



Anisotropic Flow Correlations in Heavy-Ion collisions

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Outline:

- Collective flow and flow fluctuations
- Elliptic and triangular flow
- Mixed flow cumulants
- Intrinsic cumulants: a new perspective
- Predictions and comparison with data

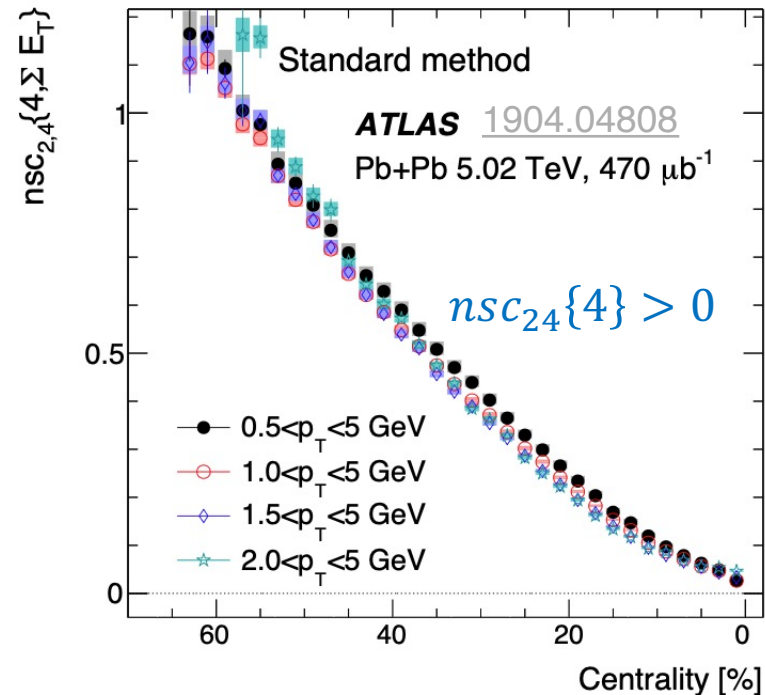
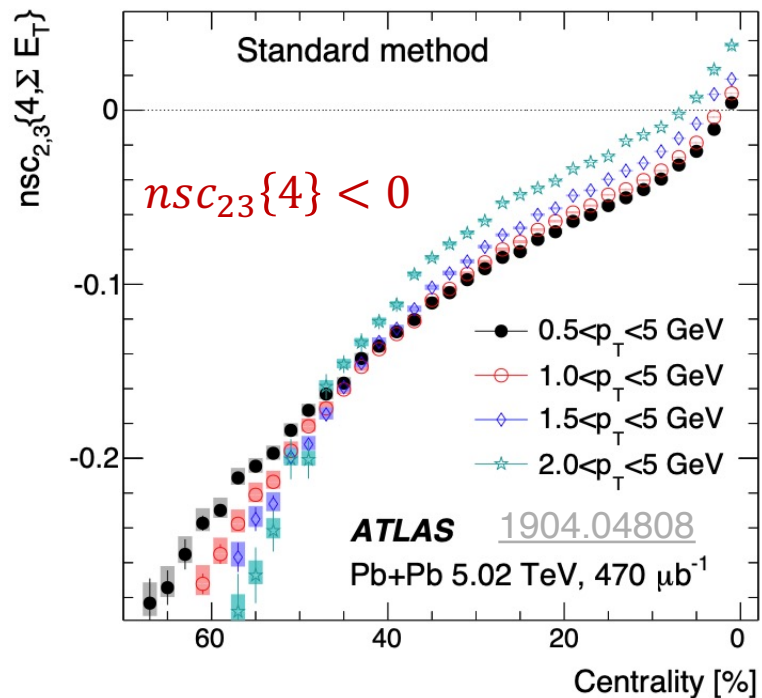
Based on:

M Alqahtani and Jean-Yves Ollitrault [*Phys.Lett.B* 872 \(2026\) 140066](#)

M Alqahtani and Jean-Yves Ollitrault [*Phys.Lett.B* 877 \(2026\) 140504](#)

Correlations between flow harmonics I

$$nsc_{n,m}\{4\} \propto \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$$



- v_2 and v_3 are anti-correlated, while v_2 and v_4 are positively correlated.

From initial geometry to collective flow

The initial overlap geometry of the collision is quantified by the eccentricities

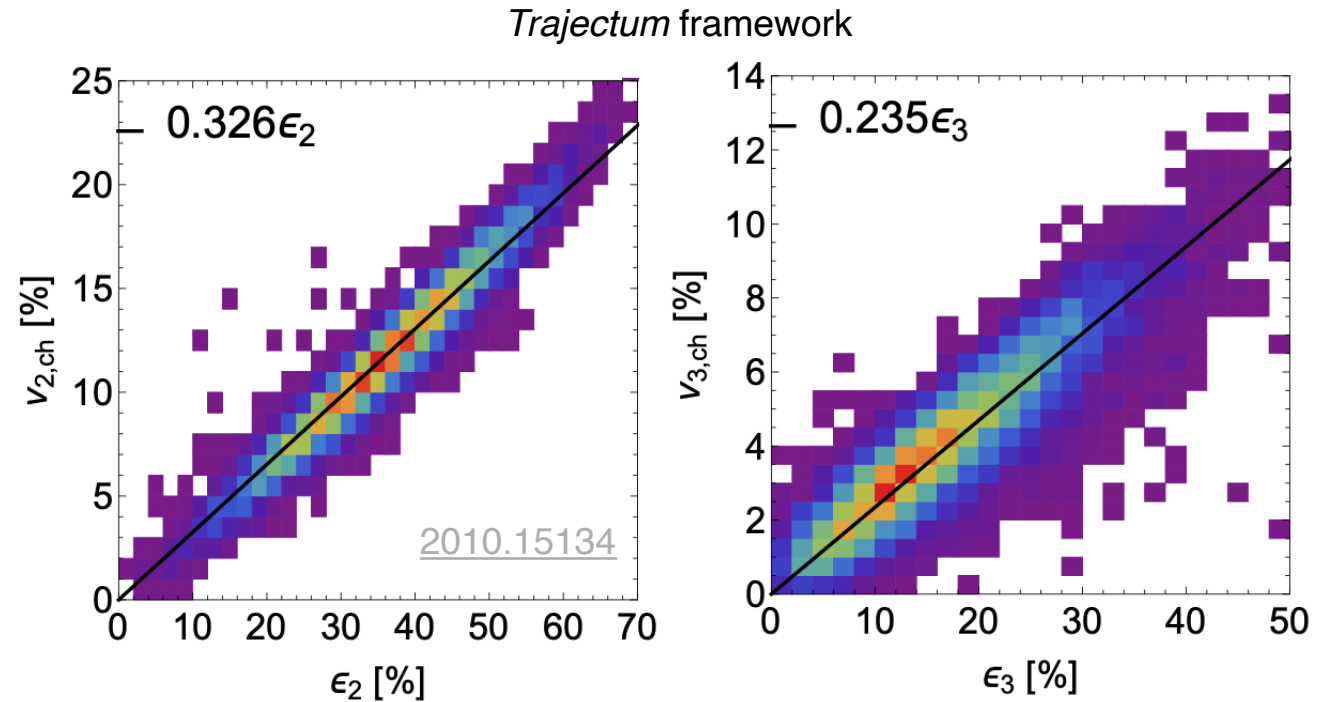
$$\boldsymbol{\varepsilon}_n = \varepsilon_n e_n^{in\Phi_n}$$

ε_2 : ellipticity ε_3 : triangularity

The final-state anisotropic flow is characterized by

$$\mathbf{V}_n = v_n e_n^{in\Psi_n}$$

v_2 : elliptic flow v_3 : triangular flow



Hydrodynamic evolution converts initial-state eccentricities ε_n into final-state collective flow.

$$v_2 \approx \kappa_2 \varepsilon_2 \quad \& \quad v_3 \approx \kappa_3 \varepsilon_3$$

See also H. Niemi et al [1212.1008](#)

Key message:

- Event-by-event fluctuations of the initial geometry are reflected in the flow harmonics v_2 and v_3 .
- We will use this approximate relations to understand relations between anisotropic flow coefficients.

Distributions of initial eccentricity

- Hydrodynamic response implies

$$P(v_2) \propto P(\varepsilon_2). \quad \&. \quad P(v_3) \propto P(\varepsilon_3).$$

What one finds (See slide 10 for a clearer illustration)

ε_2 : geometry-dominated.

ε_3 : fluctuation-dominated.

ε_4 : mixed origin (geometry + fluctuations).

