Exploring the QCD phase diagram with fluctuations and correlations

"A theory is something nobody believes, except the person who made it.

An experiment is something everybody believes, except the person who made it."

A. Einstein

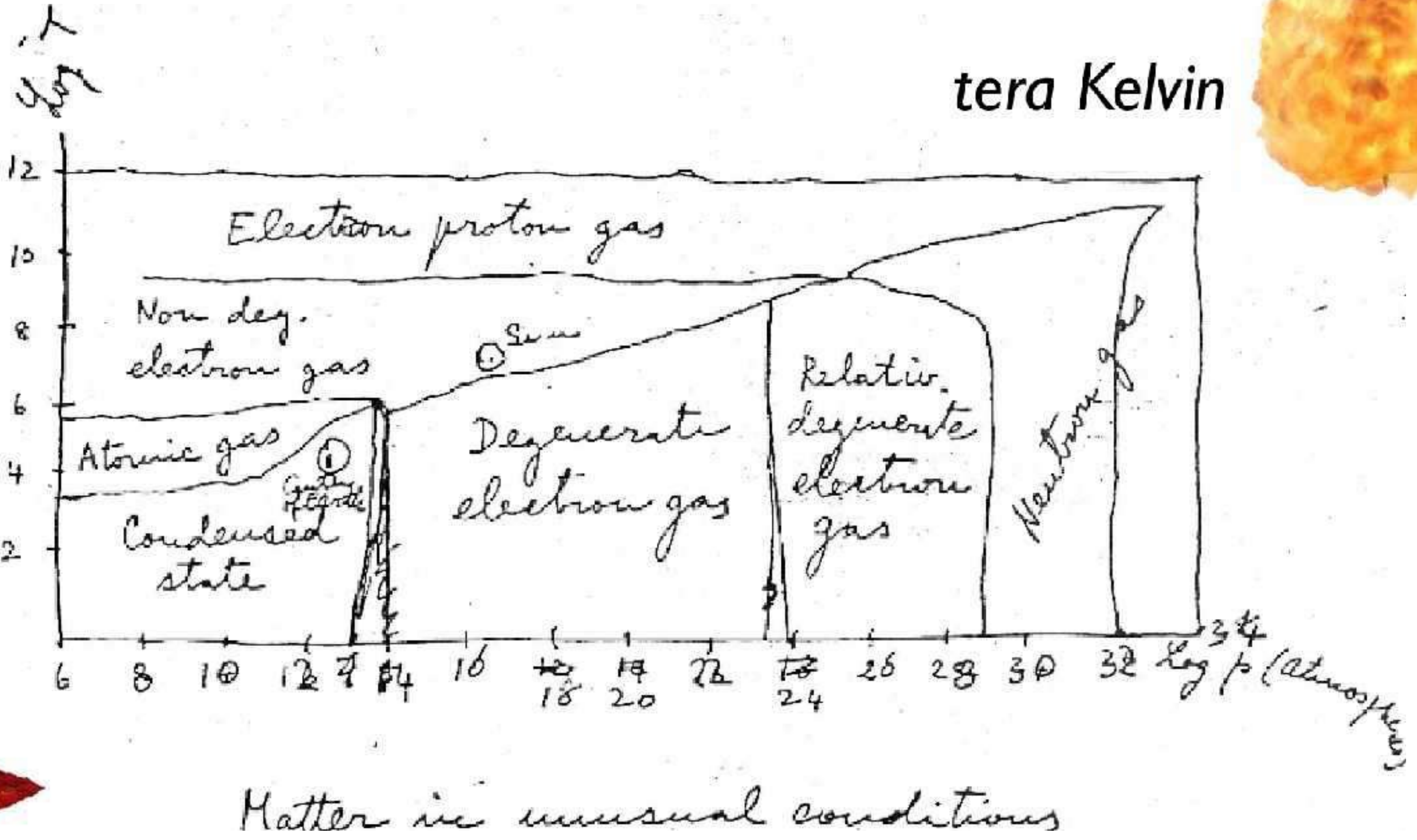
Thanks to:
D. Oliinychenko, A. Sorensen, J. Steinheimer, V. Vovchenko

