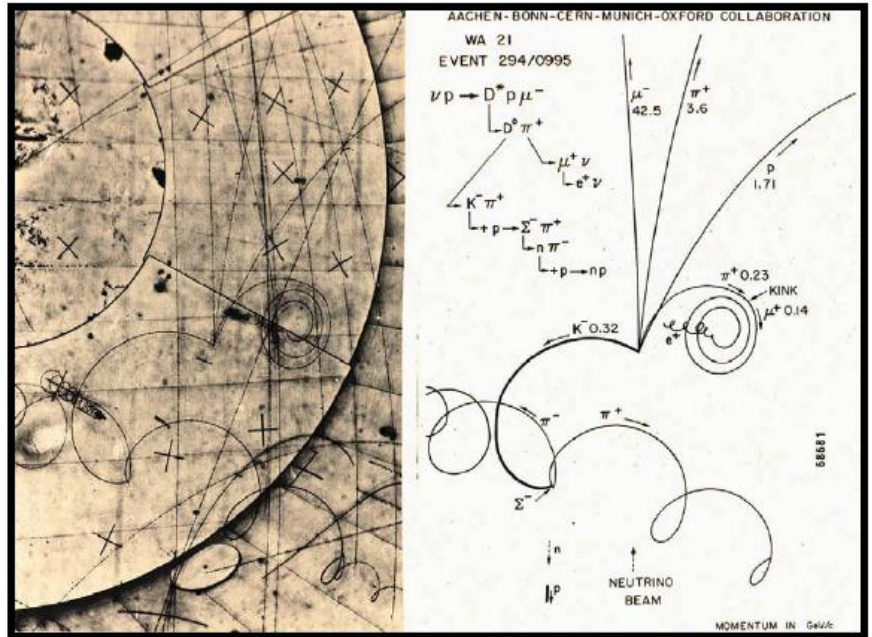


# Elementary Particle Physics: theory and experiments

## Symmetries and the Quark Model



Follow the course/slides from M. A. Thomson lectures at Cambridge University

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# Introduction/Aims

- ★ Symmetries play a central role in particle physics; one aim of particle physics is to discover the fundamental symmetries of our universe
- ★ In this handout will apply the idea of symmetry to the quark model with the aim of :
  - ♦ Deriving hadron wave-functions
  - ♦ Providing an introduction to the more abstract ideas of colour and QCD
  - ♦ Ultimately explaining why hadrons only exist as  $\bar{q}q$  (mesons)  $qqq$  (baryons) or  $\bar{q}\bar{q}\bar{q}$  (anti-baryons)
- + introduce the ideas of the  $SU(2)$  and  $SU(3)$  symmetry groups which play a major role in particle physics

# Symmetries and Conservation Laws

- ★ Suppose physics is invariant under the transformation

$$\psi \rightarrow \psi' = \hat{U} \psi \quad \text{e.g. rotation of the coordinate axes}$$

- To conserve probability normalisation require

$$\langle \psi | \psi \rangle = \langle \psi' | \psi' \rangle = \langle \hat{U} \psi | \hat{U} \psi \rangle = \langle \psi | \hat{U}^\dagger \hat{U} | \psi \rangle$$

$$\rightarrow \boxed{\hat{U}^\dagger \hat{U} = 1} \quad \text{i.e. } \hat{U} \text{ has to be unitary}$$

- For physical predictions to be unchanged by the symmetry transformation, also require all QM matrix elements unchanged

$$\boxed{\langle \psi | \hat{H} | \psi \rangle = \langle \psi' | \hat{H} | \psi' \rangle} = \langle \psi | \hat{U}^\dagger \hat{H} \hat{U} | \psi \rangle$$

i.e. require

$$\hat{U}^\dagger \hat{H} \hat{U} = \hat{H}$$

$\times \hat{U}$

$$\hat{U} \hat{U}^\dagger \hat{H} \hat{U} = \hat{U} \hat{H} \rightarrow \hat{H} \hat{U} = \hat{U} \hat{H}$$

therefore

$$\boxed{[\hat{H}, \hat{U}] = 0}$$

$\hat{U}$  commutes with the Hamiltonian

- ★ Now consider the infinitesimal transformation ( $\epsilon$  small)

$$\hat{U} = 1 + i\epsilon \hat{G}$$

( $\hat{G}$  is called the generator of the transformation)

- For  $\hat{U}$  to be unitary

$$\hat{U}\hat{U}^\dagger = (1 + i\varepsilon\hat{G})(1 - i\varepsilon\hat{G}^\dagger) = 1 + i\varepsilon(\hat{G} - \hat{G}^\dagger) + O(\varepsilon^2)$$

neglecting terms in  $\varepsilon^2$        $UU^\dagger = 1 \quad \rightarrow \quad \boxed{\hat{G} = \hat{G}^\dagger}$

i.e.  $\hat{G}$  is Hermitian and therefore corresponds to an observable quantity  $G$  !

- Furthermore,  $[\hat{H}, \hat{U}] = 0 \Rightarrow [\hat{H}, 1 + i\varepsilon \hat{G}] = 0 \Rightarrow [\hat{H}, \hat{G}] = 0$

But from QM       $\frac{d}{dt} \langle \hat{G} \rangle = i \langle [\hat{H}, \hat{G}] \rangle = 0$

i.e.  $G$  is a **conserved** quantity.

**Symmetry  $\longleftrightarrow$  Conservation Law**

★ For each symmetry of nature have an observable conserved quantity

**Example:** Infinitesimal spatial translation  $x \rightarrow x + \varepsilon$

i.e. expect physics to be invariant under  $\psi(x) \rightarrow \psi' = \psi(x + \varepsilon)$

$$\psi'(x) = \psi(x + \varepsilon) = \psi(x) + \frac{\partial \psi}{\partial x} \varepsilon = \left( 1 + \varepsilon \frac{\partial}{\partial x} \right) \psi(x)$$

but  $\hat{p}_x = -i \frac{\partial}{\partial x} \quad \rightarrow \quad \psi'(x) = (1 + i\varepsilon \hat{p}_x) \psi(x)$

The generator of the symmetry transformation is  $\hat{p}_x \rightarrow p_x$  is conserved

- **Translational invariance of physics implies momentum conservation !**

- In general the symmetry operation may depend on more than one parameter

$$\hat{U} = 1 + i\vec{\epsilon} \cdot \vec{G}$$

For example for an infinitesimal 3D linear translation :  $\vec{r} \rightarrow \vec{r} + \vec{\epsilon}$

$$\rightarrow \hat{U} = 1 + i\vec{\epsilon} \cdot \vec{p} \quad \vec{p} = (\hat{p}_x, \hat{p}_y, \hat{p}_z)$$

- So far have only considered an infinitesimal transformation, however a finite transformation can be expressed as a series of infinitesimal transformations

$$\hat{U}(\vec{\alpha}) = \lim_{n \rightarrow \infty} \left( 1 + i \frac{\vec{\alpha}}{n} \cdot \vec{G} \right)^n = e^{i\vec{\alpha} \cdot \vec{G}}$$

**Example:** Finite spatial translation in 1D:  $x \rightarrow x + x_0$  with  $\hat{U}(x_0) = e^{ix_0 \hat{p}_x}$

$$\begin{aligned} \psi'(x) = \psi(x + x_0) &= \hat{U} \psi(x) = \exp \left( x_0 \frac{d}{dx} \right) \psi(x) && \left( p_x = -i \frac{\partial}{\partial x} \right) \\ &= \left( 1 + x_0 \frac{d}{dx} + \frac{x_0^2}{2!} \frac{d^2}{dx^2} + \dots \right) \psi(x) \\ &= \psi(x) + x_0 \frac{d\psi}{dx} + \frac{x_0^2}{2} \frac{d^2\psi}{dx^2} + \dots \end{aligned}$$

i.e. obtain the expected Taylor expansion

# Symmetries in Particle Physics: Isospin

- The proton and neutron have very similar masses and the nuclear force is found to be approximately charge-independent, i.e.

$$V_{pp} \approx V_{np} \approx V_{nn}$$

- To reflect this symmetry, Heisenberg (1932) proposed that if you could “switch off” the electric charge of the proton

There would be no way to distinguish between a proton and neutron

- Proposed that the neutron and proton should be considered as two states of a single entity; the nucleon

$$p = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad n = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- ★ Analogous to the spin-up/spin-down states of a spin- $\frac{1}{2}$  particle

**ISOSPIN**

- ★ Expect physics to be invariant under rotations in this space
- The neutron and proton form an isospin doublet with total isospin  $I = \frac{1}{2}$  and third component  $I_3 = \pm \frac{1}{2}$

# Flavour Symmetry of the Strong Interaction

We can extend this idea to the quarks:

★ Assume the strong interaction treats all quark flavours equally (it does)

• Because  $m_u \approx m_d$ :

The strong interaction possesses an **approximate** flavour symmetry i.e. from the point of view of the strong interaction nothing changes if all up quarks are replaced by down quarks and *vice versa*.

• Choose the basis

$$u = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad d = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

• Express the invariance of the strong interaction under  $u \leftrightarrow d$  as invariance under “rotations” in an abstract isospin space

$$\begin{pmatrix} u' \\ d' \end{pmatrix} = \hat{U} \begin{pmatrix} u \\ d \end{pmatrix} = \begin{pmatrix} U_{11} & U_{12} \\ U_{21} & U_{22} \end{pmatrix} \begin{pmatrix} u \\ d \end{pmatrix}$$

The 2x2 **unitary** matrix depends on 4 complex numbers, i.e. 8 real parameters  
But there are four constraints from  $\hat{U}^\dagger \hat{U} = 1$

➔ **8 - 4 = 4 independent matrices**

• In the language of group theory the four matrices form the **U(2)** group

- One of the matrices corresponds to multiplying by a phase factor

$$\hat{U}_1 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} e^{i\phi}$$

not a flavour transformation and of no relevance here.

