

Spectral Triples and AF algebras

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Based on arXiv:2207.04466:

T.Masson, G.N: Lifting Bratteli Diagrams between Krajewski Diagrams

Talk for the workshop Noncommutative geometry:
metric and spectral aspects

- Real spectral triple $(\mathcal{A}, \mathcal{H}_{\mathcal{A}}, D_{\mathcal{A}}, J_{\mathcal{A}}, \gamma_{\mathcal{A}})$.
- NCGFT which reproduces the standard model Lagrangian coupled to Gravity (Connes, Chamseddine, Marcolli, Suijlekom.. 1996 \rightarrow 2012)
- Based on the model of Almost-Commutative Manifold (A-C Manifold)

$$\widehat{\mathcal{A}} := C^{\infty}(M) \otimes \mathcal{A} \xrightarrow{S_{\mathcal{A}}} \text{NCGFT}_{\mathcal{A}}$$

\mathcal{A} : finite dimensional complex algebra.

$S_{\mathcal{A}}$: spectral action.

Purpose: Set up a general formalism to build Grand Unified Theories (GUTs) beyond the Standard Model of particle physics (SMPP) in the NCGFT framework using AF-algebras for \mathcal{A} .

Non-Commutative Gauge Field Theory (NCGFT) beyond the Standard Model of Particle Physics (SMPP)

- AF -algebra \mathcal{A} : inductive limit of a sequence of finite-dimensional algebras:

$$\mathcal{A}_1 \hookrightarrow \mathcal{A}_2 \hookrightarrow \dots \hookrightarrow \mathcal{A}_n \hookrightarrow \dots$$

- Embedding structure encoded by $\phi: \mathcal{A}_n \xrightarrow{\phi_n} \mathcal{A}_{n+1}$
- Given $\widehat{\mathcal{A}} := C^\infty(M) \otimes \mathcal{A}$ and $\widehat{\mathcal{B}} := C^\infty(M) \otimes \mathcal{B}$ with $\mathcal{A} \xrightarrow{\phi} \mathcal{B}$
 \rightsquigarrow How $\text{NCGFT}_{\mathcal{A}}$ and $\text{NCGFT}_{\mathcal{B}}$ are connected?
- If $\text{NCGFT}_{\mathcal{A}}$ gives the SMPP, how $\text{NCGFT}_{\mathcal{B}}$ can represent a GUT extension of the SMPP.

$$\text{NCGFT}_{\mathcal{A}_1} \hookrightarrow \text{NCGFT}_{\mathcal{A}_2} \hookrightarrow \dots \hookrightarrow \text{NCGFT}_{\mathcal{A}_n} \hookrightarrow \dots$$

- \rightarrow New way to create models beyond the SMPP.
- T.Masson, G.N: Lifting Bratteli Diagrams between Krajewski Diagrams: Spectral Triples, Spectral Actions, and AF algebras (arXiv:2207.04466)

Essentials on Krajewski diagram

Notations for Algebras:

- $\mathcal{A} = \bigoplus_{i=1}^r M_{n_i}(\mathbb{C}) \quad \rightsquigarrow \quad \mathcal{A}^e = \bigoplus_{i,j=1}^r M_{n_i} \otimes M_{n_j}^\circ$
- $\iota^i(\mathcal{A}_i)$ the inclusion map

Notations for Modules and Bimodules:

- $M_{n_i}(\mathbb{C})$ act on the irrep \mathbb{C}^{n_i} .
- $\Lambda := \{n_1, \dots, n_r\}$ each element n_i corresponding to the irrep \mathbb{C}^{n_i}
- $\widehat{\mathcal{H}}_{n_i n_j} := \iota^i(\mathcal{A}_i) \iota^j(\mathcal{A}_j)^\circ \mathcal{H} \quad \mathcal{H} = \bigoplus_{i,j=1}^r \widehat{\mathcal{H}}_{n_i n_j}$
- $\mathcal{H}_{n_i n_j} = \mathbb{C}^{n_i} \otimes \mathbb{C}^{n_j^\circ} \quad \rightsquigarrow \quad \widehat{\mathcal{H}}_{n_i n_j} \simeq \mathcal{H}_{n_i n_j} \otimes \mathbb{C}^{\mu_{ij}} \simeq \mathbb{C}^{n_i} \otimes \mathbb{C}^{\mu_{ij}} \otimes \mathbb{C}^{n_j^\circ}$
- μ_{ij} is the multiplicity of the irrep $\mathcal{H}_{n_i n_j}$.
- orthonormal basis $\{\sigma_{ij}^p\}_{1 \leq p \leq \mu_{ij}}$ of $\mathbb{C}^{\mu_{ij}}$ with p for the irreps

