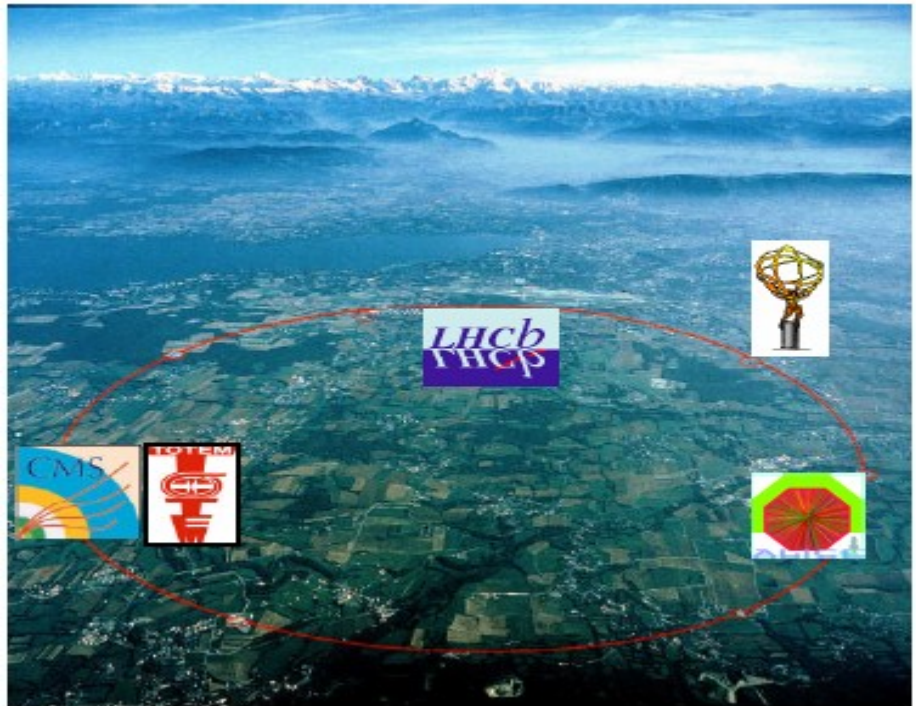


Physics with first fb^{-1} at Large Hadron Collider

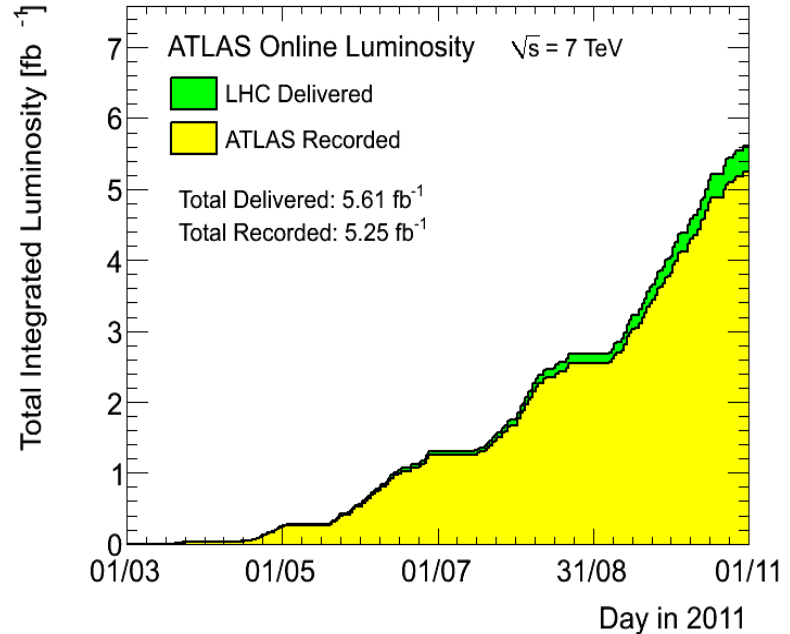
Today:

- **W,Z inclusive cross-section**
- **Asymmetry**
- **W+jets, Z + jets**



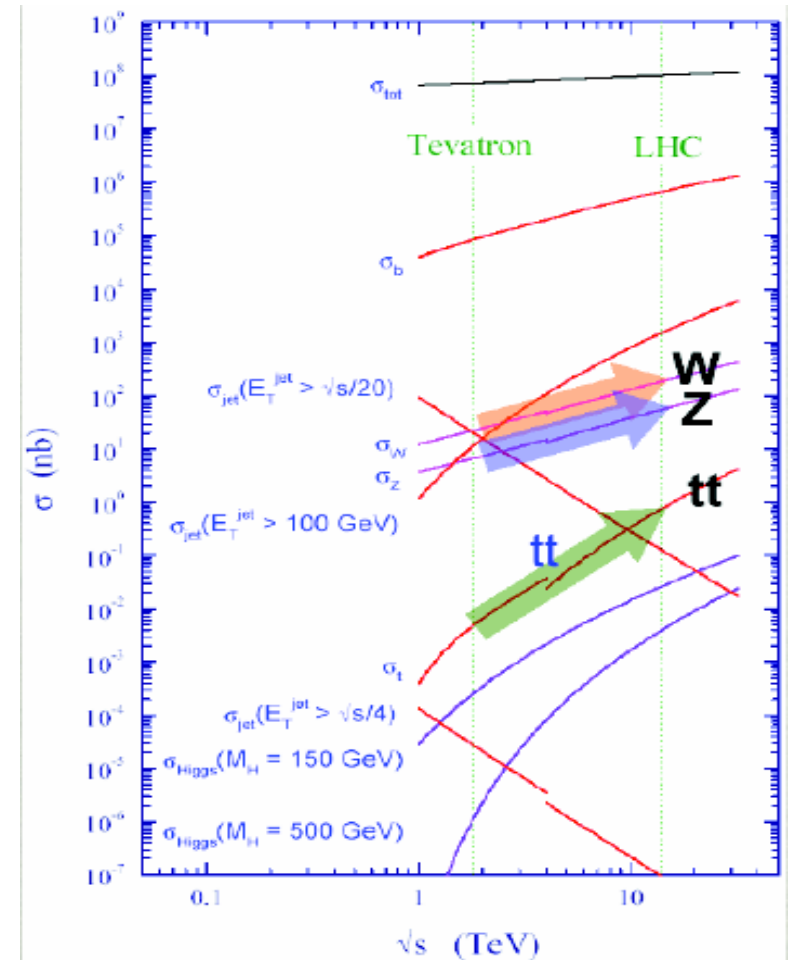
News from last week(s)

- **5.25 fb⁻¹ pp collisions recorded by ATLAS**
- **PbPb run of 2011 started**



Bosons at LHC

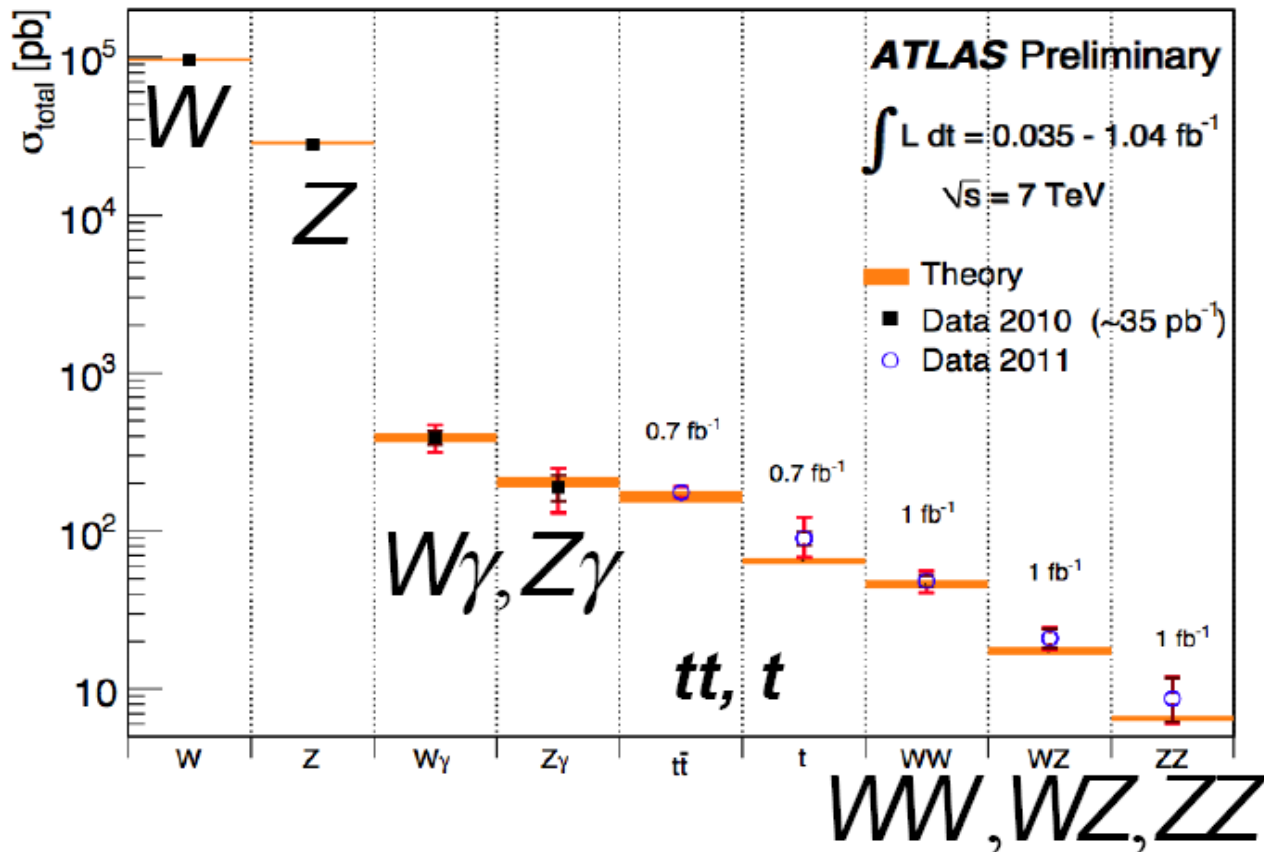
- Well measured by previous experiments
 - Inclusive cross sections, $R(W^+/W^-)$, $R(W/Z)$
 - Differential distributions, associated jet multiplicity, A_{FB} , etc.
- Yet still educational at the LHC
 - Cross sections at $\sqrt{s}=7\text{TeV}$
 - New pdf constraints possible
- “Standard candles” for high- p_T analyses
 - Calibration, alignment
 - Independent luminosity measurements



Just departure point for high- p_T
Beyond Standard Model analyses

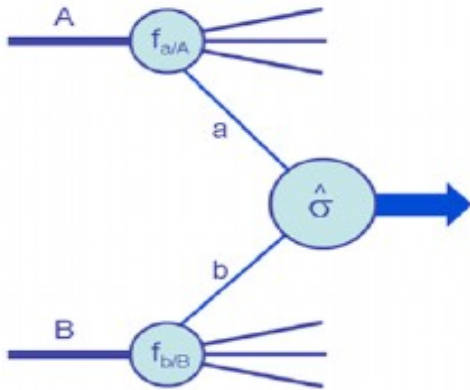
The Standard Model

- SM measurements are the foundations of all searches (33 papers on measurements to date)



Drell-Yan cross-section

- Keywords:



- factorisation μ_F and renormalisation μ_R scales
- universal parton distribution functions
- LO, NLO, NNLO matrix elements and DGLAP kernels

also depends on μ_R and μ_F , so as to cancel scale dependence in PDF's and α_s , to this order

$$\sigma_{AB} = \int dx_a dx_b f_{a/A}(x_a, Q^2) f_{b/B}(x_b, Q^2) \hat{\sigma}_{ab \rightarrow X}$$

$$\sigma_{AB} = \int dx_a dx_b f_{a/A}(x_a, \mu_F^2) f_{b/B}(x_b, \mu_F^2) \times [\hat{\sigma}_0 + \alpha_S(\mu_R^2) \hat{\sigma}_1 + \dots]_{ab \rightarrow X}$$

- All orders cross section has no dependence on μ_F and μ_R ; a residual dependence remains (to order α_s^{n+1}) for a finite order (α_s^n) calculations.

DGLAP equations

- Parton distributions used in hard scattering calculations are solutions of DGLAP equations

$$\frac{\partial q_i(x, \mu^2)}{\partial \log \mu^2} = \frac{\alpha_S}{2\pi} \int_x^1 \frac{dz}{z} \left\{ P_{q_i q_j}(z, \alpha_S) q_j\left(\frac{x}{z}, \mu^2\right) + P_{q_i g}(z, \alpha_S) g\left(\frac{x}{z}, \mu^2\right) \right\},$$

$$\frac{\partial g(x, \mu^2)}{\partial \log \mu^2} = \frac{\alpha_S}{2\pi} \int_x^1 \frac{dz}{z} \left\{ P_{g q_j}(z, \alpha_S) q_j\left(\frac{x}{z}, \mu^2\right) + P_{g g}(z, \alpha_S) g\left(\frac{x}{z}, \mu^2\right) \right\},$$

- Splitting functions have perturbative expansions

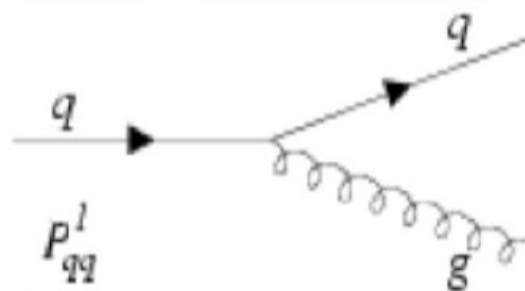
$$P_{ab}(x, \alpha_S) = P_{ab}^{(0)}(x) + \frac{\alpha_S}{2\pi} P_{ab}^{(1)}(x) + \dots$$

Thus, a full NLO calculation will contain both $\hat{\sigma}_1$ (previous slide) and $P_{ab}^{(1)}$

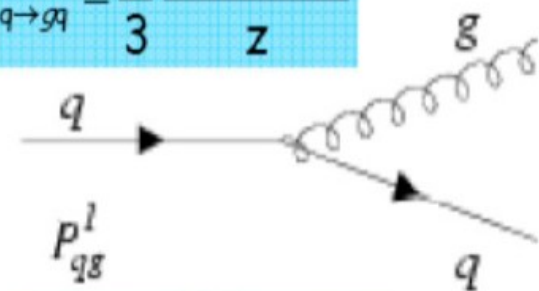
Altarelli-Parisi splitting functions

Altarelli-Parisi splitting functions:

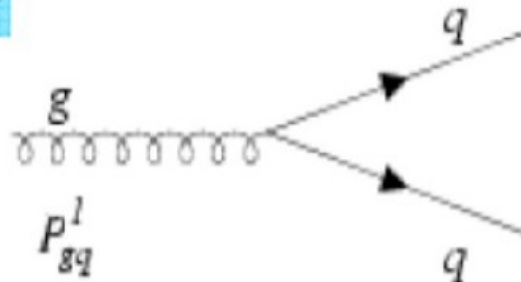
$$P_{q \rightarrow qg}^l = \frac{4}{3} \left(\frac{1+z^2}{1-z} \right)$$



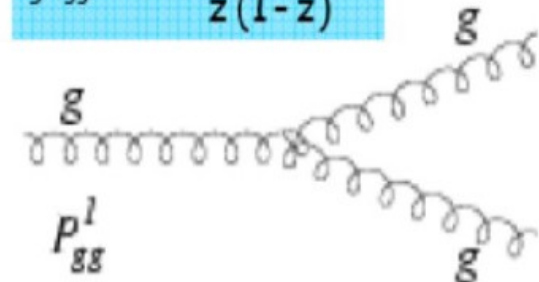
$$P_{q \rightarrow qg}^l = \frac{4}{3} \frac{1+(1-z)^2}{z}$$



$$P_{g \rightarrow qq}^l = \frac{n_f^2}{2} (z^2 + (1-z)^2)$$



$$P_{g \rightarrow gg}^l = 3 \frac{(1-z)(1-z)^2}{z(1-z)}$$



W/Z production

- Cross sections for on-shell W and Z production (in narrow width limit) given by

$$\hat{\sigma}^{q\bar{q}' \rightarrow W} = \frac{\pi}{3} \sqrt{2} G_F M_W^2 |V_{qq'}|^2 \delta(\hat{s} - M_W^2),$$

$$\hat{\sigma}^{q\bar{q} \rightarrow Z} = \frac{\pi}{3} \sqrt{2} G_F M_Z^2 (v_q^2 + a_q^2) \delta(\hat{s} - M_Z^2),$$

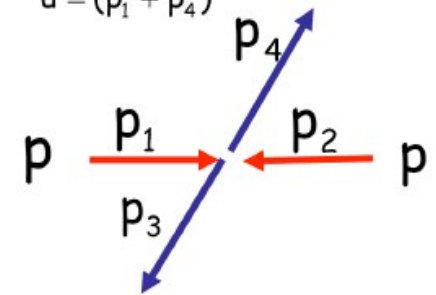
- Where $V_{qq'}$ is appropriate CKM matrix element and v_q and a_q are the vector and axial couplings of the Z to quarks
- At LO there is no α_s dependence; EW vertex only
- NLO contribution to the cross section is proportional to α_s ; NNLO to α_s^2 ; ...

