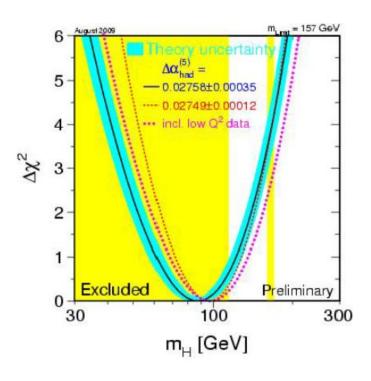
Statistics and Data Analysis (HEP)

The Dark Arts:

- Bayesian inference
- Credible intervals
- Systematics uncertainties





Follow the course/slides from M. A. Thomson lectures at Cambridge University

Lecture 1: The basics

Introduction, Probability distribution functions, Binomial distributions, Poisson distribution

Lecture 2: Treatment of Gaussian Errors

The central limit theorem, Gaussian errors, Error propagation, Combination of measurements, Multi-dimensional Gaussian errors, Error Matrix

Lecture 3: Fitting and Hypothesis Testing

The χ^2 test, Likelihood functions, Fitting, Binned maximum likelihood, Unbinned maximum likelihood

Lecture 4: The Dark Arts

Bayesian Inference, Credible Intervals

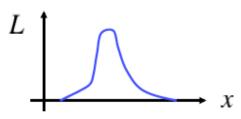
The Frequentist approach, Confidence Intervals

Systematic Uncertainties

Parameter Estimation Revisited

- ★ Let's consider more carefully the maximum likelihood method for simplicity consider a single parameter X
- ★ Construct the likelihood that our data are consistent with the model, i.e. the probability that the model would give the observed data

$$L = P(\text{data}; x)$$



- ★ We have then (very reasonably) taken the value of x which maximises the likelihood as our best estimate of the parameter
- ★ With less justification we then took our error estimate from

$$-\ln L \to -\ln L + \frac{1}{2}$$

- ★ Does this really make sense ?
- ***** What we really want to calculate is the posterior PDF for the parameter given the data, i.e. P(x; data)

"assumed"
$$P(x; data) = P(data; x)$$

Can not justify this – in general it is not the case

Conditional Probabilities and Bayes' Theory

- **★** A nice example of conditional probability (from L. Lyons)
 - In the general population, the probability of a randomly selected woman being pregnant is 2%

$$P(\text{pregnant; woman}) = 0.02$$

But

- $P(\text{woman; pregnant}) \gg 0.02$
- ★ Correct treatment of conditional probabilities requires Bayes' theorem
 - Probability of A and B can be expressed in terms of conditional probabilities

$$P(AB) = P(A;B)P(B) = P(B;A)P(A)$$

$$P(A;B) = \frac{P(B;A)P(A)}{P(B)}$$

★ Here the prior probability of selecting a woman is

$$P(\text{woman}) = 0.50$$

i.e. half population are women

and the prior probability of selecting a pregnant person is

$$P(pregnant) = 0.01$$

i.e. 1 % of population are pregnant

$$P(\text{woman; pregnant}) = \frac{P(\text{pregnant; woman})P(\text{woman})}{P(\text{pregnant})} = \frac{0.02 \times 0.5}{0.01} = 1$$
Sanity restored...

- ★ Apply Bayes' theory to our the measurement of a parameter x
 - We determine $P(\mathrm{data};x)$, i.e. the likelihood function
 - We want P(x; data), i.e. the PDF for x in the light of the data
 - Bayes' theory gives:

$$P(x; data) = \frac{P(data; x)P(x)}{P(data)}$$

 $P(\mathrm{data};x) \quad \text{the likelihood function, i.e. what we measure} \\ P(x;\mathrm{data}) \quad \text{the posterior PDF for } x\text{, i.e. in the light of the data} \\ P(\mathrm{data}) \quad \left\{ \begin{array}{l} \text{prior probability of the data. Since this doesn't depend on} \\ x \text{ it is essentially a normalisation constant} \end{array} \right. \\ P(x) \quad \left\{ \begin{array}{l} \text{prior probability of } x\text{, i.e. encompassing our knowledge of} \\ x \text{ before the measurement} \end{array} \right.$

★ Bayes' theory tells us how to modify our knowledge of x in the light of new data

Bayes' theory is the formal basis of Statistical Inference

Applying Bayes' Theorem

★ Bayes' theory provides an unambiguous prescription for going from

$$P(\text{data};x) \rightarrow P(x;\text{data})$$

- **\star** But you need to provide the PRIOR PROBABILITY P(x)
- ★ This is fine if you have an objective prior, e.g. a previous measurement

$$P(x) = \frac{1}{\sqrt{2\pi}\sigma_1} \exp\left\{-\frac{(x-x_1)^2}{2\sigma_1^2}\right\}$$