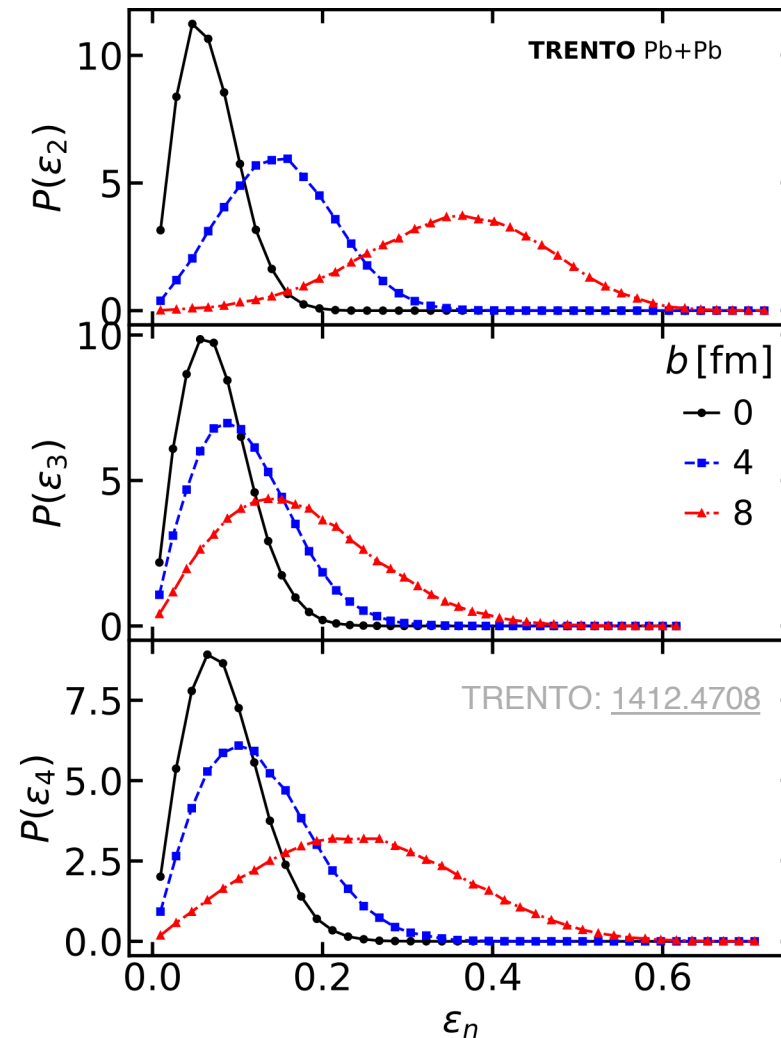
- Eccentricity fluctuations are approximately Gaussian.
- Deviations from Gaussianity:
 - Bessel-Gaussian distribution
 - Elliptic power distribution

See:

Voloshin et al [0708.0800](#)

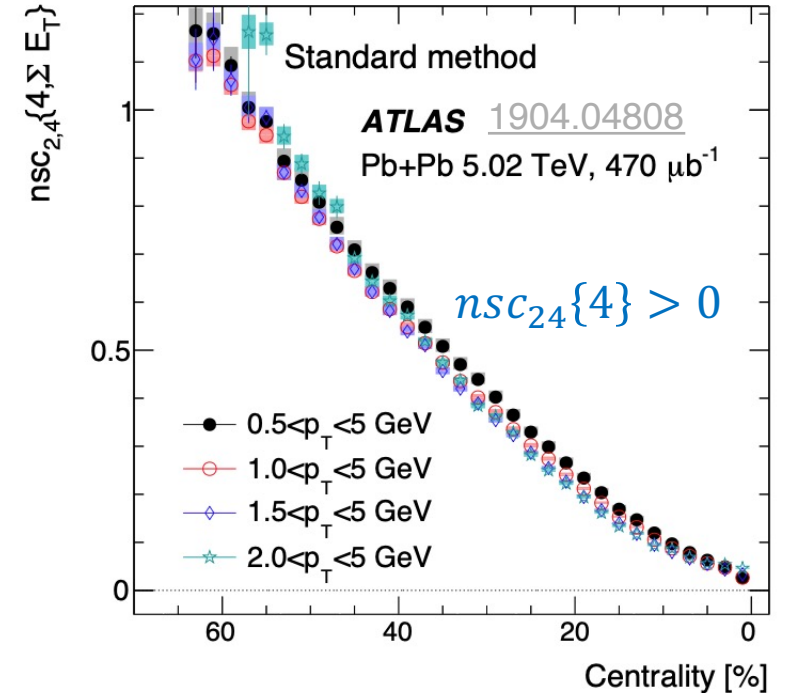
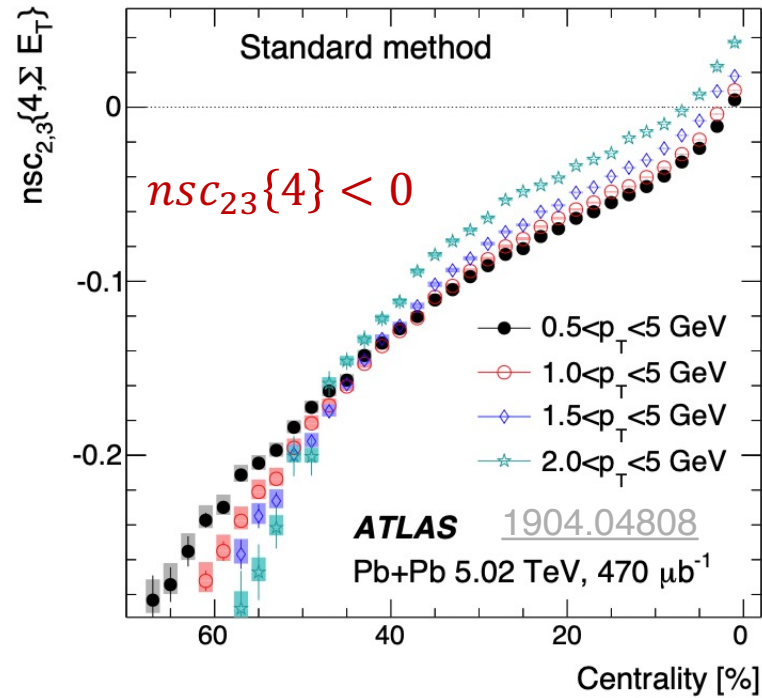
Li Yan et al [1405.6595](#)

TRENTO predictions for eccentricity distributions at different b



Correlations between flow harmonics II

$$nsc_{n,m}\{4\} \propto \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$$



- v_2 and v_3 are **anti-correlated**, while v_2 and v_4 are **positively correlated**.
- $nsc_{23}\{4\} \ll 1$, can be understood by the smallness of non-Gaussian effects.
- In this talk, we address the question: **What drives the v_2 - v_3 anti-correlation?**

Moment and cumulant generating functions

For a random variable X , the moment-generating function is defined as:

$$M(t) \equiv \langle e^{tX} \rangle = 1 + t\langle X \rangle + \frac{t^2}{2!}\langle X^2 \rangle + \dots = \sum_{n=0}^{\infty} \frac{t^n}{n!} \langle X^n \rangle.$$

Where t is an expansion parameter. One can isolate any moment by:

$$\langle X^m \rangle = \left. \frac{d^m M(t)}{dt^m} \right|_{t=0}$$

The cumulant-generating function is defined by:

$$K(t) \equiv \ln M(t) = \ln \langle e^{tX} \rangle = t \kappa_1 + \frac{t^2}{2!} \kappa_2 + \dots = \sum_{n=1}^{\infty} \frac{t^n}{n!} \kappa_n$$

$$\kappa_1 = \langle X \rangle; \text{ mean } \mu$$

$$\kappa_2 = \langle X^2 \rangle - \langle X \rangle^2; \text{ variance } \sigma^2$$

$$\kappa_3 = \langle X^3 \rangle - 3\langle X \rangle \langle X^2 \rangle + 2\langle X \rangle^3$$

$$\kappa_4 = \langle X^4 \rangle - 4\langle X \rangle \langle X^3 \rangle - 3\langle X^2 \rangle^2 + 12\langle X \rangle^2 \langle X^2 \rangle - 6\langle X \rangle^4$$

... and so on



$$\mu = \kappa_1$$

$$\sigma^2 = \kappa_2$$

$$\text{skw} = \kappa_3 / \sigma^3$$

$$\text{Excess kurt} = \frac{\kappa_4}{\sigma^4}$$

Remark:

$\kappa_n = 0 \ (n \geq 3)$
for a Gaussian Distribution

See Rajeev Bhalerao, [arXiv:2009.09586](https://arxiv.org/abs/2009.09586)

Measuring correlations between v_2 and v_3

- Correlations between v_2 and v_3 can be quantified using moments or cumulants.
- Mixed-harmonic cumulants (MHC) isolate **genuine correlations** between v_2 and v_3 .
- MHC comes in different orders: $MHC(v_2^{2m}, v_3^{2q})$

For example:

$$MHC(v_2^2, v_3^2) = \langle v_2^2 v_3^2 \rangle - \langle v_2^2 \rangle \langle v_3^2 \rangle ; \text{ Lowest-order cumulant (covariance)}$$

$$MHC(v_2^4, v_3^2) = \langle v_2^4 v_3^2 \rangle - 4 \langle v_2^2 v_3^2 \rangle \langle v_2^2 \rangle - \langle v_2^4 \rangle \langle v_3^2 \rangle + 4 \langle v_2^2 \rangle^2 \langle v_3^2 \rangle$$

...

