QCD lecture 6

November 8

Infrared divergences

$$S_F^R = \frac{i}{\not p}\left(1 + \frac{\alpha(\mu^2)}{4\pi} C_F \left(\ln\left(\frac{-p^2}{\bar\mu^2}\right) - 1\right)\right)$$

Divergent for $p^2 = 0$. This is infrared divergence (from the lower int. limit).
It can be regularized by going to the number of dimensions higher than 4.
Before expansion, change $\varepsilon \to -\kappa$

$$S_F^R(p) = \frac{i}{\not p}\left(1 - \frac{\alpha_s}{4\pi} C_F \left(\frac{\bar\mu^2}{-p^2}\right)^{\varepsilon}\left(\frac{1}{\varepsilon} + 1\right) + \frac{\alpha_s}{4\pi} C_F \frac{1}{\varepsilon}\right)$$

Infrared divergences

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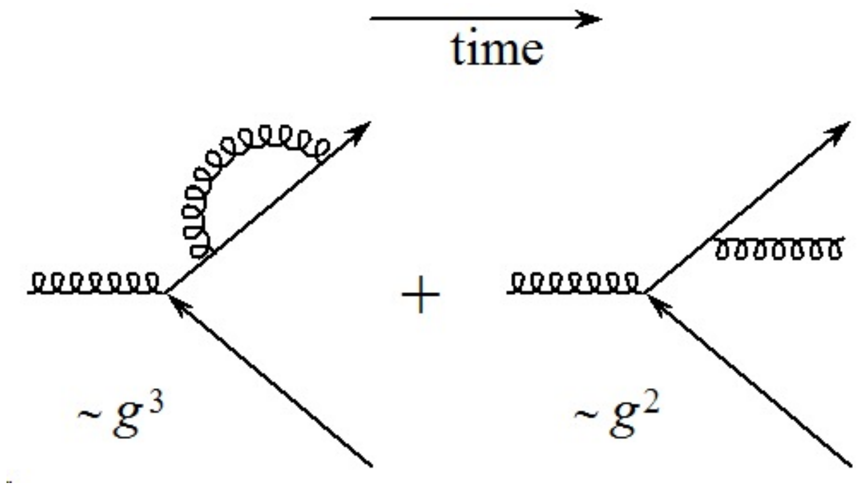
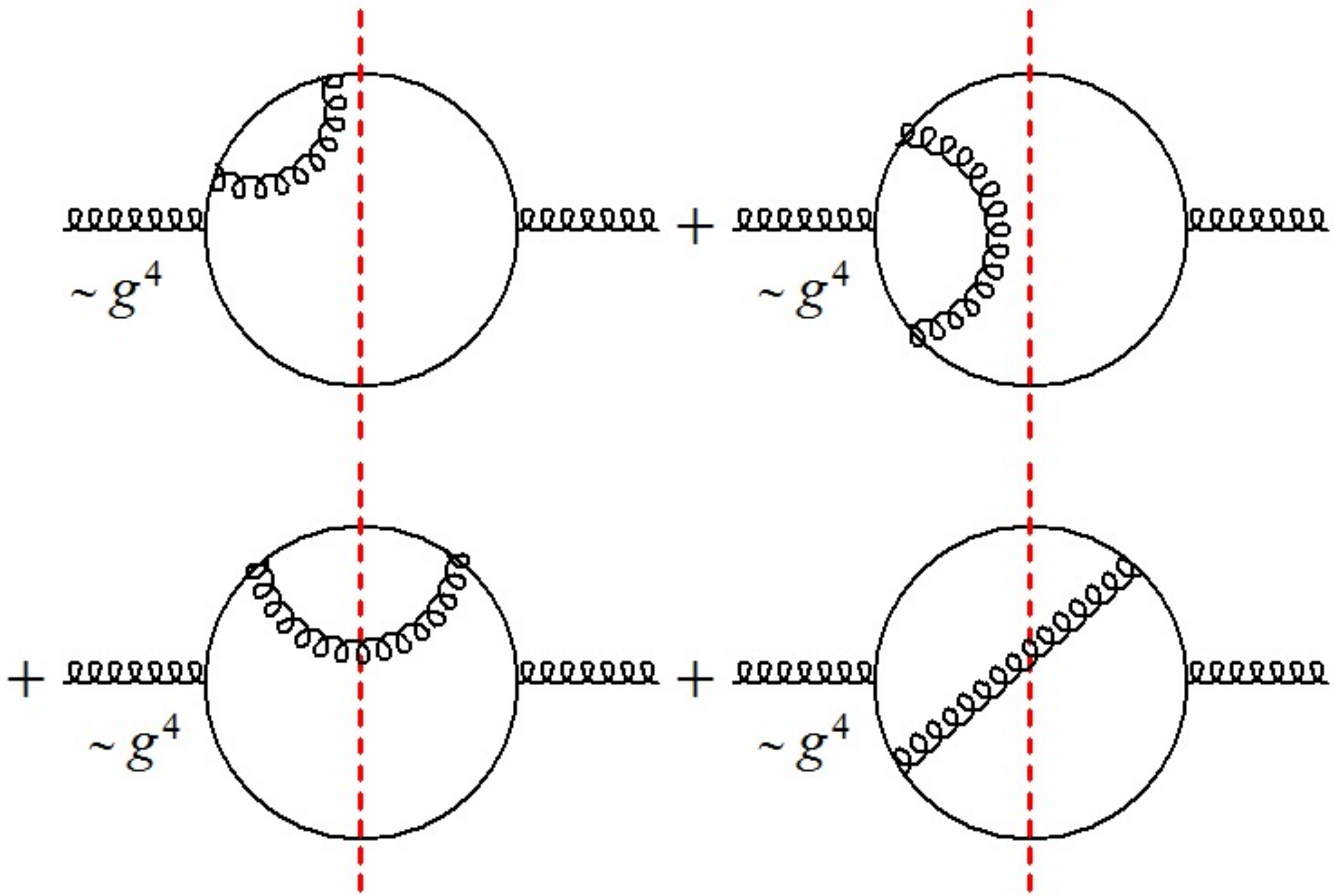
$$\begin{aligned} S_F^R(p) &= \frac{i}{\not p}\left(1 - \frac{\alpha_s}{4\pi} C_F \left(\frac{-p^2}{\bar{\mu}^2}\right)^{\kappa}\left(-\frac{1}{\kappa} + 1\right) - \frac{\alpha_s}{4\pi} C_F \frac{1}{\kappa}\right) \\ &\underset{p^2=0}{=} \frac{i}{\not p}\left(1 - \frac{\alpha_s}{4\pi} C_F \frac{1}{\kappa}\right). \end{aligned}$$

Infrared divergencies



One cannot distinguish
a single electron from
an electron accompanied
by a zero energy foton
or a collinear foton (for
massless fermion).
One has to sum over such
degenerate states.

