

Elementary Particle Physics: theory and experiments

Discovery of the Higgs boson

- Measurements: mass, spin, couplings

Searches for New Physics

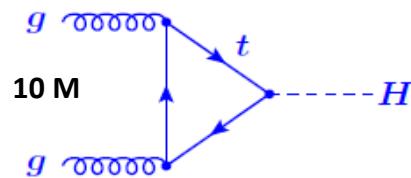
- Supersymmetry
- Exotic models
- Dark Matter
- Unconventional signatures

Discovery of the Higgs boson

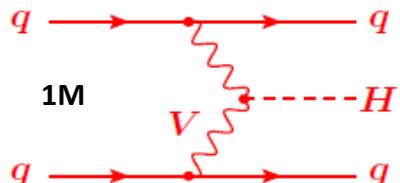
- Measurements: mass, spin, couplings

The Higgs Boson at the LHC

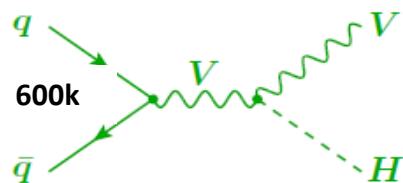
Production



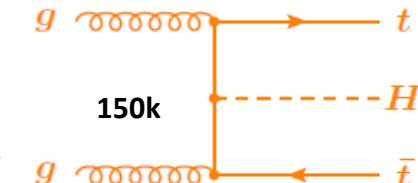
Main production channel



2 forward jets,
little central
hadronic activity



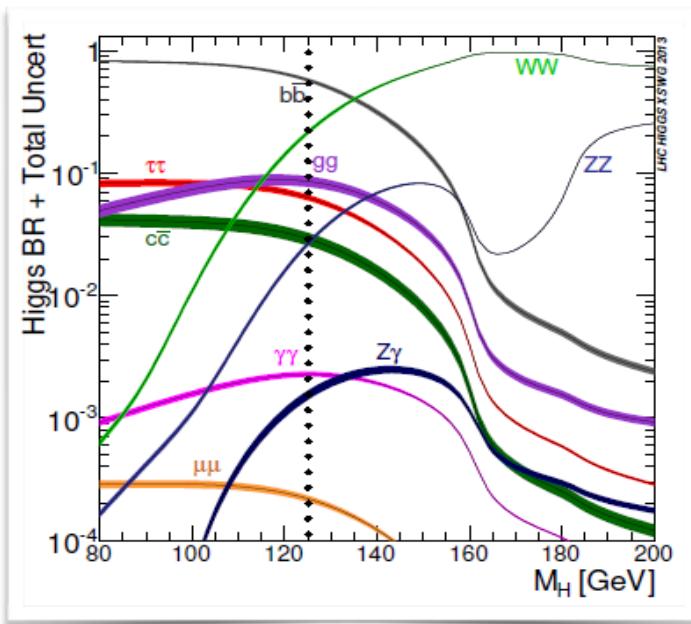
Tag W and Z
decays



Tag 2 top quarks

#Higgs produced at
13 TeV (2015-2017)

Decays



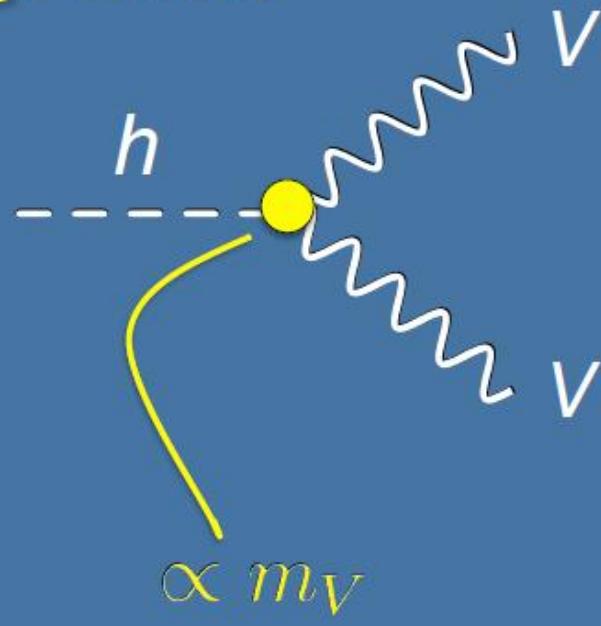
5 main channels at the LHC

Decay branching fractions for
 $m_H = 125$ GeV

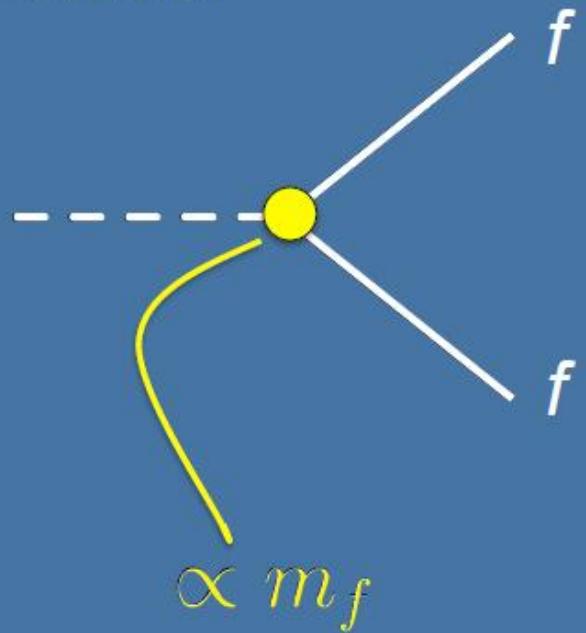
- $H \rightarrow b\bar{b}$: 58 %
- $H \rightarrow WW^*$: 21%
- $H \rightarrow \tau^+\tau^-$: 6.3%
- $H \rightarrow ZZ^*$: 2.6%
- $H \rightarrow \gamma\gamma$: 0.2%

Higgs boson couplings

Gauge bosons

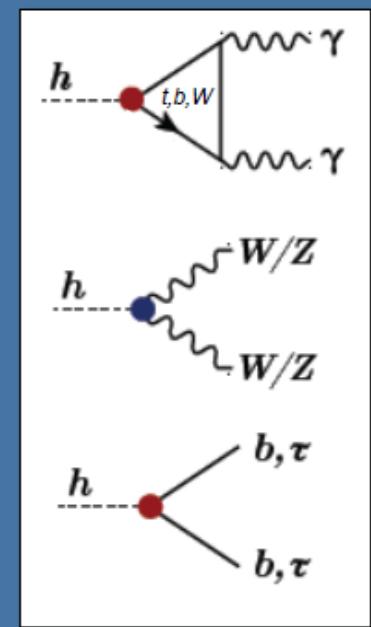
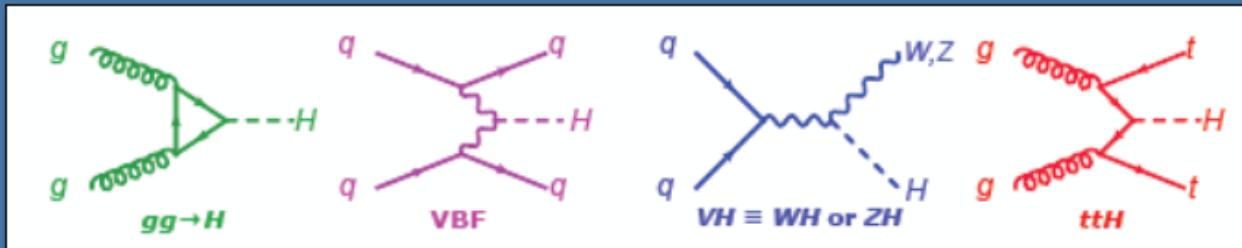


Fermions



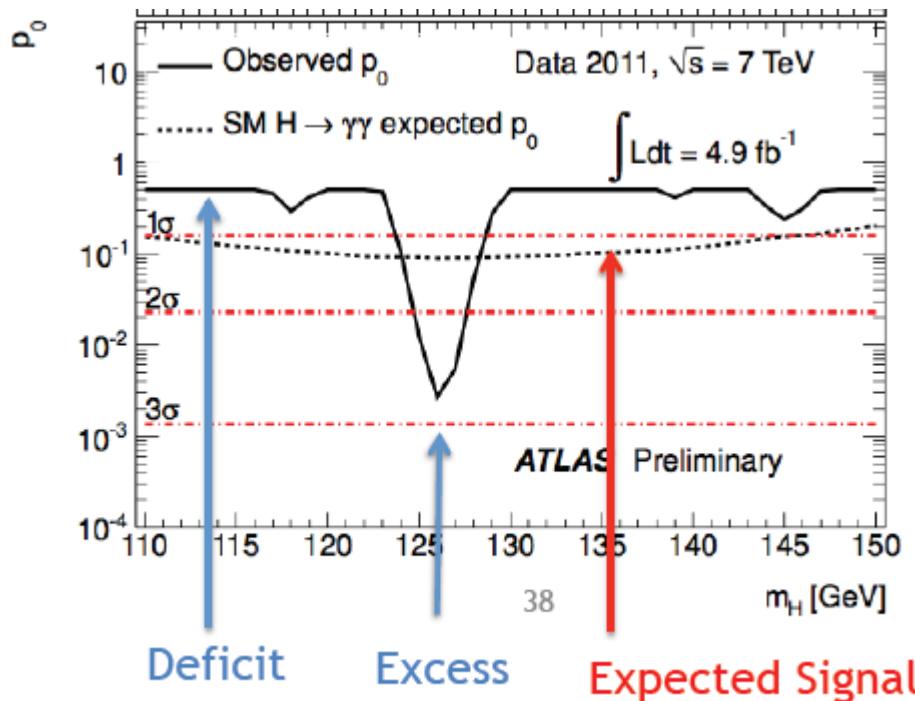
Higgs boson properties are fixed in the Standard Model (m_h)

Higgs boson phenomenology

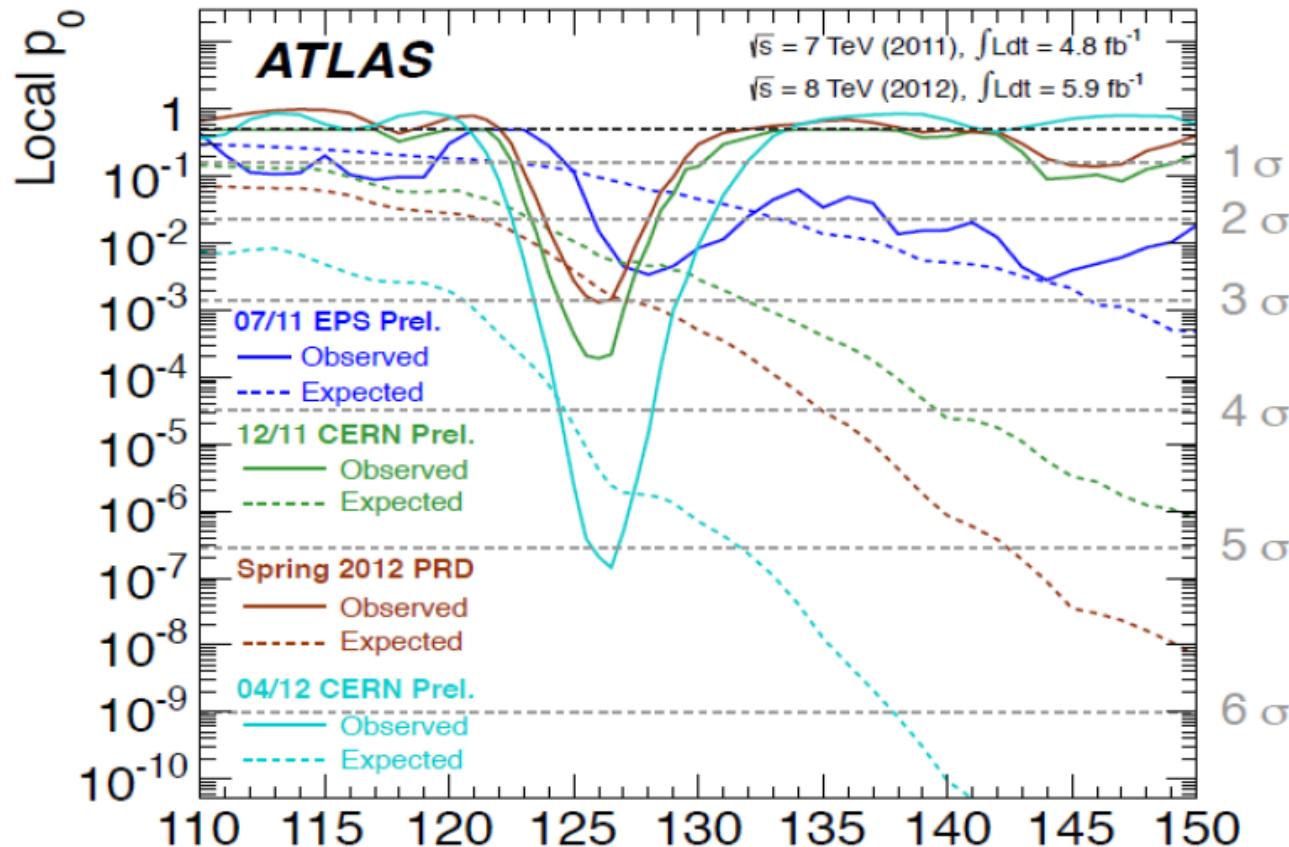


Local p_0

Probability that the background can produce a fluctuation greater than or equal to the excess observed in data. Equivalent in terms of number of standard deviations is called local significance.



Birth of a particle

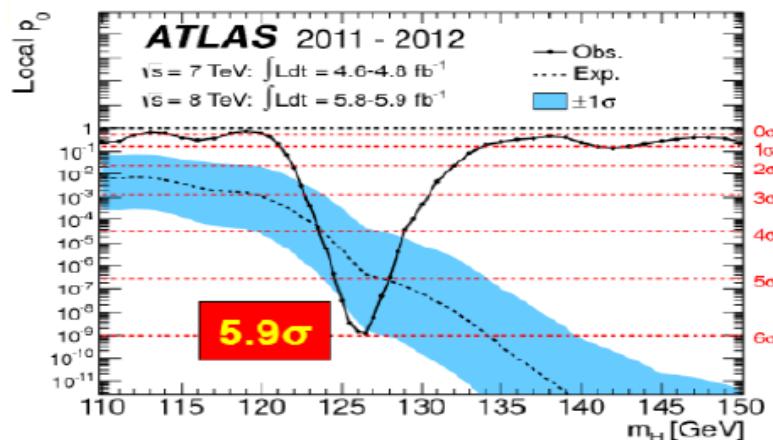


Probability of observing an excess at
one specific mass
(in absence of signal)...

m_H [GeV]

Higgs-like particle – 4 July 2012

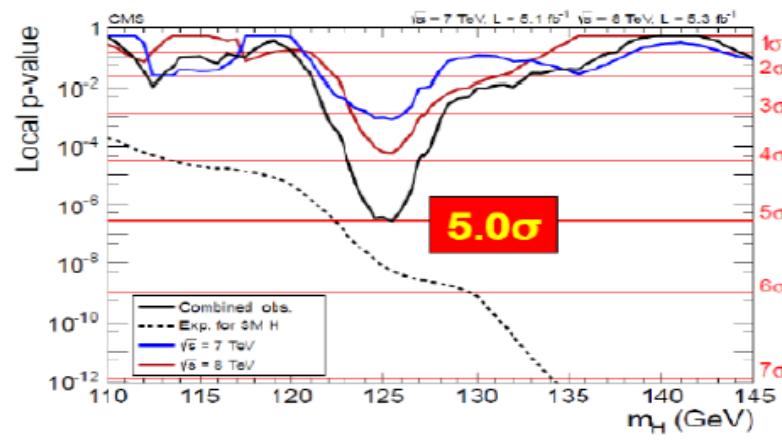
- We are living in a privileged moment in the history of High Energy Physics: **first fundamental scalar**
- The discovery came at half of the design energy, much more severe pile-up and one-third of integrated luminosity than was originally judged as



ATLAS [PLB 716 \(2012\) 1-29](#), Sept 17 (2012)

Largest local excess:
5.9 σ at $m_H = 126.5$ GeV

$H \rightarrow \gamma\gamma, b\bar{b}, \tau\bar{\tau}, WW(l\nu l\nu, l\nu q\bar{q}), ZZ(4l, ll\nu\nu, llq\bar{q})$

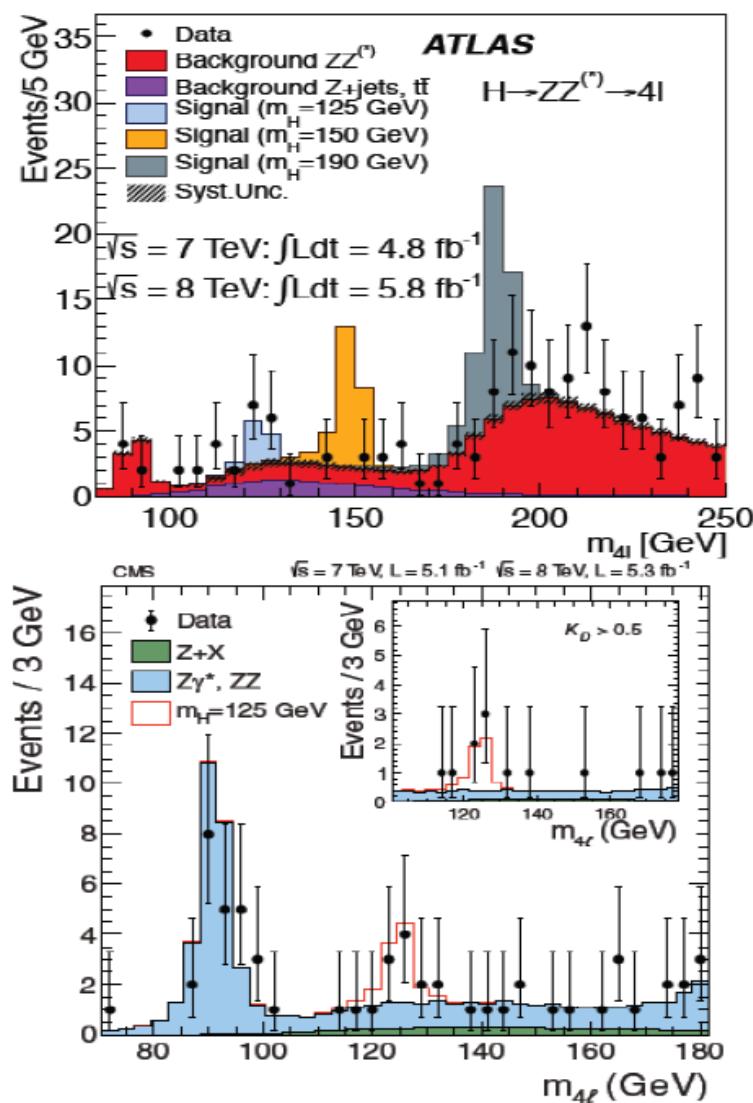
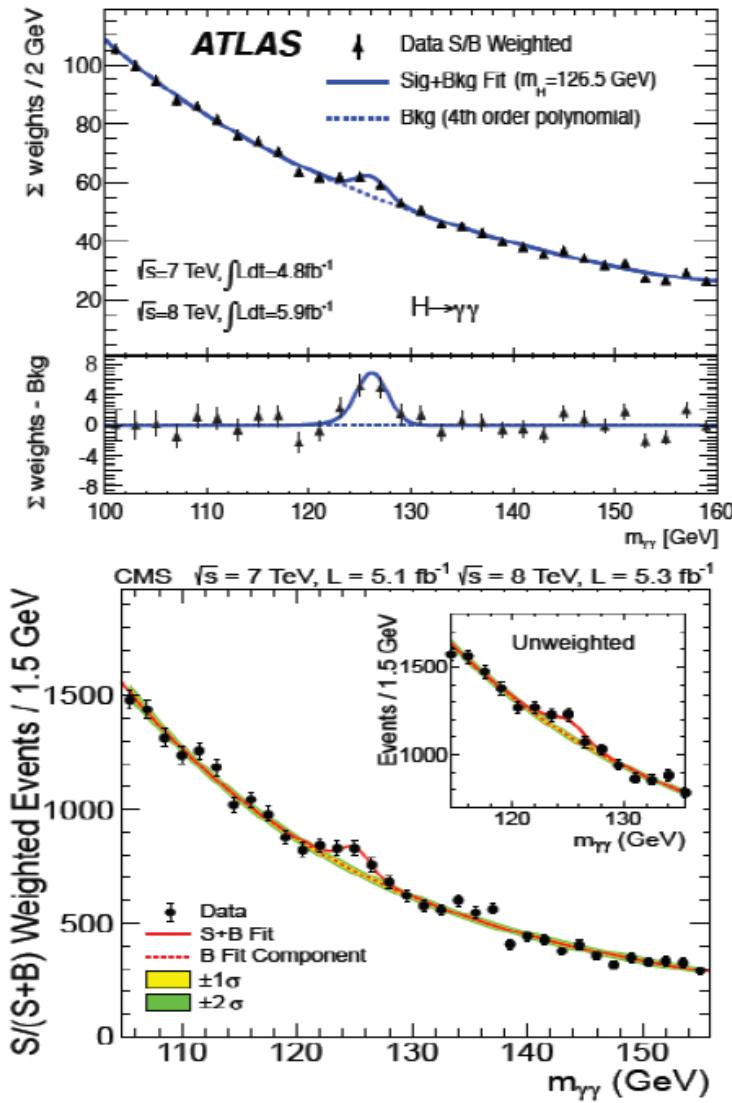


CMS [PLB 716 \(2012\) 30-61](#), Sept 17 (2012)

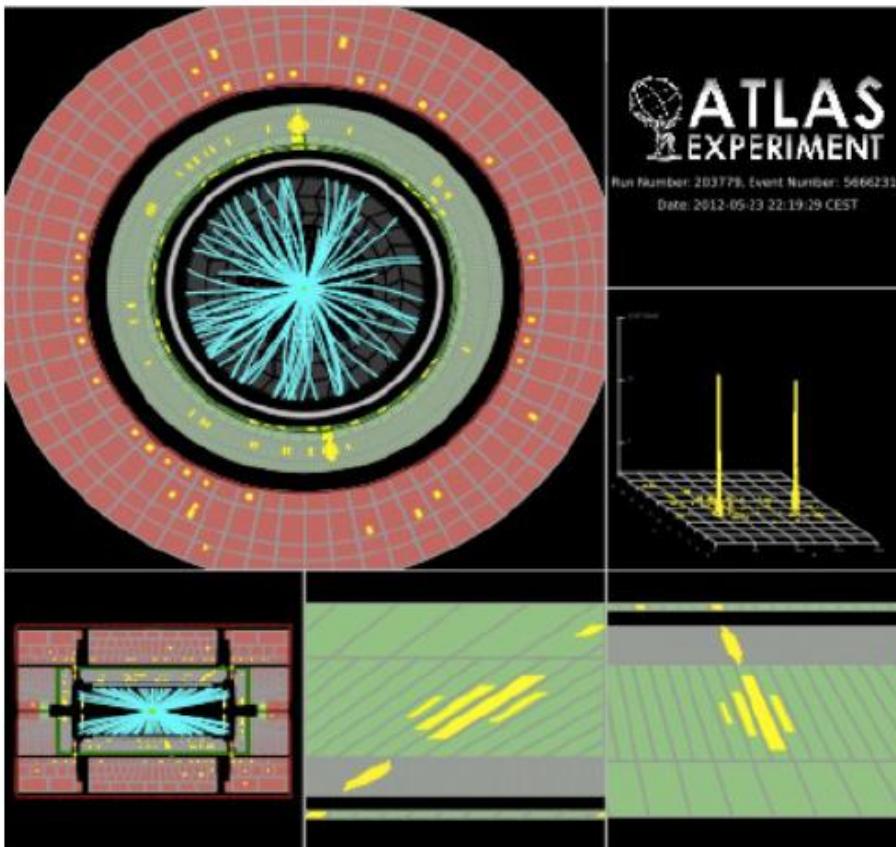
Largest local excess:
5.0 σ at $m_H = 125.5$ GeV

$H \rightarrow \gamma\gamma, b\bar{b}, \tau\bar{\tau}, WW(l\nu l\nu), ZZ(4l, ll\tau\tau, ll\nu\nu, llq\bar{q})$

Higgs-like particle – 4 July 2012



H-> $\gamma\gamma$: events signature



Simple event signature

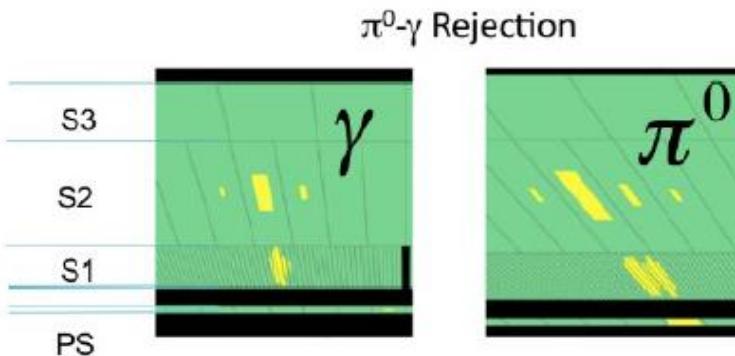
- Two high pT photons
 $pT_1 > 40 \text{ GeV}$ and $pT_2 > 30 \text{ GeV}$
- High trigger efficiency
 $\sim 99\%$
- High event selection efficiency
despite high jet-jet & γ -jet production
 $\sim 40\%$
- High signal over background
 $\sim 3\text{-}10\%$ (depending on sub-category)

$$\text{Invariant mass reconstruction } m_{\gamma\gamma}^2 = 2 * E_1 E_2 (1 - \cos \alpha)$$

- Good energy calibration
- Robust primary vertex reconstruction
- Excellent invariant mass resolution $\sim 1.6 \text{ GeV}$ with 90% of events within $\pm 2\sigma$

Shower shapes and vertex reconstr.

