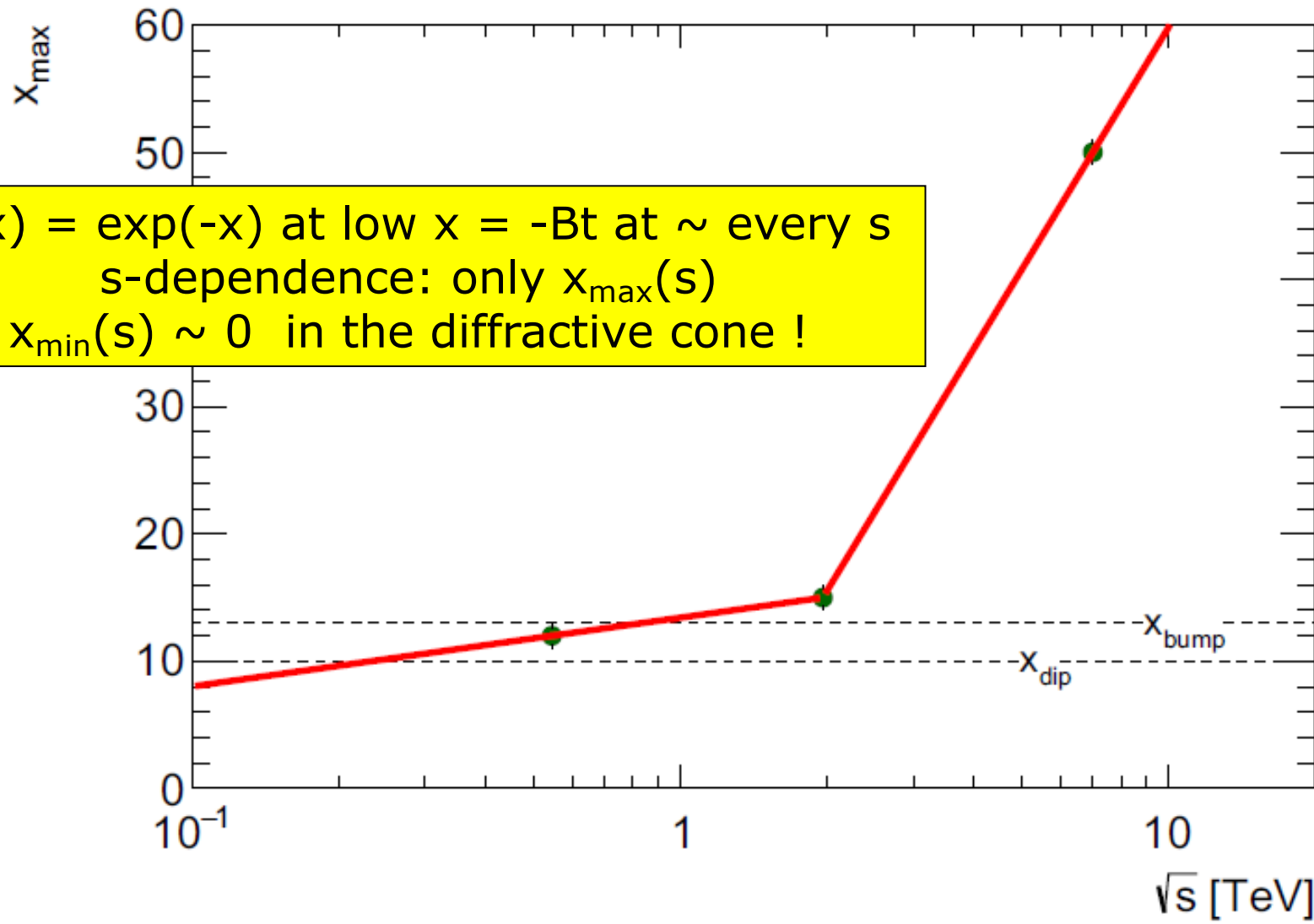
**Modulus and phase
for both Odderon and
Pomeron**

**Extracted from UA4, D0 and
TOTEM data using ReBB fit**

Best @ dip/bump: ~ 10%
Second best @ large $-t$: 1 %
Most difficult @ $t=0$: 0.1%

pp: ReBB model limit on $H(x,s) = H(x)$ scaling

$H(x) = \exp(-x)$ at low $x = -Bt$ at \sim every s
 s -dependence: only $x_{\max}(s)$
 $x_{\min}(s) \sim 0$ in the diffractive cone !



Energy range: 200 GeV – 8 TeV (nearly factor of 40)
With decreasing s , the $x = -Bt$ range for $H(x)$ scaling decreases

AN OLD STORY REVIVED

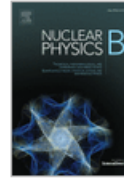
GEOMETRIC SCALING

GEOMETRIC SCALING (GS)



Nuclear Physics B

Volume 59, Issue 1, 13 August 1973, Pages 231-236



Geometric scaling, multiplicity distributions and cross sections

J. Dias De Deus



Nuclear Physics B

Volume 71, Issue 3, 25 March 1974, Pages 481-492

Scaling law for the elastic differential section in pp scattering from geometric scaling ☆

A.J. Buras, J. Dias de Deus

Vol. B6 (1975)

ACTA PHYSICA POLONICA

No 4

GEOMETRICAL SCALING, QUARKS AND THE POMERON

BY J. DIAS DE DEUS

The Niels Bohr Institute, University of Copenhagen

and

Theory Division, Rutherford Laboratory, Chilton*

(Received December 10, 1974)

From quark model additivity applied in the impact parameter plane to the inelastic overlap function we obtain a kind of scale invariant factorizable Pomeron. Quarks themselves are seen as behaving asymptotically like extended objects (quark pancakes), cross-sections and multiplicities being related to their overlap in a high energy collision. Predictions, which can be tested soon at NAL, are given for overlap functions, cross-sections, the ratio σ^{el}/σ^{tot} in the case of various reactions. Universality features of multiplicity distributions are explained in a natural way and an attempt is made to compute the modifications coming from the leading particle effect.

AN OLD STORY REVIVED

Phys. Lett. B 856 (2024) 138960



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Physics Letters B

journal homepage: www.elsevier.com/locate/physletb



Letter

Scaling laws of elastic proton-proton scattering

Cristian Baldenegro^a, Michał Przaszałowicz^{b, *}, Christophe

^a Dipartimento di Fisica, Sapienza Università di Roma, Piazzale Aldo Moro, 2, 00185 Rome, Italy

^b Institute of Theoretical Physics, Faculty of Physics, Astronomy and Applied Computer Science, Jagiello

^c Department of Physics and Astronomy, The University of Kansas, Lawrence, KS 66045, USA

^d Department of Physics, Penn State University, University Park, PA 16802, USA

ARTICLE INFO

Editor: G.F. Giudice

Keywords:

Elastic pp scattering

Geometrical scaling

ABSTRACT

We show that elastic scattering pp square $|t|$ have a universal property the ISR to the LHC, from tens of GeV scaling observed at the ISR with the present experimental uncertainties, scaling laws impose on the paramet



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Physics Letters B

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Letter

Geometric scaling of elastic pp cross section at the LHC

Michał Przaszałowicz^{}

Institute of Theoretical Physics, Faculty of Physics, Astronomy and Applied Computer Science, Jagiellonian University, S. Łojasiewicza 11 30-348 Kraków, Poland

ARTICLE INFO

Editor: Dr Francois Gelis

Keywords:

Elastic pp scattering

geometrical scaling

ABSTRACT

We show that geometric scaling conjectured and observed at the ISR more than 50 years ago, still holds at the LHC. We discuss regularities of the dip - bump structures of the differential elastic cross sections, emphasizing the fact that the ratio of bump to dip positions is constant from the ISR to the LHC. Applying crossing and analyticity we identify imaginary and real parts of the scattering amplitude and compute the ρ parameter and the ratio of bump to dip values of the differential pp cross sections. We also discuss the energy dependence of the total elastic cross section and the violation of geometrical scaling outside the dip - bump region at the LHC.



GS: DEFINITIONS

Geometric scaling

$$\Omega(s, b) = \Omega(b/R(s))$$

Opacity is a function of one variable, and $R(s)$ grows with energy. Changing variable

$$\mathbf{b} \rightarrow \mathbf{B} = \mathbf{b}/R(s)$$

$$\sigma_{\text{inel}} = R^2(s) \underbrace{\int d^2\mathbf{B} \left[1 - |e^{-\Omega(\mathbf{B})}|^2 \right]}_{\text{constant}}$$

Geometric Scaling (GS) assumes $\rho_0(s) \sim 0$

Immediate consequences

$$\sigma_{\text{el}} = R^2(s) \int d^2\mathbf{B} |1 - e^{-\Omega(\mathbf{B})}|^2$$

$$\sigma_{\text{tot}} = 2R^2(s) \int d^2\mathbf{B} \text{Re} [1 - e^{-\Omega(\mathbf{B})}]$$

$$\sigma_{\text{inel}} = R^2(s) \int d^2\mathbf{B} [1 - |e^{-\Omega(\mathbf{B})}|^2]$$

If we neglect χ (indeed ρ parameter is small), then all cross-sections have the same energy dependence.

Does not work at the LHC!

Immediate consequences

$$\sigma_{\text{el}} = \int d^2\mathbf{b} |1 - e^{-\Omega(s,b) + i\chi(s,b)}|^2,$$

$$\sigma_{\text{tot}} = 2 \int d^2\mathbf{b} \text{Re} [1 - e^{-\Omega(s,b) + i\chi(s,b)}],$$

$$\sigma_{\text{inel}} = \int d^2\mathbf{b} [1 - |e^{-\Omega(s,b)}|^2].$$

Motivating issue (Cs.T.): Opacity is, in general complex.

Why not scaling for the imaginary part of the optical opacity (here denoted by χ) too?

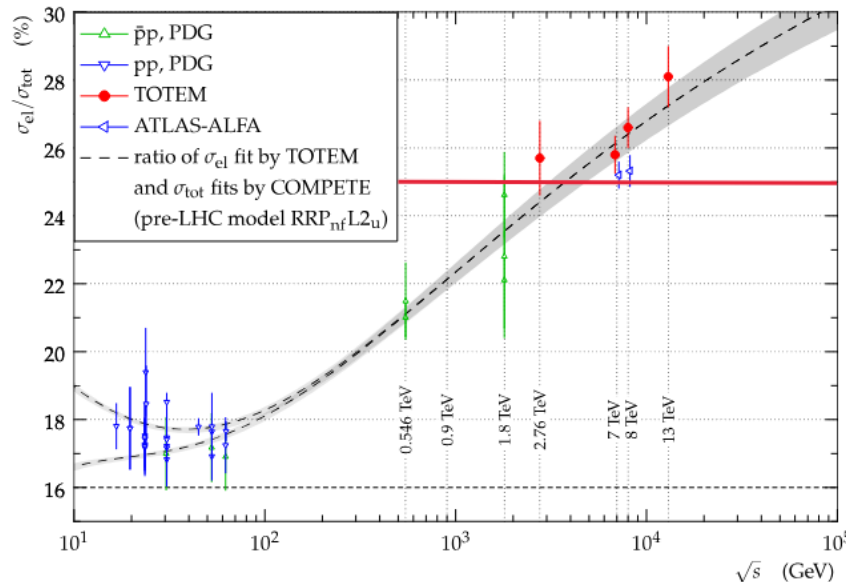
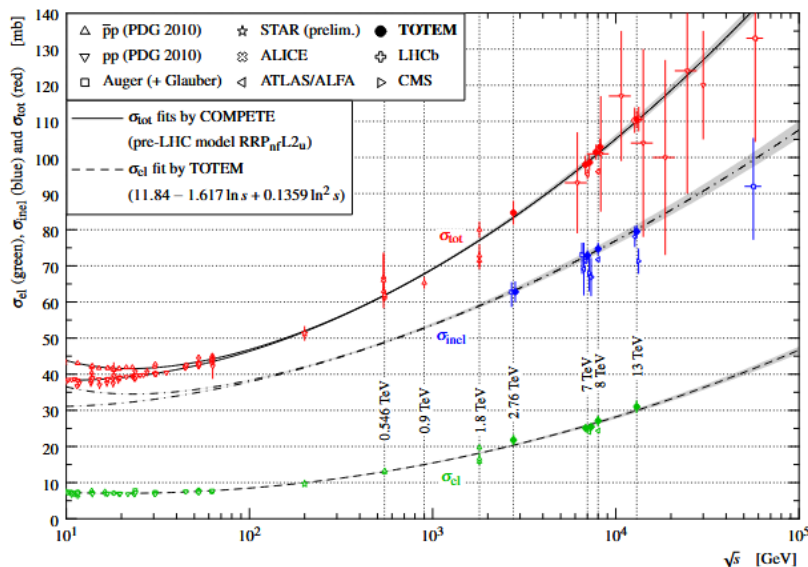
M. Praszalowicz:
PLB **869** (2025) 1398-48,

M. Praszalowicz:
Talk at Moriond QCD 2025,

GS: WHY DOES NOT WORK AT LHC?

Table 1. Summary of cross-section results of TOTEM at $\sqrt{s} = 2.76$ and 13 TeV in pp collisions [1, 3, 4, 51].

\sqrt{s} (TeV)	σ_{tot} (mb)	σ_{el} (mb)	σ_{in} (mb)	σ_{el}/σ_{tot} (%)
2.76	84.7 ± 3.3	21.8 ± 1.4	62.8 ± 2.9	25.7 ± 1.1
13.0	110.6 ± 3.4	31.0 ± 1.7	79.5 ± 1.8	28.1 ± 0.9



consequences

$$|1 - e^{-\Omega(B)}|^2$$

$$3 \operatorname{Re} [1 - e^{-\Omega(B)}]$$

$$[1 - |e^{-\Omega(B)}|^2]$$

ρ parameter is small),
have the same energy

$\sigma_{el}(s)/\sigma_{tot}(s)$ is not a constant of s at LHC energies:
Table 1 and Fig. 1 from T. Csörgő for TOTEM, Proc. ISMD 2018

Does not work at the LHC!

Recent Results from the CERN LHC Experiment TOTEM -- Implications for Odderon Exchange

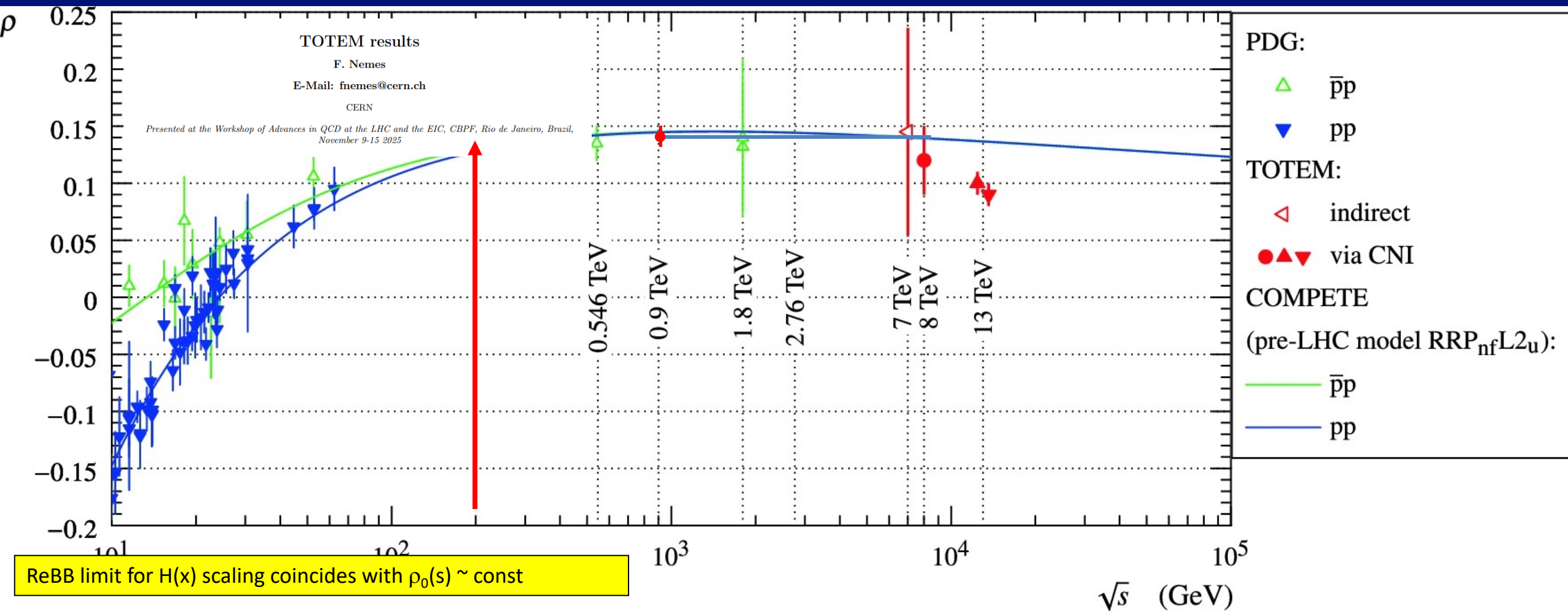
TOTEM Collaboration • T. Csörgő (Wigner RCP, Budapest and Eszterhazy Karoly U., Eger and CERN) for the collaboration. (Mar 16, 2019)