Mandelstam variables :

$$\hat{s} = (p_1 + p_1)^2$$

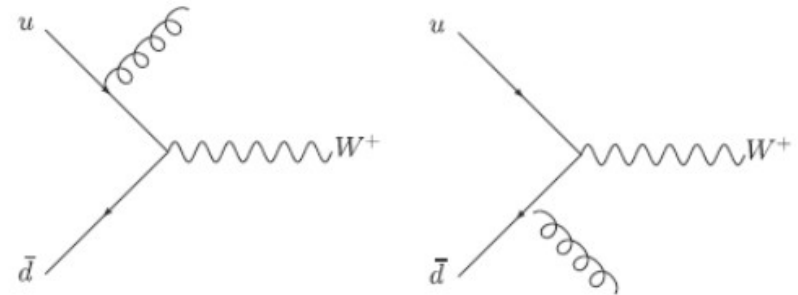
$$\hat{t} = (p_1 + p_3)^2$$

$$\hat{u} = (p_1 + p_4)^2$$



W/Z p_T distributions

- Most of W/Z produced at low p_T but can be produced at non-zero p_T due to the diagrams with emitted gluon



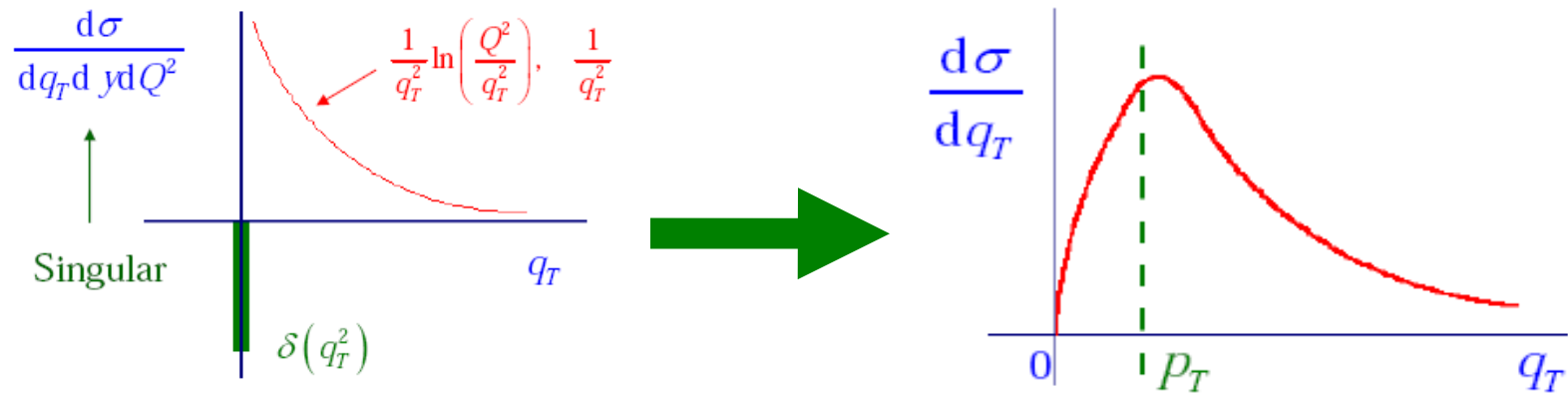
$$\sum |\mathcal{M}^{q\bar{q}' \rightarrow Wg}|^2 = \pi \alpha_S \sqrt{2} G_F M_W^2 |V_{qq'}|^2 \frac{8}{9} \frac{\hat{t}^2 + \hat{u}^2 + 2M_W^2 \hat{s}}{\hat{t}\hat{u}},$$

$$\sum |\mathcal{M}^{gq \rightarrow Wq'}|^2 = \pi \alpha_S \sqrt{2} G_F M_W^2 |V_{qq'}|^2 \frac{1}{3} \frac{\hat{s}^2 + \hat{u}^2 + 2\hat{t}M_W^2}{-\hat{s}\hat{u}},$$

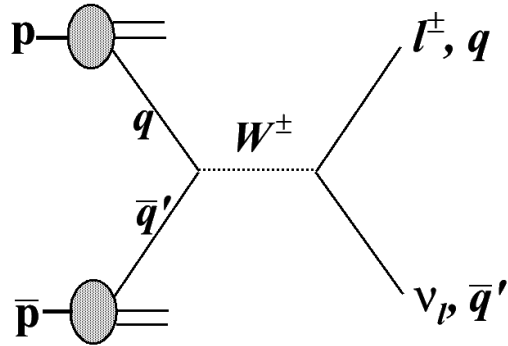
- Sum over colors and spins in initial states and average over same in final states
- Transverse momentum distribution obtained by convoluting these matrix elements with pdf's in usual way

QCD resummation

- Resummation: reorganise calculations in terms of large Logs $L(Q^2/p_T^2)$; regularised at low p_T range;
- Different schemes: CSS which includes also non-perturbative effects; Sudakov form factors; exponentiation;



Cross-section at LHC (7TeV)



$$\sigma_{W^+ \rightarrow \ell \nu}^{NNLO} = 6.15 \text{ nb}$$

$$\sigma_{W^- \rightarrow \ell \nu}^{NNLO} = 4.3 \text{ nb}$$

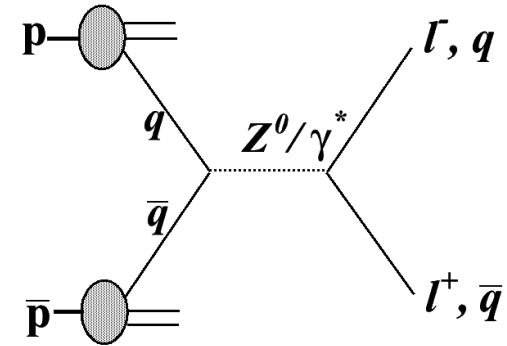
$$\sigma_{W \rightarrow \ell \nu}^{NNLO} = 10.45 \text{ nb}$$

$$\sigma(W^+) \neq \sigma(W^-)$$

W^+ production: $u\bar{d} + c\bar{s}$

W^- production: $d\bar{u} + s\bar{c}$

Z production: $u\bar{u} + d\bar{d} + s\bar{s} + c\bar{c} + b\bar{b}$

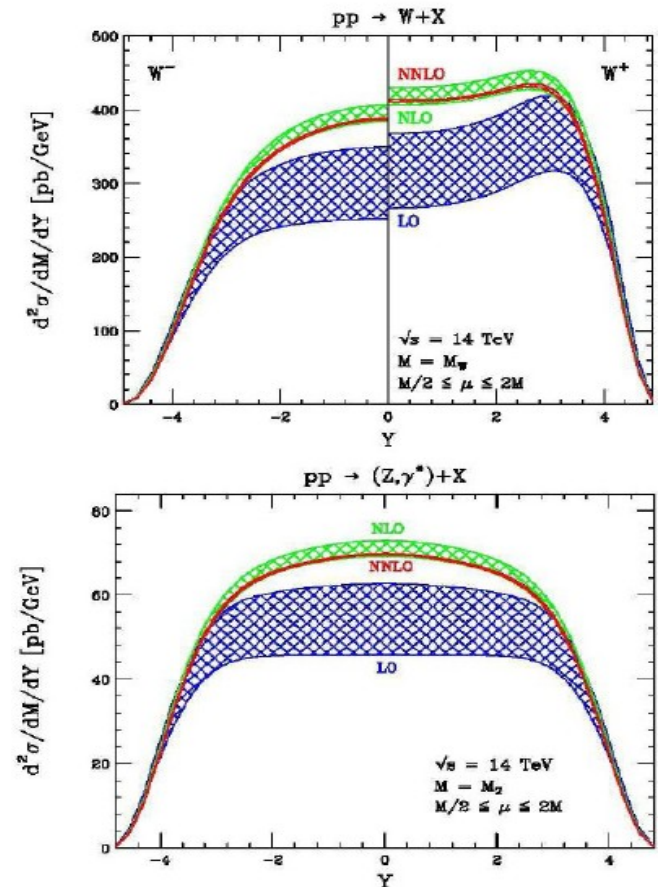


$$\sigma_{Z/\gamma^* \rightarrow \ell \bar{\ell}}^{NNLO} = 0.989 \text{ nb}$$

- Test QCD (up to NNLO) in production
 - Hard and soft gluon emission
- Sensitive to parton distribution functions
- Extract electroweak parameters
 - $\sin\Theta_W$, m_W , quark-boson couplings

Monte Carlo simulations

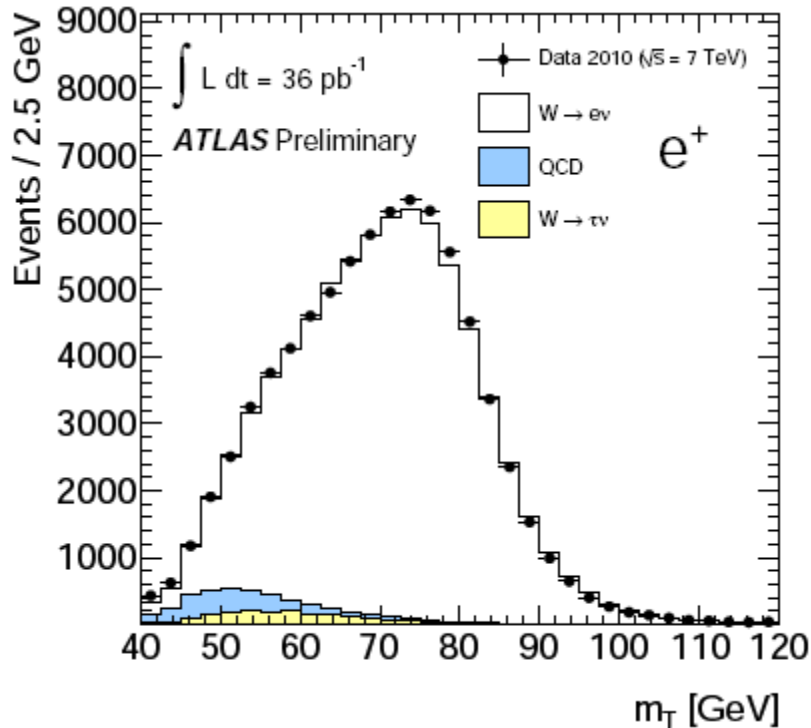
- Base-line generators:
 - Pythia, Herwig (LO),
 - MCatNLO (NLO)
 - POWHEG (NLO)
- Used as components of for cross-checks
 - FEWZ: complete NLO, NNLL
 - ResBos: NNLL resummation
 - Horace: full 1-loop electroweak
 - PHOTOS: final state QED (exponentiated)



Event selection

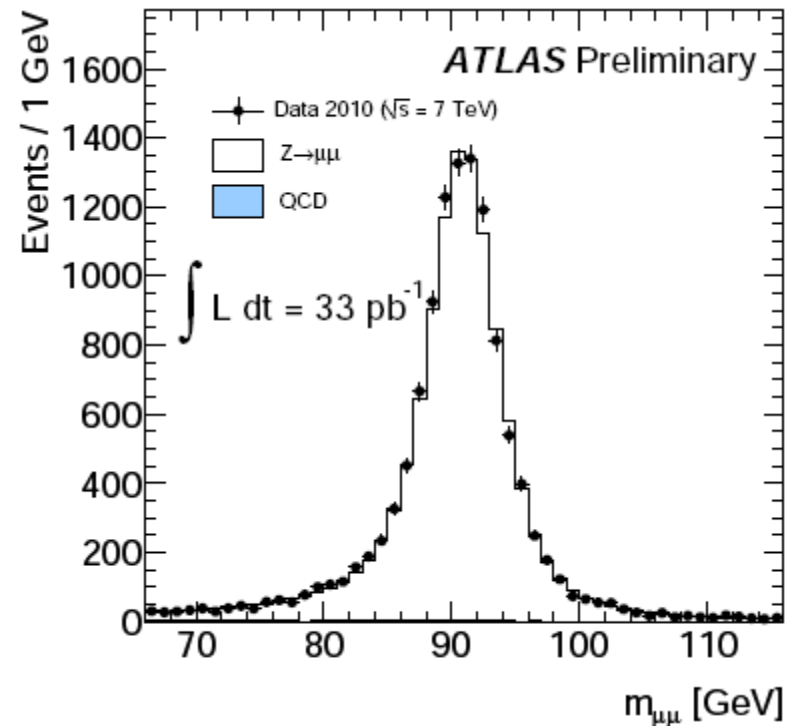
$$W \rightarrow \ell \nu$$

- One e/μ with $p_T > 20$ GeV
- $E_T^{\text{miss}} > 25$ GeV
- $m_T(\ell, E_T^{\text{miss}}) > 40$ GeV

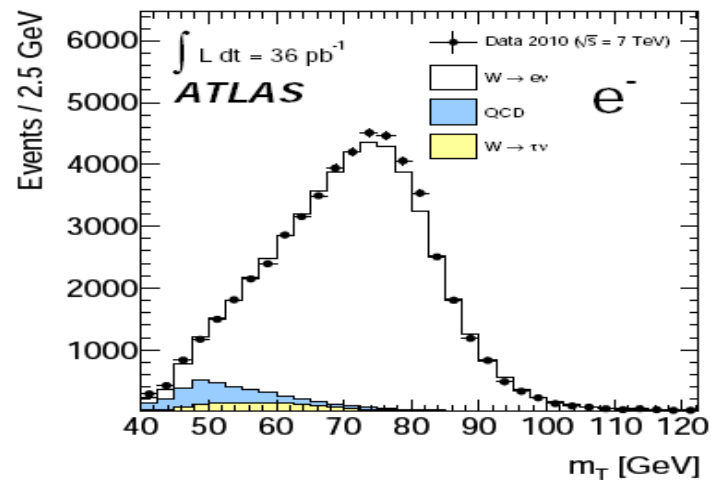
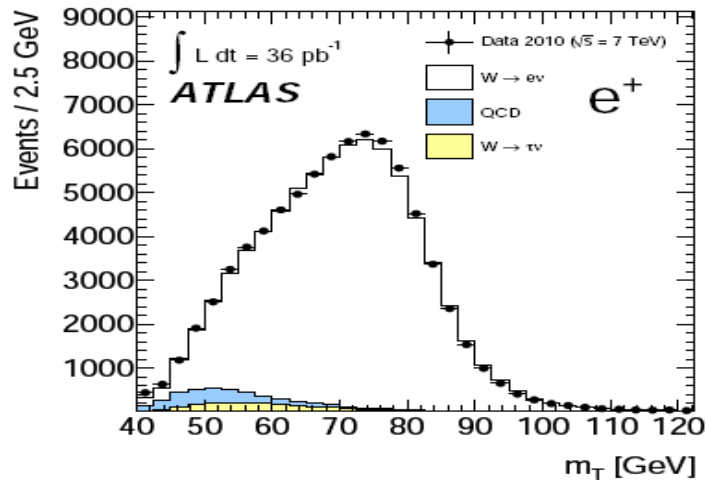
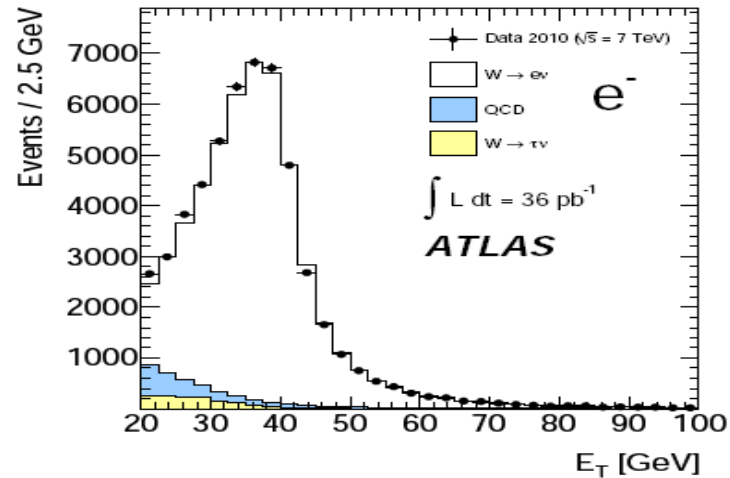
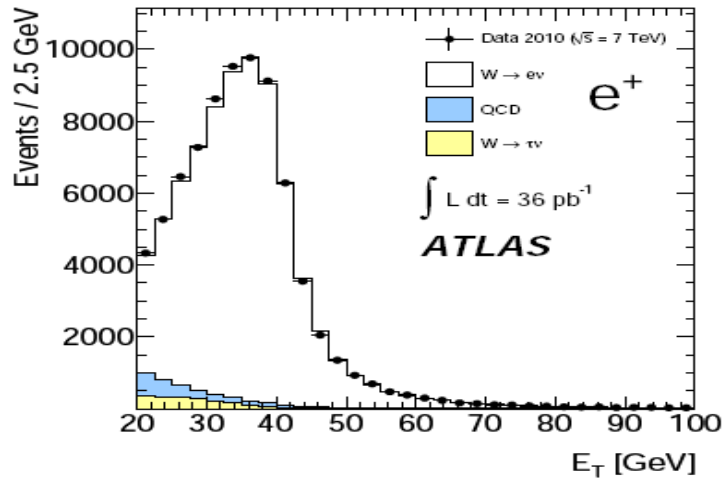