- The remaining three matrices form an **SU(2)** group (special unitary) with  $\det U = 1$
- For an infinitesimal transformation, in terms of the Hermitian generators  $\hat{G}$

$$\hat{U} = 1 + i\varepsilon \hat{G}$$

- $\det U = 1 \Rightarrow \text{Tr}(\hat{G}) = 0$

- A linearly independent choice for  $\hat{G}$  are the Pauli spin matrices

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

- The proposed flavour symmetry of the strong interaction has the same transformation properties as SPIN !

- Define **ISOSPIN**:  $\vec{T} = \frac{1}{2} \vec{\sigma} \quad \hat{U} = e^{i\vec{\alpha} \cdot \vec{T}}$

- Check this works, for an infinitesimal transformation

$$\hat{U} = 1 + \frac{1}{2} i \vec{\varepsilon} \cdot \vec{\sigma} = 1 + \frac{i}{2} (\varepsilon_1 \sigma_1 + \varepsilon_2 \sigma_2 + \varepsilon_3 \sigma_3) = \begin{pmatrix} 1 + \frac{1}{2} i \varepsilon_3 & \frac{1}{2} i (\varepsilon_1 - i \varepsilon_2) \\ \frac{1}{2} i (\varepsilon_1 + i \varepsilon_2) & 1 - \frac{1}{2} i \varepsilon_3 \end{pmatrix}$$

Which is, as required, unitary and has unit determinant

$$U^\dagger U = I + O(\varepsilon^2) \quad \det U = 1 + O(\varepsilon^2)$$



# Properties of Isospin

- Isospin has the exactly the same properties as spin

$$[T_1, T_2] = iT_3 \quad [T_2, T_3] = iT_1 \quad [T_3, T_1] = iT_2$$

$$[T^2, T_3] = 0 \quad T^2 = T_1^2 + T_2^2 + T_3^2$$

As in the case of spin, have three non-commuting operators,  $T_1, T_2, T_3$  and even though all three correspond to observables, can't know them simultaneously. So label states in terms of **total isospin**  $I$  and the third component of isospin  $I_3$

**NOTE: isospin has nothing to do with spin - just the same mathematics**

- The eigenstates are exact analogues of the eigenstates of ordinary angular momentum  $|s, m\rangle \rightarrow |I, I_3\rangle$   
with  $T^2|I, I_3\rangle = I(I+1)|I, I_3\rangle \quad T_3|I, I_3\rangle = I_3|I, I_3\rangle$
- In terms of isospin:

$$u = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \left| \frac{1}{2}, +\frac{1}{2} \right\rangle \quad d = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \left| \frac{1}{2}, -\frac{1}{2} \right\rangle$$

$$I = \frac{1}{2}, \quad I_3 = \pm \frac{1}{2}$$

- In general  $I_3 = \frac{1}{2}(N_u - N_d)$

- Can define isospin ladder operators – analogous to spin ladder operators

$$T_- \equiv T_1 - iT_2$$

$u \rightarrow d$



$$T_+ \equiv T_1 + iT_2$$

$d \rightarrow u$

$$T_+ |I, I_3\rangle = \sqrt{I(I+1) - I_3(I_3+1)} |I, I_3+1\rangle$$

$$T_- |I, I_3\rangle = \sqrt{I(I+1) - I_3(I_3-1)} |I, I_3-1\rangle$$

Step up/down in  $I_3$  until reach end of multiplet  $T_+ |I, +I\rangle = 0$   $T_- |I, -I\rangle = 0$

$$T_+ u = 0 \quad T_+ d = u \quad T_- u = d \quad T_- d = 0$$

- Ladder operators turn  $u \rightarrow d$  and  $d \rightarrow u$
- ★ Combination of isospin: e.g. what is the isospin of a system of two  $d$  quarks, is exactly analogous to combination of spin (i.e. angular momentum)

$$|I^{(1)}, I_3^{(1)}\rangle |I^{(2)}, I_3^{(2)}\rangle \rightarrow |I, I_3\rangle$$

- $I_3$  additive :  $I_3 = I_3^{(1)} + I_3^{(2)}$

- $I$  in integer steps from  $|I^{(1)} - I^{(2)}|$  to  $|I^{(1)} + I^{(2)}|$

- ★ Assumed **symmetry** of Strong Interaction under isospin transformations implies the existence of **conserved quantities**

- In strong interactions  $I_3$  and  $I$  are conserved, analogous to conservation of  $J_z$  and  $J$  for angular momentum

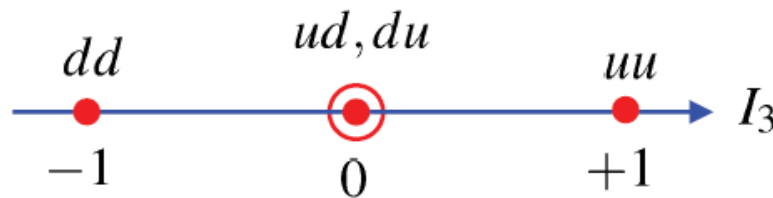
# Combining Quarks


## Goal: derive proton wave-function

- First combine two quarks, then combine the third
- Use requirement that fermion wave-functions are anti-symmetric

Isospin starts to become useful in defining states of more than one quark.

e.g. two quarks, here we have four possible combinations:



Note:  represents two states with the same value of  $I_3$

- We can immediately identify the extremes ( $I_3$  additive)

$$uu \equiv \left| \frac{1}{2}, \frac{1}{2} \right\rangle \left| \frac{1}{2}, \frac{1}{2} \right\rangle = |1, +1\rangle$$

$$dd \equiv \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \left| \frac{1}{2}, -\frac{1}{2} \right\rangle = |1, -1\rangle$$

To obtain the  $|1, 0\rangle$  state use ladder operators

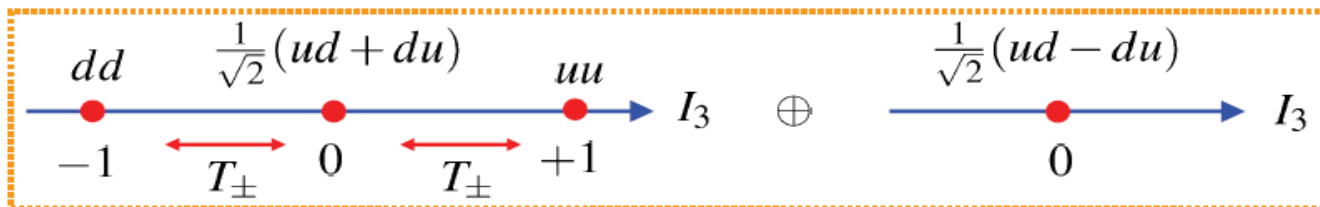
$$T_- |1, +1\rangle = \sqrt{2} |1, 0\rangle = T_-(uu) = ud + du$$

$$\rightarrow |1, 0\rangle = \frac{1}{\sqrt{2}} (ud + du)$$

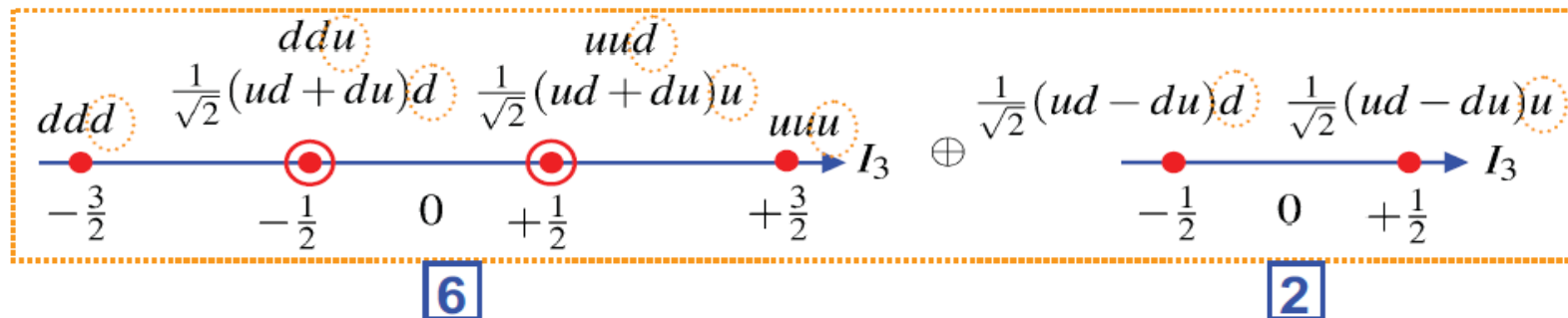
The final state,  $|0, 0\rangle$ , can be found from orthogonality with  $|1, 0\rangle$

$$\rightarrow |0, 0\rangle = \frac{1}{\sqrt{2}} (ud - du)$$

- From four possible combinations of isospin doublets obtain a **triplet** of isospin 1 states and a **singlet** isospin 0 state  $2 \otimes 2 = 3 \oplus 1$