Published in: EPJ Web Conf. 206 (2019) 06004 • Contribution to: ISMD 2018, ISMD 2018 • e-Print: 1903.06992 [hep-ex]

M. Praszalowicz:

PLB 869 (2025) 1398-48,

GS: Does NOT work at LHC, because neither $\rho(s)$ is 0, nor $\sigma_{el}(s)/\sigma_{tot}(s)$ is constant



Recent Results from the CERN LHC Experiment TOTEM -- Implications for Odderon Exchange

TOTEM Collaboration • T. Csörgő (Wigner RCP, Budapest and Eszterhazy Karoly U., Eger and CERN) for the collaboration. (Mar 16, 2019)

Published in: *EPJ Web Conf.* 206 (2019) 06004 • Contribution to: ISMD 2018, ISMD 2018 • e-Print: 1903.06992 [hep-ex]

DERIVATION 1

$H(x)$ SCALING FOR SMALL $x \ll 1$

H(x) scaling for small-x, in the cone region (x ≪ 1)

$$\frac{d\sigma}{dt} = A(s) \exp [B(s)t]$$

$$A(s) = B(s) \sigma_{\text{el}}(s) = \frac{1 + \rho_0^2(s)}{16\pi} \sigma_{\text{tot}}^2(s),$$

$$\frac{1}{B(s)\sigma_{\text{el}}(s)} \frac{d\sigma}{dt} = \exp [tB(s)]$$

$$H(x) \equiv \frac{1}{B(s)\sigma_{\text{el}}(s)} \frac{d\sigma}{dt},$$
$$x = -tB(s).$$

Advantages:

- 1) $H(x) \approx \exp(-x)$ in the diffractive cone, for $x \ll 1$
- 2) $H(0) = 1 \rightarrow$ Start from a place that you know
- 3) Measurable both for pp and pbarp

BONUS EXTRA: An Odderon signal

$H(x|pp)$ scales with energy, but for pbarp, no scaling

DERIVATION 2

$H(x)$ SCALING FOR ALL x

Derivation of $H(x)$ scaling for all x

$$t_{el}(s, \mathbf{b}) = (i + \rho_0) r(s) E(\tilde{\mathbf{x}}).$$

$$\text{Re exp} [-\Omega(s, b)] = 1 - r(s) E(\tilde{\mathbf{x}}),$$

$$\text{Im exp} [-\Omega(s, b)] = \rho_0 r(s) E(\tilde{\mathbf{x}}),$$

$$\tilde{\mathbf{x}} = \mathbf{b}/R(s),$$

$$R(s) = \sqrt{B(s)},$$

Similarly to GS, only one
s-dependent scale $R(s)$

Dissimilarly to GS,
 $\rho_0(s) = \rho_0(s_0) = \rho_0$
But it is not vanishing!

$$\frac{d\sigma}{dt} = \frac{1}{4\pi} |T_{el}(\Delta)|^2 = \frac{1 + \rho_0^2}{4\pi} r^2(s) R^2(s) |\tilde{E}(R(s)\Delta)|^2$$

$$A = \left. \frac{d\sigma}{dt} \right|_{t=0} = \frac{1 + \rho_0^2}{4\pi} r^2(s) R^2(s) |\tilde{E}(0)|^2,$$

$$\frac{1}{A} \frac{d\sigma}{dt} = \frac{|\tilde{E}(\sqrt{x})|^2}{|\tilde{E}(x=0)|^2} = H(x),$$

Dissimilarly to GS, no
singularity at minimum!

Advantages:

$H(x) \neq \exp(-x)$ arbitrary positive def. in the dip-bump region
Measurable both for pp and p-antip. Normalized as $H(0) = 1$.

Model independent $H(x|pp)$ scaling yields Odderon $> 6.26 \sigma$

Evidence of Odderon-exchange from scaling properties of elastic scattering at TeV energies #5

T. Csörgő (Wigner RCP, Budapest and CERN), J. Nyak (Unlisted, HU), R. Pasechnik (Lund U., Dept. Theor. Phys.), A. Ster (Wigner RCP, Budapest), J. Szanyi (Wigner RCP, Budapest) (Dec 26, 2019)

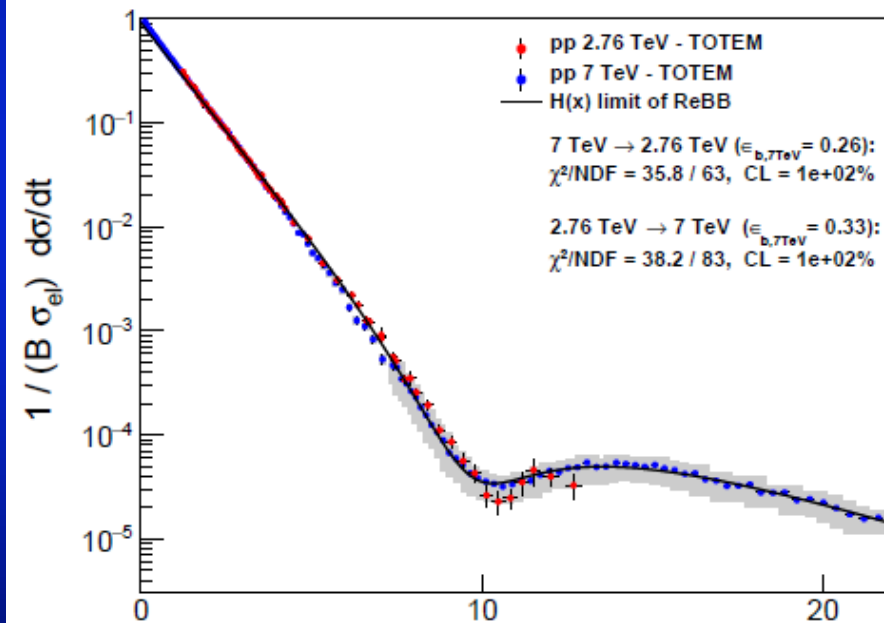
Published in: *Eur.Phys.J.C* 81 (2021) 2, 180 • e-Print: 1912.11968 [hep

Eur. Phys. J. C (2021) 81: 180, published February 2021

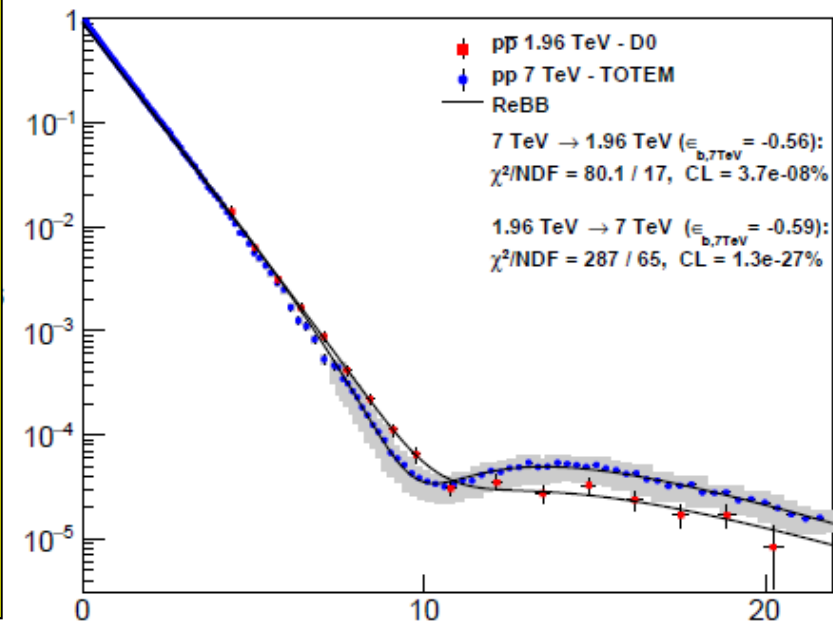
<https://doi.org/10.1140/epjc/s10052-021-08867-6>

pdf DOI cite

15 citations



$H(x) = 1/(B \sigma_{el}) d\sigma/dt \text{ vs } x = -Bt$



$B \equiv B_0(s)$ from now on

$-Bt$

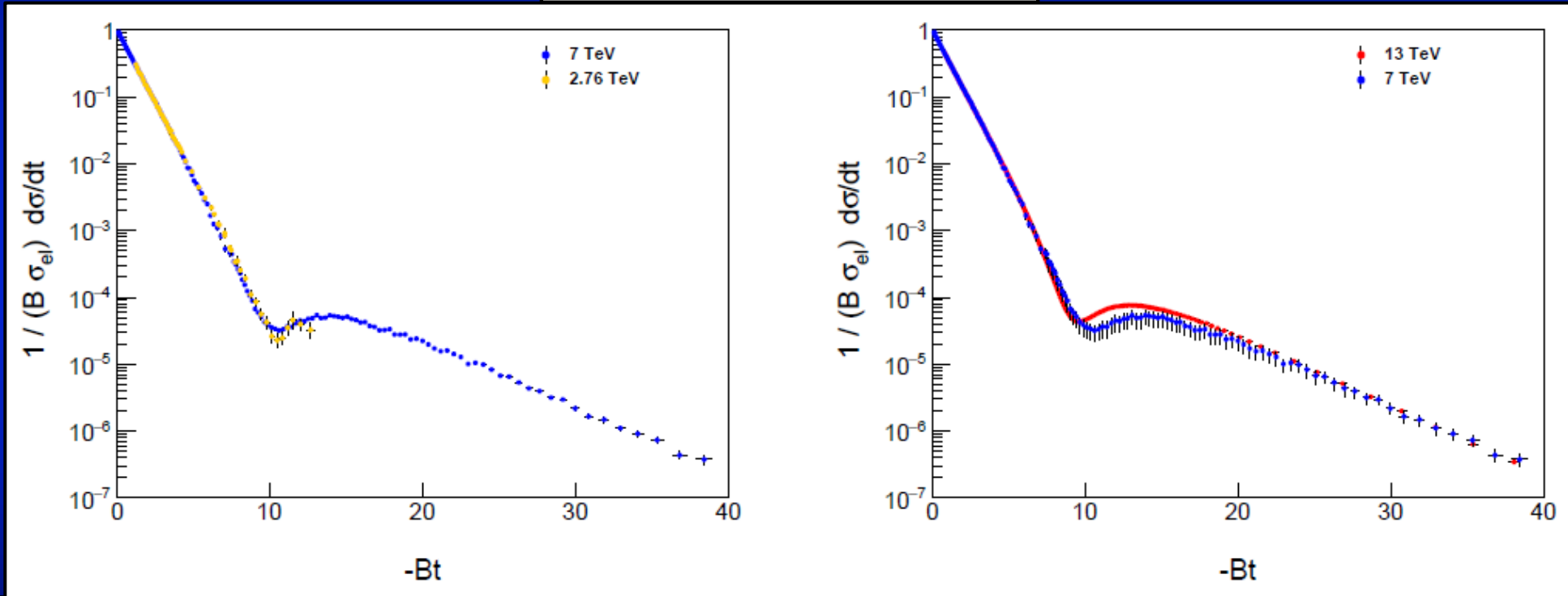
$x = -Bt = -B_0(s)t$

$-Bt$

S: Model independent Odderon significance $\geq 6.26 \sigma$
C1: All D0 and TOTEM published data at 1.96, 2.76 and 7.0 TeV
C2: domain of validity is still determined model dependently.

Test of the $H(x)$ scaling at 7 vs 2.76 TeV

$$H(x) \equiv \frac{1}{B(s)\sigma_{\text{el}}(s)} \frac{d\sigma}{dt},$$
$$x = -tB(s).$$



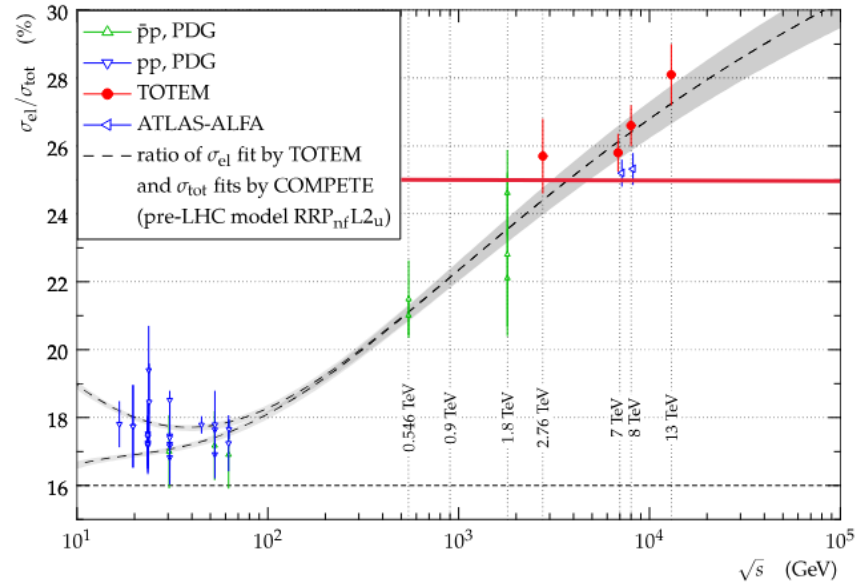
Valid between 2.76 and 7 TeV, even with stat errors only,

$H(x)$ scaling valid even in the bump/tail region!

Between 8 and 13 TeV, scaling limited to the cone,

as $\rho(s)$ starts to decrease, so $H(x|pp)$ scaling **violated** beyond stat+syst errors in dip/bump!

H(x) vs GS: Summary and conclusions

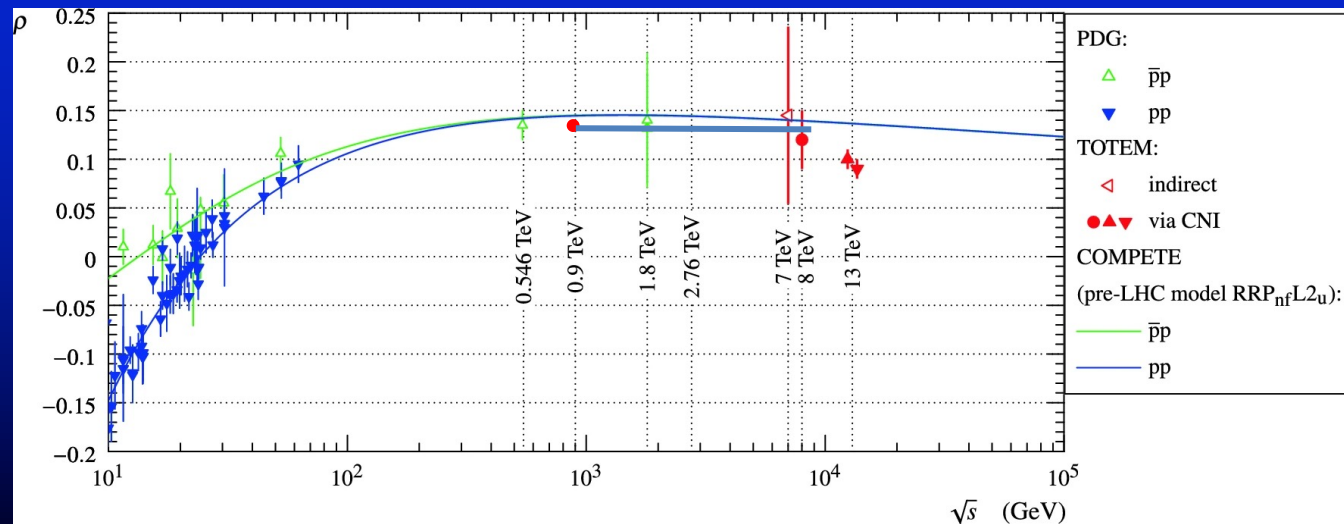


Geometric scaling is revived but does NOT work at LHC,

because at LHC, $\sigma_{el}(s)/\sigma_{tot}(s)$ is not a constant of s and $\rho_0(s) \neq 0$

But H(x) is valid if $\rho_0(s) = \text{const}$

→ reason for H(x) violations at 13 TeV



If $\rho_0(s) = 0$, geometric scaling is recovered and $\sigma_{el}(s)/\sigma_{tot}(s)$ becomes $\text{const}(s)$

Thus H(x) scaling is a generalized geometric scaling !
Foundation of Odderon discovery extended to $0.2 \leq \sqrt{s} \leq 8$ TeV for H(x) scaling of the ReBB model. TOTEM/ATLAS data support this.

Summary and conclusions

Real Extended Bialas Model is beautiful



because it is a Minimal Odderon Model
it has a $H(x)$ scaling limit
describes elastic pp and pbarp data
from 23 GeV to 8 TeV

It predicted successfully ATLAS $\sigma_{\text{tot}}(s)$
and combined ATLAS low-t
+ TOTEL large-t $d\sigma/dt$ data

Foundation of Odderon discovery
extended to $0.2 \leq \sqrt{s} \leq 8$ TeV.

Within the ReBB model, Odderon
exchange is a certainty!