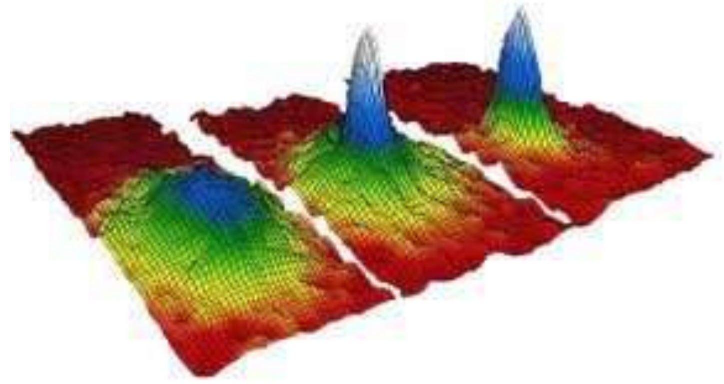
EMMI

An old question

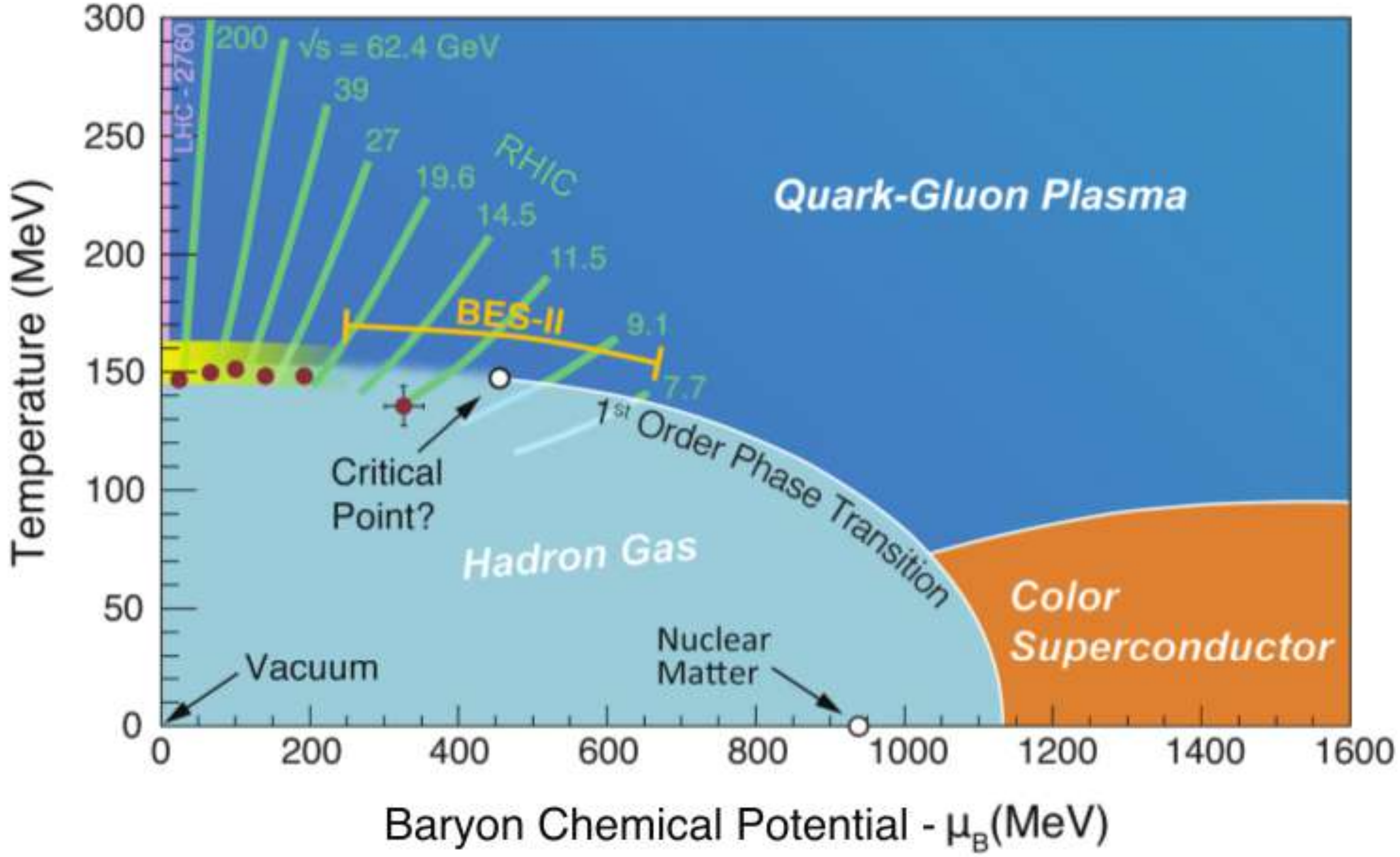
Fermi 1953

tera Kelvin





The phase diagram



Increase chemical potential by lowering the beam energy

In reality, we add baryons (nucleons) from target and projectile to mid-rapidity

What we know about the Phase Diagram

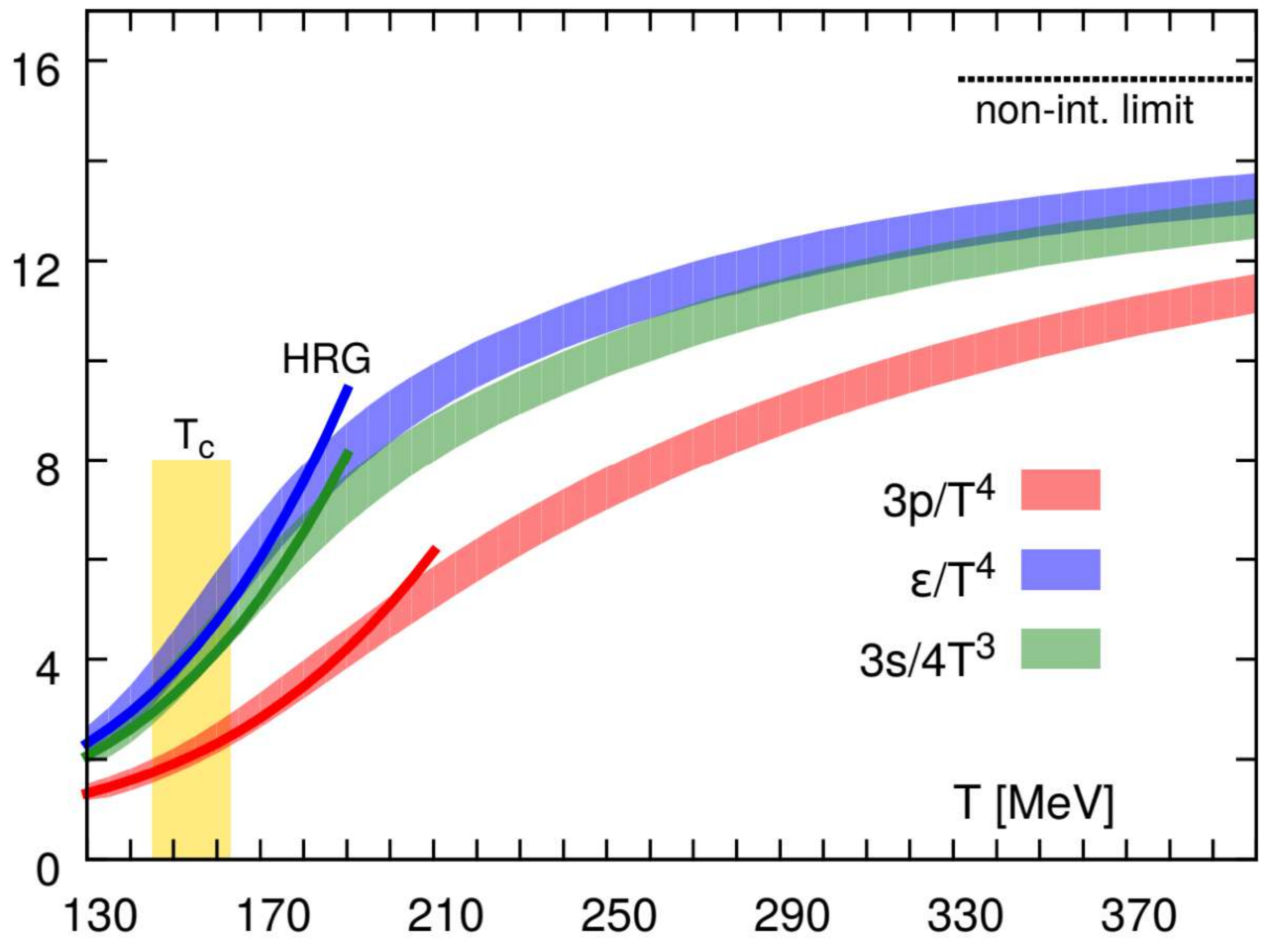
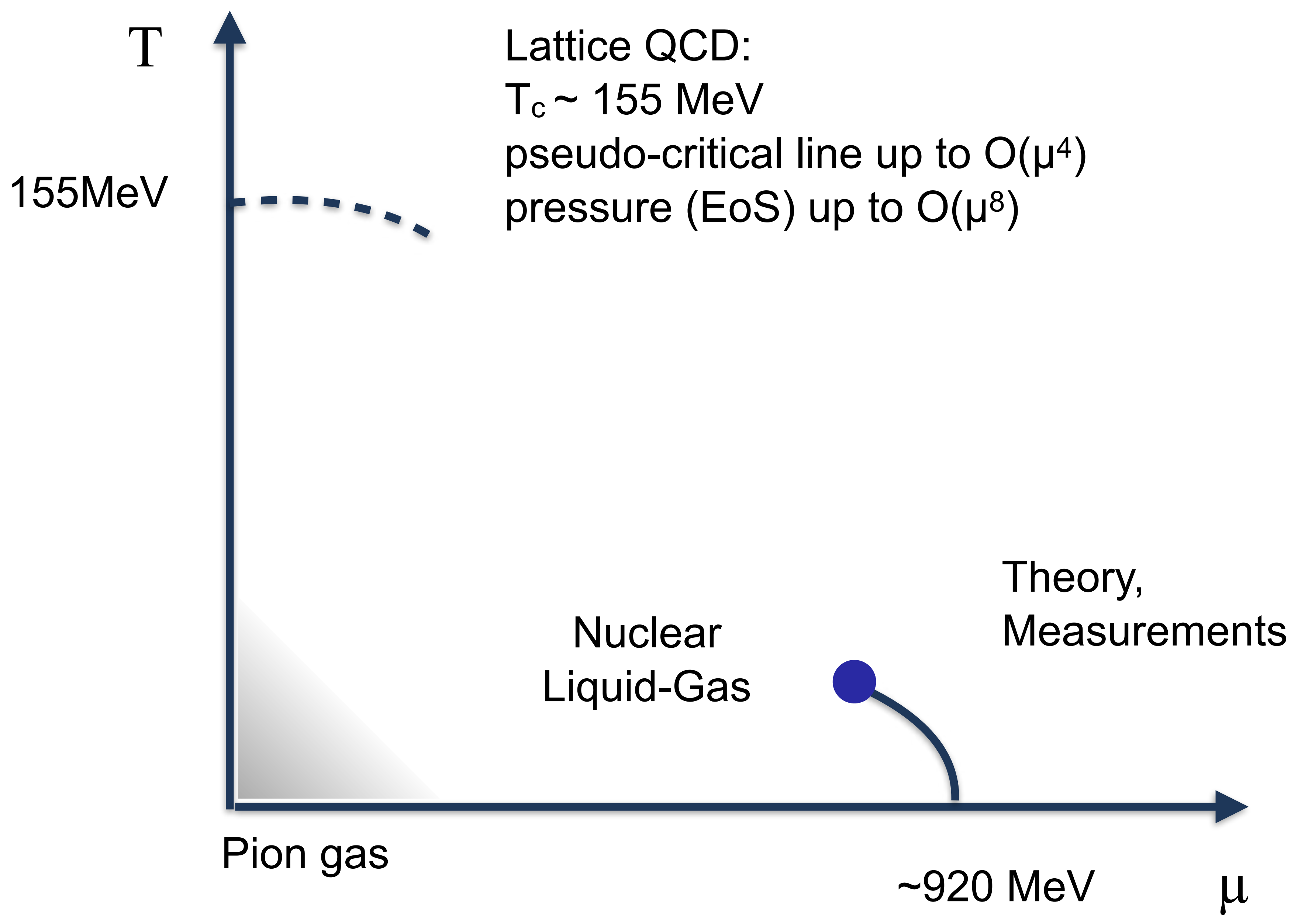
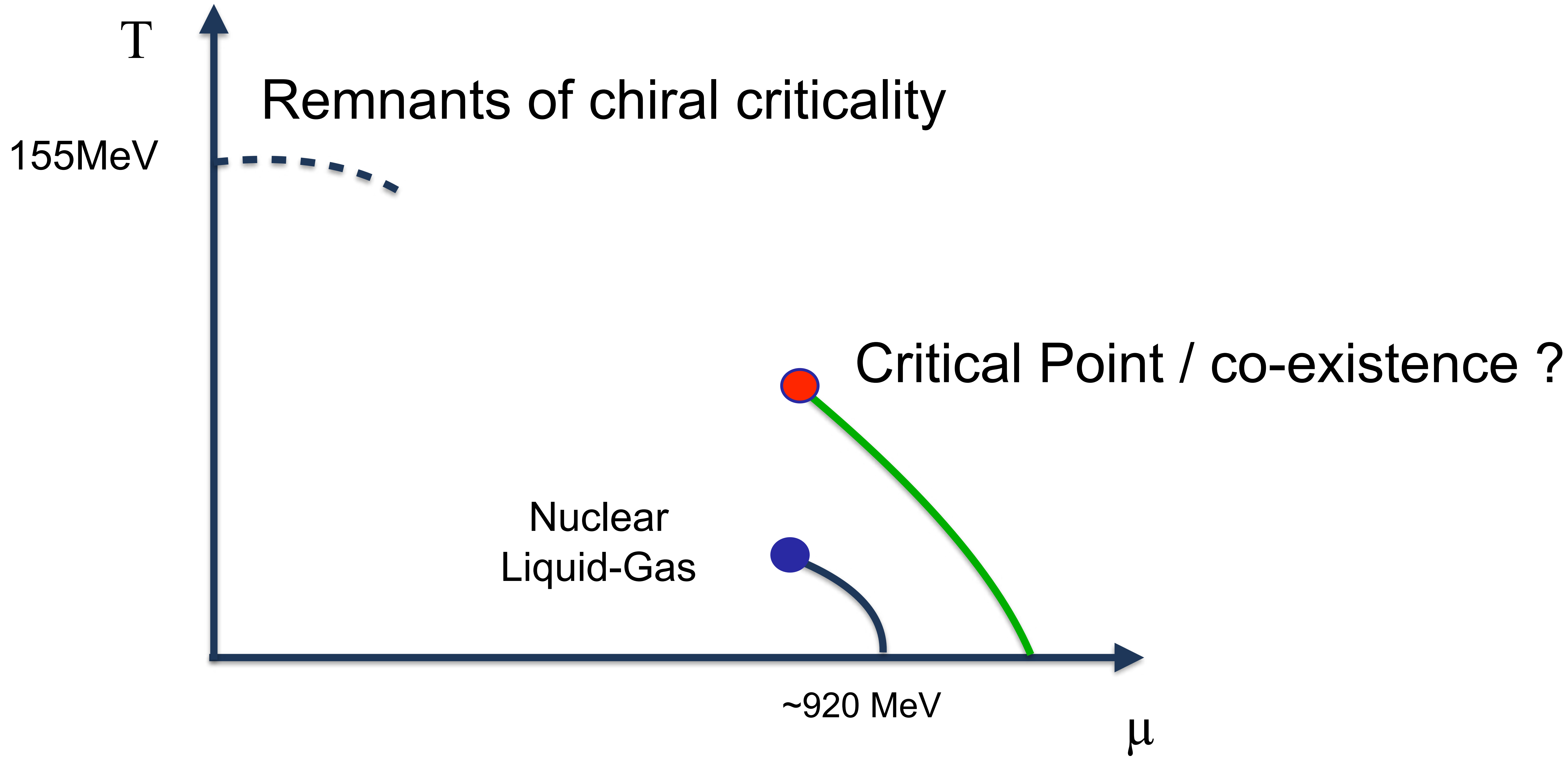


Figure from HotQCD coll., PRD ‘14



What we are looking for



We are dealing with small system of finite lifetime

NO real singularities!