- To compare different orders and systems, cumulants are often normalized.

Normalized correlation strength

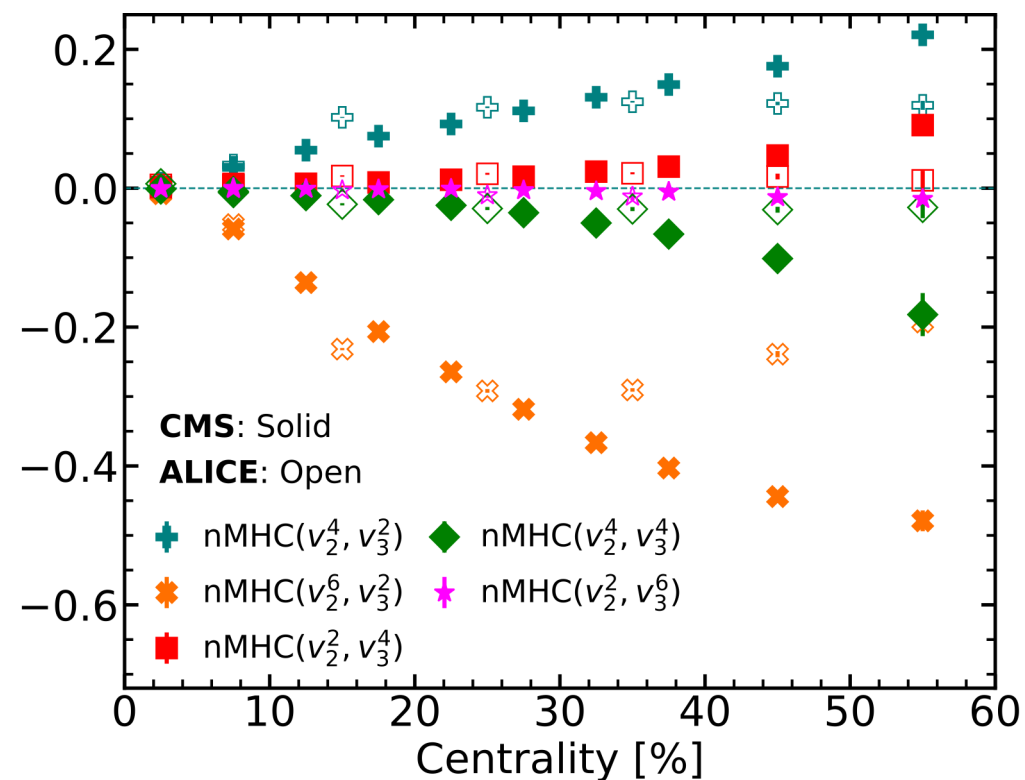
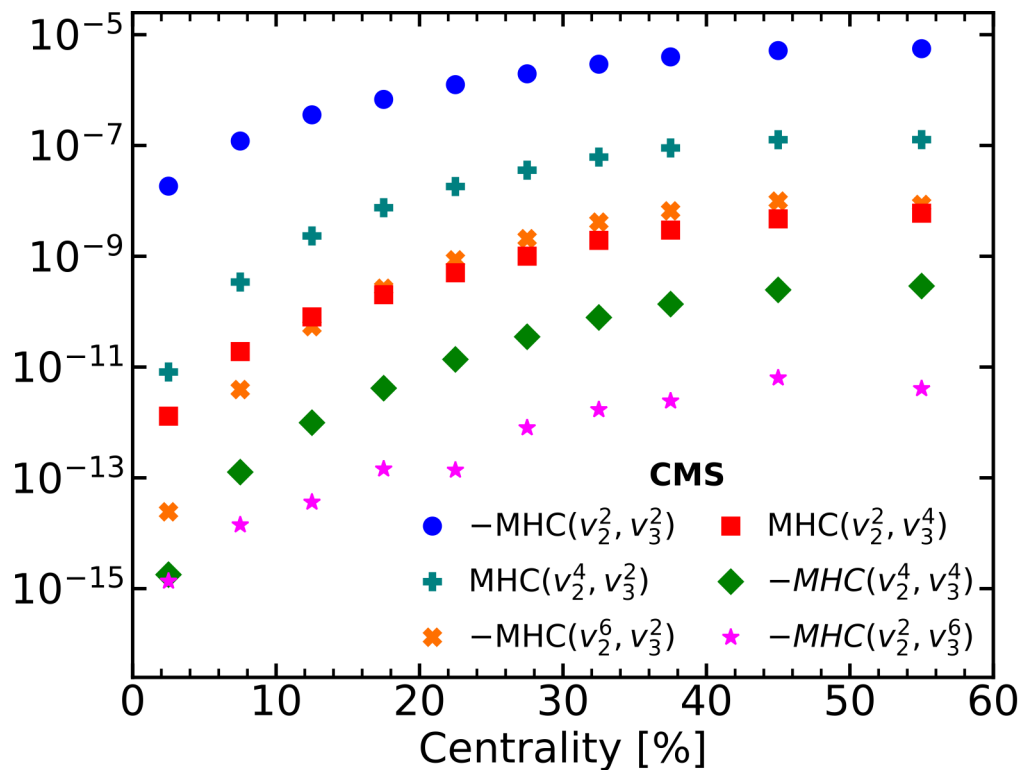
$$nMHC(v_2^{2m}, v_3^{2q}) \equiv \frac{MHC(v_2^{2m}, v_3^{2q})}{\langle v_2^{2m} \rangle \langle v_3^{2q} \rangle}$$

0 → no correlation

+1 → correlated

-1 → anti-correlated

Measured v_2-v_3 mixed-harmonic cumulants

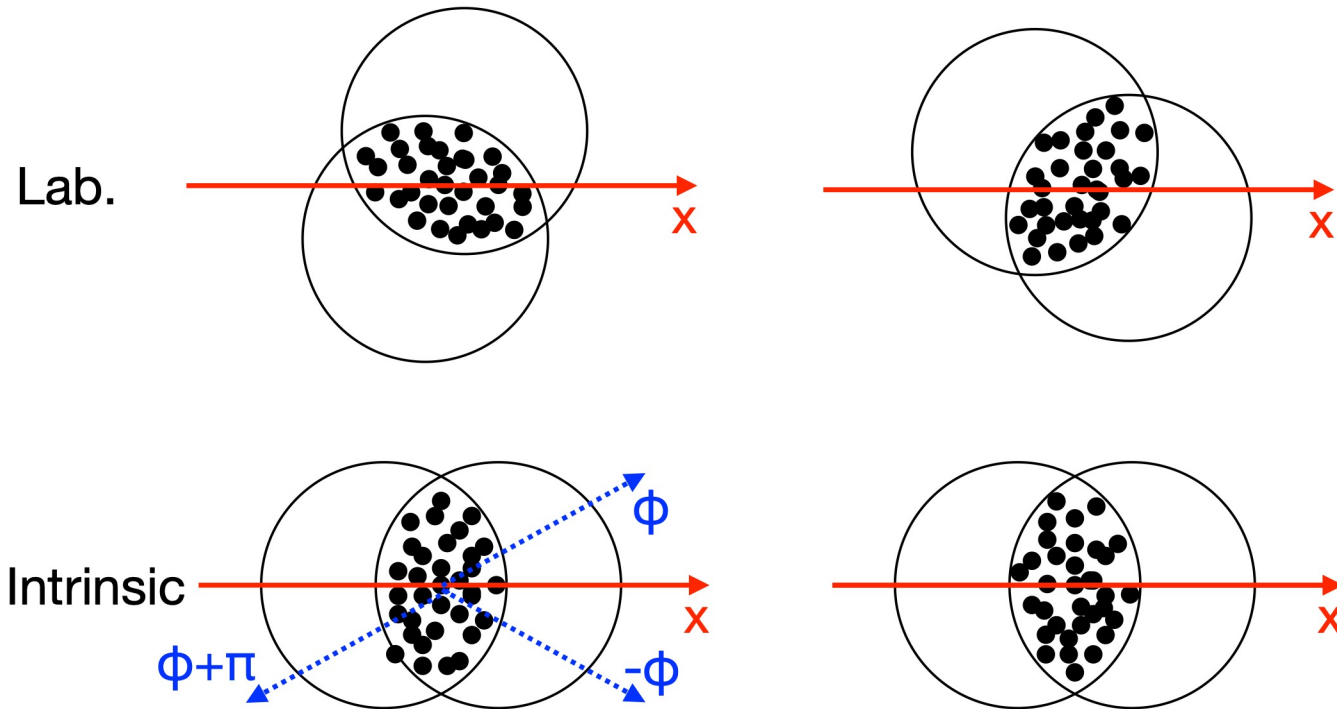


- Higher-order MHCs are suppressed by several orders of magnitude.