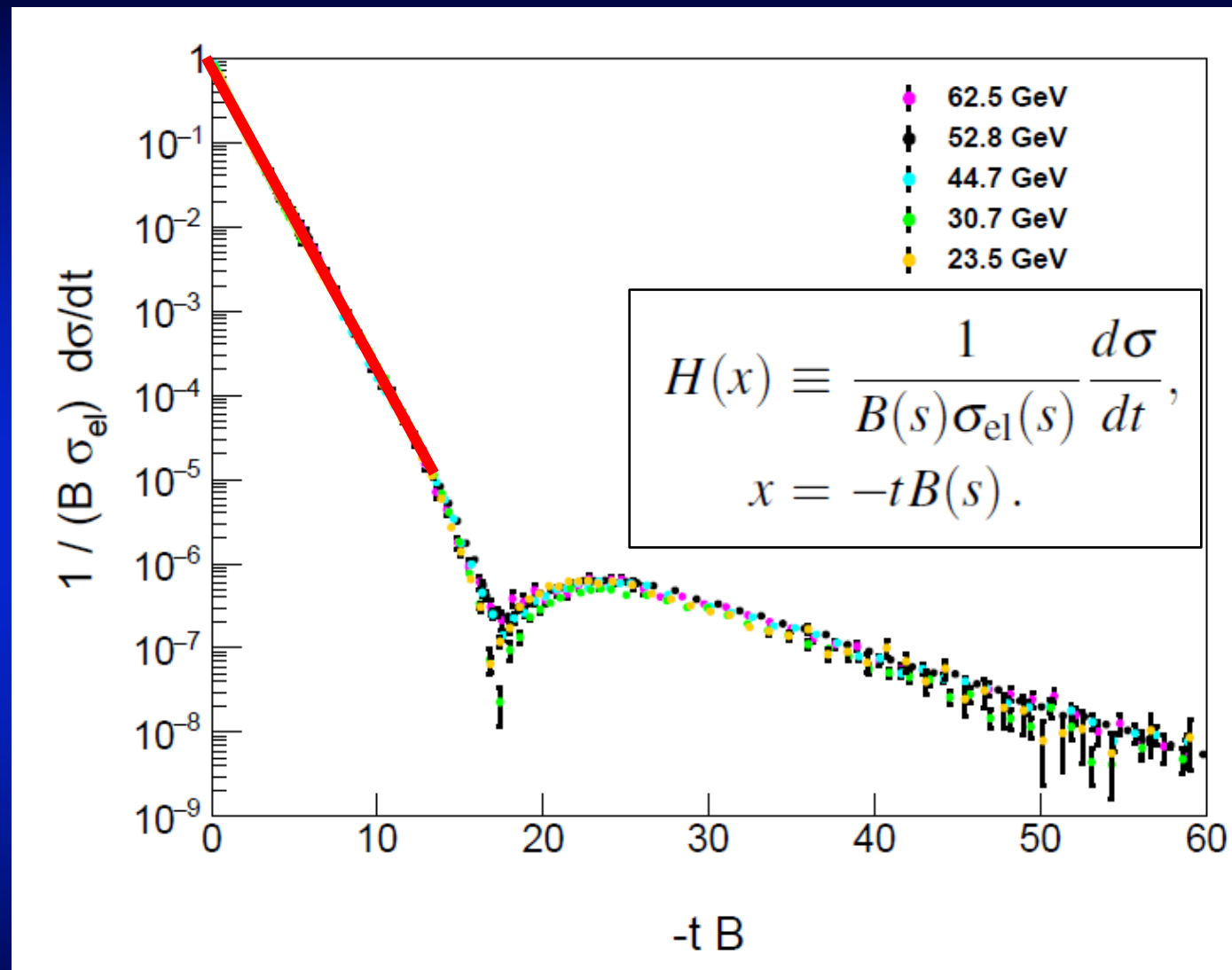
Dear Andrzej,
Happy 90!

THANK YOU FOR YOUR ATTENTION !



BACKUP SLIDE

Test of the $H(x)$ scaling at ISR

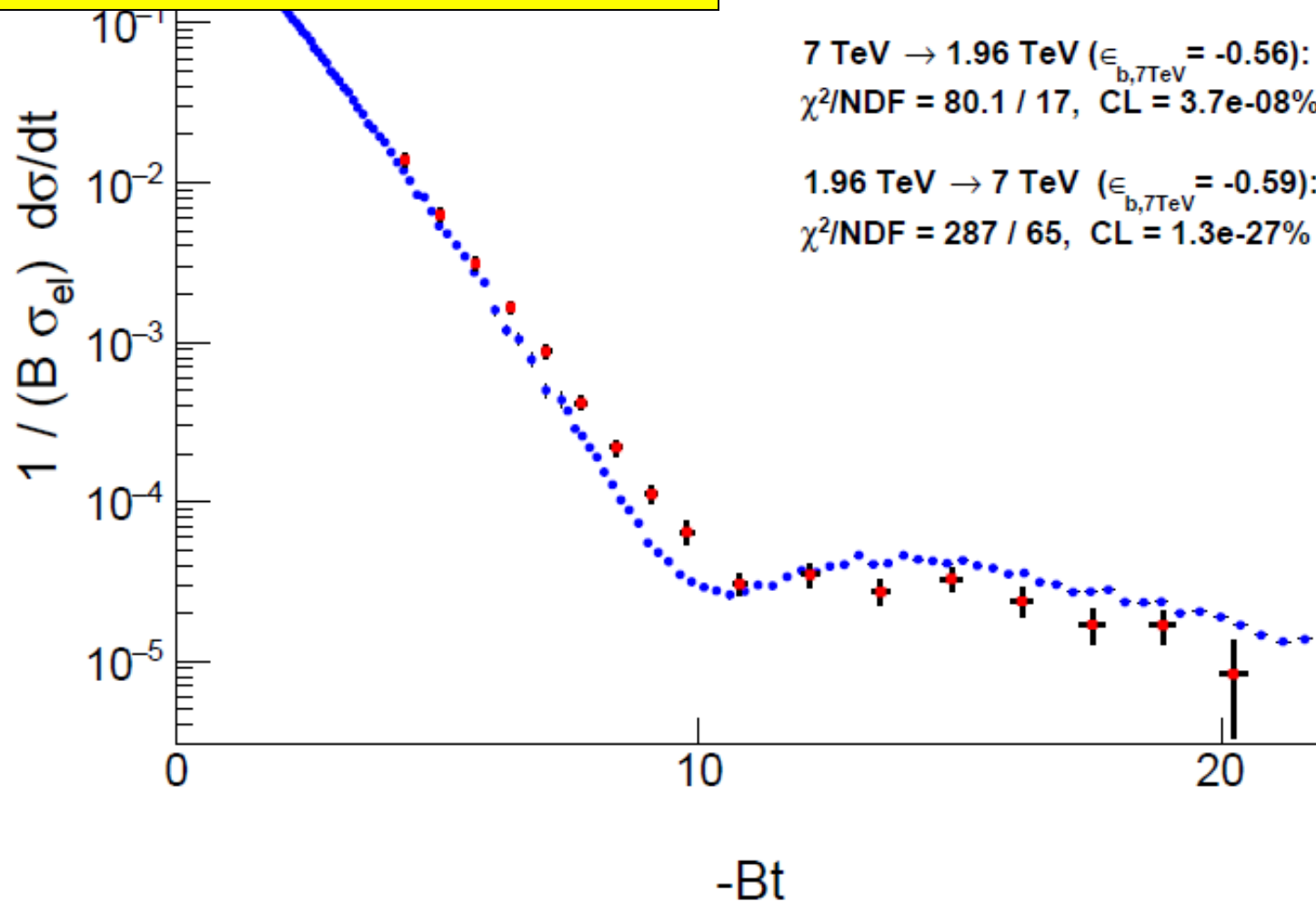


Hungarian-Swedish team:
Eur. Phys. J. C (2021) **81**: 180,

$H(x) = \exp(-x)$ in the cone
Works better than expected, up to $x \sim 10$, even in the bump/tail, $x > 25$ region!

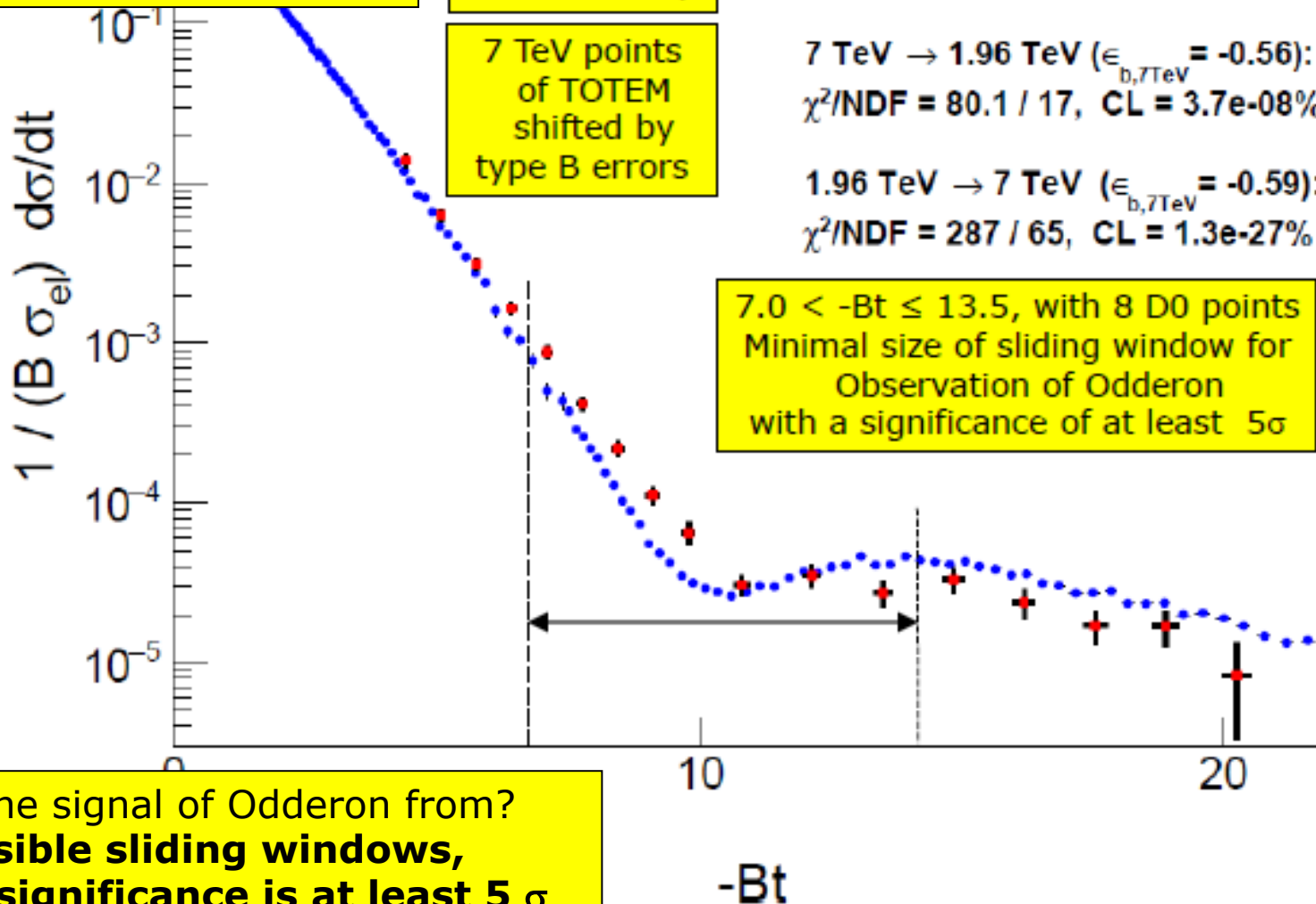
OBSERVATION OF ODDERON

7 TeV data shifted
by $\epsilon_{B7,TeV}$ to minimize χ^2
Type A errors are shown only
Both swing and dip regions important!



SLIDING WINDOW for 5σ

Model independent results:
only datapoints,
without s-dependent
extrapolations !



Where is the signal of Odderon from?
All possible sliding windows,
where the significance is at least 5σ

Is $H(x,s) = H(x)$ at 1.96 TeV?

MODEL INDEPENDENTLY: YES!

In the background of the Odderon signal,
defined as $x \leq 7.0$ in union with $x > 13.5$

$H(x|pp, 7 \text{ TeV}) \sim H(x|pbarp, 1.96 \text{ TeV})$

Agreement: a significance of 2.39σ

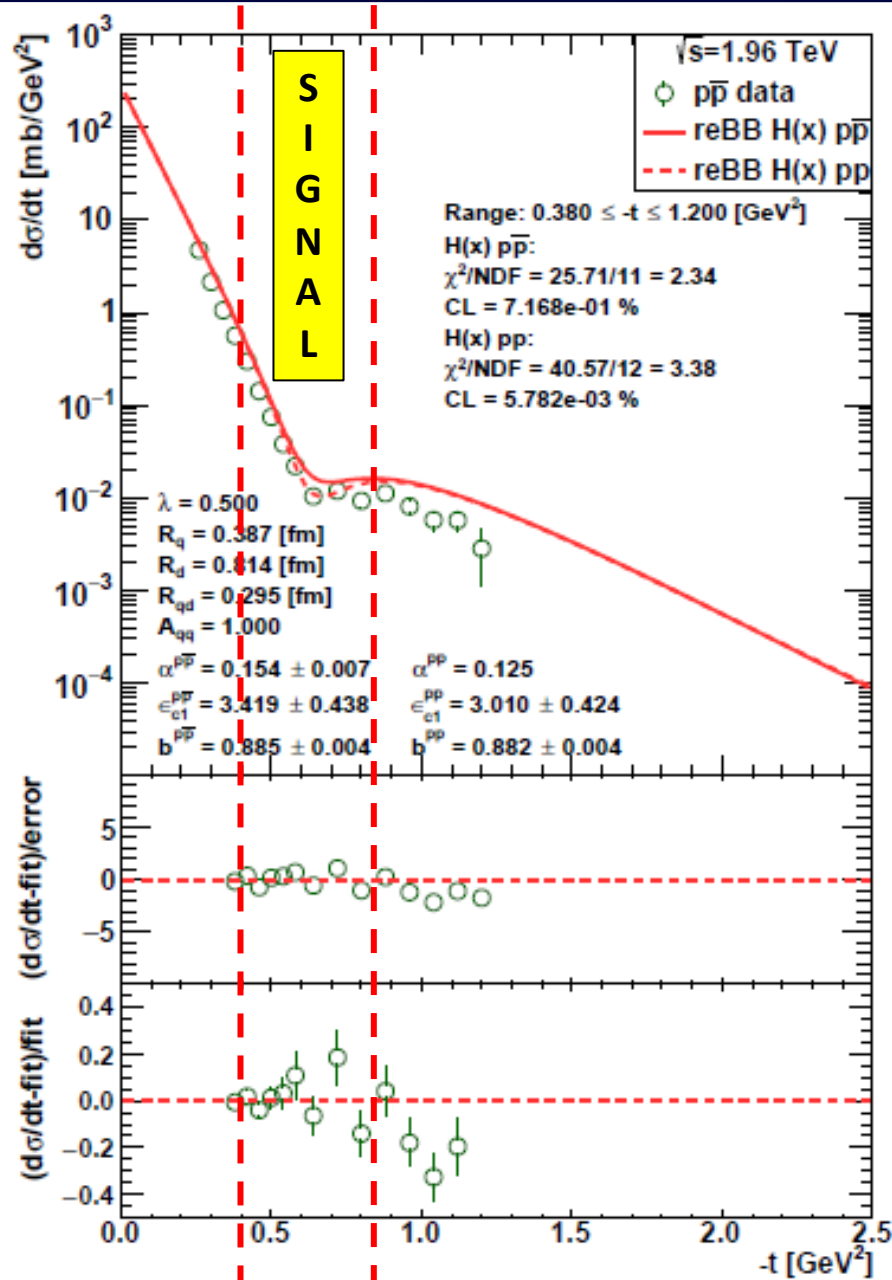
Results for the background: $x \leq 7.0$ in union with $x > 13.5$

for $\epsilon_{B21}(7 \text{ TeV}) = -1.1$ that minimizes signal in the background

x_{\max}	ϵ_{B21} of $\min[\chi^2(\text{background})]$	$\Delta\chi^2(\text{background})$	NDF(background)	σ (background)
20.2	-1.10	20.20	9	2.39

New MODEL INDEPENDENT result

Is $H(x,s) = H(x)$ at 1.96 TeV?



MODEL DEPENDENTLY: Yes
1.96 TeV

Highest energy where p+antip data are available

H(x) scaling limit:
in the Bialas-Bzdak model

Fits pbarp data up to largest -t
(red line, dashed line: pp)

Pull plots:
(data-fit)/error
(data-fit)/fit

$t_{\text{max}}(1.96 \text{ TeV}, pp) > 1.2 \text{ GeV}^2$

$\rightarrow x_{\text{max}}(1.96 \text{ TeV}, pp) > 20$

Safely above the 5 σ threshold

Role of the H(x) scaling violations
Do they decrease the signal or not?