Infrared divergencies



Here IR singularities cancel out

Infrared singularities

IR singulariteis arise when the theory has massless particles (photon, gluon)

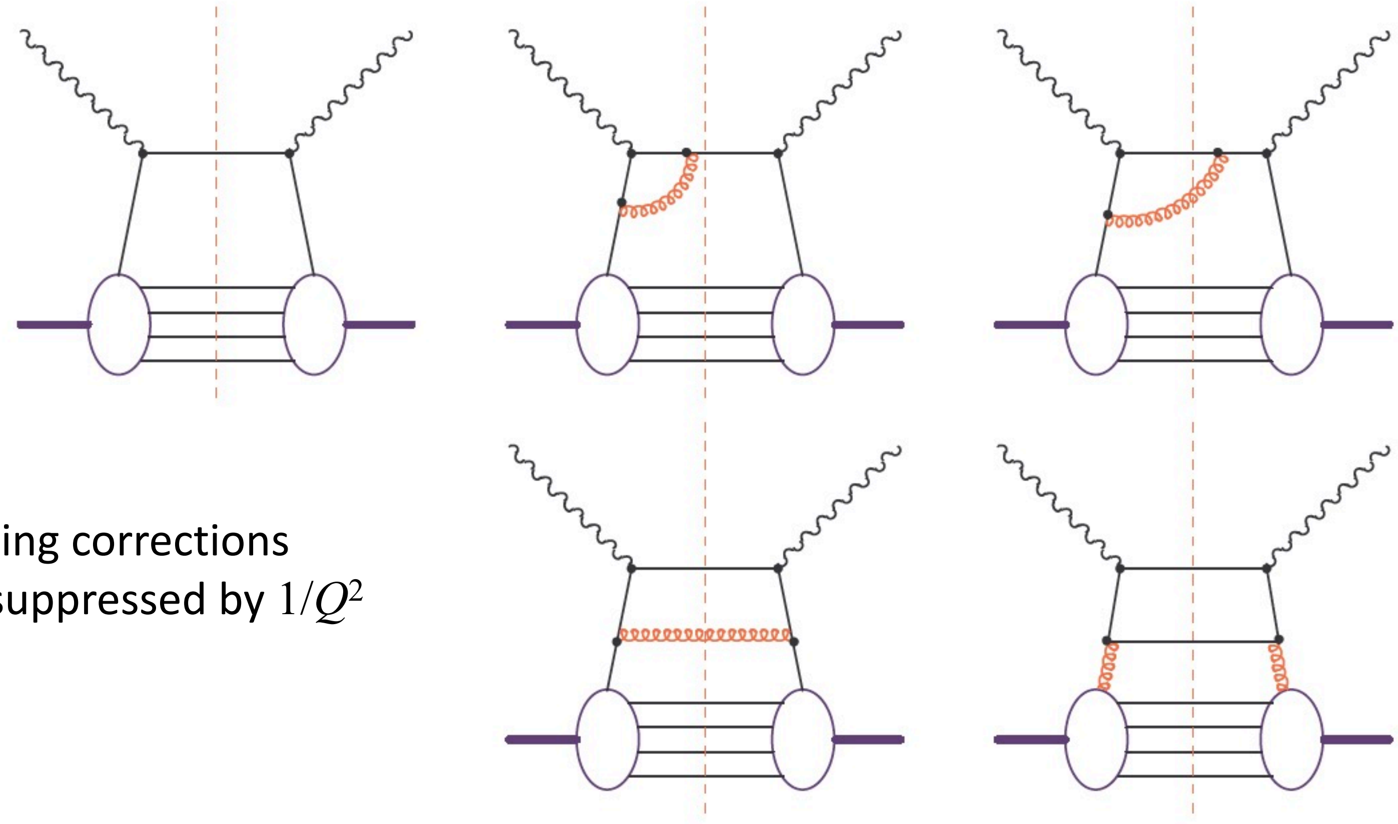
- when energy of photon (gluon) is small – soft singularity
- when for massless fermion photon (gluon) is parallel to that fermion – collinear singularity

Bloch – Nordsieck theorem (baically derived for QED)
Kinoshita – Lee – Nauenberg theorem (generalized to QCD)

Kinoshita-Lee-Nauenberg (KLN) theorem assures that a summation over degenerate initial and final states removes all infrared (IR) divergences in QCD.

This very broad topic, beyond the scope of this lecture

QCD corrections to parton model

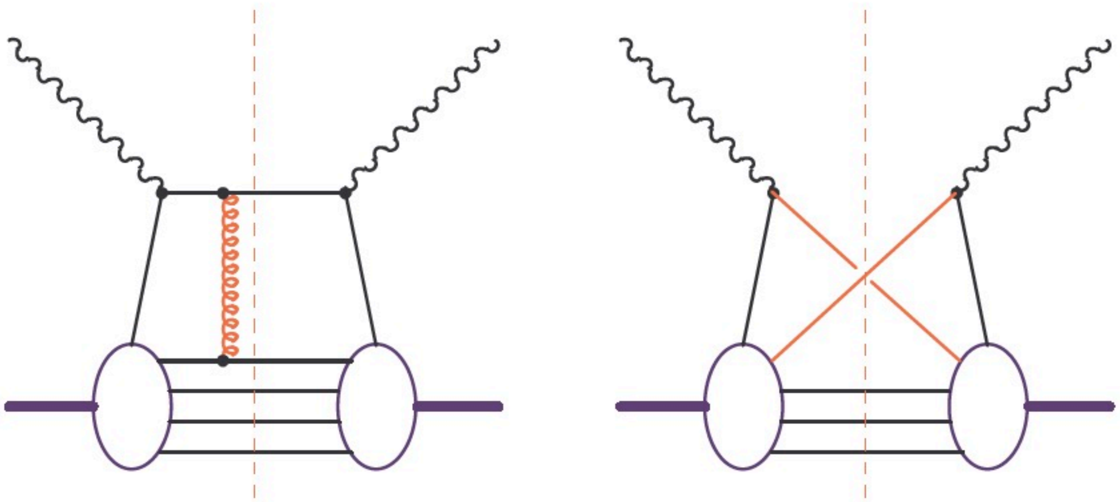


Leading corrections
not suppressed by $1/Q^2$

photon scatters off the gluon

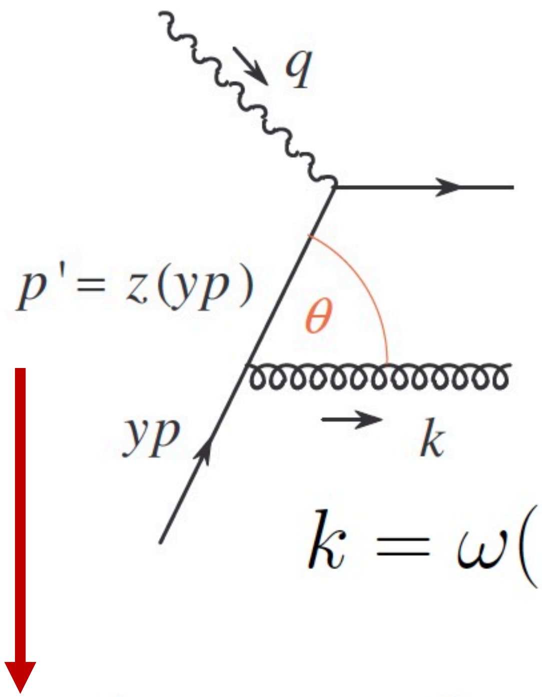
QCD corrections to parton model

Non-leading corrections suppressed by $1/Q^2$



QCD corrections to parton model

$$yp = E(1,0,0,1)$$

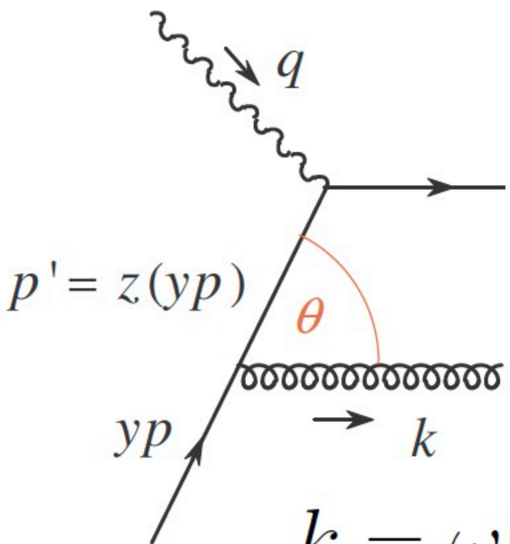


$$d^4k\,\delta(k^2) \sim \frac{d^3\boldsymbol{k}}{\omega} \sim \omega d\omega d\cos\theta d\varphi$$