$$Z \rightarrow \ell\ell$$

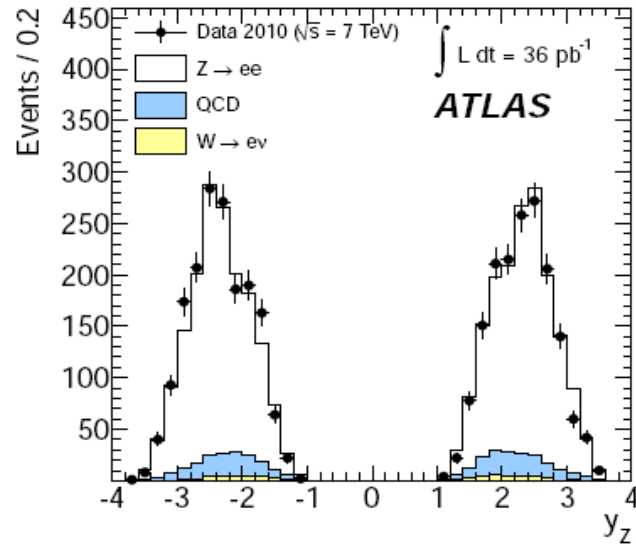
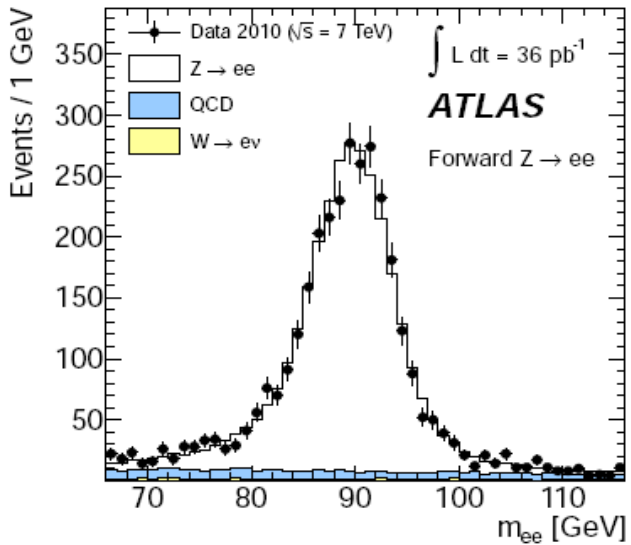
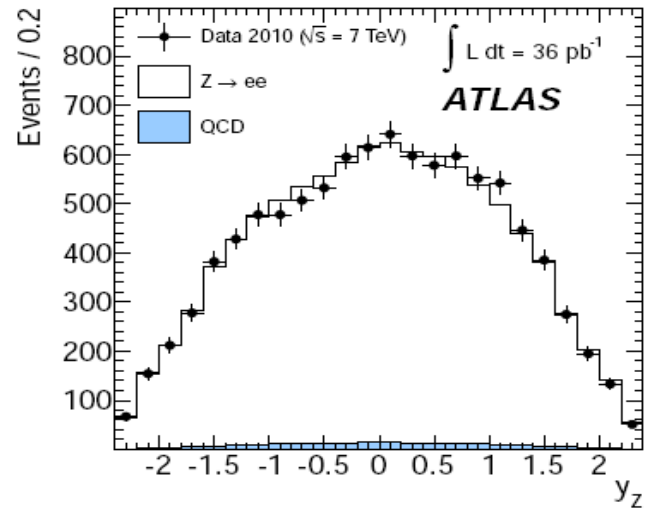
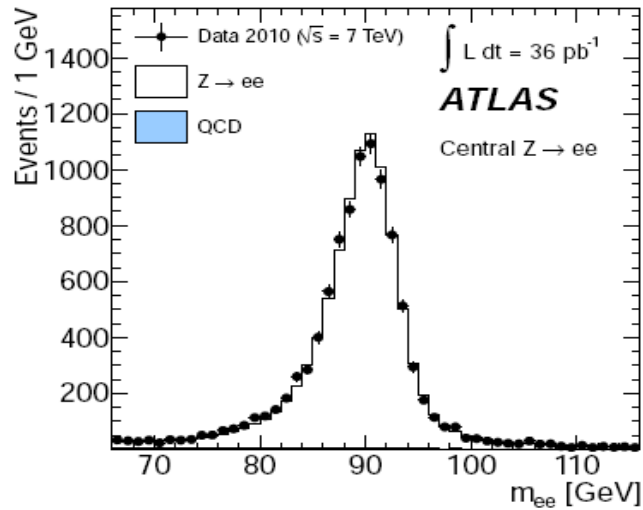
- Two e/μ with $p_T > 20$ GeV
- $m_{\ell\ell} = 66\text{--}116$ GeV



Event selection

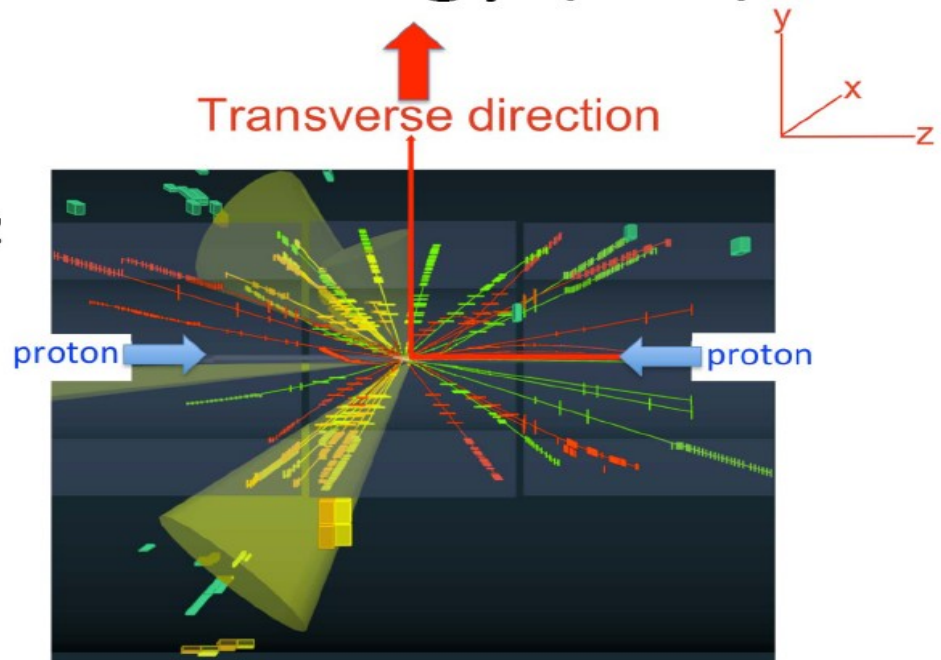


Event selection



Missing transverse energy

- In pp collisions at the LHC, a significant, unmeasured amount of energy escapes in z direction
- Total initial and final momentum is zero in transverse direction
- Imbalance of energy in transverse direction signals presence of weakly or non-interacting particles such as neutrinos



$$E_T^{miss} = \sqrt{(E_x^{miss})^2 + (E_y^{miss})^2}$$

$$E_{x(y)}^{miss} = - \sum_{particles} E_{x(y)}$$

$$\sum E_T = \sum_{particles} E_T$$

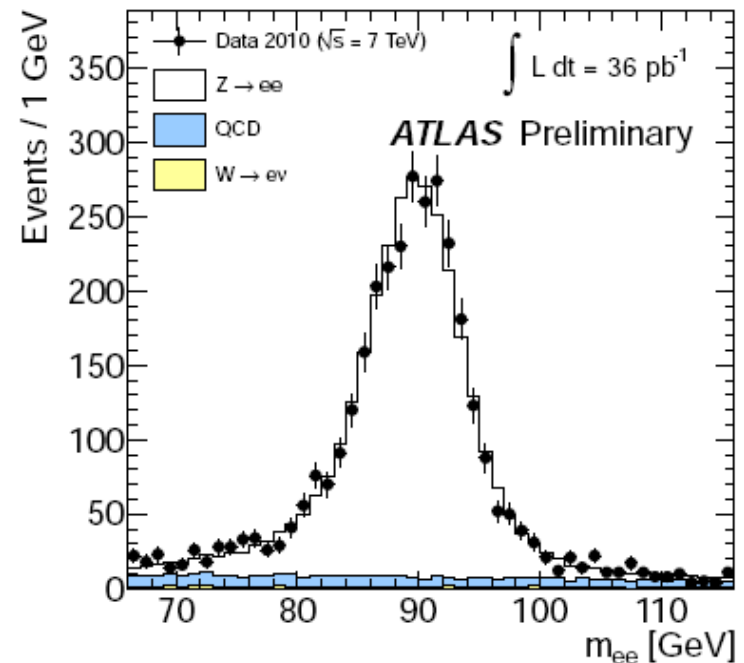
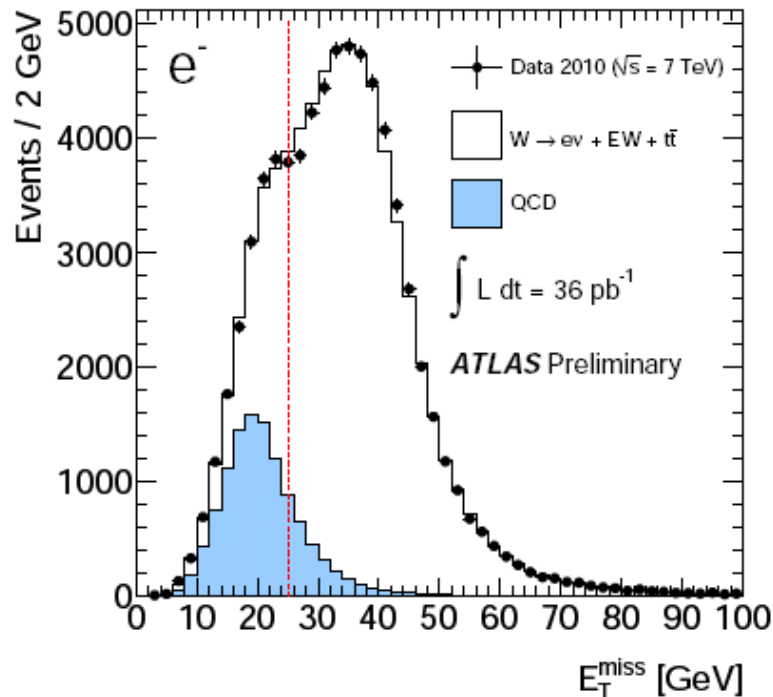
QCD background estimation

$W \rightarrow e\nu$: template fit to E_T^{miss} . Template derived from data with inverted electron ID and isolation.

$Z \rightarrow ee$: template fit to m_{ee} to a sample with looser electron ID, extrapolated to the signal region.

$W \rightarrow \mu\nu$: matrix method using track isolation.

$Z \rightarrow \mu\mu$: ABCD method with track isolation in $m_{\mu\mu}$ side-band.



Cross-section measurement

$$\sigma = \frac{N_{\text{obs}} - N_{\text{bkg}}}{A \cdot C \cdot \int dt \mathcal{L}}$$

N_{obs} : number of observed events in the signal region

N_{bkg} : estimated number of background events

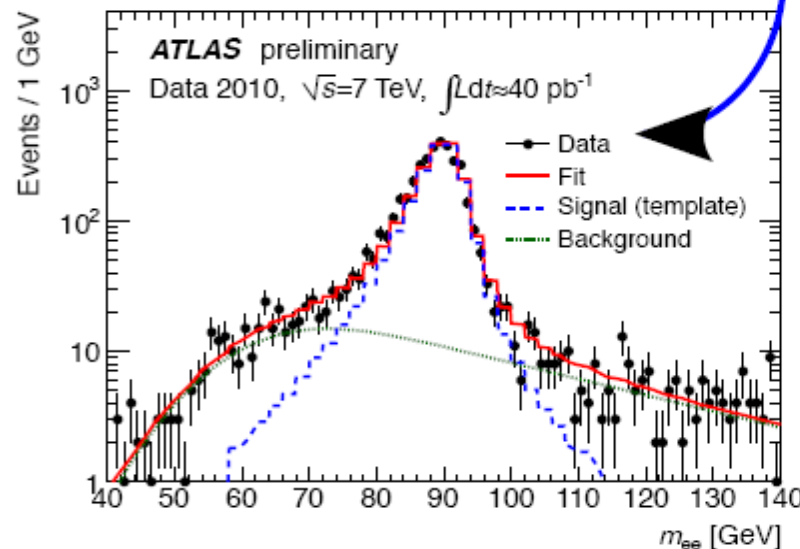
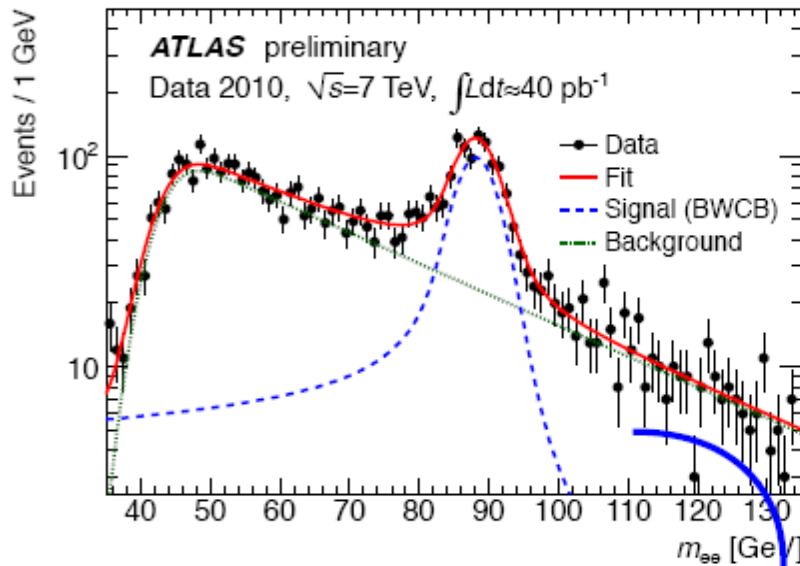
- EW backgrounds are estimated with Monte Carlo, constrained to data with performance scale factors.
- QCD backgrounds are estimated with **data-driven** methods.

A : kinematic acceptance factor, estimated with generator-level Monte Carlo.

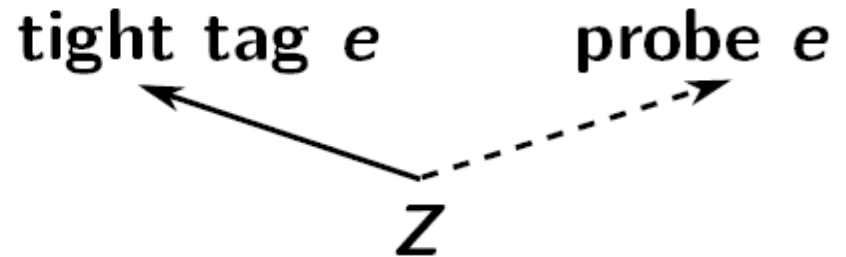
C : summarizes reconstruction efficiency, estimated with reconstructed Monte Carlo, corrected with **scale factors**.

$\int dt \mathcal{L}$: integrated luminosity.

Scale factor: tag-and-probe studies



apply ID

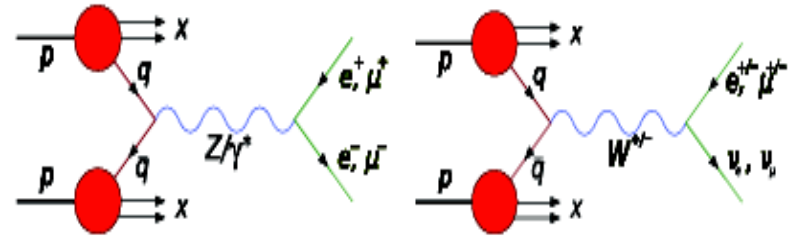


- “Tag” events with sufficient purity, leaving an unbiased “probe” object.
- Measure probe ID efficiency *in situ*.
- Constrains the performance of our object identification.
- Derive **scale factors** for correcting our simulation.