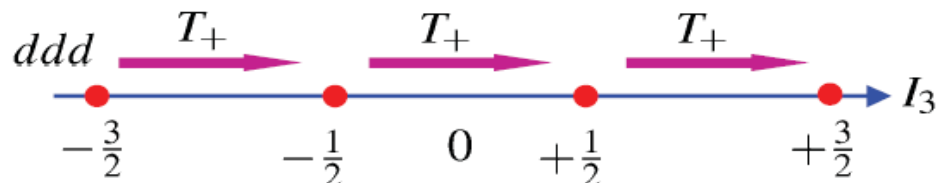


- Can move around **within** multiplets using ladder operators
- note, as anticipated  $I_3 = \frac{1}{2}(N_u - N_d)$
- States with different total isospin are physically different - the isospin 1 triplet is **symmetric** under interchange of quarks 1 and 2 whereas singlet is **anti-symmetric**
- ★ Now add an additional up or down quark. From **each of the above 4 states** get two new isospin states with  $I'_3 = I_3 \pm \frac{1}{2}$



- Use ladder operators and orthogonality to group the 6 states into isospin multiplets, e.g. to obtain the  $I = \frac{3}{2}$  states, step up from  $ddd$

★ Derive the  $I = \frac{3}{2}$  states from  $ddd \equiv |\frac{3}{2}, -\frac{3}{2}\rangle$



$$T_+|\frac{3}{2}, -\frac{3}{2}\rangle = T_+(ddd) = (T_+d)dd + d(T_+d)d + dd(T_+)d$$

$$\sqrt{3}|\frac{3}{2}, -\frac{1}{2}\rangle = udd + dud + ddu$$

$$|\frac{3}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(udd + dud + ddu)$$

$$T_+|\frac{3}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}T_+(udd + dud + ddu)$$

$$2|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(uud + udu + uud + duu + udu + duu)$$

$$|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(uud + udu + duu)$$

$$T_+|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}T_+(uud + udu + duu)$$

$$\sqrt{3}|\frac{3}{2}, +\frac{3}{2}\rangle = \frac{1}{\sqrt{3}}(uuu + uuu + uuu)$$

$$|\frac{3}{2}, +\frac{3}{2}\rangle = uuu$$

★ From the **6** states on previous page, use orthogonality to find  $|\frac{1}{2}, \pm\frac{1}{2}\rangle$  states

★ The **2** states on the previous page give another  $|\frac{1}{2}, \pm\frac{1}{2}\rangle$  doublet

- ★ The eight states  $uuu, uud, udu, udd, duu, dud, ddu, ddd$  are grouped into an **isospin quadruplet** and two **isospin doublets**

$$2 \otimes 2 \otimes 2 = 2 \otimes (3 \oplus 1) = (2 \otimes 3) \oplus (2 \otimes 1) = 4 \oplus 2 \oplus 2$$

- Different multiplets have different symmetry properties

$$|\frac{3}{2}, +\frac{3}{2}\rangle = uuu$$

$$|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(uud + udu + duu)$$

$$|\frac{3}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(ddu + dud + udd)$$

$$|\frac{3}{2}, -\frac{3}{2}\rangle = ddd$$

**S**

A quadruplet of states which are symmetric under the interchange of any two quarks

$$|\frac{1}{2}, -\frac{1}{2}\rangle = -\frac{1}{\sqrt{6}}(2ddu - udd - dud)$$

$$|\frac{1}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{6}}(2uud - udu - duu)$$

**M<sub>S</sub>**

Mixed symmetry.  
Symmetric for  $1 \leftrightarrow 2$

$$|\frac{1}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(udd - dud)$$

$$|\frac{1}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(udu - duu)$$

**M<sub>A</sub>**

Mixed symmetry.  
Anti-symmetric for  $1 \leftrightarrow 2$

- Mixed symmetry states have no definite symmetry under interchange of quarks  $1 \leftrightarrow 3$  etc.

- ★ The eight states  $uuu, uud, udu, udd, duu, dud, ddu, ddd$  are grouped into an **isospin quadruplet** and two **isospin doublets**

$$2 \otimes 2 \otimes 2 = 2 \otimes (3 \oplus 1) = (2 \otimes 3) \oplus (2 \otimes 1) = 4 \oplus 2 \oplus 2$$

- Different multiplets have different symmetry properties

$$|\frac{3}{2}, +\frac{3}{2}\rangle = uuu$$

$$|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(uud + udu + duu)$$

$$|\frac{3}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(ddu + dud + udd)$$

$$|\frac{3}{2}, -\frac{3}{2}\rangle = ddd$$

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A quadruplet of states which are symmetric under the interchange of any two quarks

$$|\frac{1}{2}, -\frac{1}{2}\rangle = -\frac{1}{\sqrt{6}}(2ddu - udd - dud)$$

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**M<sub>S</sub>**

Mixed symmetry.  
Symmetric for  $1 \leftrightarrow 2$

$$|\frac{1}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(udd - dud)$$

$$|\frac{1}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(udu - duu)$$

**M<sub>A</sub>**

Mixed symmetry.  
Anti-symmetric for  $1 \leftrightarrow 2$

- Mixed symmetry states have no definite symmetry under interchange of quarks  $1 \leftrightarrow 3$  etc.

# Combining Spin

- Can apply exactly the same mathematics to determine the possible spin wave-functions for a combination of 3 spin-half particles

$$|\frac{3}{2}, +\frac{3}{2}\rangle = \uparrow\uparrow\uparrow$$

$$|\frac{3}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(\uparrow\uparrow\downarrow + \uparrow\downarrow\uparrow + \downarrow\uparrow\uparrow)$$

$$|\frac{3}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}(\downarrow\downarrow\uparrow + \downarrow\uparrow\downarrow + \uparrow\downarrow\downarrow)$$

$$|\frac{3}{2}, -\frac{3}{2}\rangle = \downarrow\downarrow\downarrow$$

**S**

A quadruplet of states which are symmetric under the interchange of any two quarks

$$|\frac{1}{2}, -\frac{1}{2}\rangle = -\frac{1}{\sqrt{6}}(2\downarrow\downarrow\uparrow - \uparrow\downarrow\downarrow - \downarrow\uparrow\downarrow)$$

$$|\frac{1}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{6}}(2\uparrow\uparrow\downarrow - \uparrow\downarrow\uparrow - \downarrow\uparrow\uparrow)$$

**M<sub>S</sub>**

Mixed symmetry.  
Symmetric for 1 ↔ 2

$$|\frac{1}{2}, -\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\downarrow - \downarrow\uparrow\downarrow)$$

$$|\frac{1}{2}, +\frac{1}{2}\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\uparrow - \downarrow\uparrow\uparrow)$$

**M<sub>A</sub>**

Mixed symmetry.  
Anti-symmetric for 1 ↔ 2

- Can now form total wave-functions for combination of three quarks



# Bayron Wave-functions (ud)

★ Quarks are fermions so require that the total wave-function is anti-symmetric under the interchange of any two quarks

★ the total wave-function can be expressed in terms of:

$$\Psi = \phi_{\text{flavour}} \chi_{\text{spin}} \xi_{\text{colour}} \eta_{\text{space}}$$

★ The colour wave-function for all bound qqq states is anti-symmetric

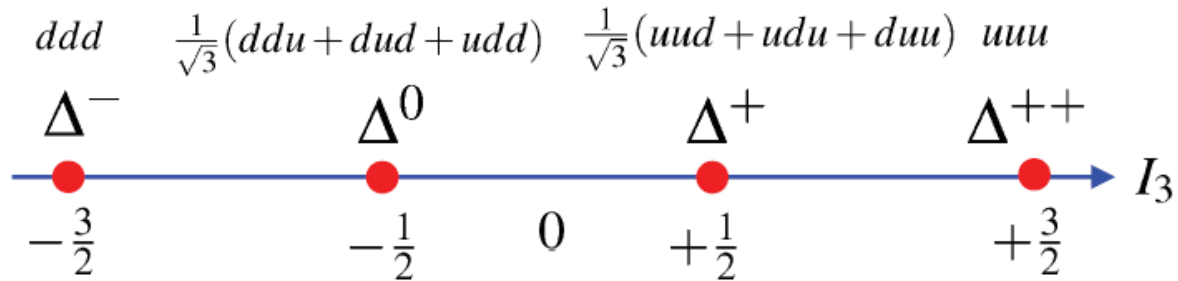
• Here we will only consider the lowest mass, ground state, baryons where there is no internal orbital angular momentum.

• For  $L=0$  the spatial wave-function is symmetric  $(-1)^L$ .