Essentials on Krajewski diagram : vertex

- The set of vertex $\Gamma^{(0)}$ of the graph is equipped with a map $\pi_{\lambda\rho} : \Gamma^{(0)} \rightarrow \Lambda \times \Lambda$ such that $\pi_{\lambda\rho}(v) = (n_i, n_j)$
- $\Gamma_{n_i n_j}^{(0)} := \{v \in \Gamma^{(0)} \mid \pi_{\lambda\rho}(v) = (n_i, n_j)\} \rightsquigarrow \Gamma^{(0)} := \cup_{i,j=1}^r \Gamma_{n_i n_j}^{(0)}$
- $\sigma_{ij}^p \in \mathbb{C}^{\mu_{ij}} \rightarrow \mathcal{H}_v := \text{Span}\{\xi_i \otimes \sigma_{ij}^p \otimes \eta_j^\circ \mid \xi_i \in \mathbb{C}^{n_i}, \eta_j^\circ \in \mathbb{C}^{n_j^\circ}\} \simeq \mathcal{H}_{n_i n_j}$
 $\rightsquigarrow \mathcal{H} = \oplus_{v \in \Gamma^{(0)}} \mathcal{H}_v$
- The element $\pi_{\lambda\rho}(v) \in \Lambda \times \Lambda$ is a decoration of the vertex v .
- γ induce sign decoration $s(v) = \pm 1$ of $v \rightsquigarrow \mathcal{H}_A = \mathcal{H}_A^+ \oplus \mathcal{H}_A^-$
- $J : \mathcal{H}_v \rightarrow \mathcal{H}_{\kappa(v)} \quad \kappa : \Gamma_{n_i n_j}^{(0)} \rightarrow \Gamma_{n_j n_i}^{(0)} \quad \widehat{\kappa}_v : \mathcal{H}_v \rightarrow \mathcal{H}_{\kappa(v)}$

Essentials on Krajewski diagram : edges

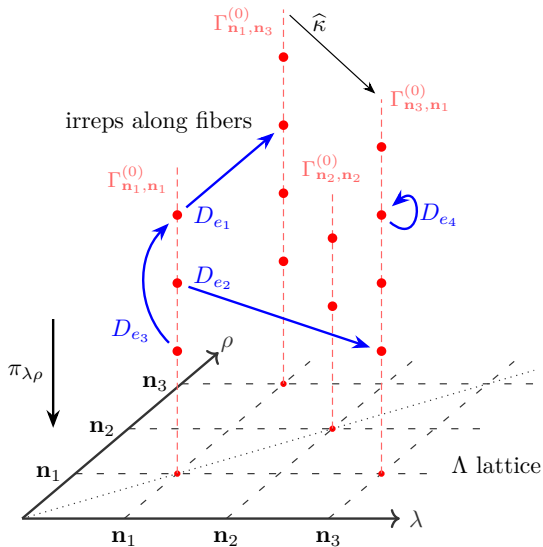
- The space $\Gamma^{(1)} \subset \Gamma^{(0)} \times \Gamma^{(0)}$ of edges \equiv couples $e = (v_1, v_2)$
- Given $e = (v_1, v_2) \in \Gamma^{(0)} \times \Gamma^{(0)}$ we have $D_e : \mathcal{H}_{v_1} \rightarrow \mathcal{H}_{v_2}$ such that:

$$\widehat{D}_e(\xi \otimes \sigma_v \otimes \eta^\circ) = \begin{cases} 0 & \text{if } v \neq v_1 \\ (D_{L,e}^{(1)}\xi) \otimes \sigma_{v_2} \otimes (D_{R,e}^{(2)}\eta^\circ) & \text{if } v = v_1 \end{cases}$$