Where is the critical point? Ask AI

UNIVERSITY OF HOUSTON

ChatGPT struggles

Is there a QCD critical point and where is it?

As of my knowledge cutoff in September 2021, the existence and precise location of a critical point in the phase diagram of Quantum Chromodynamics (QCD), the theory of the strong interaction between quarks and gluons, was not definitively established.

...

However, finding the QCD critical point using lattice QCD is a particularly challenging problem due to the so-called "sign problem." The sign problem arises when trying to extend lattice QCD calculations from zero to non-zero baryon chemical potential, which is necessary to locate the critical point on the phase diagram. At non-zero chemical potential, the action of the QCD path integral becomes complex, which makes the standard Monte Carlo techniques used in lattice QCD simulations ineffective.

...

Experimental searches for the QCD critical point in heavy-ion collisions generally rely on identifying fluctuations and correlations in the properties of the particles produced in the collisions, which are expected to exhibit unique features near the critical point.

Courtesy of V. Vovchenko

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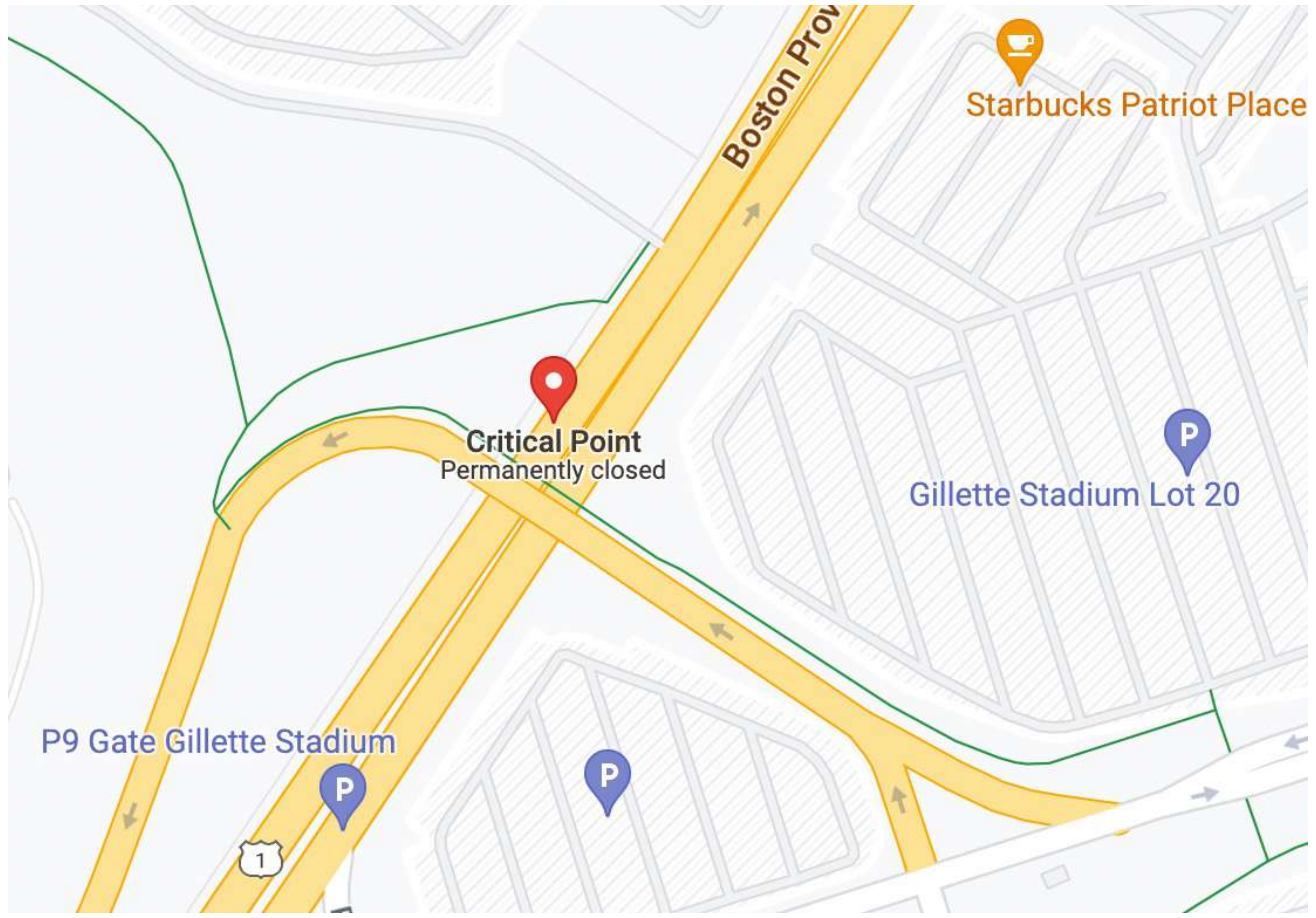
No issues for Google



Courtesy of V. Vovchenko

Phase Transitions

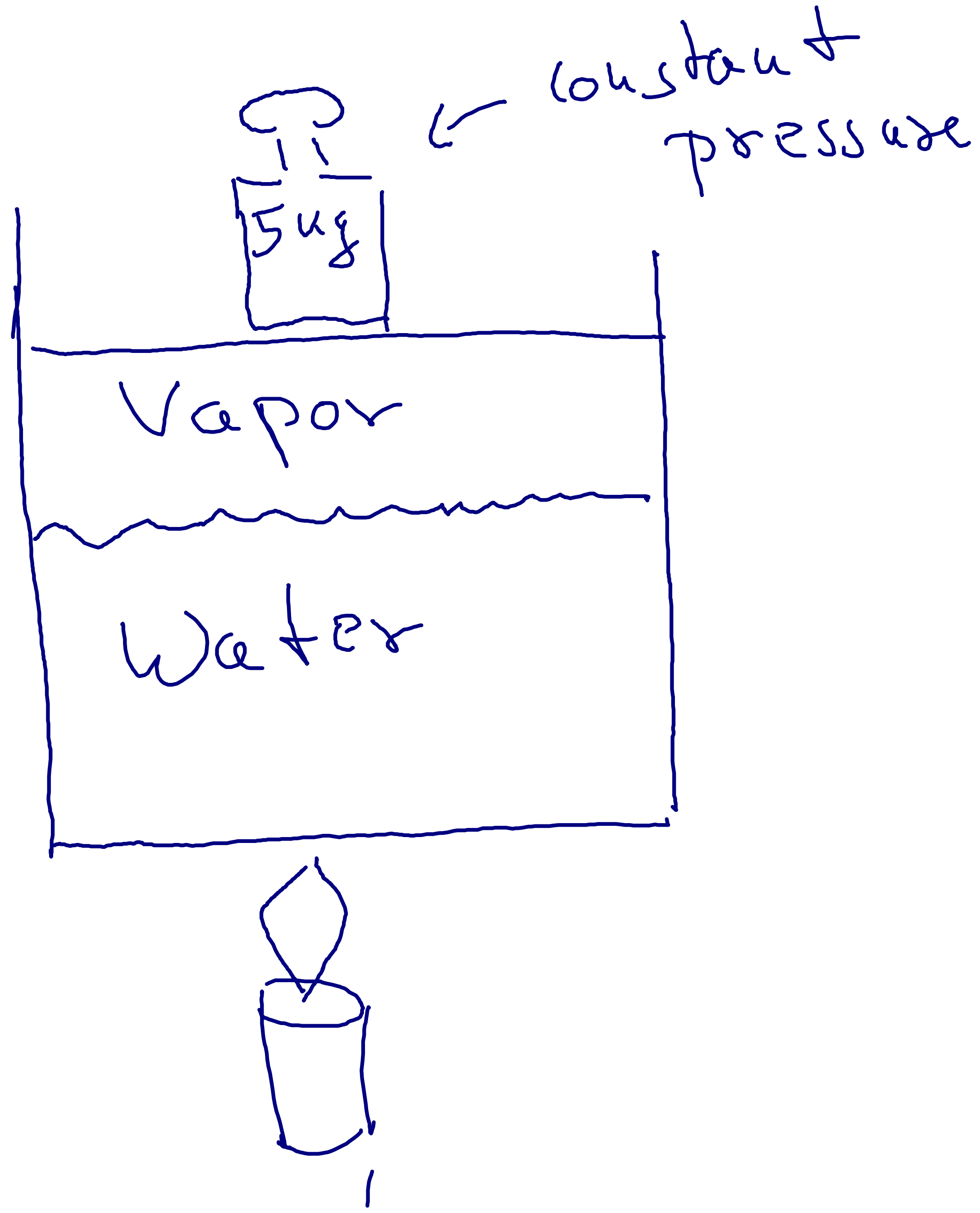
Examples:
Water - vapor (liquid - gas)
Water - ice
Ferromagnet
….

Order parameter: Tells in which phase the system is
Examples ?