- Non-trivial higher-order v_2-v_3 correlation patterns.
- The sign changes with increasing m or q .

Lab vs. intrinsic frames

Two collision events with the same impact parameter b



Intrinsic frame reveals the underlying anisotropy washed out in the lab frame

LAB Frame:

- Random reaction-plane Ψ_{RP}
- Azimuthal averaging

$$\langle v_{2x} \rangle_{\text{lab}} = 0.$$

Intrinsic Frame:

- All events are aligned, $b \parallel x$
- Preserves reaction-plane information
- Nonzero mean elliptic flow in the reaction plane

$$\langle v_{2x} \rangle_{\text{int}} \equiv \bar{V}_2 \neq 0,$$

- At $b=0$, rotational symmetry is exact:

$$\langle v_{2x} \rangle = 0,$$

intrinsic and lab cumulants coincide.

Initial ellipticity and triangularity

TRENTO initial conditions at fixed impact parameter

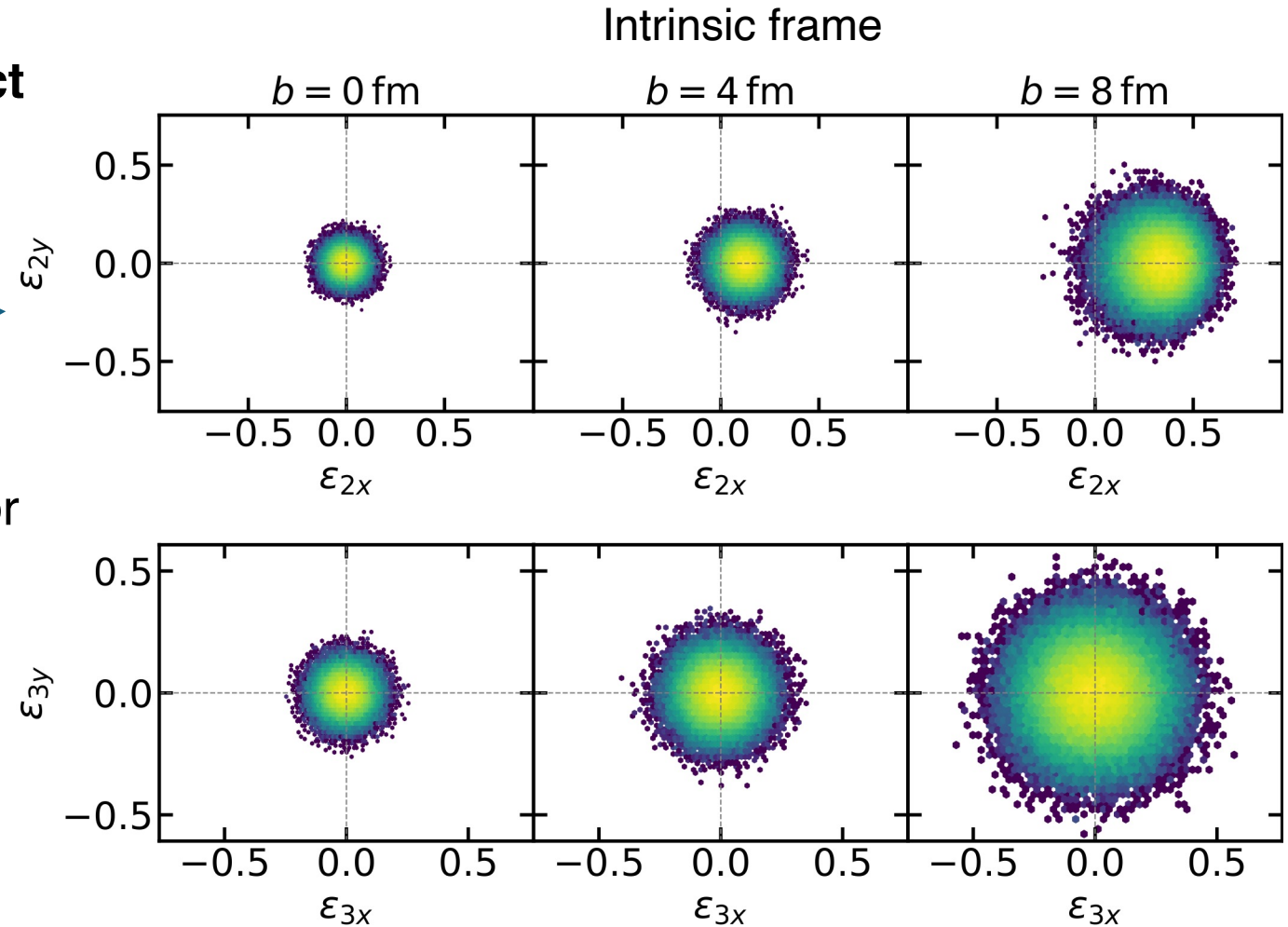
- Probability distribution of the eccentricity vector $\vec{\epsilon}_n = (\epsilon_{nx}, \epsilon_{ny})$
 $\epsilon_{nx} = \epsilon_n \cos(n\Phi_n)$ & $\epsilon_{ny} = \epsilon_n \sin(n\Phi_n)$

- Only ellipticity develops a nonzero mean for $b > 0$.

$$\langle \epsilon_{2x} \rangle \neq 0, \quad \langle \epsilon_{2y} \rangle = 0$$

$$\langle \epsilon_{3x} \rangle = \langle \epsilon_{3y} \rangle = 0$$

- v_2 is geometry-driven, whereas v_3 is fluctuation-driven.



Predictions from intrinsic cumulants

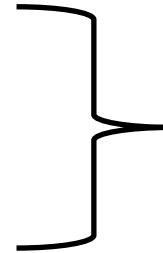
Intrinsic cumulants are defined as *

$$c_{mpqr} = \langle V_2^m (V_2^*)^p V_3^q (V_3^*)^r \rangle_c$$

where $\langle \dots \rangle_c$ denotes the connected cumulant

Lab-frame mixed cumulants can be expanded in terms of intrinsic cumulants.

$$\begin{aligned} MHC(v_2^2, v_3^2) &= 2\bar{V}_2 c_{0111} + \dots \\ MHC(v_2^4, v_3^2) &= -4\bar{V}_2^3 c_{0111} + \dots \\ MHC(v_2^6, v_3^2) &= 24\bar{V}_2^5 c_{0111} + \dots \end{aligned}$$



- Higher-order terms are neglected (see slide 11).
- $\bar{V}_2 \approx v_2\{4\}$ See: [arXiv:0708.0800](https://arxiv.org/abs/0708.0800)

c_{0111} quantifies the linear correlation between elliptic flow in the reaction plane and the triangularity.

$$c_{0111} = \langle v_{2x} v_3^2 \rangle - \langle v_{2x} \rangle \langle v_3^2 \rangle \text{ mixed skewness}$$