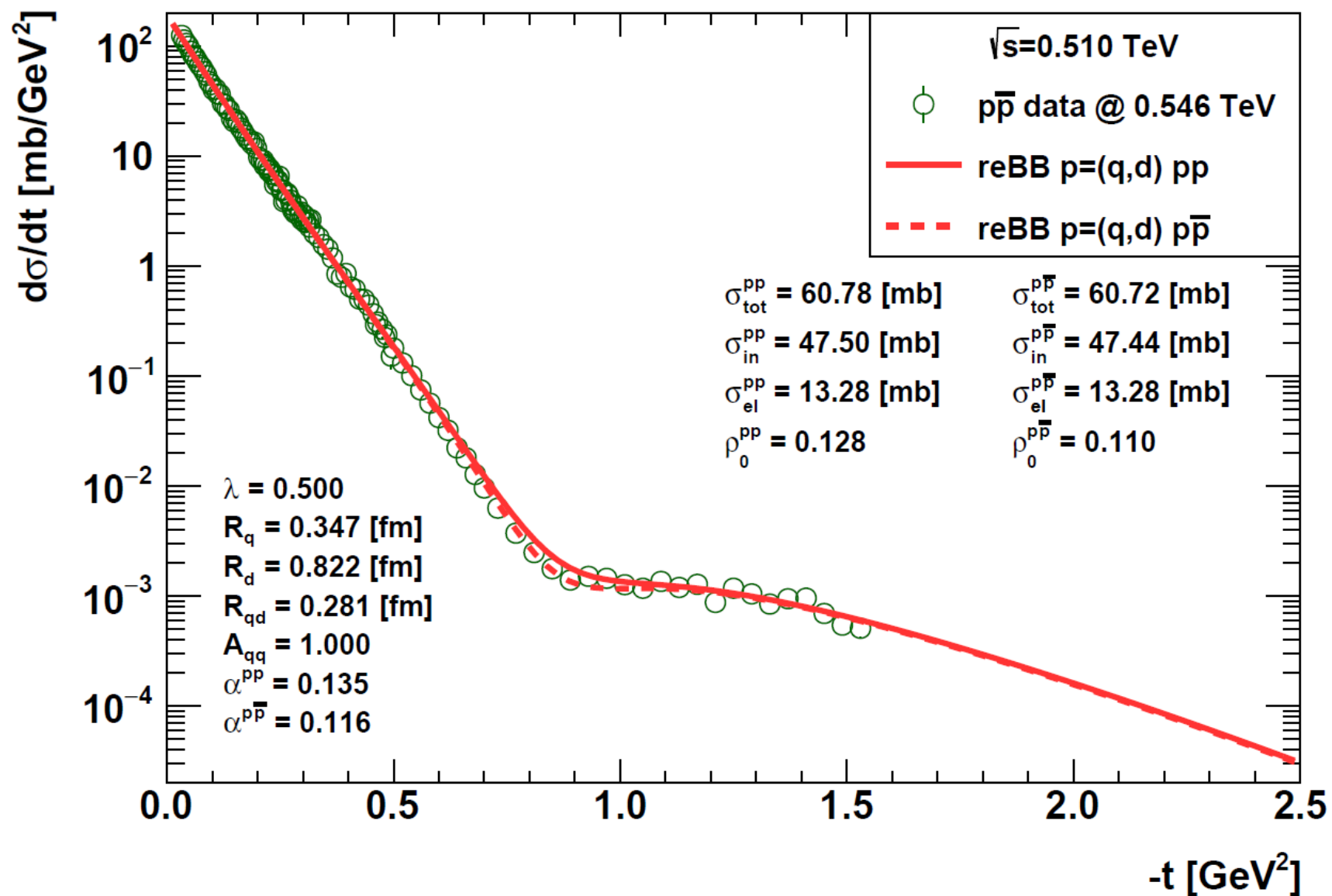
\sqrt{s} (TeV)	χ^2	NDF (ReBB)	σ (ReBB)
1.96	24.28	13	2.19
2.76	100.35	20	7.12
1.96 and 2.76	124.63	33	7.08

H(x) scaling: allows to project pp data ONLY
Scaling violations decrease significance at 1.96 TeV
BUT
Also allow to evaluate pbarp data at 2.76 TeV

Trade-off effect!

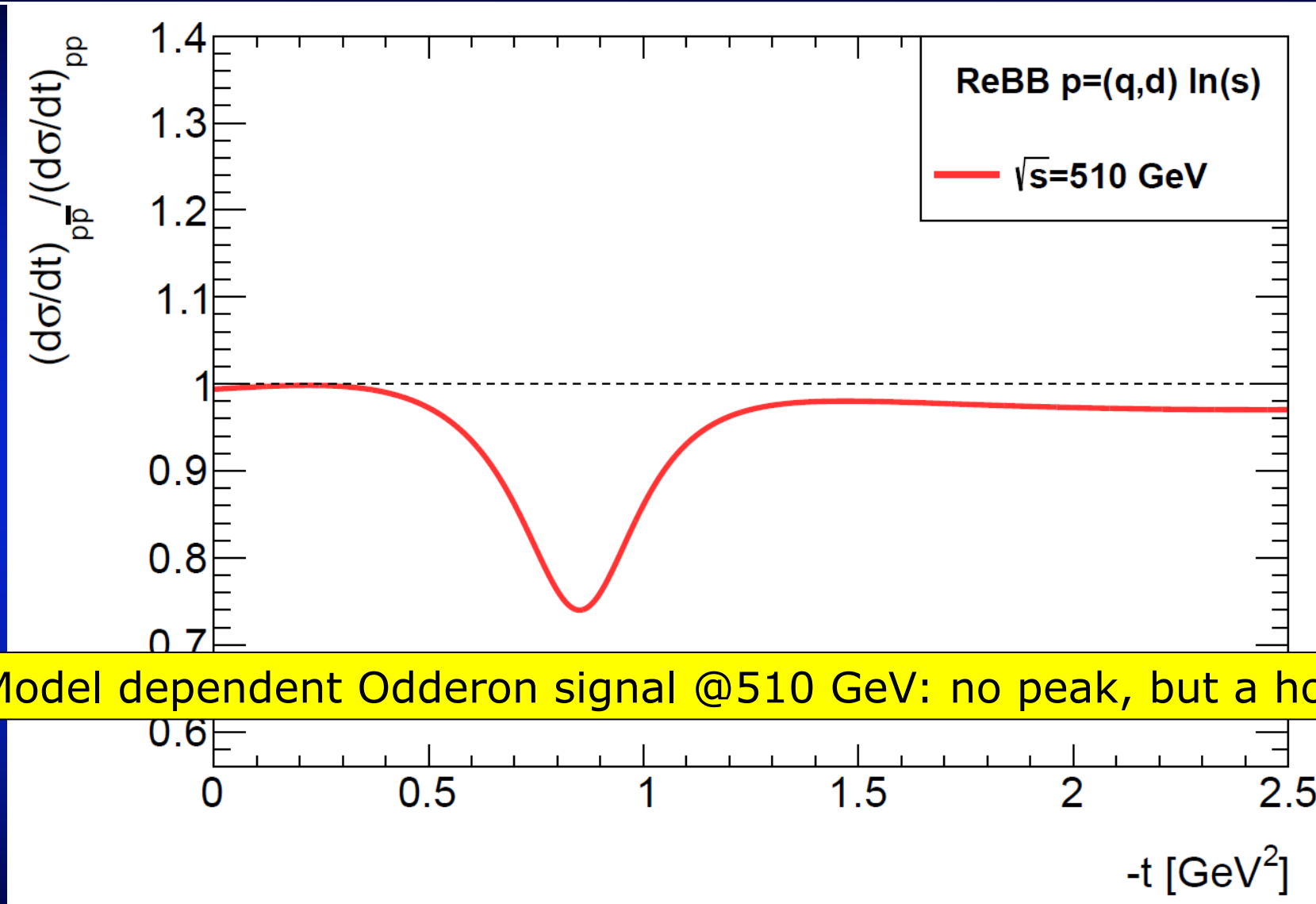
**Odderon significance increases
from 6.26 to 7.08 σ .**

Predictions for pp and pbarp dσ/dt @ 510 GeV



Model dependent Odderon signal @510 GeV: pp above pbarp !

Ratio of pbarp to pp cross-sections @ 510 GeV

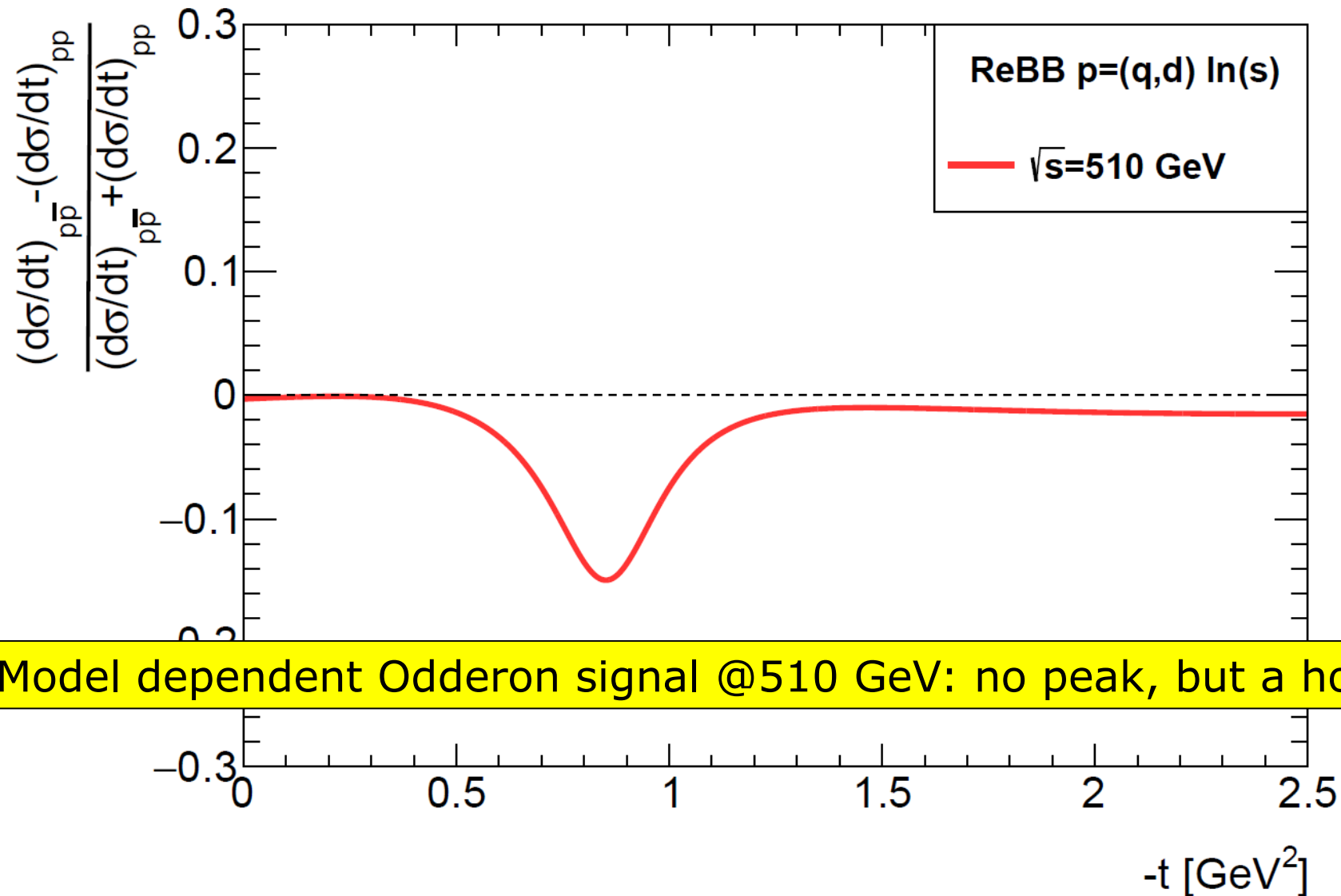


Model dependent Odderon signal @510 GeV: no peak, but a hole

Scaling violations: dominant @510 GeV

Model dependent Odderon signal: pbarp $d\sigma/dt \sim 25\%$ below pp !!

Asymmetry parameter @ 510 GeV



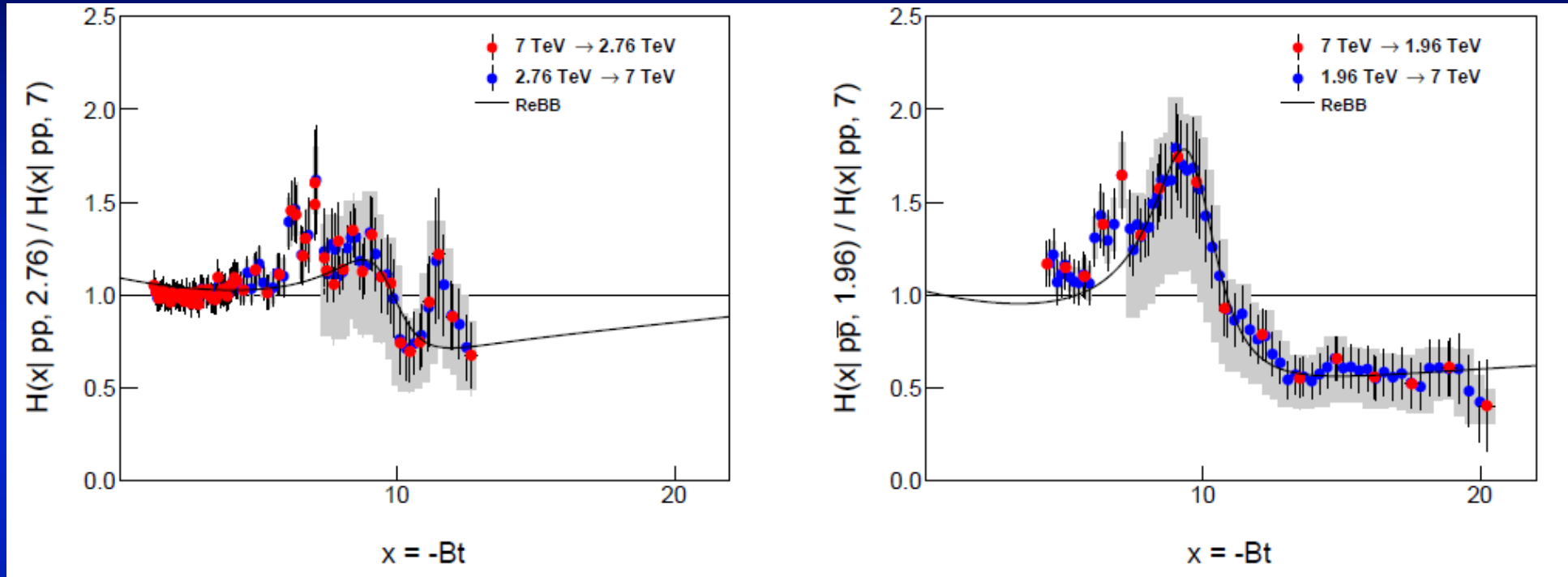
Model dependent Odderon signal @510 GeV: no peak, but a hole

Scaling violations: dominant @510 GeV

Model dependent Odderon signal: $A \sim -15\%$ @ $-t \sim 0.85$ GeV²

SUMMARY: AT LEAST 6.26 σ ODDERON

An at least 6.26 σ Odderon effect



A discovery level, **model independent** Odderon effect at TeV scale.

Published: Eur. Phys. J. C **81**, 180 (2021).

<https://doi.org/10.1140/epjc/s10052-021-08867-6>

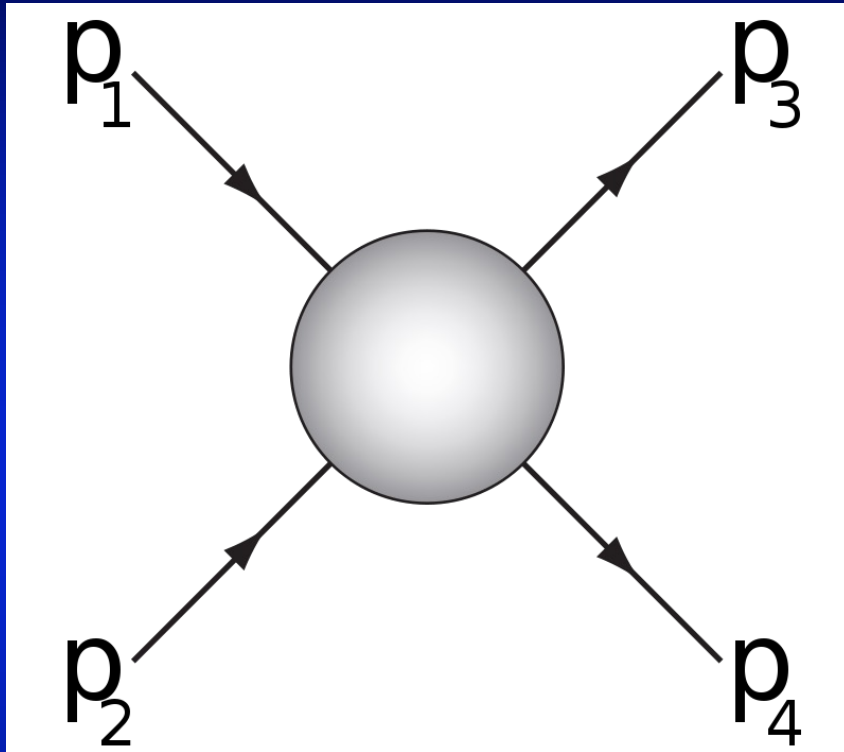
Domain of validity of $H(x)$ scaling: full $x = -tB$ range of D0 at 1.96 TeV.

Published result confirmed with NEW, model INDEPENDENT result !