$$k = \omega(1,\ \sin\theta\sin\varphi,\ \sin\theta\cos\varphi,\ \cos\theta)$$

$$\frac{1}{p'^{\,2}} = \frac{1}{(yp-k)^2} = \frac{1}{2yp\,k} = \frac{1}{2E\omega(1-\cos\theta)}$$

QCD corrections to parton model



$$yp = E(1,0,0,1)$$

$$d^4k\,\delta(k^2) \sim \frac{d^3\boldsymbol{k}}{\omega} \sim \omega d\omega d\cos\theta d\varphi$$

$$k = \omega(1,\ \sin\theta\sin\varphi,\ \sin\theta\cos\varphi,\ \cos\theta)$$

$$|\mathcal{M}|^2\, d^4k\,\delta(k^2) \sim \sin^2\theta\,\frac{\omega d\omega\, d\cos\theta}{\omega^2(1-\cos\theta)^2} \sim \frac{d\omega}{\omega}\frac{d\theta^2}{\theta^2}$$

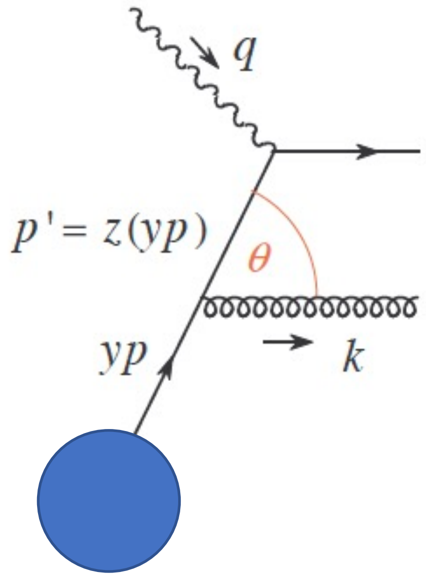
QCD corrections to parton model

$$\frac{d\omega}{\omega}\frac{d\theta^2}{\theta^2}$$

- soft (cancel) $\omega \longrightarrow 0$
- collinear (remain) $\theta \longrightarrow 0$

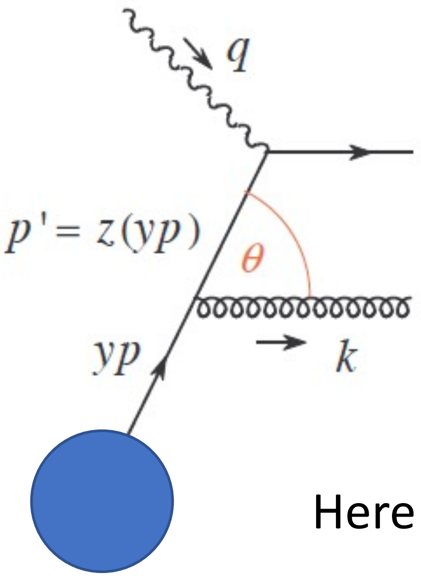
In dimensional regularization:

$$\left(\frac{Q^2}{\mu^2}\right)^{\kappa}\frac{1}{\kappa} = \frac{1}{\kappa} + \log\left(\frac{Q^2}{\mu^2}\right)$$



Poles can be absorbed into bare parton densities.
Logs can be summed up to all orders. Factrozation.
Coefficients of the poles are universal functions of z

Altarelli-Parisi probabilities



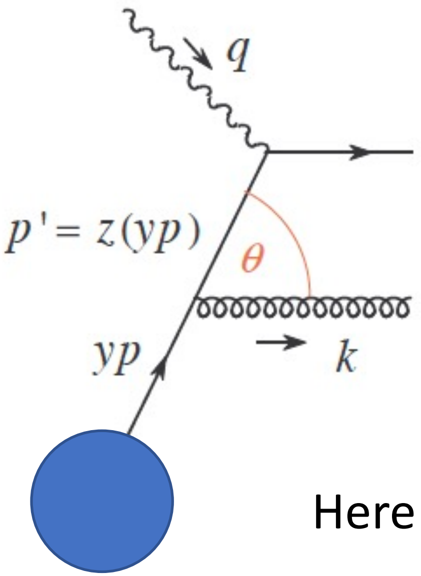
$$\log\left(\frac{Q^2}{\mu^2}\right) P_{qq}(z)$$

It turns out that potentially large logs are multiplied by universal functions of the momentum fraction *z* (with respect to the emitting parton)

Here $P_{qq}(z) = P_{q\leftarrow q}(z)$ is a probability of “finding”

a quark of the longitudinal momentum fraction z in initial quark

Altarelli-Parisi probabilities



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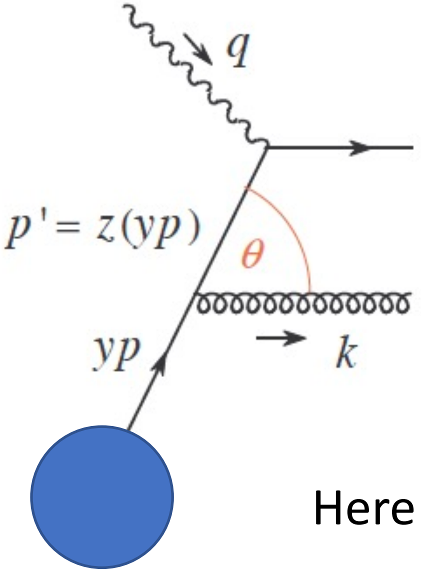
$$P_{qq}(z) = C_F\left(\frac{1+z^2}{1-z}\right)_+$$

Altarelli-Parisi probabilities



sample diagram

$$\log\left(\frac{Q^2}{\mu^2}\right) P_{qq}(z)$$

It turns out that potentially large logs are multiplied by universal functions of the momentum fraction *z* (with respect to the emitting parton)