Photon ID 2 – Photon shower shapes and background rejection



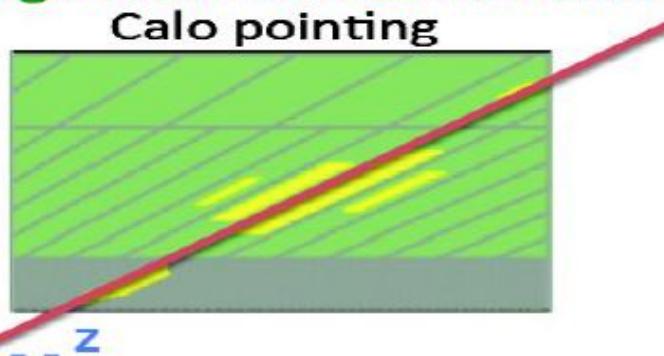
- Photons shower shape distributions in LAr sampling layers - different for signal and background (π^0)

Vertex Reconstruction

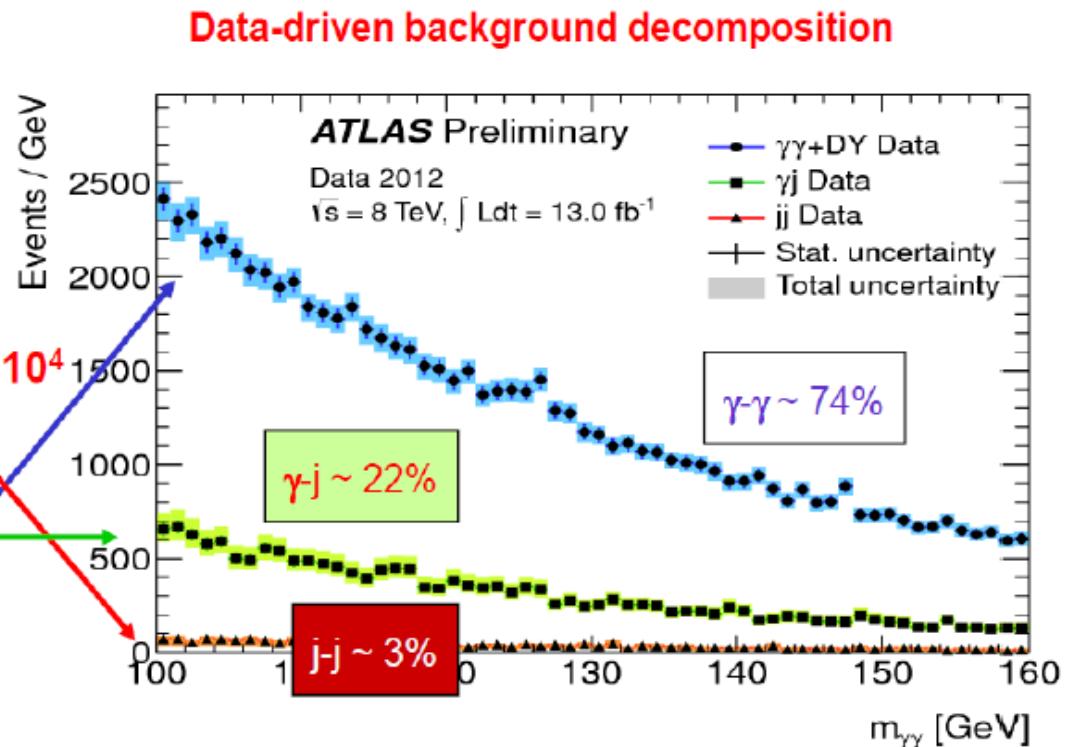
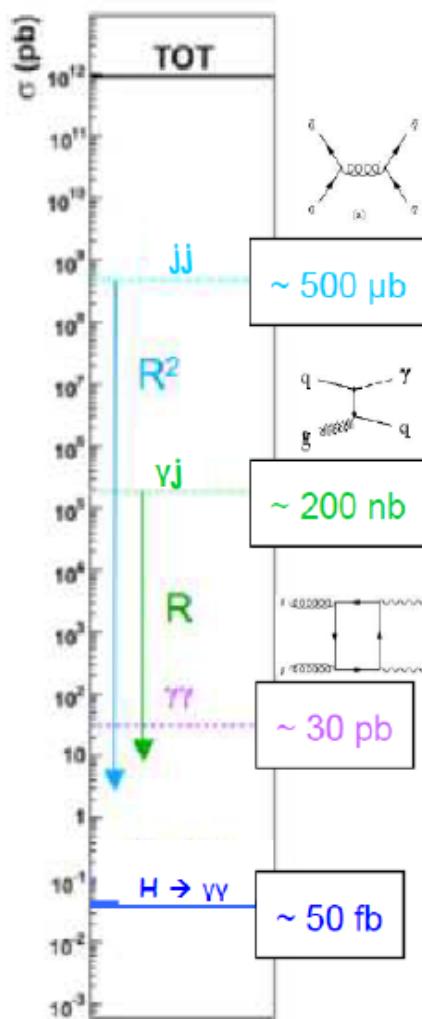
$$m_{\gamma\gamma}^2 = 2 * E_1 E_2 (1 - \cos \alpha)$$

- Vertex reconstructed through likelihood combination

- Calorimeter ‘pointing’
- Σ tracks pT^2
- Conversion vertex
- Mean vertex position



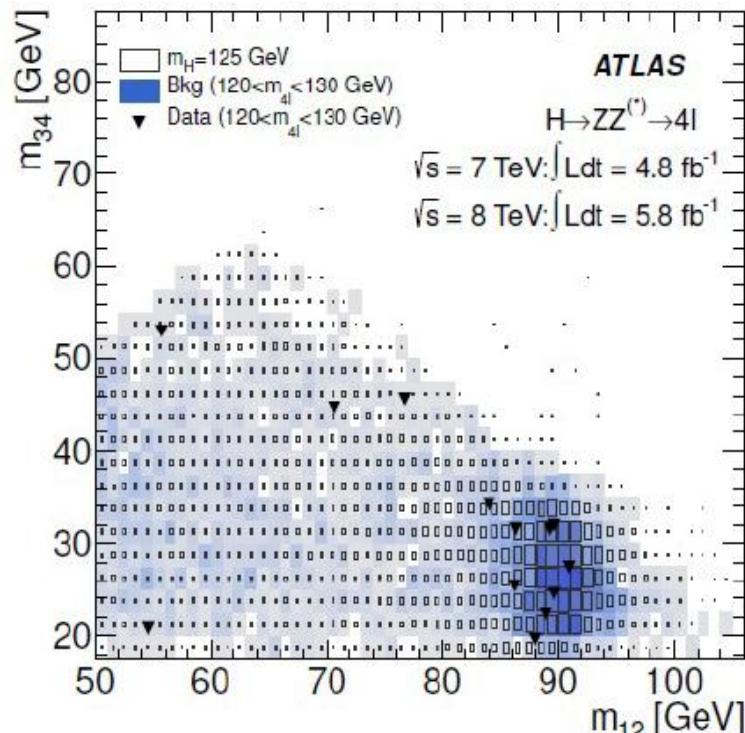
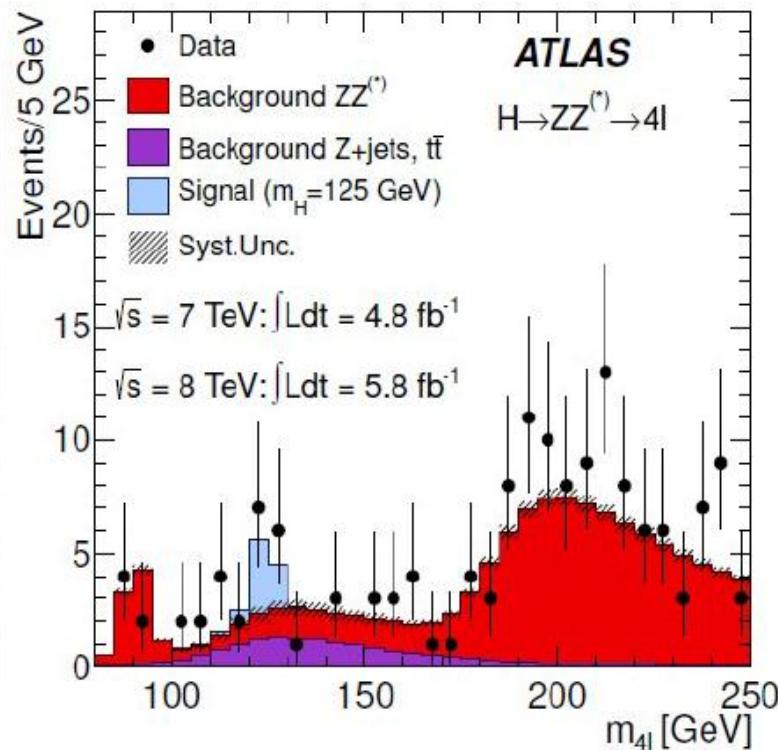
H $\rightarrow\gamma\gamma$: background rejection



→ Reducible background γ -jet and jet-jet largely below irreducible background $\gamma\gamma$

- ▶ Remark 1: this decomposition is not directly used in the Higgs search: the background is extrapolated from data sidebands
- ▶ Remark 2: Drell-Yan \sim negligible for $m_{\gamma\gamma} > 100 \text{ GeV}$ ($\sim 1\%$)

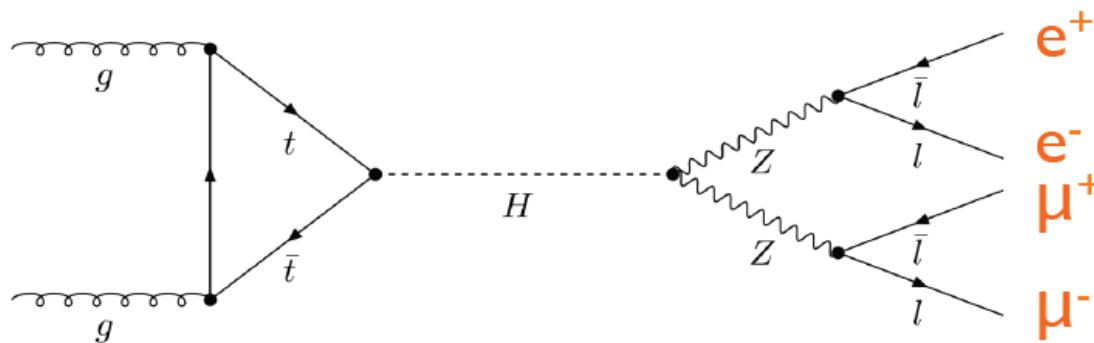
The golden channel: H->ZZ, Z->ll



In a m_{4l} window
around 120-130 GeV:

	Signal	$ZZ^{(*)}$	$Z + \text{jets}, t\bar{t}$	Observed
4μ	2.09 ± 0.30	1.12 ± 0.05	0.13 ± 0.04	6
$2e2\mu/2\mu2e$	2.29 ± 0.33	0.80 ± 0.05	1.27 ± 0.19	5
$4e$	0.90 ± 0.14	0.44 ± 0.04	1.09 ± 0.20	2

Signal and background

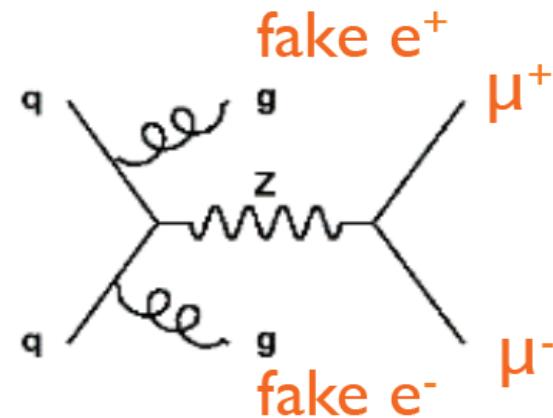
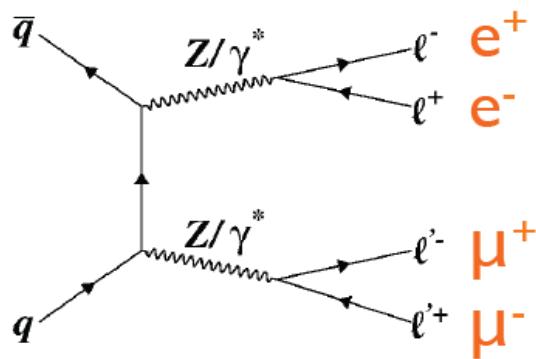


Irreducible background

The final state is exactly the same, but it does not come from the particle you are looking for

Reducible background

The final state looks like the same, but some f the particle fakes what you are looking for



4e candidate. $m_{4\ell} = 124.6$ GeV, $m_{12} = 70.6$ GeV, $m_{34} = 44.7$ GeV.

$e_1: P_T = 24.9$ GeV, $\eta = -0.33$, $\phi = 1.98$

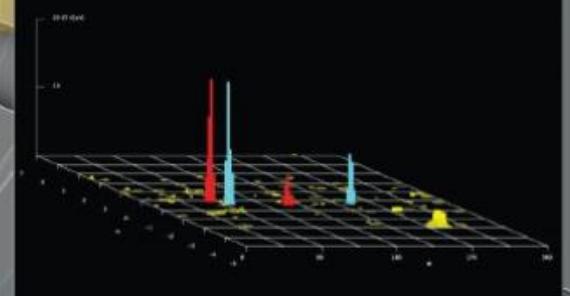
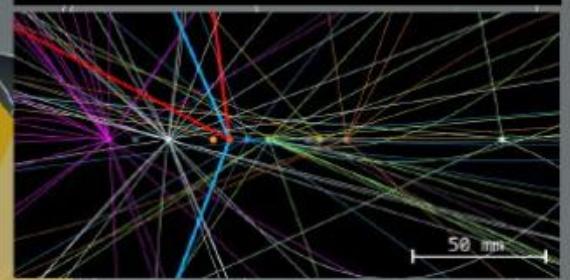
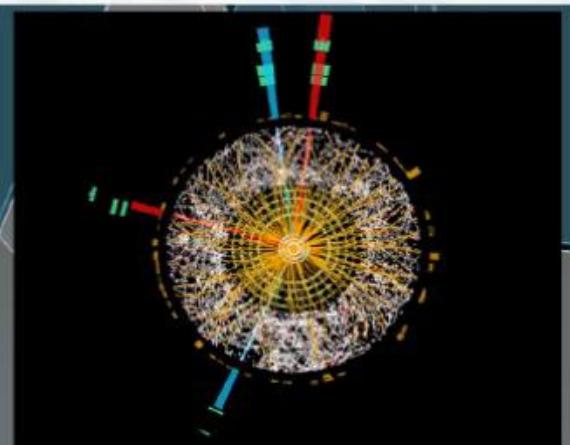
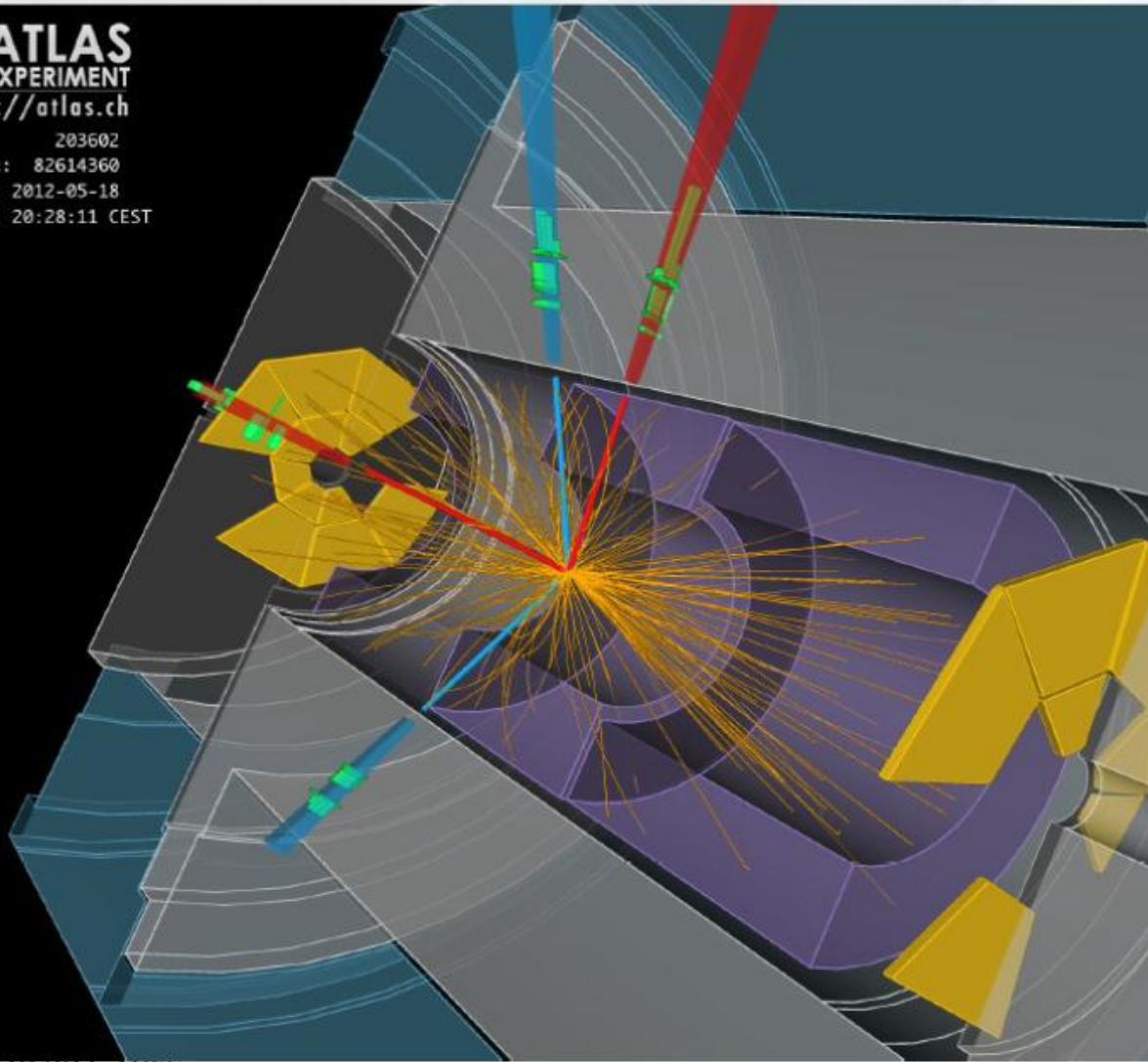
$e_2: P_T = 53.9$ GeV, $\eta = -0.40$, $\phi = 1.69$

$e_3: P_T = 61.9$ GeV, $\eta = -0.12$, $\phi = 1.45$

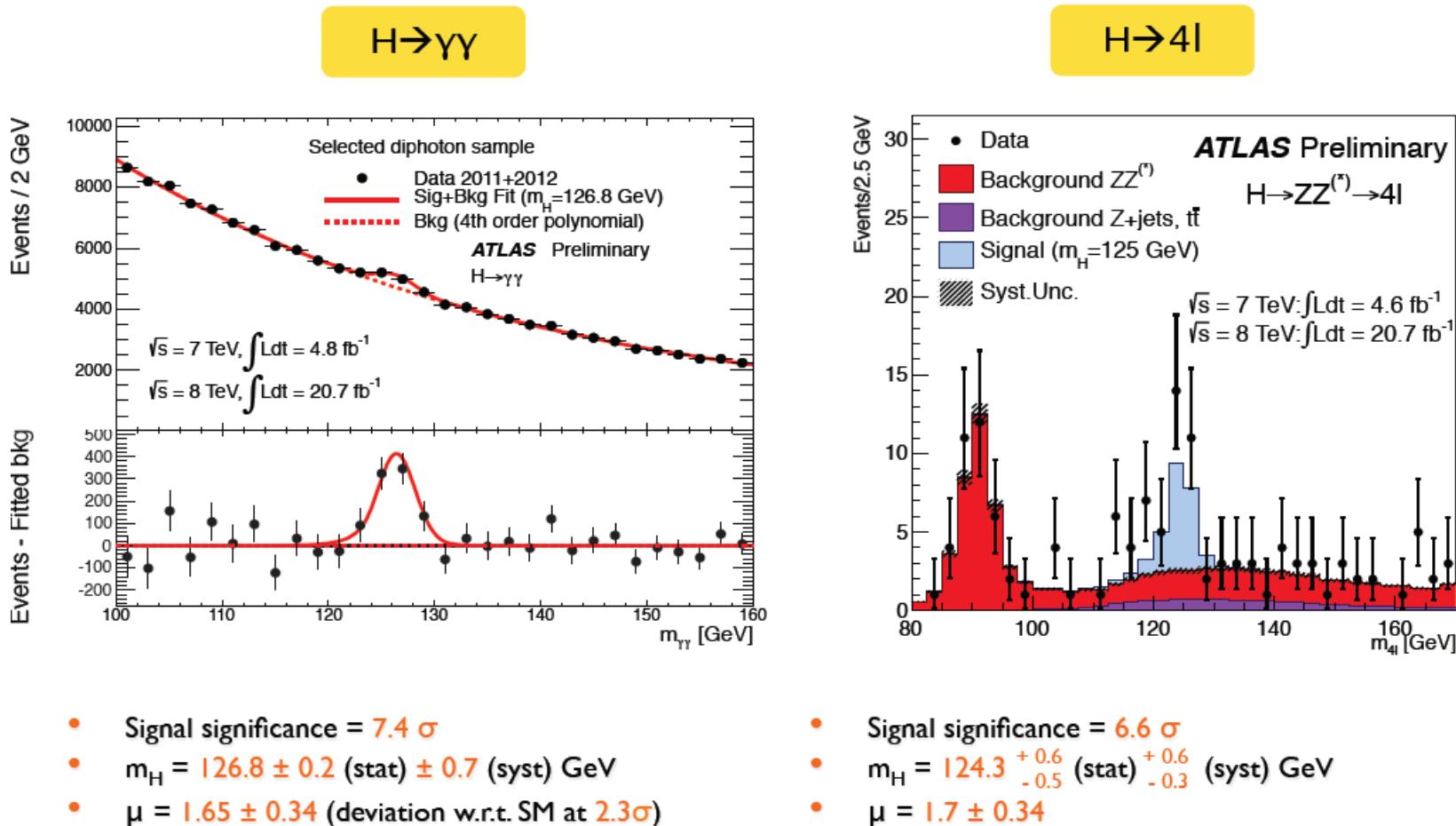
$e_4: P_T = 17.8$ GeV, $\eta = -0.51$, $\phi = 2.84$



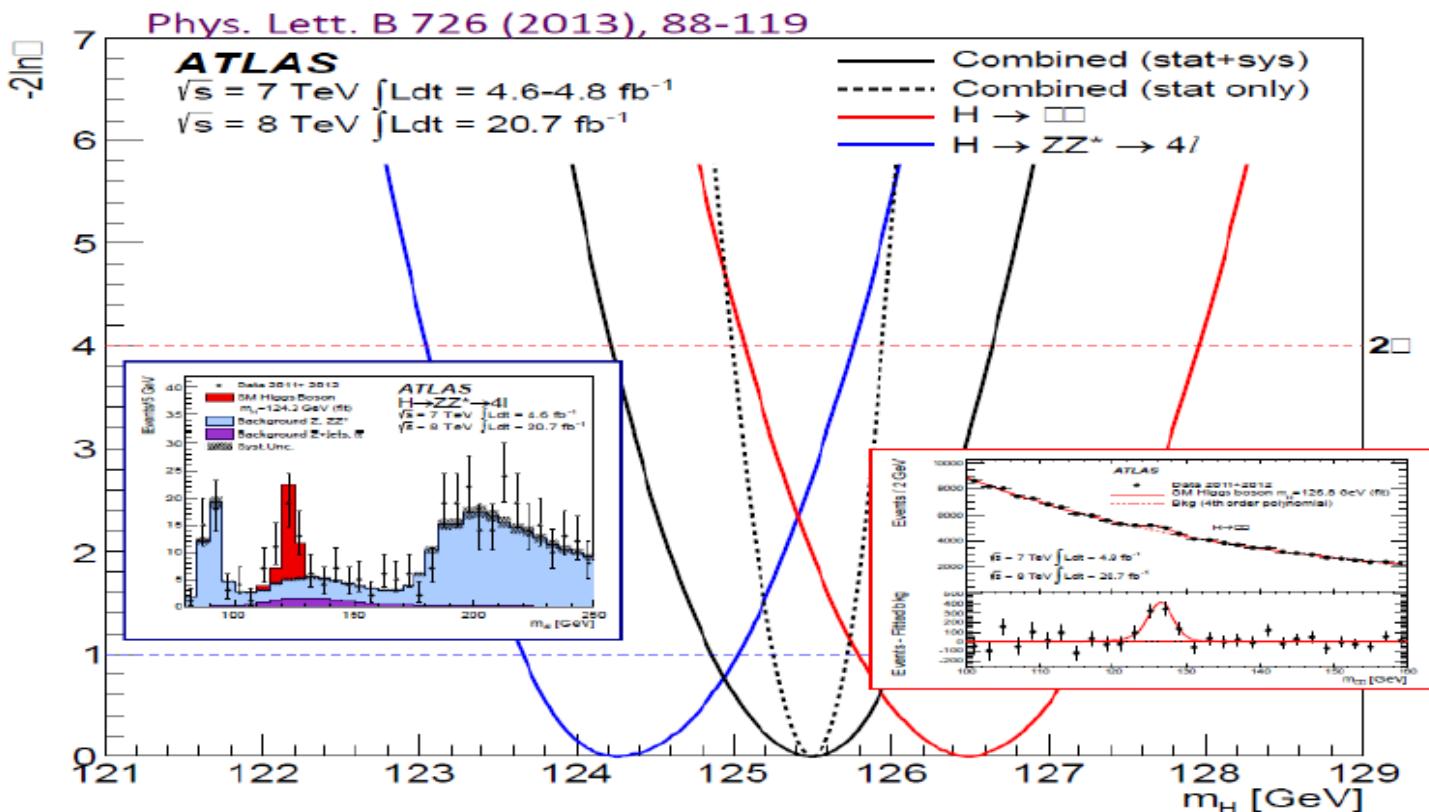
Run: 203602
Event: 82614360
Date: 2012-05-18
Time: 20:28:11 CEST



Higgs like signal with 7 TeV and 8 TeV data



Mass measurement

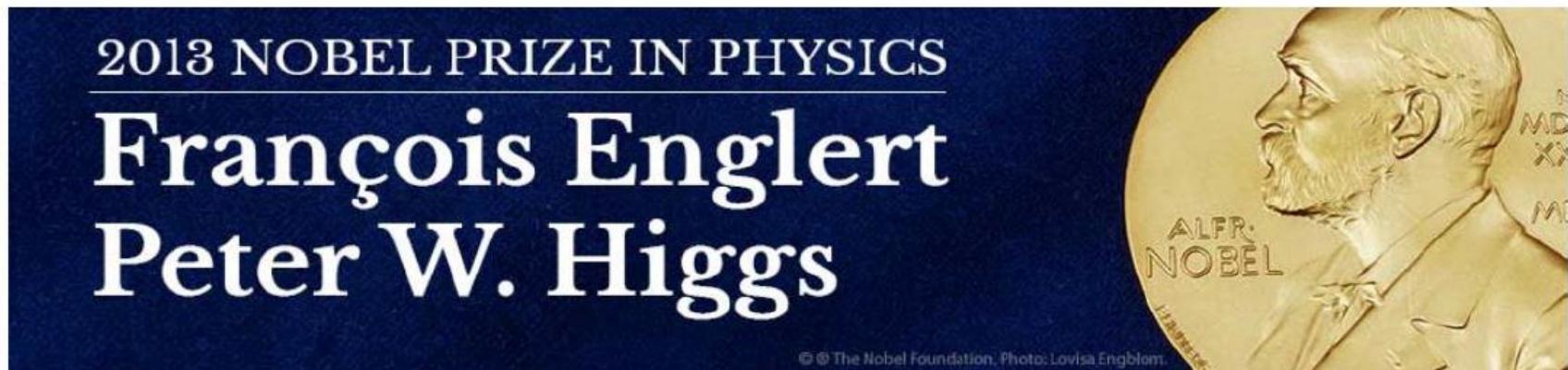


$$4\ell: M_H = 124.3 \pm 0.6_{\text{stat}} \pm 0.4_{\text{sys}} \text{ GeV}$$

$$\gamma\gamma: M_H = 126.8 \pm 0.2_{\text{stat}} \pm 0.7_{\text{sys}} \text{ GeV}$$

$$\text{Combined: } M_H = 125.5 \pm 0.2_{\text{stat}} \pm 0.6_{\text{sys}} \text{ GeV}$$

Nobel price for predicting Higgs particle



THE BEH-MECHANISM, INTERACTIONS WITH SHORT RANGE FORCES
AND SCALAR PARTICLES



8 October 2013

The Royal Swedish Academy of Sciences
has decided to award the Nobel Prize in
Physics for 2013 to

François Englert and Peter Higgs

"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

Entrance of the Higgs into PDG

2013

Higgs Bosons — H^0 and H^\pm

A REVIEW GOES HERE – Check our WWW List of Reviews

CONTENTS:

- H^0 (Higgs Boson)
 - H^0 Mass
 - H^0 Spin
 - H^0 Decay Width
 - H^0 Decay Modes
 - H^0 Signal Strengths in Different Channels
 - Combined Final States
 - $W^+ W^-$ Final State
 - $Z Z^*$ Final State
 - $\gamma \gamma$ Final State
 - $b\bar{b}$ Final State
 - $\pi^+ \pi^-$ Final State
- Standard Model H^0 (Higgs Boson) Mass Limits
 - H^0 Direct Search Limits
 - H^0 Indirect Mass Limits from Electroweak Analysis
- Searches for Other Higgs Bosons
 - Mass Limits for Neutral Higgs Bosons in Supersymmetric Models
 - H^0 (Higgs Boson) Mass Limits in Supersymmetric Models
 - A^0 (Pseudoscalar Higgs Boson) Mass Limits in Supersymmetric Models
 - H^0 (Higgs Boson) Mass Limits in Extended Higgs Models
 - Limits in General two-Higgs-doublet Models
 - Limits for H^0 with Vanishing Yukawa Couplings
 - Limits for H^0 Decaying to Invisible Final States
 - Limits for Light A^0
 - Other Limits
 - H^\pm (Charged Higgs) Mass Limits
 - Mass limits for $H^{\pm\pm}$ (doubly-charged Higgs boson)
 - Limits for $H^{\pm\pm}$ with $T_3 = \pm 1$
 - Limits for $H^{\pm\pm}$ with $T_3 = 0$

H^0 (Higgs Boson)

The observed signal is called a Higgs Boson in the following, although its detailed properties and in particular the role that the new particle plays in the context of electroweak symmetry breaking need to be further clarified. The signal was discovered in searches for a Standard Model (SM)-like Higgs. See the following section for mass limits obtained from those searches.

H^0 MASS	<small>VALUE (GeV)</small>	DOCUMENT ID	TECH	COMMENT
125.0 ± 0.4 OUR AVERAGE				
$125.0 \pm 0.4 \pm 0.4$		1 CHATRCHYAN13J	CMS	p.p. T and 8 TeV
$126.0 \pm 0.4 \pm 0.4$		2 AAD	12A	ATLAS p.p. T and 8 TeV
*** We do not use the following data for averages, fits, limits, etc. ***				
$126.2 \pm 0.6 \pm 0.2$		3 CHATRCHYAN13J	CMS	p.p. T and 8 TeV
$125.5 \pm 0.4 \pm 0.5$		4 CHATRCHYAN12N	CMS	p.p. T and 8 TeV

¹ Combined value from $Z Z$ and $\gamma \gamma$ final states.
² AAD 12A obtain results based on $4.6\text{--}4.8 \text{ fb}^{-1}$ of $p\bar{p}$ collisions at $E_{\text{cm}} = 7 \text{ TeV}$ and $5.0\text{--}5.9 \text{ fb}^{-1}$ at $E_{\text{cm}} = 8 \text{ TeV}$. An excess of events over background with a local significance of 3.9σ is observed at $m_{H^0} = 125 \text{ GeV}$. See also AAD 12D.
³ Result based on $Z Z \rightarrow 4\ell$ final state in 5.1 fb^{-1} of $p\bar{p}$ collisions at $E_{\text{cm}} = 7 \text{ TeV}$ and 12.3 fb^{-1} at $E_{\text{cm}} = 8 \text{ TeV}$.
⁴ CHATRCHYAN 12N obtain results based on $4.9\text{--}5.1 \text{ fb}^{-1}$ of $p\bar{p}$ collisions at $E_{\text{cm}} = 7 \text{ TeV}$ and $5.1\text{--}5.3 \text{ fb}^{-1}$ at $E_{\text{cm}} = 8 \text{ TeV}$. An excess of events over background with a local significance of 5.0σ is observed at about $m_{H^0} = 125 \text{ GeV}$. See also CHATRCHYAN 12D.