(1, 2) \equiv notation for finite summation

- $D_e : \mathcal{H}_{v_1} \rightarrow \mathcal{H}_{v_2}$ defines a decoration of e .
- $\gamma D = -\gamma D \rightarrow s(v_2) = -s(v_1)$
- $D^\dagger = D$ is equivalent to $D_{\bar{e}} = D_e^\dagger$ with $\bar{e} := (v_2, v_1)$

The Krajewski diagram



$$\lambda : \Gamma_{n_i}^{(0)} \rightarrow \Lambda$$

$$\lambda(v) := n_i$$

$$\rho(v) := n_j \circ$$

$$\mathbb{C}^{\lambda(v)} \otimes \mathbb{C}^{\rho(v)} = \mathbb{C}^{n_i} \otimes \mathbb{C}^{n_j \circ}$$

AF-algebras

- $\mathcal{A} = \bigoplus_{i=1}^r M_{n_i}$

- $\forall a \in \mathcal{A} : a = \bigoplus_{i=1}^r a_i$

$$\mathcal{B} = \bigoplus_{k=1}^s M_{m_k}$$

$$a_i \in M_{n_i}$$

$$\phi_k(a) := \pi_k^{\mathcal{B}} \circ \phi(a) = \begin{pmatrix} a_1 \otimes \mathbb{1}_{\alpha_{k1}} & 0 & \cdots & 0 & 0 \\ 0 & a_2 \otimes \mathbb{1}_{\alpha_{k2}} & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & a_r \otimes \mathbb{1}_{\alpha_{kr}} & 0 \\ 0 & 0 & \cdots & 0 & 0_{n_{0,k}} \end{pmatrix} \subset M_{m_k}$$

$$a_i \otimes \mathbb{1}_{\alpha_{ki}} = \left. \begin{pmatrix} a_i & 0 & 0 & 0 \\ 0 & a_i & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & a_i \end{pmatrix} \right\} \alpha_{ki} \text{ times.}$$

- α_{ki} is the multiplicity of the embedding.

$$\phi_k^i(a_i) = \begin{pmatrix} 0 & \cdots & 0 & 0 & 0 & \cdots & 0 \\ \vdots & \ddots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & \cdots & 0 & 0 & 0 & \cdots & 0 \\ 0 & \cdots & 0 & a_i \otimes \mathbb{1}_{\alpha_{ki}} & 0 & \cdots & 0 \\ 0 & \cdots & 0 & 0 & 0 & \cdots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & \cdots & 0 & 0 & 0 & \cdots & 0 \end{pmatrix} = \sum_{\alpha=1}^{\alpha_{ki}} \phi_{k,\alpha}^i(a_i) : M_{n_i} \rightarrow M_{m_k}$$

- The n_i 's designate vertices in Bratteli \rightsquigarrow matrix blocs M_{n_i} .
- Edges (n_i, m_k) between two vertices of $(\mathcal{A}, \mathcal{B})$ exist iff $\alpha_{ki} \neq 0$.
- **The multiplicities α_{ki} define the Bratteli diagram of the AF-algebra.**

ϕ -compatibility condition

- ϕ -compatibility condition \rightarrow link between $\text{NCGFT}_{\mathcal{A}}$'s and $\text{NCGFT}_{\mathcal{B}}$'s structure at the level of their spectral triples.
 \rightsquigarrow **Physical meaning: conservation of actions on embedded Hilberts spaces**
- ϕ -compatibility of Hilbert spaces: $\phi_{\mathcal{H}} : \mathcal{H}_{\mathcal{A}} \rightarrow \mathcal{H}_{\mathcal{B}}$ is ϕ -compatible if $\phi_{\mathcal{H}}(a\psi) = \phi(a)\phi_{\mathcal{H}}(\psi)$ for any $(a, \psi) \in (\mathcal{A}, \mathcal{H}_{\mathcal{A}})$
- $\mathcal{H}_{\mathcal{B}} = \phi_{\mathcal{H}}(\mathcal{H}_{\mathcal{A}}) \oplus \phi_{\mathcal{H}}(\mathcal{H}_{\mathcal{A}})^{\perp} \rightarrow \forall B \text{ on } \mathcal{H}_{\mathcal{B}}: B = \begin{pmatrix} B_{\phi}^{\phi} & B_{\perp}^{\phi} \\ B_{\phi}^{\perp} & B_{\perp}^{\perp} \end{pmatrix}.$
- Given A on $\mathcal{H}_{\mathcal{A}}$ and B on $\mathcal{H}_{\mathcal{B}}$ we said that they are:
 - ϕ -compatible if $\forall \psi \in \mathcal{H}_{\mathcal{A}}: \phi_{\mathcal{H}}(A\psi) = B_{\phi}^{\phi}\phi_{\mathcal{H}}(\psi)$
 \equiv equality in $\phi_{\mathcal{H}}(\mathcal{H}_{\mathcal{A}})$
 - Strong ϕ -compatible if $\forall \psi \in \mathcal{H}_{\mathcal{A}}: \phi_{\mathcal{H}}(A\psi) = B\phi_{\mathcal{H}}(\psi)$
 \equiv equality in $\mathcal{H}_{\mathcal{B}}$
- ϕ -compatibility looks more natural for physics since it is based on constraints on inherited degrees of freedom (dofs) only.