Model dependent results, using the ReBB model

Significance $\geq 7.08 \sigma$, see e-Print: [2005.14319](https://arxiv.org/abs/2005.14319) [hep-ph]

Mandelstam variables



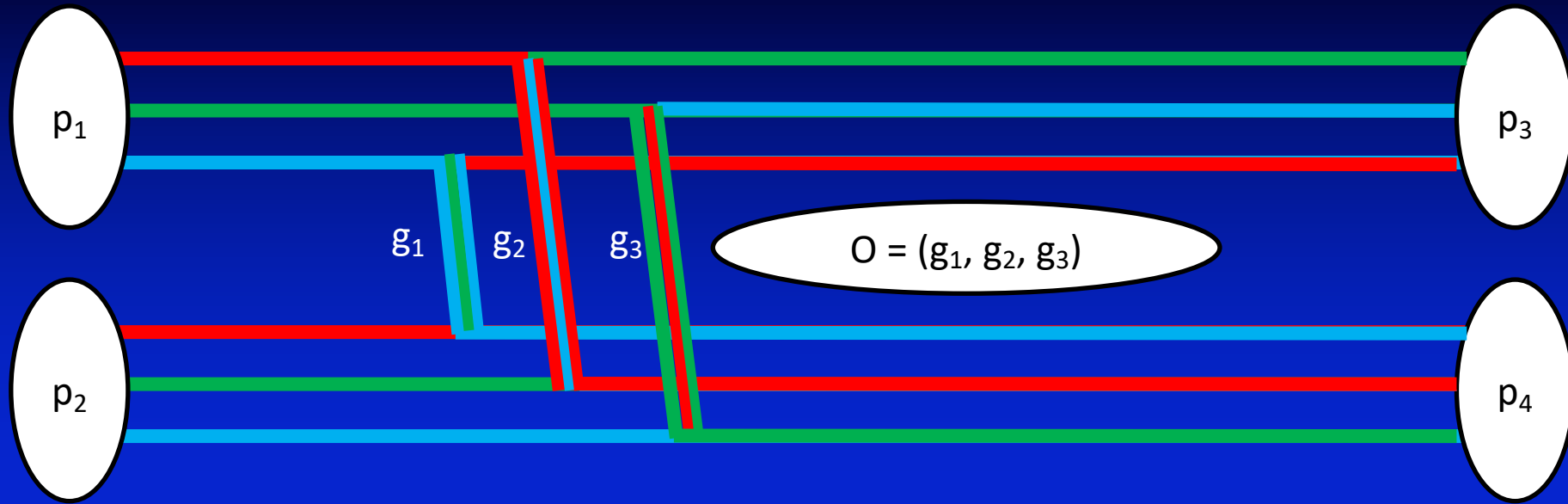
p_1, p_2 : four-momenta
before elastic scattering

p_3, p_4 : four-momenta
after elastic scattering

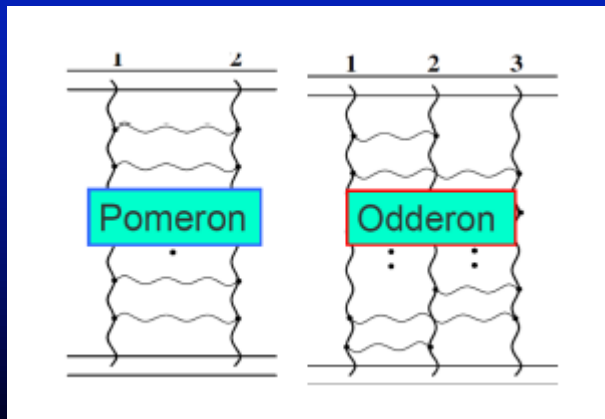
$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2$$
$$t = (p_1 - p_3)^2 = (p_4 - p_2)^2$$
$$u = (p_1 - p_4)^2 = (p_3 - p_2)^2$$

s : square of the cms energy
 t : square of four-momentum
transfer

Odderon and QCD in Laymen's Terms



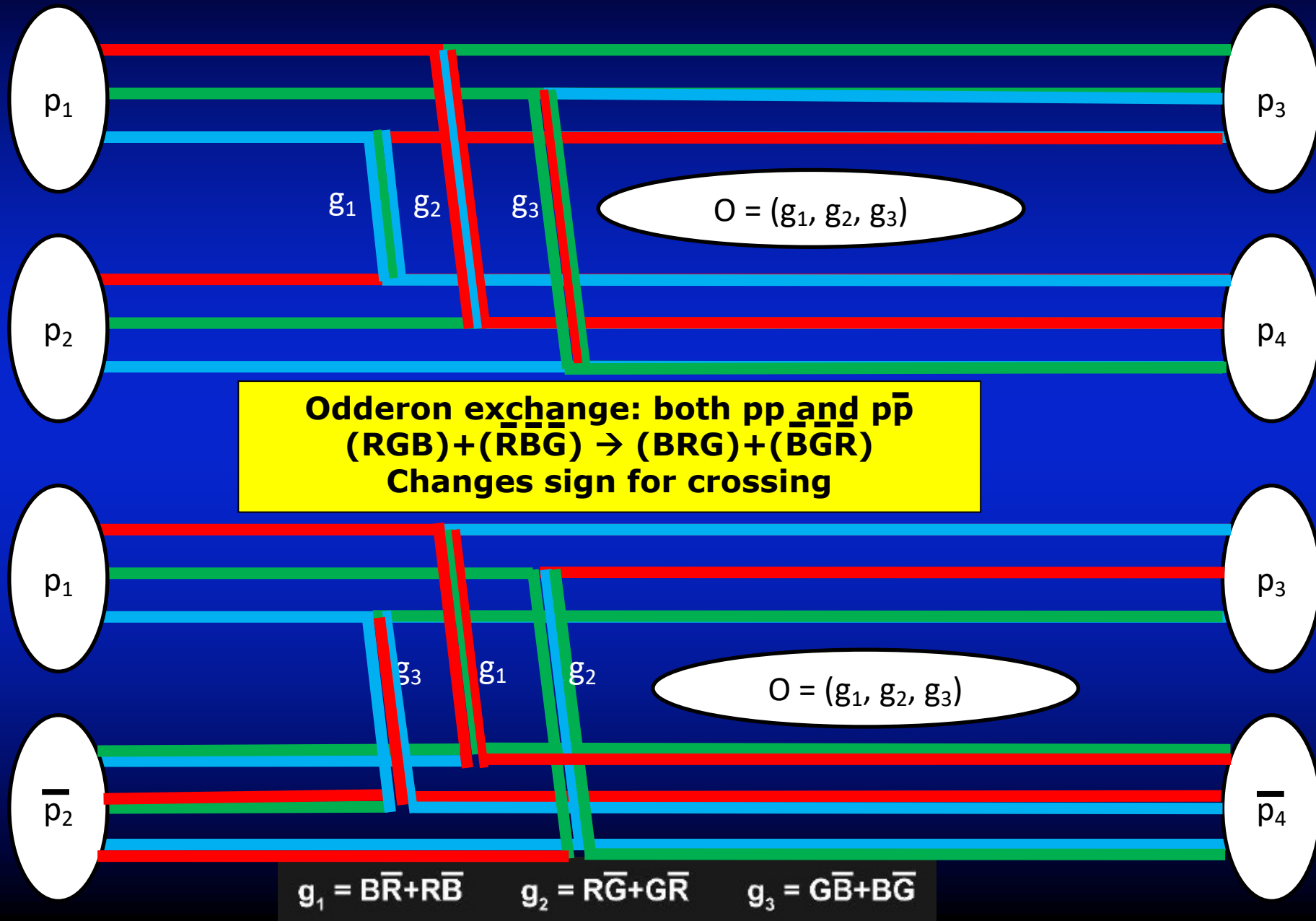
$$g_1 = B\bar{R} + R\bar{B} \quad g_2 = R\bar{G} + G\bar{R} \quad g_3 = G\bar{B} + B\bar{G}$$



Pomeron (2+4+...) gluon in pp:
 $(RGB) + (RGB) \rightarrow (GRB) + (GRB)$

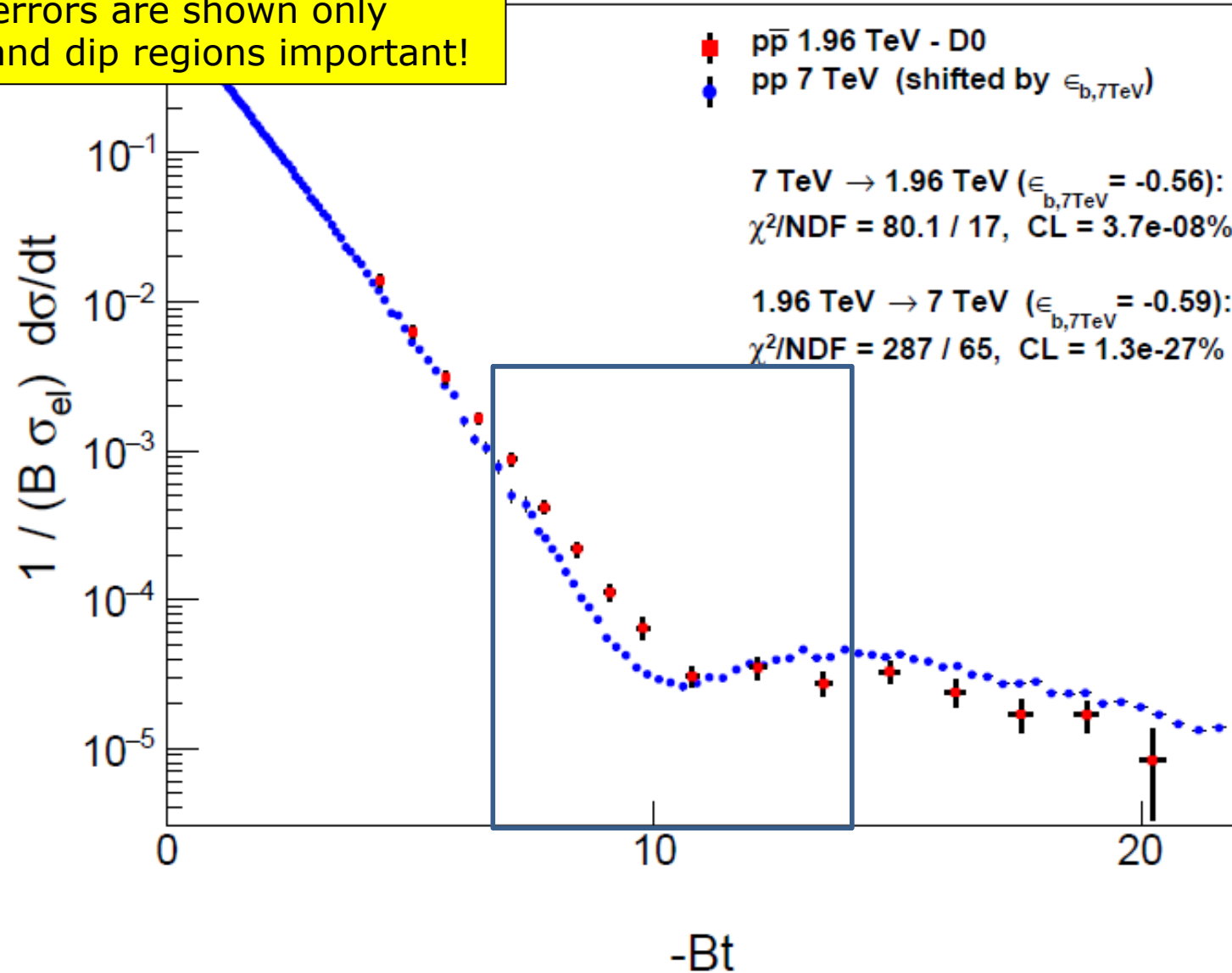
Odderon (3+5+... gluon) in pp:
 $(RGB) + (RGB) \rightarrow (GBR) + (BRG)$
 Well established in QCD

Odderon and elastic collisions



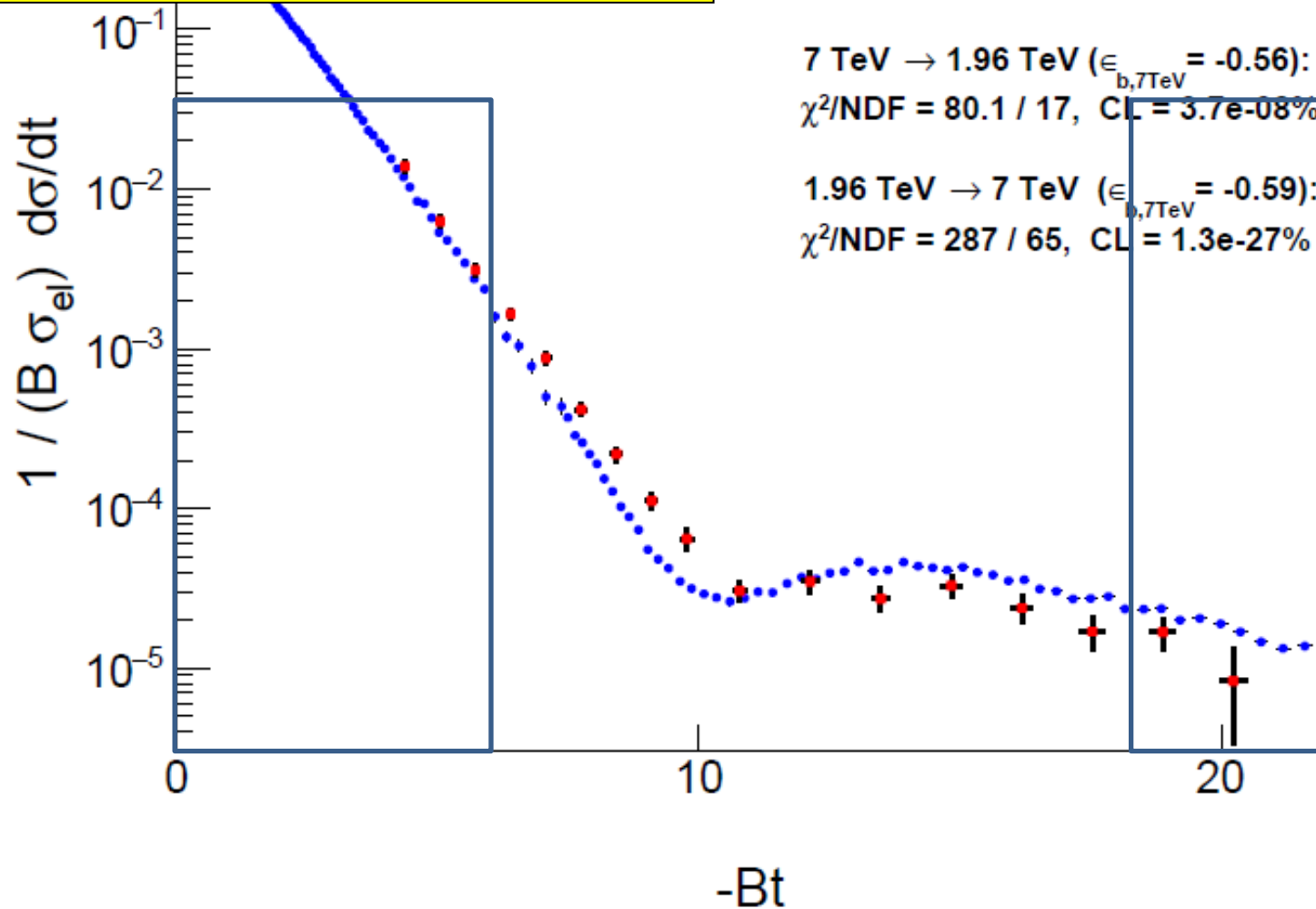
SLIDING WINDOWS

7 TeV data shifted
by $\epsilon_{B7,7\text{TeV}}$ to minimize χ^2
Type A errors are shown only
Both swing and dip regions important!



CLOSING DOORS

7 TeV data shifted
by $\epsilon_{B7,7\text{TeV}}$ to minimize χ^2
Type A errors are shown only
Both swing and dip regions important!