NODE=S655
NODE=S655

NODE=S655CNT
NODE=S655CNT

NODE=S655CNT

NODE=S655210

NODE=S655210

NODE=S655HBM
NODE=S655HBM

OCCUR=2

NODE=S655HBM;LINKAGE=CA
NODE=S655HBM;LINKAGE=AA

NODE=S655HBM;LINKAGE=CT

NODE=S655HBM;LINKAGE=CH

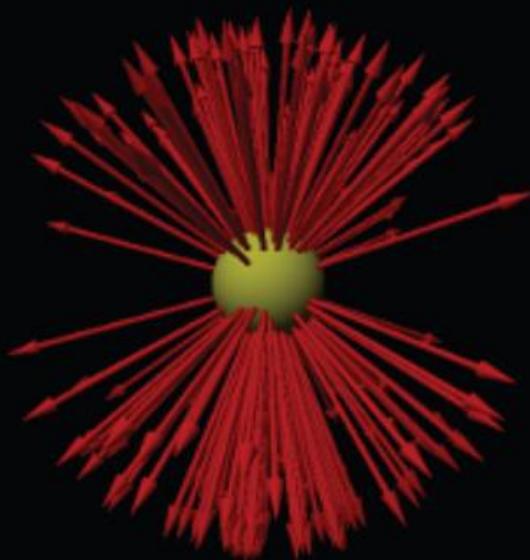
Inaugural entrance of
the Higgs boson in the
PDG particle listing !

H^0

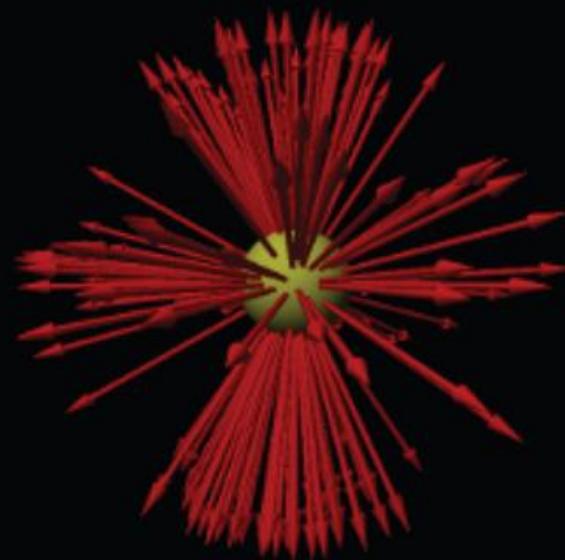
How we can recognize spin?



spin 0



spin 1

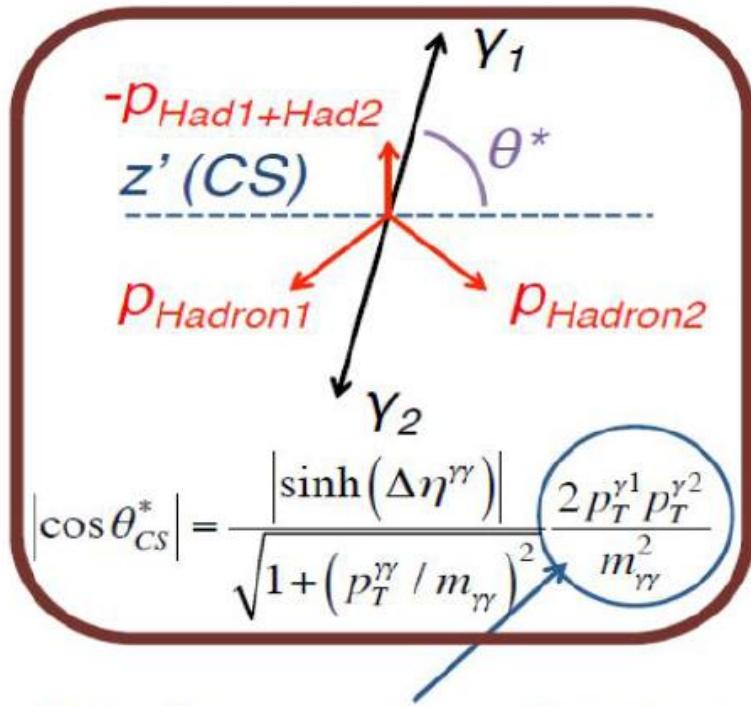


spin 2

Spin-0 decays in all directions with equal probability; spin-1 prefers decaying toward or away from the direction of spin; spin-2 prefers the poles and the equator to the region in between. These pictures exaggerate the real distributions for clarity.

Spin observables for H-> $\gamma\gamma$

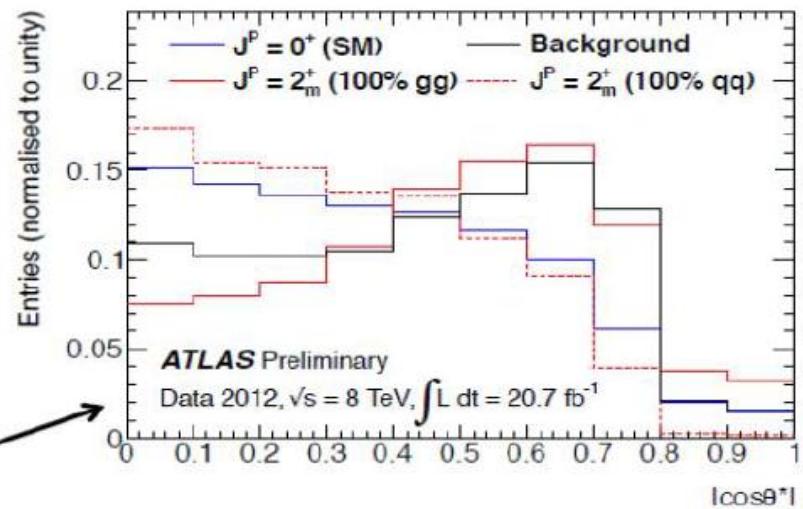
Separate 0⁺ and 2⁺ spin hypotheses using the angular correlation of the two photons



Relative p_T cuts on the photons remove most correlation with $m_{\gamma\gamma}$
 $qq \rightarrow 2^+$ very similar to SM $gg \rightarrow 0^+$

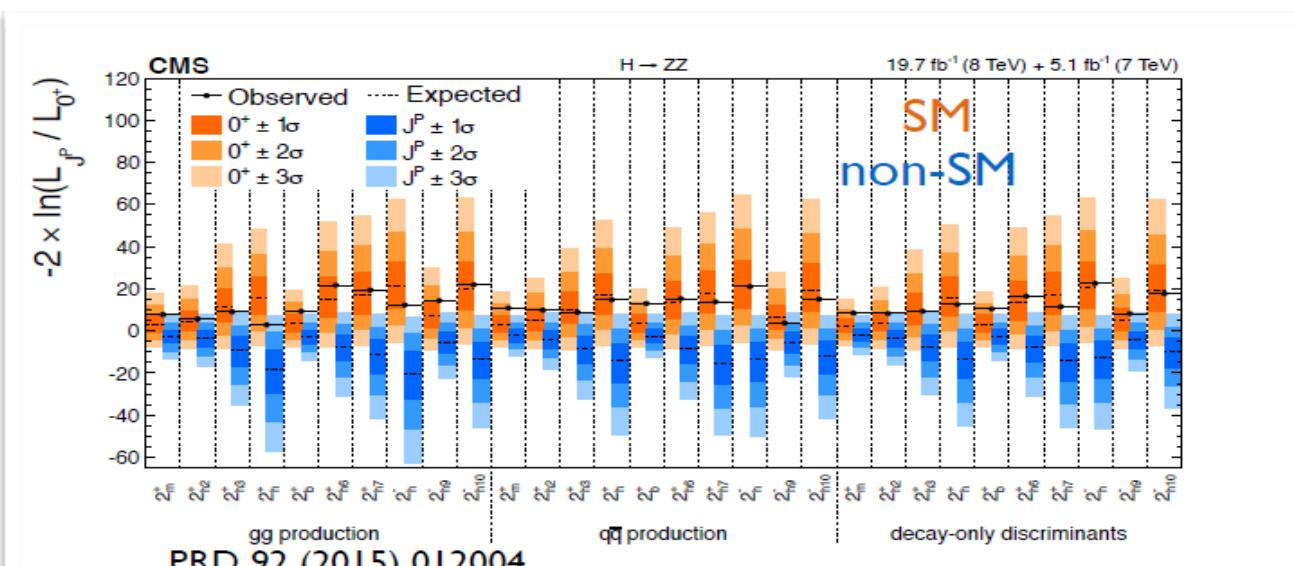
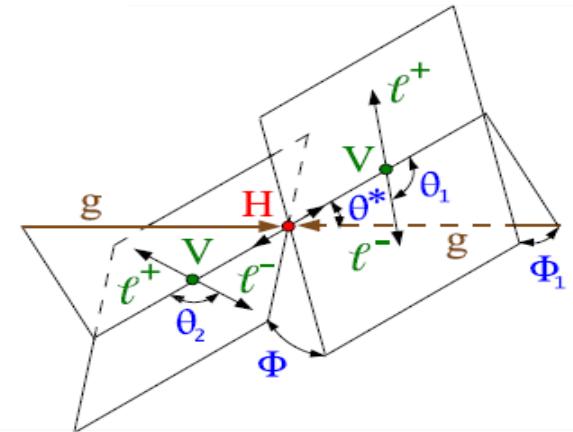
Collins-Soper frame used to get reference axis z' for $\cos(\theta^*)$

- z-axis bisects angle between the momenta of colliding hadrons
- Minimizes impact of ISR
- Better 0⁺ / 2⁺ discrimination



Spin study with H->4l

- SM predicts $J^{PC} = 0^{++}$
- Angular distributions sensitive to JP
- Wide range of alternative quantum numbers excluded at $>99\% \text{ CL}$
- All observations consistent with expectations for the SM Higgs boson



Tests of
alternative J^P
hypotheses in ZZ

Higgs boson decay width

$$m_h = 125 \text{ GeV} \rightarrow \Gamma_h = 4.07 \text{ MeV}$$

$$\tau_h = 1.62 \cdot 10^{-22} \text{ [s]}$$

A deviation would imply a decay to non-SM particles

Differential Higgs production cross-section

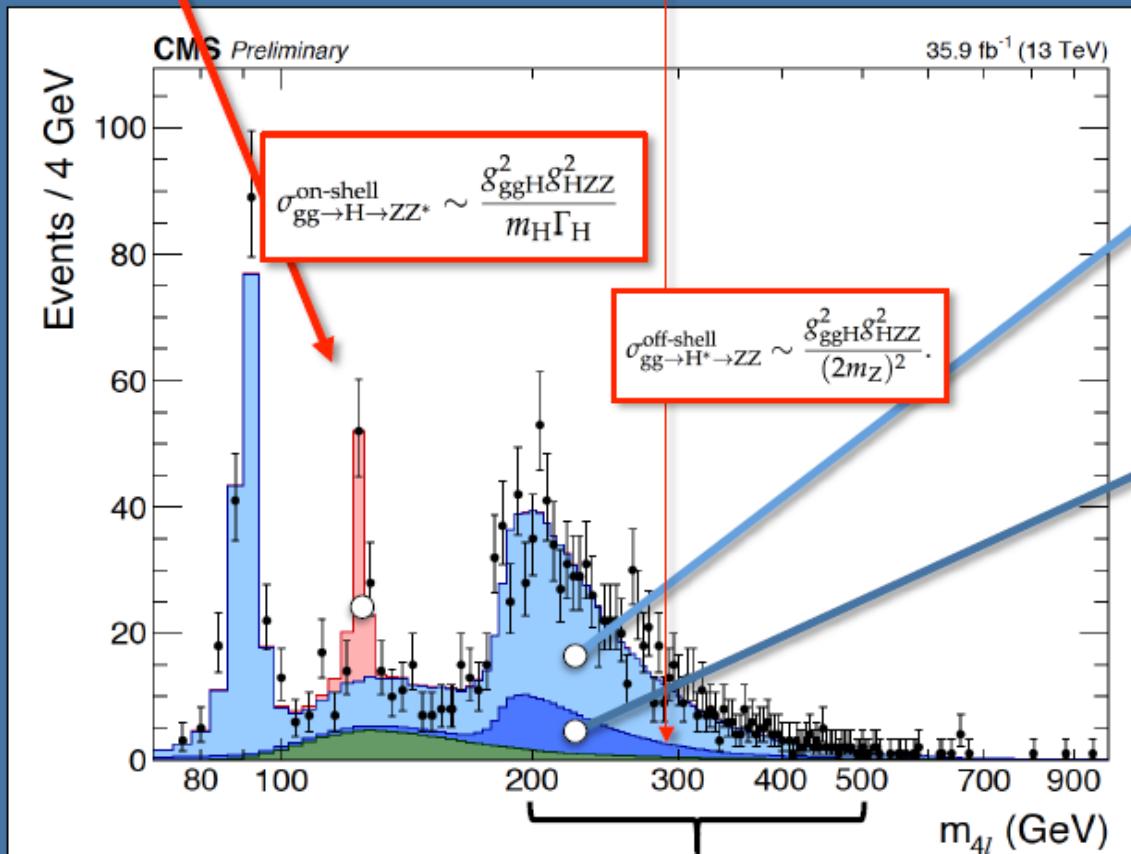
$$\frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}}{dm_{ZZ}^2} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{(m_{ZZ}^2 - m_H^2)^2 + m_H^2 \Gamma_H^2},$$

Γ_h cannot be accessed directly (experimental resolution $\sim 1-2 \text{ GeV}$)

A

higgs \rightarrow ZZ

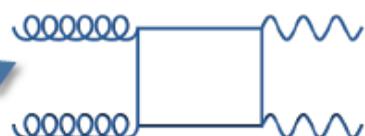
Indirect measurement



B

qq \rightarrow ZZ

C

gg \rightarrow ZZ

Interference between A and C : $(A+C)^2 = A^2 + C^2 + 2AC$

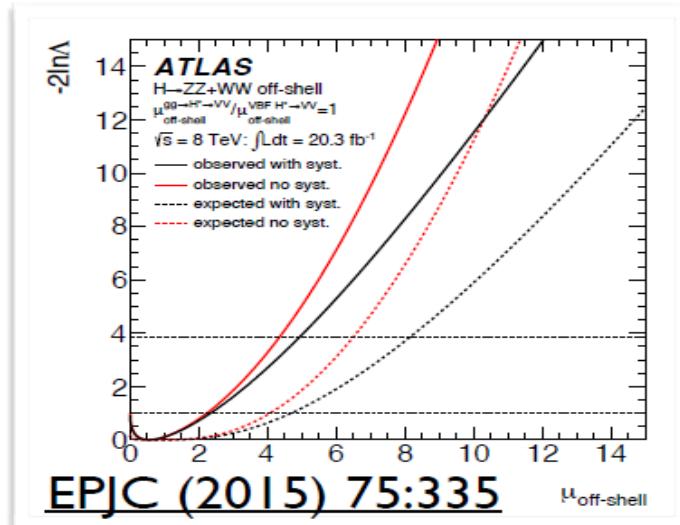
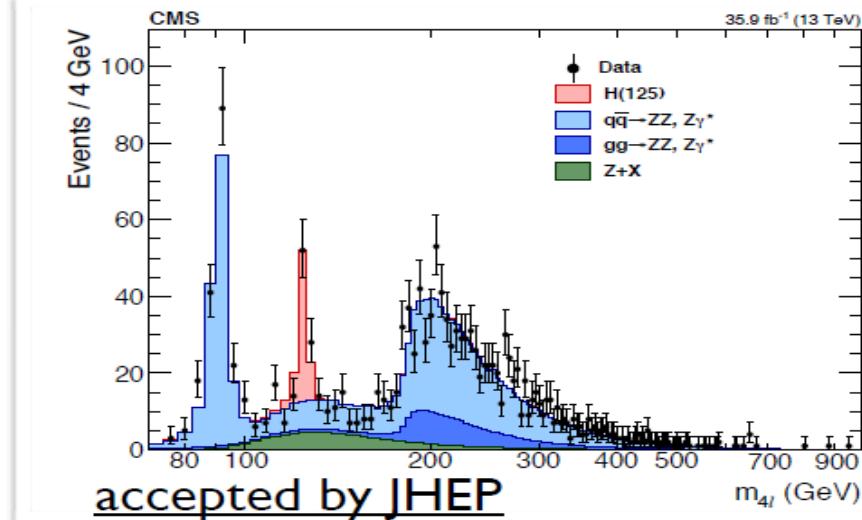
tiny ←

→ *accessible*

Higgs boson decay width

Total width

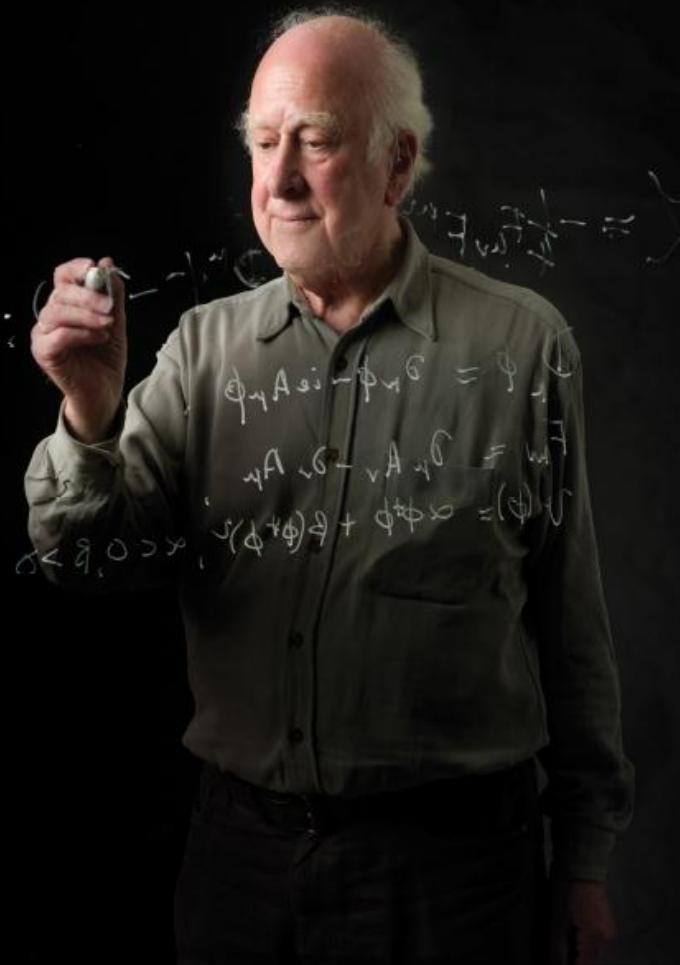
- Lower bound on total width from decay measurements
- Direct experimental measurements probe 3 orders of magnitude larger than SM width ($\Gamma=4$ MeV)
- Indirect constraint* on the width via measurement of ratio of off-peak to on-peak cross-section
 - CMS: $\Gamma < 13$ MeV
 - ATLAS: $\Gamma < 22$ MeV



*N. Kauer and G. Passarino, JHEP (2012) 2012: 116

*F. Caola and K. Melnikov, PRD88 (2013) 054024

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
 & + i \bar{\psi} D \psi + h.c. \\
 & + \bar{\psi}_i \gamma_{ij} \psi_j \phi + h.c. \\
 & + |D_m \phi|^2 - V(\phi)
 \end{aligned}$$



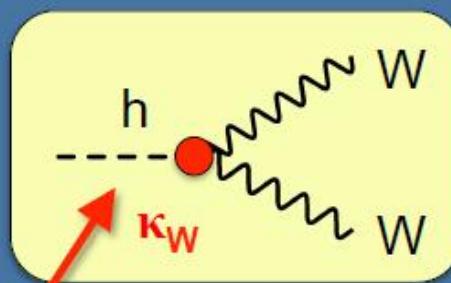
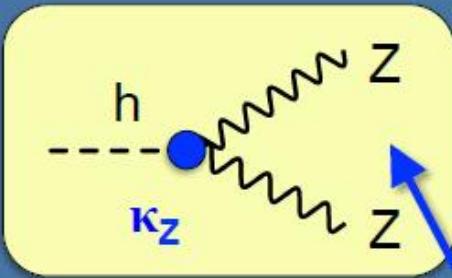
Couplings: kappa-framework

Scale factor for each (fundamental) coupling:

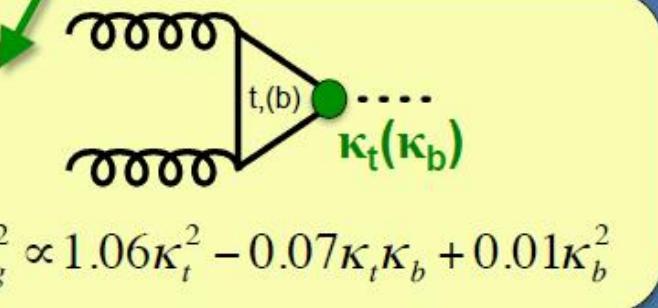
$$\sigma(i \rightarrow h \rightarrow f) = \kappa_i^2 \sigma_i^{SM} \frac{\kappa_f^2 \Gamma_f^{SM}}{\kappa_h^2 \Gamma_h^{SM}}$$

$$\begin{aligned}\mathcal{L} = & \kappa_3 \frac{m_H^2}{2v} H^3 + \kappa_Z \frac{m_Z^2}{v} Z_\mu Z^\mu H + \kappa_W \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} H \\ & + \kappa_g \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} H + \kappa_\gamma \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H + \kappa_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H \\ & - \left(\kappa_t \sum_{f=u,c,t} \frac{m_f}{v} f \bar{f} + \kappa_b \sum_{f=d,s,b} \frac{m_f}{v} f \bar{f} + \kappa_\tau \sum_{f=e,\mu,\tau} \frac{m_f}{v} f \bar{f} \right) H\end{aligned}$$

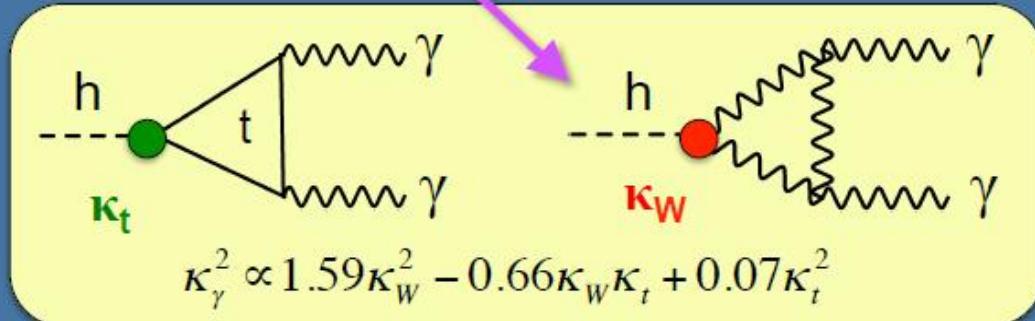
Scale Higgs boson couplings (wrt SM): production & decay



$$\mathcal{L} = \kappa_3 \frac{m_H^2}{2v} H^3 + \kappa_Z \frac{m_Z^2}{v} Z_\mu Z^\mu H + \kappa_W \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} H - \kappa_g \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} H - \kappa_\gamma \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H + \kappa_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H - \left(\kappa_t \sum_{f=u,c,t} \frac{m_f}{v} f \bar{f} + \kappa_b \sum_{f=d,s,b} \frac{m_f}{v} f \bar{f} + \kappa_\tau \sum_{f=e,\mu,\tau} \frac{m_f}{v} f \bar{f} \right) H$$



$$\kappa_g^2 \propto 1.06\kappa_t^2 - 0.07\kappa_t\kappa_b + 0.01\kappa_b^2$$



$$\kappa_\gamma^2 \propto 1.59\kappa_W^2 - 0.66\kappa_W\kappa_t + 0.07\kappa_t^2$$

What do we measure?

We measure event yields

We want to derive couplings
and signal strengths

The first thing we want to
measure is the the “signal
strength” per channel

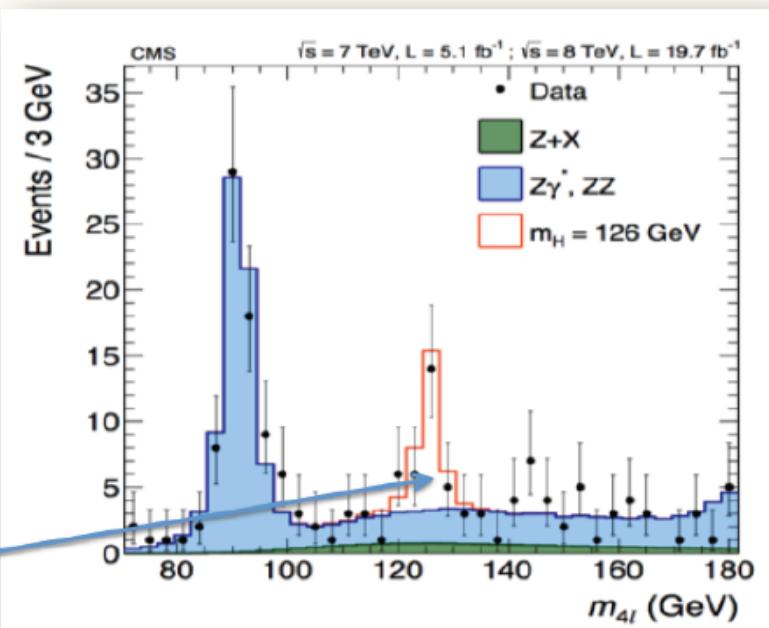
The analysis is using
discriminators (usually
reconstructed mass related)
to increase S/B

$$n_s^i = \mu^i \times \sum_p (\sigma^p \times Br^i)_{SM} \times A_p^i \times \epsilon_p^i \times Lumi$$

$p \in (ggF, VBF, VH, ttH)$ $i \in (\gamma\gamma, ZZ, WW, bb, \tau\tau)$

$$\mu^{ZZ}(@125.5 \text{ GeV}) = 1.44^{+0.40}_{-0.35}$$

$$\mu^{ZZ}(@125.6 \text{ GeV}) = 0.93^{+0.26+0.13}_{-0.23-0.09}$$



6.6 σ (4.4 exp) ATLAS

6.8 σ (6.7 exp) CMS

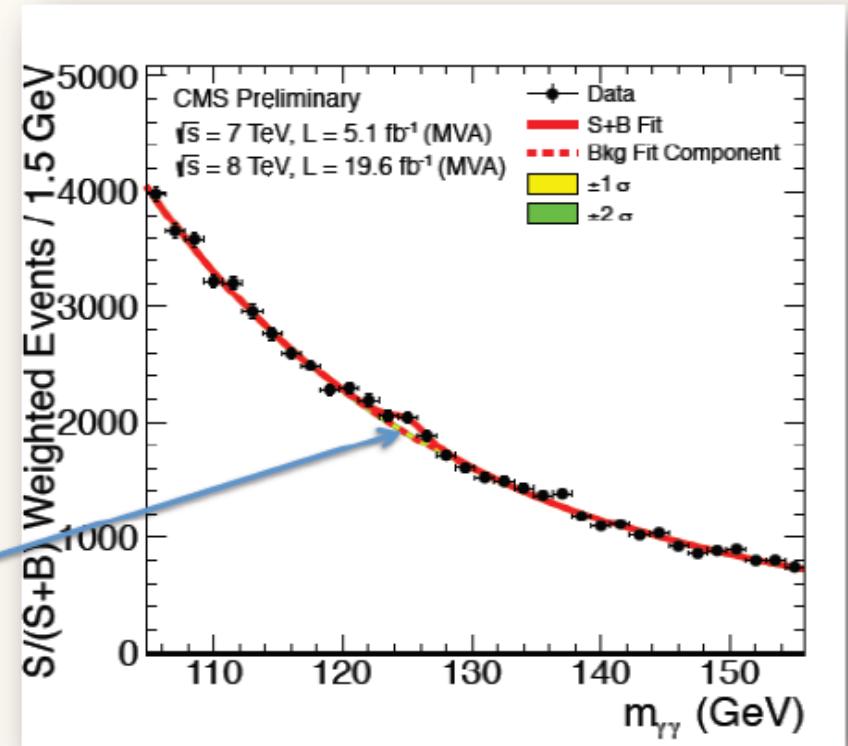
What do we measure?