Consequences of ϕ -compatibility on real spectral Triples

- Strong ϕ -compatibility implies ϕ -compatibility.
- $(\mathcal{A}, \mathcal{H}_{\mathcal{A}}, D_{\mathcal{A}}, J_{\mathcal{A}}, \gamma_{\mathcal{A}})$ and $(\mathcal{B}, \mathcal{H}_{\mathcal{B}}, D_{\mathcal{B}}, J_{\mathcal{B}}, \gamma_{\mathcal{B}})$ are said to be ϕ -compatible if each of its structure are.
→ lift of arrows in a Bratteli diagram to arrows between Krajewski diagrams.
- If two real spectral triples are strong ϕ -compatible → same KO -dimension (mod 8).
- If two real spectral triples are ϕ -compatible and $J_{\mathcal{B}}$ is strong ϕ -compatible with $J_{\mathcal{A}}$ → same KO -dimension (mod 8).
- If A on $\mathcal{H}_{\mathcal{A}}$ and B on $\mathcal{H}_{\mathcal{B}}$ are strong ϕ -compatible then $B_{\perp}^{\phi} = 0$.
- $\forall a \in \mathcal{A}$, $\pi_{\mathcal{B}} \circ \phi(a)$ on $\mathcal{H}_{\mathcal{B}}$ reduces to
$$\begin{pmatrix} \pi_{\mathcal{B}} \circ \phi(a)_{\phi}^{\phi} & 0 \\ 0 & \pi_{\mathcal{B}} \circ \phi(a)_{\perp}^{\phi} \end{pmatrix}$$
- Gauge transformations are preserved under certain circumstances.

Embedding structure of AF algebras

- Let $\mathcal{A} = \mathcal{A}_n$ and $\mathcal{B} = \mathcal{A}_{n+1}$, and (v, w) be nodes in $(\Gamma_{\mathcal{A}}^{(0)}, \Gamma_{\mathcal{B}}^{(0)})$.
- $\phi_{\mathcal{H}}$ is injective.
- $\phi_{\mathcal{H}}$ is totally specified by the $\phi_{\mathcal{H},w}^v$'s.
- Taking $\pi_{\lambda\rho}(v) = (n_i, n_j)$ and $\pi_{\lambda\rho}(w) = (m_k, m_l)$.

Embedding structure of AF algebras

$\phi_{\mathcal{H},w}^v$ reduces to $\mathbb{C}^{n_i} \otimes \mathbb{C}^{n_j^\circ} \rightarrow \mathbb{C}^{n_i} \otimes \mathbb{C}^{\alpha_{ki}} \otimes \mathbb{C}^{\alpha_{lj}} \otimes \mathbb{C}^{n_j^\circ}$.

Then it reduces to a linear map $\mathbb{C} \rightarrow \mathbb{C}^{\alpha_{ki}} \otimes \mathbb{C}^{\alpha_{lj}}$.