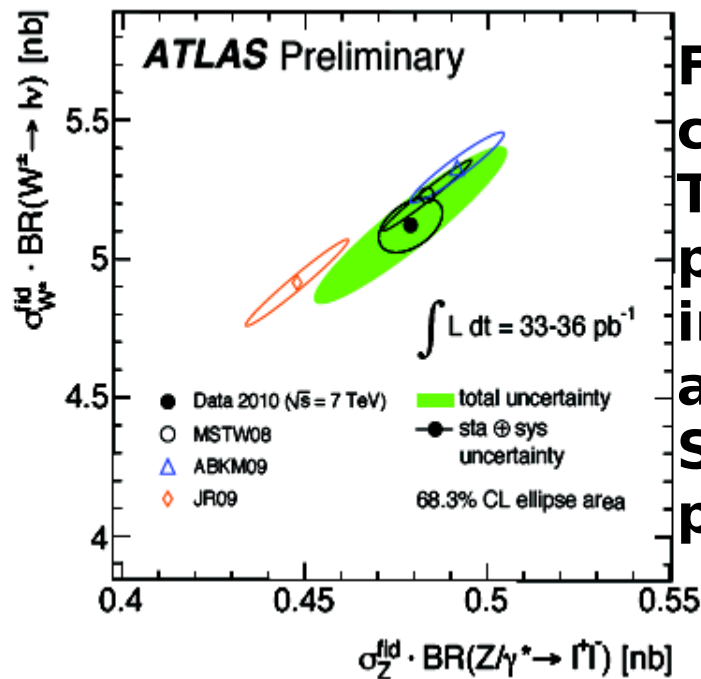
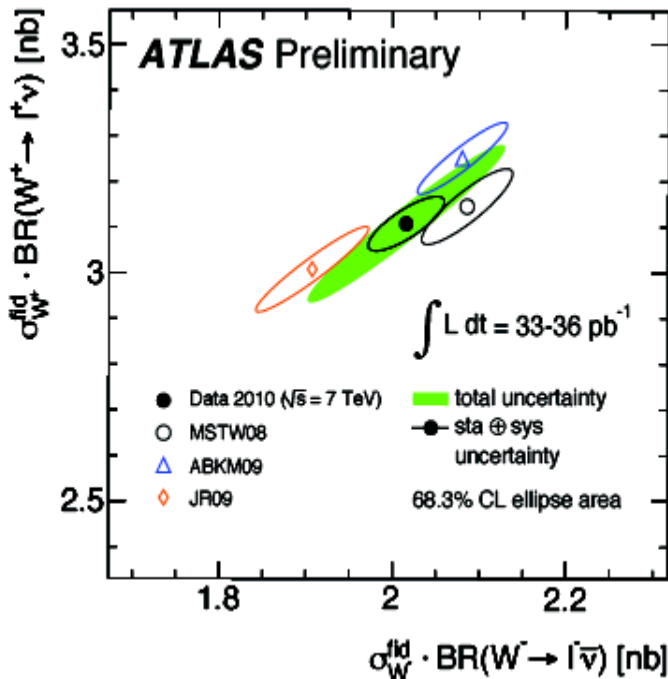
Systematic error

	$\delta\sigma_{W\pm}$	$\delta\sigma_{W+}$	$\delta\sigma_{W-}$	$\delta\sigma_Z$
Trigger	0.4	0.4	0.4	<0.1
Electron reconstruction	0.8	0.8	0.8	1.6
Electron identification	0.9	0.8	1.1	1.8
Electron isolation	0.3	0.3	0.3	—
Electron energy scale and resolution	0.5	0.5	0.5	0.2
Non-operational LAr channels	0.4	0.4	0.4	0.8
Charge misidentification	0.0	0.1	0.1	0.6
QCD background	0.4	0.4	0.4	0.7
Electroweak+ $t\bar{t}$ background	0.2	0.2	0.2	<0.1
E_T^{miss} scale and resolution	0.8	0.7	1.0	—
Pile-up modeling	0.3	0.3	0.3	0.3
Vertex position	0.1	0.1	0.1	0.1
$C_{W/Z}$ theoretical uncertainty	0.6	0.6	0.6	0.3
Total experimental uncertainty	1.8	1.8	2.0	2.7
$A_{W/Z}$ theoretical uncertainty	1.5	1.7	2.0	2.0
Total excluding luminosity	2.3	2.4	2.8	3.3
Luminosity	3.4			

W, Z inclusive measurements



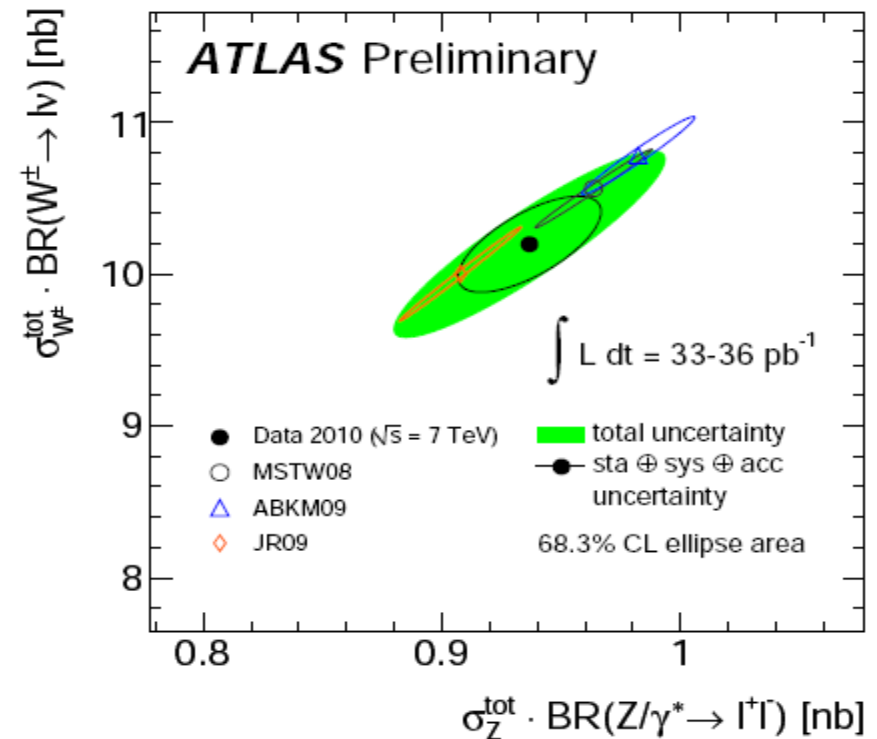
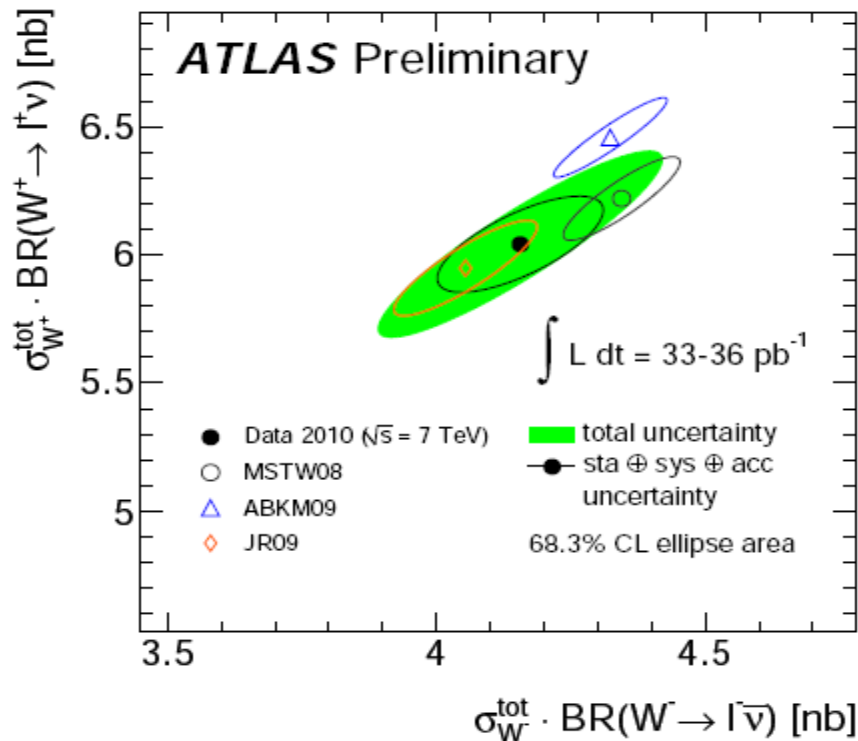
- Inclusive in number of jets, allow to reach high accuracy in QCD predictions.
- Computation available at NNLO on the total cross-section, at NNLO error dominated by PDF's



Fiducial cross-section. Theory predictions including acceptance. Sensitive to pdf's set.

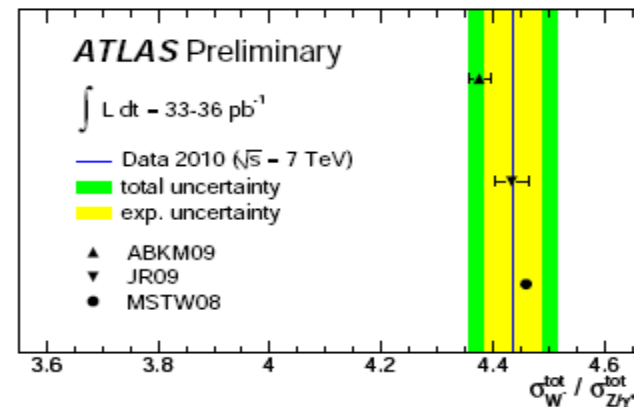
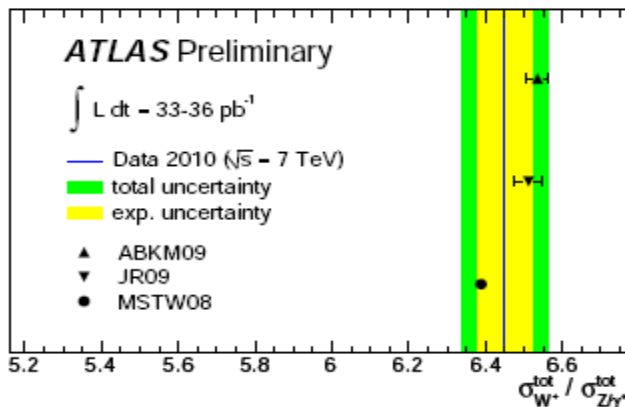
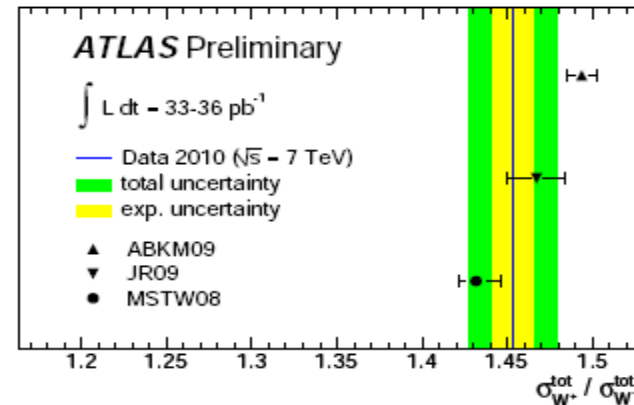
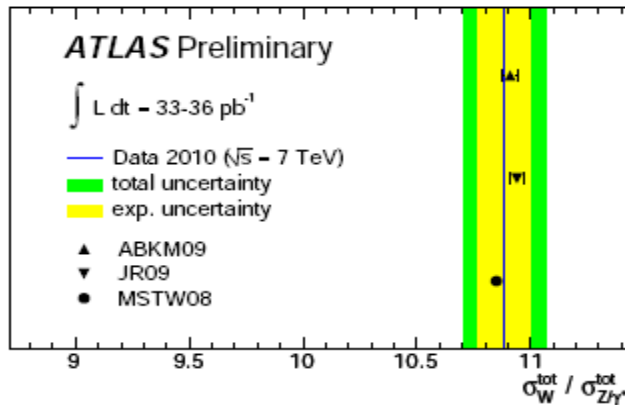
Theory comparison: total cross-section

- Overall remarkable agreement with NNLO PDF predictions
- A few differences between different PDFs (w/ only 68% CL PDF errors)
- Comparing total cross sections, the acceptance uncertainty accounts for effect of different PDFs on the unmeasured phase space ...



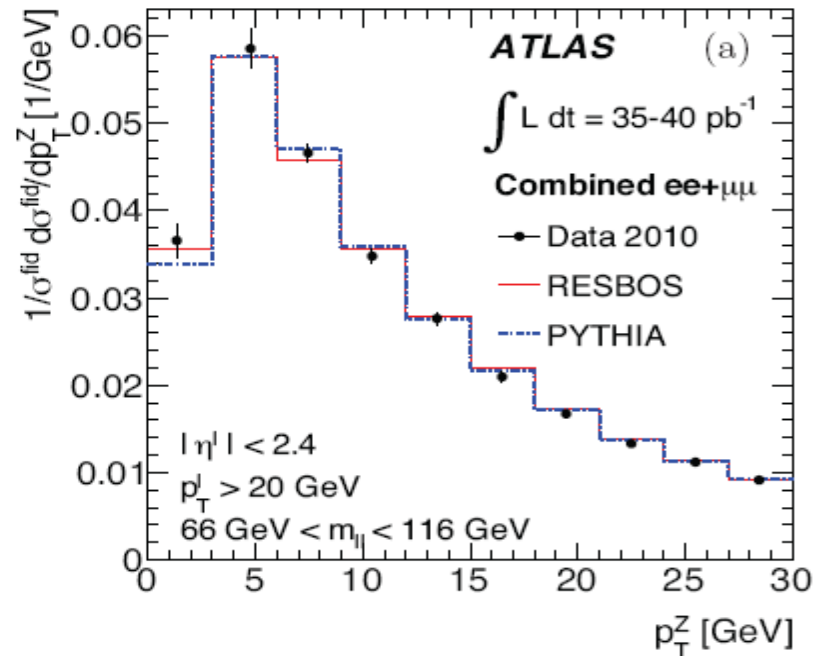
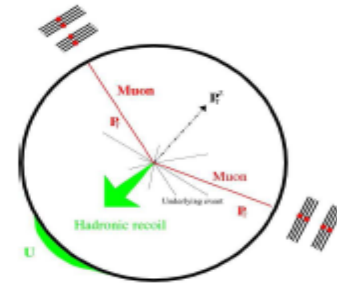
Theory comparison: total cross-section ratios

- W^\pm/Z , W^+/W^- ratios profit from exp. and theor. systematics cancellation
- W^\pm/Z ratio measured with total uncert. of 1.5%, W^+/W^- with 1.7%