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$$n_s^i = \mu^i \times \sum_p (\sigma^p \times Br^i)_{SM} \times A_p^i \times \epsilon_p^i \times \text{Lumi}$$

$p \in (ggF, VBF, VH, ttH)$ $i \in (\gamma\gamma, ZZ, WW, bb, \tau\tau)$

Higgs boson decay channels

Significance

7.4σ (4.3σ)

6.6σ (4.4σ)

3.8σ (3.8σ)

4.1σ (3.2σ)

0.36σ (1.64σ)

Obs. (*Exp.*) Combined

3.2σ (4.2σ)

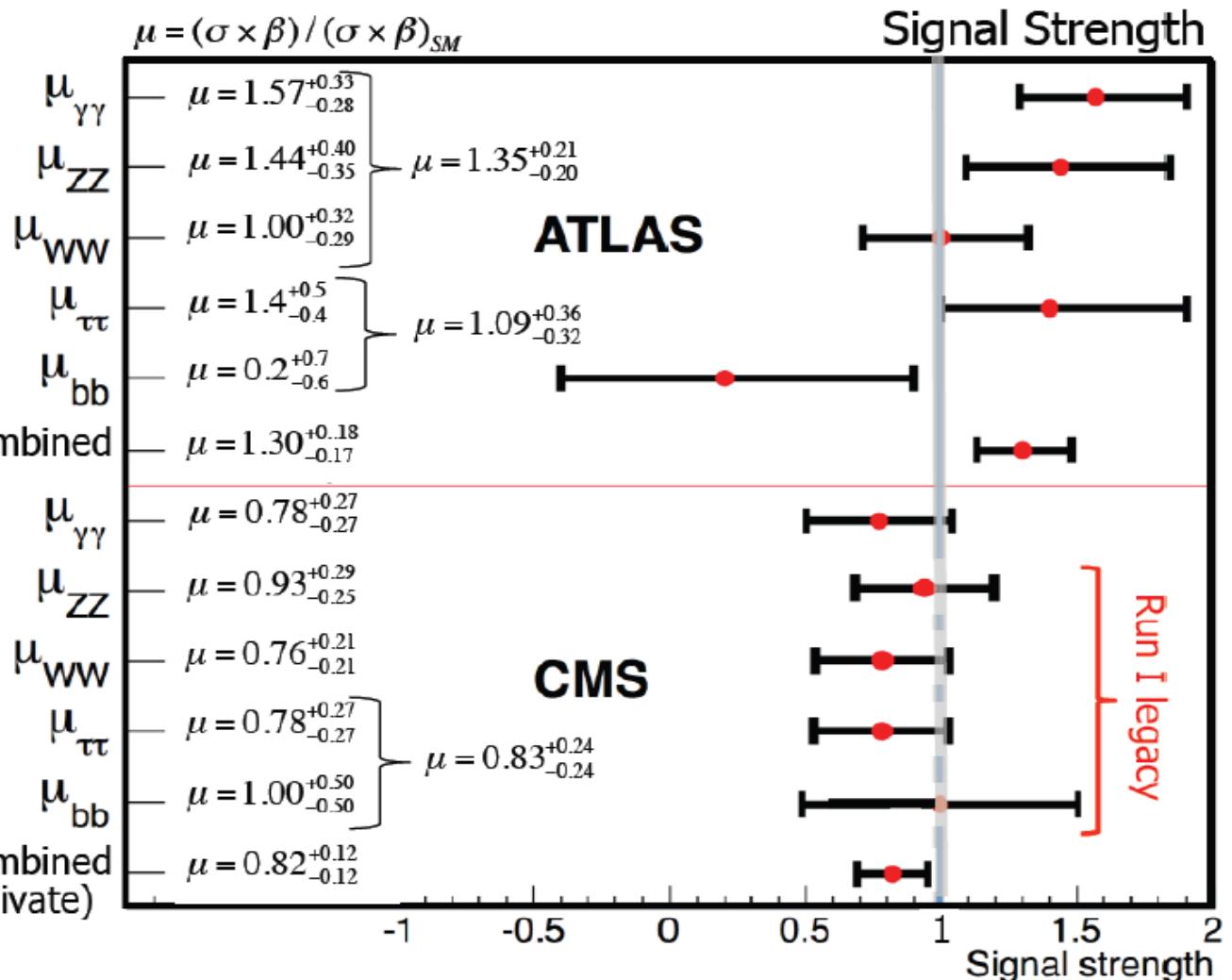
6.8σ (6.7σ)

4.3σ (5.8σ)

3.3σ (3.7σ)

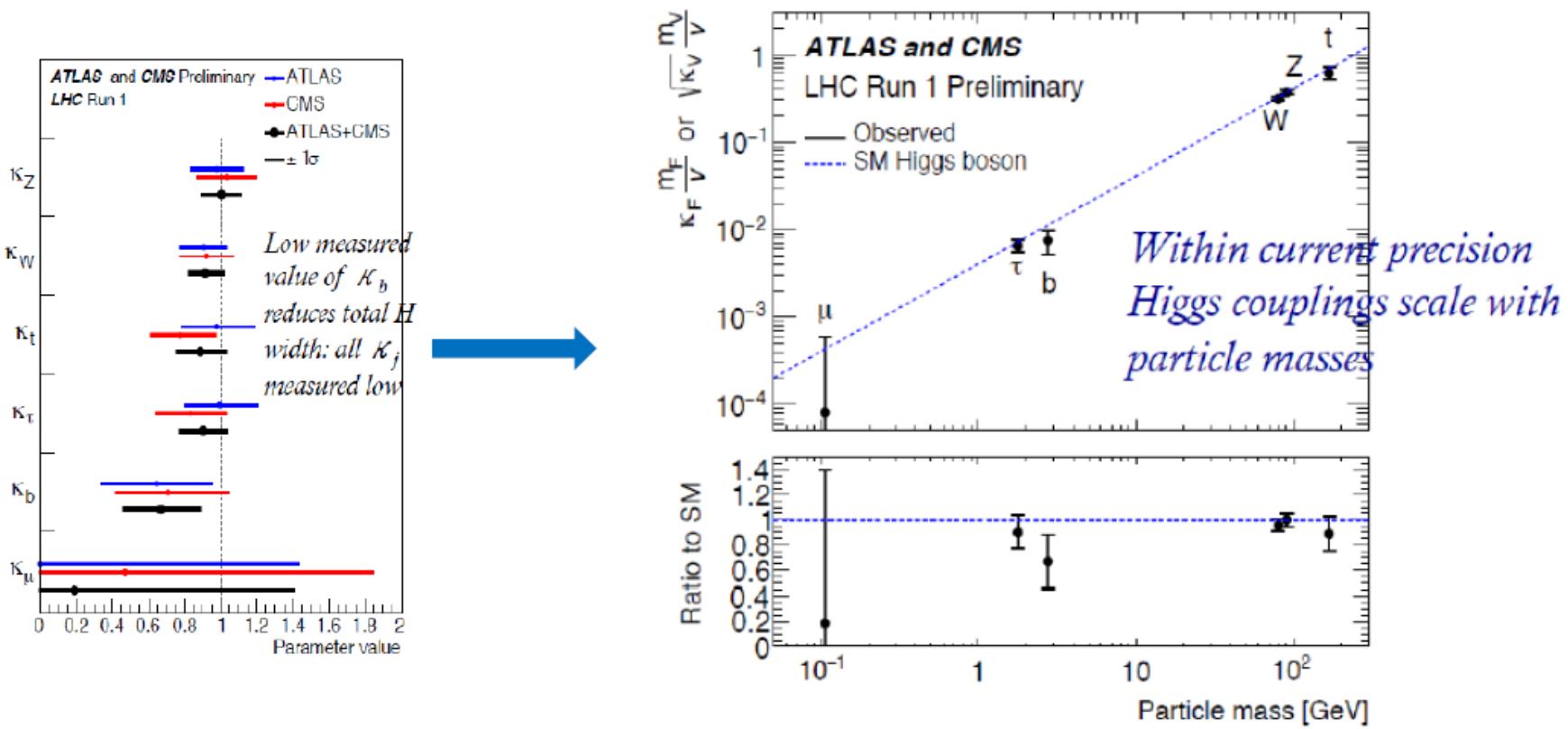
2.1σ (2.1σ)

Obs. (*Exp.*) Combined
(private)

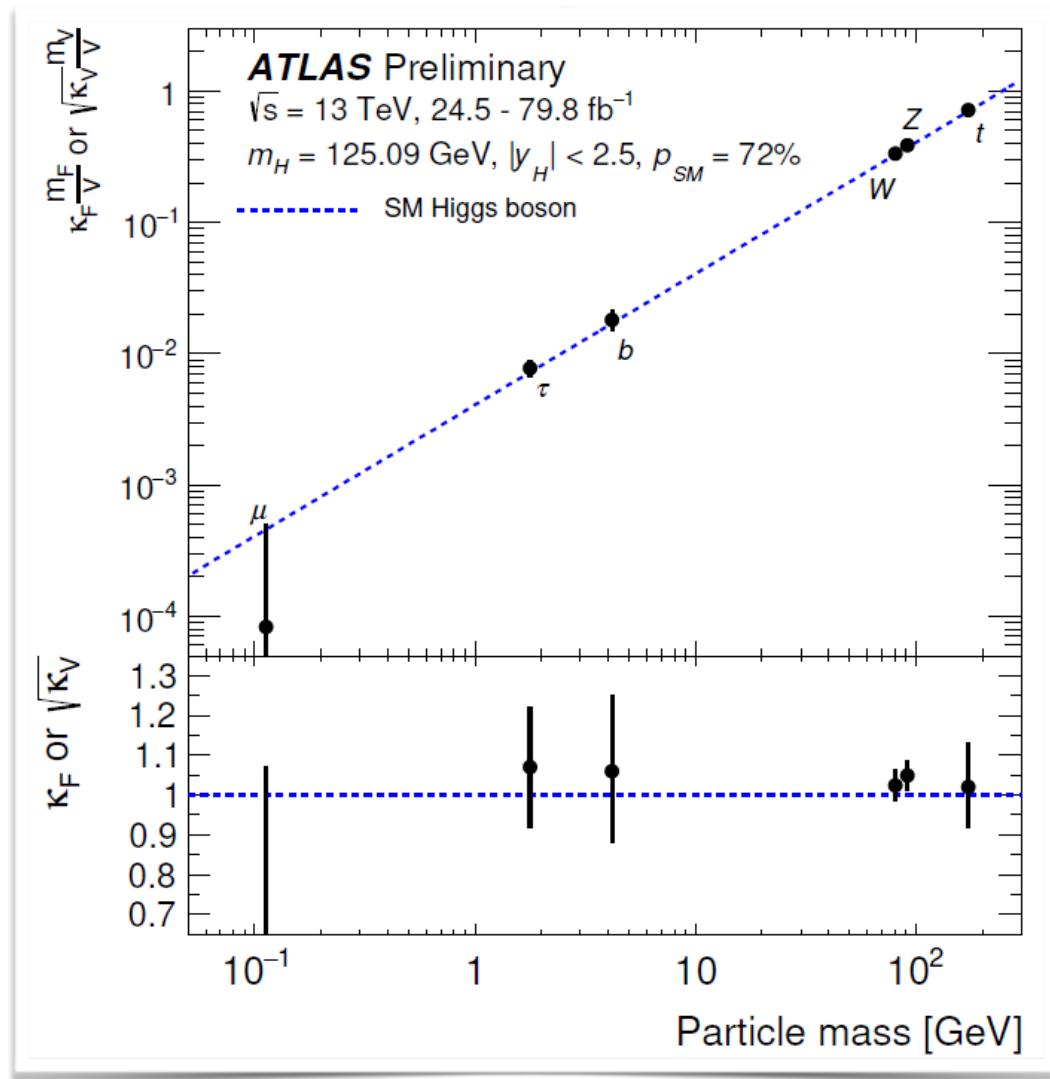


Constraints on three-level Higgs couplings

- Assume only SM physics in loops, no invisible or unseen BSM Higgs decays
- Fit for scaling parameters for Higgs couplings to W, Z, b, t, τ , μ



Couplings vs mass (2019)



Interpret the results in the κ framework as a function of the particle mass assuming no BSM contributions to the total width

$$\kappa_j^2 = \frac{\sigma_j}{\sigma_j^{\text{SM}}} \quad \text{or} \quad \kappa_j^2 = \frac{\Gamma^j}{\Gamma_{\text{SM}}^j}$$

[ATLAS-CONF-2019-005](#)

ATLAS combination

	$H \rightarrow \gamma\gamma$	$H \rightarrow ZZ^*$	$H \rightarrow WW^*$	$H \rightarrow \tau\tau$	$H \rightarrow bb$	
$t\bar{t}H$	$t\bar{t}H$ leptonic (3 categories) $t\bar{t}H$ hadronic (4 categories)	$t\bar{t}H$ multilepton 1 $\ell + 2 \tau_{\text{had}}$ $t\bar{t}H$ multilepton 2 opposite-sign $\ell + 1 \tau_{\text{had}}$ $t\bar{t}H$ multilepton 2 same-sign ℓ (categories for 0 or 1 τ_{had}) $t\bar{t}H$ multilepton 3 ℓ (categories for 0 or 1 τ_{had}) $t\bar{t}H$ multilepton 4 ℓ (except $H \rightarrow ZZ^* \rightarrow 4\ell$) $t\bar{t}H$ leptonic, $H \rightarrow ZZ^* \rightarrow 4\ell$ $t\bar{t}H$ hadronic, $H \rightarrow ZZ^* \rightarrow 4\ell$				$t\bar{t}H$ 1 ℓ , boosted $t\bar{t}H$ 1 ℓ , resolved (11 categories) $t\bar{t}H$ 2 ℓ (7 categories)
VH	VH 2 ℓ VH 1 ℓ , $p_T^{\ell+E_T^{\text{miss}}} \geq 150$ GeV VH 1 ℓ , $p_T^{\ell+E_T^{\text{miss}}} < 150$ GeV VH $E_T^{\text{miss}}, E_T^{\text{miss}} \geq 150$ GeV VH $E_T^{\text{miss}}, E_T^{\text{miss}} < 150$ GeV $VH + VBF$ $p_T^{j_1} \geq 200$ GeV VH hadronic (2 categories)	VH leptonic 0-jet, $p_T^{4\ell} \geq 100$ GeV 2-jet, $m_{jj} < 120$ GeV			$2 \ell, 75 \leq p_T^V < 150$ GeV, $N_{\text{jets}} = 2$ $2 \ell, 75 \leq p_T^V < 150$ GeV, $N_{\text{jets}} \geq 3$ $2 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} = 2$ $2 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} \geq 3$ $1 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} = 2$ $1 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} = 3$ $0 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} = 2$ $0 \ell, p_T^V \geq 150$ GeV, $N_{\text{jets}} = 3$	
VBF	$VBF, p_T^{\gamma\gamma jj} \geq 25$ GeV (2 categories) $VBF, p_T^{\gamma\gamma jj} < 25$ GeV (2 categories)	2-jet VBF, $p_T^{j_1} \geq 200$ GeV 2-jet VBF, $p_T^{j_1} < 200$ GeV	2-jet VBF	$VBF p_T^{\tau\tau} > 140$ GeV ($\tau_{\text{had}}\tau_{\text{had}}$ only) VBF high- m_{jj} VBF low- m_{jj}	VBF , two central jets VBF , four central jets $VBF + \gamma$	
ggF	2-jet, $p_T^{\gamma\gamma} \geq 200$ GeV 2-jet, 120 GeV $\leq p_T^{\gamma\gamma} < 200$ GeV 2-jet, 60 GeV $\leq p_T^{\gamma\gamma} < 120$ GeV 2-jet, $p_T^{\gamma\gamma} < 60$ GeV 1-jet, $p_T^{\gamma\gamma} \geq 200$ GeV 1-jet, 120 GeV $\leq p_T^{\gamma\gamma} < 200$ GeV 1-jet, 60 GeV $\leq p_T^{\gamma\gamma} < 120$ GeV 1-jet, $p_T^{\gamma\gamma} < 60$ GeV 0-jet (2 categories)	1-jet, $p_T^{4\ell} \geq 120$ GeV 1-jet, 60 GeV $\leq p_T^{4\ell} < 120$ GeV 1-jet, $p_T^{4\ell} < 60$ GeV 0-jet, $p_T^{4\ell} < 100$ GeV	1-jet, $m_{\ell\ell} < 30$ GeV, $p_T^{\ell_2} < 20$ GeV 1-jet, $m_{\ell\ell} < 30$ GeV, $p_T^{\ell_2} \geq 20$ GeV 1-jet, $m_{\ell\ell} \geq 30$ GeV, $p_T^{\ell_2} < 20$ GeV 1-jet, $m_{\ell\ell} \geq 30$ GeV, $p_T^{\ell_2} \geq 20$ GeV 0-jet, $m_{\ell\ell} < 30$ GeV, $p_T^{\ell_2} < 20$ GeV 0-jet, $m_{\ell\ell} < 30$ GeV, $p_T^{\ell_2} \geq 20$ GeV 0-jet, $m_{\ell\ell} \geq 30$ GeV, $p_T^{\ell_2} < 20$ GeV 0-jet, $m_{\ell\ell} \geq 30$ GeV, $p_T^{\ell_2} \geq 20$ GeV	Boosted, $p_T^{\tau\tau} > 140$ GeV Boosted, $p_T^{\tau\tau} \leq 140$ GeV		

$$\mu = 1.11^{+0.09}_{-0.08} = 1.11 \pm 0.05 \text{ (stat.)} {}^{+0.05}_{-0.04} \text{ (exp.)} {}^{+0.05}_{-0.04} \text{ (sig. th.)} {}^{+0.03}_{-0.03} \text{ (bkg. th.)}$$

Searches for New Physics

- **Supersymmetry**
- **Exotic models**
- **Dark Matter**
- **Unconventional signatures**

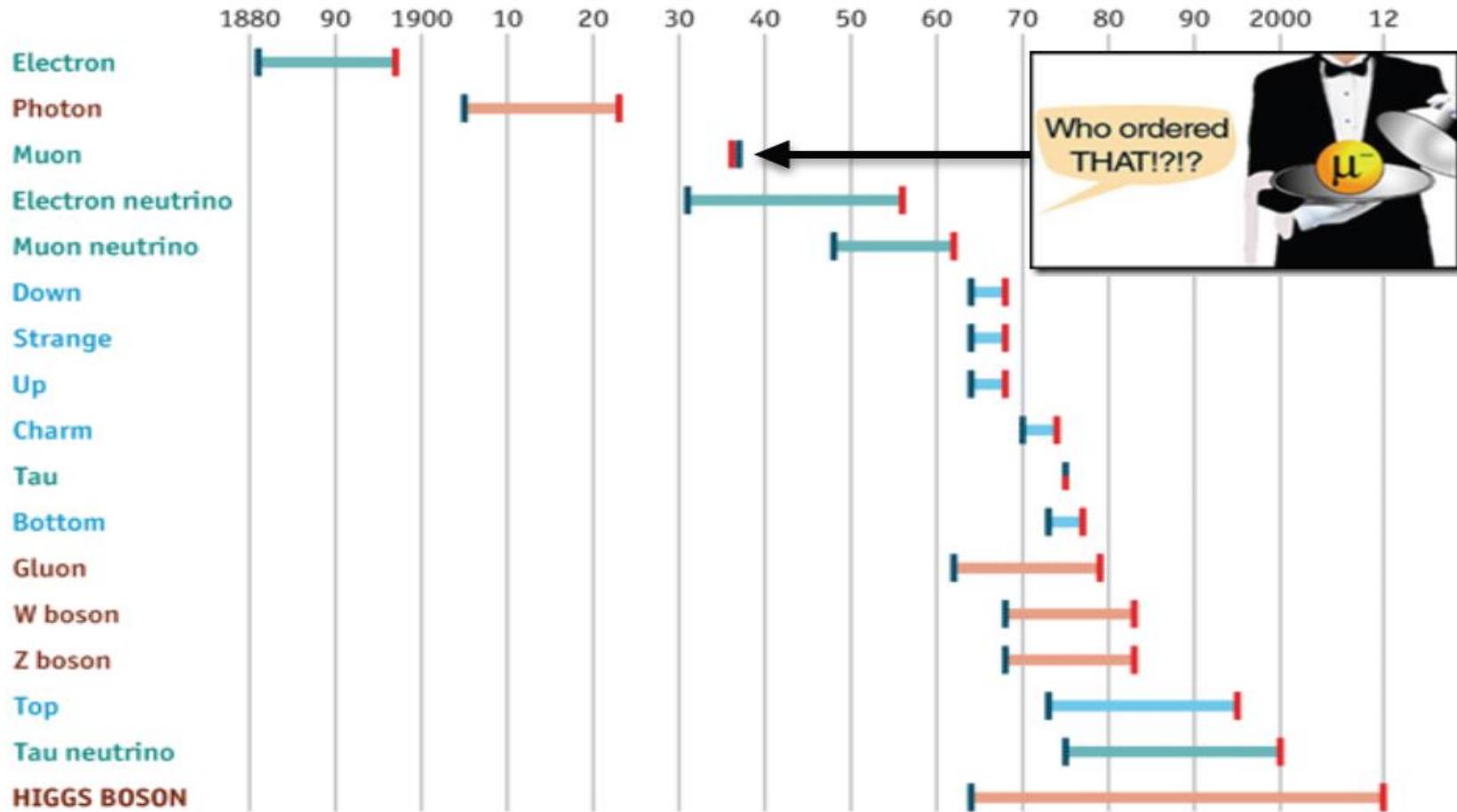
Uncharted discoveries?

The Standard Model of particle physics

Years from concept to discovery

Leptons
Bosons
Quarks

Theorised/explained
Discovered



Source: *The Economist*

Many unanswered questions ...

Why there are 3 families of particles? Are there more?

Why is the top quark so heavy?

Why there's more matter than anti-matter?

How do neutrinos get mass?

1960: SLAC u up quark	1984: Brookhaven & SLAC c charm quark	1985: Fermilab t top quark	1973: DESY g gluon
1966: SLAC d down quark	1987: Manchester University s strange quark	1977: Fermilab b bottom quark	1923: Washington University γ photon
1980: Savannah River Plant ν_e electron neutrino	1982: Brookhaven ν_μ muon neutrino	2000: Fermilab ν_τ tau neutrino	1983: CERN W W boson
1987: Cernvud Laboratory e electron	1987: Caltech and Harvard μ muon	1990: SLAC τ tau	1983: CERN Z Z boson
			2012: CERN H Higgs boson

Are there more forces?

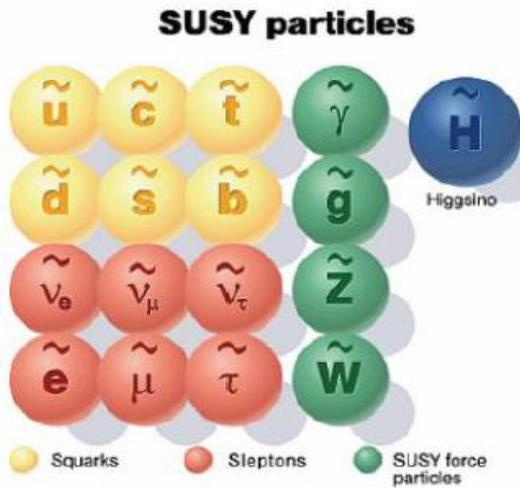
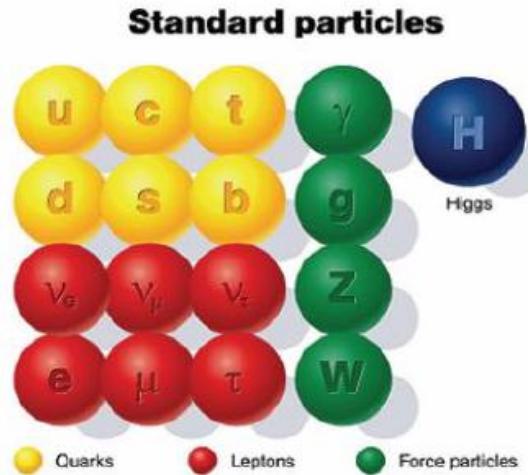
How do we incorporate gravity?

What is Dark Matter?

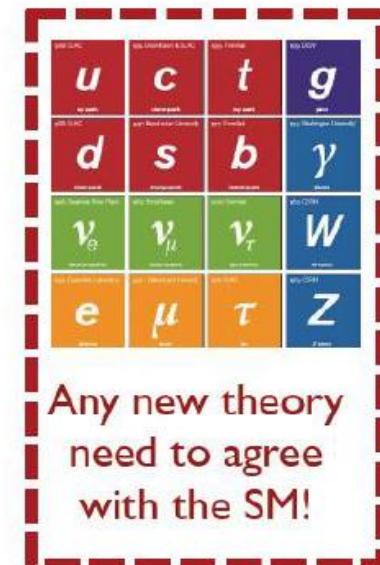
What keeps the Higgs mass so small?

... and as many possible answers to probe!

- Super-symmetry?



- Composite quark and/or leptons?
- New Heavy bosons?
- Gravitons?
- Dark Matter particles?
- ...

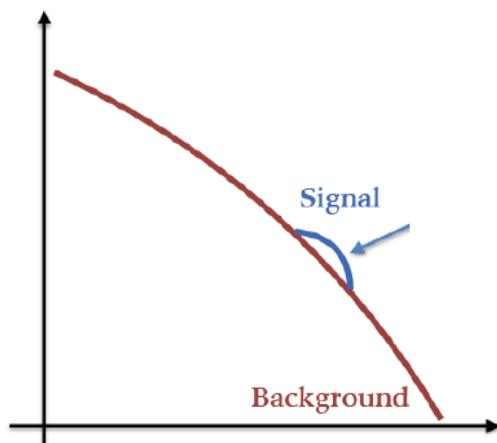


How would new phenomena manifest itself?

New particles:

resonant excess (bump) over Standard Model background

Number of events

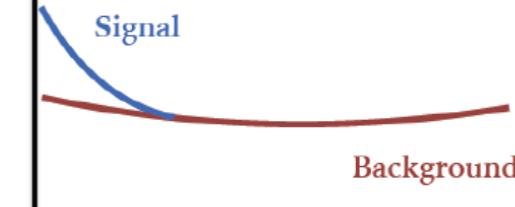


Mass of di-jet system
(~new particle mass)

New interactions:

more central production (~Rutherford experiment)

Normalized number of events

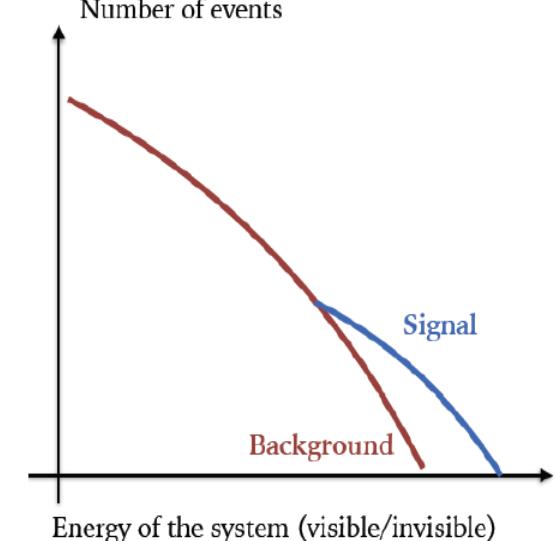


Angle between two jets

New particles and states:

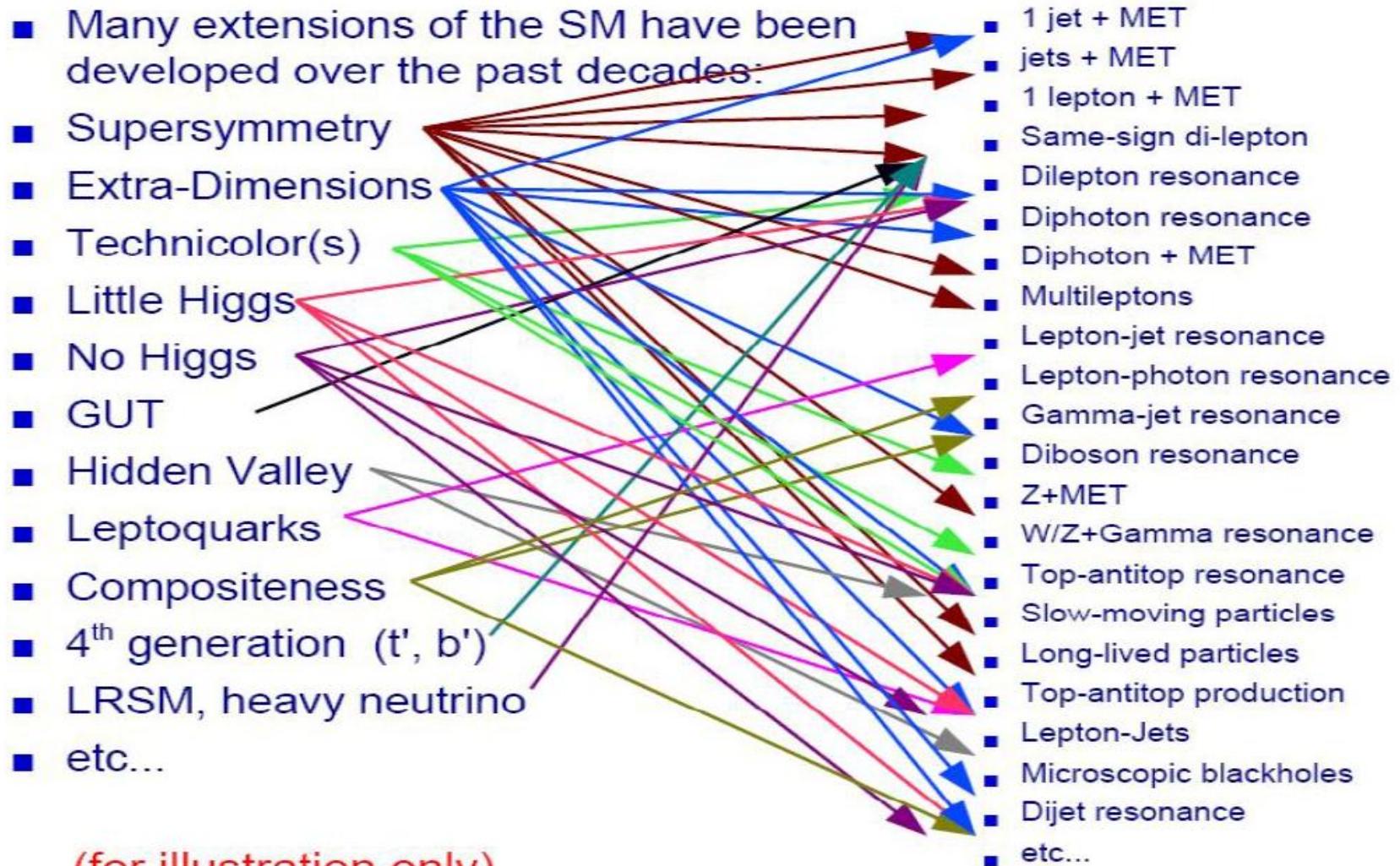
larger multiplicity of objects at high masses

Number of events



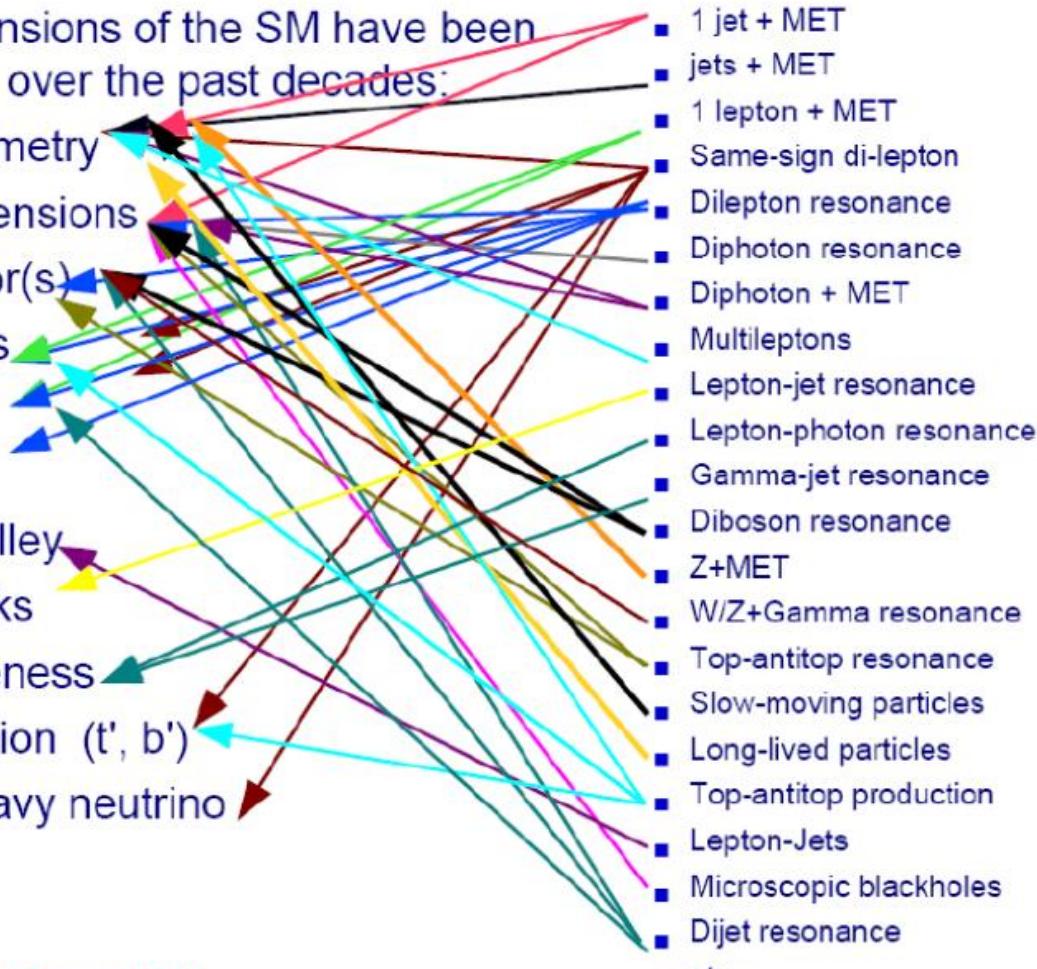
Energy of the system (visible/invisible)

Long list of models and signatures



Long list of models and signatures

- Many extensions of the SM have been developed over the past decades:
- Supersymmetry
- Extra-Dimensions
- Technicolor(s)
- Little Higgs
- No Higgs
- GUT
- Hidden Valley
- Leptoquarks
- Compositeness
- 4th generation (t', b')
- LRSM, heavy neutrino
- etc...