★ Two ways to form a totally symmetric wave-function from spin and isospin states:

① combine totally symmetric spin and isospin wave-functions  $\phi(S)\chi(S)$



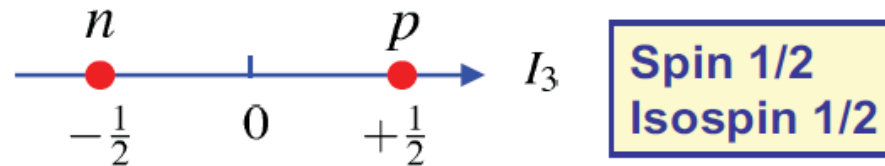
Spin 3/2  
Isospin 3/2

② combine mixed symmetry spin and mixed symmetry isospin states

- Both  $\phi(M_S)\chi(M_S)$  and  $\phi(M_A)\chi(M_A)$  are sym. under inter-change of quarks  $1 \leftrightarrow 2$
- Not sufficient, these combinations have no definite symmetry under  $1 \leftrightarrow 3, \dots$
- However, it is not difficult to show that the (normalised) linear combination:

$$\frac{1}{\sqrt{2}}\phi(M_S)\chi(M_S) + \frac{1}{\sqrt{2}}\phi(M_A)\chi(M_A)$$

is **totally symmetric** (i.e. symmetric under  $1 \leftrightarrow 2$ ;  $1 \leftrightarrow 3$ ;  $2 \leftrightarrow 3$  )



- The spin-up proton wave-function is therefore:

$$|p \uparrow\rangle = \frac{1}{6\sqrt{2}}(2uud - udu - duu)(2\uparrow\uparrow\downarrow - \uparrow\downarrow\uparrow - \downarrow\uparrow\uparrow) + \frac{1}{2\sqrt{2}}(udu - duu)(\uparrow\downarrow\uparrow - \downarrow\uparrow\uparrow)$$

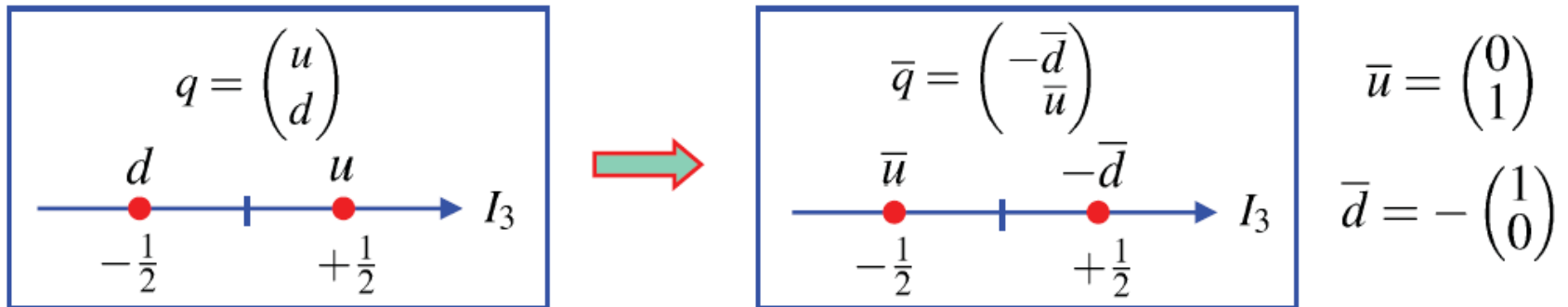


$$|p \uparrow\rangle = \frac{1}{\sqrt{18}}( 2u \uparrow u \uparrow d \downarrow - u \uparrow u \downarrow d \uparrow - u \downarrow u \uparrow d \uparrow + \\ 2u \uparrow d \downarrow u \uparrow - u \uparrow d \uparrow u \downarrow - u \downarrow d \uparrow u \uparrow + \\ 2d \downarrow u \uparrow u \uparrow - d \uparrow u \downarrow u \uparrow - d \uparrow u \uparrow u \uparrow )$$

**NOTE:** not always necessary to use the fully symmetrised proton wave-function, e.g. the first 3 terms are sufficient for calculating the proton magnetic moment

# Anti-quarks and Mesons (u and d)

- ★ The **u, d quarks** and  **$\bar{u}, \bar{d}$  anti-quarks** are represented as isospin doublets



- **Subtle point:** The ordering and the minus sign in the anti-quark doublet ensures that anti-quarks and quarks transform in the same way (see Appendix I). This is necessary if we want physical predictions to be invariant under  $u \leftrightarrow d; \bar{u} \leftrightarrow \bar{d}$
- Consider the effect of ladder operators on the anti-quark isospin states

e.g. 
$$T_+ \bar{u} = T_+ \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = -\bar{d}$$

- The effect of the ladder operators on anti-particle isospin states are:

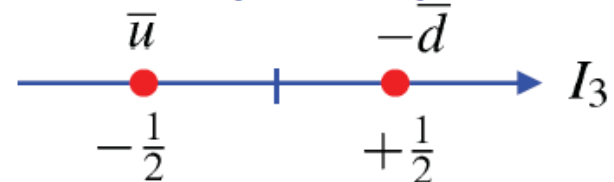
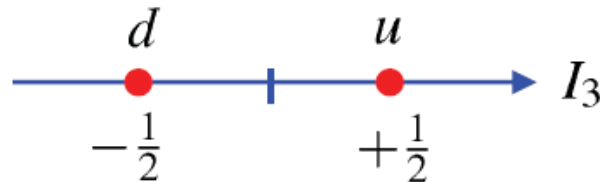
$$T_+ \bar{u} = -\bar{d} \quad T_+ \bar{d} = 0 \quad T_- \bar{u} = 0 \quad T_- \bar{d} = -\bar{u}$$

Compare with

$$T_+ u = 0 \quad T_+ d = u \quad T_- u = d \quad T_- d = 0$$

# Light ud Mesons

- ★ Can now construct meson states from combinations of up/down quarks



- Consider the  $q\bar{q}$  combinations in terms of isospin

$$|1, +1\rangle = \left| \frac{1}{2}, +\frac{1}{2} \right\rangle \overline{\left| \frac{1}{2}, +\frac{1}{2} \right\rangle} = -u\bar{d}$$

$$|1, -1\rangle = \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \overline{\left| \frac{1}{2}, -\frac{1}{2} \right\rangle} = d\bar{u}$$

The bar indicates this is the isospin representation of an anti-quark

To obtain the  $I_3 = 0$  states use ladder operators and orthogonality

$$T_- |1, +1\rangle = T_- [-u\bar{d}]$$

$$\sqrt{2}|1, 0\rangle = -T_- [u]\bar{d} - uT_- [\bar{d}]$$

$$= -d\bar{d} + u\bar{u}$$

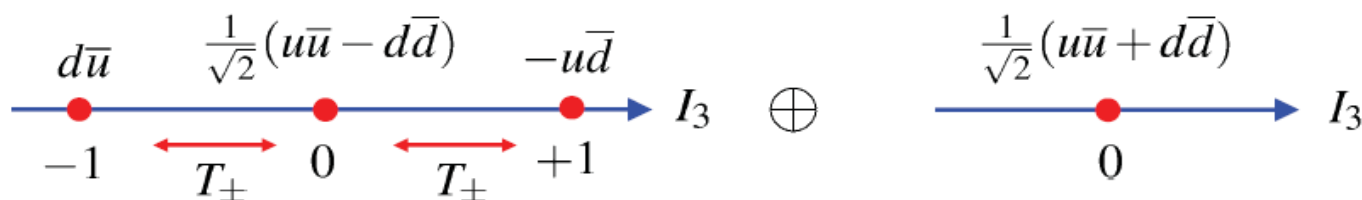
$$\Rightarrow |1, 0\rangle = \frac{1}{\sqrt{2}} (u\bar{u} - d\bar{d})$$

- Orthogonality gives:  $|0, 0\rangle = \frac{1}{\sqrt{2}} (u\bar{u} + d\bar{d})$

★ To summarise:



➔ Triplet of  $I = 1$  states and a singlet  $I = 0$  state



• You will see this written as  $2 \otimes \bar{2} = 3 \oplus 1$

Quark doublet

Anti-quark doublet

• To show the state obtained from orthogonality with  $|1, 0\rangle$  is a singlet use ladder operators

$$T_+ |0, 0\rangle = T_+ \frac{1}{\sqrt{2}} (u\bar{u} + d\bar{d}) = \frac{1}{\sqrt{2}} (-u\bar{d} + u\bar{d}) = 0$$

similarly  $T_- |0, 0\rangle = 0$

★ A singlet state is a “dead-end” from the point of view of ladder operators

# SU(3) Flavour

- ★ Extend these ideas to include the strange quark. Since  $m_s > m_u, m_d$  don't have an exact symmetry. But  $m_s$  not so very different from  $m_u, m_d$  and can treat the strong interaction (and resulting hadron states) as if it were symmetric under  $u \leftrightarrow d \leftrightarrow s$

- **NOTE:** any results obtained from this assumption are only **approximate** as the symmetry is not exact.

- The assumed uds flavour symmetry can be expressed as

$$\begin{pmatrix} u' \\ d' \\ s' \end{pmatrix} = \hat{U} \begin{pmatrix} u \\ d \\ s \end{pmatrix} = \begin{pmatrix} U_{11} & U_{12} & U_{13} \\ U_{21} & U_{22} & U_{23} \\ U_{31} & U_{32} & U_{33} \end{pmatrix} \begin{pmatrix} u \\ d \\ s \end{pmatrix}$$

- The 3x3 **unitary** matrix depends on **9** complex numbers, i.e. **18** real parameters  
There are **9** constraints from  $\hat{U}^\dagger \hat{U} = 1$

➡ Can form  $18 - 9 = 9$  linearly independent matrices

These 9 matrices form a U(3) group.