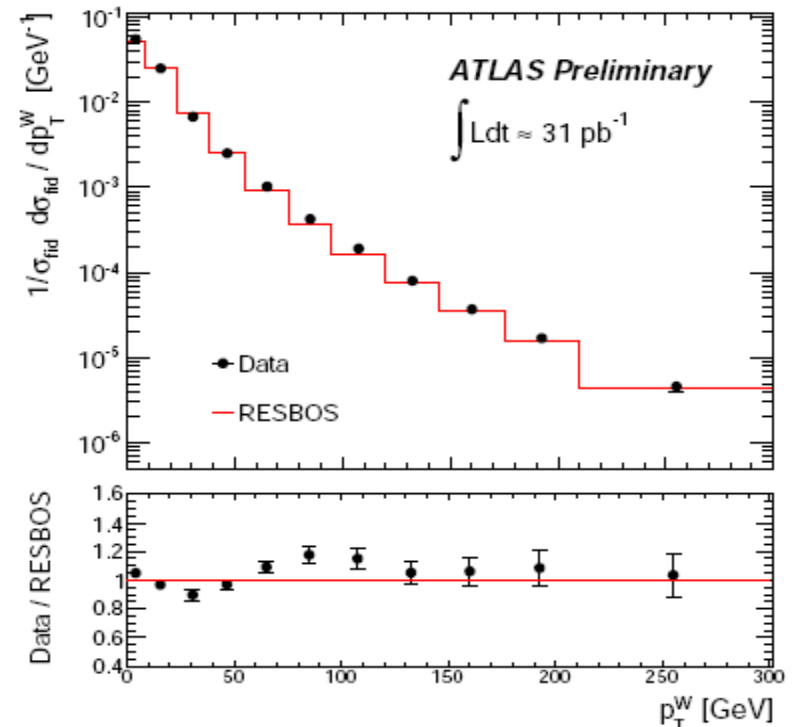
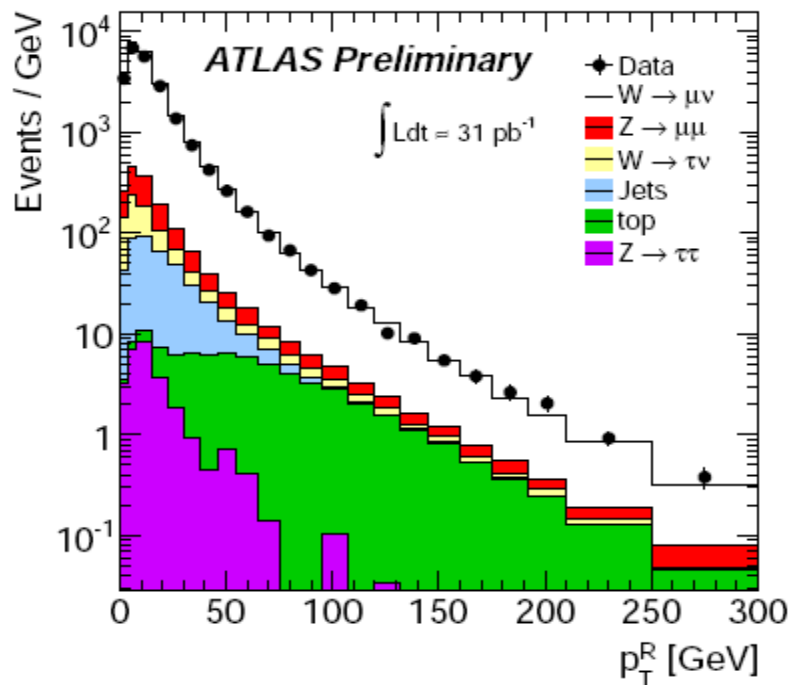
Z boson p_T measurement

- Important for modeling high- p_T lepton kinematics.
- At leading order, $p_T^{W/Z} = 0$
- Non-zero $p_T^{W/Z}$ is generated through the hadronic recoil of ISR, p_T^R .
- p_T^Z reconstructed directly from $p_T(\mu_1) + p_T(\mu_2)$, while p_T^W reconstructs p_T^R .
- Detector and FSR effects removed with a bin-by-bin unfolding.
- 3-4% precision per bin.



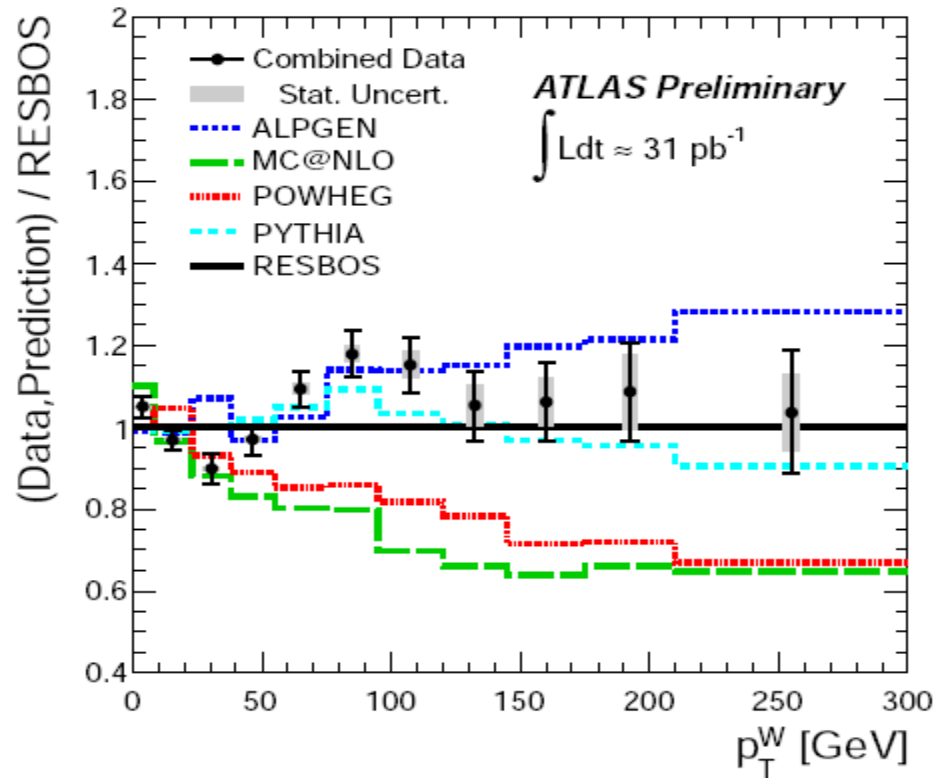
W boson p_T measurement

- Necessary for a future precision W mass measurement.
- Detector and FSR effects removed by inverting a response matrix parametrizing the probabilistic mapping of p_T^R to p_T^W .



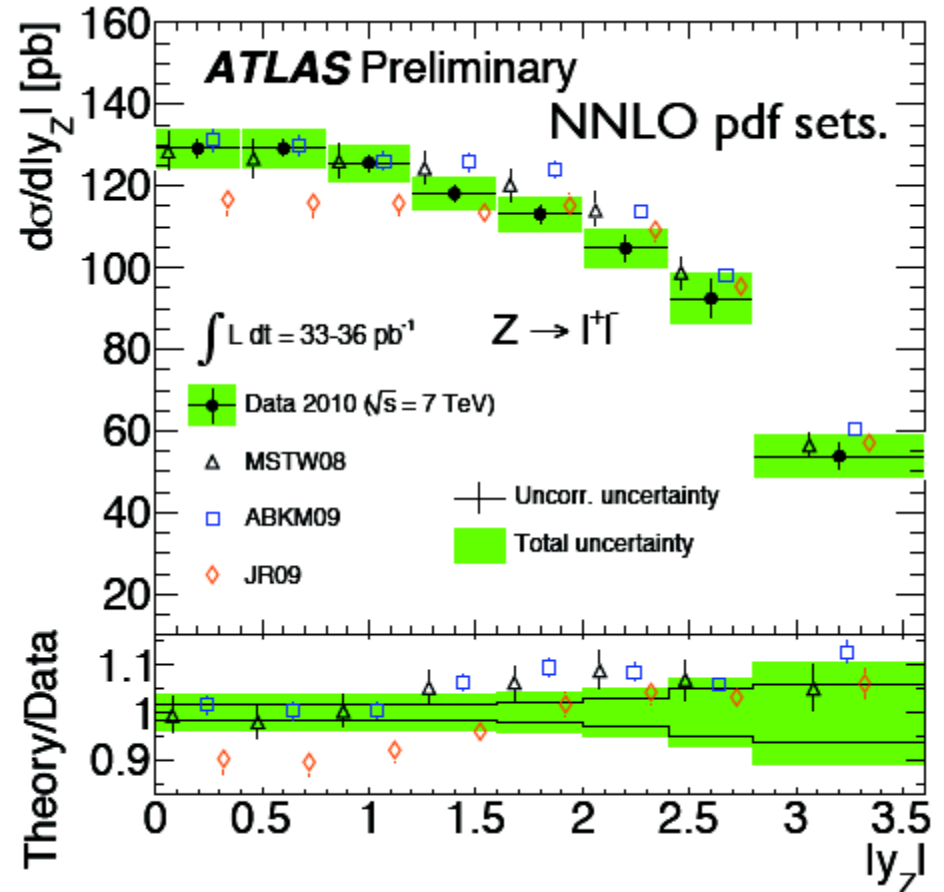
W, Z boson p_T reweighting

- The modeling of $d\sigma/dp_T^{W/Z}$ can have significant effects on the expected efficiency and acceptance.
- NLO generators MC@NLO and POWHEG have deficits at high $p_T^{W/Z}$.
- NLO effects are important at high $p_T^{W/Z}$ because the W/Z is polarized by higher order QCD.
- $W \rightarrow l\nu$ and $Z \rightarrow ll$ cross section measurements use MC@NLO reweighted to match $p_T^{W/Z}$ for LO Pythia, which agrees with the data because it has been tuned well to the Tevatron data.



Z differential

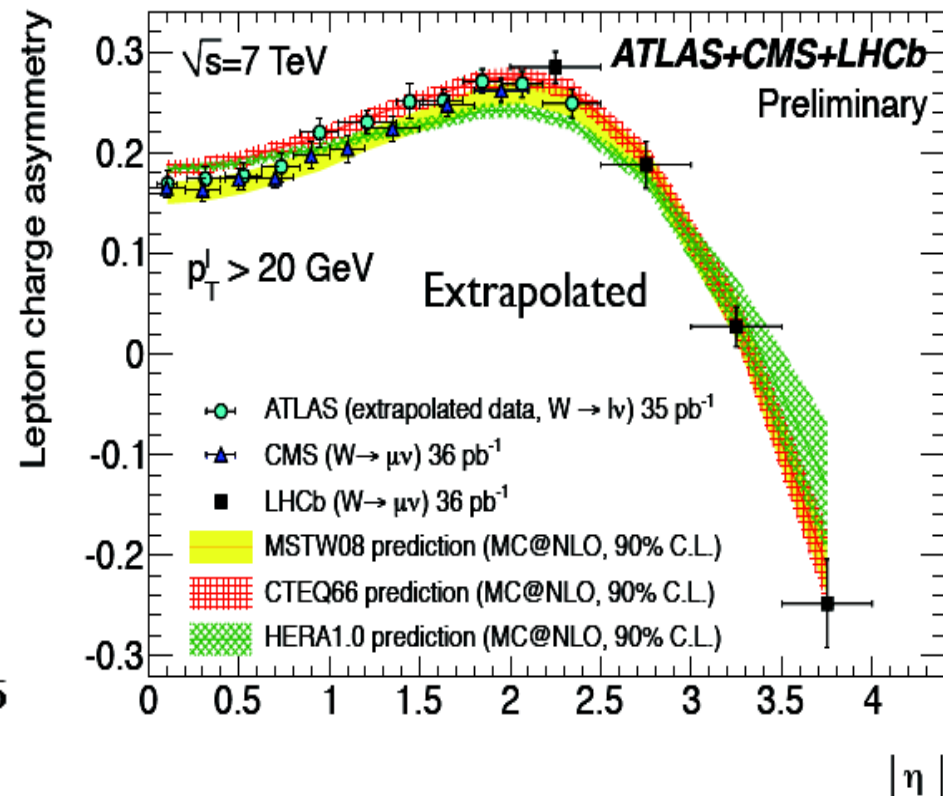
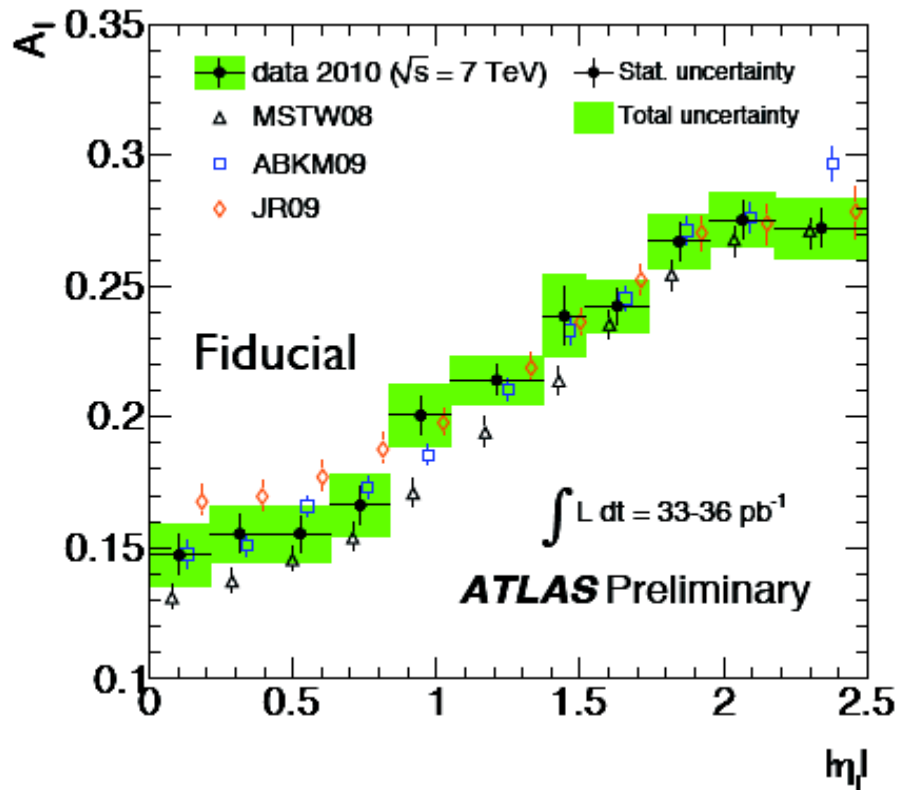
- Inclusive production as a function of the Z pseudorapidity.
- Lepton flavours combined together taking into account all correlations.
- Z rapidity reaches $|y| < 3.5$ with special electron reconstruction outside tracking volume ($|y| < 2.5$)



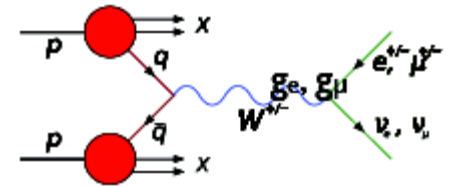
W[±] asymmetry

$$A(\eta_l) = \frac{\sigma^{W^+}(\eta_l) - \sigma^{W^-}(\eta_l)}{\sigma^{W^+}(\eta_l) + \sigma^{W^-}(\eta_l)}$$

- Asymmetry induced by the different flavours contributing to W⁺ and W⁻ production and by asymmetry in flavour content of pp interaction



Lepton universality



$$R_W = \frac{\sigma_W^e}{\sigma_W^\mu} = \frac{Br(W \rightarrow e\nu)}{Br(W \rightarrow \mu\nu)} = 1.006 \pm 0.004 \text{ (sta)} \pm 0.006 \text{ (unc)} \pm 0.023 \text{ (cor)} = 1.006 \pm 0.024$$

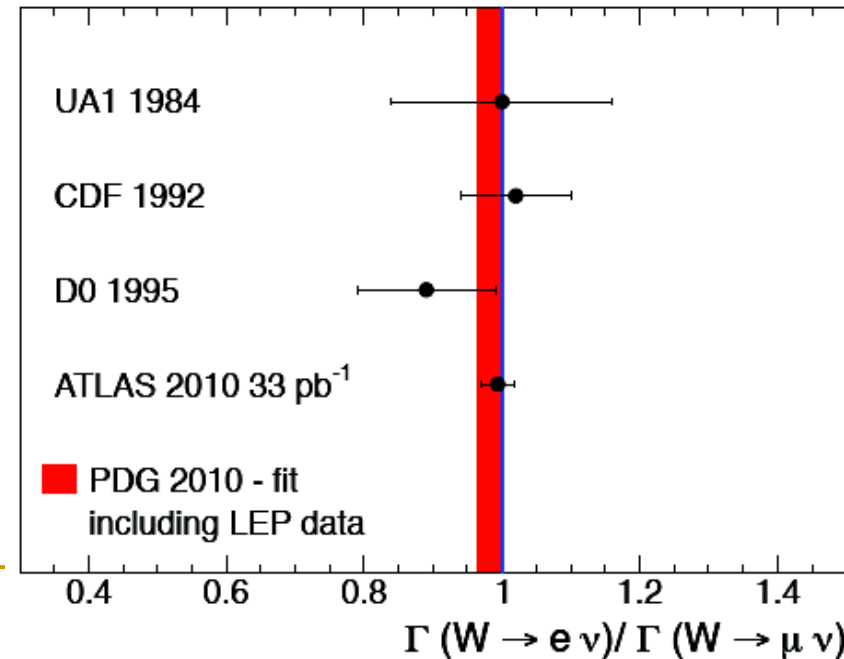
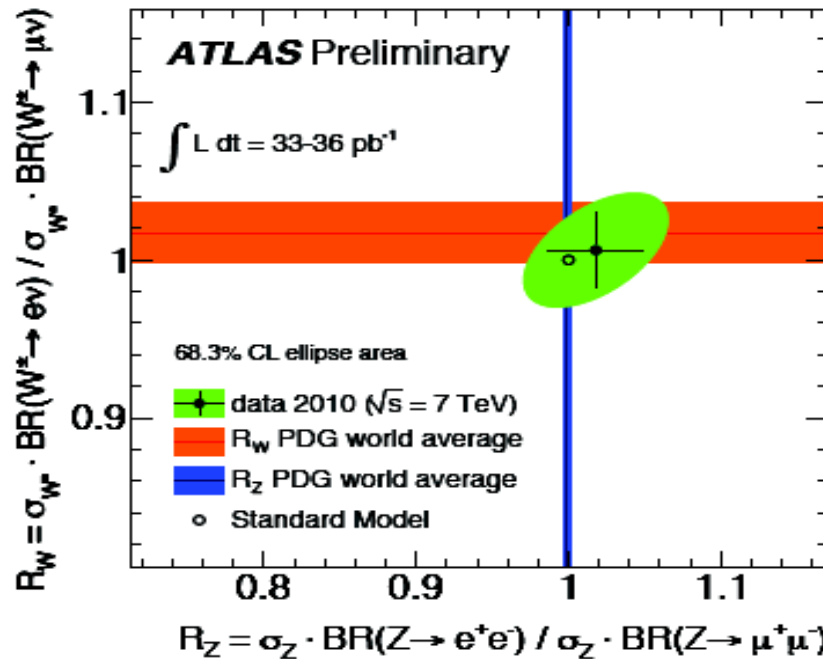
$$R_Z = \frac{\sigma_Z^e}{\sigma_Z^\mu} = \frac{Br(Z \rightarrow ee)}{Br(Z \rightarrow \mu\mu)} = 1.018 \pm 0.014 \text{ (sta)} \pm 0.016 \text{ (unc)} \pm 0.028 \text{ (cor)} = 1.018 \pm 0.031$$

$$R_W^{\text{PDG}} = 1.017 \pm 0.019$$

$$R_Z^{\text{PDG}} = 0.991 \pm 0.0024$$

$$R_W = g_e^2 / g_\mu^2 \text{ (neglecting loop correction)}$$

$$R_Z = \frac{g_e^2 (2\sin\theta_{eW} - 1) + 4\sin^2\theta_{eW}}{g_\mu^2 (2\sin\theta_{\mu W} - 1) + 4\sin^2\theta_{\mu W}}$$