(for illustration only)

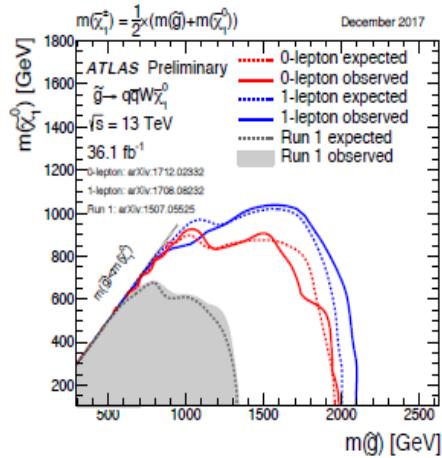
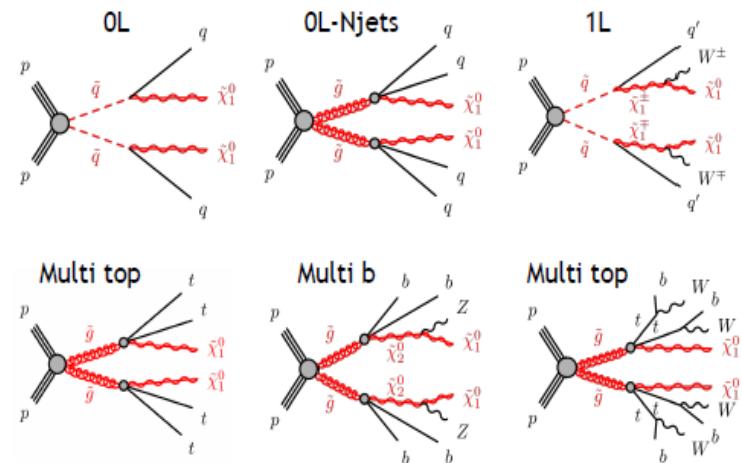
A complex 2D problem

Experimentally,
a **signature**
standpoint
makes a lot of
sense:

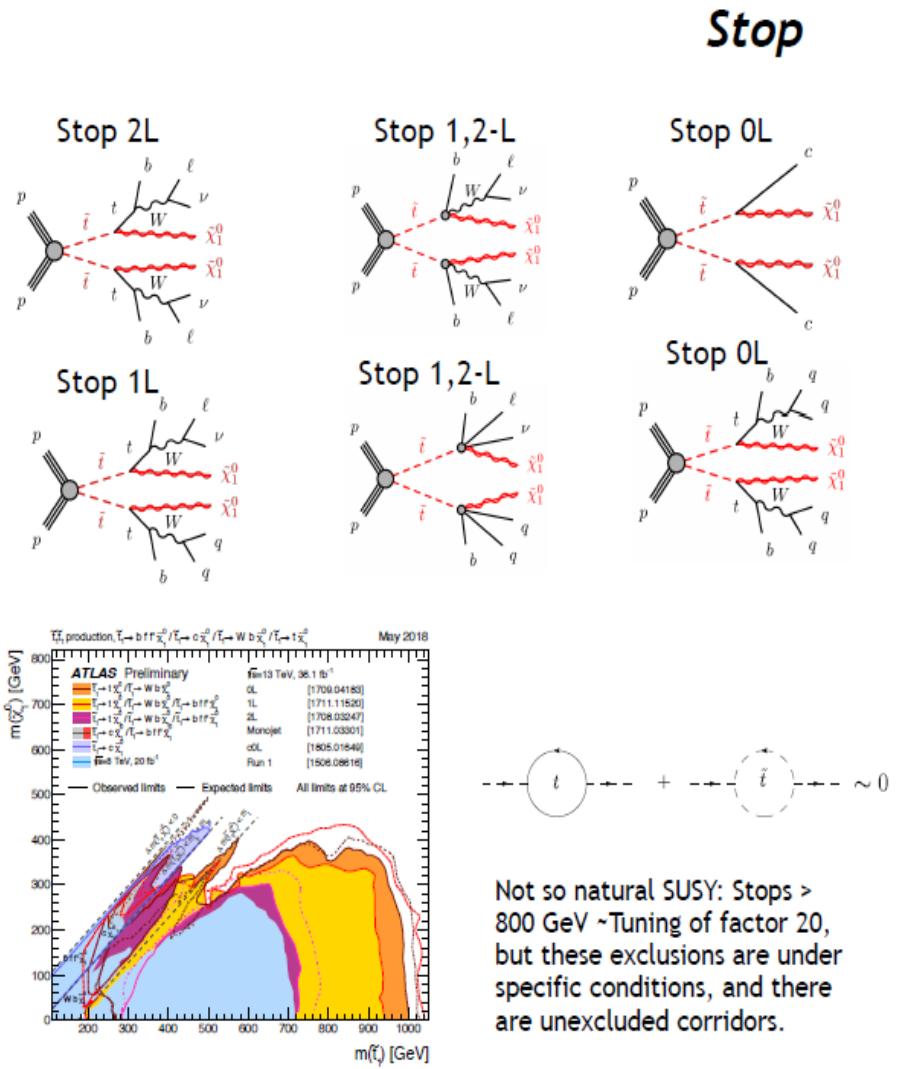
- Practical
- Less model-dependent
- Important to cover every possible signature

Strongly produced SUSY searches

Squarks and gluinos



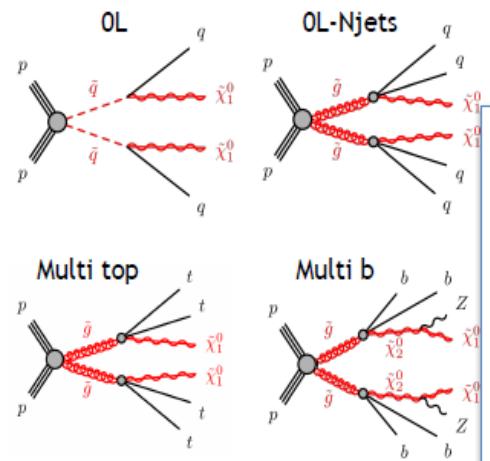
Stop also a scalar requires light gluinos to be light enough: for gluinos > 1.8 TeV ~tuning of Factor of 30



Not so natural SUSY: Stops > 800 GeV ~Tuning of factor 20, but these exclusions are under specific conditions, and there are unexcluded corridors.

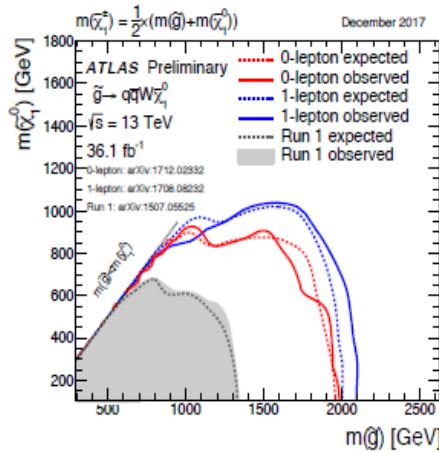
Strongly produced SUSY searches

Squarks and gluinos

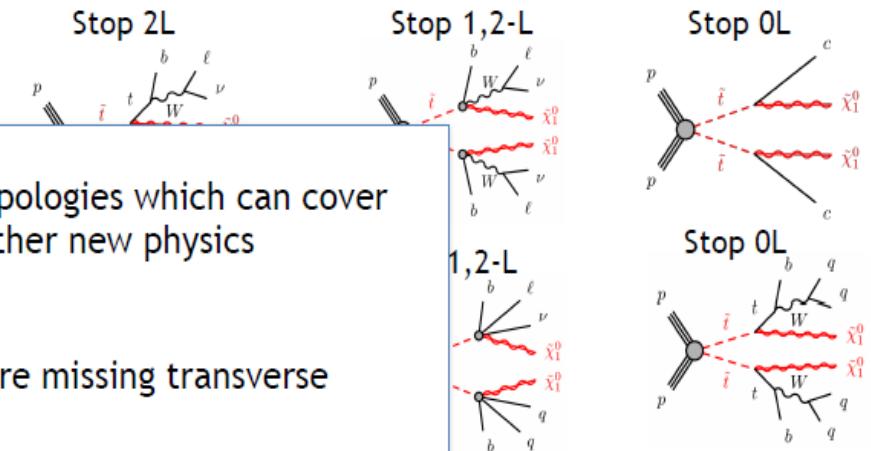
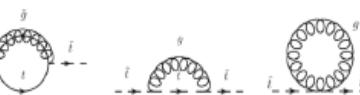


Large number of topologies which can cover different SUSY or other new physics scenarios

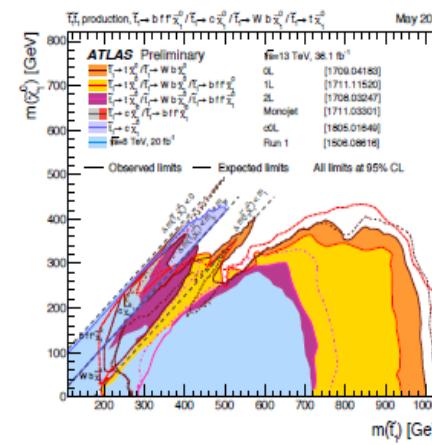
All signatures feature missing transverse energy!



Stop also a scalar requires light gluinos to be light enough: for gluinos > 1.8 TeV ~tuning of Factor of 30



$$-\cdots - t \cdots - + \cdots - \tilde{t} \cdots - \sim 0$$



Not so natural SUSY: Stops > 800 GeV ~Tuning of factor 20, but these exclusions are under specific conditions, and there are unexcluded corridors.

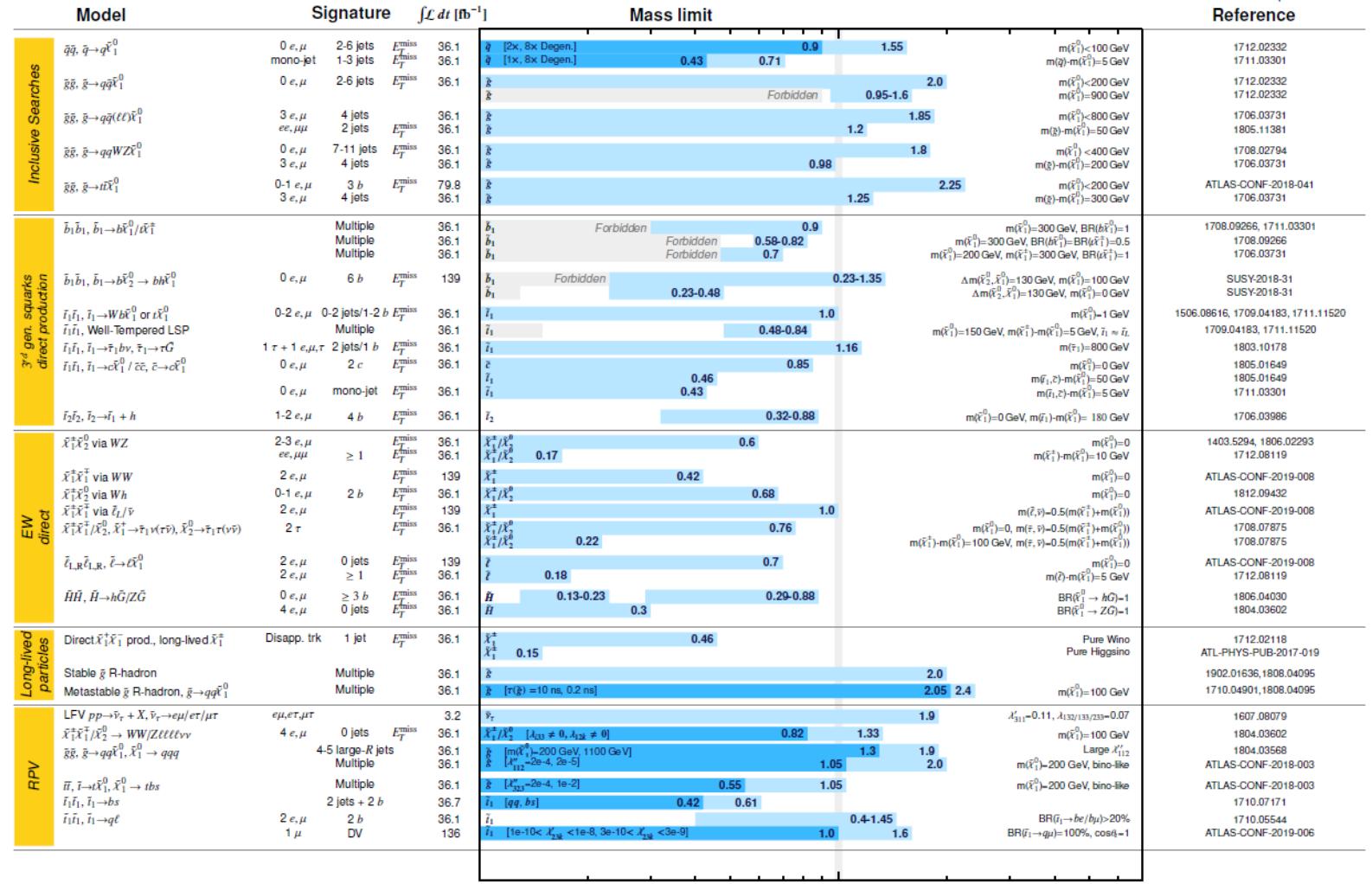
SUSY Searches Overview (2019)

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2019

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$

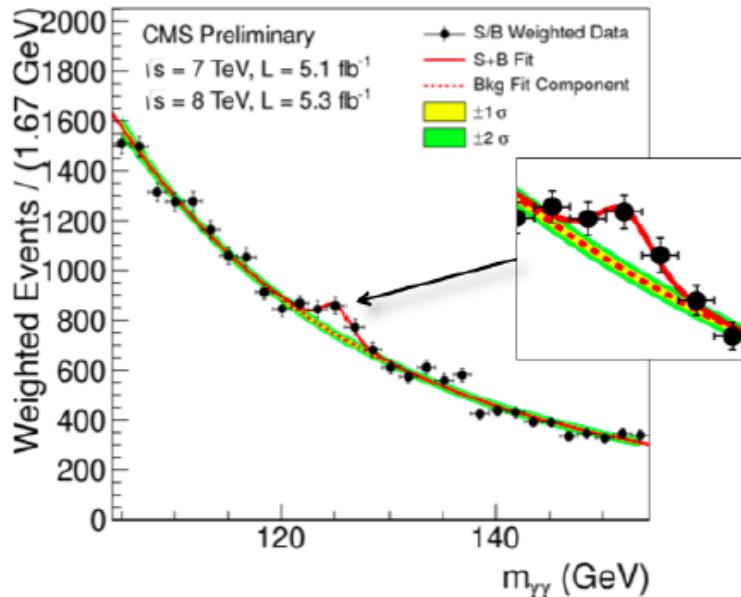


*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

What characterizes Exotics Searches

No precise model to guide us

Standard Model:
Predicted Higgs boson

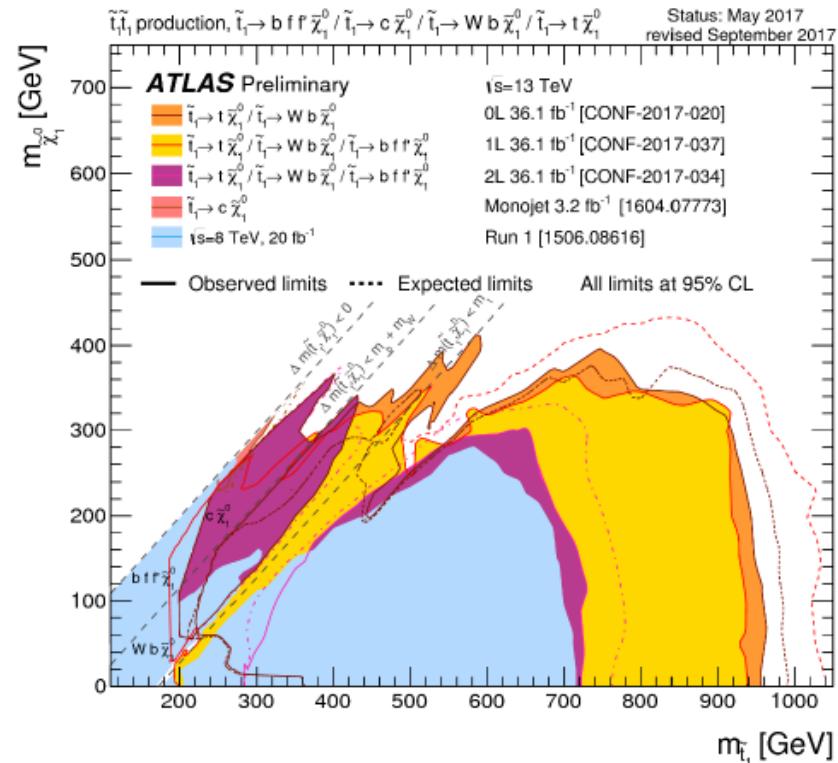


[Phys. Lett. B 716 \(2012\) 1-29](#)

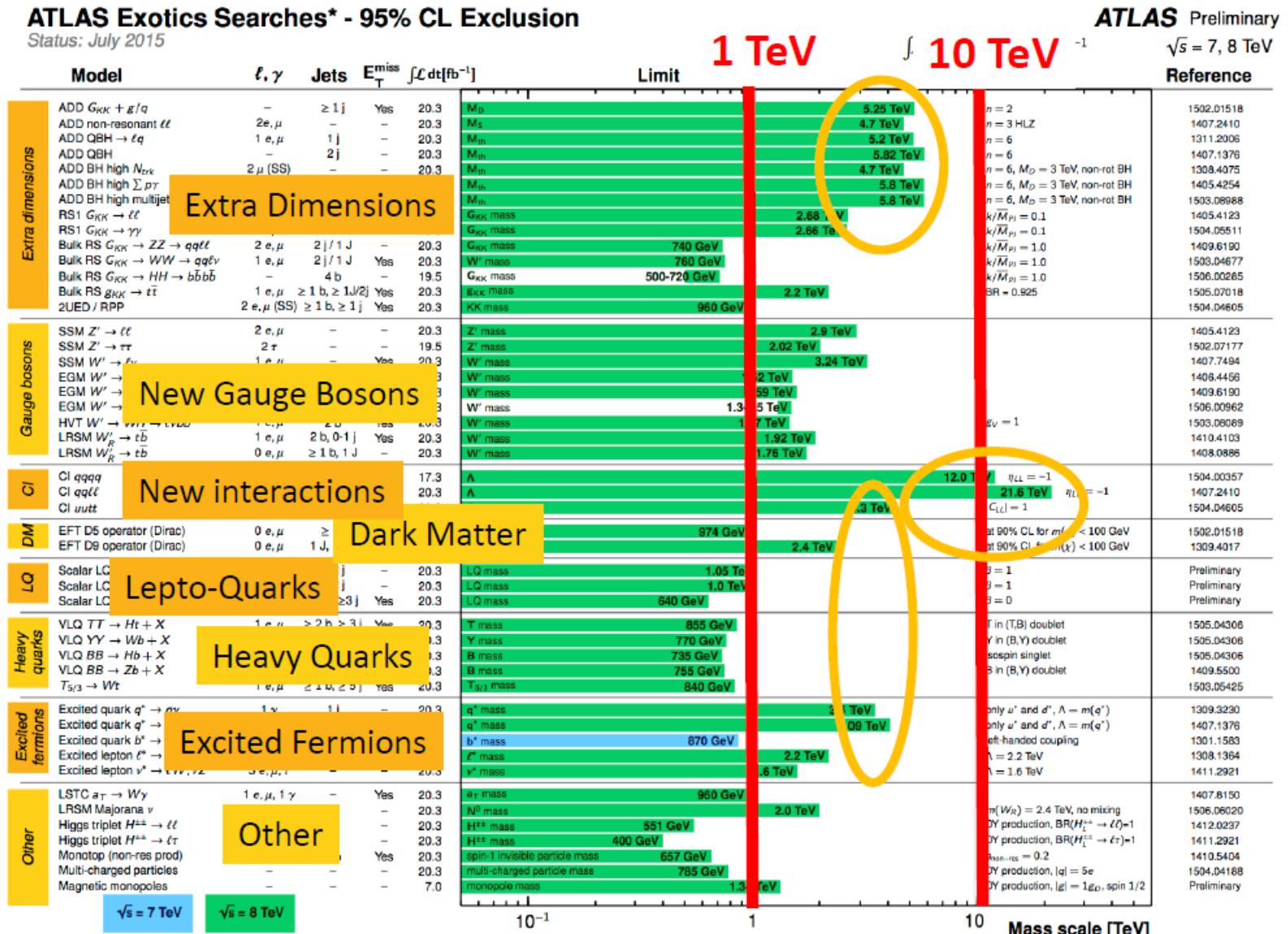
[Phys. Lett. B 716 \(2012\) 30-61](#)

No unified parameter phase space

Supersymmetry Searches:



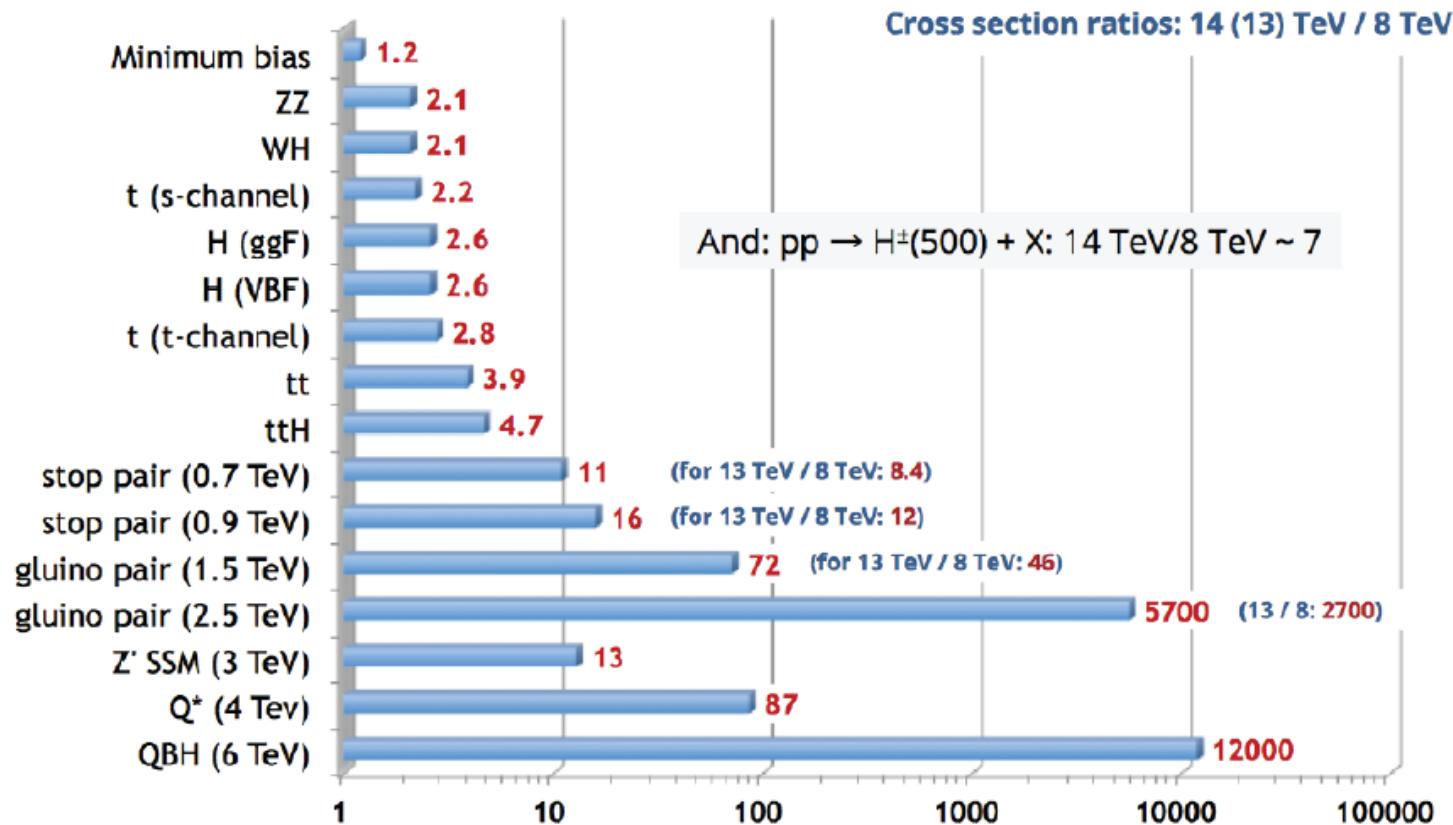
Exploration range of LHC by mid 2015



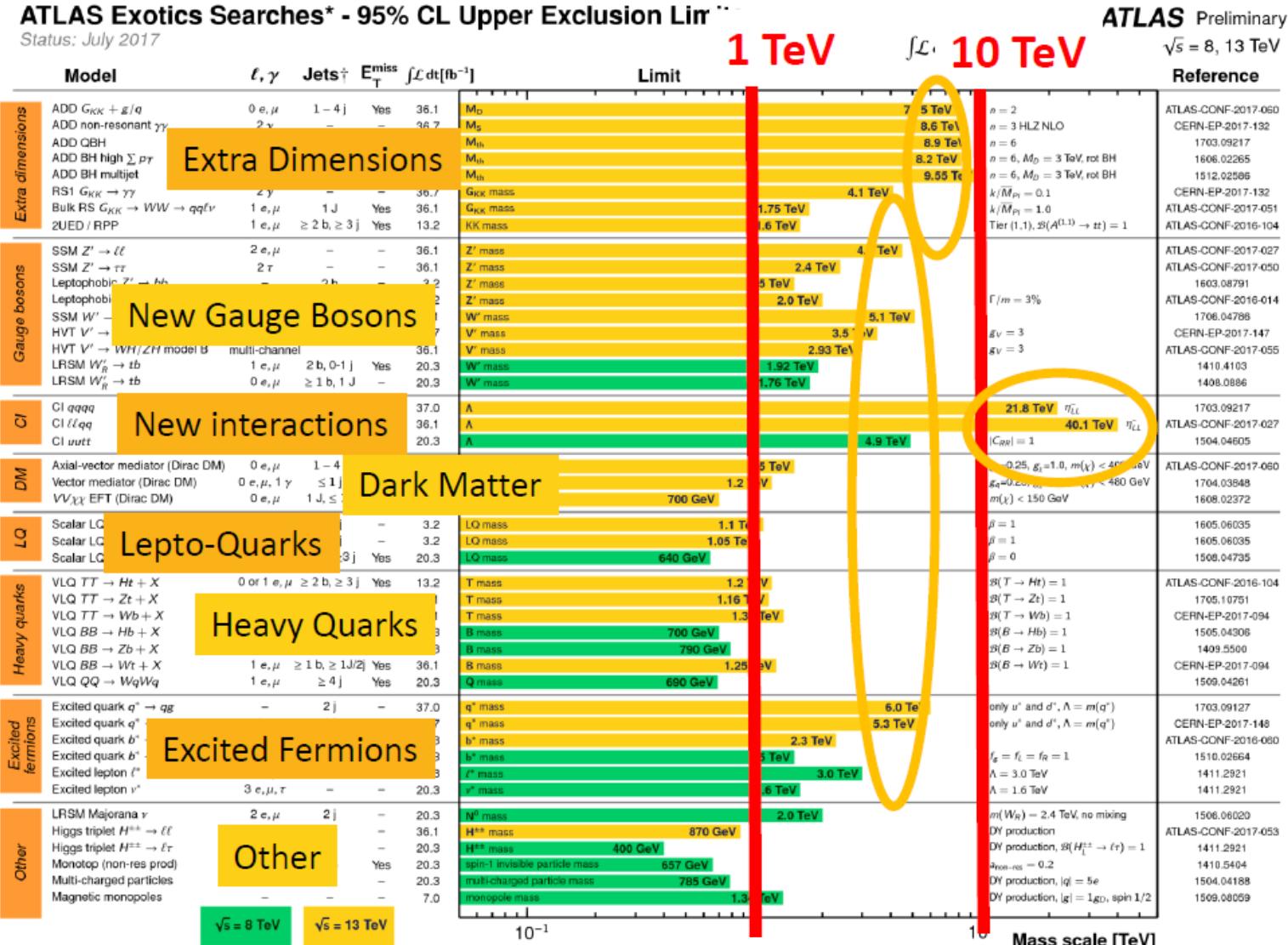
*Only a selection of the available mass limits on new states or phenomena is shown.