• If we now make a new measurement, i.e. determine the likelihood function

$$P(\text{data};x) = \frac{1}{\sqrt{2\pi}\sigma_2} \exp\left\{-\frac{(x-x_2)^2}{2\sigma_2^2}\right\}$$

Bayes' theory then gives

$$P(x; \text{data}) = \frac{P(\text{data}; x)P(x)}{P(\text{data})} = \frac{1}{P(\text{data})} \frac{1}{2\pi\sigma_1\sigma_2} \exp\left\{-\frac{(x - x_1)^2}{2\sigma_1^2} - \frac{(x - x_2)^2}{2\sigma_2^2}\right\}$$

$$P(x; \mathrm{data}) = \frac{1}{P(\mathrm{data})} \frac{1}{2\pi\sigma_1\sigma_2} \exp\left\{-\frac{(x-\overline{x})^2}{2\sigma^2}\right\} \quad \begin{array}{l} \text{Where \overline{x} and σ are the usual mean and variance for combining two measurements} \end{array}$$

For this to be a (normalised) PDF can infer (although it isn't of any interest):

$$P(\text{data}) = [2\pi(\sigma_1^2 + \sigma_2^2)]^{-\frac{1}{2}}$$

The Problem with Applying Bayes' Theorem

- ★ The problem arises when there is no objective prior
- ★ For example, in a hypothetical background free search for a Z', observe no events
 - No problem in calculating the likelihood function (a conditional probability)

$$P(\text{data};x) = P(0;x) = e^{-x}$$
 Poisson prob. for observing 0

x is the true number of expected events

- What is the best estimate of x and the 90 % "confidence level upper limit" ?
- Depends on the choice of prior probability:

$$P(x; data) = P(x)e^{-x}$$

- What to do about the prior ?
- i.e. how do we express our knowledge (none) of x prior to the measurement
- **★** In general there is no objective answer, always putting in some extra information
 - i.e. a subjective bias
 - could argue that a flat prior, i.e. P(x) = constant, is objective
 - but why not choose a prior that is flat in ln x?
 - for some limits/measurements (e.g. a mass) a flat prior in In x is more natural
 - the arbitrariness in the choice of prior is a problem for the Bayesian approach
 - it can make a big difference...

Choice of Prior, example I

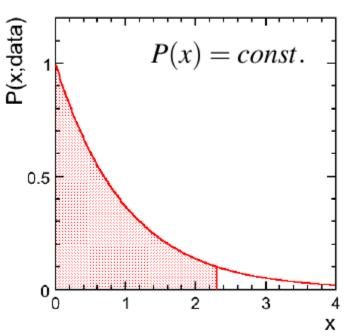
★ See no events...

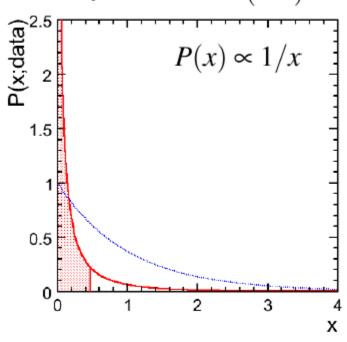
$$P(\text{data}; x) = P(0; x) = e^{-x}$$

Poisson prob. for observing 0

Prior flat prior in x : P(x) = const.

Prior flat prior in $\ln x : P(\ln x) = const$.





★ The Conclusions are very different. Compare regions containing 90 % of probability

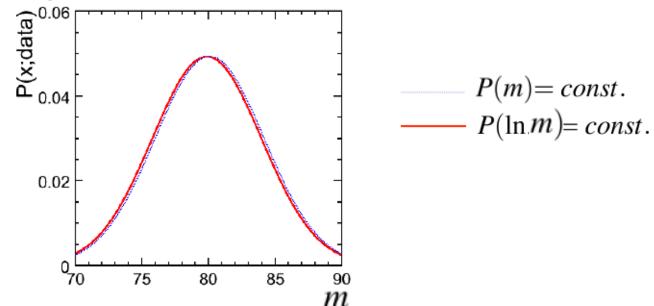
In this case, the choice of prior is important

Choice of Prior, example II

***** Suppose we measure the W-boson mass: $80.1 \pm 4.1\,\mathrm{GeV}$

$$P(\mathrm{data};m) = G(80.1;m) \propto \exp\left\{-\frac{(80.1-m)^2}{2\times4.1^2}\right\}$$
 ** We want $P(m;\mathrm{data}) = P(m)P(\mathrm{data};m)$

