

QCD lecture 6b

November 15

Axial anomaly

In massless QCD there are two conserved currents

$$\bar{u}(p')\gamma^\mu u(p) \quad \bar{u}(p')\gamma^\mu\gamma_5 u(p)$$

We will show that when loop corrections are included only one current remains conserved. We know from gauge invariance that the vector current must be conserved, so it is the axial-vector current that is not conserved.

This phenomenon is called axial anomaly.

Conserved currents

$$\begin{aligned}q_\mu j^\mu(q) &= q_\mu \bar{u}(p') \gamma^\mu u(p) \\ &= \bar{u}(p') (\not{p}' - \not{p}) u(p) \\ &= 0\end{aligned}$$

$$\gamma_5 = i\gamma^0\gamma^1\gamma^2\gamma^3 \quad \{\gamma^\mu, \gamma_5\} = 0$$

$$\begin{aligned}q_\mu j^\mu(q) &= q_\mu \bar{u}(p') \gamma^\mu \gamma_5 u(p) \\ &= \bar{u}(p') (\not{p}' - \not{p}) \gamma_5 u(p) \\ &= \bar{u}(p') (\not{p}' \gamma_5 + \gamma_5 \not{p}) u(p) \\ &= 2m \bar{u}(p') \gamma_5 u(p)\end{aligned}$$

Chiral symmetry

$$\gamma_5 = i\gamma^0\gamma^1\gamma^2\gamma^3$$

$$\{\gamma^\mu, \gamma_5\} = 0$$

Dirac equation in chiral representation for gamma matrices

$$\gamma^0 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \gamma^i = \begin{bmatrix} 0 & \sigma_i \\ -\sigma_i & 0 \end{bmatrix}, \quad \gamma_5 = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$$

splits into two equations

$$(i\partial_t - i\boldsymbol{\sigma} \cdot \boldsymbol{\nabla})\psi_L - m\psi_R = 0, \quad (i\partial_t + i\boldsymbol{\sigma} \cdot \boldsymbol{\nabla})\psi_R - m\psi_L = 0,$$

where $\psi = \begin{bmatrix} \psi_L \\ \psi_R \end{bmatrix}$. Note that for massless fermions these eqs. are *independent*.

Projection operators: $P_L = \frac{1}{2}(1 - \gamma_5)$, $P_R = \frac{1}{2}(1 + \gamma_5)$ project solutions of

given chirality (eigen value of γ_5)

$$\psi_- = \begin{bmatrix} \psi_L \\ 0 \end{bmatrix}, \quad \psi_+ = \begin{bmatrix} 0 \\ \psi_R \end{bmatrix}$$

Axial anomaly

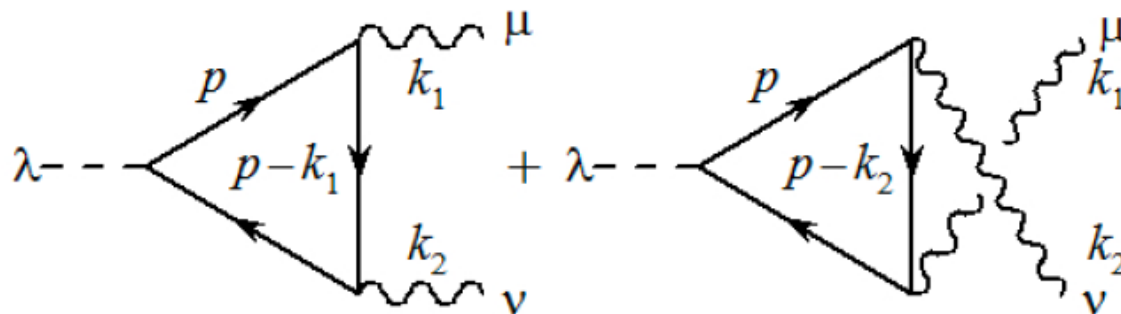
pseudoscalar
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Gauge invariance of QED (and QCD): $q_\mu j^\mu(q) = q_\mu \bar{u}(p') \gamma^\mu u(p) = 0$

divergence of axial-vector current: $q_\mu j_5^\mu(q) = q_\mu \bar{u}(p') \gamma^\mu \gamma_5 u(p) = 2m \bar{u}(p') \gamma_5 u(p)$