- As before, one matrix is simply the identity multiplied by a complex phase and is of no interest in the context of flavour symmetry
- The remaining **8** matrices have  $\det U = 1$  and form an **SU(3)** group
- The **eight** matrices (the Hermitian generators) are:  $\vec{T} = \frac{1}{2} \vec{\lambda}$        $\hat{U} = e^{i\vec{\alpha} \cdot \vec{T}}$

★ In SU(3) flavour, the three quark states are represented by:

$$u = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad d = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad s = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

★ In SU(3) uds flavour symmetry contains SU(2) ud flavour symmetry which allows us to write the first three matrices:

$$\lambda_1 = \begin{pmatrix} \sigma_1 & 0 \\ 0 & 0 \end{pmatrix} \quad \lambda_2 = \begin{pmatrix} \sigma_2 & 0 \\ 0 & 0 \end{pmatrix} \quad \lambda_3 = \begin{pmatrix} \sigma_3 & 0 \\ 0 & 0 \end{pmatrix}$$

i.e.  $\boxed{u \leftrightarrow d}$   $\lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

▪ The third component of isospin is now written  $I_3 = \frac{1}{2}\lambda_3$

$$\text{with } I_3 u = +\frac{1}{2}u \quad I_3 d = -\frac{1}{2}d \quad I_3 s = 0$$

▪  $I_3$  "counts the number of up quarks - number of down quarks in a state"

▪ As before, ladder operators  $T_{\pm} = \frac{1}{2}(\lambda_1 \pm i\lambda_2)$   $d \bullet \longleftarrow T_{\pm} \longrightarrow \bullet u$

- Now consider the matrices corresponding to the  $u \leftrightarrow s$  and  $d \leftrightarrow s$

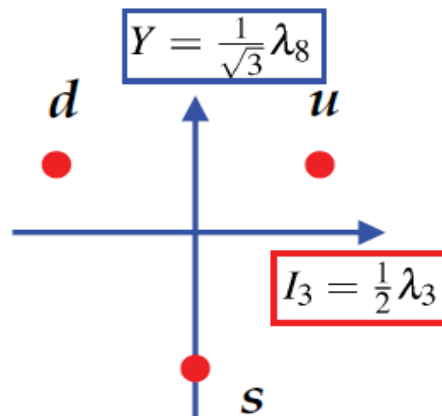
$u \leftrightarrow s$	$\lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$	$\lambda_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$
$d \leftrightarrow s$	$\lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$	$\lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$

- Hence in addition to  $\lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$  have two other traceless diagonal matrices
- However the three diagonal matrices are not be independent.
- Define the eighth matrix,  $\lambda_8$ , as the linear combination:

$$\lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix} + \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

which specifies the "vertical position" in the 2D plane

"Only need two axes (quantum numbers) to specify a state in the 2D plane":  $(I_3, Y)$





★ The other six matrices form six ladder operators which step between the states

$$T_{\pm} = \frac{1}{2}(\lambda_1 \pm i\lambda_2)$$

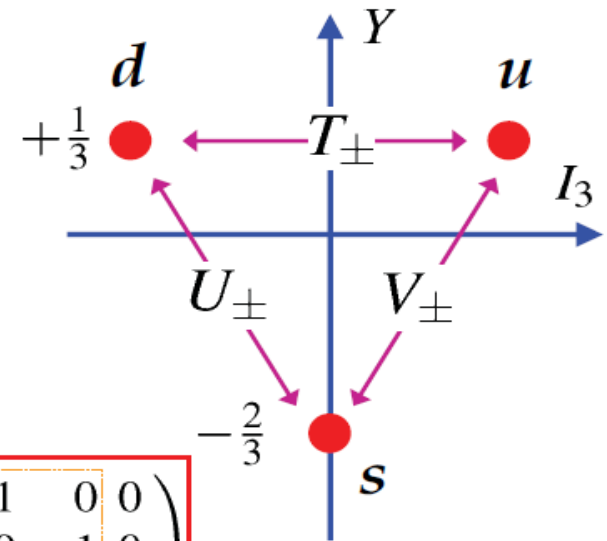
$$V_{\pm} = \frac{1}{2}(\lambda_4 \pm i\lambda_5)$$

$$U_{\pm} = \frac{1}{2}(\lambda_6 \pm i\lambda_7)$$

with

$$I_3 = \frac{1}{2}\lambda_3 \quad Y = \frac{1}{\sqrt{3}}\lambda_8$$

and the eight Gell-Mann matrices



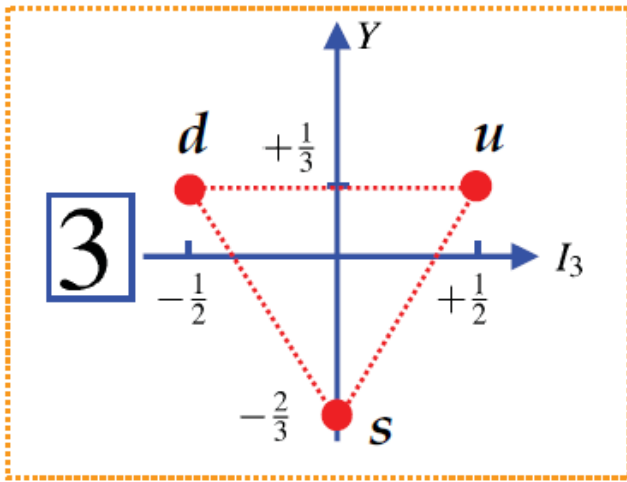
$$\boxed{u \leftrightarrow d} \quad \lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\boxed{u \leftrightarrow s} \quad \lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \lambda_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}$$

$$\boxed{d \leftrightarrow s} \quad \lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad \lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}$$

$$\lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

# Quarks and anti-quarks in SU(3) Flavour

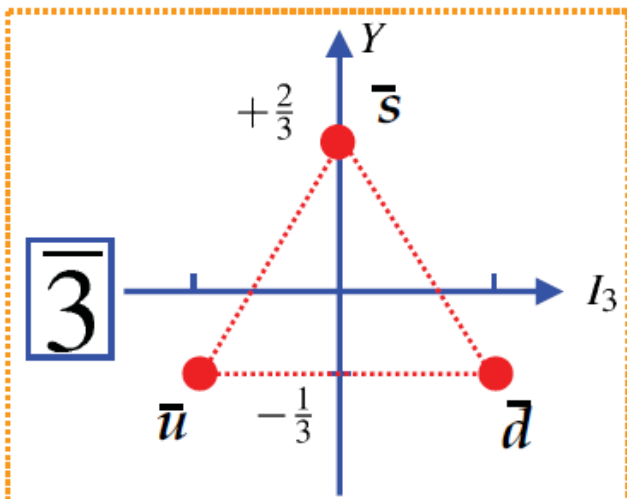


## Quarks

$$I_3 u = +\frac{1}{2}u; \quad I_3 d = -\frac{1}{2}d; \quad I_3 s = 0$$

$$Y u = +\frac{1}{3}u; \quad Y d = +\frac{1}{3}d; \quad Y s = -\frac{2}{3}s$$

- The anti-quarks have opposite SU(3) flavour quantum numbers



## Anti-Quarks

$$I_3 \bar{u} = -\frac{1}{2}\bar{u}; \quad I_3 \bar{d} = +\frac{1}{2}\bar{d}; \quad I_3 \bar{s} = 0$$

$$Y \bar{u} = -\frac{1}{3}\bar{u}; \quad Y \bar{d} = -\frac{1}{3}\bar{d}; \quad Y \bar{s} = +\frac{2}{3}\bar{s}$$

# SU(3) Ladder Operators

- **SU(3)** *uds* flavour symmetry contains *ud*, *us* and *ds* **SU(2)** symmetries
- Consider the  $u \leftrightarrow s$  symmetry “V-spin” which has the associated  $s \rightarrow u$  ladder operator

$$V_+ = \frac{1}{2}(\lambda_4 + i\lambda_5) = \frac{1}{2} \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} + \frac{i}{2} \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

with

$$V_+ s = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = +u$$

- ★ The effects of the six ladder operators are:

$T_+ d = u;$	$T_- u = d;$	$T_+ \bar{u} = -\bar{d};$	$T_- \bar{d} = -\bar{u}$
$V_+ s = u;$	$V_- u = s;$	$V_+ \bar{u} = -\bar{s};$	$V_- \bar{s} = -\bar{u}$
$U_+ s = d;$	$U_- d = s;$	$U_+ \bar{d} = -\bar{s};$	$U_- \bar{s} = -\bar{d}$

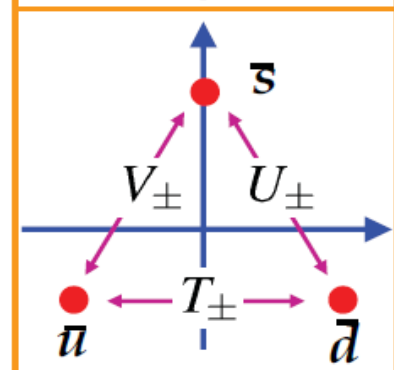
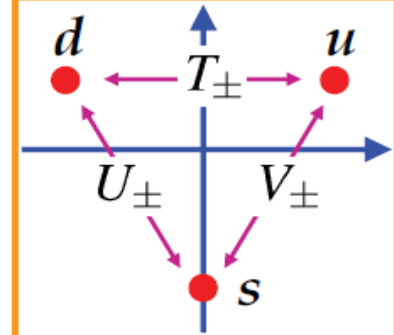
all other combinations give zero