- - Again consider two priors



- Here the choice of prior is NOT important
- The data are "strong enough" to overcome our prior assumptions (subjective bias)
- Here, can interpret the measurement as a Gaussian PDF for m

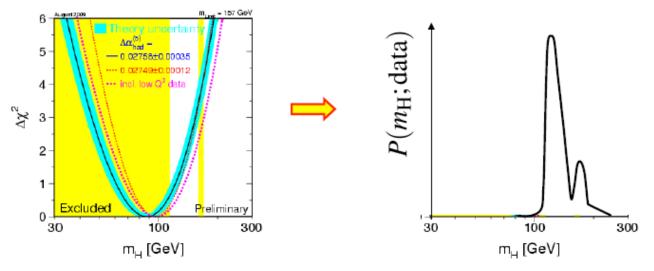
- ★ So what do we learn from this ?
 - Whilst we know how to apply Bayesian statistical inference, we have insufficient data, i.e. we don't know the prior
 - Unless the data are "strong", i.e. override the information in the reasonable range of prior probabilities, we cannot expect to know

 Applies equally to our experiment where we saw zero events and wanted to arrive at a PDF for the expected mean number of events...

Don't have enough information to answer this question

Bayesian Credible Intervals

★ Ideally, (I) would like to work with probabilities, i.e. a PDF which encompasses all our knowledge of a particular parameter, e.g. $P(m_H; data)$



- ★ Could then integrate PDF to contain 95 % of probability. Can then define the "95 % Credible Interval*: m_H < 186 GeV"</p>
- ★ To do this need to go from $P(\text{data}; m_{\text{H}})$, i.e. from $\Delta \ln \mathcal{L}$, to $P(m_{\text{H}}; \text{data})$ requires "subjective" choice of prior probability
- ★ Hence Bayesian Credible Intervals necessarily include some additional input beyond the data alone...

*This is not what is done.

Bayesian Credible Intervals - example

- **★** Trying to estimate a selection efficiency using MC events. All N events pass cuts.
 - what statement can we make about the efficiency?
- **★** Binomial distribution...

$$P(\text{data};x) \longrightarrow P(N;\varepsilon) = {}^{N}C_{N}\varepsilon^{N}(1-\varepsilon)^{0} = \varepsilon^{N}$$

★ Apply Bayes' theorem:

$$P(x; \mathrm{data}) \longrightarrow P(\varepsilon; N) = \frac{P(N; \varepsilon) P(\varepsilon)}{P(N)} \longleftarrow \begin{array}{c} \mathrm{Prior} \\ \mathrm{Constant} \end{array}$$

★ Choose prior, e.g. $P(\varepsilon) = 1$

$$P(\varepsilon;N) = \kappa \varepsilon^N$$

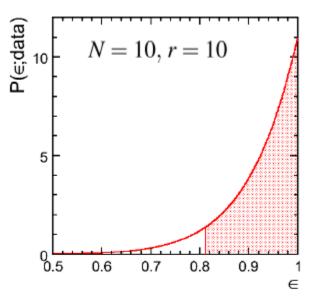
* Normalise $\int_0^1 P(\varepsilon;N)d\varepsilon=1 \implies \kappa=(N+1)$ $P(\varepsilon;N)=(N+1)\varepsilon^N$

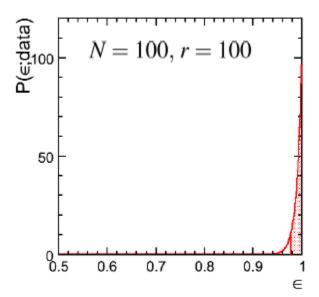
$$P(\varepsilon;N) = (N+1)\varepsilon^N$$

$$P(\varepsilon;N) = (N+1)\varepsilon^N$$

to find region containing 90% of probability

$$\Rightarrow \quad \varepsilon_{90\%} = (1 - 0.90)^{\frac{1}{N+1}}$$





90 % Credible Interval:

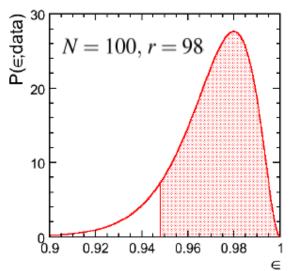
$$\varepsilon > 0.81$$

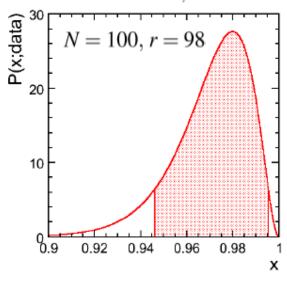
$$\varepsilon > 0.977$$

(with a flat prior probability)