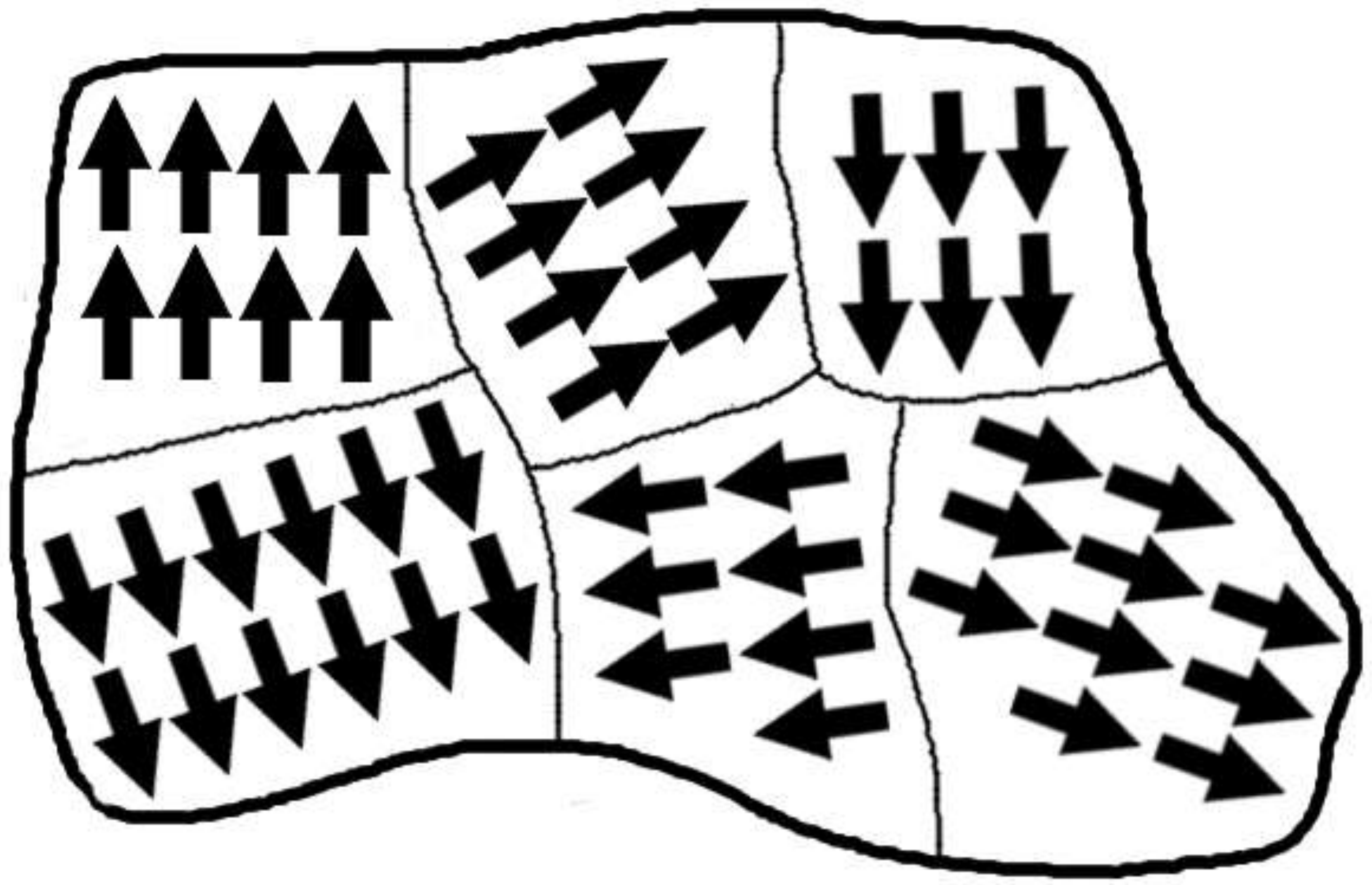
Control parameter: Moves system from one phase to another
Examples ?

Phase co-existence: Two or more phases can exist together
Examples ?