Taking ratios yields

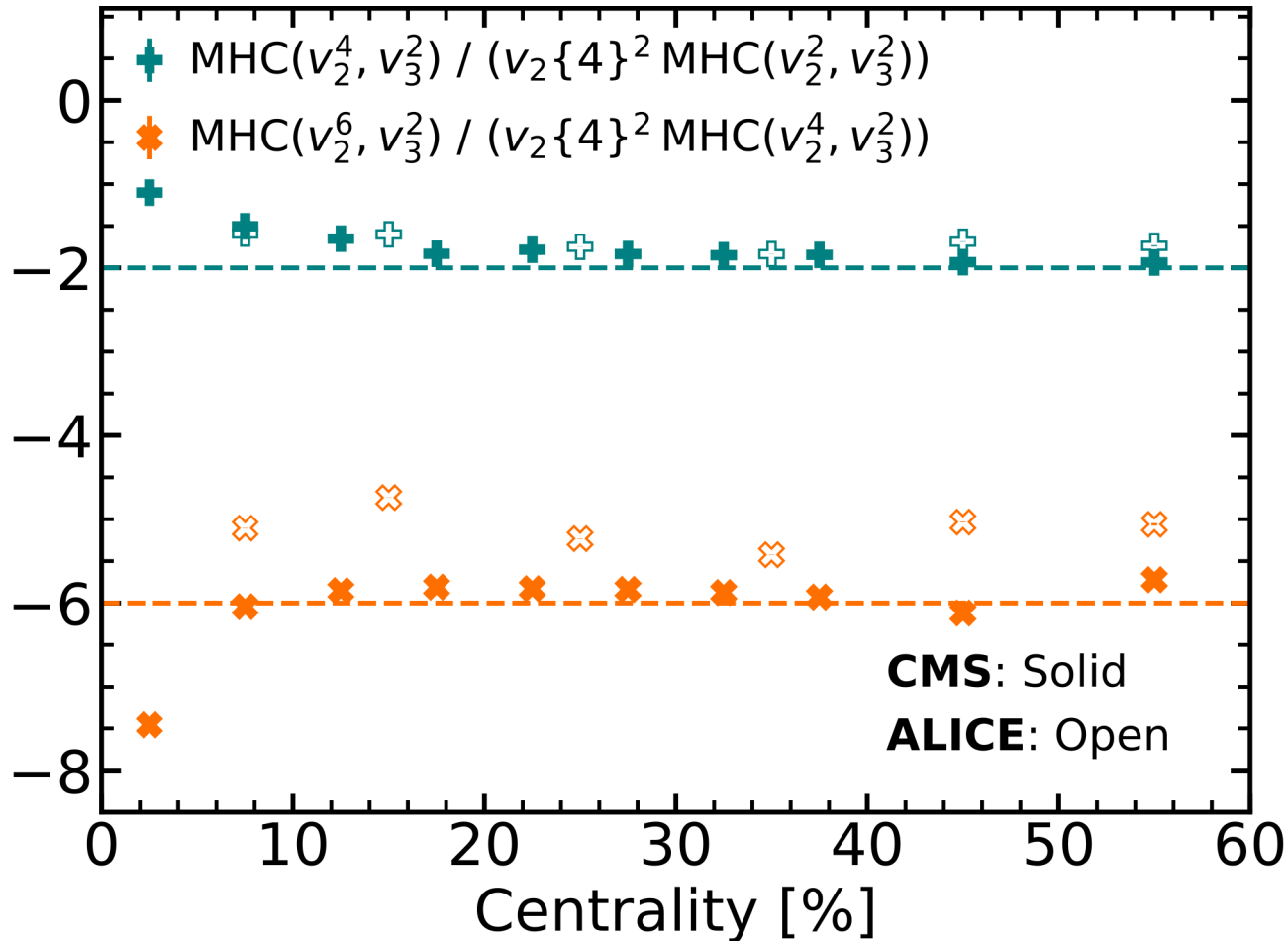
$$\frac{MHC(v_2^6, v_3^2)}{v_2\{4\}^2 MHC(v_2^4, v_3^2)} = -6 \quad \dots \quad \frac{MHC(v_2^4, v_3^2)}{v_2\{4\}^2 MHC(v_2^2, v_3^2)} \approx -2$$

Parameter-free predictions from c_{0111} dominance

*They can be generated systematically using the cumulant generating function

Predicted ratios vs. data

e.g.



- The predicted ratios are in good agreement with CMS data.
- ALICE data show deviations, possibly due to wider centrality bins.
- Data are broadly consistent with the intrinsic-cumulant hierarchy.

ALICE: [Phys. Lett. B 818 \(2021\) 136354](#).

CMS: [Phys. Lett. B 876 \(2026\) 140359](#).

F. G. Gardim et al [1608.02982](#)

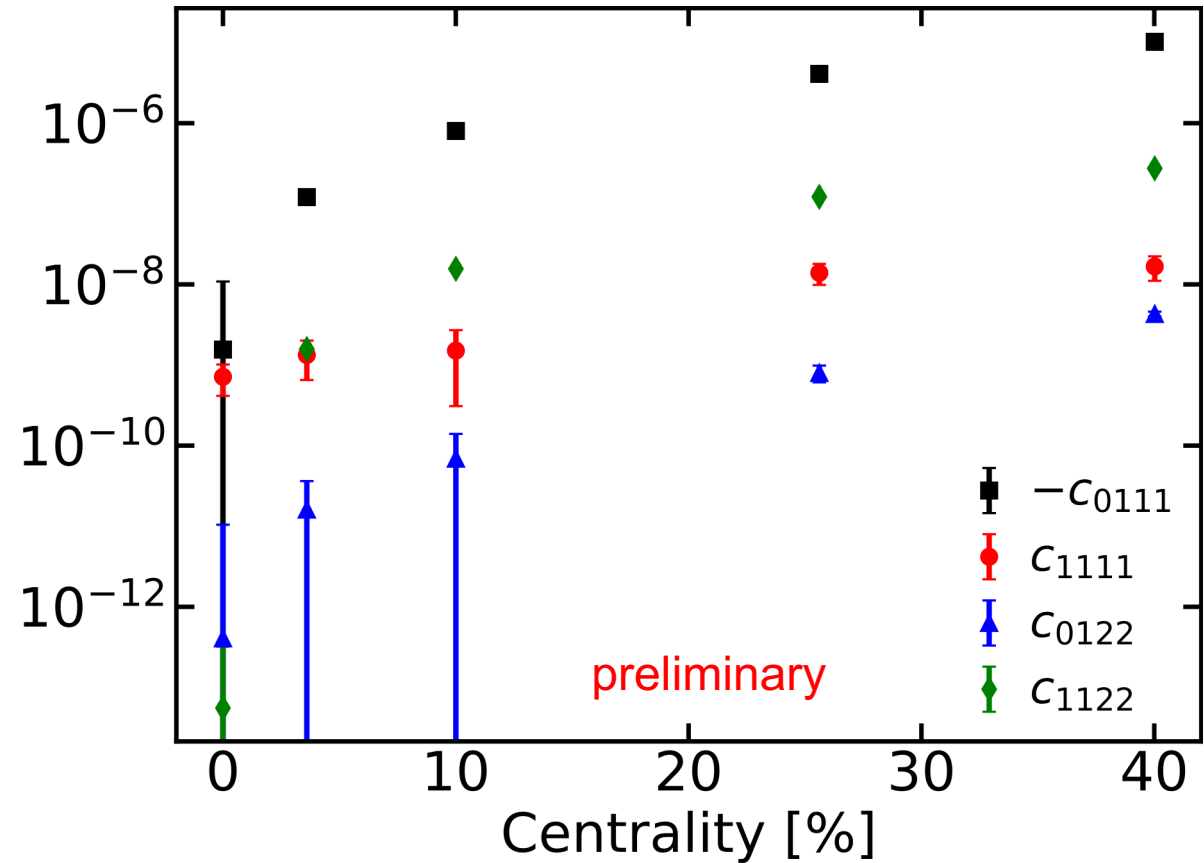
Hierarchy of intrinsic cumulants

How important are higher-order intrinsic cumulants?

- Fix b to eliminate centrality fluctuations.
- Run hydrodynamics at fixed b .
- Compute intrinsic cumulants c_{mpqr}^{int} .

Main Result:

- c_{0111} is the Dominant cumulant for $b > 0$.
- Higher-order cumulants are strongly suppressed by several orders of magnitude.



Work in progress with Andreas Kirchner and Jean-Yves Ollitrault.