LHC Run II

Hugely increased potential for discovery of heavy particles at 13 TeV
Perfect occasion for young motivated physicists: join the search!



Exploration range of LHC by mid 2017



*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

Exploration range of LHC by mid 2019

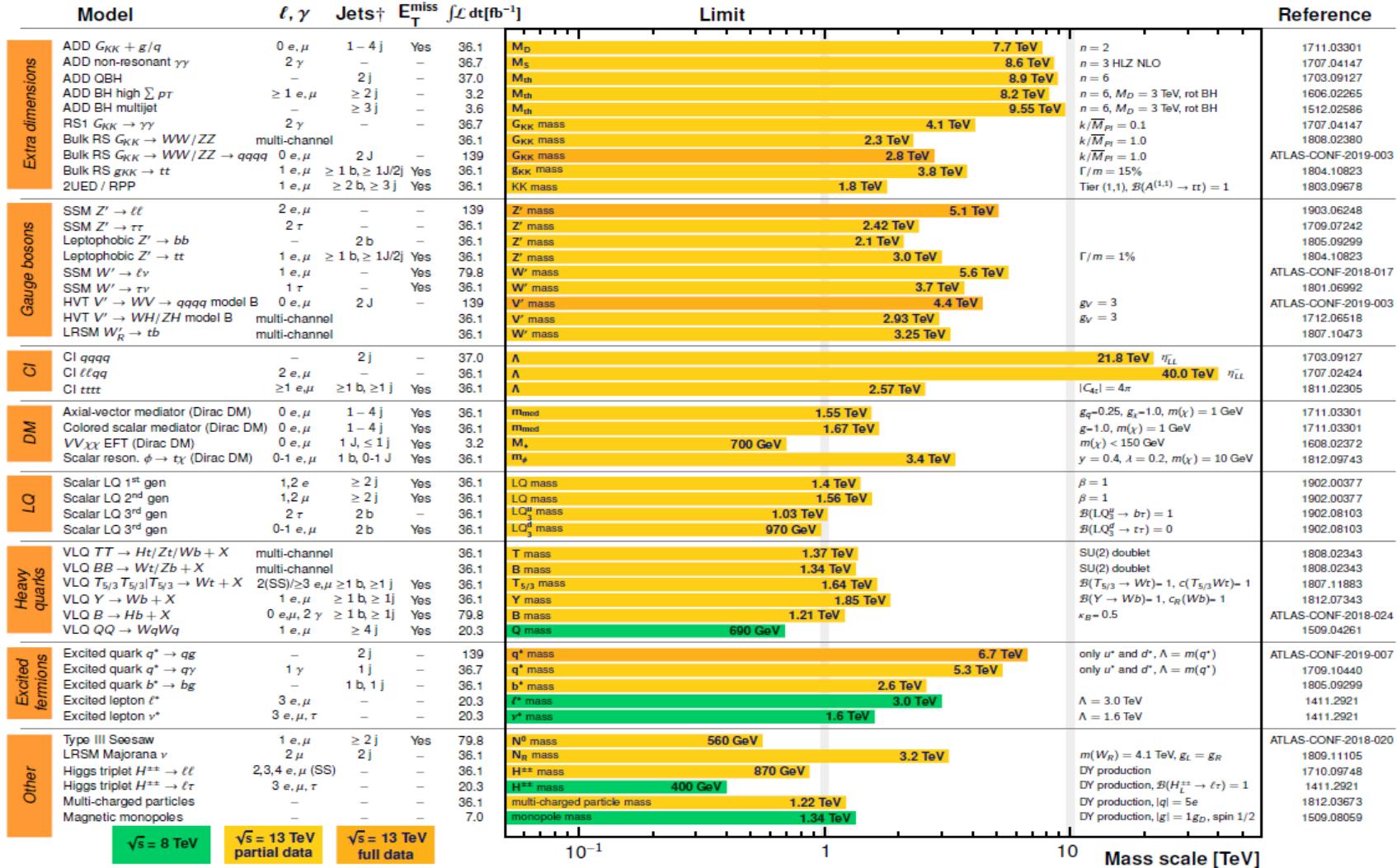
ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: March 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$



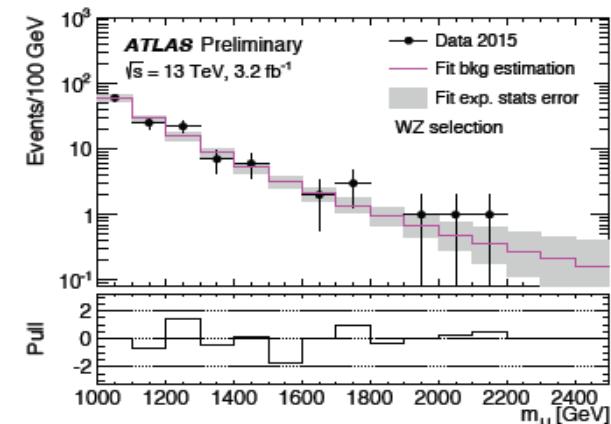
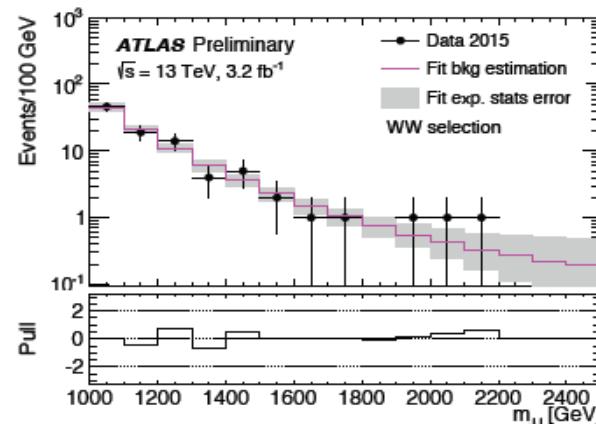
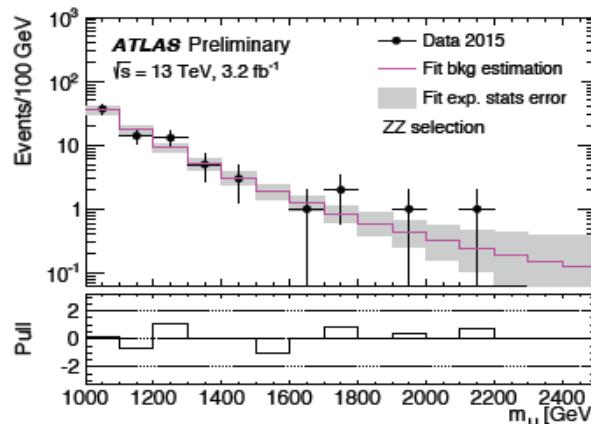
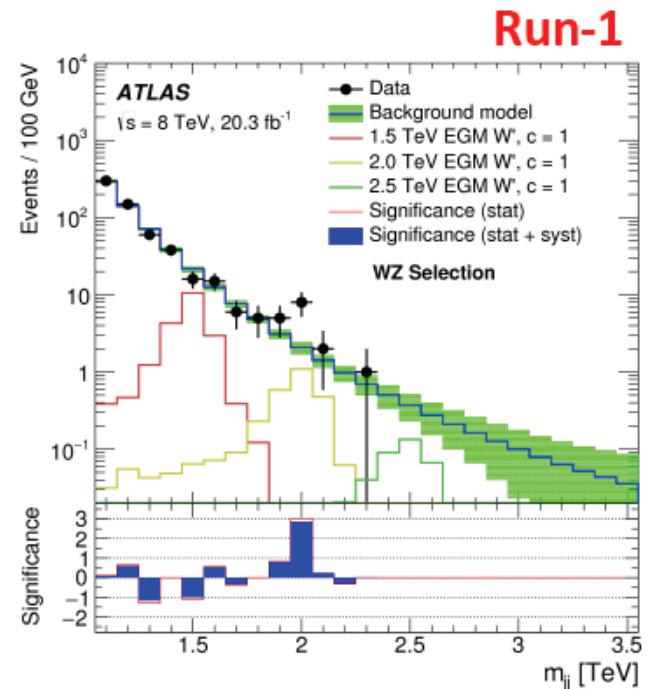
*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

Fully hadronic Diboson Searches

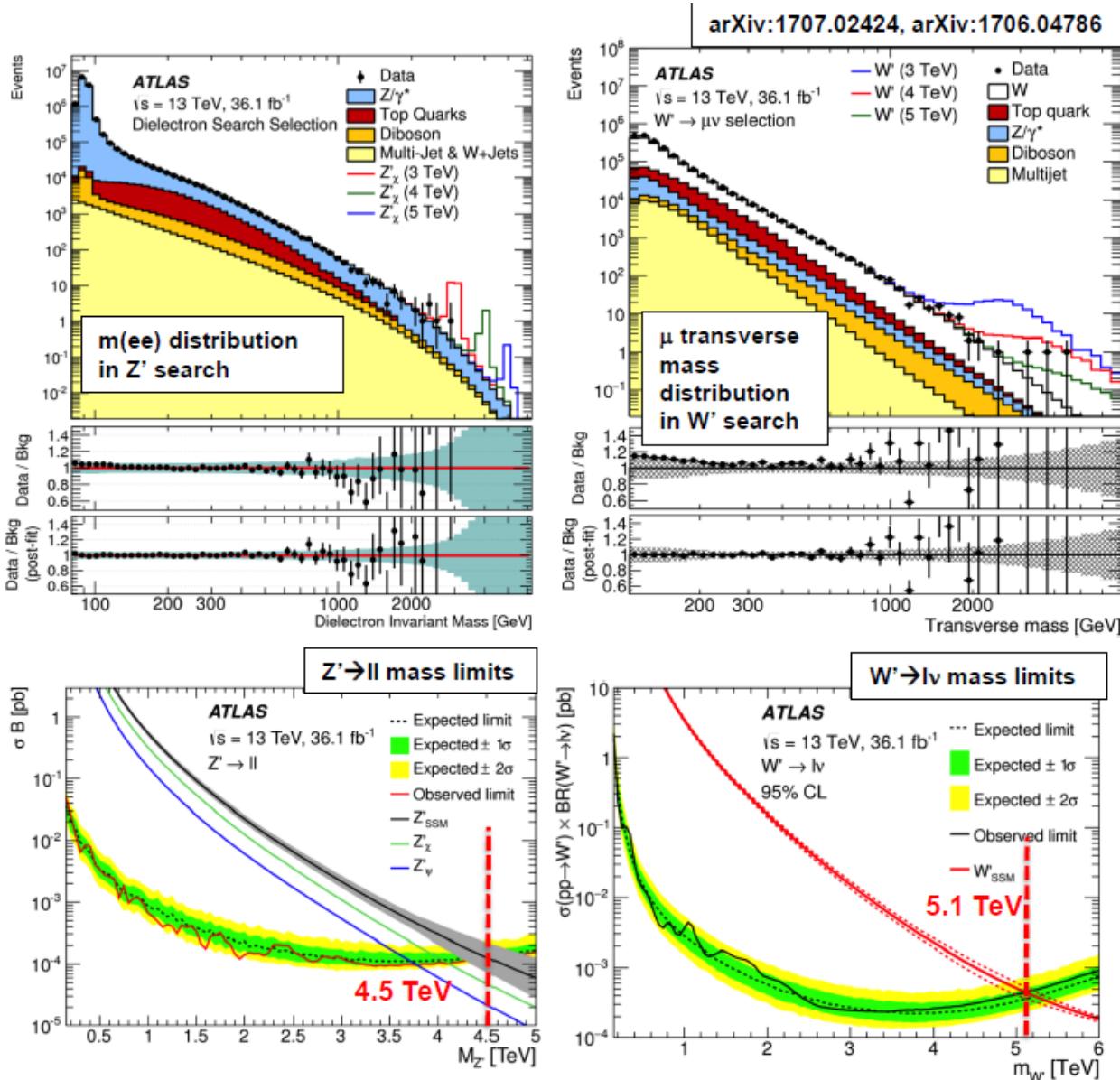
ATLAS-CONF-2015-073

- Modest excess at Run-1: 3.4σ local / 2.5σ global
- Analysis very similar to Run 1, with functional fit of the background
- No significant excess is observed
however sensitivity not high enough for conclusive probe of the Run 1 excess



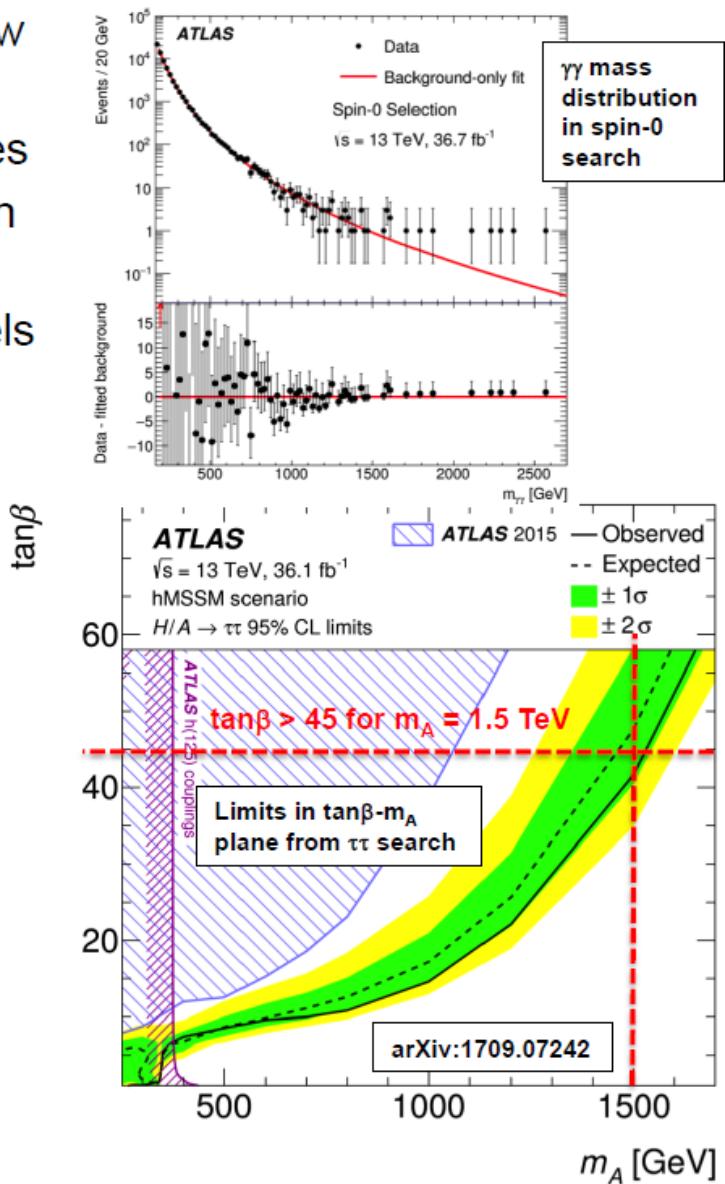
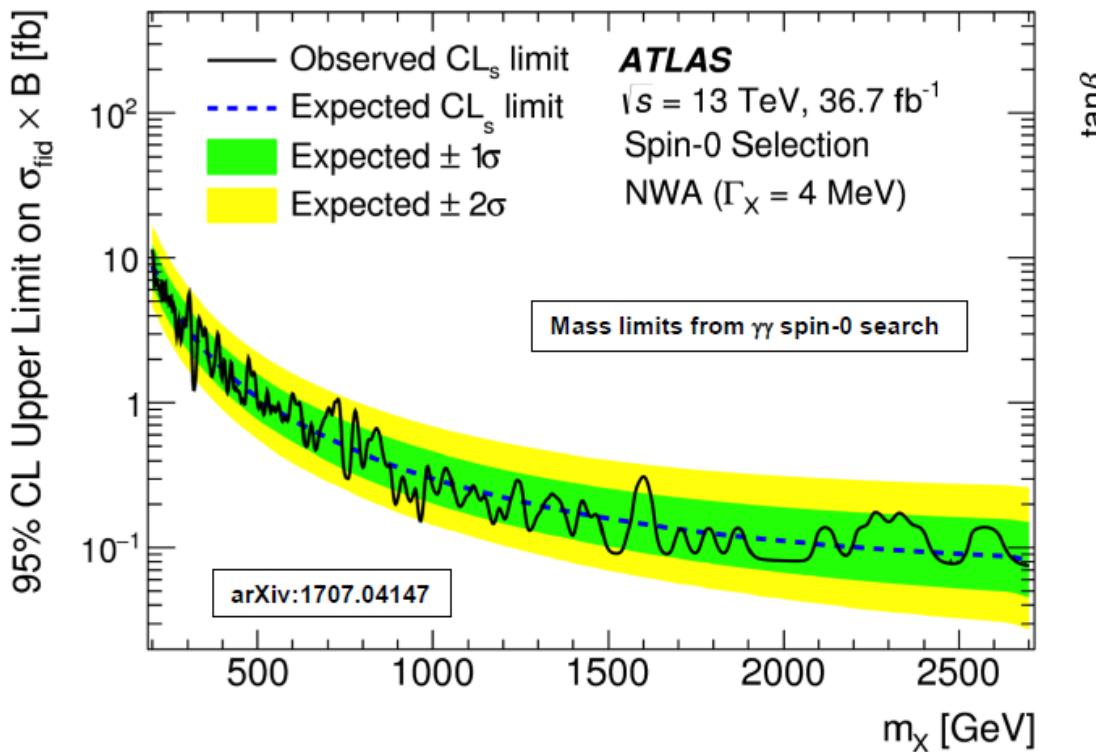
Resonance Searches (Dilepton, Lepton+ETmiss)

- Searches for new resonances decaying to lepton pairs (e.g. Z') or lepton+ E_T^{miss} (e.g. W')
- Signature is peak in invariant mass distribution (dilepton) or transverse mass distributions (lepton+ E_T^{miss})
- No significant excess over SM expectation
- 95% CL exclusion limits extracted in various new physics Z' and W' scenarios, e.g. the Sequential Standard Model (SSM)

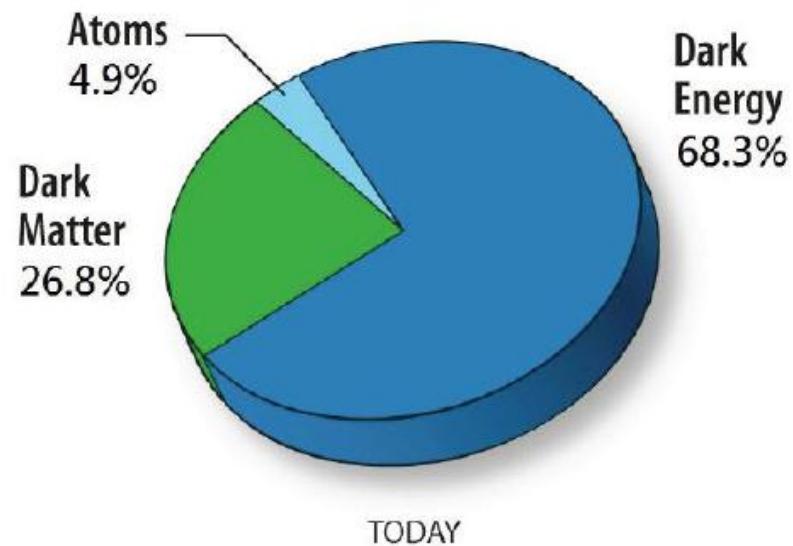
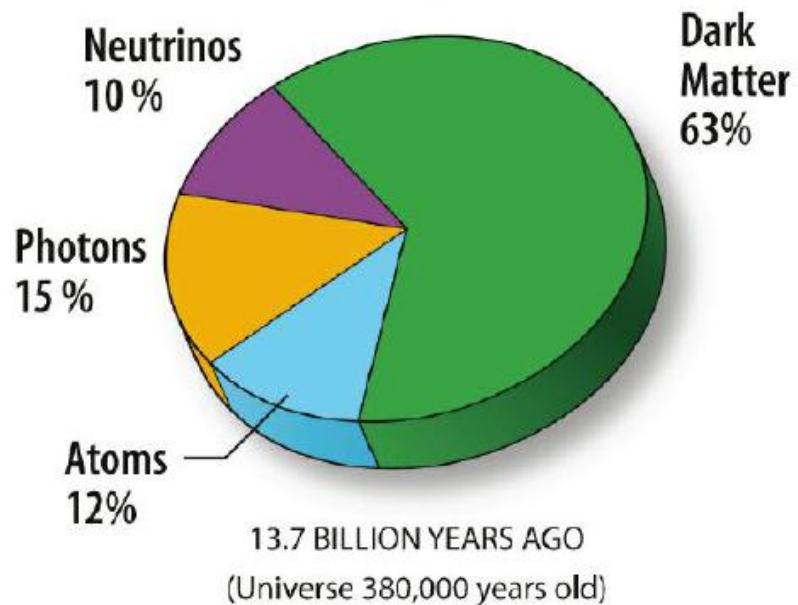


Resonance Searches ($\gamma\gamma$, $\tau\tau$)

- Diboson resonance searches also sensitive to new heavy scalars, e.g. Higgs bosons.
- Searches also conducted with $\gamma\gamma$ and $\tau\tau$ final states
- $\gamma\gamma$ search also targets spin-2 (graviton) production with a dedicated selection
- $\tau\tau$ searches sensitive to SUSY Higgs (H/A) models
- No significant excesses over SM expectation



What do we know about Dark Matter

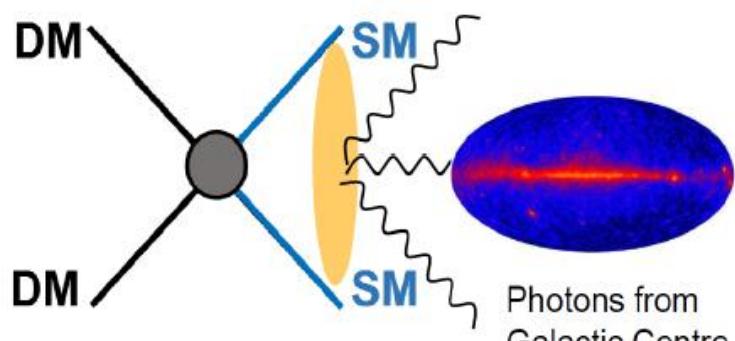


Strong astrophysical evidence for the existence of dark matter

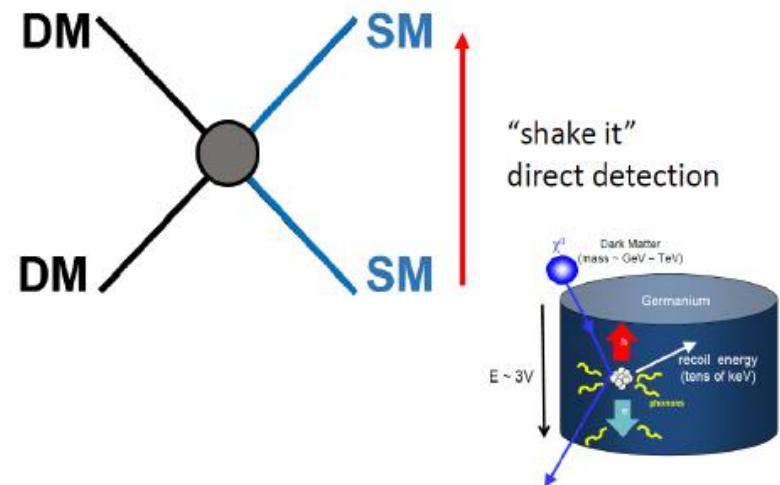
What do we know about Dark Matter

- Massive
- Non-relativistic (slow)
- Long lived (old)
- No electric or colour charge
- Very weakly interacting with ordinary matter
- Subject to gravity interactions

Experimental detection of Dark Matter

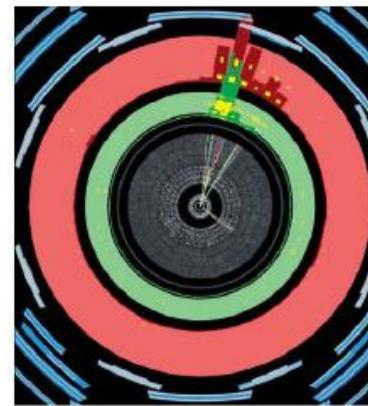
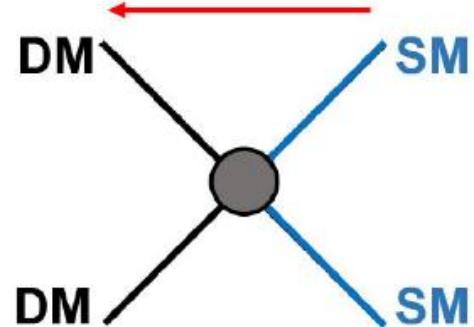