Likelihood Ordering

- ★ Note, 90 % credible interval is not uniquely defined
 - more than one interval contains 90 % probability, e.g. N = 100, r = 98





90 % Credible Interval: $\varepsilon > 0.9482$

 $0.9456 < \varepsilon < 0.9956$

- ★ Natural, to choose the interval such that all points in the excluded region are lower in likelihood than those in the credible interval : likelihood ordering
- ★ Credible intervals provide an intuitive way of interpreting data, but:
 - Rarely used in Particle Physics as a way of presenting data
 - Because they represent the "data" and "prior" combined
 - NOTE: all information from the experiment is in the likelihood $P(\mathrm{data};x)$

C.I. vs C. L.

- ***** From data obtain P(data;x)
- ★ Bayes' theorem provides the mathematical framework for statistical inference
- \star To go from $P({
 m data};x) o P(x;{
 m data})$ requires a (usually) subjective choice of Prior probability
- ★ For "weak" data, the choice of Prior can drive the interpretation of the data
- ★ Credible intervals are a useful way of interpreting data, but are generally not used in Particle Physics as a way of presenting the conclusions of an experiment.
- ★ Particle Physics to use Frequentist "Confidence limits" which are not P(x;data) [and do not form a mathematically consistent basis for statistical inference]
- ★ Finally, never forget that credible intervals (or confidence limits) are an interpretation of the data

The experimental result is the likelihood function P(data;x)

A Few words on Systematics Uncertainties

- **★ Systematic Uncertainties are often associated with an internal unknown bias, e.g.**
 - How well do you know your calibration
 - How well does MC model the data, e.g. jet fragmentation parameters
- **★ Parametric Uncertainties associated with uncertain parameters**
 - How does the uncertainty on the Higgs mass impact the interpretation of a a measurement
- **★ No over-riding principle just some general guidelines**
 - Once a result is published, systematic errors will be treated as if they are Gaussian

$$x = a \pm b$$
 (stat.) $\pm c$ (syst.)

- Some systematic errors are Gaussian: e.g. energy scale determined from data e.g. $Z \rightarrow e^+e^-$ to determine electron energy scale
- Others are not: e.g. impact of different jet hadronisation models, where one might compare PYTHIA with HERWIG – here one obtains a single estimate of the scale of the uncertainties
- Theoretical uncertainties: e.g. missing HO corrections. Again these are estimates – should not be treated as Gaussian (although they are)
- **★ Systematic dominated measurements**
 - Beware if there is a single dominating systematic error and it is inherently non-Gaussian, this is a problem

Estimating Systematics Uncertainties

★ No rules – just guidelines

- Remember syst. errors will be treated as Gaussian, so try to evaluate them on this basis, e.g. suppose use 3 alternative MC jet fragmentation models and result changes by $+\Delta_1$, $+\Delta_2$ and $-\Delta_3$ (where Δ_2 is the largest):
 - i) take largest shift as systematic error estimate: Δ_2 ?
 - ii) assume error distributed uniformly in "box" of width $2\Delta_2$ giving an rms of $2\Delta_2 /\sqrt{12}$?
- Cut variation is evil (i.e. vary cuts and see how results change)
 - at best, introduces statistical noise
 - at worst, hides away lack of understanding of some data MC discrepancy
 understand the origin of the discrepancy
- Wherever possible use data driven estimates, energy scales, control samples, etc.
- Remember that you are estimating the scale of a possible systematic bias

Incorporating Systematics into Fits

- ★ Two commonly used approaches
 - Error matrix with (correlated) systematic uncertainties
 - Nuisance parameters
- ★ Nuisance parameter example:
 - Suppose we are looking at WW decays and count numbers of events in three different decay channels qqqq, qqlv and lvlv
 - Want to measure cross section and hadronic branching fractions accounting for common luminosity uncertainty
 - i) build physics model

$$N_{\text{qqqq}}^{\text{exp}}(\sigma_{\text{WW}}, B_{\text{qq}}, \mathcal{L}) = \sigma_{\text{WW}} B_{\text{qq}}^2 \, \varepsilon \mathcal{L}$$

