## From Skyrmions to holographic exotics

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## Dedicated to Michał Praszałowicz, on the occasion of His 70th birthday



- 1974 Close encounter with Michał
- 1977-1983 Period of unlimited appreciation of pQCD
- 1983-1985 Lost illusions, farewell to pQCD
- [MAN, M Praszałowicz, W Słomiński; A practical guide to the next-to-leading order of the perturbation expansion in QCD, Annals of Physics 166, 433, 1986]
- 1983 Revelation: M. Peskin brings preprints by E. Witten on Skyrme model to Zakopane School

## Skyrmion revisited

• Skyrme idea (1961): **Baryon as topological soliton** (isospin-space hedgehog) of the mesonic (pionic) fields

$$L = L_{\sigma} + L_{Skyrme}$$

- Witten, Adkins, Nappi (1982-83):
  - Link to large N
  - Anomalous term

$$L = L_{\sigma} + L_{Skyrme} + L_{WZW}$$

- Quantization of collective modes from symmetries (moduli space) leads to

$$H = M_0 + \frac{\vec{I}^2}{2\Omega_{sol}}$$

where  $\vec{l} (= \vec{J})$  are isospin (angular momentum) respectively. (Technically, equivalent to algebraic hydrogen atom quantization done W. Pauli in 1921 (!))

[PO Mazur, MAN, M Praszałowicz, SU (3) extension of the Skyrme model, Physics Letters B 147, 137, 1984]

- Quantization of the symmetric top with some local symmetry included
- States superselected by **triality zero** condition for representations of SU(3) (octet, decuplet)
- Challenges: third triality zero representation is antidecuplet: Michał notes that leads to a light pentaquark. Consequences for elusive Θ<sup>+</sup>.
- Challenges: Explicit chiral symmetry breaking does not work.

- Already at the level of strange quarks symmetry breaking is so bad, so collective method fails.
- Proper way is to use Born-Oppenheimer approximation fast vibration of the kaon in the SU(2) solitonic background.
   Then slow rotation of the bound state happens in the presence of non-Abelian Berry phase originating from kaons.

$$H = M_0 + \text{binding} + \frac{(\vec{J} - (1 - c_k)tr(\vec{K}\vec{I}\vec{K}^{\dagger}))^2}{2\Omega_{sol}}$$

([isospin-spin transmutation in Callan-Klebanov 1985])

• Bound kaon behaves as s-quark, modulo baryon number

 For c(b) baryons, similar picture, but with two Berry phases from D and D\*, so one gets

$$H_{1} = \frac{[(\vec{J} - \vec{S}_{H}) - (1 - c_{D})tr(D\vec{I}D^{\dagger}) - (1 - c_{D^{*}})tr(D^{*}\vec{I}D^{*\dagger})]^{2}}{2\Omega_{sol}}$$

• In the infinitely heavy mass Berry phases *exactly* **cancel** Then

$$\mathcal{H}_1 = rac{(ec{J}-ec{S}_H)^2}{2\Omega_{sol}} = rac{ec{I}^2}{2\Omega_{sol}}$$

(Realization of Isgur-Wise symmetry at the baryonic level)

• Also, soliton can capture more than one meson (double heavy baryons, exotica).

Combining chiral symmetry with heavy-spin symmetry leads to novel feature [MAN, Rho, Zahed (1992), Bardeen-Hill (1993)] Both symmetries enforce the **presence of opposite parity**  $(0^+, 1^+)$  multiplet  $G = \frac{1+\psi}{2} (\tilde{D} + \gamma^{\mu} \gamma_5 \tilde{D}^*_{\mu})$ in addition to standard  $(0^-, 1^-)$  one  $H = \frac{1+\psi}{2} (\gamma_5 D + \gamma^{\mu} D^*_{\mu})$ 

- Consequence of chiral symmetry  $[{\it v},\gamma_5]_+=0$
- Doublers communicate only through axial current
- Physical split in axial couplings and masses,  $m_G - m_H \sim O(\Sigma_I) \sim 350 MeV$
- Chiral doublers do not double the number of states in quark model, but *reorganize* them in different way.