Summary and outlook

- Correlations between v_2 and v_3 are driven by non-Gaussian properties of the joint distribution $P(v_2, v_3)$.
- The intrinsic frame reveals anisotropic information hidden by azimuthal averaging in the lab frame.
- Intrinsic cumulants exhibit a strong hierarchy:

$c_{0111} \gg$ higher-order intrinsic cumulants

- This hierarchy leads to simple relations between mixed-harmonic cumulants, e.g.

$$\frac{MHC(v_2^4, v_3^2)}{v_2\{4\}^2 MHC(v_2^2, v_3^2)} \approx -2$$

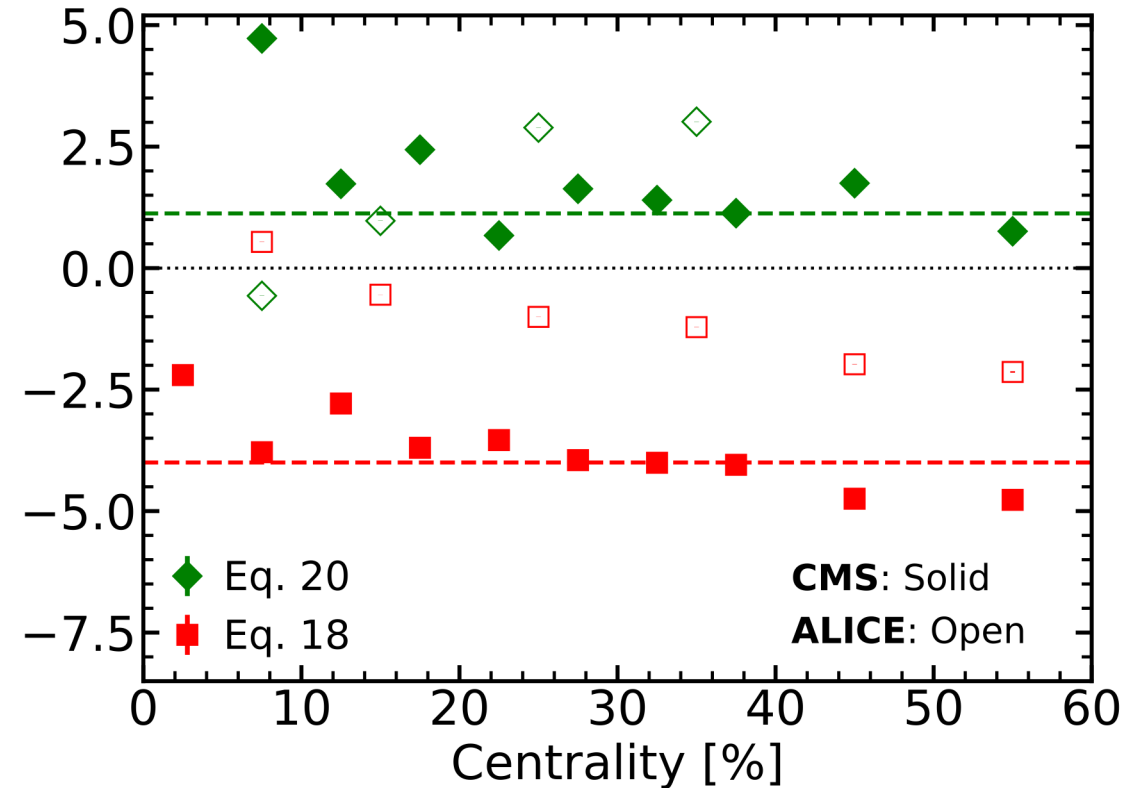
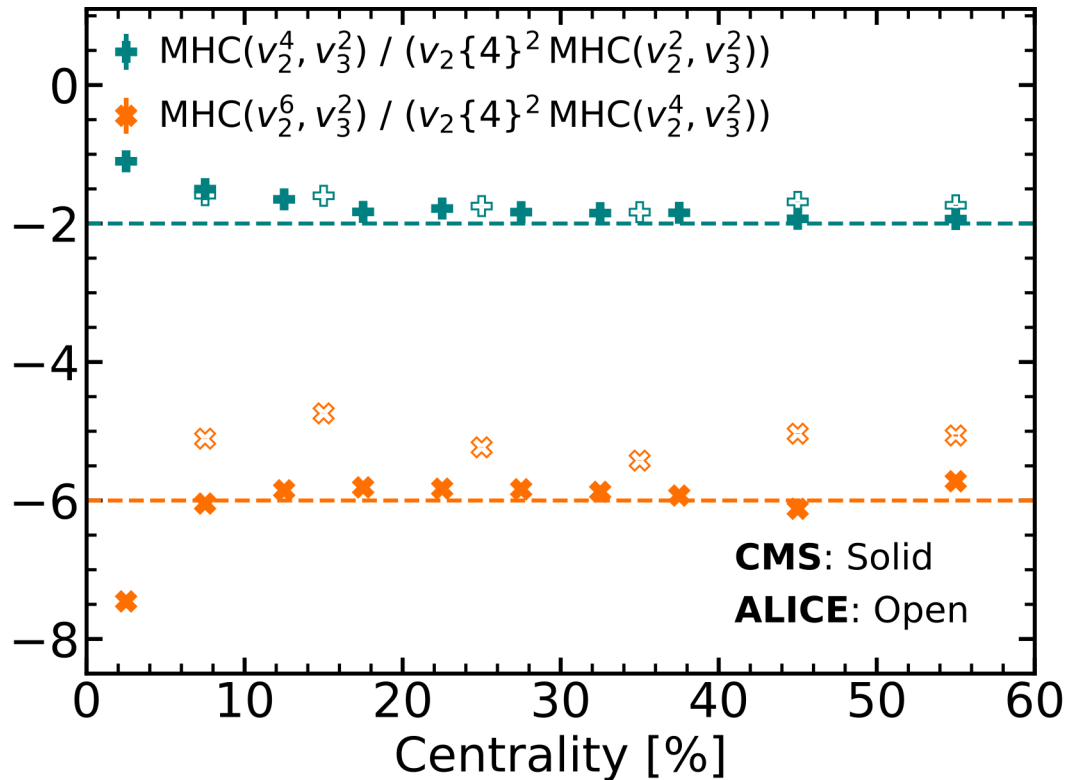
- Further explore intrinsic cumulants using hydrodynamic simulations at fixed b .

Work in progress with Andreas Kirchner and Jean-Yves Ollitrault.

- Extend the analysis to v_2 - v_4 correlations (not trivial!) Left for future work