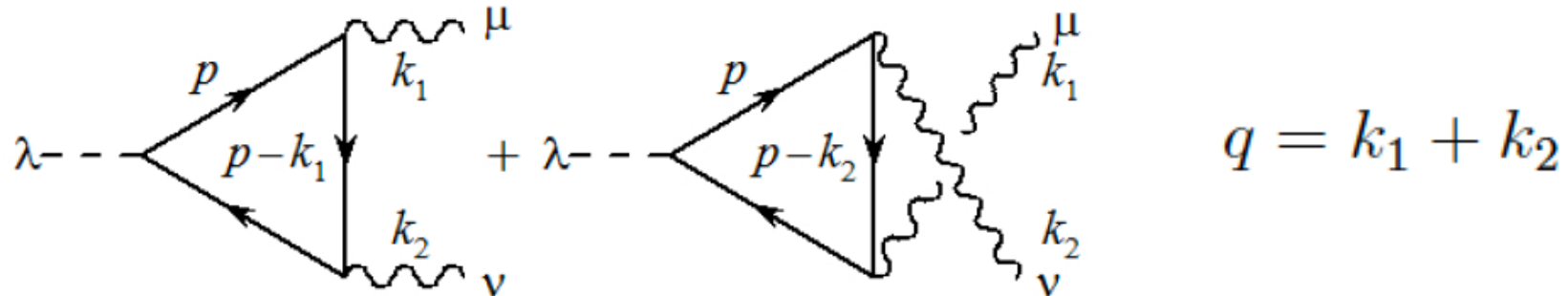
Axial current is conserved for massless fermions: chiral symmetry

It is not possible to maintain both symmetries when loop corrections are included. This is called: AXIAL ANOMALY



photons are bosons and they are not distinguishable hence amplitude has to be symmetrized

Naïve current conservation



Skipping coupling constants (charges) the amplitude reads:

$$T_{\mu\nu\lambda} = -i \int \frac{d^4 p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_\mu \right] \\ -i \int \frac{d^4 p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\mu \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu \right]$$

Naively we expect:

$$k_1^\mu T_{\mu\nu\lambda} = k_2^\nu T_{\mu\nu\lambda} = 0 \quad q^\lambda T_{\mu\nu\lambda} = 2m T_{\mu\nu}$$

Naïve current conservation

Vector current, first diagram:

$$k_1^\mu T_{\mu\nu\lambda} \sim \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \not{k}_1 \frac{i}{\not{p} - m} \right]$$

use trick:

$$\not{k}_1 = (\not{p} - m) - ((\not{p} - \not{k}_1) - m)$$

we get:

$$= i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \right] - i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{\not{p} - m} \right]$$

Naïve current conservation

Vector current, first diagram:

$$k_1^\mu T_{\mu\nu\lambda} \sim \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \not{k}_1 \frac{i}{\not{p} - m} \right]$$

use trick:

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$$= i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \right] - i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{\not{p} - m} \right]$$

same trick with the second diagram gives

$$= i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{\not{p} - m} \right] - i \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu \frac{i}{\not{p} - m} \right]$$

Naïve current conservation

$$k_1^\mu T_{\mu\nu\lambda} \sim \int \frac{d^4 p}{(2\pi)^4}$$

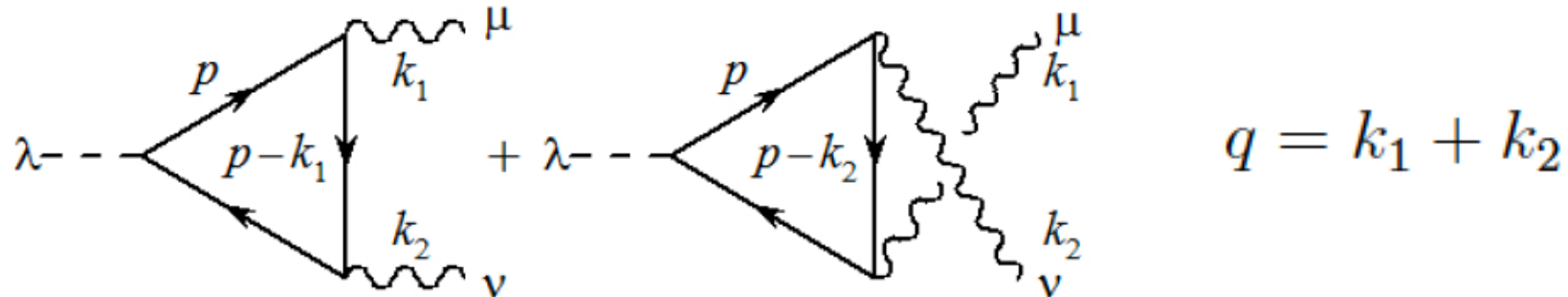
$$\left\{ \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \right] - \text{Tr} \left[\gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu \frac{i}{\not{p} - m} \right] \right\}$$