## SU(3) LADDER OPERATORS

$$T_{\pm} = \frac{1}{2}(\lambda_1 \pm i\lambda_2)$$

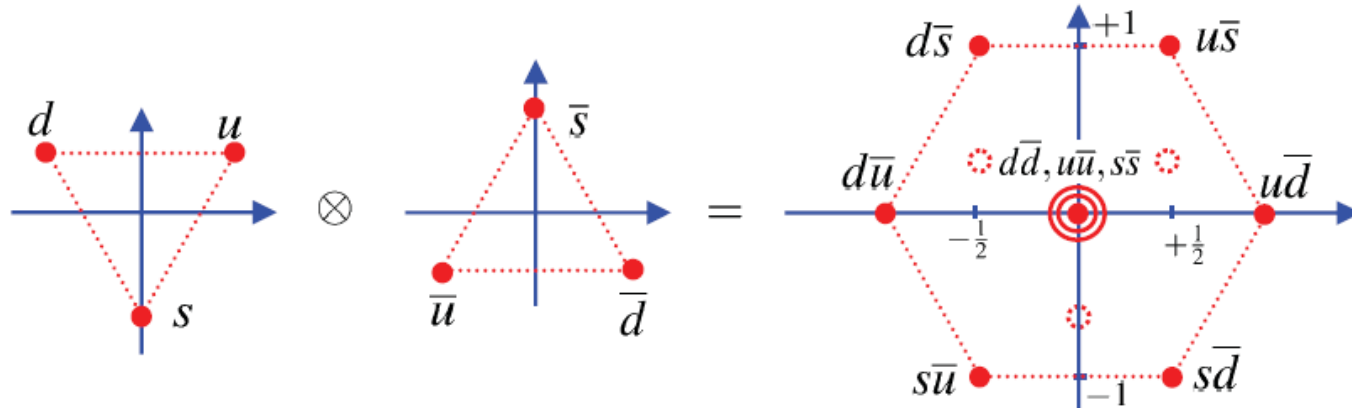
$$V_{\pm} = \frac{1}{2}(\lambda_4 \pm i\lambda_5)$$

$$U_{\pm} = \frac{1}{2}(\lambda_6 \pm i\lambda_7)$$

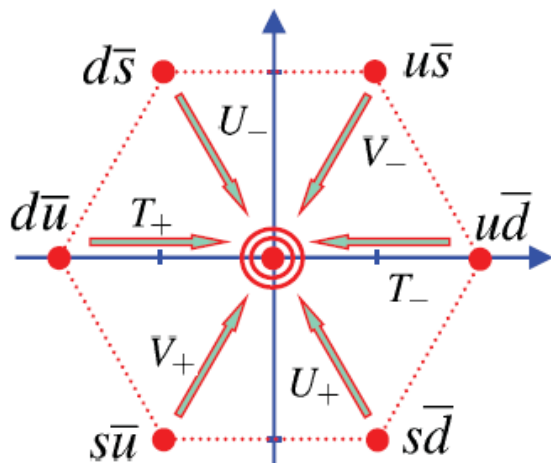


# Light (uds) Mesons

- Use ladder operators to construct uds mesons from the nine possible  $q\bar{q}$  states



- The three central states, all of which have  $Y = 0$ ;  $I_3 = 0$  can be obtained using the ladder operators and orthogonality. Starting from the outer states can reach the centre in six ways



$$\begin{aligned}
 T_+ |d\bar{u}\rangle &= |u\bar{u}\rangle - |d\bar{d}\rangle & T_- |u\bar{d}\rangle &= |d\bar{d}\rangle - |u\bar{u}\rangle \\
 V_+ |s\bar{u}\rangle &= |u\bar{u}\rangle - |s\bar{s}\rangle & V_- |u\bar{s}\rangle &= |s\bar{s}\rangle - |u\bar{u}\rangle \\
 U_+ |s\bar{d}\rangle &= |d\bar{d}\rangle - |s\bar{s}\rangle & U_- |d\bar{s}\rangle &= |s\bar{s}\rangle - |d\bar{d}\rangle
 \end{aligned}$$

- Only **two** of these six states are linearly independent.
- But there are **three** states with  $Y = 0$ ;  $I_3 = 0$
- Therefore one state is not part of the same multiplet, i.e. cannot be reached with ladder ops.

- First form two linearly independent orthogonal states from:

$$|u\bar{u}\rangle - |d\bar{d}\rangle \quad |u\bar{u}\rangle - |s\bar{s}\rangle \quad |d\bar{d}\rangle - |s\bar{s}\rangle$$

- ★ If the SU(3) flavour symmetry were exact, the choice of states wouldn't matter. However,  $m_s > m_{u,d}$  and the symmetry is only approximate.

- **Experimentally** observe three light mesons with  $m \sim 140$  MeV:  $\pi^+$ ,  $\pi^0$ ,  $\pi^-$
- **Identify one state** (the  $\pi^0$ ) with the isospin triplet (derived previously)

$$\psi_1 = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d})$$

- The second state can be obtained by taking the linear combination of the other two states which is orthogonal to the  $\pi^0$

$$\psi_2 = \alpha(|u\bar{u}\rangle - |s\bar{s}\rangle) + \beta(|d\bar{d}\rangle - |s\bar{s}\rangle)$$

with orthonormality:  $\langle \psi_1 | \psi_2 \rangle = 0$ ;  $\langle \psi_2 | \psi_2 \rangle = 1$

$$\rightarrow \psi_2 = \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s})$$

- The final state (which is not part of the same multiplet) can be obtained by requiring it to be orthogonal to  $\psi_1$  and  $\psi_2$

$$\rightarrow \psi_3 = \frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s})$$

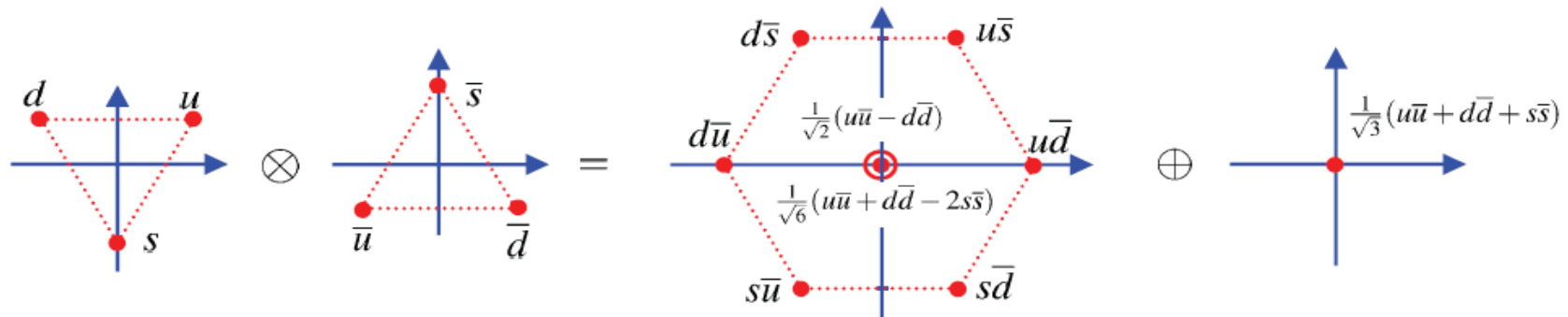
**SINGLET**

- ★ It is easy to check that  $\psi_3$  is a singlet state using ladder operators

$$T_+ \psi_3 = T_- \psi_3 = U_+ \psi_3 = U_- \psi_3 = V_+ \psi_3 = V_- \psi_3 = 0$$

which confirms that  $\psi_3 = \frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s})$  is a “flavourless” singlet

- Therefore the combination of a quark and anti-quark yields nine states which breakdown into an **OCTET** and a **SINGLET**



- In the language of group theory:  $3 \otimes \bar{3} = 8 \oplus 1$

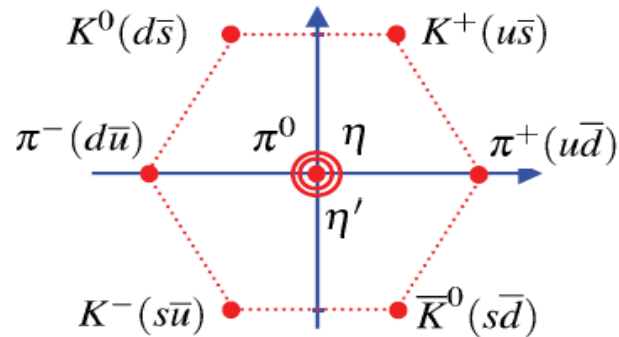
- ★ Compare with combination of two spin-half particles  $2 \otimes 2 = 3 \oplus 1$

**TRIPLET** of spin-1 states:  $|1, -1\rangle, |1, 0\rangle, |1, +1\rangle$

spin-0 **SINGLET**:  $|0, 0\rangle$

- These spin triplet states are connected by ladder operators just as the meson *uds* octet states are connected by **SU(3)** flavour ladder operators
- The singlet state carries no angular momentum - in this sense the **SU(3) flavour singlet** is “flavourless”

## PSEUDOSCALAR MESONS (L=0, S=0, J=0, P= -1 )



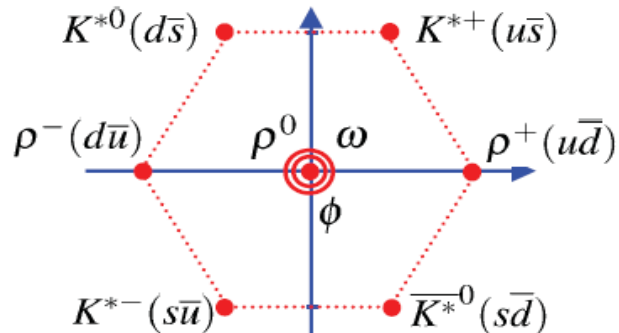
- Because SU(3) flavour is only approximate the physical states with  $I_3 = 0, Y = 0$  can be mixtures of the octet and singlet states.