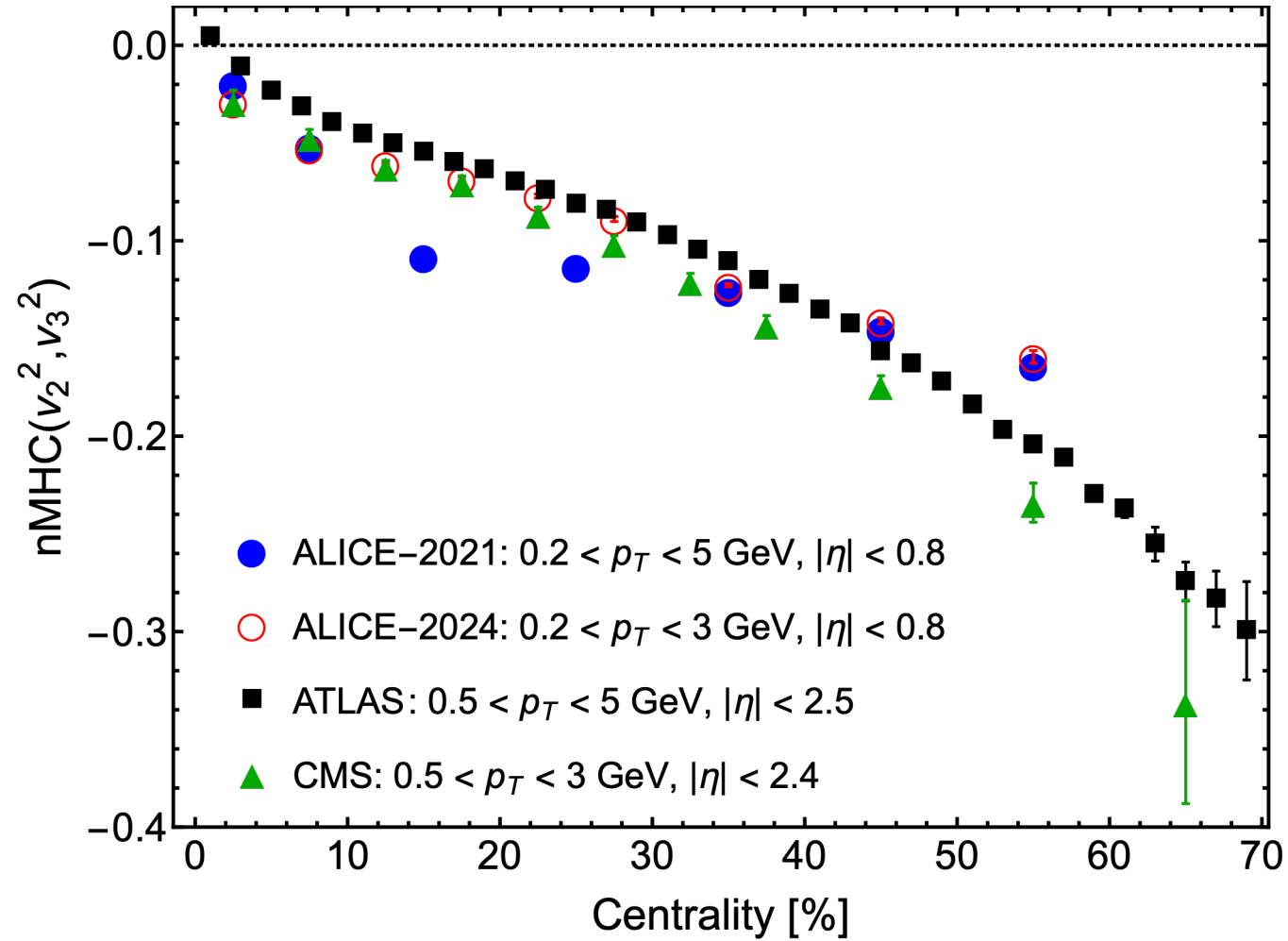
Backup Slides

Predicted ratios vs. data



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- ALICE data show deviations, possibly due to wider centrality bins.
- Data are broadly consistent with the intrinsic-cumulant hierarchy.

Comparison of experimental measurements



Estimating the hierarchy of intrinsic cumulants

- Two relevant sources of fluctuations: size and shape.
- **Size:** a cumulant of order k varies with N like N^{1-k} (Here $N \sim N_{part}$)
- **Shape:** for $\bar{V}_2 \ll 1$, intrinsic cumulants scale with powers of \bar{V}_2 according to

$$c_{mpqr} \sim \mathcal{O} \left(\mathcal{F}^{2(m+p+q+r-1)} \bar{V}_2^{\left| m-p + \frac{3}{2}(q-r) \right|} \right)$$

\mathcal{F} : typical magnitude of fluctuations

Consistency check: $c_{1100} \sim c_{0011} \sim \mathcal{O}(\mathcal{F}^2)$ (variance)

Predicted Hierarchy:

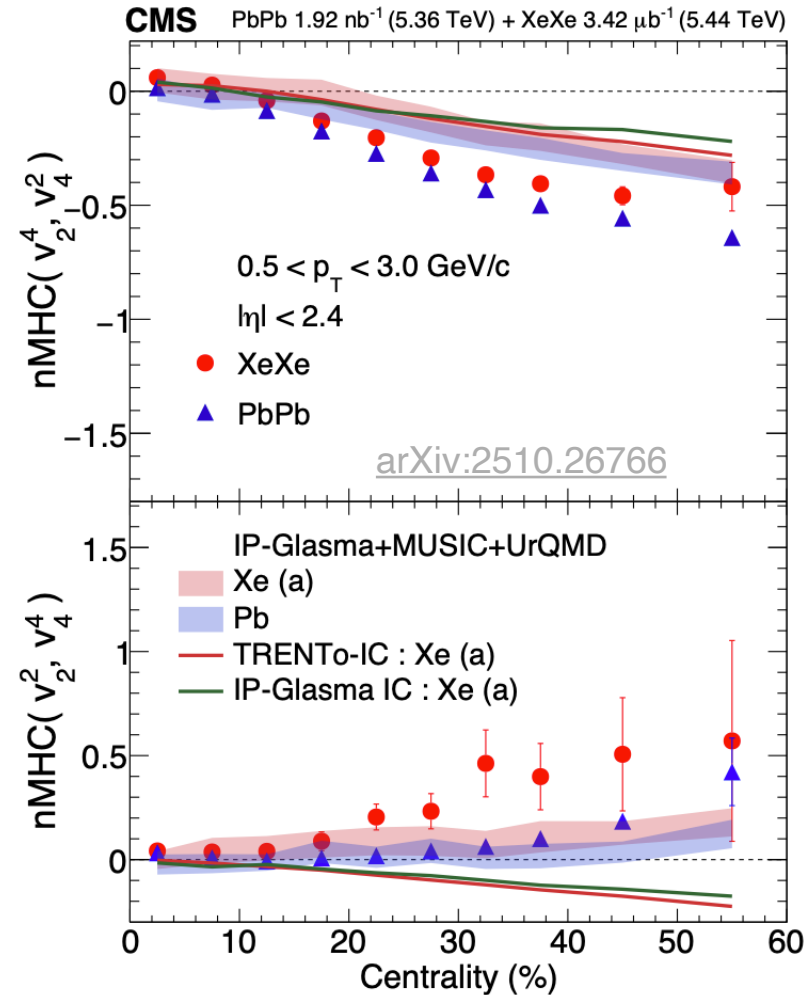
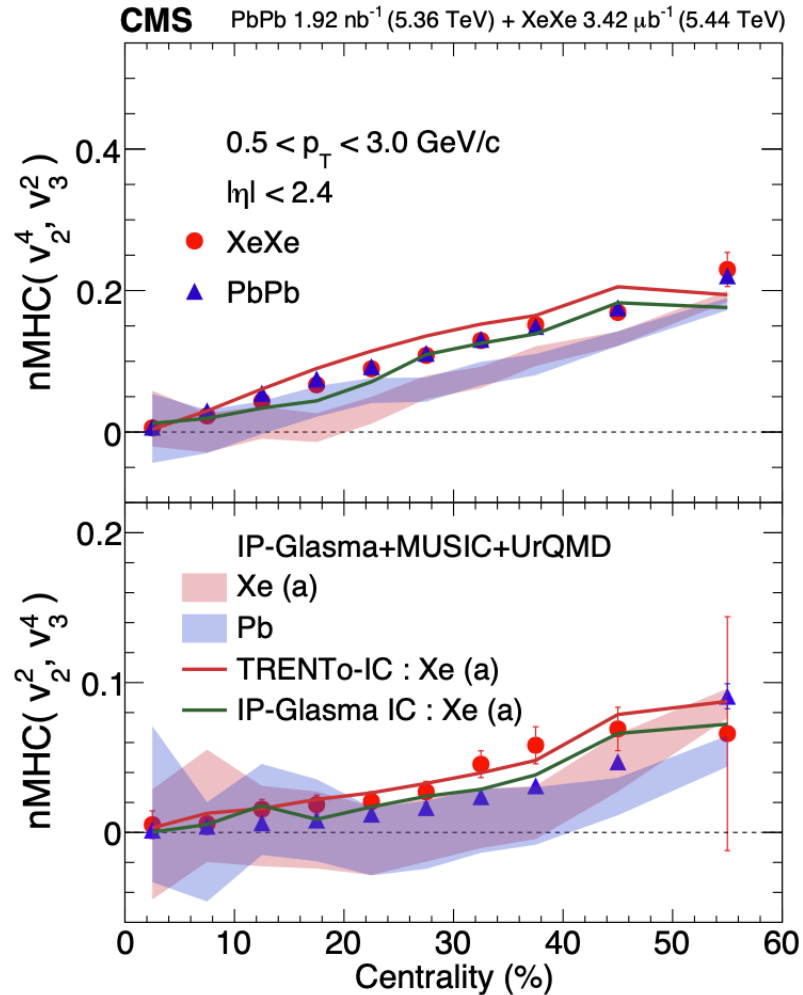
$$c_{0111} \sim \mathcal{O}(\mathcal{F}^4 \bar{V}_2) \sim \mathcal{O}(\mathcal{F}^5)$$

$$c_{0122} \sim \mathcal{O}(\mathcal{F}^8 \bar{V}_2)$$

...