"break it": indirect detection

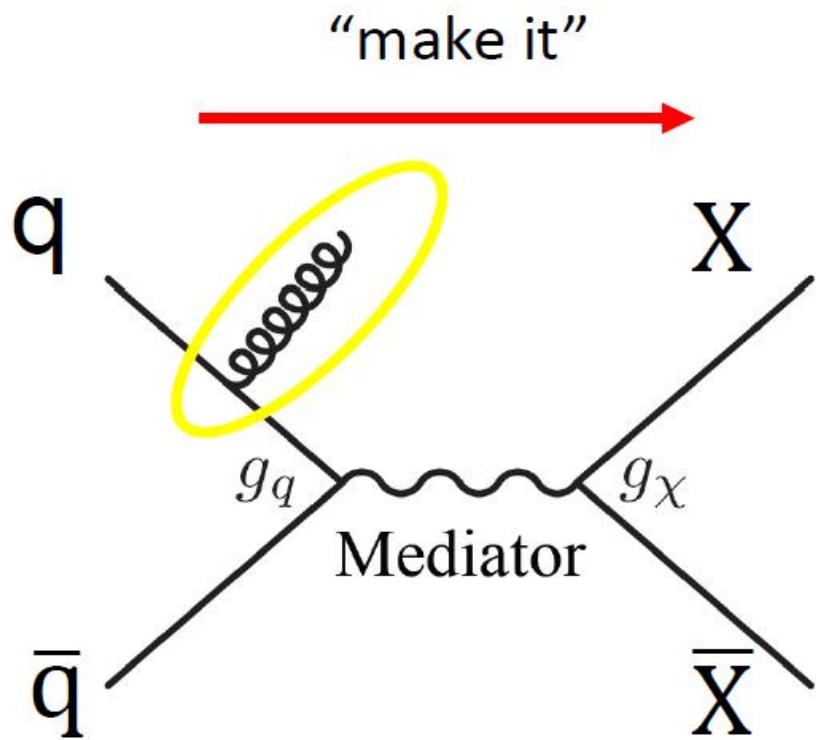


"shake it"
direct detection

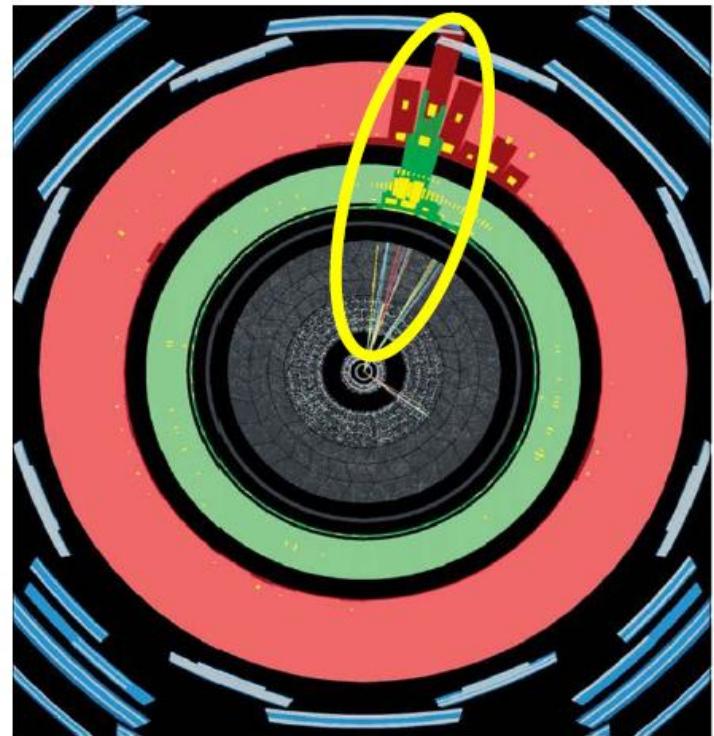
"make it": Collider Production



Dark Matter searches at Colliders

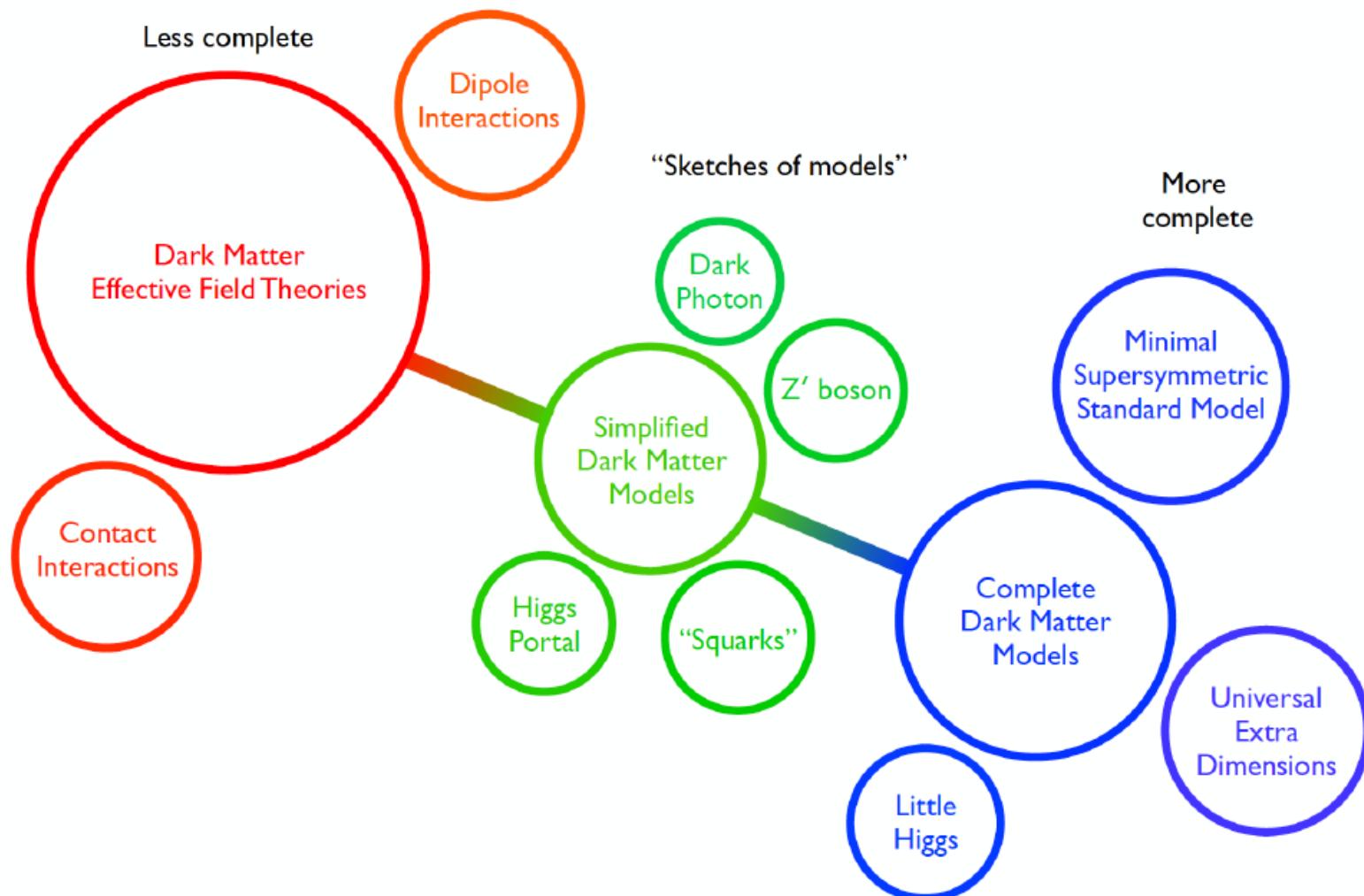


g_q and g_χ coupling strengths

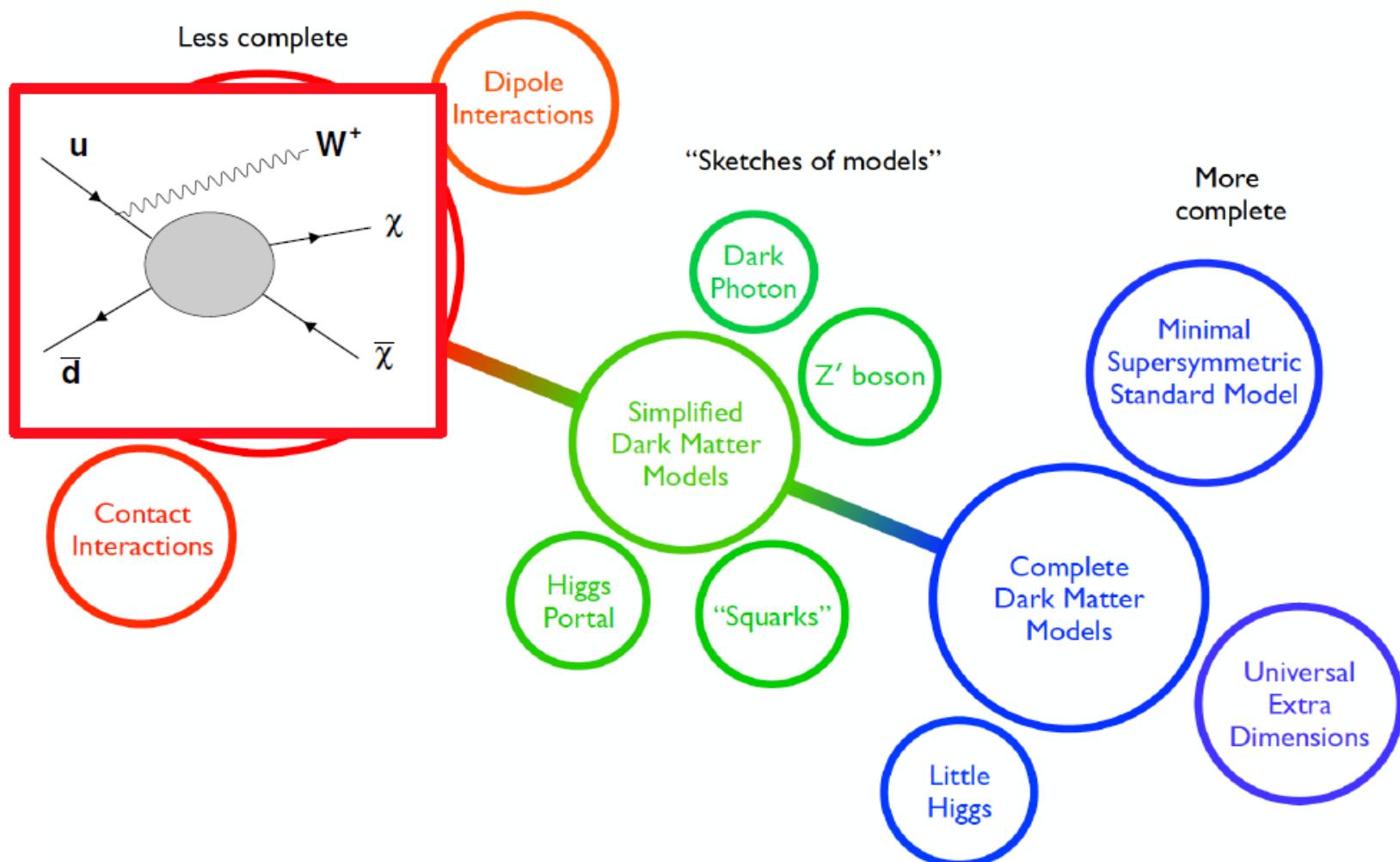


Empty detector + something

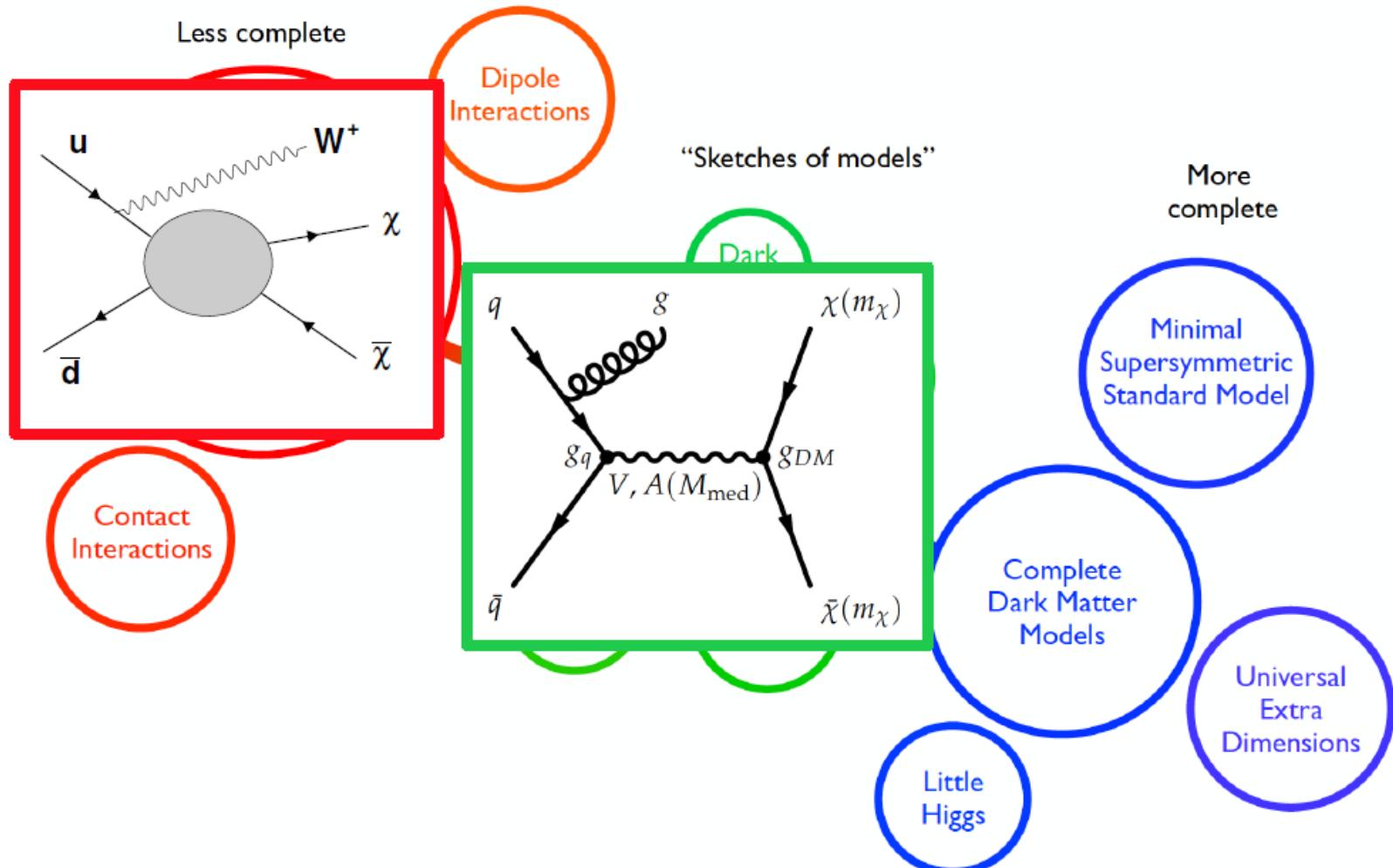
Dark Matter theory space



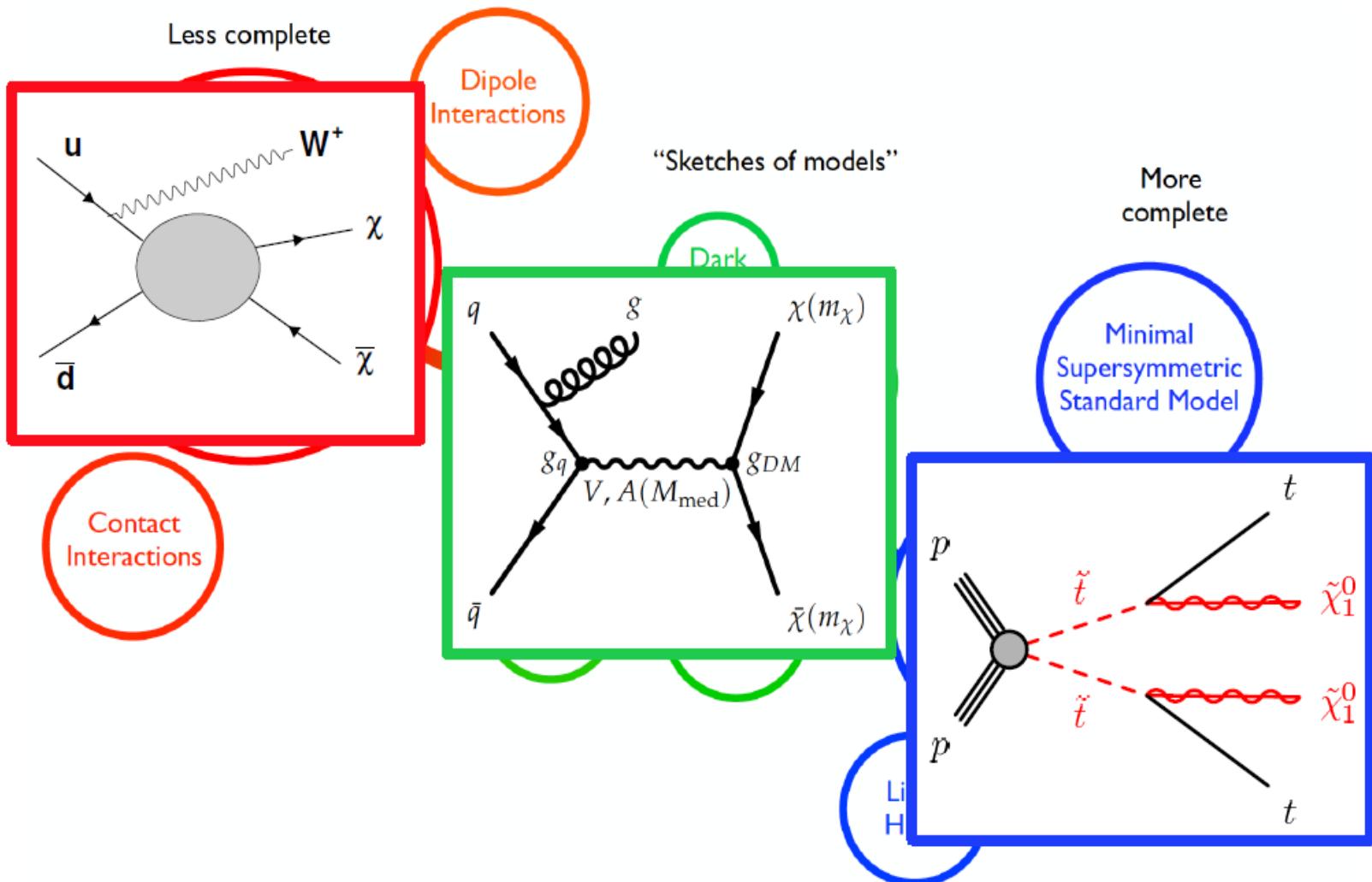
Dark Matter theory space



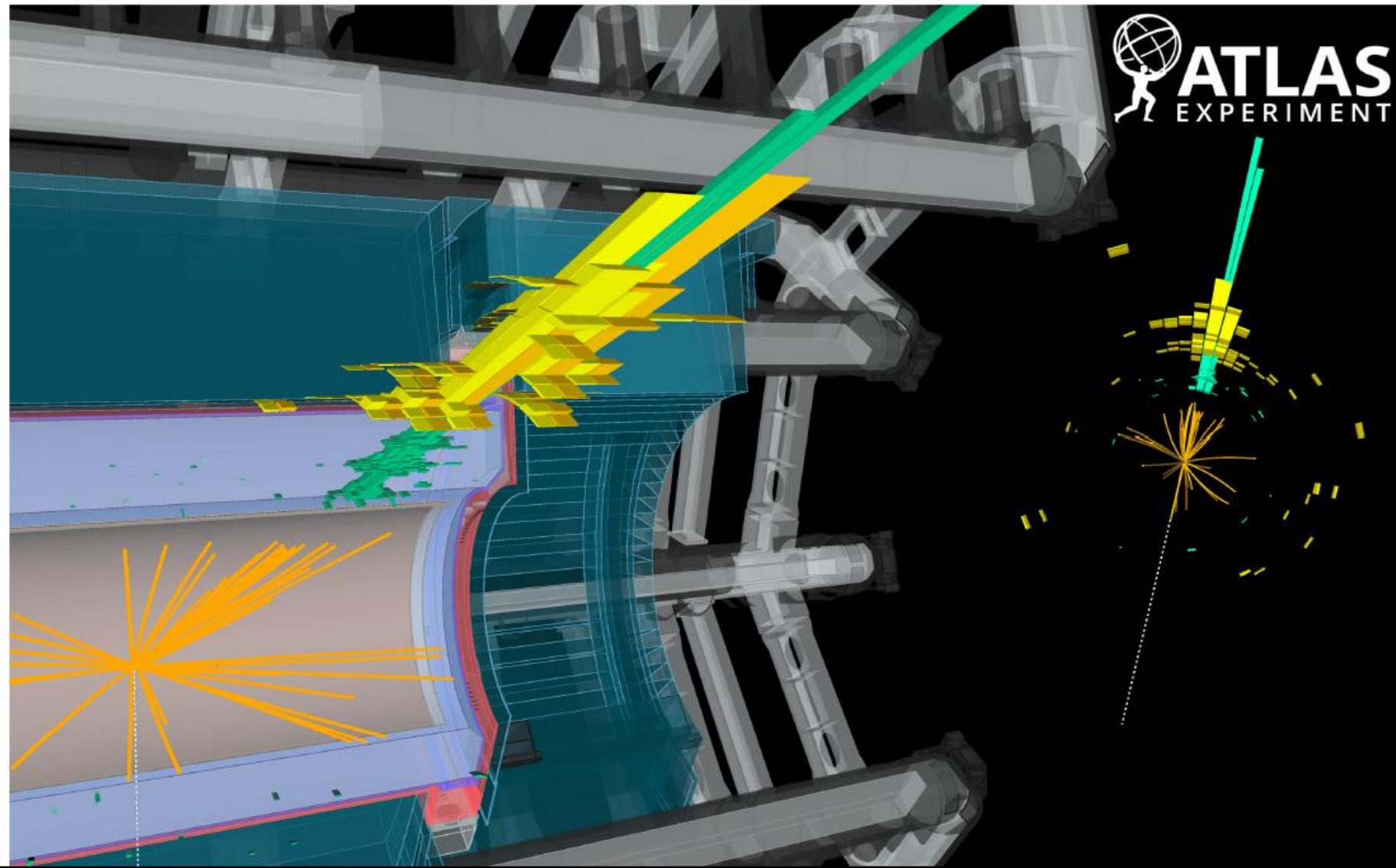
Dark Matter theory space



Dark Matter theory space

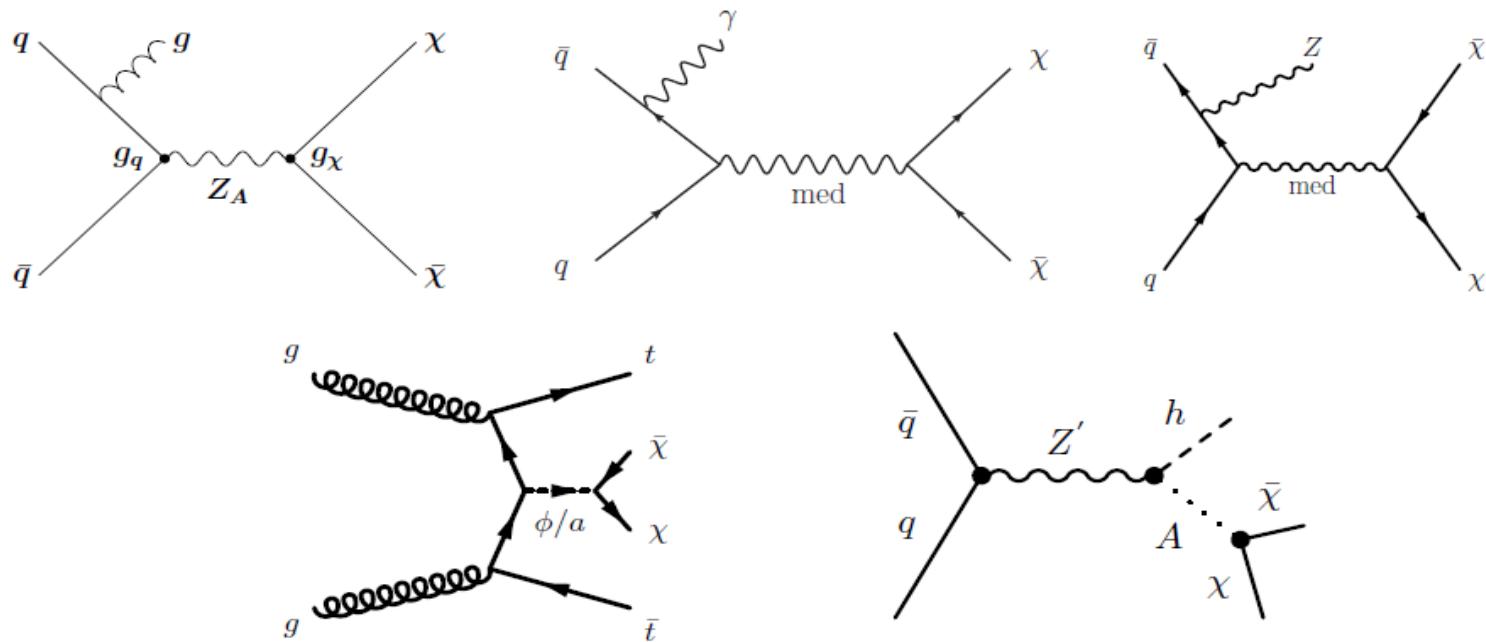


Searches for DM with ($E_T^{\text{miss}} + X$) Signatures



$(E_T^{\text{miss}} + X)$ or Mono-X Signatures

SM-DM mediator decays to DM pairs $\chi\bar{\chi}$ when $M_{\text{med}}/2 > m_\chi$.
DM escapes detection $\Rightarrow E_T^{\text{miss}} + X$ signature where
 $X = \text{SM particles that tag the event}, X = \text{jet}, \gamma, V, t, b, h \dots$

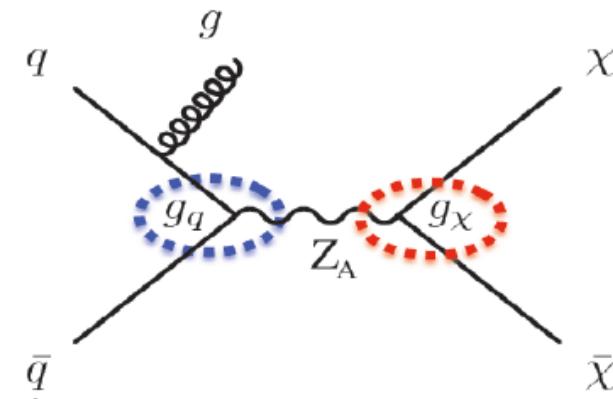


Mediators: vector, axial-vector, scalar, pseudoscalar
Parameters: $m_{\text{med}}, m_\chi, g_q, g_\chi$

Simplified Model

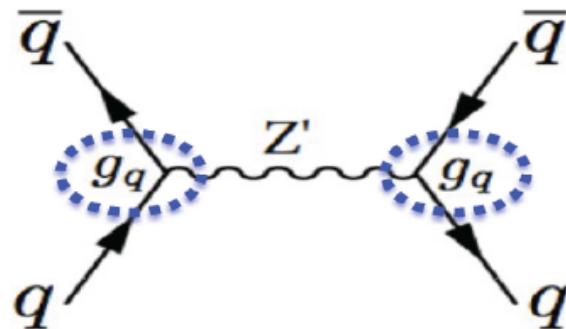
Simplified Model

$SM \rightarrow \text{mediator} \rightarrow DM$



\rightarrow Mono-X signature
 E_T^{miss} + jet, W/Z/H, γ , ...