ii) build likelihood function

$$\chi^{2}(\sigma_{\text{WW}}, B_{\text{qq}}, \mathcal{L}) = -2 \ln L = \frac{(N_{\text{qqqq}}^{\text{exp}} - N_{\text{qqqq}}^{\text{obs}})^{2}}{N_{\text{qqqq}}^{\text{exp}}} + \frac{(N_{\text{qqlv}}^{\text{exp}} - N_{\text{qqlv}}^{\text{obs}})^{2}}{N_{\text{qqlv}}^{\text{exp}}} + \frac{(N_{\text{lvlv}}^{\text{exp}} - N_{\text{lvlv}}^{\text{obs}})^{2}}{N_{\text{lvlv}}^{\text{exp}}}$$

iii) add penalty term for nuisance parameters, here integrated lumi. Known to be L_0 with uncertainty σ_L

$$\chi^{2}(\sigma_{WW}, B_{qq}, \mathcal{L}) = -2 \ln L = \frac{(N_{qqqq}^{exp} - N_{qqqq}^{obs})^{2}}{N_{qqqq}^{exp}} + \frac{(N_{qqlv}^{exp} - N_{qqlv}^{obs})^{2}}{N_{qqlv}^{exp}} + \frac{(N_{lvlv}^{exp} - N_{lvlv}^{obs})^{2}}{N_{lvlv}^{exp}} + \frac{(\mathcal{L} - \mathcal{L}_{0})^{2}}{\sigma_{\mathcal{L}}^{2}}$$

Incorporating Systematics into Fits

★ Let's consider this more closely

$$\chi^{2}(\sigma_{\text{WW}}, B_{\text{qq}}, \mathcal{L}) = -2 \ln L = \frac{(N_{\text{qqqq}}^{\text{exp}} - N_{\text{qqqq}}^{\text{obs}})^{2}}{N_{\text{qqqq}}^{\text{exp}}} + \dots \frac{(N_{\text{qqlv}}^{\text{exp}} - N_{\text{qqlv}}^{\text{obs}})^{2}}{N_{\text{qqlv}}^{\text{exp}}} + \frac{(\mathcal{L} - \mathcal{L}_{0})^{2}}{\sigma_{\mathcal{L}}^{2}}$$

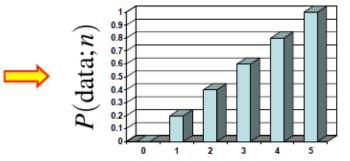
- We are now fitting 3 parameters
 - the number of degrees of freedom has not changed, since we have added one parameter, but also one additional "data point"
- Of the 3 parameters, we are "not interested" in the fitted value of the lumi.
- The penalty term constrains the luminosity to be consistent with the externally measured value
- The presence of the nuisance parameters will flatten the fitted likelihood surface – increasing the uncertainties on the fitted parameters
- Also have some measure of the tension in the fit
 - if the data pull the nuisance parameter away from the expected value, could indicate a problem

Appendix: Choice of Prior, example III

- An example (apparently due to Newton), e.g. see CERN Yellow Report 2000-005
- ★ Suppose you are in the Tower of London facing execution.
- ★ The Queen arrives carrying a small bag and says "This bag contains 5 balls; the balls are either white or black. If you correctly guess the number of black balls, I will spare your life and set you free."
- ★ The Queen is in a good mood and continues "To give you a better chance, you can take one of the balls from the bag."

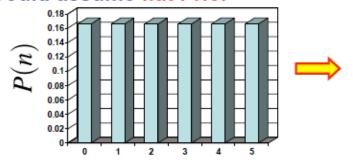
It's BLACK

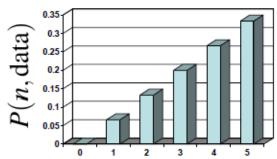
- ★ The Queen points her pistol at you "Time to choose, sucker..."
- ★ What do you guess to maximise your chance of survival ?
- ★ Use statistical inference to analyse the problem.
 - Let n be the number of black balls in the bag.
 - The data are "that you picked out a black ball"
 - Can calculate P(data;n) e.g. if there were two black balls chance of picking out a black ball from the five in the bag was 2/5.



- **★** But we want $P(n; \text{data}) \propto P(n)P(\text{data}; n)$
- ★ Answer depends on choice of Prior

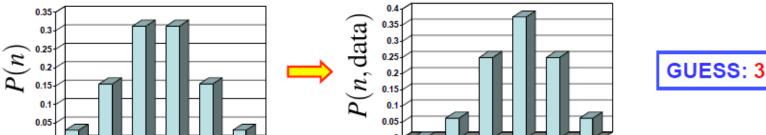
★ Could assume flat Prior





GUESS: 5

★ Could assume balls drawn randomly from a large bag containing equal nos. B & W



★ Oh dear... answer depends on Prior (unknown) assumptions