Therefore to an element $\mathbf{u}(\mathbf{v},\mathbf{w}) \in \mathbb{C}^{\alpha_{ki}} \otimes \mathbb{C}^{\alpha_{lj}}$, such that we have:

$$\phi_{\mathcal{H},w}^v(\xi_i \otimes \eta_j^\circ) = I_{k,\ell}^{i,j}(\xi_i \otimes \mathbf{u}(\mathbf{v},\mathbf{w}) \otimes \eta_j^\circ) \quad \text{for any } \xi_i \otimes \eta_j^\circ \in \mathcal{H}_{A,v}.$$

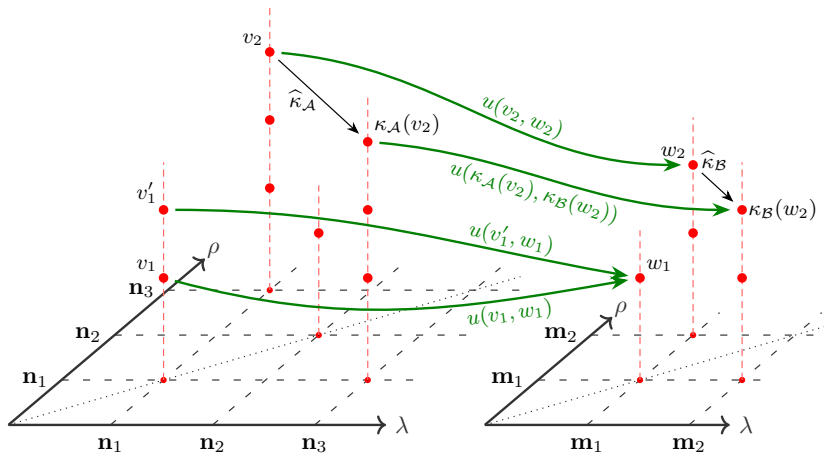
with the inclusion:

$$I_{k,\ell}^{i,j} : \mathbb{C}^{n_i} \otimes \mathbb{C}^{\alpha_{ki}} \otimes \mathbb{C}^{\alpha_{lj}} \otimes \mathbb{C}^{n_j^\circ} \hookrightarrow \mathbb{C}^{m_k} \otimes \mathbb{C}^{m_\ell^\circ}$$

Then $\phi_{\mathcal{H},w}^v$ is completely determined by $\mathbf{u}(\mathbf{v},\mathbf{w}) \in M_{\alpha_{ki} \times \alpha_{lj}} \simeq \mathbb{C}^{\alpha_{ki}} \otimes \mathbb{C}^{\alpha_{lj}}$

Embedding structure of AF algebras

The $u(v, w)$ contain all the data of edges in the Bratteli Diagram, thus linking Krajewski and Bratteli diagram structures:



Link between scalar products

Scalar products on \mathcal{H}_A and \mathcal{H}_B are linked by the relation:

$$\langle \phi_{\mathcal{H}}^{v_1}(\psi_{v_1}), \phi_{\mathcal{H}}^{v_2}(\psi'_{v_2}) \rangle_{\mathcal{H}_B} = \langle \psi_{v_1}, \iota_{v_1}^{v_2}(\psi'_{v_2}) \rangle_{\mathcal{H}_{A,v_1}} T^{v_1, v_2}$$

$$\text{with } T^{v_1, v_2} := \begin{cases} 0 & \text{if } \pi_{\lambda\rho}(v_1) \neq \pi_{\lambda\rho}(v_2) \\ \sum_{w \in \Gamma_B^{(0)}} \text{tr}(u(v_1, w)^* u(v_2, w)) & \text{if } \pi_{\lambda\rho}(v_1) = \pi_{\lambda\rho}(v_2) \end{cases}$$

There are orthonormal bases $\{\sigma_{ij}^p\}_{1 \leq p \leq \mu_{ij}}$ of $\mathbb{C}^{\mu_{ij}}$ such that:

- $\langle \phi_{\mathcal{H}}^{v_1}(\psi_{v_1}), \phi_{\mathcal{H}}^{v_2}(\psi'_{v_2}) \rangle_{\mathcal{H}_B} = 0$ if $v_1 \neq v_2$.
- $\langle \phi_{\mathcal{H}}^v(\psi_v), \phi_{\mathcal{H}}^v(\psi'_v) \rangle_{\mathcal{H}_B} = t_v \langle \psi_v, \psi'_v \rangle_{\mathcal{H}_{A,v}} \quad \forall v.$
- $t_{\kappa_A(v)} = t_v$ is a real number $\quad \forall v.$

Normalized $\phi_{\mathcal{H}}$ map: $\phi_{\mathcal{H}}(\bigoplus_{v \in \Gamma_A^{(0)}} \psi_v) := \sum_{v \in \Gamma_A^{(0)}} t_{v(\tilde{v})}^{-1/2} \phi_{\mathcal{H}}^{0,v}(\psi_v)$

\rightsquigarrow More natural because it preserve scalar products.