$SM \rightarrow \text{mediator} \rightarrow SM$

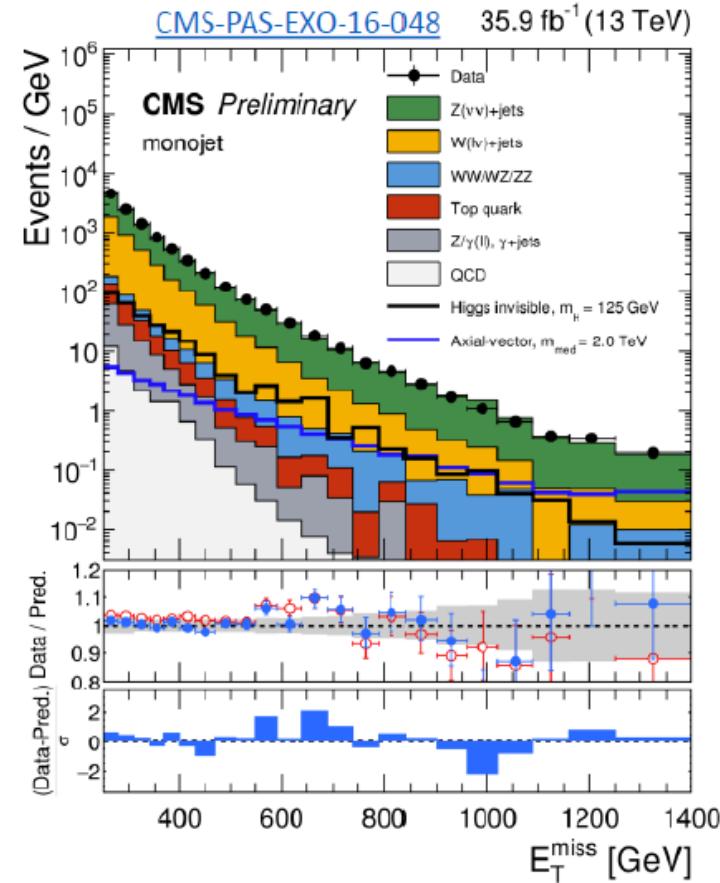
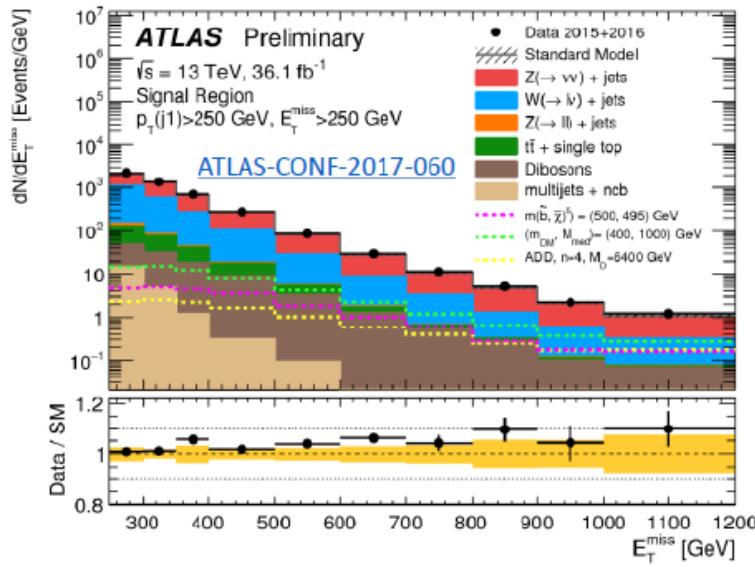


\rightarrow resonant production
Dijet, ditop, dilepton.....

	spin 0	spin 1
Charge	Q=0 for s-channel	
Lorentz structure	Scalar	Vector
	$g_q \frac{\phi}{\sqrt{2}} \sum_f y_f \bar{f} f$	$g_q \sum_q V_\mu \bar{q} \gamma^\mu q$
Coupling	Pseudoscalar	Axial-vector
	$g_q \frac{iA}{\sqrt{2}} \sum_f y_f \bar{f} \gamma^5 f$	$g_q \sum_q A_\mu \bar{q} \gamma^\mu \gamma^5 q$
	\propto mass	\propto charge

Mono-X searches

Mono-jet



ATLAS

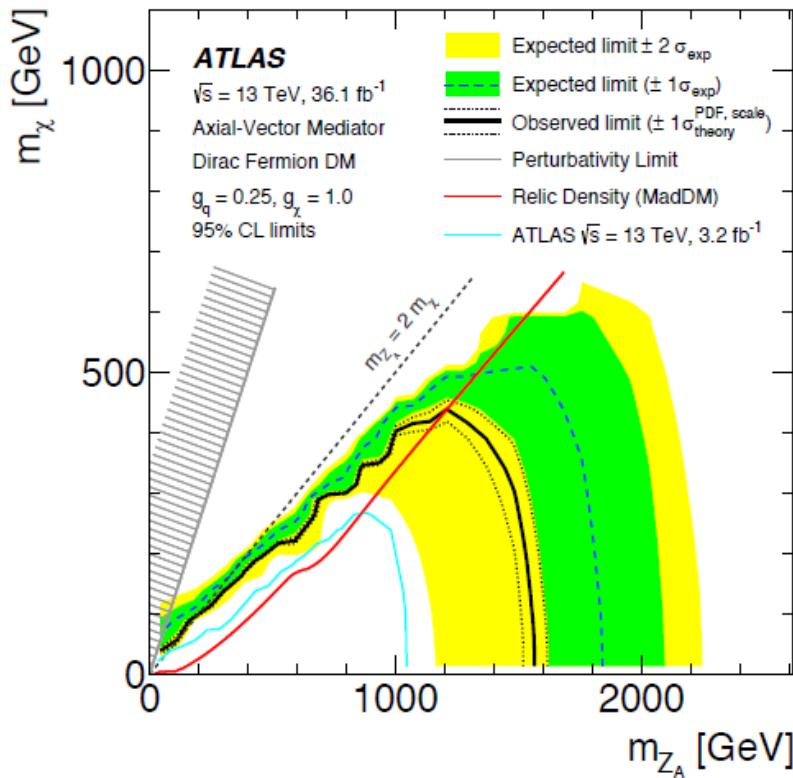
- $E_T^{\text{miss}} > 250 \text{ GeV}, \Delta\phi(\text{jet}, p_T^{\text{miss}}) > 0.4$
- $\text{Jet } p_T > 250 \text{ GeV}, |\eta| < 2.4$
- $N_{\text{jets}} \leq 4$

CMS

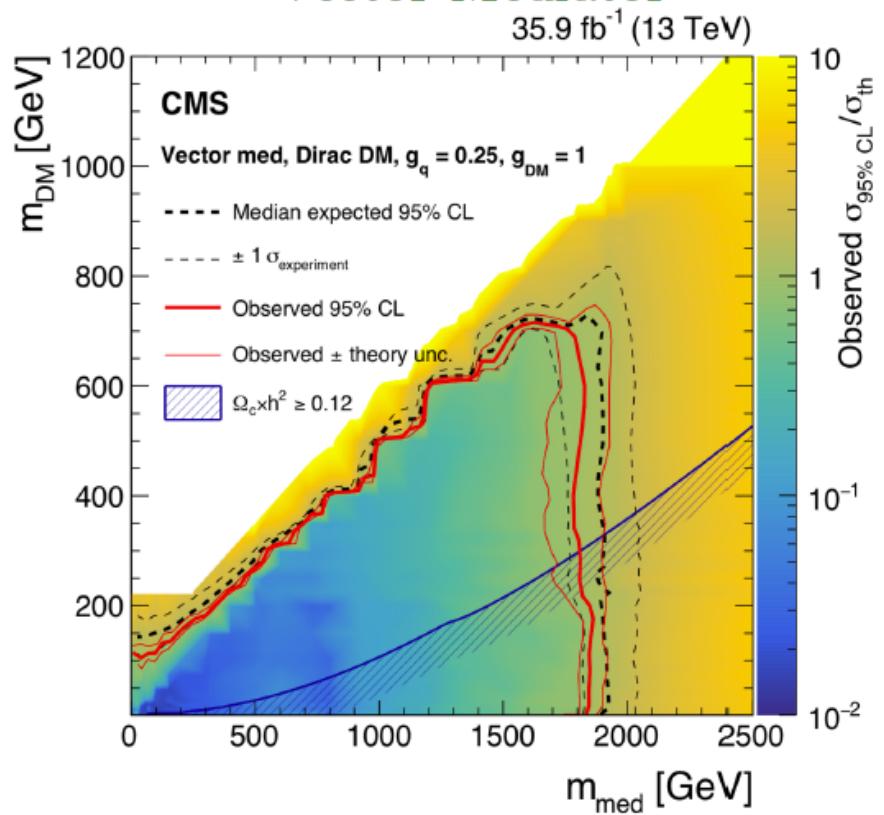
- $E_T^{\text{miss}} > 250 \text{ GeV}$
- $\text{Jet } p_T > 100 \text{ GeV}, |\eta| < 2.5$

Mono-X searches

Axial-Vector Mediator



Vector Mediator

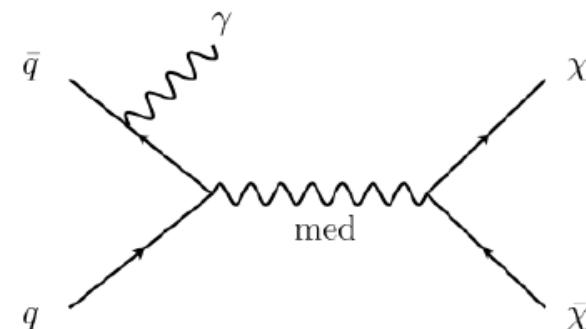
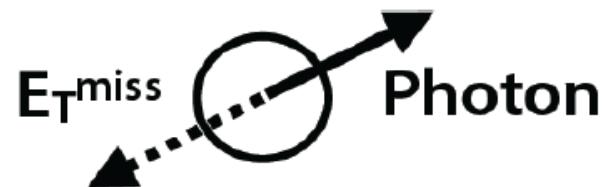
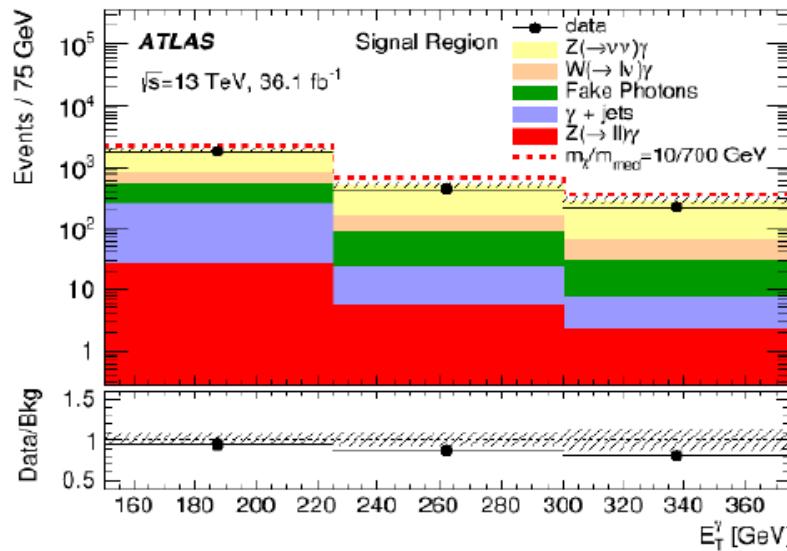


For couplings $g_q = 0.25$, $g_\chi = 1.0$, axial-vector and vector mediators excluded up to 1.8 TeV (1.55 TeV) by CMS (ATLAS) for $m_\chi \sim 1 \text{ GeV}$.

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Mono-X searches

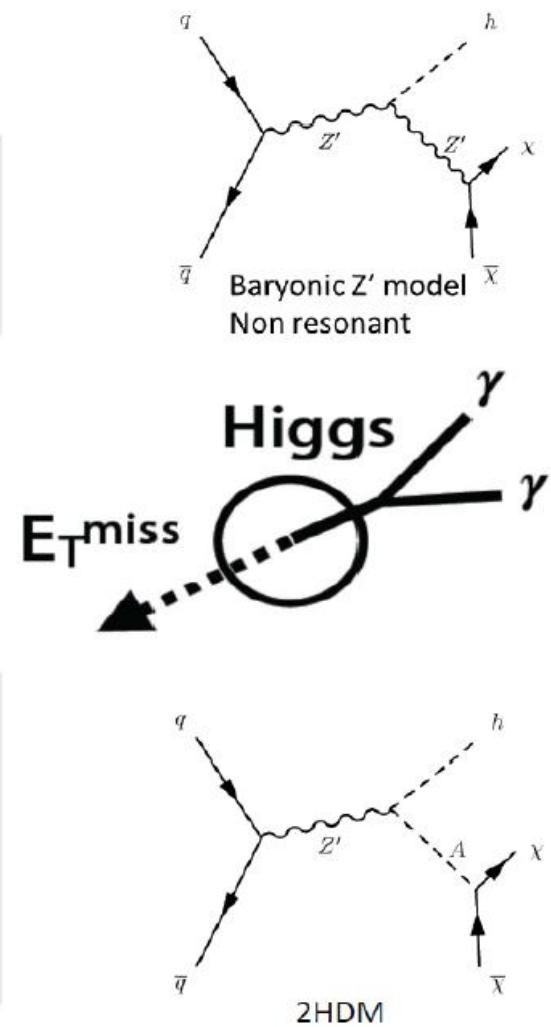
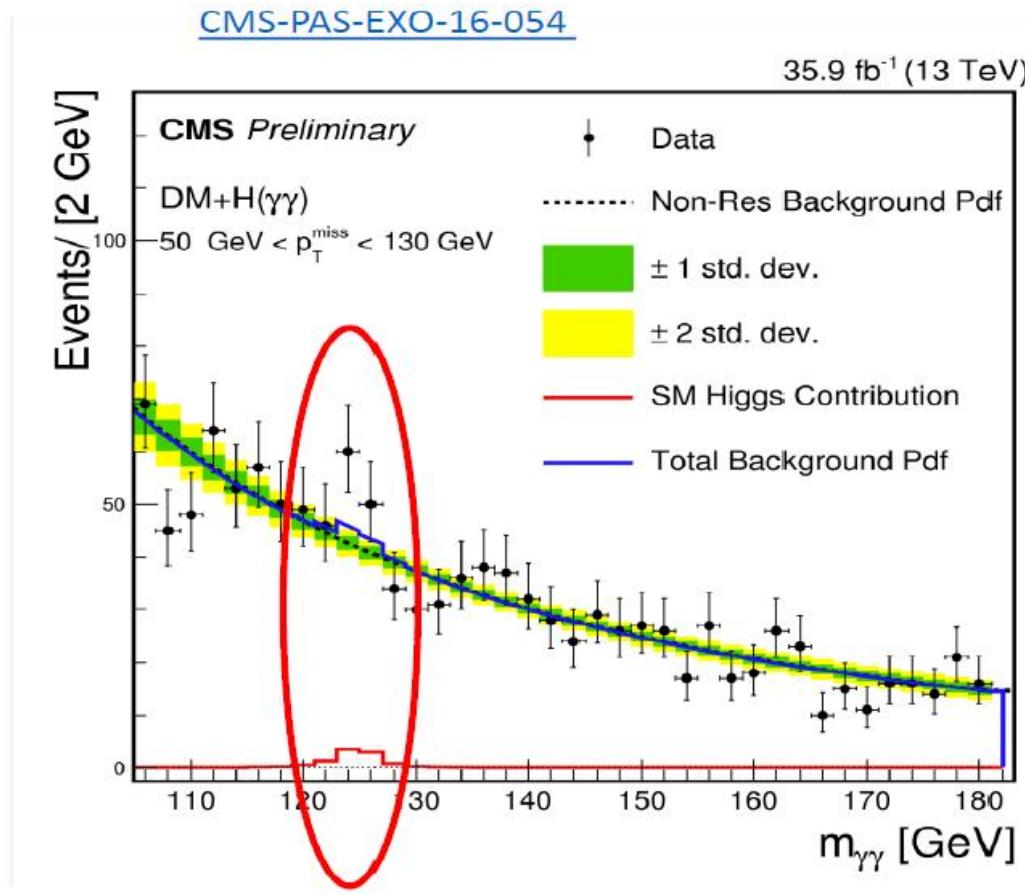
Mono-photon



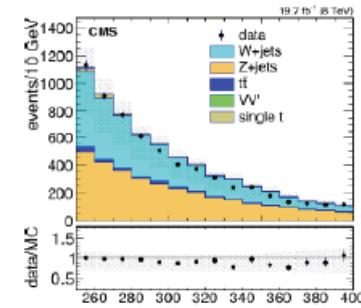
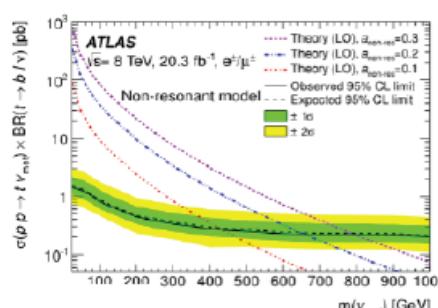
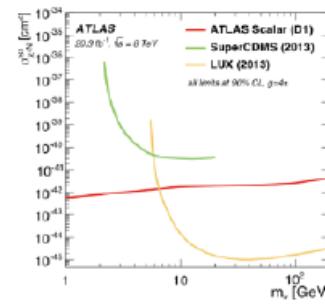
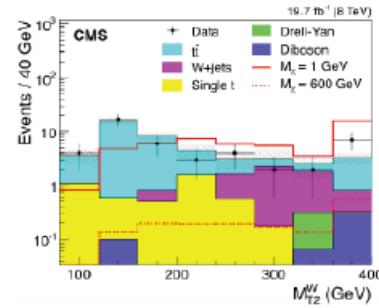
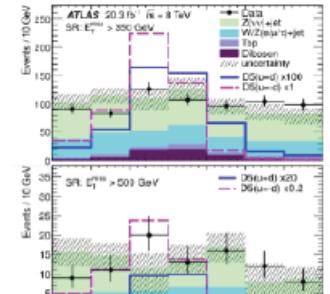
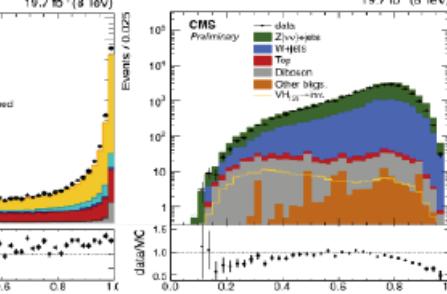
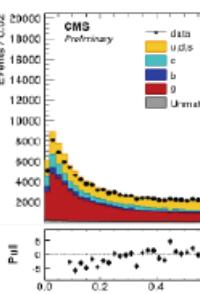
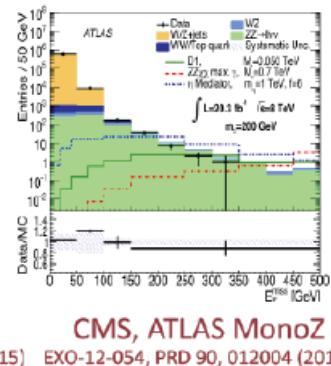
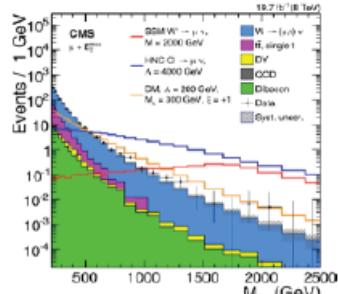
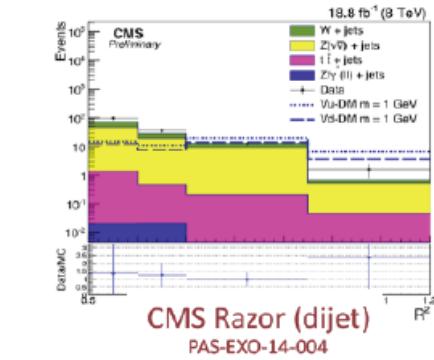
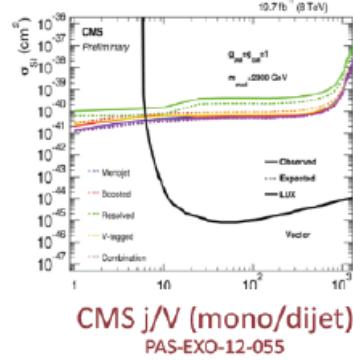
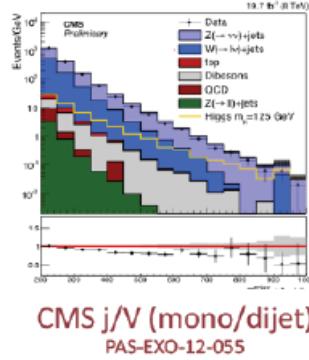
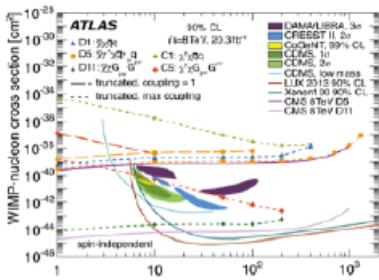
- Photon $E_T > 150 \text{ GeV}, |\eta| < 2.37$
- $E_T^{\text{miss}}/\sqrt{\sum E_T} > 8.5 \text{ GeV}^{1/2}$
- $\Delta\phi(\text{photon}, E_T^{\text{miss}}) > 0.4$
- $N_{\text{jets}}(p_T > 30 \text{ GeV}, |\eta| < 4.5) \leq 1$

Mono-X searches

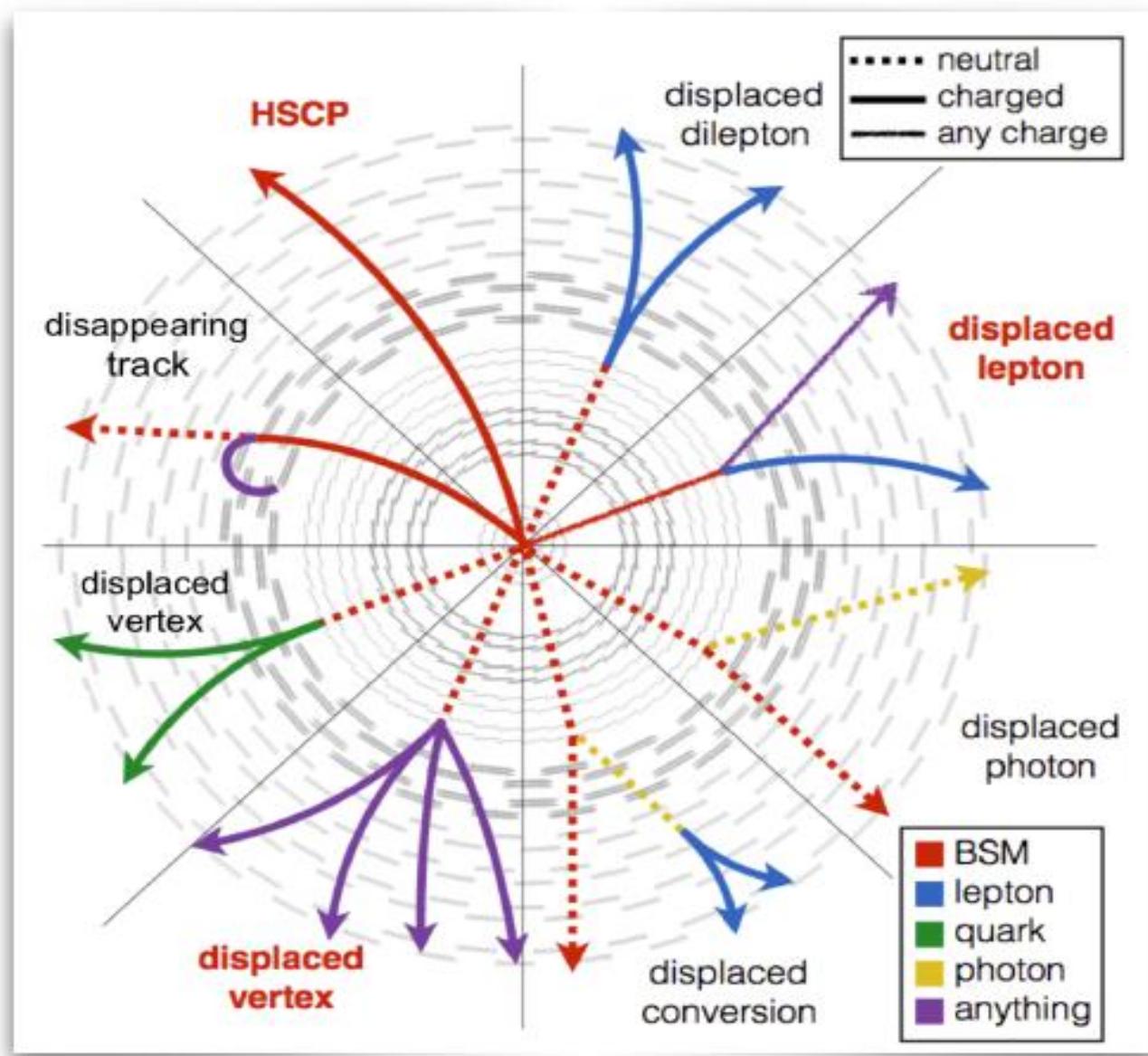
Mono-Higgs



Plenty of mono-signatures

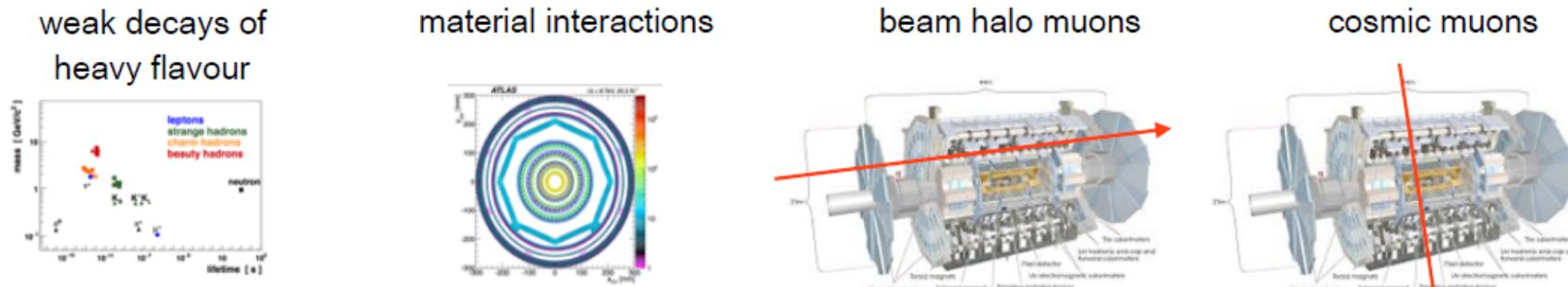


Unconventional signatures



Unconventional signatures: challenges

- ▶ **Trigger:** combination of hardware + software that must decide very quickly whether to save an event or lose it forever
 - First step in every search for LLPs: make sure that interesting events are saved!
 1. In associated production, trigger on prompt particle (Eg. WH prod. trigger on mu; ISR trigger on MET)
 2. Design and develop a new trigger. Need to keep trigger rates under control and within budget
- ▶ **Object identification** algorithms assume prompt particles. Need to adapt them
- ▶ **Backgrounds:** usually instrumental background such as miss-identified leptons (“fakes”) and non-collision backgrounds (NCB) have to be taken into account

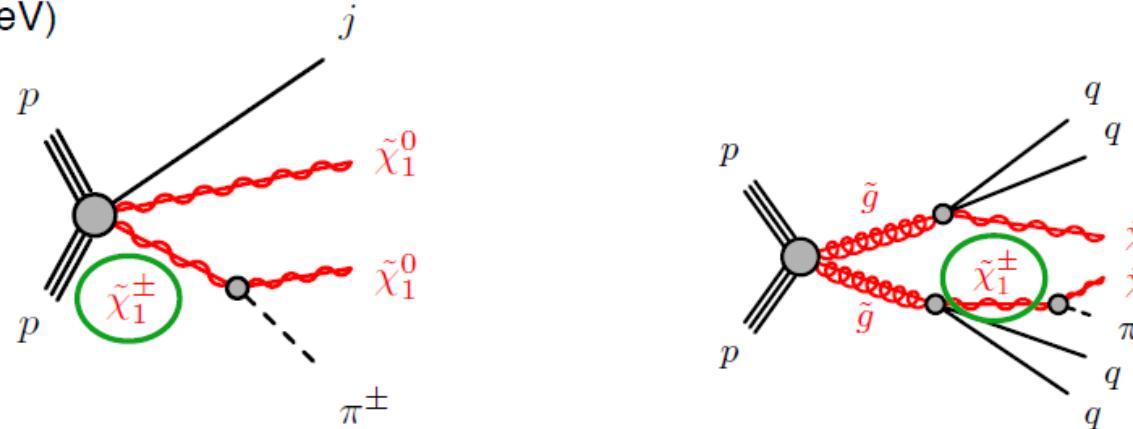


- ▶ Systematic uncertainties: can't use standard recommendations for object reconstruction nor trigger

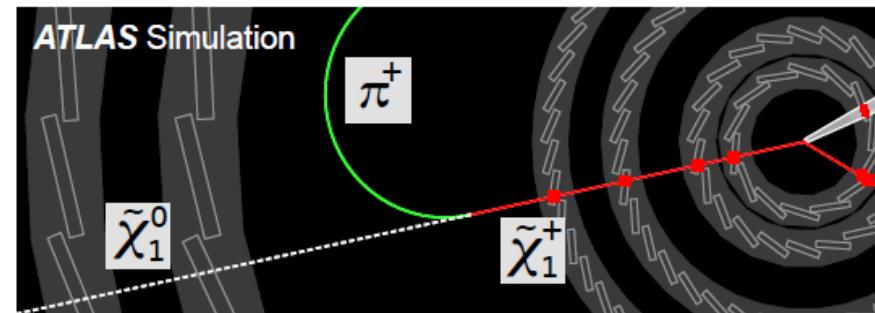
Unconventional signatures: disappearing tracks

- ▶ Search for disappearing track + MET + jets
- ▶ Signature: Chargino track “disappears” when it decays, into MET
 - ▶ Low momentum pion track (~ 0.1 GeV) is hard to reconstruct
 - ▶ Challenge to identify the legitimate real tracklets (non-fake) using only a few measurement tracks
- ▶ Benchmark model: AMSB model with almost degenerate neutralino and chargino
 $\Delta m \sim \mathcal{O}(100 \text{ MeV})$

LLP

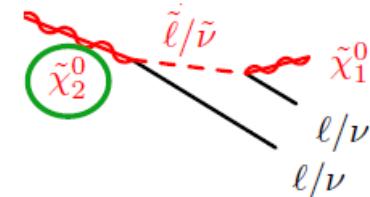
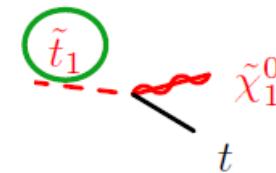
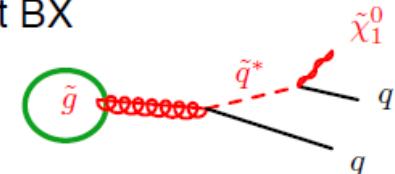


- ▶ Search sensitive to LLP lifetime of 10ps to 10 ns



Unconventional signatures: stopped particles

- ▶ Search for stopped LLPs decaying during non-collision bunch crossings (BX)
- ▶ Signature: LLPs come to rest in the detector and decays after the current BX
 - ▶ most likely to stop in the densest detector materials:
 - ▶ Calorimeters (ECAL, HCAL):
 - a) Split SUSY: two-body and three-body decays of a gluino
 - b) top squark decay
 - ▶ Steel yoke in the muon system:
 - a) three-body decay of the gluino ($g \rightarrow q\bar{q}\chi_2$, $\chi_2 \rightarrow \mu^+\mu^-\chi_1$)
 - b) MCHAMPs, with charge $|Q| = 2e$ decays into two same-sign muons ($MCHAMP \rightarrow \mu^\pm\mu^\pm$)
 - ▶ Search sensitive to wide range of LLP lifetime: 10^{-5} to 10^6 s



Summary (Exotic Searches)

- Searches for Exotic searches
 - All major search channels reached 1 TeV scales
 - Quite a few at 10 TeV
 - New probe: Higgs boson → emerging field
- Dark Matter Searches are thriving at the LHC
- For vector and axial vector interactions
 - Dark Matter masses up to 400 GeV – 700 GeV (mono-jet) excluded
 - Mediator mass up to 1.6 – 1.8 TeV (mono-jet) excluded
 - Mediator mass up to 1.2 TeV (mono-photon) excluded
 - Mediator mass up to 0.7 TeV (mono-Z) excluded
- LHC searches complement DD experiments
 - $m_{DM} < O(10 \text{ GeV})$

Glimpse at Future Hadron Colliders

The candidate machines in a tiny nutshell

Project	HL-LHC	HE-LHC	FCC-hh	SppC
Location	CERN	CERN	CERN	China TBD
Circ.	27 km	27 km	100 km	55 - 100 km
COM energy	14 (15?) TeV	27 TeV	100 TeV	70 -140 TeV
Luminosity	3 ab-1	15 ab-1	20-30 ab-1	TBD
PU	up to 200	up to 800	up to 1000	TBS
Bunch sp.	25 ns	25 ns	25 ns	25 ns
Field	8T	16T	16T	20T
When?	Until 2037	After 2037?	After 2037	TBS

Much much more in Lecture by R. Corsini

Detector, Trigger DAQ, Reconstruction challenges

Two challenges: higher PU (1000) higher pT

Answer: granularity and resolution

Decay products of a Z at 10 TeV are separated by $\Delta R = 0.01$

A b at 5 TeV can travel 50cm and a tau 10 cm