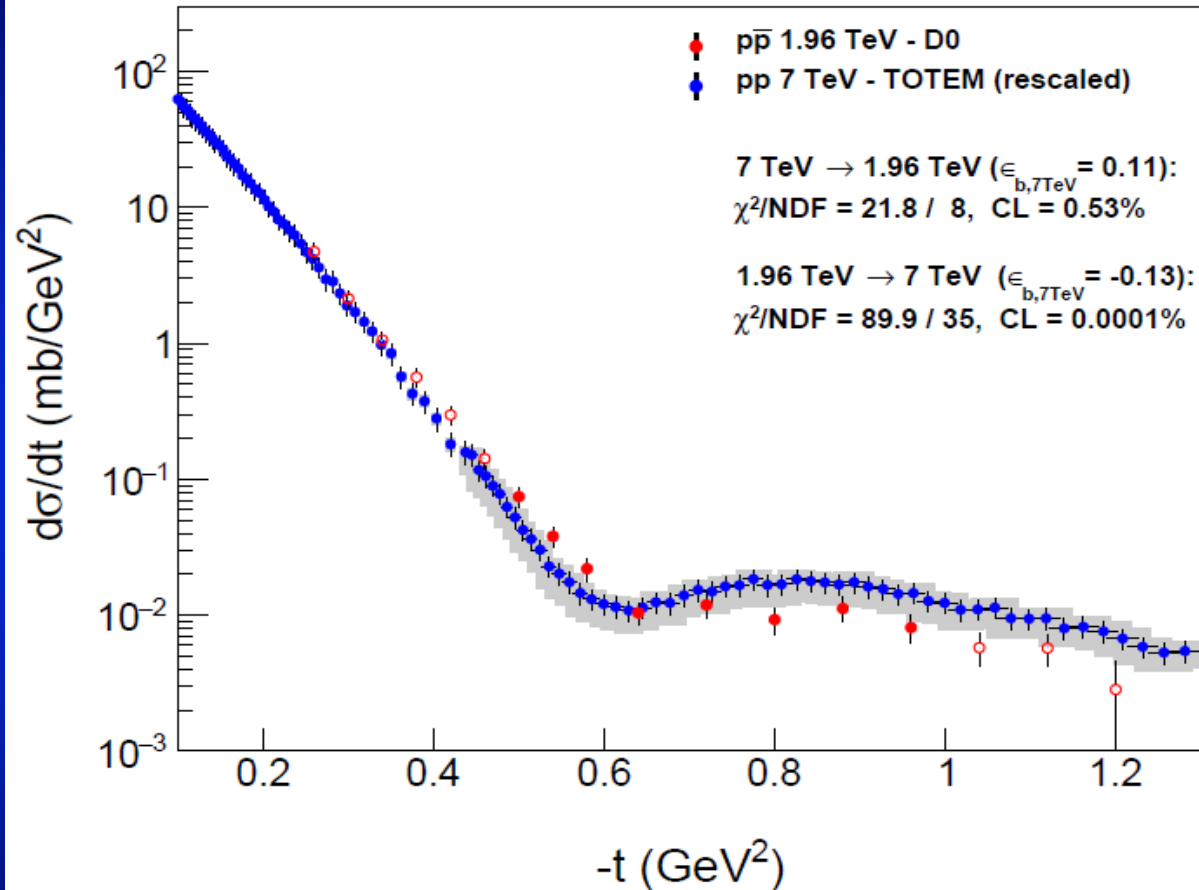
RESULTS FOR CLOSING DOORS

Two sliding doors of size n and size m:
(n,m): Leaving out the first n and last m D0 point

Sliding door technique with two wings (n,m)				
Left door excludes the first n, right door excludes the last m D0 points				
	n	m	Odderon signal	Background
	2	2	6.27 σ	1.68 σ
	3	2	6.33 σ	1.70 σ
	4	2	6.21 σ	2.37 σ
	2	1	6.11 σ	TBD
	2	2	6.27 σ	TBD
	2	3	5.90 σ	TBD

New MODEL INDEPENDENT RESULT
Odderon signal at least 6.33 σ

D0/TOTEM FIRMS UP OUR RESULTS



If we study $d\sigma/dt$ and limit **our analysis to the same range as D0/TOTEM**:
Significance reduces to **5.01 σ effect**, due to leaving out 9 D0 points

If we add D0's 14.4 % overall correlated error to fluctuating errors, for all D0 data:
Our *published* value is **3.27 σ**

If we conservatively optimize coefficient $\epsilon_{B,7\text{TeV}}$ of point-to-point correlated errors: **2.79 σ**
Significance of D0/TOTEM for $d\sigma/dt$: 3.4 σ

Recent results from D0/TOTEM

including our contributions

Comparison of pp and $p\bar{p}$ differential elastic cross sections and observation of the exchange of a #1
colorless C -odd gluonic compound

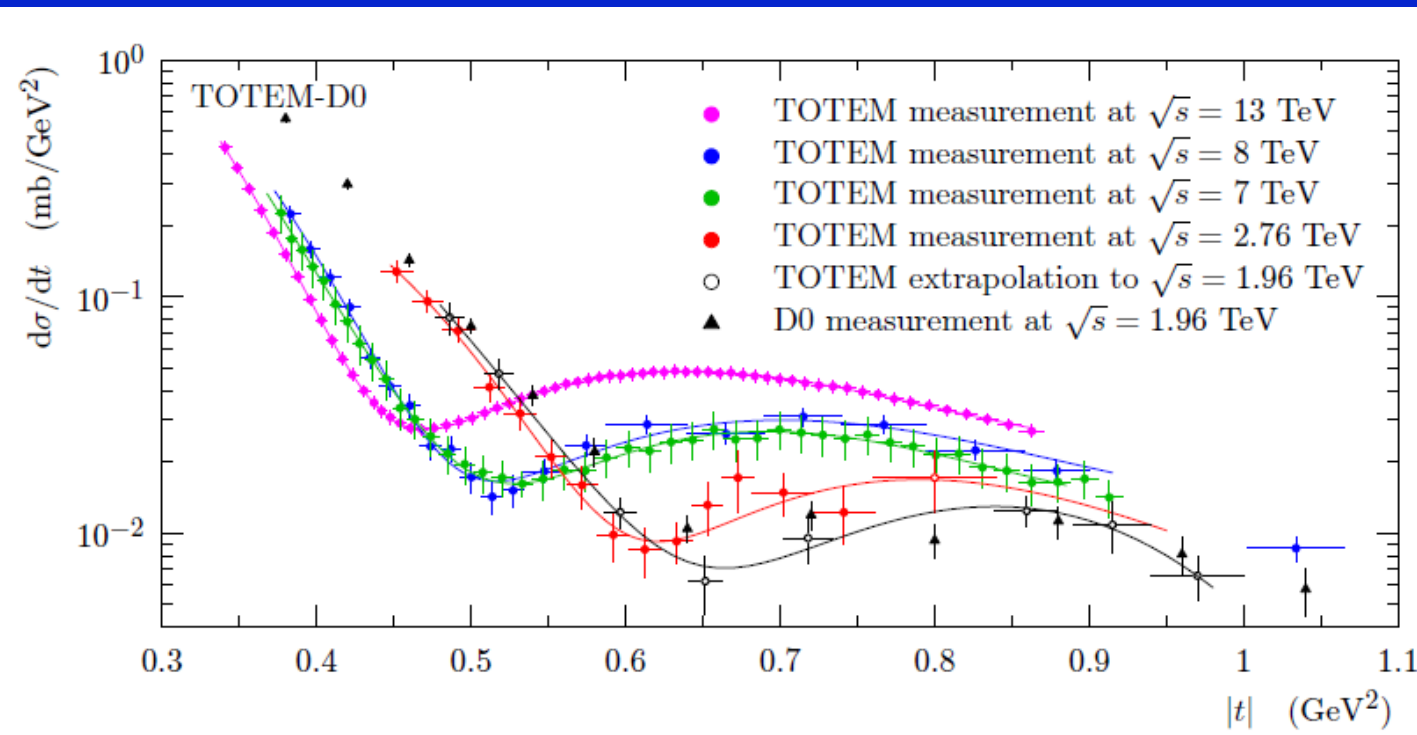
D0 and TOTEM Collaborations • V.M. Abazov (Dubna, JINR) et al. (Dec 7, 2020)

e-Print: 2012.03981 [hep-ex]

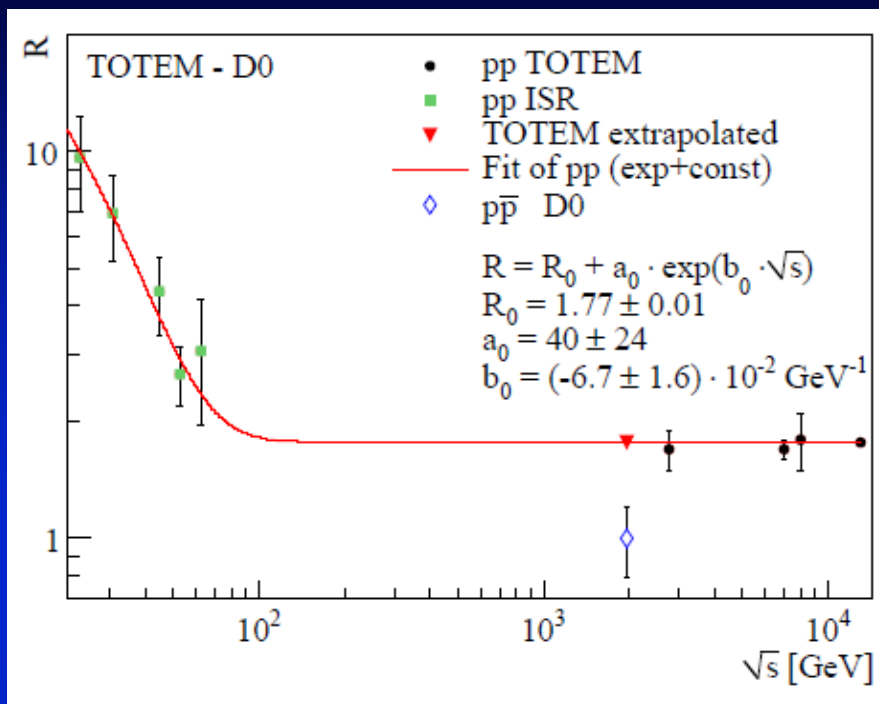
pdf links cite

1 citation

Submitted to PRL in December 2020.
Uses 13, 8, 7 and 2.76 TeV TOTEM data,
limited in $-t$ to the dip-bump structure.



APPENDIX: D0/TOTEM Fig. 2 OK



Our cross-test of Fig. 2 of [arxiv:2012.03981](https://arxiv.org/abs/2012.03981):

Fits ISR and LHC data with separate lines

$$p_1^{\text{LHC}} = 0.034 \pm 0.050$$

Consistent with 0 \rightarrow fix it to 0!

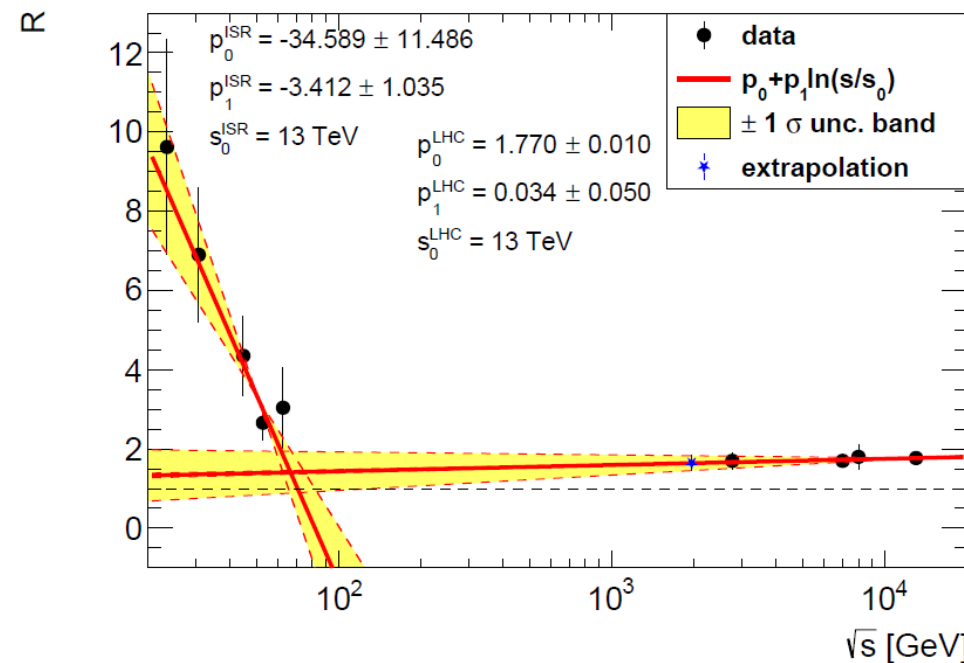
$$R(\text{pp}) = 1.77 \pm 0.01 \text{ @ } 1.96 \text{ TeV}$$

\rightarrow Reggeon effects negligible @ 1.96 TeV, OK.

Fig. 2 of [arxiv:2012.03981](https://arxiv.org/abs/2012.03981):
Fits ISR and LHC data with same curve

$$R(\text{pp}) = 1.77 \pm 0.01 \text{ @ } 1.96 \text{ TeV}$$

Reggeon effects from ISR? Test this!



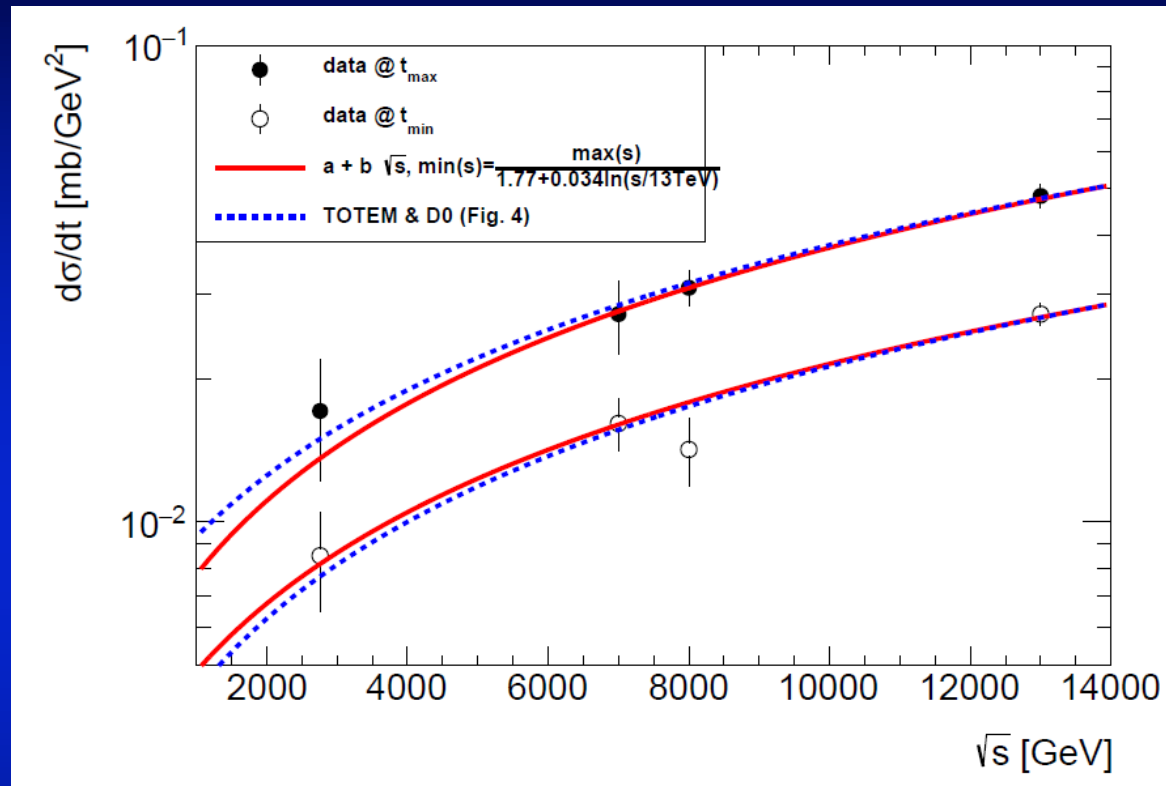
APPENDIX: D0/TOTEM FIG. 3 OK

Our cross-test of
Fig. 3 of [arxiv:2012.03981](https://arxiv.org/abs/2012.03981):
Fits to $\max(s)$ and $\min(s)$ neglect
the constraint of Fig. 2:

$$R(s|pp) = \max(s|pp)/\min(s|pp)$$

measured to be 1.77 ± 0.01 !

What about constrained fits?