Incredible number of V+jets events at 7 TeV

$$E_T^{n^{\text{th}} \text{ jet}} > 25 \text{ GeV}$$

plus standard cuts on jet (anti- k_T , $R=0.4$)
rapidity, lepton, missing E_T

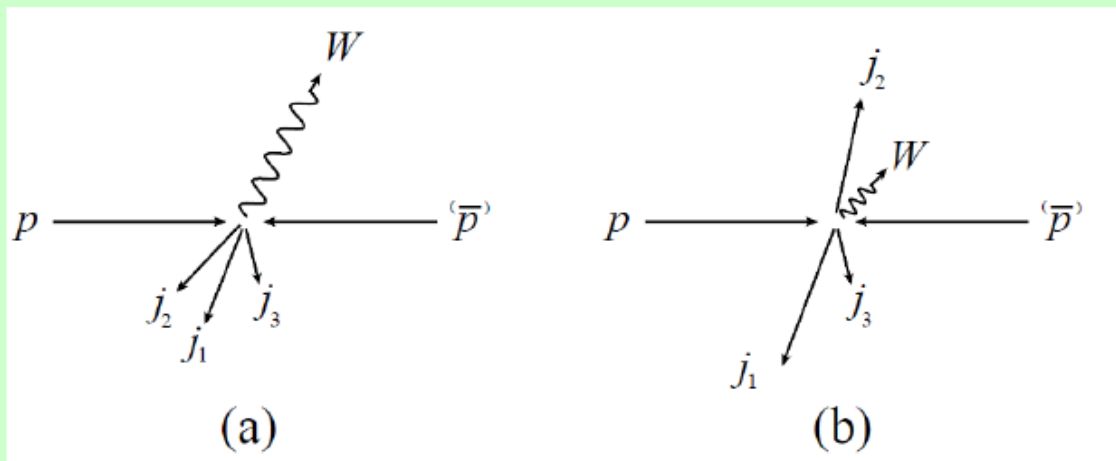
$$W^\pm \rightarrow e^\pm \nu \text{ or } \mu^\pm \nu$$

$$Z^0 \rightarrow e^+ e^- \text{ or } \mu^+ \mu^-$$

$\int \mathcal{L} dt = 2 \text{ fb}^{-1}$	
$\sigma(W^\pm j) \approx 800 \text{ pb}$	1,600,000
$\sigma(W^\pm jj) \approx 200 \text{ pb}$	400,000
$\sigma(W^\pm jjj) \approx 45 \text{ pb}$	90,000
$\sigma(Zj) \approx 80 \text{ pb}$	160,000
$\sigma(Zjj) \approx 20 \text{ pb}$	40,000
$\sigma(Zjjj) \approx 4.8 \text{ pb}$	9,600

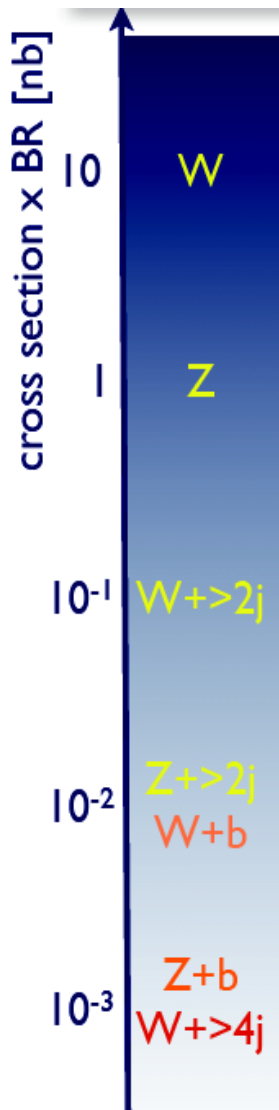
V+jets dynamics at 7 TeV

- Can now easily get events with $E_T^{\text{jet}} \gg M_W$.
- “Soft W” enhancement \rightarrow tend to be jet dominated:
(b) preferred over (a).



- Wide dynamic range at LHC
 \rightarrow Importance of choosing reasonable scales $\mu_{R,F}$

W/Z + jets physics

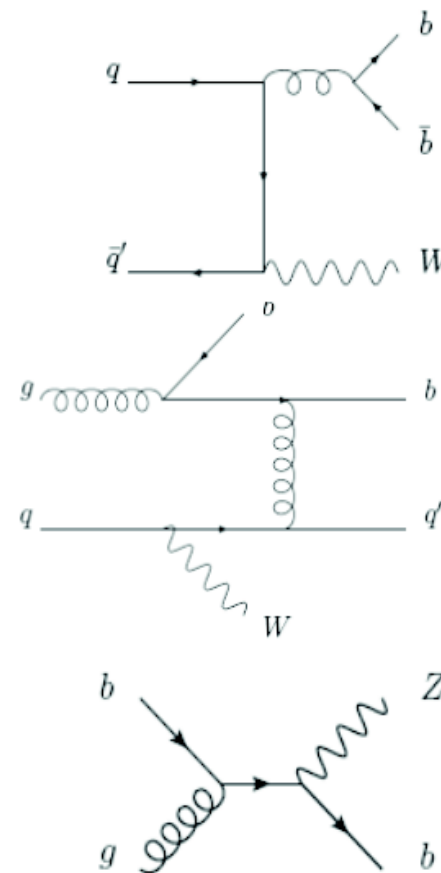


- **W/Z+jets physics is a fundamental ingredient for reestablishing the Standard Model (SM) in pp collisions at 7 TeV**

✓ larger available energy than at Tevatron:
 \Rightarrow more jets; larger kinematic reach
 \Rightarrow cross sections spanning several orders of magnitude

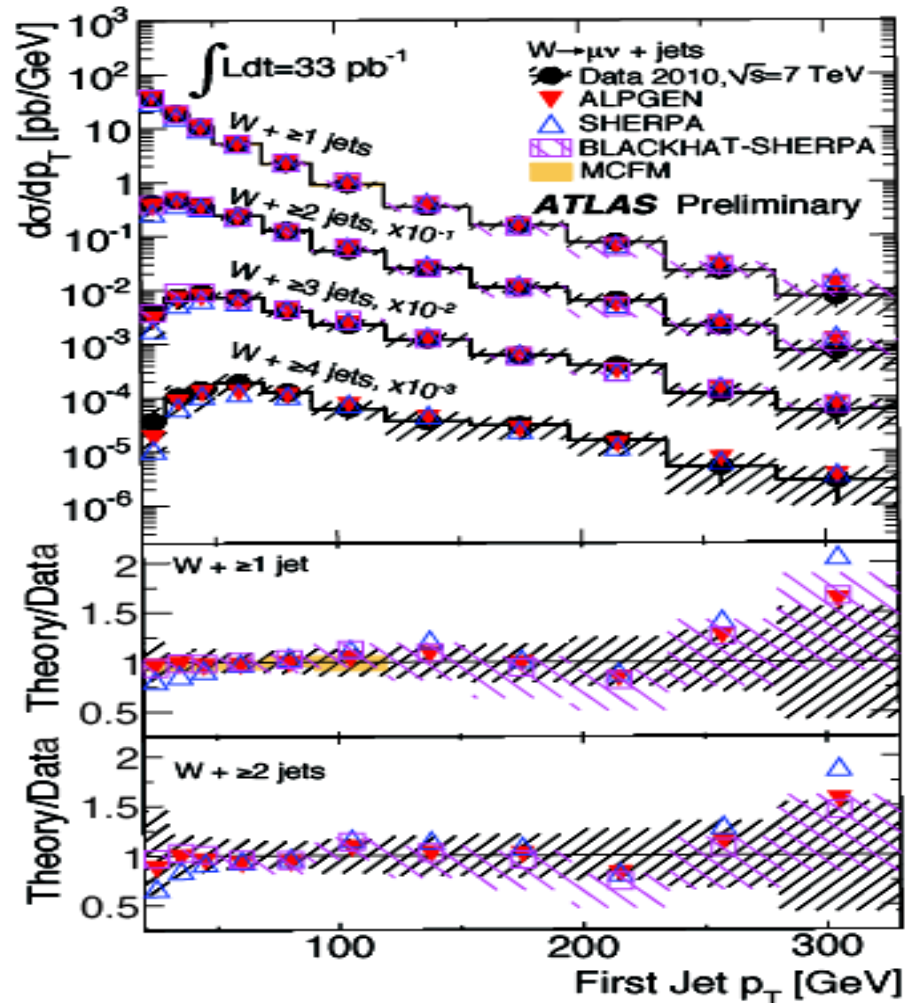
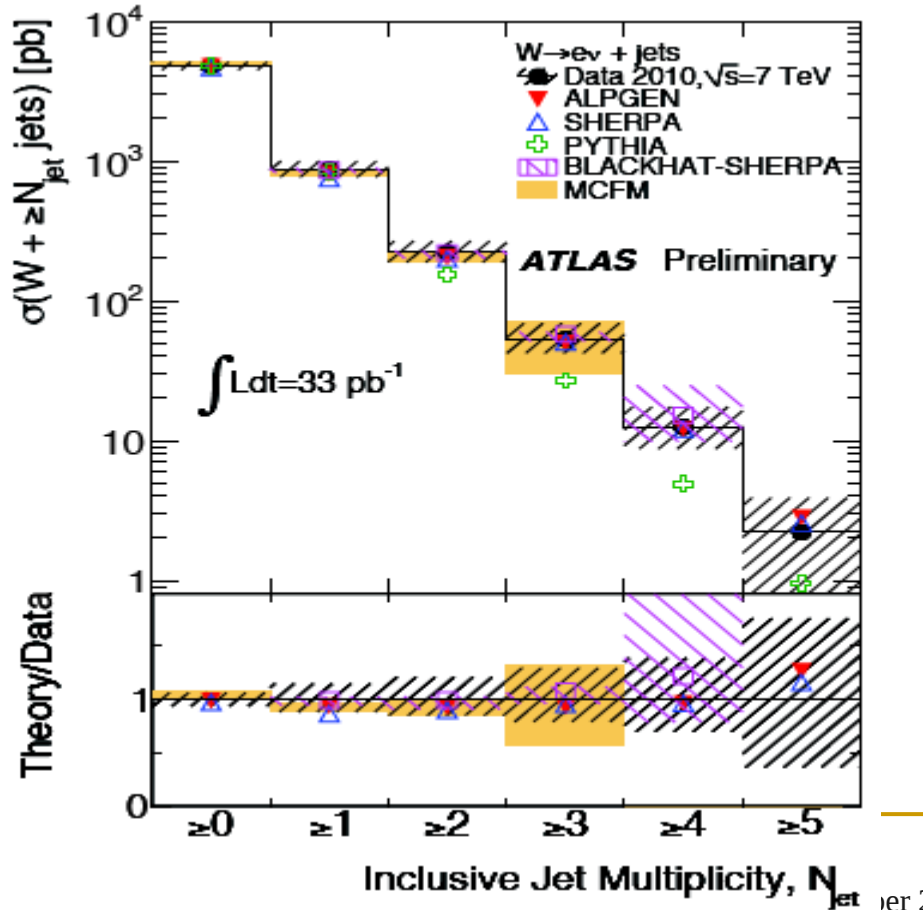
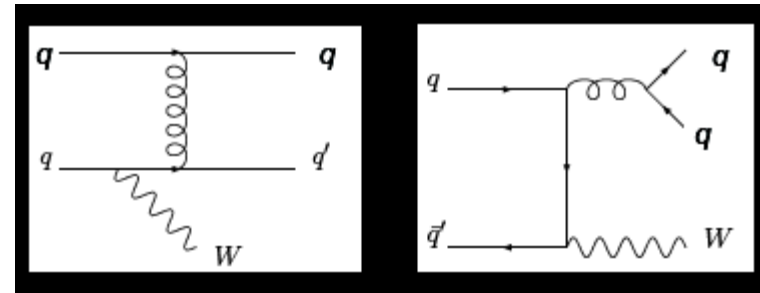
✓ higher relevance of processes initiated by qg and gg scattering
 \Rightarrow different contribution to the cross section compared to Tevatron
 \Rightarrow processes with heavy flavour in the initial state become important

- **compelling test for the new NLO pQCD calculations of W/Z+(b)jets (up to 4jet for light- and 2 for b-jets)**



W/Z + jets

- Measure up to 4-5 jets and up to $E_T = 300$ GeV

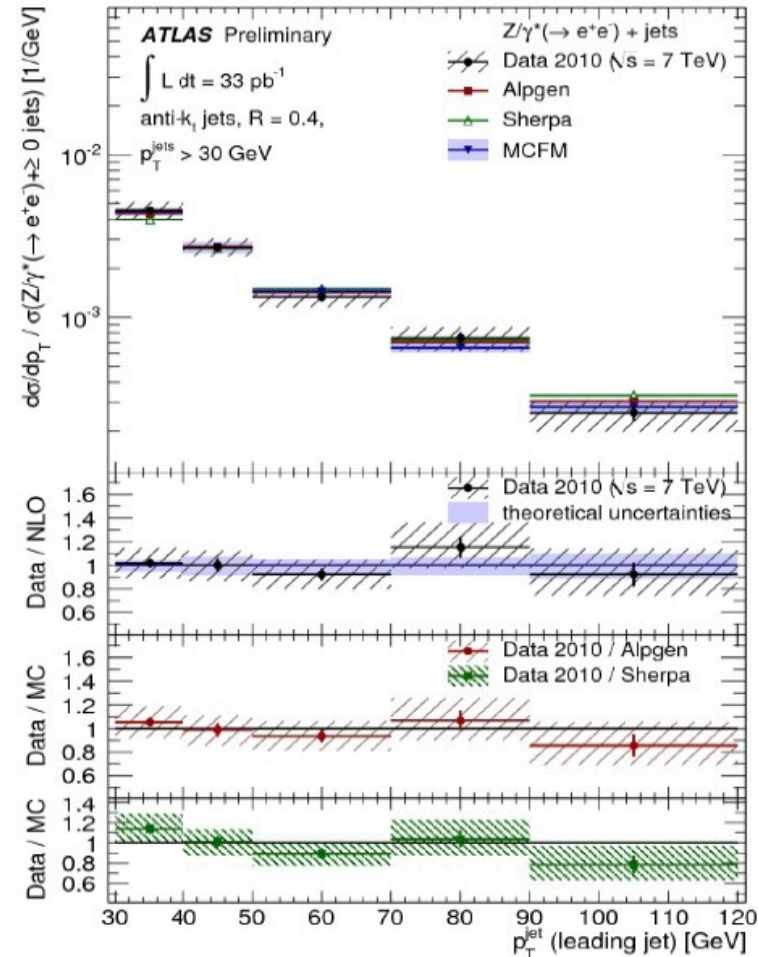


W/Z+jets

- cross section measured as a function of several kinematic variables (see end of this talk)
- **very good agreement with NLO** predictions from MCFM and Blackhat-Sherpa in the total and differential cross sections
- good agreement with matched LO prediction from AlpGen and Sherpa once normalized to the NNLO prediction
- Poor agreement with LO PYTHIA in the high jet multiplicity