\rightsquigarrow Idea of dilution of the dofs.

The Corresponding AC-Manifold

- $\widehat{\mathcal{A}} := C^\infty(M) \otimes \mathcal{A}$ and $\widehat{\mathcal{B}} := C^\infty(M) \otimes \mathcal{B}$ are said ϕ -compatible if \mathcal{A} and \mathcal{B} are.
- $ST_{\mathcal{A}} = (\widehat{\mathcal{A}} := C^\infty(M) \otimes \mathcal{A}, \mathcal{H}_{\widehat{\mathcal{A}}} := L^2(S) \otimes \mathcal{H}_{\mathcal{A}}, D_{\widehat{\mathcal{A}}} := D_M \otimes 1 + J_M \otimes D_{\mathcal{A}}, J_{\widehat{\mathcal{A}}} := J_M \otimes J_{\mathcal{A}}, \gamma_{\widehat{\mathcal{A}}} := \gamma_M \otimes \gamma_{\mathcal{A}})$
- $ST_{\mathcal{B}} = (\widehat{\mathcal{B}}, \mathcal{H}_{\widehat{\mathcal{B}}}, D_{\widehat{\mathcal{B}}} := D_M \otimes 1 + J_M \otimes D_{\mathcal{B}}, J_M \otimes J_{\mathcal{B}}, \gamma_M \otimes \gamma_{\mathcal{B}})$
- Fluctuated Dirac operator:

$$D_{\widehat{\mathcal{A}}, \omega} = D_M \otimes 1 + \gamma^\mu \otimes B_\mu + \gamma_M \otimes \Phi$$

- B_μ and Φ are the usual gauge connections and Higgs fields.
- ϕ -compatibility condition is taken on ω (it's equivalent to take it on B_μ and Φ).

Bosonic and Fermionic's Lagrangian

- Given the spectral triple $ST_{\mathcal{A}}$ the associated action is:

$$S_{\mathcal{A}}[\omega, \tilde{\psi}] = S_{b,\mathcal{A}}[\omega] + S_{f,\mathcal{A}}[\omega, \tilde{\psi}]$$

with the Bosonic spectral action and Fermionic action given by:

$$S_{b,\mathcal{A}}[\omega] := \text{Tr} f(D_{\hat{\mathcal{A}},\omega}/\Lambda) \qquad S_{f,\mathcal{A}}[\omega, \tilde{\psi}] := \langle J_{\hat{\mathcal{A}}}\tilde{\psi}, D_{\hat{\mathcal{A}},\omega}\tilde{\psi} \rangle_{\tilde{\mathcal{H}}_{\hat{\mathcal{A}}}}$$

Remark

NCSMPP: $\dim(M) = 4$, and KO dimension of \mathcal{A} (then \mathcal{B}) is 6.

Comparison between the Lagrangian of \mathcal{A} and \mathcal{B}

If $(\omega, \tilde{\psi})$ and $(\omega', \tilde{\psi}')$ are ϕ -compatible, the action of $ST_{\mathcal{B}}$ is:

$$S_{\mathcal{B}}[\omega', \tilde{\psi}'] = S_{\mathcal{A}}[\omega, \tilde{\psi}] + \text{TNIC}$$

TNIC \equiv Terms with Non-Inherited Components.

The Spectral action of \mathcal{A} is recovered, new terms that correspond to the mixing with new dofs appear.

This structural result works in general KO dimension for \mathcal{A} (then \mathcal{B}).

\equiv Constraints of a NCGFT on \mathcal{B} with another on \mathcal{A}

Perspective to create models beyond the SMPP

- General framework for building phenomenological models.
- Explore the way in which the interplay between old and new dofs permits to extend the NCSMPP.
- In arXiv:2106.08358 (Derivation-based noncommutative field theories on AF algebras) we study the phenomenology of these embeddings:
 - $M_2 \hookrightarrow M_3$
 - $M_2 \oplus M_2 \hookrightarrow M_4$
 - $M_2 \oplus M_2 \hookrightarrow M_5$
 - $M_2 \oplus M_3 \hookrightarrow M_5 \quad \sim \sim \quad$ Georgi-Glashow model
- Many possibilities to do unified theories in different ways.