Here $P_{qq}(z) = P_{q\leftarrow q}(z)$ is a probability of "finding"

a quark of the longitudinal momentum fraction z in initial quark

$$P_{qq}(z) = C_F\left(\frac{1+z^2}{1-z}\right)_+$$

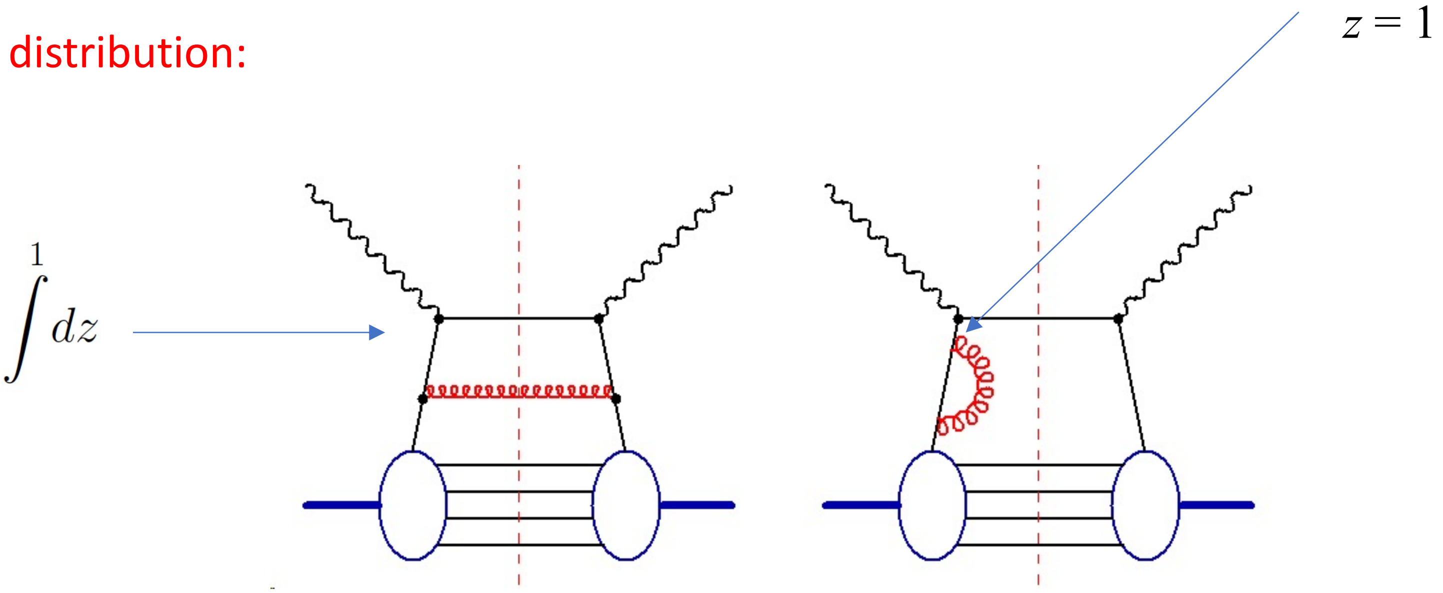
"Plus" distribution:

$$\int^1 dz\,(\ldots)_+\,g(z) = \int^1 dz\,(\ldots)\,[g(z)-g(1)]$$

appears because of the virtual diagram for which $z = 1$

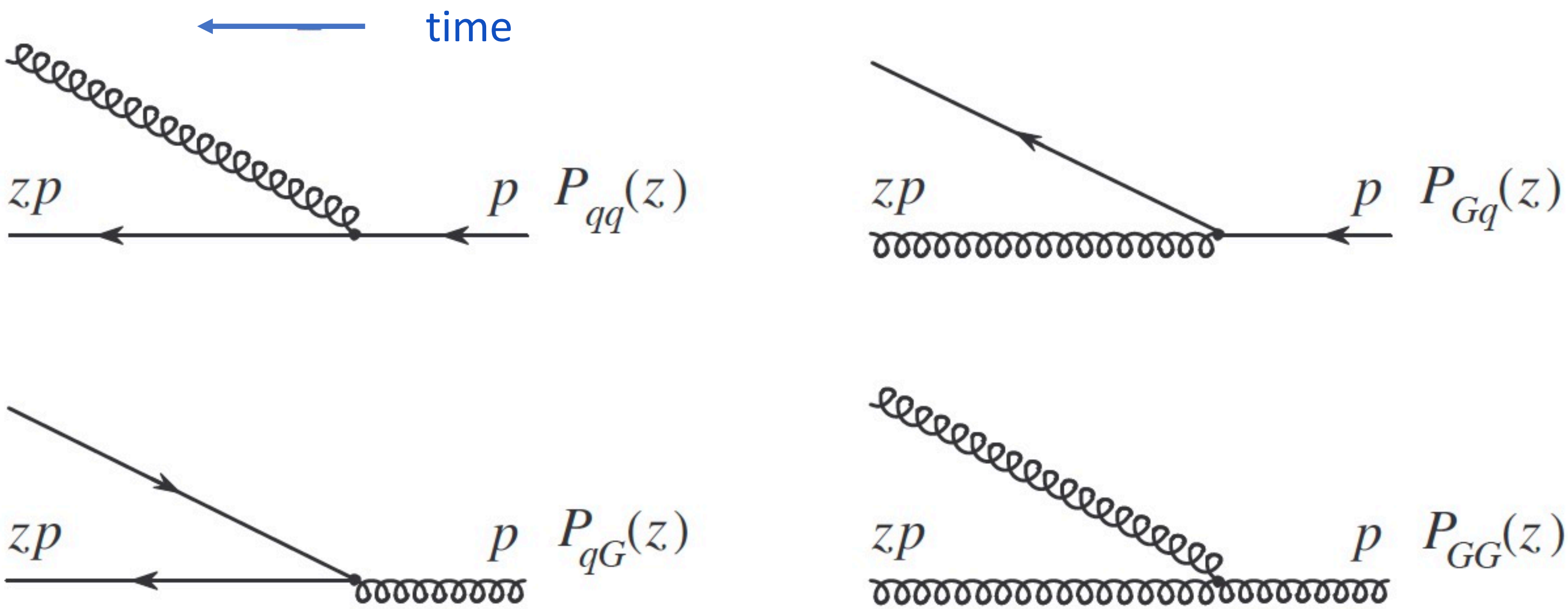
Altarelli-Parisi probabilities

“Plus” distribution:



Different diagrams give extra contribution at $z = 1$ in different gauges.
The result is the same: no singularity at $z = 1$.