dominant systematics

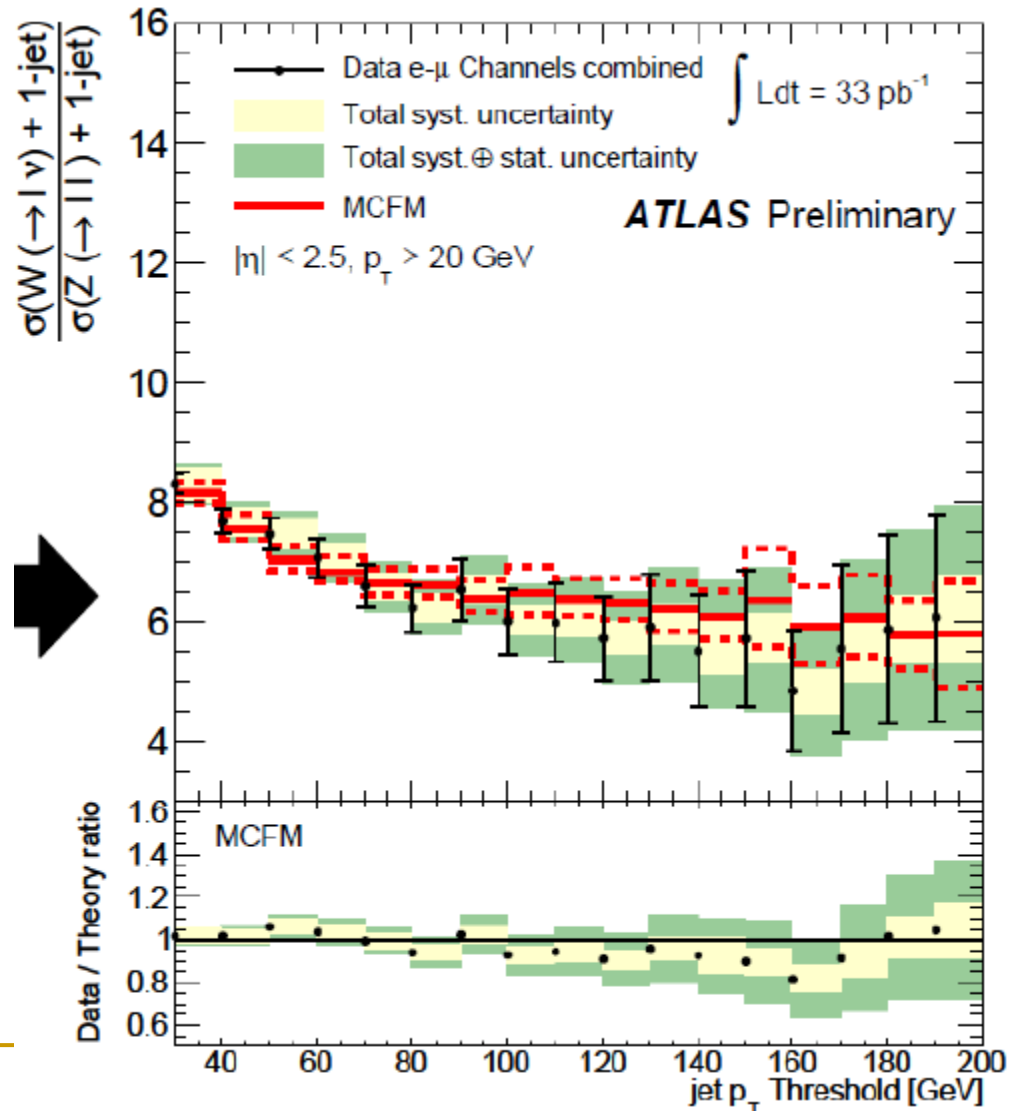
- ▶ **JES: 8(26)%** for $N_j \geq 1$ (4)
- ▶ jets from pile-up $\approx 7\%$
- ▶ lep. reco. $\approx 2\%$
- ▶ QCD bkgd $\approx 2\%$
- ▶ unfolding $\approx 2\%$



Rjets = ratio W+1jet/Z+1jet

- This is the first time this ratio is measured
 - Sensitive to new physics
 - Very small sensitivity to PDG
 - CTEQ6.6: 0.5%
 - MSTW2008: 0.3%

$$R_{\text{jets}}(p_T > x) = \frac{\sigma_{W+1\text{-jet}}(p_T > x)}{\sigma_{Z+1\text{-jet}}(p_T > x)}$$

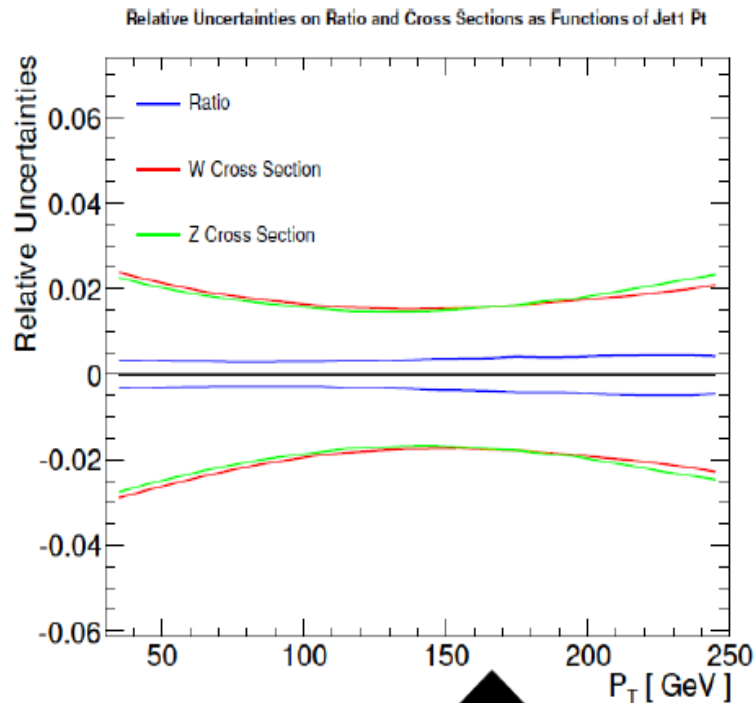


PDF uncertainty on Rjets

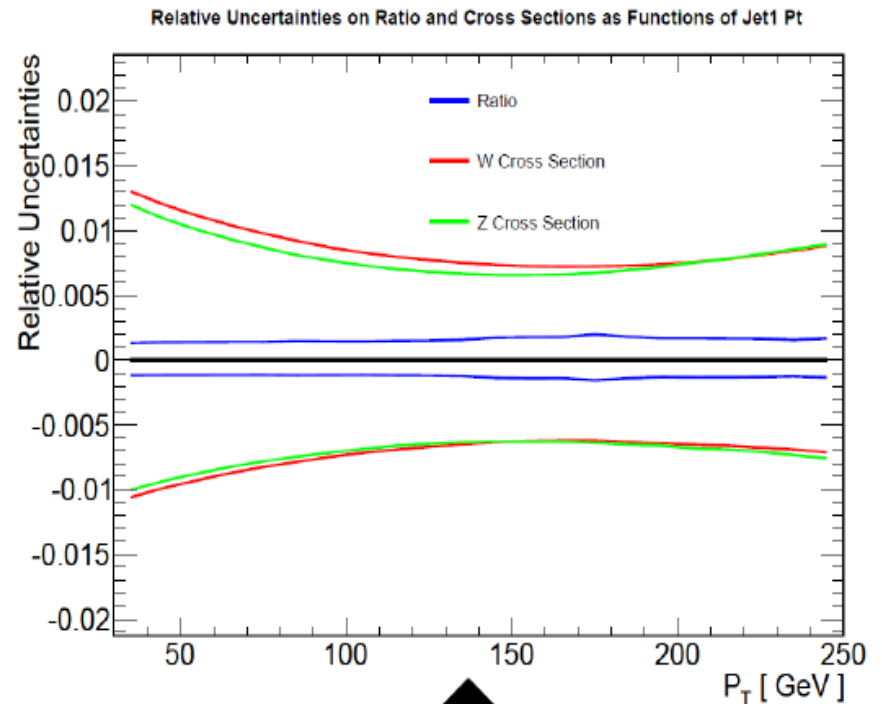
- Very small uncertainty on PDF's

CTEQ6.6: 0.5%,

MSTW2008: 0.3%



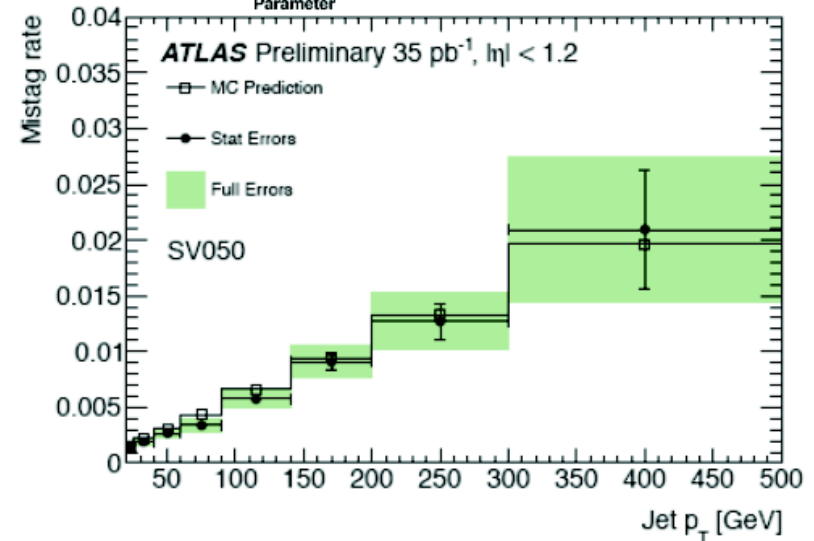
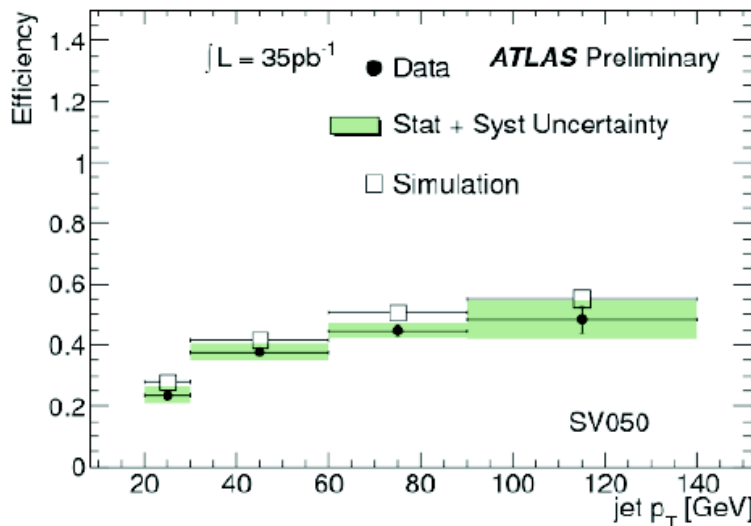
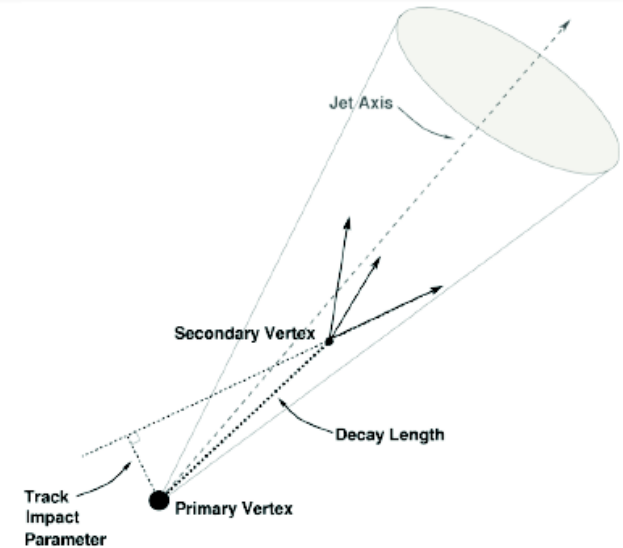
CTEQ6.6



MSTW2008

W/Z + b-jets: b-tagging

- b-jets are selected exploiting the long lifetime (1.5 ps) and the large mass of B-hadrons
- The SV0 b-tagging algorithm is based on requiring a displaced secondary vertex reconstructed within a jet with a decay length significance > 5.85
- The b-tagging efficiency and its systematics is estimated by studying semi-leptonic B decays in QCD multi-jet events, and top events



W/Z+bjets: backgrounds

The b-tag changes the composition of backgrounds with respect to W/Z+jet measurements

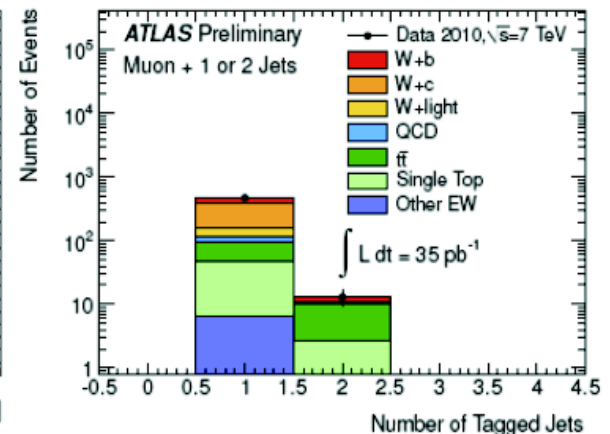
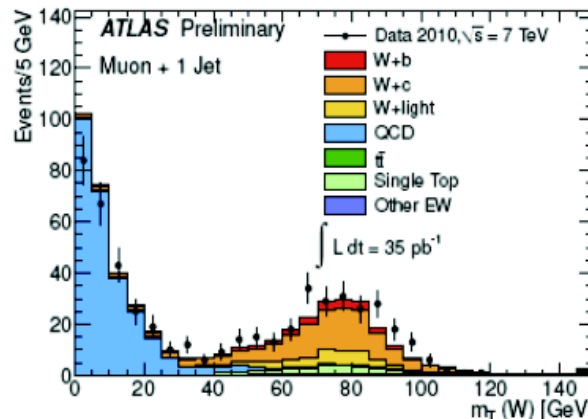
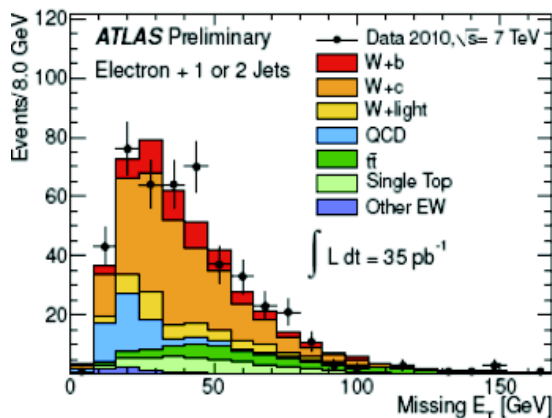
Z+b-jet signal / background > 10
 W+b-jet signal / background ≈ 0.5

non-W background

- ▶ comes from QCD multi-jet heavy flavour production.
- ▶ Tighten lep. ID to keep bkgd < 30% with 50% uncertainty

top background

- ▶ t-tbar largest background (partially irreducible)
- ▶ single-top estimated with MC (partially irreducible)
- ▶ top yield measured in the ≥ 4 jet bin and extrapolated to 1, 2 jet wit MC
- ▶ uncertainty $\approx 30\%$ (JES)
- ▶ b-tag uncertainty reduced to 2%



W/Z+bjets: extraction of b-jets fraction

A maximum likelihood fit to the SV0 mass distribution is used to separate b-jets from c- and light-jets, and extract the flavour fraction on a statistical basis.