Thank you

Suggested readings:

- WD Van Suijlekom: Noncommutative geometry and particle physics.
- KR Davidson: C^* -algebras by example.
- T Masson, G Nieuviarts: Derivation-based noncommutative field theories on algebras.
- T Masson, G Nieuviarts: Lifting Bratteli Diagrams between Krajewski Diagrams: Spectral Triples, Spectral Actions, and AF algebras.

Backup slides

The operator J

- $J_v := \epsilon(v, d) J_0 \widehat{\kappa}_v = \epsilon(v, d) \widehat{\kappa}_v J_0 : \mathcal{H}_v \rightarrow \mathcal{H}_{\kappa(v)}$

$$\epsilon(v, d) := \begin{cases} 1 & \text{for } i(v) < j(v), \\ \epsilon & \text{for } i(v) > j(v), \\ 1 & \text{for } i(v) = j(v) \text{ and } d = 0, 1, 7, \\ \epsilon^{\chi(v)} & \text{for } i(v) = j(v) \text{ and } d = 2, 3, 4, 5, 6. \end{cases} \quad (1)$$

- $J\gamma = \epsilon''\gamma J \quad \rightsquigarrow \quad s \circ \kappa = \epsilon''s \quad J^2 = \epsilon$

- $\phi^e : \mathcal{A}^e \rightarrow \mathcal{B}^e$ as $\phi^e := \phi \otimes \phi^\circ$, i.e. $\phi^e(a_1 \otimes a_2^\circ) = \phi(a_1) \otimes \phi^\circ(a_2^\circ)$

- $\phi_{\mathcal{H}}$ restricts to maps $\mathcal{H}_{\mathcal{A}}^\pm \rightarrow \mathcal{H}_{\mathcal{B}}^\pm$.

- $J_{\mathcal{B}}$ is strong ϕ -compatible with $J_{\mathcal{A}}$ iff $\forall (v, w) \in (\Gamma_{\mathcal{A}}^{(0)}, \Gamma_{\mathcal{B}}^{(0)})$:

$$u(\kappa_{\mathcal{A}}(v), \kappa_{\mathcal{B}}(w)) = \frac{\epsilon_{\mathcal{A}}(v, d_{\mathcal{A}})}{\epsilon_{\mathcal{B}}(w, d_{\mathcal{B}})} u(v, w)^* \quad (2)$$

where $d_{\mathcal{A}}$ (resp. $d_{\mathcal{B}}$) is the KO -dimension of \mathcal{A} (resp. \mathcal{B}).

- ϕ -compatibility for Hilbert spaces :

$$\phi_{\mathcal{H},w}^v(a_i b_j^\circ \psi_v) = \phi_k^i(a_i) \phi_\ell^j(b_j)^\circ \phi_{\mathcal{H},w}^v(\psi_v)$$

- ϕ -compatibility for operators :

$$\sum_{v_2 \in \Gamma_A^{(0)}} \phi_{\mathcal{H},w_2}^{v_2}(A_{v_2}^{v_1} \psi_{v_1}) = \sum_{w_1 \in \Gamma_B^{(0)}} B_{\phi,w_2}^{\phi,w_1} \phi_{\mathcal{H},w_1}^{v_1}(\psi_{v_1})$$

- $S_{b,A}[\omega] \sim \int_M \mathcal{L}(B_\mu, \Phi), d^4x + \mathcal{O}(\Lambda^{-1})$
- Bosonic Lagrangian : $\mathcal{L}(B_\mu, \Phi) = \mathcal{L}_B(B_\mu) + \mathcal{L}_\varphi(B_\mu, \Phi)$
 - $\mathcal{L}_B(B_\mu) = \frac{f(0)}{24\pi^2} \text{tr}(F_{\mu\nu}F^{\mu\nu})$
 - $\mathcal{L}_\varphi(B_\mu, \Phi) = -\frac{2f_2\Lambda^2}{4\pi^2} \text{tr}(\Phi^2) + \frac{f(0)}{8\pi^2} \text{tr}(\Phi^4) + \frac{f(0)}{8\pi^2} \text{tr}((D_\mu\Phi)(D^\mu\Phi))$