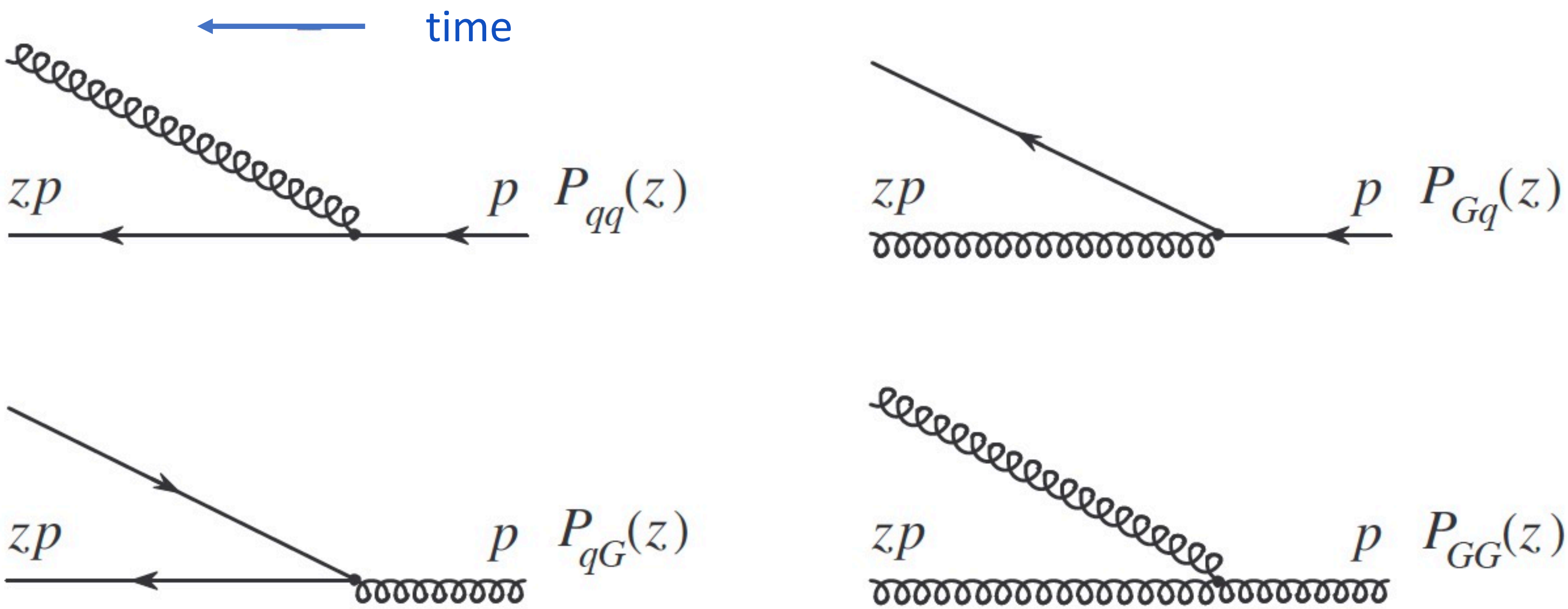
Altarelli-Parisi probabilities



$$P_{qq}(z) = C_F \left(\frac{1+z^2}{1-z}\right)_+ , \quad P_{Gq}(z) = C_F \frac{1+(1-z)^2}{z}, \quad P_{qG}(z) = \frac{1}{2}\left[z^2 + (1-z)^2\right]$$

$$P_{GG}(z) = 2C_A \left[\frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z)\right] + \frac{1}{2}\left(\frac{11}{3}C_A - \frac{2}{3}n_f\right)\delta(1-z)$$

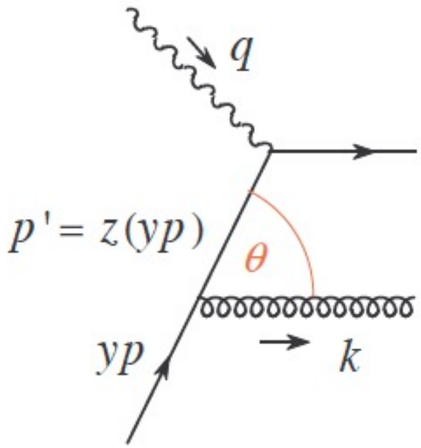
Altarelli-Parisi probabilities



$$P_{qG}(z) = P_{\bar{q}G}(z), \qquad P_{Gq}(z) = P_{G\bar{q}}(z),$$

$$P_{qq}(z) = P_{Gq}(1-z), \; P_{GG}(z) = P_{GG}(1-z), \quad P_{qG}(z) = P_{qG}(1-z)$$

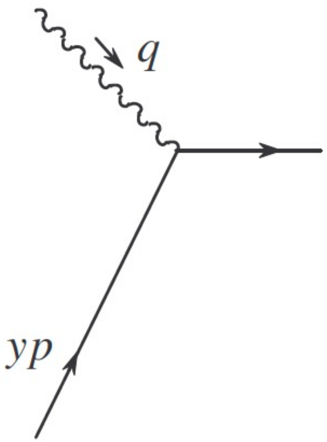
QCD corrections to parton model



on-shell condition

$$0 = (zyp + q)^2 = 2zy\,pq + q^2 = 2M\nu\,zy - Q^2$$

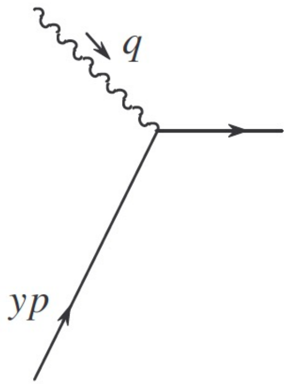
$$zy = \frac{Q^2}{2M\nu} = x$$



Recall F_1: $$F_1(x) = \frac{1}{2}\sum_i e_i^2\, f_i(x)$$

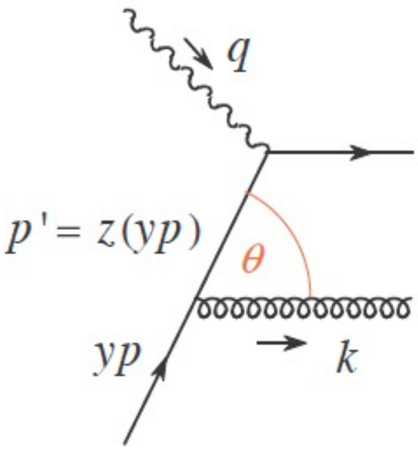
$$2F_1(x) = e_q^2 \int_0^1 dyq(y)\delta(y - x)$$

QCD corrections to parton model



Recall F_1:

$$2F_1(x) = e_q^2 \int_0^1 dy q(y) \delta(y - x)$$

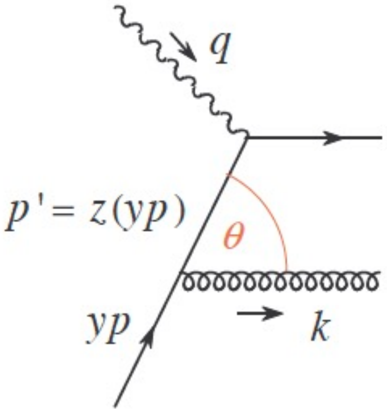


$$2\Delta F_1(x) = e_q^2 \frac{\alpha_s}{2\pi} \int_0^1 dy q(y) \int_0^1 dz \delta(zy - x) \left[P_{qq}(z) \ln \frac{Q^2}{\mu^2} + C(z) \right]$$

Correction to F_1 large logs

$$q(x,Q^2) = q(x,\mu^2) + \frac{\alpha_s}{2\pi}\ln\frac{Q^2}{\mu^2}\int_x^1 \frac{dy}{y} P_{qq}\left(\frac{x}{y}\right) q(y,\mu^2) + \dots$$

$$= q(x,\mu^2) + \frac{\alpha_s}{2\pi}\ln\frac{Q^2}{\mu^2} \underline{P_{qq}\otimes q(\mu^2)} +$$



Convolution:

$$P_{qq}\otimes q = \int_0^1 dz \int_0^1 dy\,\delta(zy - x)\underline{P_{qq}(z)}q(y)$$

Integration over $d\theta$ gave a pole

DGLAP Evolution Equation

$$\frac{d}{d\ln Q^2} = Q^2 \frac{d}{dQ^2} \quad \Rightarrow \quad q(x,Q^2) = q(x,\mu^2) + \frac{\alpha_s}{2\pi}\ln\frac{Q^2}{\mu^2} P_{qq} \otimes q(\mu^2) + \ldots$$

Evolution eq.
Dokshitzer,
Gribov, Lipatov
Altarelli, Parisi

$$\frac{d}{d\ln Q^2} q(x,Q^2) = \frac{\alpha_s}{2\pi} P_{qq} \otimes q(Q^2)$$

Such equation sums up all powers $\frac{\alpha_s}{2\pi}\ln\frac{Q^2}{\mu^2}$.