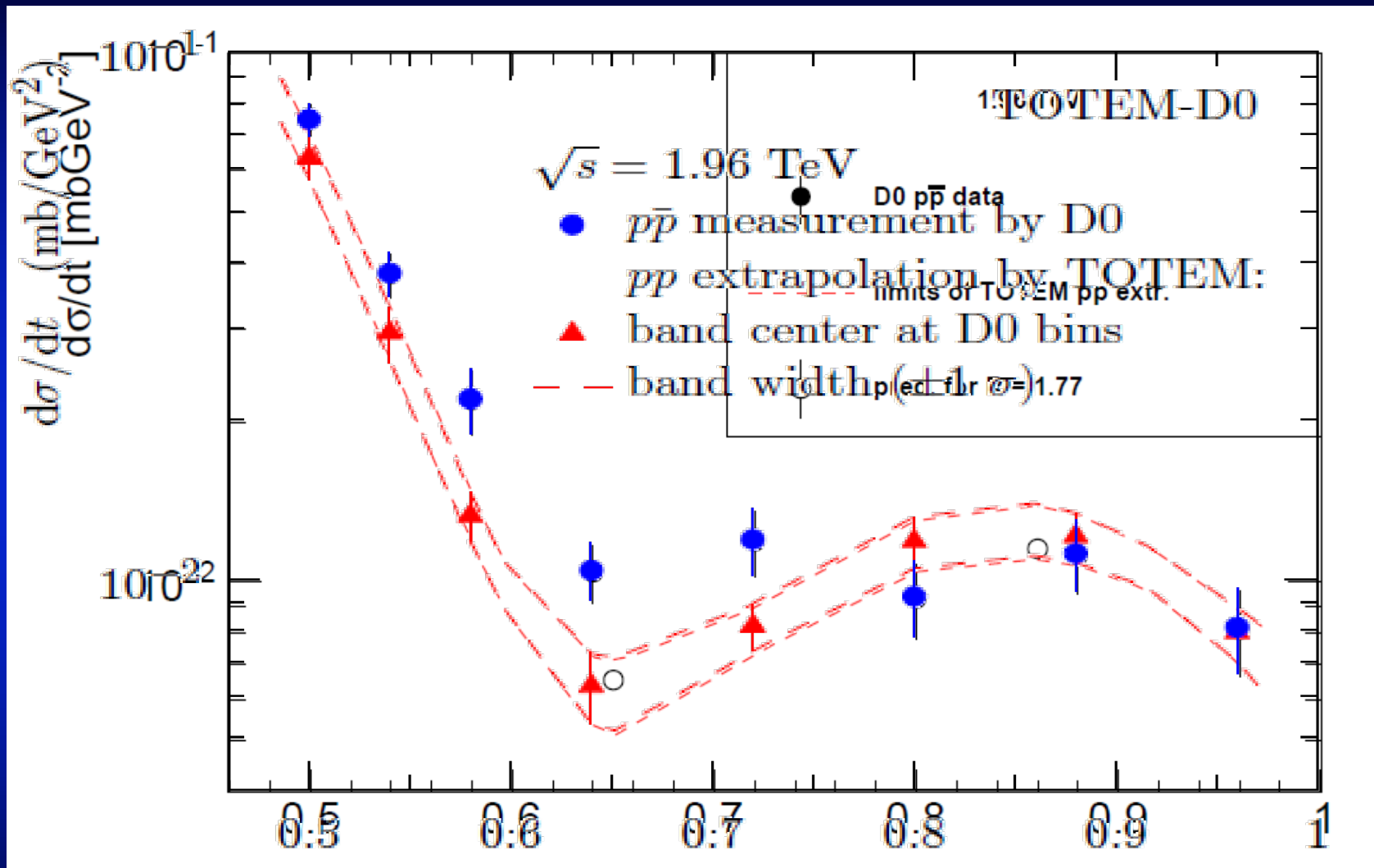
Only two out of three quantities can be fitted independently :

$$\max(s), \min(s) \text{ and } R(s) = \max(s) / \min(s)$$

Red lines: $\min(s|pp) = \max(s|pp)/R(s|pp)$ constrained fits

→ Fig. 3. of D0/TOTEM OK within 1σ

CROSS-CHECK OF D0/TOTEM FIG. 5



Empty circles from $\min(s | pp) = \max(s | pp) / R(s | pp)$ constrained fits

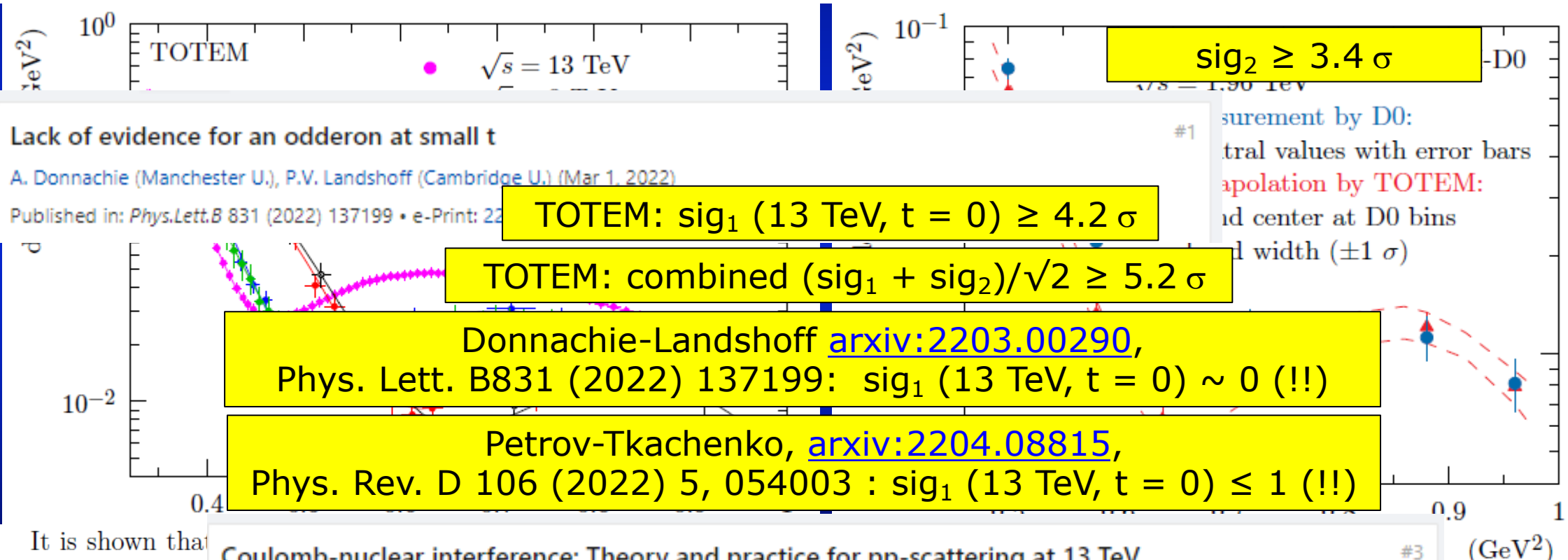
→ Fig. 5. of D0/TOTEM OK within 1σ

Status of D0-TOTEM Odderon search

Odderon Exchange from Elastic Scattering Differences between pp and $p\bar{p}$ Data at 1.96 TeV and from pp Forward Scattering Measurements #1

TOTEM and D0 Collaborations • V.M. Abazov (Dubna, JINR) et al. (Dec 7, 2020)
 Published in: *Phys.Rev.Lett.* 127 (2021) 6, 062003 • e-Print: 2012.03981 [hep-ex]

Phys. Rev. Lett. **127** (2021) 6, 062003, Published: 4 August 2021
<https://doi.org/10.1103/PhysRevLett.127.062003>



It is shown that

Coulomb-nuclear interference: Theory and practice for pp -scattering at 13 TeV #3

TOTEM – D0 detailed response :
 in preparation, see also

[K. Österberg's talk](#) at LHC Forward Physics meeting, Oct 2022

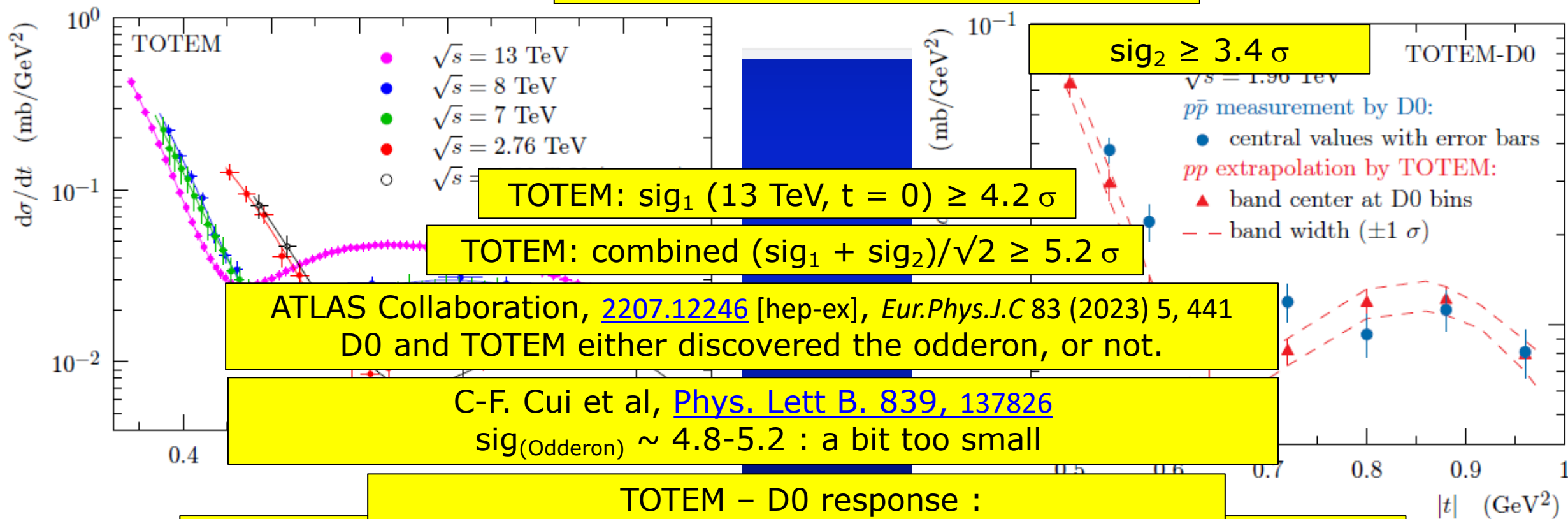
Status of D0-TOTEM Odderon 2.0

Odderon Exchange from Elastic Scattering Differences between pp and $p\bar{p}$ Data at 1.96 TeV and from pp Forward Scattering Measurements

#1

TOTEM and D0 Collaborations • V.M. A
Published in: *Phys.Rev.Lett.* 127 (2021) 6

Phys. Rev. Lett. **127** (2021) 6, 062003, [Published: 4 August 2021](https://doi.org/10.1103/PhysRevLett.127.062003)
<https://doi.org/10.1103/PhysRevLett.127.062003>



ATLAS Collaboration, [2207.12246](https://arxiv.org/abs/2207.12246) [hep-ex], *Eur.Phys.J.C* 83 (2023) 5, 441
D0 and TOTEM either discovered the odderon, or not.

C-F. Cui et al, [Phys. Lett B. 839, 137826](https://arxiv.org/abs/2301.13782)
 $\text{sig}_{(\text{Odderon})} \sim 4.8-5.2$: a bit too small

TOTEM – D0 response :

Paper in preparation, see also
[K. Österberg's talk at ISMD 2023](#)

in August 2023 and at Diffraction and Low-x 2024

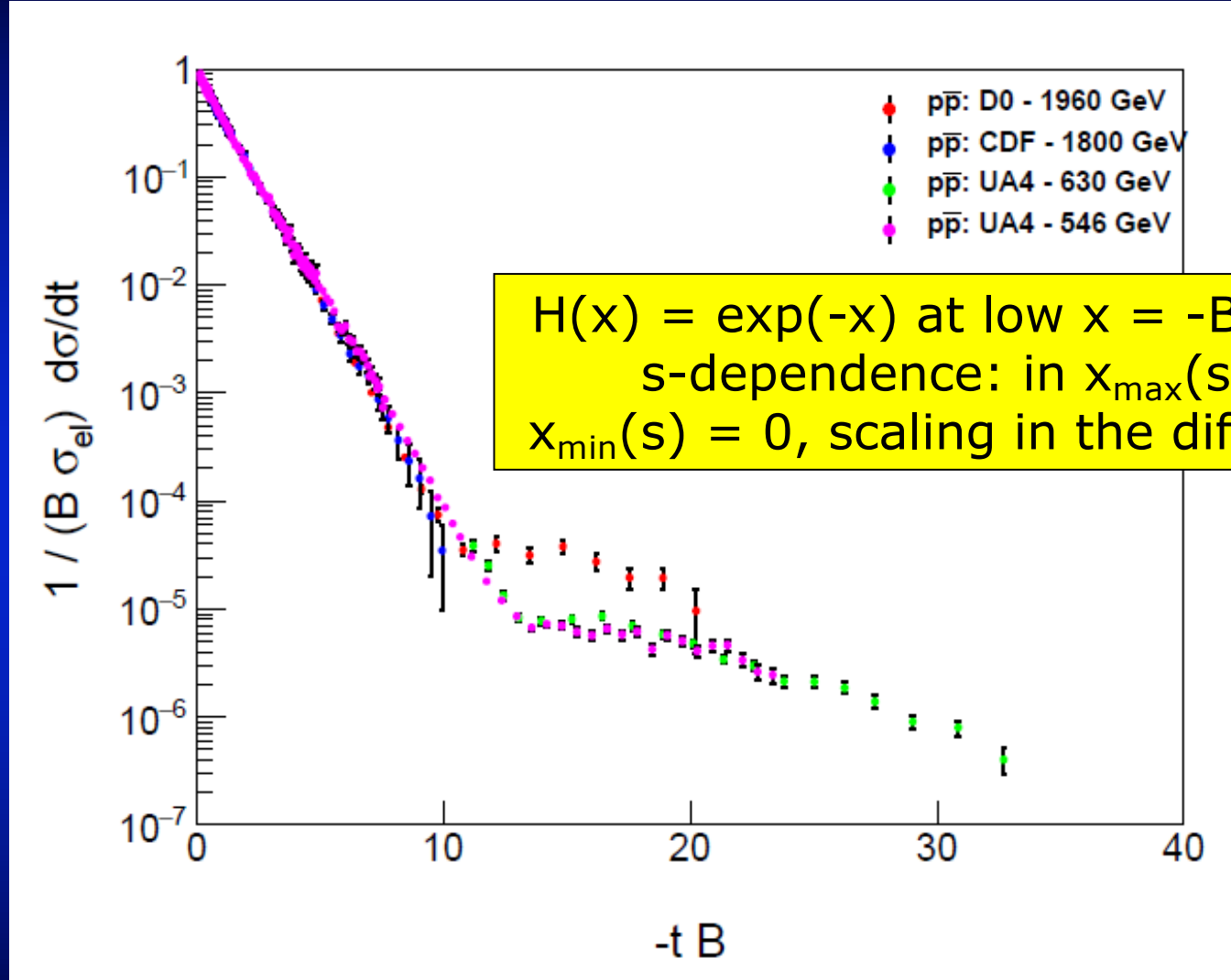
C0: i

TeV

NEW RESULTS 1

H(x) SCALING, USING 8 TeV

H(x) scaling for p antip scattering



Energy range: 546 GeV – 1.96 TeV
Qualitatively different from pp: scaling in the cone only for p+antip

Where is the Odderon signal from?

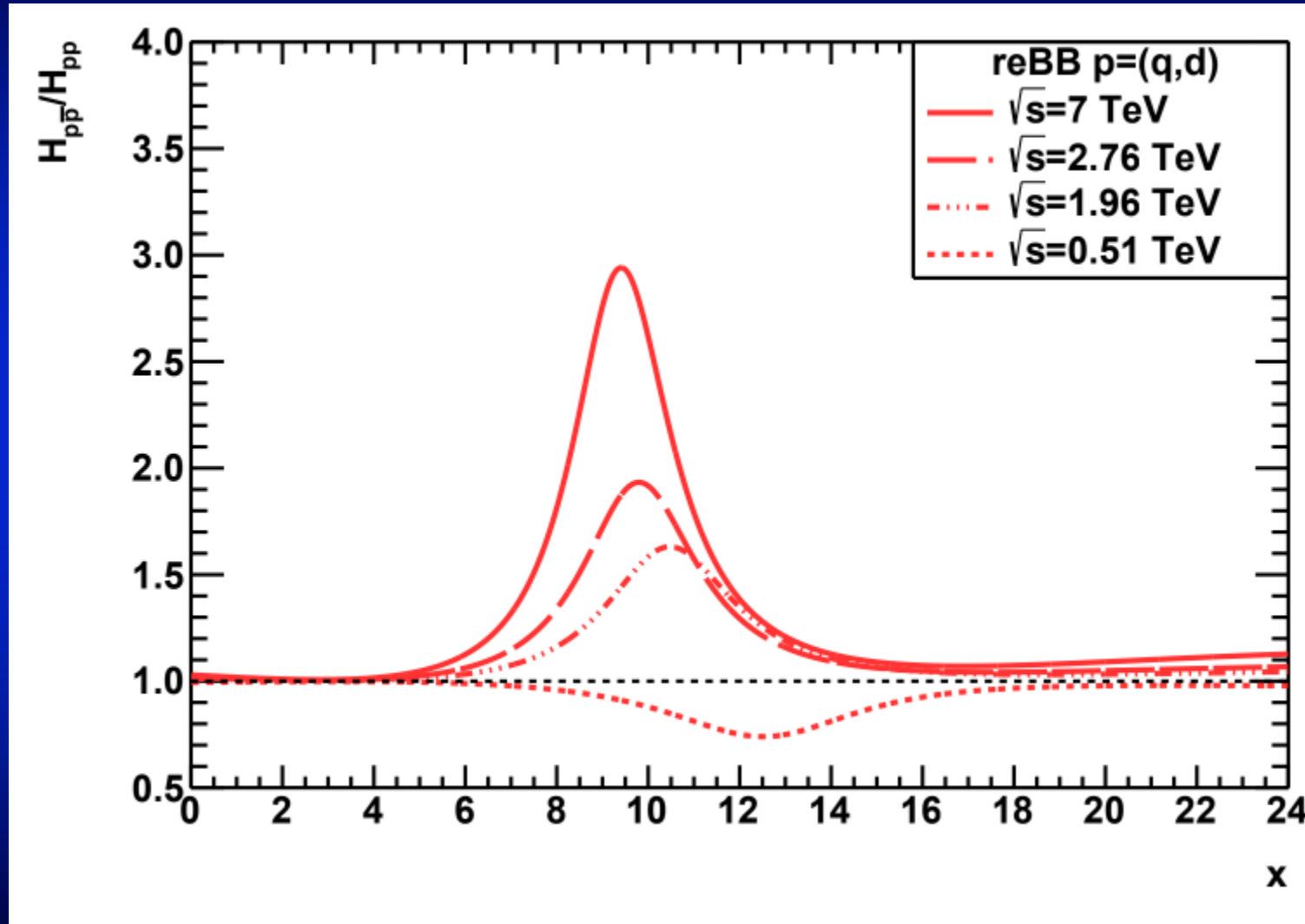
Swing, interference, tail regions
Interference region is dominant