$\Rightarrow c_{0111}$ is the dominant cumulant

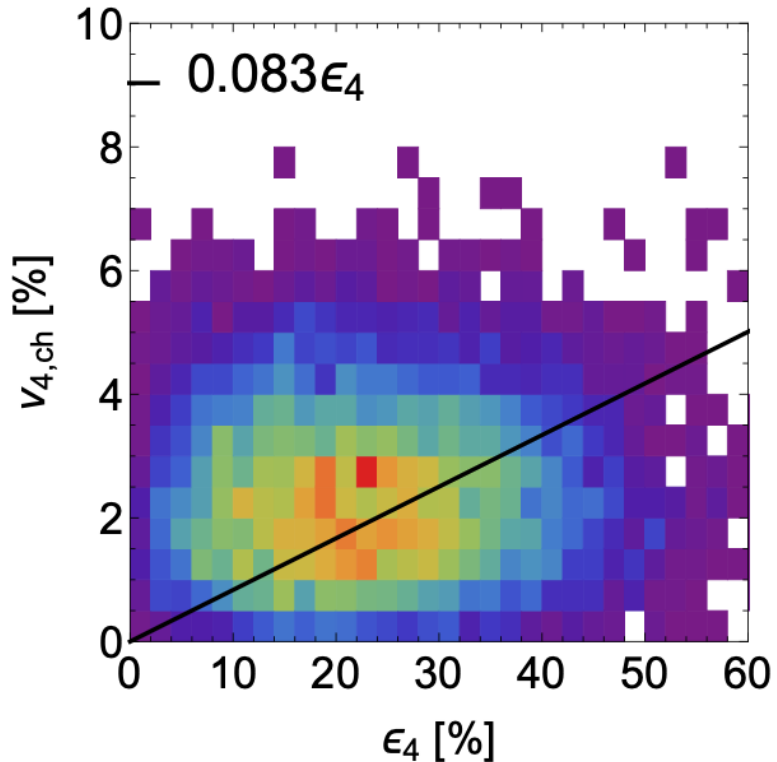
Mixed cumulants: Hydrodynamics vs Data



Overall agreement between hydrodynamics and CMS data for both XeXe and PbPb.

The v_4 - ϵ_4 relation is not purely linear due to nonlinear mode coupling.

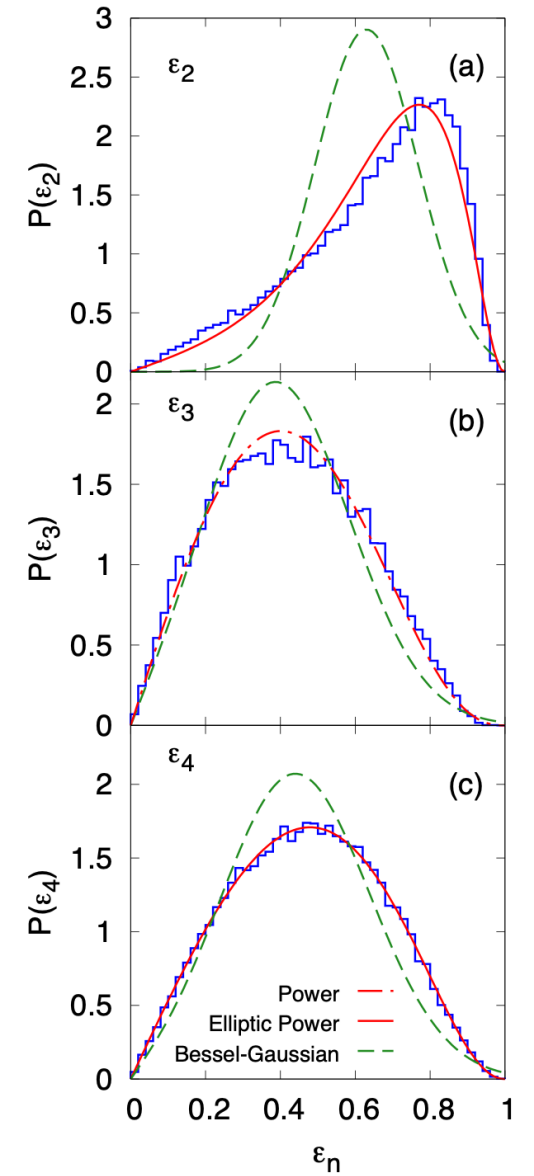
$$V_4 = V_{4L} + \chi_4 V_2^2$$



See e.g.

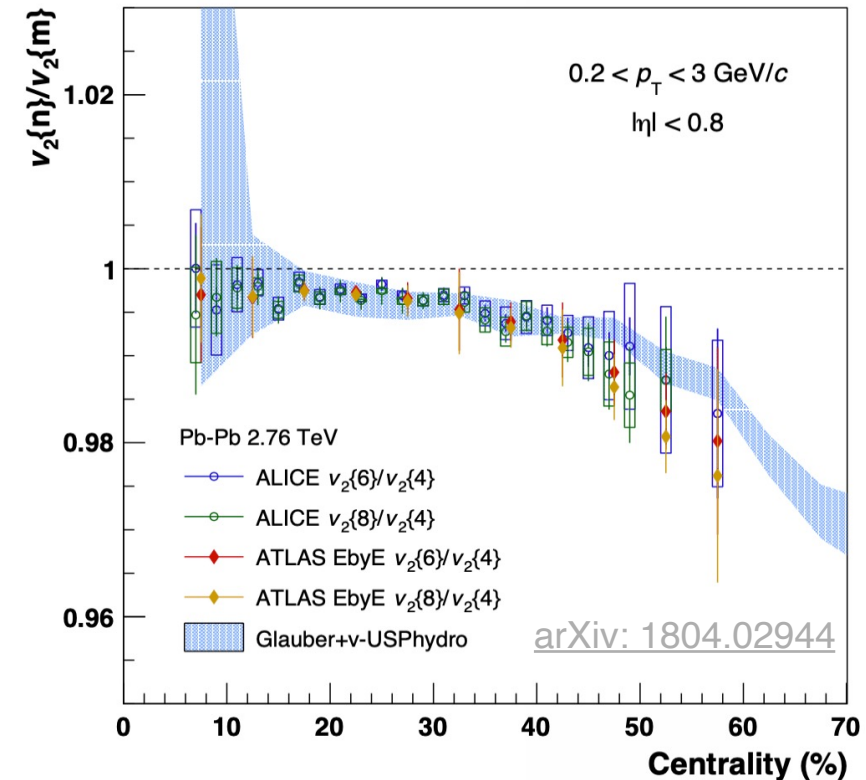
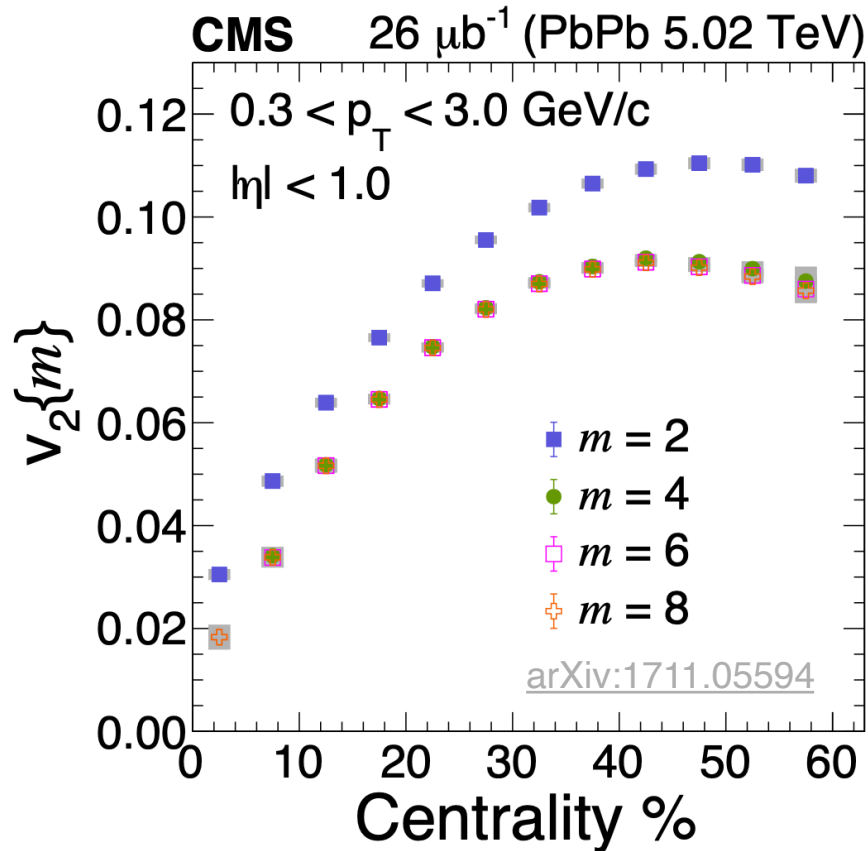
Teaney and Yan: [1206.1905](#)

Jacquelyn Noronha-Hostler et al. [1511.03896](#)



See Li Yan et al [1405.6595](#)

Non-Gaussian fluctuations



- Gaussian fluctuations imply $v_2\{4\} \simeq v_2\{6\} \simeq v_2\{8\}$
- Small deviations provide a sensitive probe

- A small splitting is observed ($< 1\%$).
- The splitting increases towards peripheral collisions
- Offers a sensitive test for hydrodynamic models