Leading Log Approximation (LLA)

DGLAP Evolution Equations

Full set of DGLAP equations:

$$Q^2 \frac{d}{dQ^2} q_i(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} \left[P_{qq} \otimes q_i(Q^2) + P_{qG} \otimes G(Q^2) \right]$$

$$Q^2 \frac{d}{dQ^2} G(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} \left[P_{Gq} \otimes \sum_i q_i(Q^2) + P_{GG} \otimes G(Q^2) \right]$$

We need an input at one scale Q_0^2 and then we can evolve them up to some other Q^2
note that index i runs over quarks and antiquarks
when we construct a difference, called non-singlet, gluons cancel

$$q_i^{NS}(x, Q^2) = q_i(x, Q^2) - \overline{q}_i(x, Q^2)$$

DGLAP Evolution Equations

Define:

singlet
$$q^S(x,Q^2) = \sum_i \left(q_i(x,Q^2) + \overline{q}_i(x,Q^2)\right)$$

nonsinglet
$$q_i^{NS}(x,Q^2) = q_i(x,Q^2) - \overline{q}_i(x,Q^2)$$

DGLAP Evolution Equations

$$Q^2 \frac{d}{dQ^2} q^{NS}(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} P_{qq} \otimes q^{NS}(Q^2)$$

$$Q^2 \frac{d}{dQ^2} q^S(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} \left[P_{qq} \otimes q^S(Q^2) + 2n_f P_{qG} \otimes G(Q^2) \right]$$

$$Q^2 \frac{d}{dQ^2} G(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} \left[P_{Gq} \otimes q^S(Q^2) + P_{GG} \otimes G(Q^2) \right]$$

DGLAP for Mellin moments

Moments of the convolution

$$M_n = \int_0^1 dx\, x^{n-1} P \otimes f = \int_0^1 dx\, x^{n-1} \int_0^1 dz \int_0^1 dy \delta(zy - x) P(z) f(y)$$

$$= \int_0^1 dz\, z^{n-1} P(z) \int_0^1 dy\, y^{n-1} f(y) = P_n\, f_n = \gamma^n f_n$$