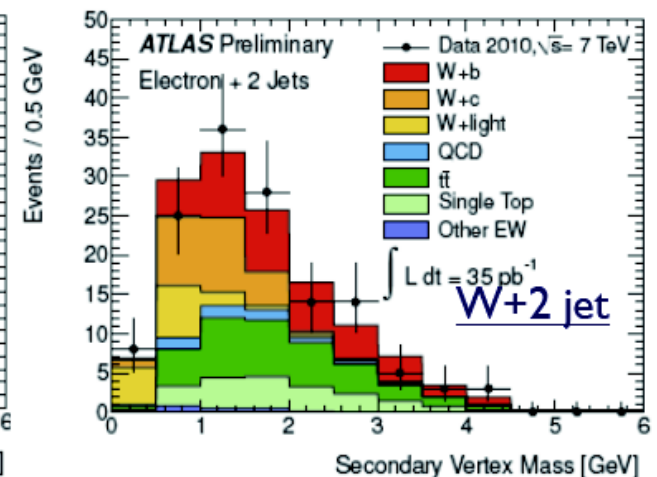
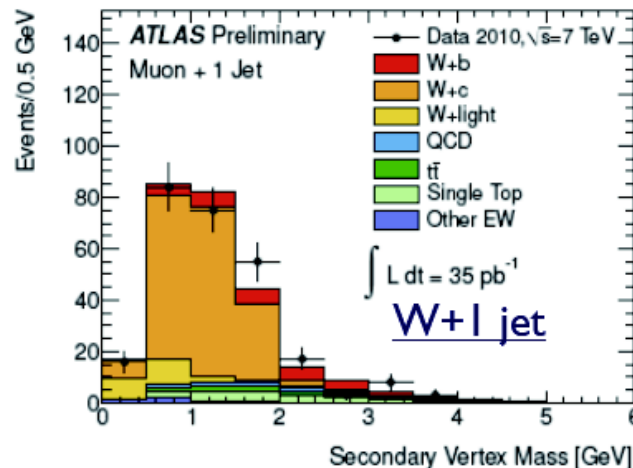
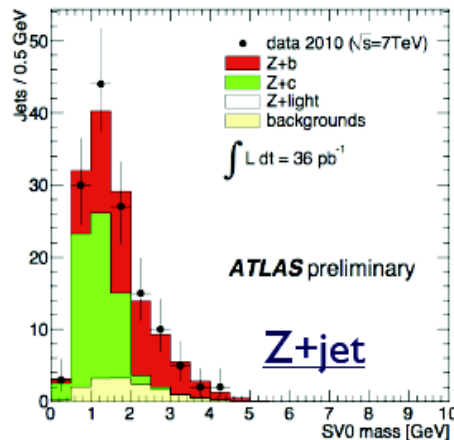
- SV0 mass template are modeled with MC
- template systematics: data vs. MC in multi-jet events enriched in light-, c-, and b-jets.

Z+bjet

- ▶ Fit the combined e and μ samples and each b-tagged jet in the event
- ▶ At least 1 b-tagged jet

W+bjet

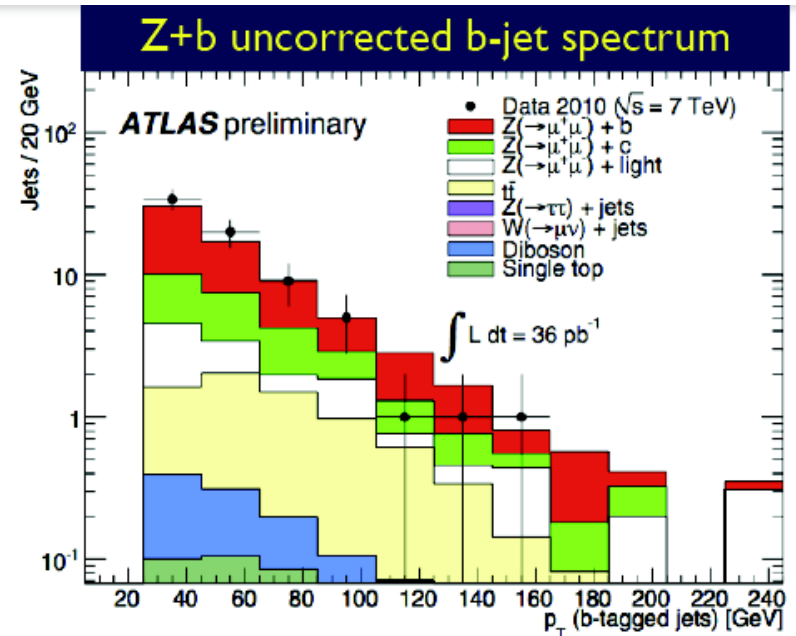
- ▶ 1 b-tagged jet
- ▶ about 10% of events have 2 b
- ▶ 1 or 2 jet
- ▶ fit each jet bin separately for e and μ



Z+b-jets: results

$$\sigma = \frac{N_b}{C_e \times \mathcal{L}_e + C_\mu \times \mathcal{L}_\mu}$$

- Inclusive b-jet production cross section in association with a Z boson
- Jet fitted yield is corrected for all detector effects with MC LO matched prediction for Zjet (including heavy flavour) from ALPGEN and SHERPA
- **uncertainty: $\approx 20\%$ stat. and $\approx 23\%$ syst.**
- dominant systematics:
 - b-tagging & SV mass template $\approx 10\%$
 - Z+b-jet modeling $\approx 10\%$
 - Jet + bjet energy scale $\approx 4\%$
- **MCFM in good agreement with data within uncertainty**



Experiment $3.55^{+0.82}_{-0.74}(\text{stat})^{+0.73}_{-0.55}(\text{syst}) \pm 0.12(\text{lumi}) \text{ pb}$

MCFM $3.40 \pm 0.44 \text{ pb}$

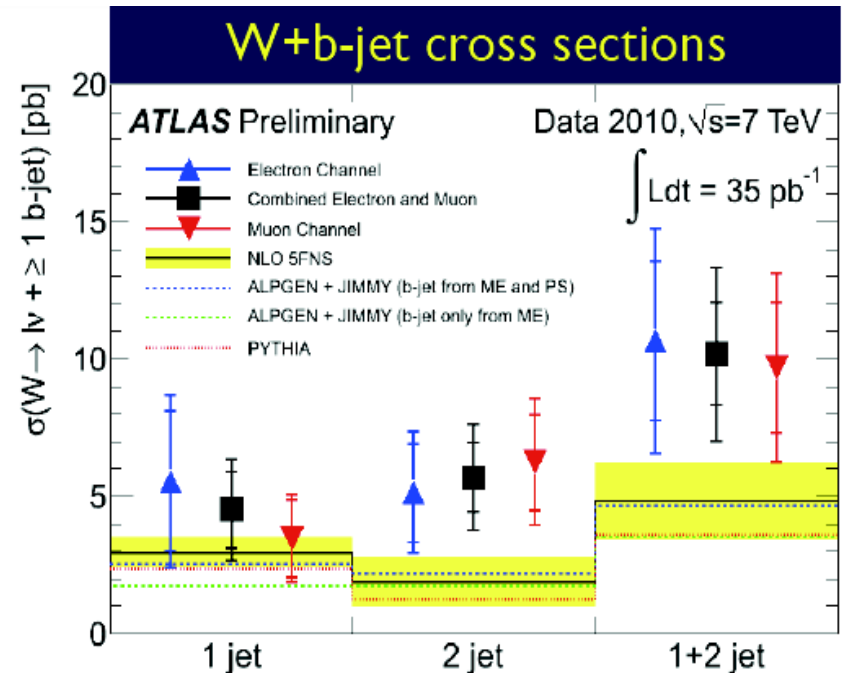
ALPGEN $2.23 \pm 0.01(\text{stat only}) \text{ pb}$

SHERPA $3.33 \pm 0.04(\text{stat only}) \text{ pb}$

W+b-jets: results

$$\sigma_{W+b\text{-jet}} \times \mathcal{B}(W \rightarrow \ell\nu) = \frac{n^{\text{tag}} \cdot f_{W+b\text{-jet}}}{\int Ldt \cdot \mathcal{U}}$$

- W+b-jet cross section (event level)
- First measurement in exclusive jet bins
- event fitted yield is corrected for all detector effects with MC LO matched prediction for Wjet (including heavy flavour) from ALPGEN
- **uncertainty: $\approx 20\%$ stat. and $\approx 25\%$ syst.**
- dominant systematics:
 - b-tagging & SV mass template $\approx 16\%$
 - top background $\approx 12\%$
 - QCD background $\approx 7\%$
 - W+b-jet modeling $\approx 10\%$
 - Jet + bjet energy scale $\approx 7\%$



- NLO prediction obtained in the 5 flavour number scheme [F. Caola et al. arXiv:1107.3714]
- **NLO agrees within 1.5σ with the measurements**

	σ_{vis} [pb]			
1 jet	$2.9^{+0.40}_{-0.36}$ (scale)	$^{+0.18}_{-0.02}$ (PDF)	$^{+0.19}_{-0.10}$ (m_b)	± 0.20 (non-pert)
2 jet	$1.9^{+0.81}_{-0.37}$ (scale)	$^{+0.14}_{-0.02}$ (PDF)	$^{+0.06}_{-0.05}$ (m_b)	± 0.13 (non-pert.)
1+2 jet	$4.8^{+1.20}_{-0.73}$ (scale)	$^{+0.32}_{-0.03}$ (PDF)	$^{+0.25}_{-0.15}$ (m_b)	± 0.34 (non-pert.)

Summary

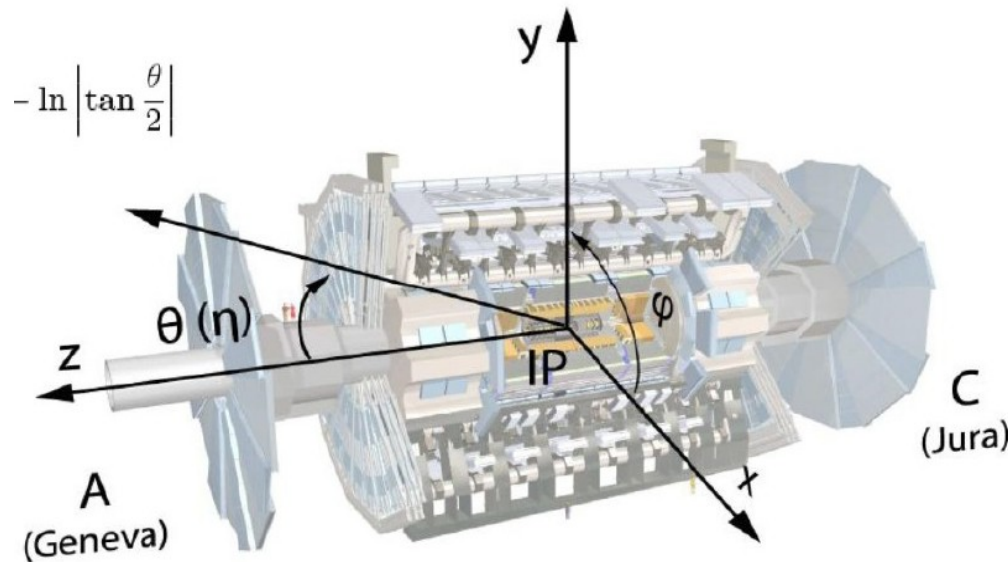
- The LHC era allowed us to verify QCD in new kinematic regimes, good testing ground for predictions
- Current understanding of detectors allows to do precision measurements in W/Z sector
- Extensive set of measurements also in W/Z+jets differential cross-sections, also with b-tagging
- Overall impressive agreement with MC predictions

Plan

- Amazing how much could have been done with only 36pb^{-1} data accumulated in 2010: numbers of results are still in the pile-line but already theory is being tested quantitatively.... and is holding its own (unfortunately)
 - 23.11** Top physics
 - 30.11** Diboson production and TGS couplings
 - 7.12** B-physics and heavy ions
 - 4.01** Higgs boson... where we are?
 - 18.01** What's new from New Physics searches?

ATLAS Detector

THE ATLAS DETECTOR IS
REALLY BIG!

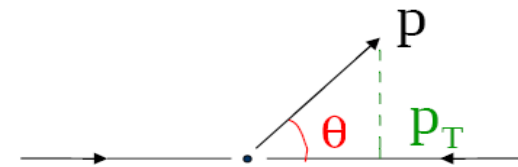


- Length : ~ 46 m
- Radius : ~ 12 m
- Weight : ~ 7000 tons
- $\sim 10^8$ electronic channels
- 3000 km of cables

Transverse momentum

(in the plane perpendicular to the beam)

$$p_T = p \sin\theta$$



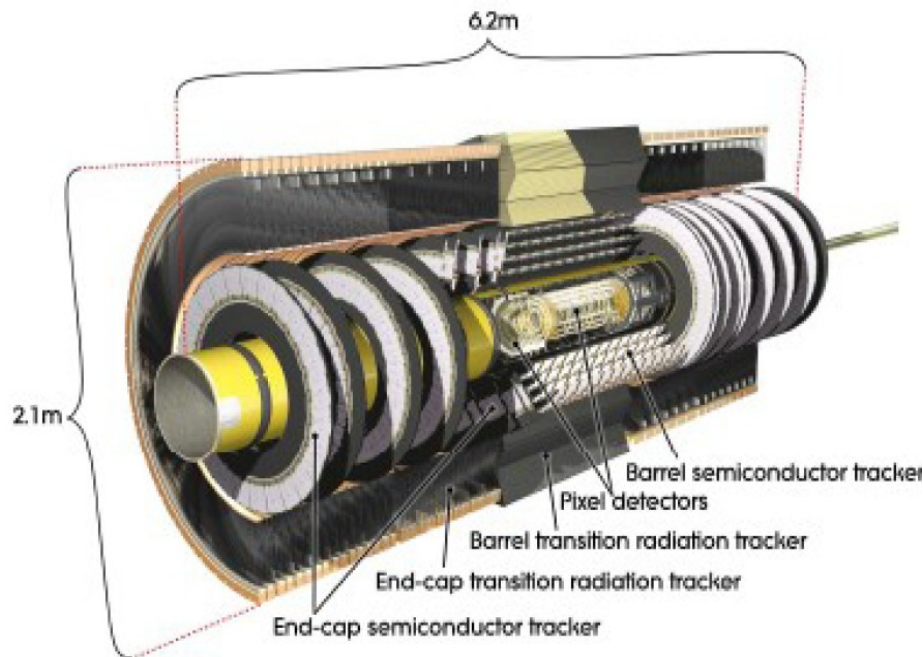
Rapidity: $\eta = -\log(\operatorname{tg} \frac{\theta}{2})$

$$\theta = 90^\circ \rightarrow \eta = 0$$

$$\theta = 10^\circ \rightarrow \eta \cong 2.4$$

$$\theta = 170^\circ \rightarrow \eta \cong -2.4$$

ATLAS Inner Detector



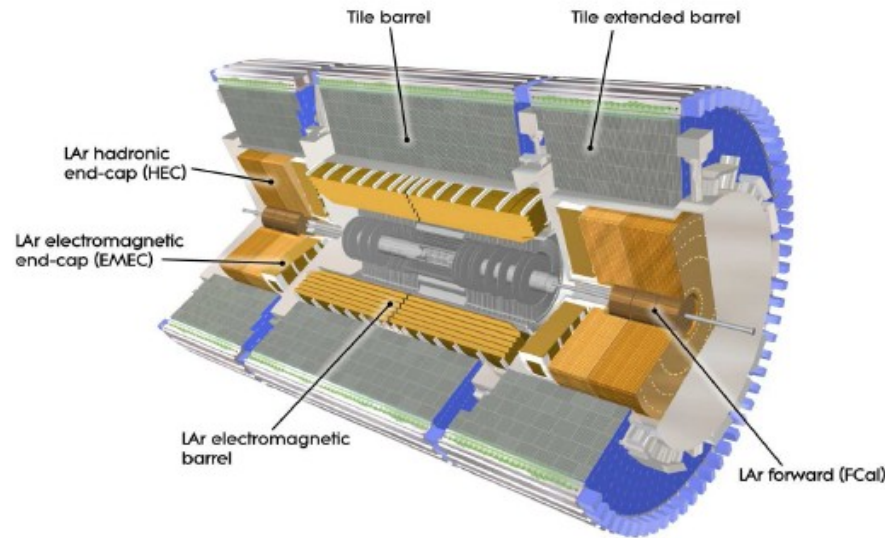
The inner detector $|\eta| < 2.5$ consists of

- Pixel detectors, semi-conductor tracker (SCT), transition radiation tracker

- ≈ 87 million readout channels
- Immersed in 2T solenoidal magnetic field

- Resolution of $\sigma/p_T = 5 \times 10^{-4} \oplus 0.015$

ATLAS Calorimeters



Electromagnetic and hadronic calorimeters

- Subsystem technology and granularity \leftrightarrow shower characteristics
- Transverse and longitudinal sampling \approx 200000 readout cells up to $|\eta| < 4.9$

Electromagnetic Calorimeters:

- Fine granularity
 $\Delta\eta \times \Delta\phi = 0.025 \times 0.025$ in central region
- Energy resolution $10\%/\sqrt{E}$

Hadronic Calorimeters:

- Granularity
 $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$ in central region, less segmented in forward region
- Energy resolution $50\%/\sqrt{E} \oplus 0.03$