change variable in the first integral $p \rightarrow p + k_1$

$$q = k_1 + k_2$$

It seems we get zero

Naïve current conservation



Skipping coupling constants (charges) the amplitude reads:

$$T_{\mu\nu\lambda} = -i \int \frac{d^4p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_\mu \right] \\ - i \int \frac{d^4p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\mu \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu \right]$$

Naively we expect:

$$q^\lambda T_{\mu\nu\lambda} = 2m T_{\mu\nu}$$

Axial current

To calculate $q^\lambda T_{\mu\nu\lambda}$

we use the following trick:

$$\begin{aligned} \not{q}\gamma_5 &= -\gamma_5\not{q} \\ &= \gamma_5 [(\not{p} - \not{q}) - m] - \gamma_5 [\not{p} - m] \\ &= \gamma_5 [(\not{p} - \not{q}) - m] + [\not{p} - m] \gamma_5 + 2m\gamma_5 \end{aligned}$$

and for the first diagram we obtain

$$q^\lambda \left[\frac{i}{\not{p} - m} \gamma_\lambda \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \right] = 2m \frac{i}{\not{p} - m} \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} + i \frac{i}{\not{p} - m} \gamma_5 + i \gamma_5 \frac{i}{(\not{p} - \not{q}) - m}$$

Axial current

Sum from the two diagrams $q^\lambda T_{\mu\nu\lambda} = 2mT_{\mu\nu} + \Delta_{\mu\nu}^{(1)} + \Delta_{\mu\nu}^{(2)}$

$$\begin{aligned}
 & \Delta_{\mu\nu}^{(1)} + \Delta_{\mu\nu}^{(2)} \\
 = & \int \frac{d^4 p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_5 \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_\mu + \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_\mu \right] \\
 + & \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_5 \gamma_\mu \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu + \gamma_5 \frac{i}{(\not{p} - \not{q}) - m} \gamma_\mu \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu \right]
 \end{aligned}$$

Axial current

$$\Delta_{\mu\nu}^{(1)} = \int \frac{d^4p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_5 \gamma_\nu \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_\mu - \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_5 \gamma_\nu \frac{i}{(\not{p} - \not{q}) - m} \gamma_\mu \right]$$

$$\Delta_{\mu\nu}^{(2)} = \int \frac{d^4p}{(2\pi)^4} \text{Tr} \left[\frac{i}{\not{p} - m} \gamma_5 \gamma_\mu \frac{i}{(\not{p} - \not{k}_2) - m} \gamma_\nu - \frac{i}{(\not{p} - \not{k}_1) - m} \gamma_5 \gamma_\mu \frac{i}{(\not{p} - \not{q}) - m} \gamma_\nu \right]$$

The question is: are $\Delta_{\mu\nu}^{(1,2)}$ equal zero?

Changing variables

seems to nullify $\Delta_{\mu\nu}^{(1,2)}$.



$$p \rightarrow p + k_2$$

$$p \rightarrow p + k_1$$

However, $\Delta_{\mu\nu}^{(1,2)} \sim \int dp p^3 \frac{1}{p^2} \sim \int dp p$ are UV divergent

Due to the minus sign the divergence is only linear

The same applies to vector current conservation.

Mathematical diggression

Consider the integral that is naively zero:

$$\int_{-\infty}^{\infty} dx [f(x+a) - f(x)]$$

However, if

$$f(\pm\infty) \neq 0.$$

we can calculate this integral by Taylor expansion:

$$\int_{-\infty}^{\infty} dx [f(x+a) - f(x)] = a [f(\infty) - f(-\infty)] + \frac{a^2}{2} [f'(\infty) - f'(-\infty)] + \dots$$

it may happen that $\neq 0$