γ^n anomalous dimension

convolution is replaced by a product

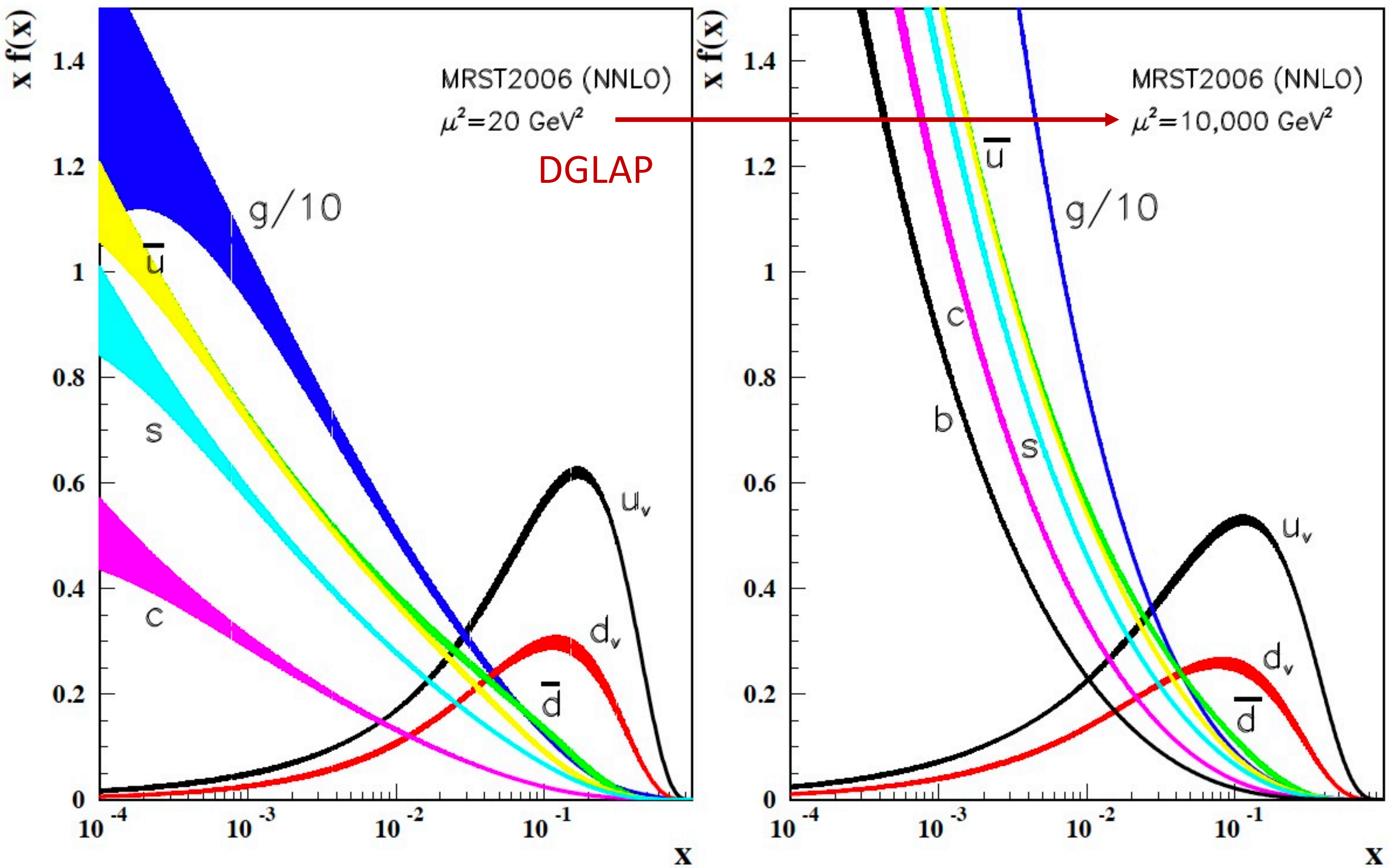
DGLAP for Mellin moments

$$\frac{dq_n^{NS}(t)}{dt} = \frac{\alpha_s(t)}{2\pi}\gamma_{qq}^n q_n^{NS}(t)$$

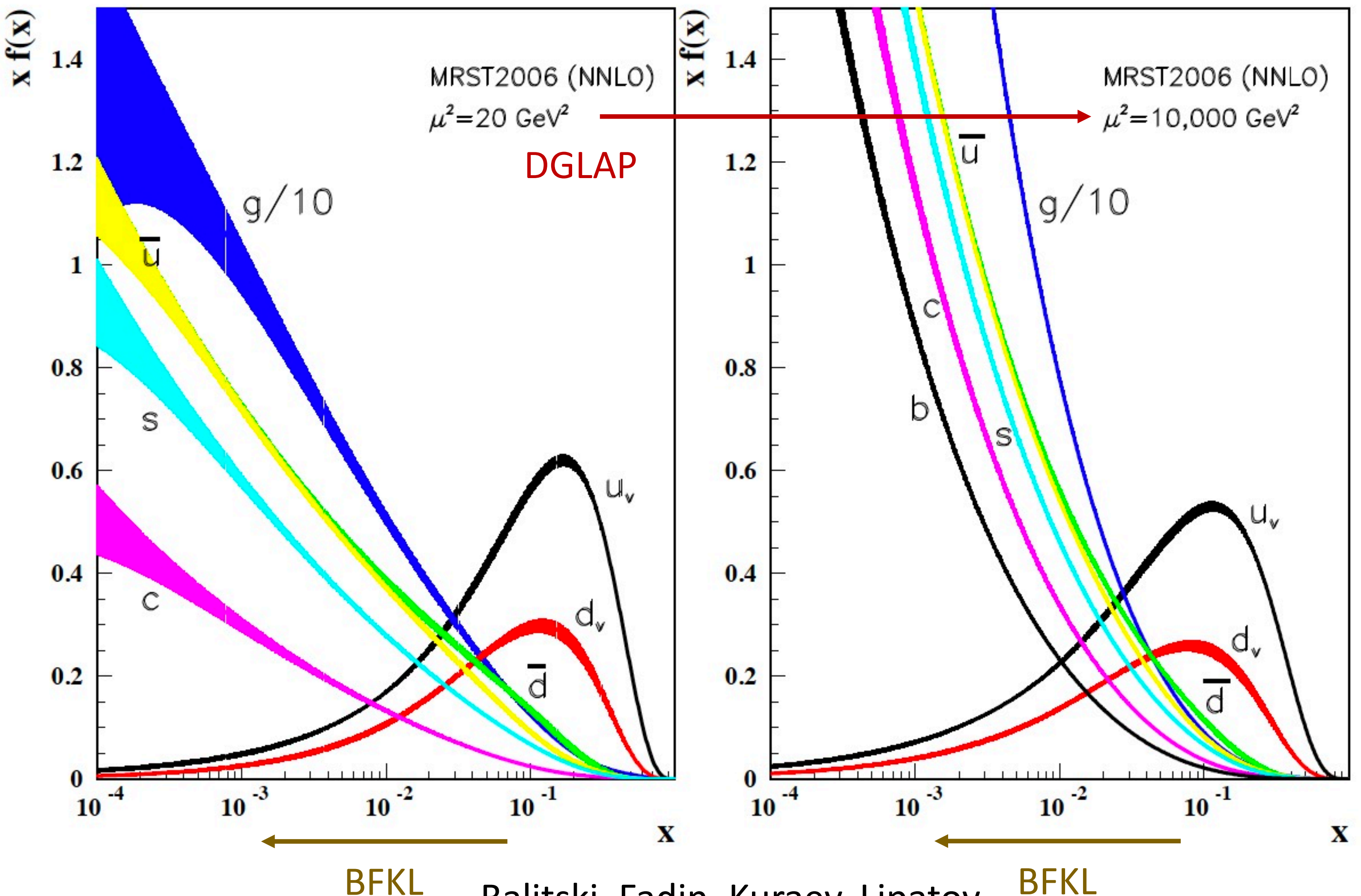
$$\frac{d}{dt}\begin{bmatrix} q_n^S(t) \\ G_n(t) \end{bmatrix} = \frac{\alpha_s(t)}{2\pi}\begin{bmatrix} \gamma_{qq}^n & 2n_f\gamma_{qG}^n \\ \gamma_{Gq}^n & \gamma_{GG}^n \end{bmatrix}\begin{bmatrix} q_n^S(t) \\ G_n(t) \end{bmatrix}$$

