

Effective QCD - problem set 1
 10.10.2017. Tuesday 14:00
 room D-02-2

1. Canonical representation of Dirac matrices (or Bjorken-Drell) takes the following form:

$$\gamma^0 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}, \quad \gamma^k = \begin{bmatrix} 0 & \sigma^k \\ -\sigma^k & 0 \end{bmatrix}$$

where σ^k are 2×2 Pauli matrices. It is useful to use different representations. Weyl (or chiral) representation differs from the canonical one only by γ^0 :

$$\gamma_{\text{ch}}^0 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \gamma_{\text{ch}}^k = \gamma^k.$$

Another useful representation, known as Majorana representation, takes the following form:

$$\gamma_{\text{M}}^0 = \begin{bmatrix} 0 & \sigma^2 \\ \sigma^2 & 0 \end{bmatrix}, \quad \gamma_{\text{M}}^1 = i \begin{bmatrix} 0 & \sigma^3 \\ \sigma^3 & 0 \end{bmatrix}, \quad \gamma_{\text{M}}^2 = i \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}, \quad \gamma_{\text{M}}^3 = -i \begin{bmatrix} 0 & \sigma^1 \\ \sigma^1 & 0 \end{bmatrix}.$$

Note that $(\gamma_{\text{M}}^\mu)^* = -\gamma_{\text{M}}^\mu$. Check by explicit calculation that all these reps. satisfy anticommutation relations:

$$\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}.$$

Calculate $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$ in all these representations. Find transformation matrices X and Y defined as:

$$\gamma_{\text{ch}}^\mu = X\gamma^\mu X^{-1}, \quad \gamma_{\text{M}}^\mu = Y\gamma^\mu Y^{-1}.$$

2. Remember that originally Dirac equation has been written in a form

$$i\hbar \frac{\partial \psi}{\partial t} = (c\vec{\alpha} \cdot \vec{p} + \beta mc^2)\psi$$

where $\vec{p} = -i\hbar\vec{\nabla}$ is the momentum operator. Relation between γ matrices and β , $\vec{\alpha}$ matrices reads as follows:

$$\gamma^0 = \beta, \quad \gamma^k = \beta\alpha^k.$$

Calculate α^k in all the above representations of Dirac matrices.

3. Spin operator is related to the J generators of the Lorentz transformations

$$J^{\mu\nu} = \frac{1}{2}\sigma^{\mu\nu} = \frac{i}{4}[\gamma^\mu, \gamma^\nu]$$

and reads

$$S_i = \frac{1}{2}\varepsilon_{ijk}J^{jk}.$$

Calculate \vec{S} in chiral representation. Calculate \vec{S}^2 .

4. Helicity operator is defined as

$$h = 2 \frac{\vec{S} \cdot \vec{k}}{|\vec{k}|}.$$

Prove that

$$\gamma^5 h = \frac{\gamma^0 \vec{\gamma} \cdot \vec{k}}{|\vec{k}|}.$$

5. Let's denote Lorentz transformations (including boosts, rotations, space and time reflections) in a usual way

$$x'^{\mu} = \Lambda^{\mu}_{\nu} x^{\nu}.$$

In order to calculate Lorentz transformation of the spinors consider matrix X related to the space-time point x^{μ} :

$$X(x) = x_{\mu} \tilde{\sigma}^{\mu} = x^0 \sigma^0 + x^1 \sigma^1 + x^2 \sigma^2 + x^3 \sigma^3.$$

To this end it is useful to introduce the following notation

$$\sigma^{\mu} = (1, \vec{\sigma}), \quad \tilde{\sigma}^{\mu} = (1, -\vec{\sigma})$$

remembering that:

$$\partial_{\mu} = (\partial_t, \vec{\nabla}), \quad \partial^{\mu} = (\partial_t, -\vec{\nabla}).$$

Show that

$$\det X = x^{\mu} x_{\mu}.$$

Therefore Lorentz transformation Λ generates $SL(2, C)$ transformation of matrix X :

$$M^{\dagger} X' M = X$$

where $X' = X(x')$. These transformations preserve determinant. Show that

$$\Lambda^{\mu}_{\nu} = \frac{1}{2} \text{Tr} (\tilde{\sigma}^{\nu} M^{\dagger} \tilde{\sigma}^{\mu} M) = \frac{1}{2} \text{Tr} (\sigma_{\mu} M^{\dagger} \tilde{\sigma}^{\mu} M).$$

Note that M are defined up to a phase. The same reasoning can be repeated for matrix $Y(x)$:

$$Y(x) = x_{\mu} \sigma^{\mu}$$

where the pertinent transformation matrix N is defined as

$$N^{\dagger} Y' N = Y.$$

Express Lorentz transformation Λ in terms of N . Prove that $N M^{\dagger} = 1$.

Calculate Λ for

$$M = N = \begin{bmatrix} e^{i\theta/2} & 0 \\ 0 & e^{-i\theta/2} \end{bmatrix}$$

and

$$M = \begin{bmatrix} e^{\theta/2} & 0 \\ 0 & e^{-\theta/2} \end{bmatrix}, \quad N = \begin{bmatrix} e^{-\theta/2} & 0 \\ 0 & e^{\theta/2} \end{bmatrix}$$

where $v/c = \tanh \theta$.