**Empirically find:**

$$\begin{aligned}\pi^0 &= \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \\ \eta &\approx \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s}) \\ \eta' &\approx \frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s})\end{aligned}$$

← **singlet**

## VECTOR MESONS (L=0, S=1, J=1, P= -1 )



- For the vector mesons the physical states are found to be approximately **“ideally mixed”**:

$$\begin{aligned}\rho^0 &= \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \\ \omega &\approx \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) \\ \phi &\approx s\bar{s}\end{aligned}$$

## MASSSES

$\pi^\pm$ : 140 MeV	$\pi^0$ : 135 MeV
$K^\pm$ : 494 MeV	$K^0/\bar{K}^0$ : 498 MeV
$\eta$ : 549 MeV	$\eta'$ : 958 MeV

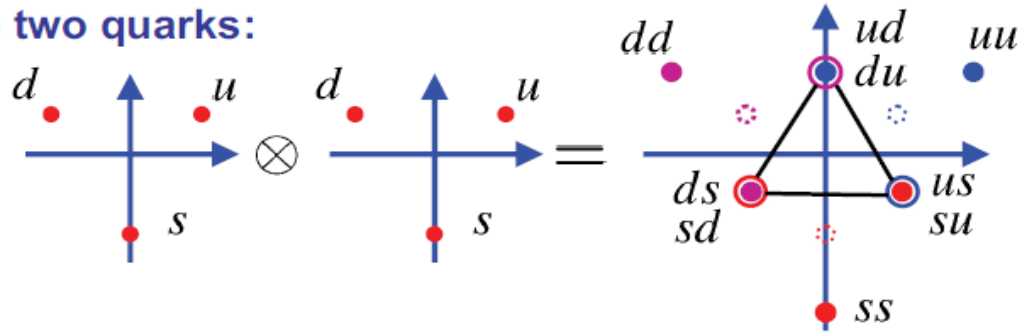
$\rho^\pm$ : 770 MeV	$\rho^0$ : 770 MeV
$K^{*\pm}$ : 892 MeV	$K^{*0}/\bar{K}^{*0}$ : 896 MeV
$\omega$ : 782 MeV	$\phi$ : 1020 MeV

# Combining uds Quarks to form Baryons

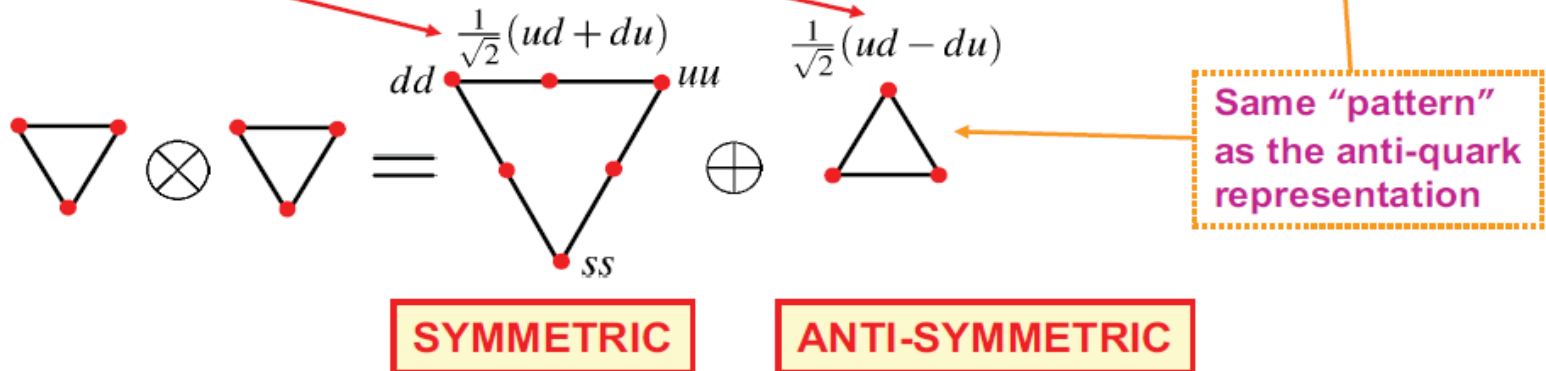
★ Have already seen that constructing Baryon states is a fairly tedious process when we derived the proton wave-function. Concentrate on multiplet structure rather than deriving all the wave-functions.

★ Everything we do here is relevant to the treatment of colour

• First combine two quarks:

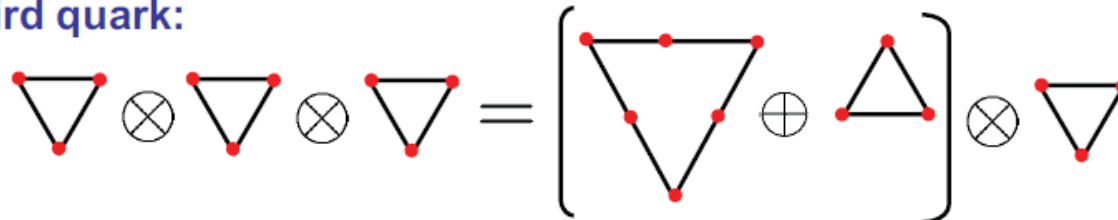


★ Yields a symmetric sextet and anti-symmetric triplet:  $3 \otimes 3 = 6 \oplus \bar{3}$



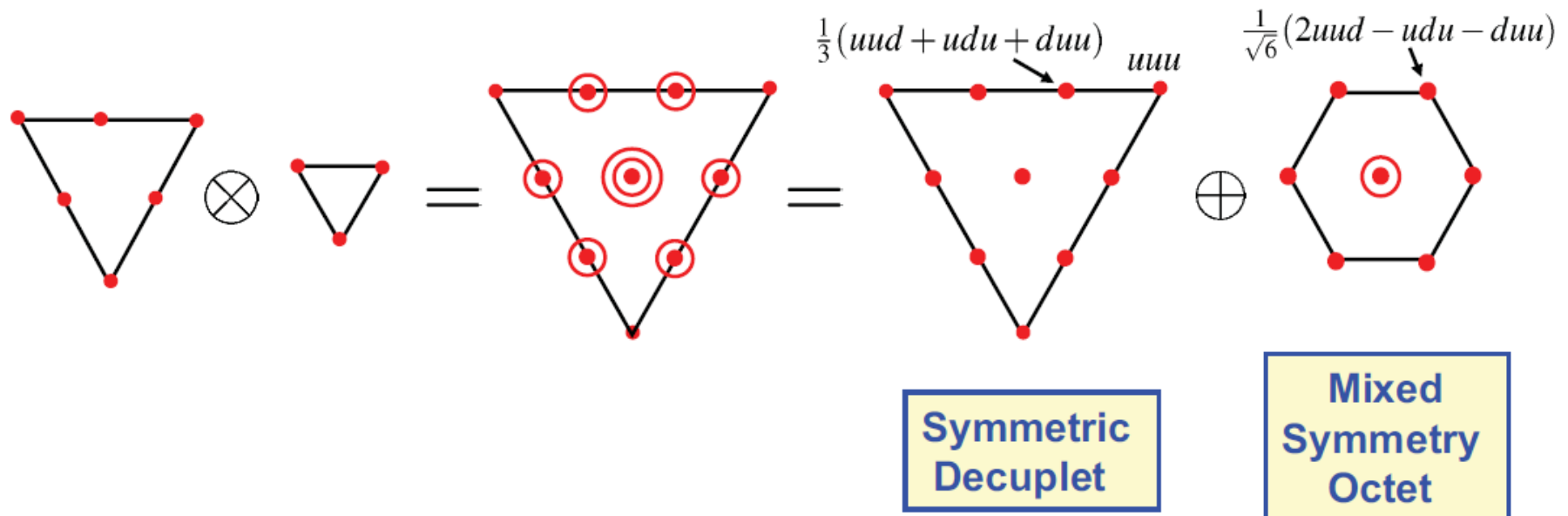


- Now add the third quark:



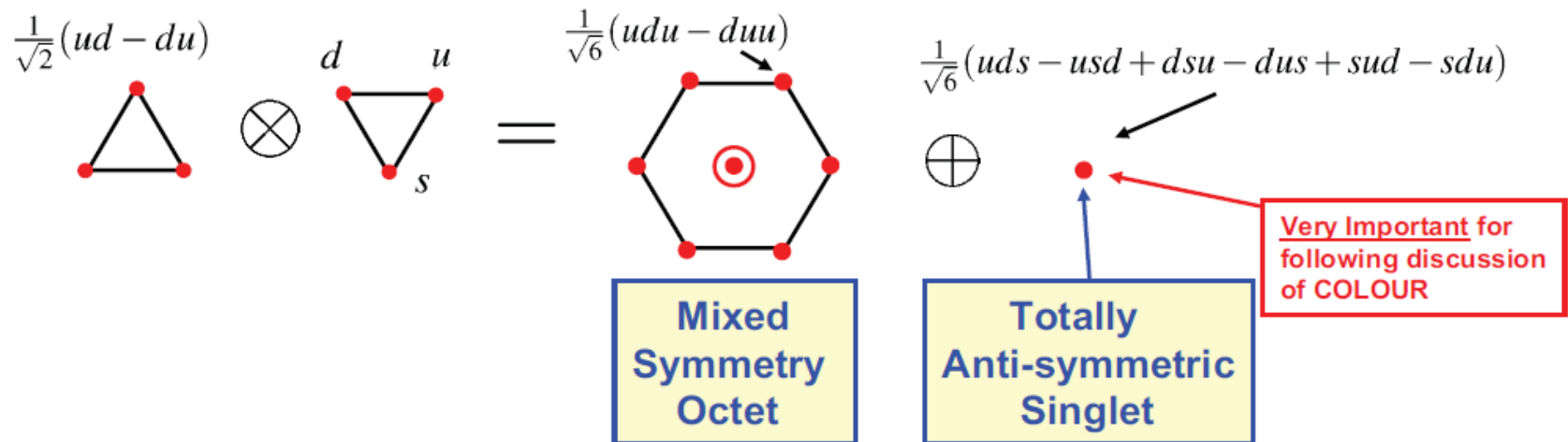
- Best considered in two parts, building on the **sextet** and **triplet**. Again concentrate on the multiplet structure (for the wave-functions refer to the discussion of proton wave-function).