$$\frac{\alpha_s(t)}{2\pi} = 2\,a_s(t) = 2\,\frac{1}{\beta_0 t}$$

Numerical solutions



Numerical solutions



Balitski, Fadin, Kuraev, Lipatov

HERA F_2: data vs. theory

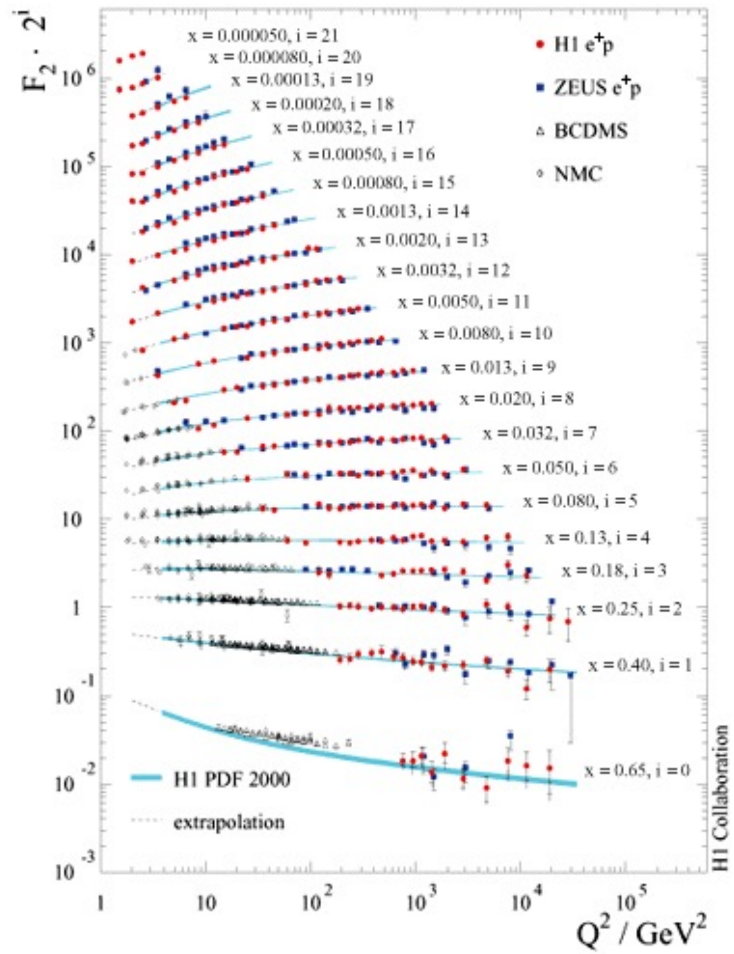


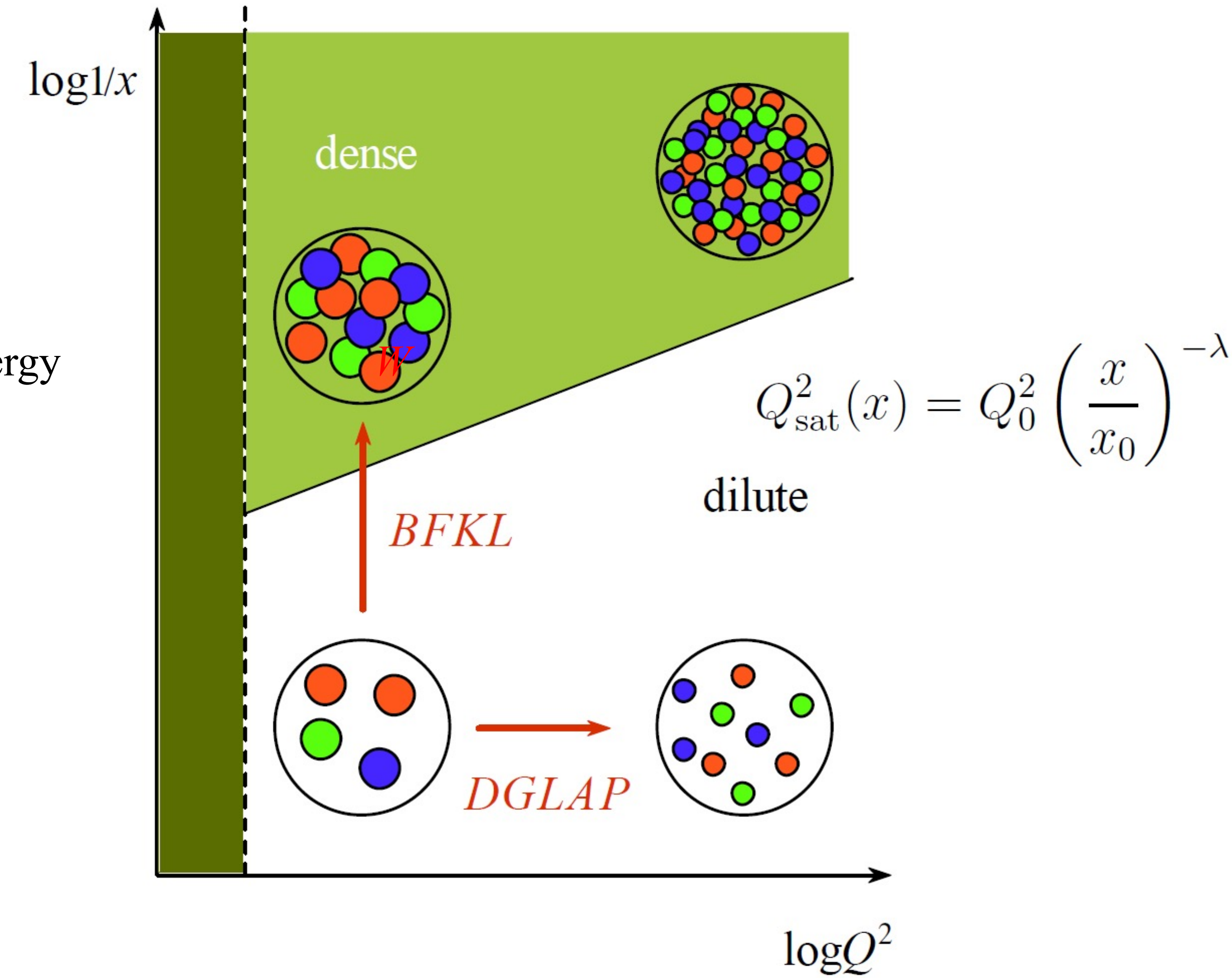
FIG. 2: Structure function F_2 as a function of Q^2 based on HERA-I measurements of H1 [2, 3] and ZEUS [4] collaboration compared to results from fixed target experiments BCDMS [5] and NMC [6].

DGLAP vs. BFKL

small x
large W

W – gamma+proton energy

large x
small W



Anomalous dimensions

$$\gamma_{qq}^{n} = C_F \left[-2\sum_{k=1}^{n+1} \frac{1}{k} + \frac{3}{2} + \frac{1}{n} + \frac{1}{n+1} \right],$$

$$\gamma_{qG}^{n} = \frac{1}{2} \frac{2+n+n^2}{n(n+1)(n+2)},$$

$$\gamma_{Gq}^{n} = C_F \frac{2+n+n^2}{n(n^2-1)}$$

$$\gamma_{GG}^{n} = 2C_A \left[\frac{11}{12} - \sum_{k=1}^{n+2} \frac{1}{k} + \frac{1}{n-1} - \frac{1}{n} + \frac{2}{n+1} \right] - \frac{n_f}{3}$$

Valnce quark # conservation

$$\gamma_{qq}^{n} = C_F \left[-2 \sum_{k=1}^{n+1} \frac{1}{k} + \frac{3}{2} + \frac{1}{n} + \frac{1}{n+1} \right]$$

$$\gamma_{qq}^{1} = 0 \quad \rightarrow \quad \frac{dq_n^{NS}(t)}{dt} = 0$$

$$\int dx \left[q_i(x, Q^2) - \overline{q}_i(x, Q^2) \right] = \text{const.} = \int dx q_{Vi}(x, Q^2)$$

Momentum conservation

consider moment $n = 2$ for the singlet eqs.

$$\frac{d}{dt}q_2^S(t) = -\frac{2}{\beta_0 t}\left[\frac{4C_F}{3}q_2^S(t) - \frac{n_f}{3}G_2(t)\right] = -\frac{2}{\beta_0 t}f(t)$$

$$\frac{d}{dt}G_2(t) = +\frac{2}{\beta_0 t}\left[\frac{4C_F}{3}q_2^S(t) - \frac{n_f}{3}G_2(t)\right] = +\frac{2}{\beta_0 t}f(t)$$

$$q_2^S(t) + G_2(t) = \text{const.}$$

$$= \int dx\, x\left[\sum_i \left(q_i(x,Q^2) + \overline{q}_i(x,Q^2)\right) + G(x,Q^2)\right] = 1$$

value of 1 is a requirement for a proper normalization

Gluon momentum

$$\frac{d}{dt}q_2^S(t) = -\frac{2}{\beta_0 t}\left[\frac{4C_F}{3}q_2^S(t) - \frac{n_f}{3}G_2(t)\right] = -\frac{2}{\beta_0 t}f(t)$$

$$\frac{d}{dt}G_2(t) = +\frac{2}{\beta_0 t}\left[\frac{4C_F}{3}q_2^S(t) - \frac{n_f}{3}G_2(t)\right] = +\frac{2}{\beta_0 t}f(t)$$

Form a linear combination

$$\frac{4C_F}{3}\frac{d}{dt}q_2^S(t) - \frac{n_f}{3}\frac{d}{dt}G_2(t) = \frac{d}{dt}f(t) = -\frac{2}{\beta_0 t}\left[\frac{4C_F}{3} + \frac{n_f}{3}\right]f(t)$$

since $c = \frac{4C_F}{3} + \frac{n_f}{3} > 0$

the solution is trivial and tends to 0 $f(t) = f(t_0)\left(\frac{t}{t_0}\right)^{-2c/\beta_0} \underset{t\to\infty}{\longrightarrow} 0$

Gluon momentum

We have two asymptotic constraints:

$$f(t) = \frac{4C_F}{3} q_2^S(t) - \frac{n_f}{3} G_2(t) = 0 \qquad\qquad q_2^S(t) + G_2(t) = 1$$

which give

$$q_2^S(t) = \frac{n_f}{4C_F} G_2(t) \qquad \rightarrow \qquad \left[\frac{n_f}{4C_F} + 1\right] G_2(t) = 1$$

numerically we have

$$G_2(t) = \frac{1}{\frac{n_f}{4C_F} + 1} = \frac{16}{16 + 3n_f} = \underset{n_f=3}{0.64}, \underset{n_f=4}{0.57}, \underset{n_f=5}{0.52}, \underset{n_f=6}{0.47}$$

asymptotically gluons carry around 50% of total momentum!