(i) soliton captures H meson (heavy baryon) (ii) soliton captures G meson (doubler of heavy baryon) (iii) soliton captures  $\overline{H}$  (heavy pentaquark) (iv) soliton captures  $\overline{G}$ (doubler of heavy pentaquark) (v) soliton captures more mesons.....

- Short, unhappy life of  $\tilde{\Theta}_c(3099)$
- [MAN, M. Praszałowicz, M. Sadzikowski, J. Wasiluk; Chiral doublers of heavy-light baryons, Phys. Rev. D70, 031502, 2004]

... In particular, we interpret the state recently reported by the H1 experiment at HERA as a chiral partner  $\tilde{\Theta}_c(3099)$  of yet undiscovered ground state pentaquark  $\Theta_c(2700)$ .

## Desperately seeking exotics ...



## Experimental revolution: Post-Babar-ian era

Abundance of exotic heavy-light particles

- **cdus** :*X*(2866), *X*<sub>1</sub>(2904)
- $c\bar{c}q\bar{q}$  :  $\chi_{c1}(3872)$
- cc̄uđ:  $Z_c(3900)$ ,  $Z_c(4020)$ ,  $Z_c(4050)$ , X(4100),  $Z_c(3985)$ ,  $Z_c(4430)$ ,  $R_{c0}(4240)$
- ccus: Z<sub>cs</sub>(3985), Z<sub>cs</sub>(4000), Z<sub>cs</sub>(4220)
- **bbud**: Z<sub>b</sub>(10610), Z<sub>b</sub>(10650)
- **cccc**: *X*(6900)
- ccūd:  $T_{cc}^+(3875)$ , also  $T_{cc}^0$ ,  $T_{cc}^{++}$  (preliminary)
- Pentaquarks cc̄uud: P<sub>c</sub>((4380), (4450) → [(4440), (4457)], (4312)], P<sub>c</sub>(4337) (3σ) significance
- cc̄uds: *P*<sub>cs</sub>(4459)
- possibilities of further heavy-light "chemistry": many more expected (?)

Input from lattice

# Theoretical revolution: Gravity/Gauge duality (holography, AdS/CFT)

- In the 70' QCD became fundamental theory of quarks and gluons, strings (flux tubes) appear as effective, e.g. Lund model
- Maldacena pointed that gauge theory is 4 dim is equivalent to string theory in higher dimensions: Towards two fundamental theories of strong interactions (?!)
- Various versions conformal window, lower-dimensions (solid state physics)
- Witten (1998) applied duality to QCD: pure YM in 3 + 1 at large N and  $\lambda = g_{YM}^2 N$ Surprising similarity to spectrum of glueballs at large N lattice

## Sakai-Sugimoto breakthrough 2005

- Adding  $N_f$  massless fermions geometric SB $\chi$ S
- Low energy limit  $S = S_{YM} + S_{CS}$  where  $S_{YM} \sim \int d^4x dz \operatorname{Tr} \left(\frac{1}{2}k(z)^{-1/3}F_{\mu\nu}^2 + k(z)F_{\mu z}^2\right)$  where  $k(z) = 1 + z^2$
- Mode expansion:  $A_{\mu}(x^{\mu}, z) = \sum_{n} B_{\mu}^{(n)}(x_{\mu}) \Psi_{n}(z)$ ,  $A_{5}(x^{\mu}, z) = \sum_{n} \phi^{(n)}(x^{\mu}) \Phi_{n}(z)$
- Keeping only  $\phi^{(0)}$  yields  $L = L_{\sigma} + L_{Skyrme} + L_{WZW}$ . Adding  $B^{(1)}_{\mu} \sim \rho$  and  $B^{(2)}_{\mu} \sim a_1$  yield hidden gauge model
- Successful phenomenology with very few parameters

- 4 dim pion  $\rightarrow$  Skyrmion (static solution)
- 5 dim gauge field  $\rightarrow$  **BPST** instanton in  $x_1, x_2, x_3, z$  in flavor
- Topological number  $\equiv$  baryon number
- Direct realization of 1989 Atiyah-Manton idea  $U(\vec{x}) = P \exp(i \int dz A_z(\vec{x}, z))$
- 8 zero modes lead to moduli space quantization  $M = M_0 + \left(\sqrt{\frac{(l+1)^2}{6} + \frac{2}{15}N^2} + \sqrt{\frac{2}{3}}(n_\rho + n_Z + 1)\right)M_{KK}$ where l = 2I = 2J = 1, 3, 5...

