$\begin{array}{c} {\rm QCD} \\ {\rm problem \ set \ 11} \end{array}$

1. Flavor SU(3) left and right currents read

$$L^a_\mu = \bar{q}_L \gamma_\mu \frac{\lambda_a}{2} q_L, \ R^a_\mu = \bar{q}_R \gamma_\mu \frac{\lambda_a}{2} q_R$$

where

$$q_L_R = \frac{1}{2} \left(1 \mp \gamma_5 \right) q.$$

Show that vector and axial currents take the following form

$$V^a_\mu = R^a_\mu + L^a_\mu = \bar{q}\gamma_\mu \frac{\lambda_a}{2}q,$$

$$A^a_\mu = R^a_\mu - L^a_\mu = \bar{q}\gamma_\mu \gamma_5 \frac{\lambda_a}{2}q.$$

2. In the case of fermion fields, commutation relations of scalar fields are replaced by anticommutation relations:

$$\left\{q_{\alpha,k}(t,\vec{x}),q_{\beta,l}^{\dagger}(t,\vec{x}')\right\} = \delta^{(3)}(\vec{x}-\vec{x}')\delta_{\alpha\beta}\delta_{kl}$$

where α, β stand for Dirac indices and k, l denote SU(3) indices. Relevant charges are defined as

$$\hat{Q}^{a}_{L,R}(t) = \int d^{3}\vec{x} \, q^{\dagger}_{L,R}(t,\vec{x}) T^{a} q_{L,R}(t,\vec{x}), \hat{Q}_{V}(t) = \int d^{3}\vec{x} \, \left[q^{\dagger}_{L}(t,\vec{x}) q_{L}(t,\vec{x}) + q^{\dagger}_{R}(t,\vec{x}) q_{R}(t,\vec{x}) \right]$$

where $T^a = \lambda^a/a$ are SU(3) generators (Gell-Mann matrices). Making use of the identity (prove it!)

$$\{ab, cd\} = a \{b, c\} d - ac\{b, d\} + \{a, c\} bd - c\{a, d\}b$$

show that

$$\begin{bmatrix} \hat{Q}_L^a, \hat{Q}_L^b \end{bmatrix} = i f^{abc} \hat{Q}_L^c,$$

$$\begin{bmatrix} \hat{Q}_R^a, \hat{Q}_R^b \end{bmatrix} = i f^{abc} \hat{Q}_R^c,$$

$$\begin{bmatrix} \hat{Q}_L^a, \hat{Q}_R^b \end{bmatrix} = 0,$$

$$\begin{bmatrix} \hat{Q}_{L,R}^a, \hat{Q}_V \end{bmatrix} = 0.$$

3. Effective Lagrangian describing Goldstone boson interactions reads

$$\mathcal{L}_{\text{eff}} = \frac{F^2}{4} \operatorname{Tr} \left(\partial_{\mu} U \partial^{\mu} U^{\dagger} \right)$$

where $U = \exp(i\phi/F)$ can be expressed in terms of the physical meson fields:

$$\phi(x) = \sum_{a} \lambda_{a} \phi^{a}(x) = \begin{pmatrix} \pi^{0} + \eta/\sqrt{3} & \sqrt{2}\pi^{+} & \sqrt{2}K^{+} \\ \sqrt{2}\pi^{-} & -\pi^{0} + \eta/\sqrt{3} & \sqrt{2}K^{0} \\ \sqrt{2}K^{-} & \sqrt{2}\bar{K}^{0} & -2\eta/\sqrt{3} \end{pmatrix}$$

Expand \mathcal{L}_{eff} up to 4-field interactions and calculate flavor trace for the case of SU(2) (see lecture 11) and SU(3).

4. Decay width Γ (inverse of the particle life-time) of a particle of mass M (at rest) into two particles of masses $m_{1,2}$ reads

$$\Gamma = \frac{1}{2M} \int |\mathcal{M}|^2 (2\pi)^4 \,\delta^{(4)}(p - k_1 - k_2) \,\frac{d^4k_1}{(2\pi)^4} (2\pi) \,\delta(k_1^2 - m_1^2) \frac{d^4k_2}{(2\pi)^4} (2\pi) \,\delta(k_2^2 - m_2^2)$$

where \mathcal{M} is the relevant matrix element. Show that one can perform all integrations in the above formula. As a first step perform integrations over $dE_{1,2}$. Then using $\delta^{(4)}(p-k_1-k_2) = \delta(M-E_1-M_2)\delta^{(3)}(\mathbf{k}_1+\mathbf{k}_2)$ perform integration over $d^3\mathbf{k}_2$. Using

$$d^3 \boldsymbol{k}_1 = k_1^2 dk_1 d\Omega$$

and assuming that matrix element \mathcal{M} does not depend on angles perform integration over $d\Omega$. Finally, using the energy δ function perform the integration over dk_1 with the help of the following change of variables:

$$u = \sqrt{k_1^2 + m_1^2} + \sqrt{k_1^2 + m_2^2}.$$

Solve the energy δ function assuming that particle 2 is massless.