Partial significances from the swing, interference, tail and all regions,
characterized by $x_{\min} < x \leq x_{\max}$

x_{\min}	x_{\max}	ϵ_{B21} of $\min \Delta \chi^2$ in $x_{\min} < x \leq x_{\max}$	$\Delta \chi^2$ in $x_{\min} < x \leq x_{\max}$	NDF in $x_{\min} < x \leq x_{\max}$	σ in $x_{\min} < x \leq x_{\max}$
5.1	8.4	1.90	4.19	5	0.64
8.4	13.5	-0.49	25.31	5	3.84
13.5	20.2	-1.39	1.79	5	0.15
5.1	13.5	0.28	48.27	10	5.01
8.4	20.2	-0.96	35.79	10	3.91
5.1	20.2	-0.60	75.41	15	6.23

OBSERVATION OF ODDERON

2020 → 2020

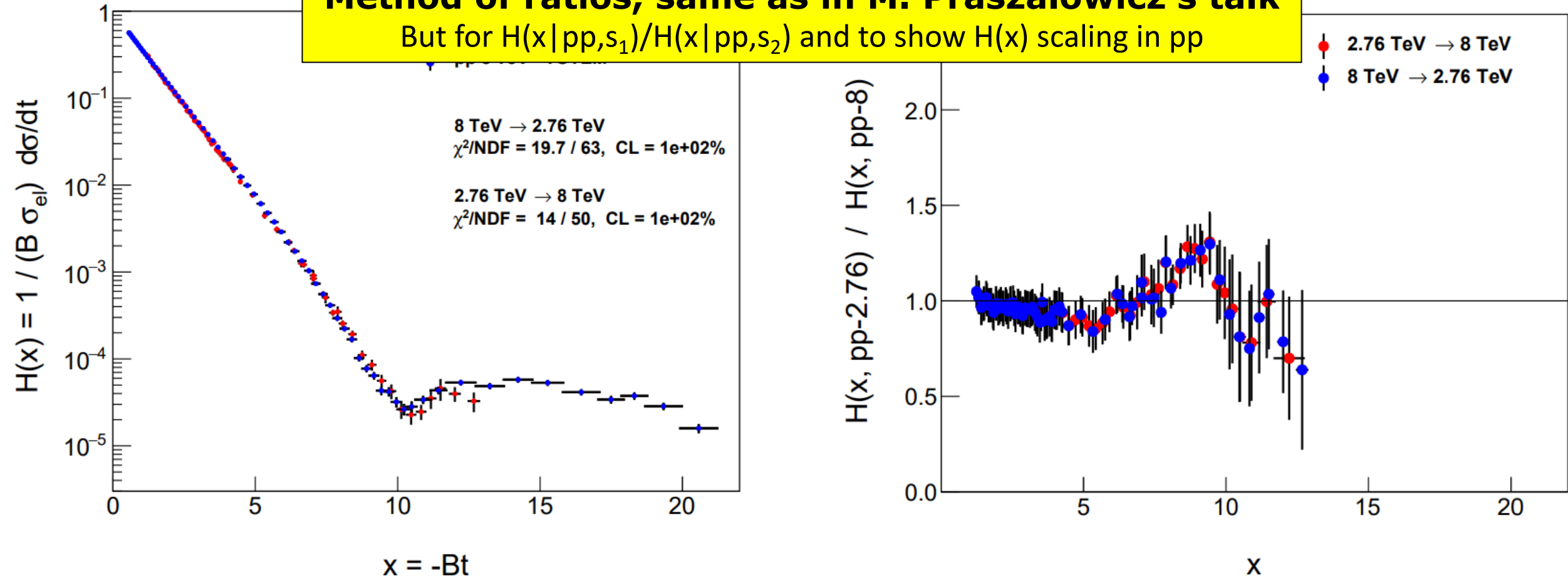


Prediction for 510 GeV pp @ RHIC: scaling violations

Test of $H(x)$ scaling: 8 vs 2.76 TeV

Method of ratios, same as in M. Praszalowicz's talk

But for $H(x|pp,s_1)/H(x|pp,s_2)$ and to show $H(x)$ scaling in pp



Between 2.76 and 8 TeV, $H(x)$ scaling observed!

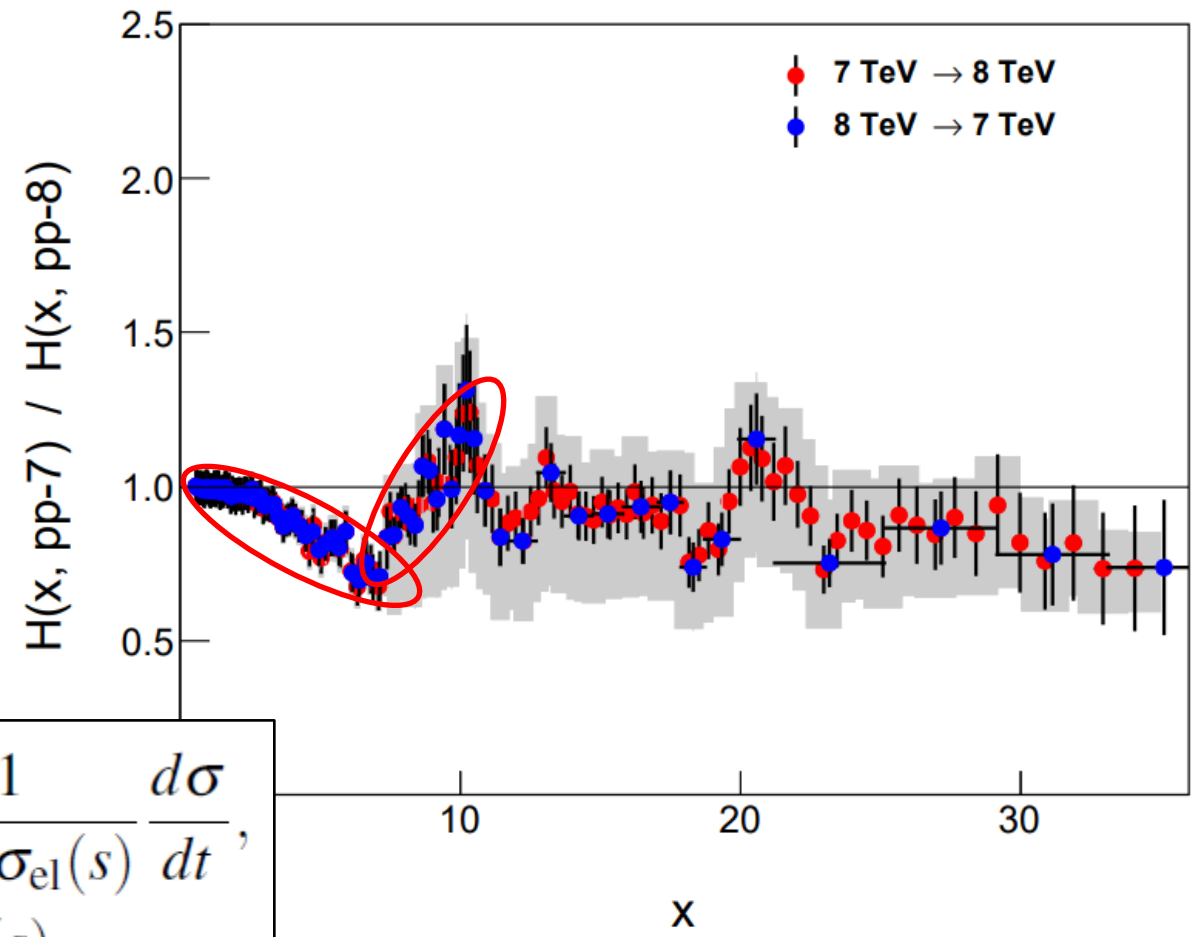
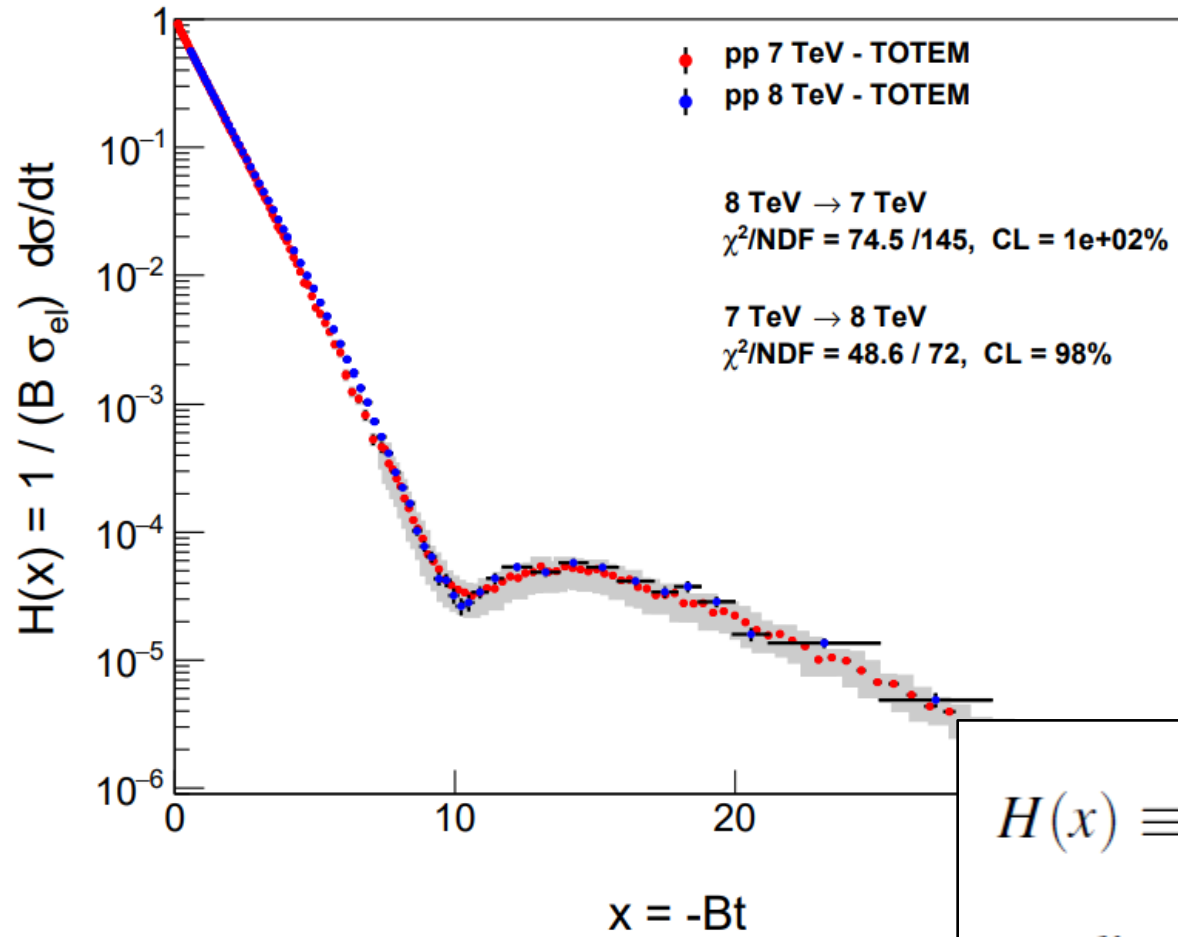
Hungarian-Swedish team, e-Print: [2405.06733](https://arxiv.org/abs/2405.06733) [hep-ph],

MDPI Universe 2024, 10(6), 264

$$H(x) \equiv \frac{1}{B(s)\sigma_{el}(s)} \frac{d\sigma}{dt},$$

$$x = -tB(s).$$

Test of $H(x)$ scaling: 8 vs 7 TeV TOTEM



$$H(x) \equiv \frac{1}{B(s)\sigma_{\text{el}}(s)} \frac{d\sigma}{dt},$$

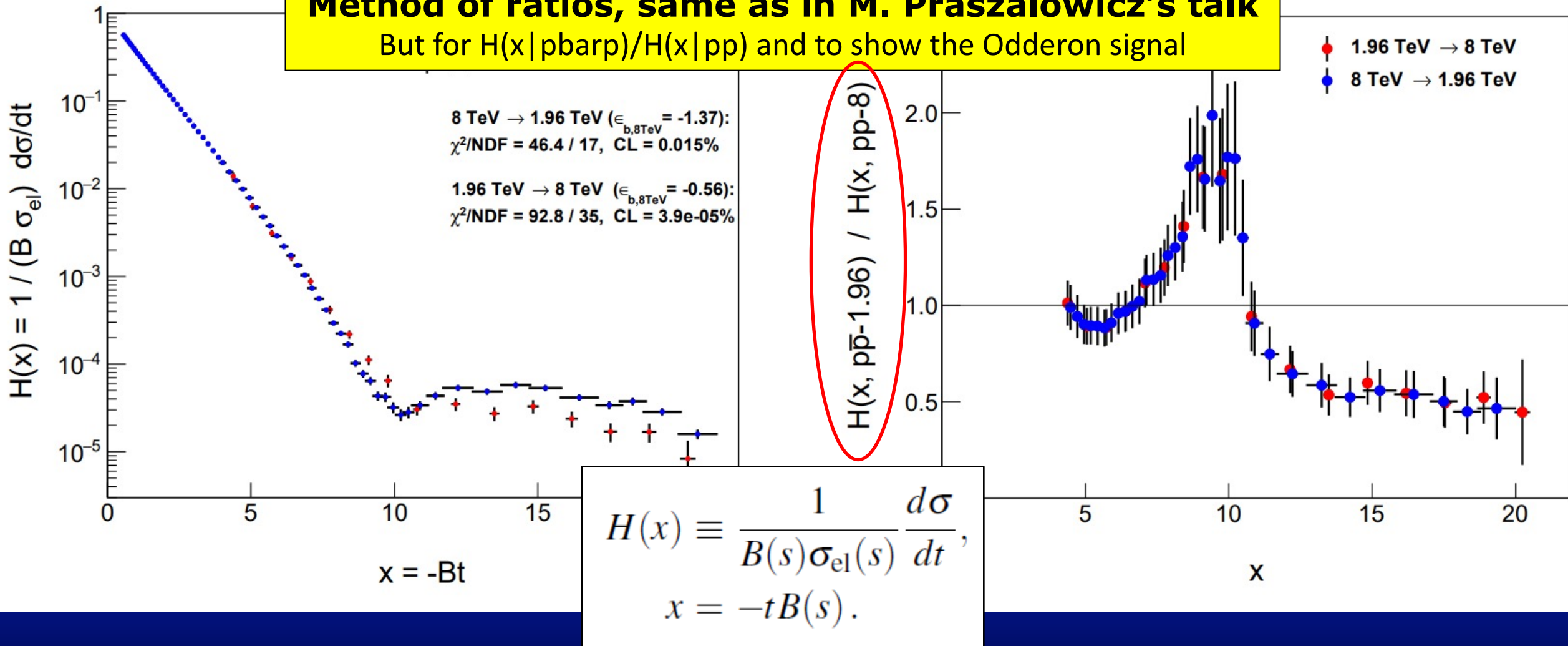
$$x = -tB(s).$$

Between 7 and 8 TeV, $H(x)$ scaling observed, but ...
 Hungarian-Swedish team, e-Print: [2405.06733](https://arxiv.org/abs/2405.06733) [hep-ph],
 MDPI Universe 2024, 10(6), 264

Odderon of H(x) scaling: 8 vs 1.96 TeV

Method of ratios, same as in M. Praszalowicz's talk

But for $H(x|p\bar{p})/H(x|pp)$ and to show the Odderon signal



Between 1.96 and 8 TeV, $H(x|s,pp)$ and $H(x|s,p\bar{p})$ are clearly different, with $3 < 3.79 < 5 \sigma$

Odderon significances from H(x) scaling

\sqrt{s} (TeV)	χ^2	NDF	CL	significance (σ)
1.96 vs. 2.76	3.85	11	9.74×10^{-1}	0.03
1.96 vs. 7	80.1	17	3.681×10^{-10}	6.26
1.96 vs. 8	46.4	17	1.502×10^{-4}	3.79

\sqrt{s} (TeV)	χ^2	NDF	CL	χ^2 /NDF method	combined σ Stouffer's method
1.96 vs 2.76 & 8	50.25	28	6.064×10^{-3}	2.74	2.70
1.96 vs 2.76 & 7	83.95	28	1.698×10^{-7}	5.22	4.44
1.96 vs 2.76 & 7 & 8	130.35	45	2.935×10^{-10}	6.30	5.81
1.96 vs 7 & 8	126.5	34	1.415×10^{-12}	7.08	7.10

If 1.96, 2.76, 7 and 8 TeV data are combined, H(x) significances on all data results in $5 < 5.8 \sigma$
If 1.96, 7 and 8 TeV data are combined, at least 7.08 σ .

Hungarian-Swedish team, e-Print: [2405.06733](https://arxiv.org/abs/2405.06733) [hep-ph],

MDPI Universe **2024**, *10*(6), 264