## HL baryons in holographic scenario: Liu, Zahed, 2017

• CK-like scheme with heavy-spin symmetry

• 
$$M = M_0 + (N_Q + N_{\bar{Q}})m_H + (\sqrt{\frac{(l+1)^2}{6} + \frac{2}{15}N^2(1 - \frac{15(N_Q - N_{\bar{Q}})}{4N} + \frac{5(N_Q - N_{\bar{Q}})^2}{3N^2})^2})M_{KK} + \sqrt{\frac{2}{3}}(n_\rho + n_Z + 1))M_{KK}$$

• Various combinations of q-numbers give all types of HL hadrons

$$N_Q = 1, N_{\bar{Q}} = 0$$
 yield **hll**  
 $N_Q = 2, N_{\bar{Q}} = 0$  yield **hhl**  
 $N_Q = N_{\bar{Q}} = 1$  yield pentaquarks **hhlll**  
 $n_Z \neq 0$  yield excited (Roper-like),  $n_\rho \neq 0$  yield odd parity

• Three parameters:  $M_{KK}, M_0, m_H \sim M_D(M_B)$ 

# Adding spin effect (subleading in $m_H^{-1}$ ) Liu, MAN, Zahed, 2021

- 3 parameters,  $M_0 
  ightarrow m_N, \ M_{KK} 
  ightarrow m_{\Lambda_c}, \ m_H \sim M_D(M_B)$  for c(b)
- 3 pentaquarks  $\frac{1}{2}, \frac{1}{2}^{-}(S = 1), \frac{1}{2}, \frac{1}{2}^{-}(S = 0), \frac{1}{2}, \frac{3}{2}^{-}(S = 1), IJ^{\pi}$  $(\frac{1}{2}, \frac{5}{2}^{\pm} \text{ ruled out})$ , consistent with  $P_c(4312, 4440, 4457)$ [LHCb]
- Recently reported  $P_c(4337)$  at  $3\sigma$  significance is not supported
- Open and hidden decay widths (Liu, MAN, Zahed, 2021) e.g.  $P_c \rightarrow \Lambda_c + \overline{D}$ ,  $\Gamma(S = 0, J = \frac{1}{2}) : \Gamma(S = 1, J = \frac{1}{2}) : \Gamma(S = 1, J = \frac{3}{2}) = \frac{1}{2} : \frac{5}{6} : \frac{1}{3}$
- Formfactors (Liu, Mamo, MAN, Zahed; 2021), consistent with recent GLUEX results on  $\gamma p \rightarrow (P_c^+) \rightarrow J/\psi p$

## Tetraquark puzzle

- **hhll** several predictions (positive/negative ±200 *MeV*)
- Measurement of Ξ<sup>++</sup><sub>cc</sub>(3621) fixed the normalization for bbūd̄ tetraquark (Karliner, Rosner (2017); Eichten, Quigg (2017)), bound up

to 200*MeV* (!)

- ${\cal T}^+_{cc}$  (01<sup>+</sup>) (LHCb), narrow, bounded at  $-360 \, keV$ ,  $\Gamma \sim 50 \, keV$
- Holographic picture: instanton-antiinstanton "molecule" binds two mesons (Liu, MAN, Zahed (2019)), b(c) tetraquarks bounded by 80(40) MeV
- Normalizing mass to  $T_{cc}^+$ , (Liu, MAN, Zahed (2022)) predict mass of  $T_{bc}$  and  $T_{bb}$ , and estimate very narrow width.
- Stronger bindings in chiral quark soliton model [Michał Praszałowicz, 2023-2024]

- Strongly coupled QCD could be approached via duality from string theory in large N and large  $\lambda$  limit, including spectra of heavy-light hadrons
- Few parameters and very restrictive predictions, so models are confutable
- Approach based on confinement, SB $\chi$ S and heavy spin symmetry, in the limit of large N and  $\lambda$ .
- "High brow" theory boils down to relatively simple QM in moduli space (top-down approach)
- Astonishing and deep analogies to "old" physics, including joint work with Michał!

## Happy Birthday, Michał !!!!!!