- 1 Building on the sextet:  $3 \otimes 6 = 10 \oplus 8$



## 2 Building on the triplet:

- Just as in the case of  $uds$  mesons we are combining  $\bar{3} \times 3$  and again obtain an octet and a singlet



- Can verify the wave-function  $\Psi_{\text{singlet}} = \frac{1}{\sqrt{6}}(uds - usd + dsu - dus + sud - sdu)$  is a singlet by using ladder operators, e.g.

$$T_+ \Psi_{\text{singlet}} = \frac{1}{\sqrt{6}}(uus - usu + usu - uus + suu - suu) = 0$$

- ★ In summary, the combination of three  $uds$  quarks decomposes into

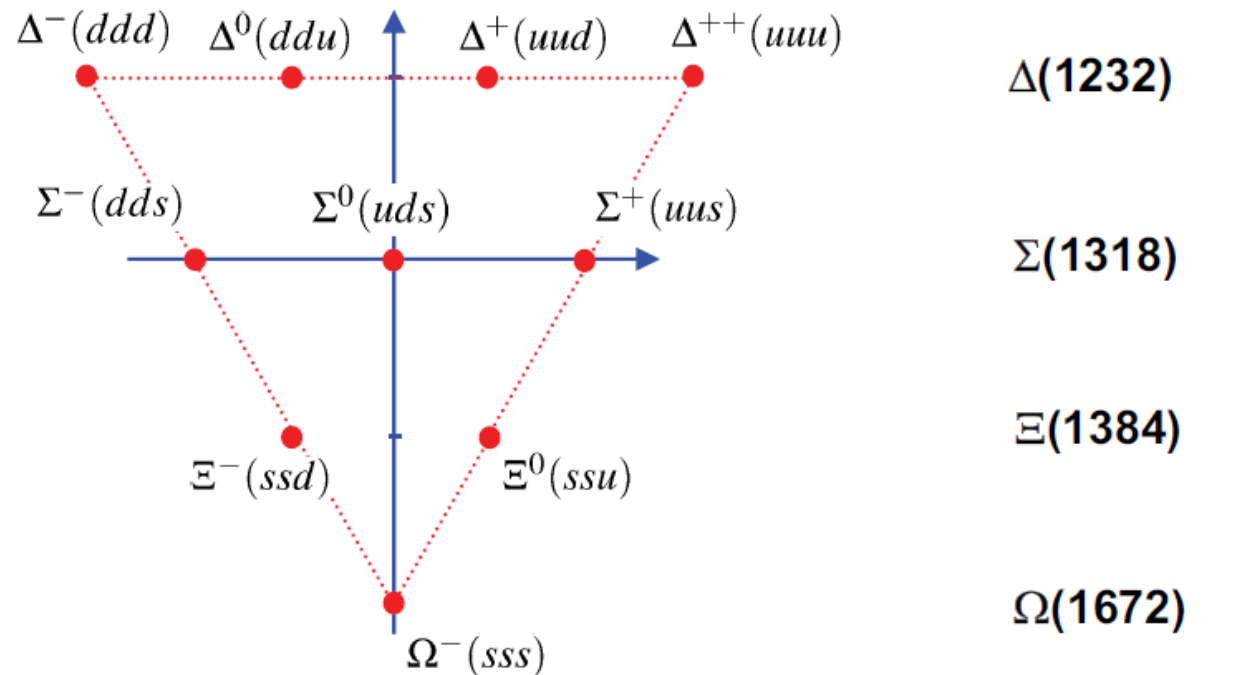
$$3 \otimes 3 \otimes 3 = 3 \otimes (6 \oplus \bar{3}) = 10 \oplus 8 \oplus 8 \oplus 1$$

# Baryon Decuplet

★ The baryon states ( $L=0$ ) are:

- the **spin 3/2 decuplet** of symmetric flavour and symmetric spin wave-functions  $\phi(S)\chi(S)$

**BARYON DECUPLET** ( $L=0$ ,  $S=3/2$ ,  $J=3/2$ ,  $P=+1$ )



★ If SU(3) flavour were an exact symmetry all masses would be the same (broken symmetry)

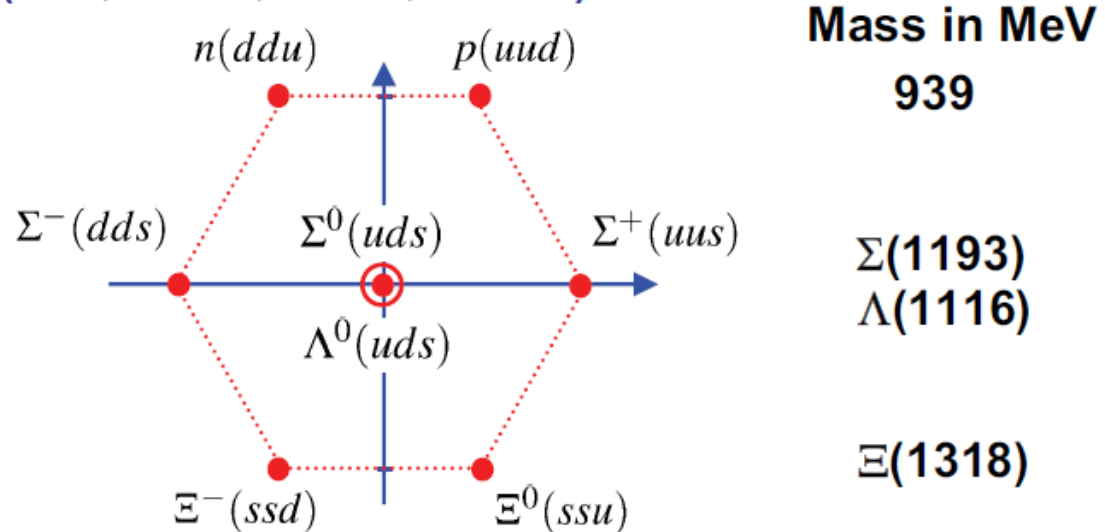
# Baryon Octet

- ★ The **spin 1/2 octet** is formed from mixed symmetry flavour and mixed symmetry spin wave-functions

$$\alpha\phi(M_S)\chi(M_S) + \beta\phi(M_A)\chi(M_A)$$

See previous discussion proton for how to obtain wave-functions

**BARYON OCTET** ( $L=0$ ,  $S=1/2$ ,  $J=1/2$ ,  $P=+1$ )



- ★ **NOTE:** Cannot form a totally symmetric wave-function based on the anti-symmetric flavour singlet as there no totally anti-symmetric spin wave-function for 3 quarks

# Summary

- ★ Considered SU(2) **ud** and SU(3) **uds** flavour symmetries
- ★ Although these flavour symmetries are only approximate can still be used to explain observed multiplet structure for mesons/baryons
- ★ In case of SU(3) flavour symmetry results, e.g. predicted wave-functions should be treated with a pinch of salt as  $m_s \neq m_{u/d}$
- ★ Introduced idea of singlet states being “spinless” or “flavourless”
- ★ In the next handout apply these ideas to colour and QCD...

# Appendix

- Define anti-quark doublet  $\bar{q} = \begin{pmatrix} -\bar{d} \\ \bar{u} \end{pmatrix} = \begin{pmatrix} -d^* \\ u^* \end{pmatrix}$

- The quark doublet  $q = \begin{pmatrix} u \\ d \end{pmatrix}$  transforms as  $q' = Uq$

$$\begin{pmatrix} u' \\ d' \end{pmatrix} = U \begin{pmatrix} u \\ d \end{pmatrix} \xrightarrow[\text{conjugate}]{\text{Complex}} \begin{pmatrix} u'^* \\ d'^* \end{pmatrix} = U^* \begin{pmatrix} u^* \\ d^* \end{pmatrix}$$

- Express in terms of anti-quark doublet

$$\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \bar{q}' = U \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \bar{q}$$

- Hence  $\bar{q}$  transforms as

$$\bar{q}' = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} U^* \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \bar{q}$$

- In general a 2x2 unitary matrix can be written as

$$U = \begin{pmatrix} c_{11} & c_{12} \\ -c_{12}^* & c_{11}^* \end{pmatrix}$$

- Giving

$$\begin{aligned} \bar{q}' &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} c_{11}^* & c_{12}^* \\ -c_{12} & c_{11} \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \bar{q} \\ &= \begin{pmatrix} c_{11} & c_{12} \\ -c_{12}^* & c_{11}^* \end{pmatrix} \\ &= U\bar{q} \end{aligned}$$

- Therefore the anti-quark doublet  $\bar{q} = \begin{pmatrix} -\bar{d} \\ \bar{u} \end{pmatrix}$

transforms in the same way as the quark doublet  $q = \begin{pmatrix} u \\ d \end{pmatrix}$

- ★ **NOTE:** this is a special property of SU(2) and for SU(3) there is no analogous representation of the anti-quarks