Phase Co-Existence

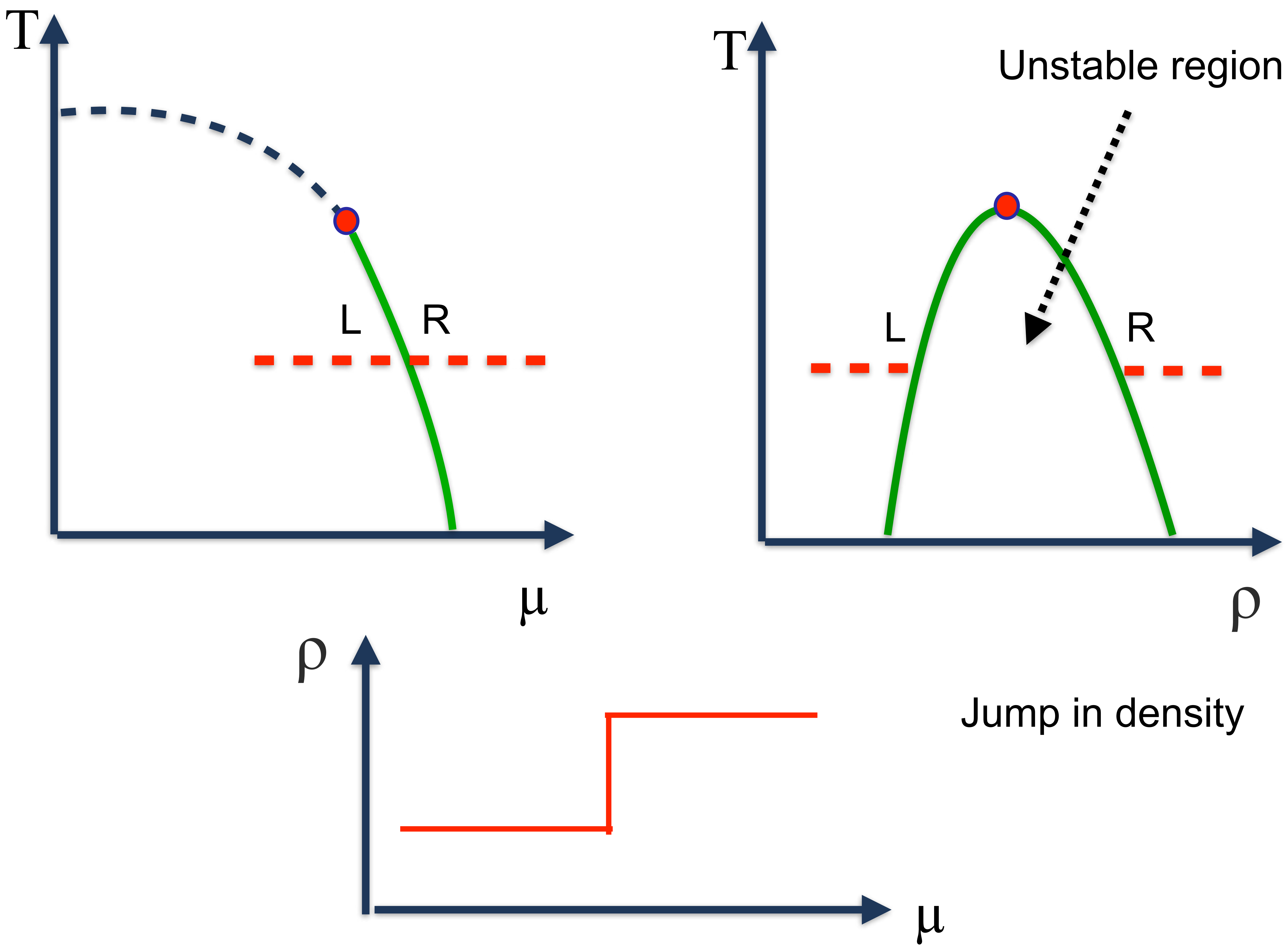


Water-vapor co-existence
a.k.a your water kettle

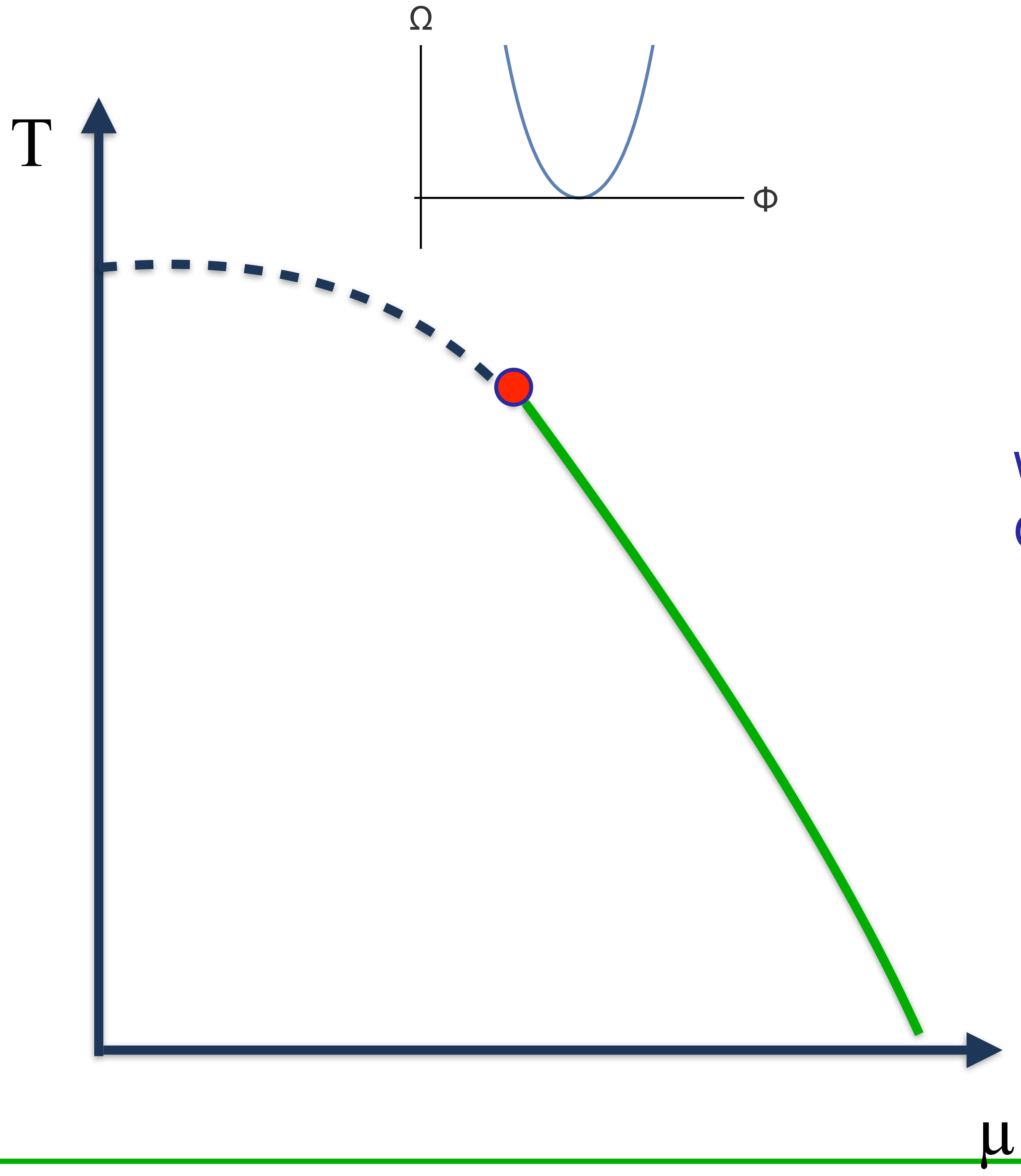


Ferro-magnet
Weiss domains

Phase diagrams



Free Energy



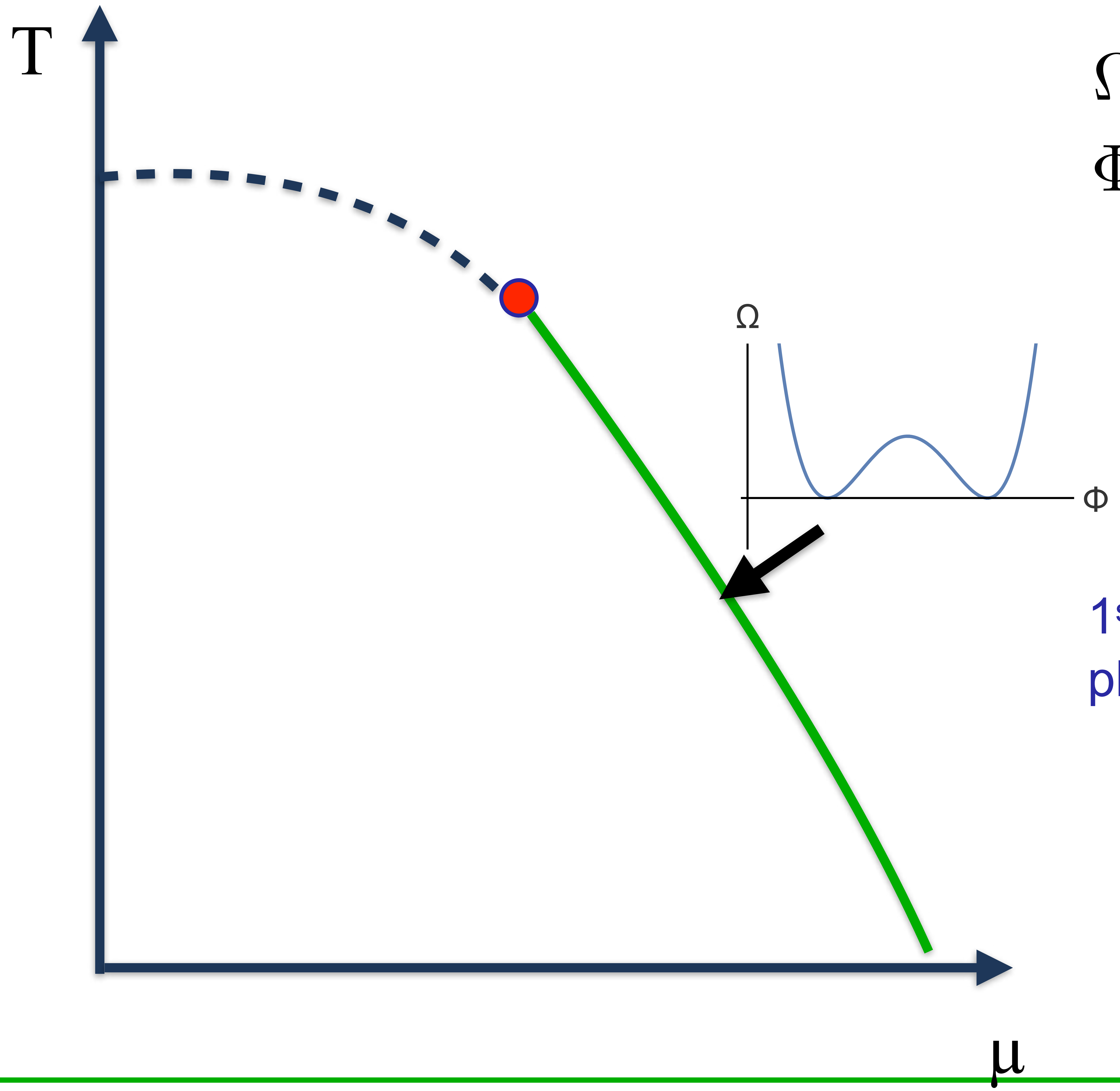
Free Energy:

$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

What we are used to:
One minimum

Free Energy



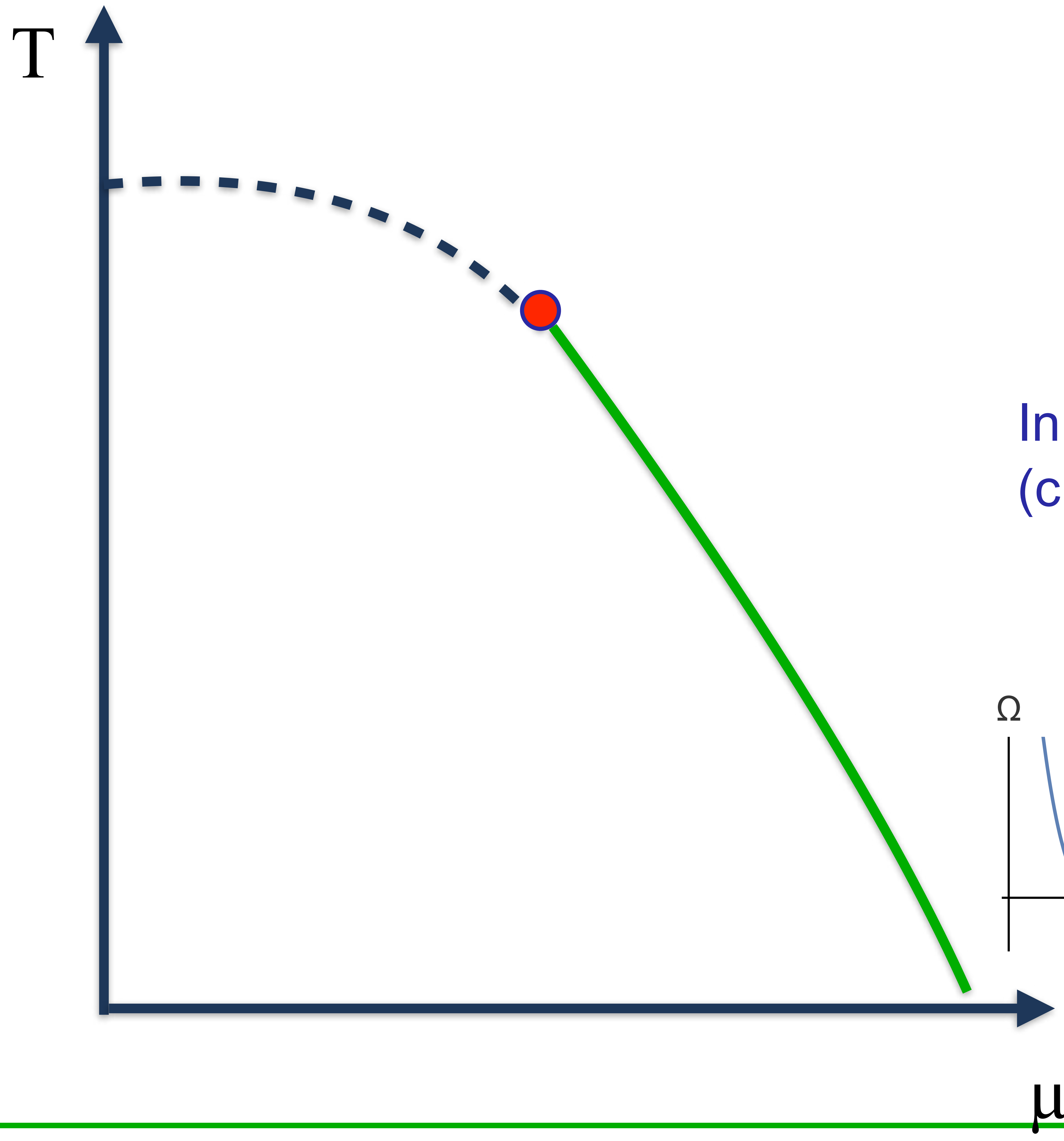
Free Energy:

$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

1st order
phase co-existence

Free Energy

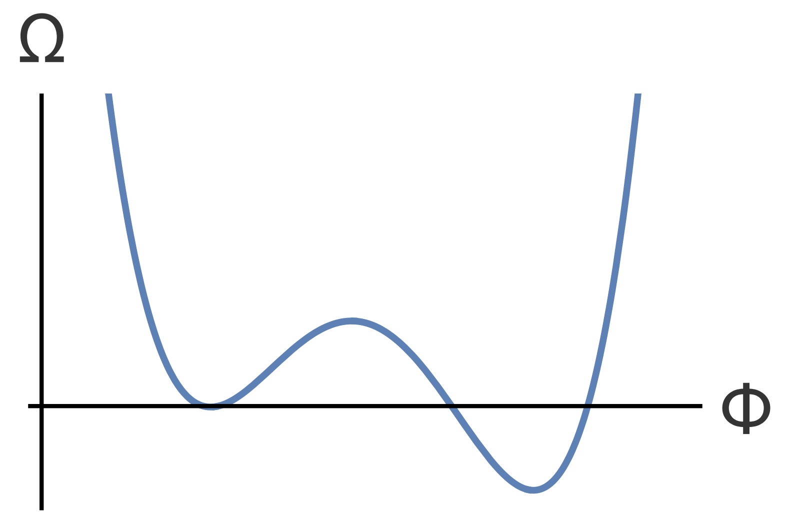


Free Energy:

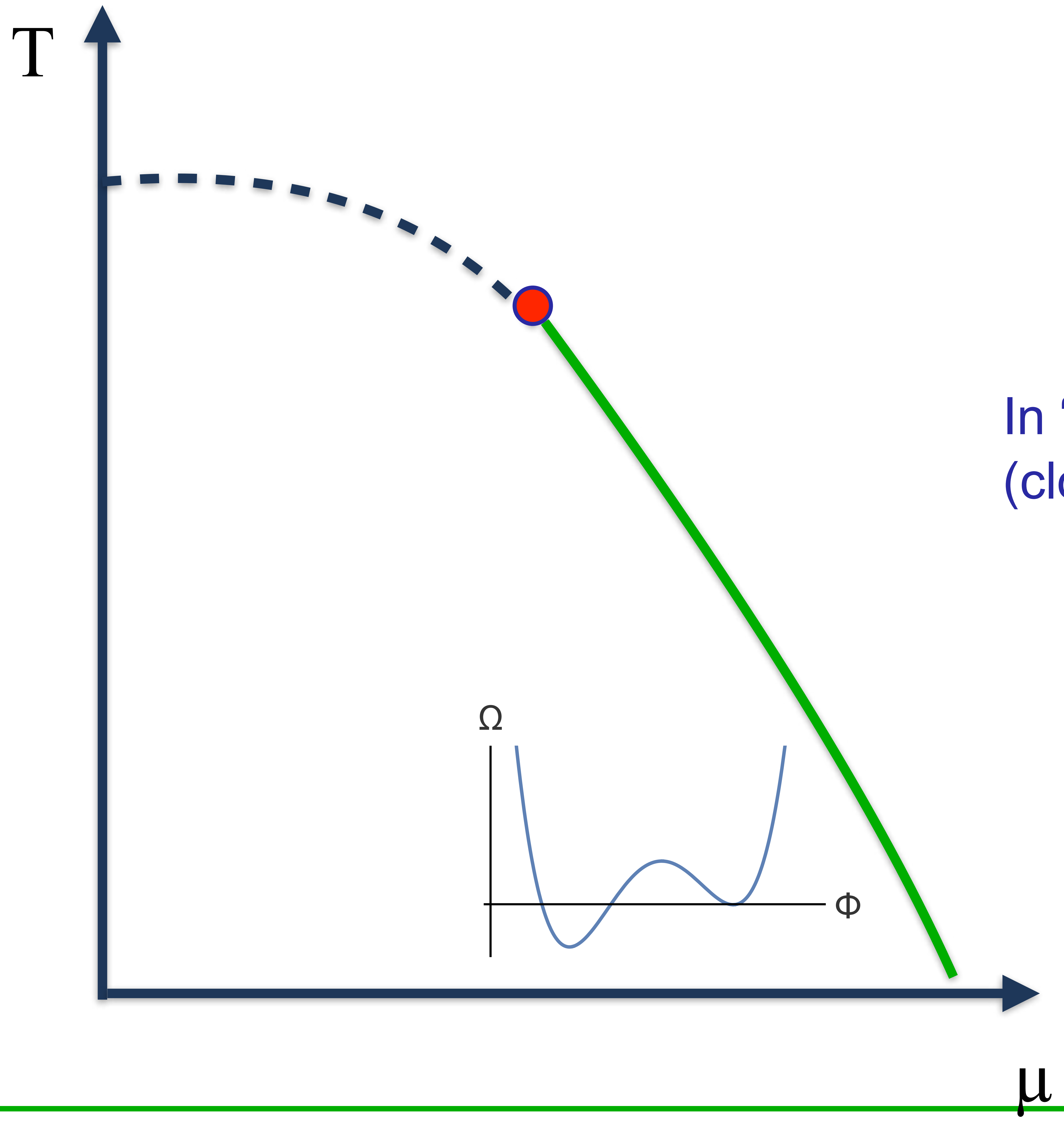
$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

In “dense” phase
(close to transition)



Free Energy



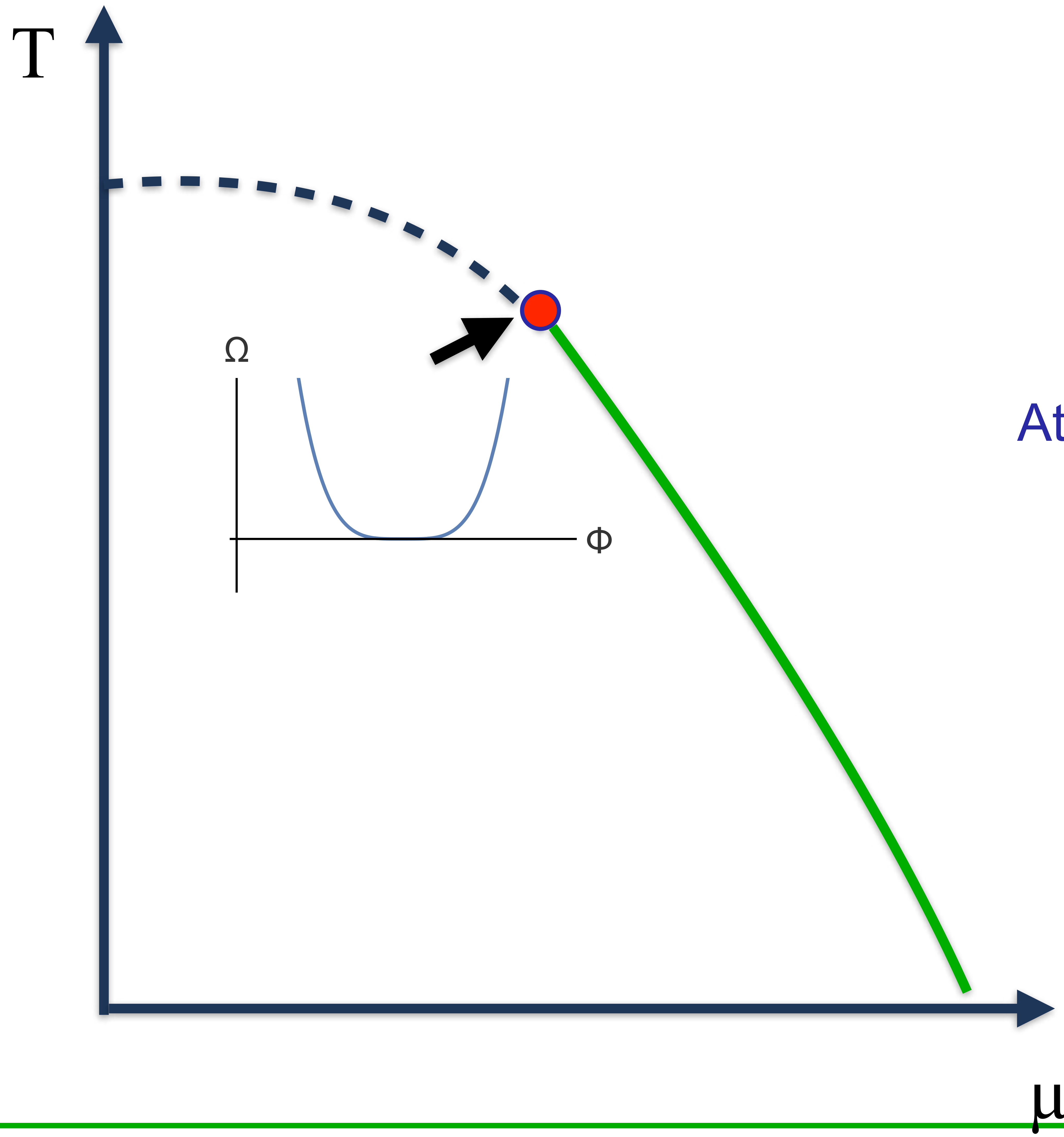
Free Energy:

$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

In “dilute” phase
(close to transition)

Free Energy



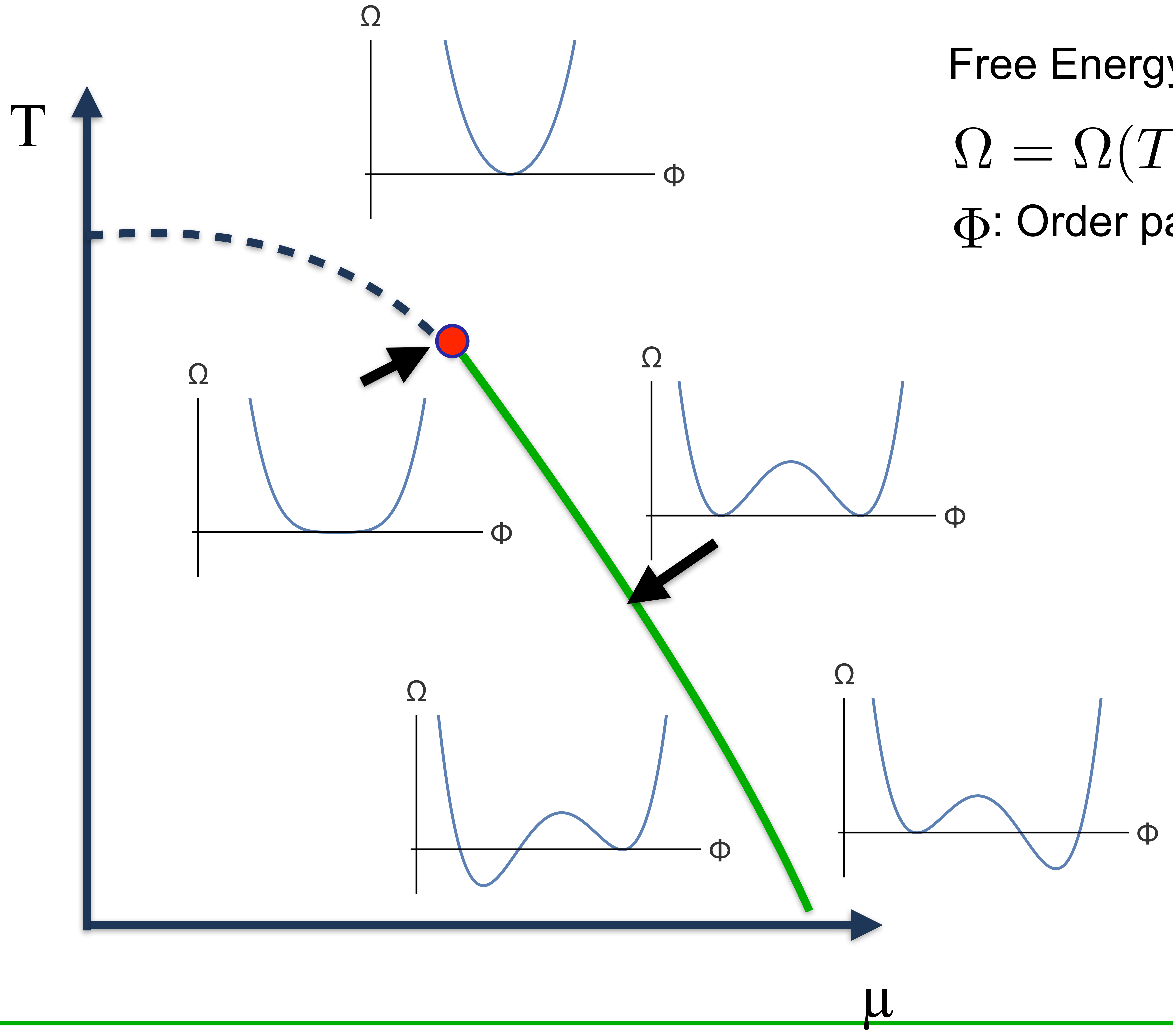
Free Energy:

$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

At the critical point

Free Energy

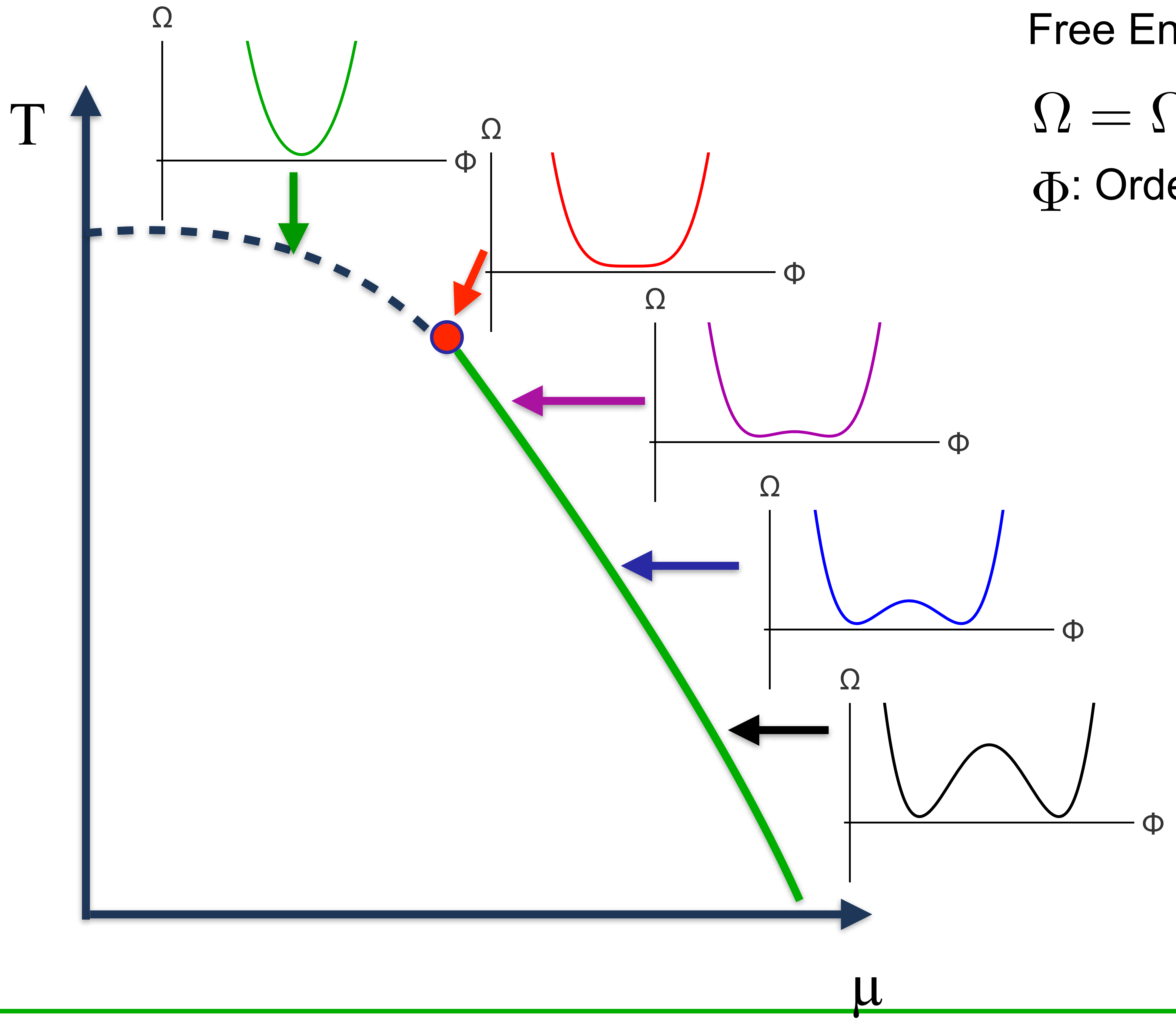


Free Energy:

$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

Free Energy

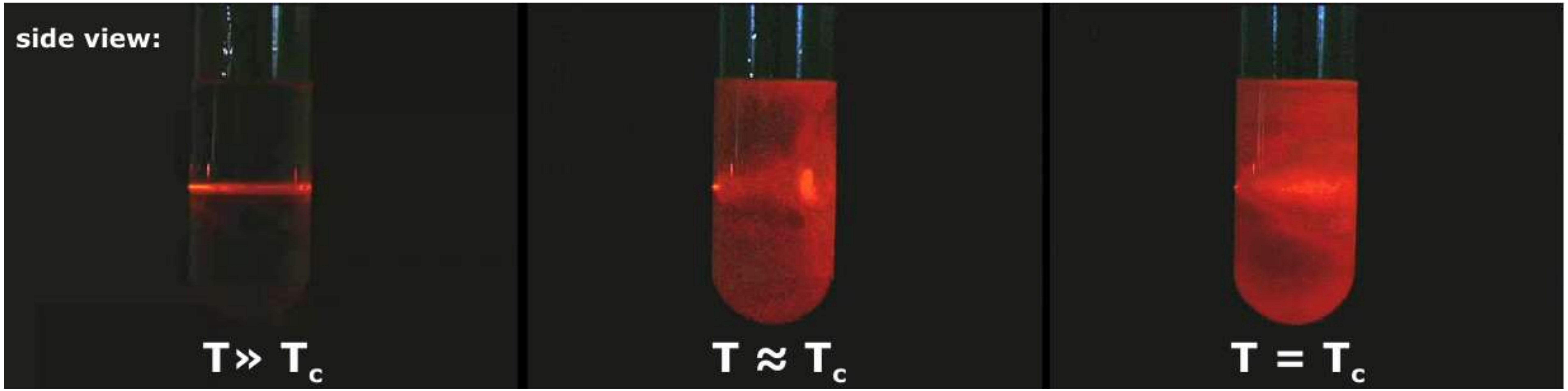


Free Energy:

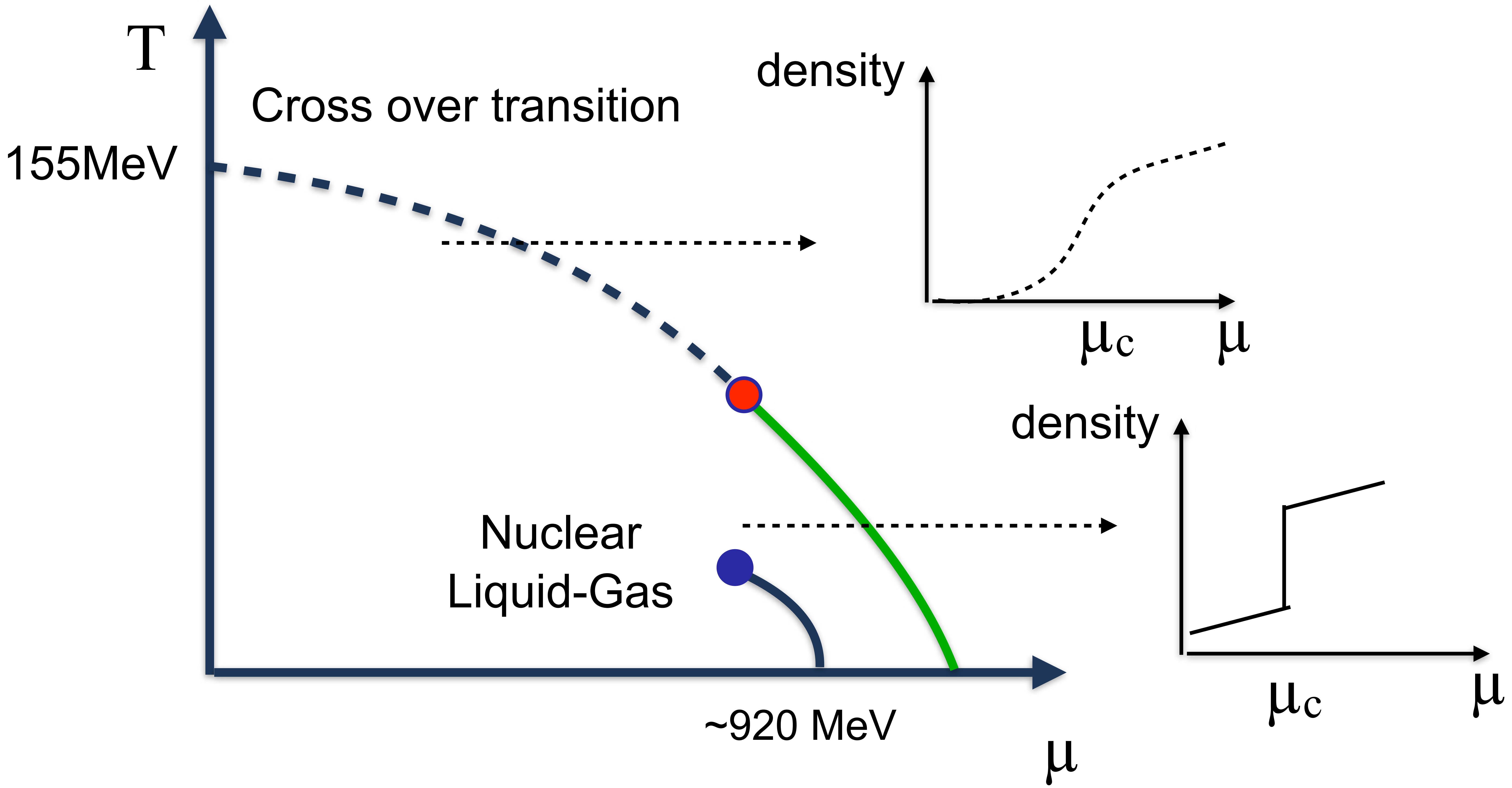
$$\Omega = \Omega(T, \mu;\ \Phi)$$

Φ: Order parameter

Fluctuations at all length scales
critical opalescence

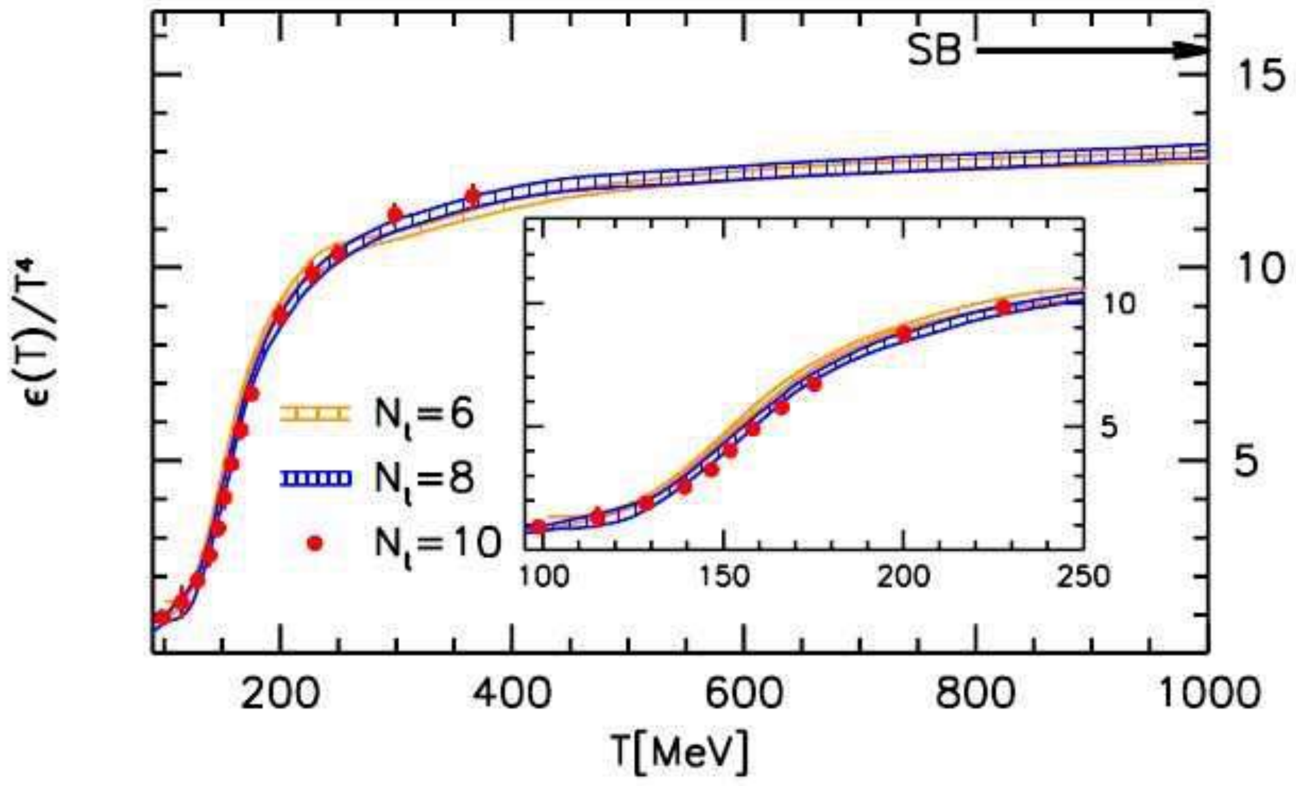


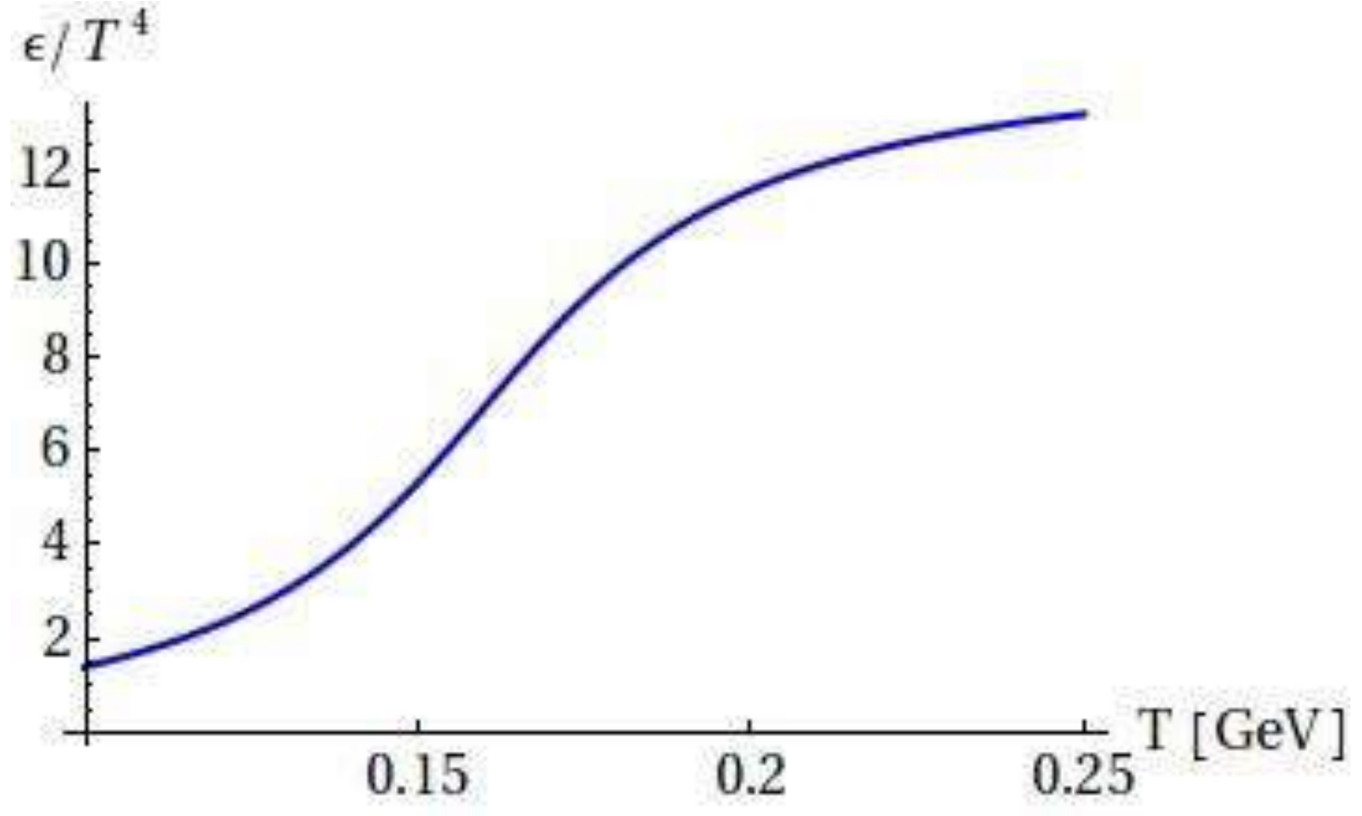
Looking for signs of a transition



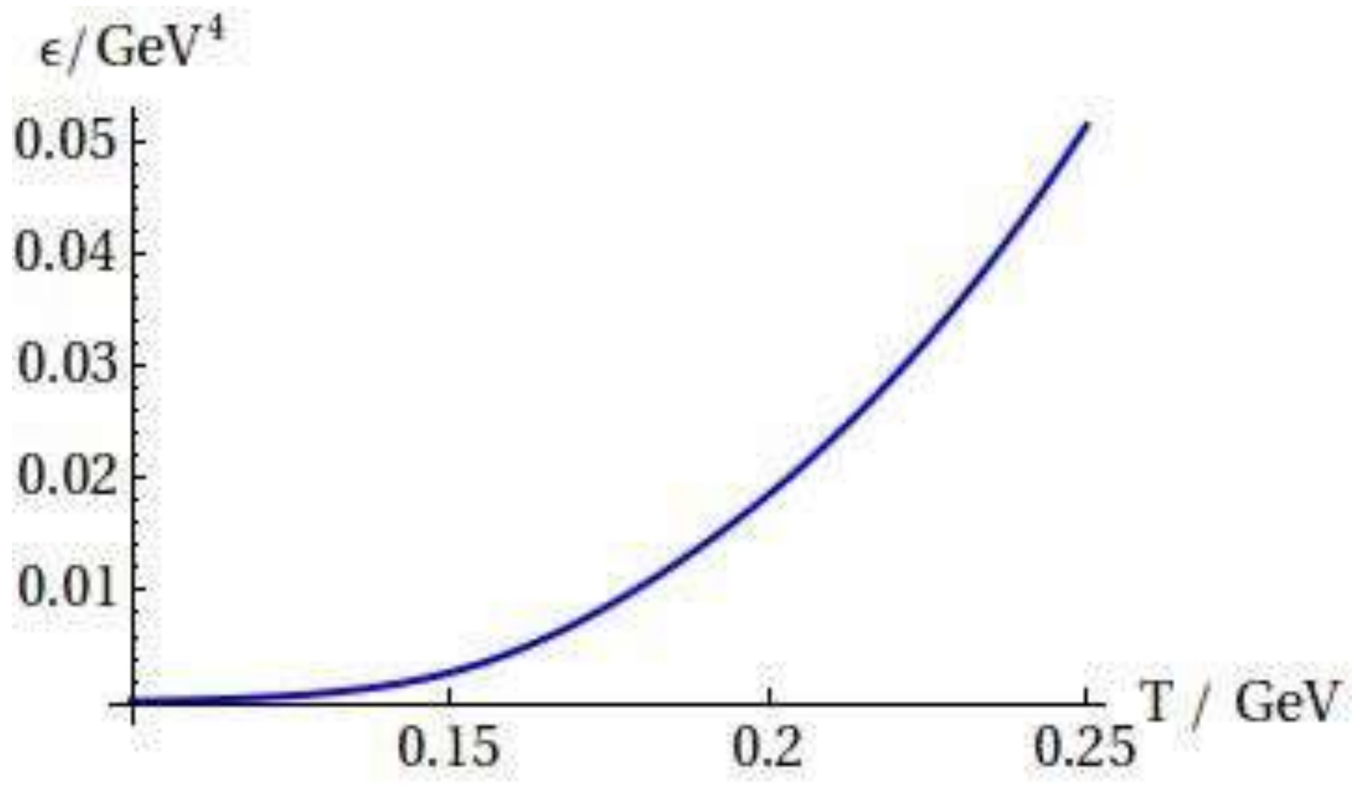
Cumulants and Phase structure

S. Borsanyi et al, JHEP 1011 (2010) 077





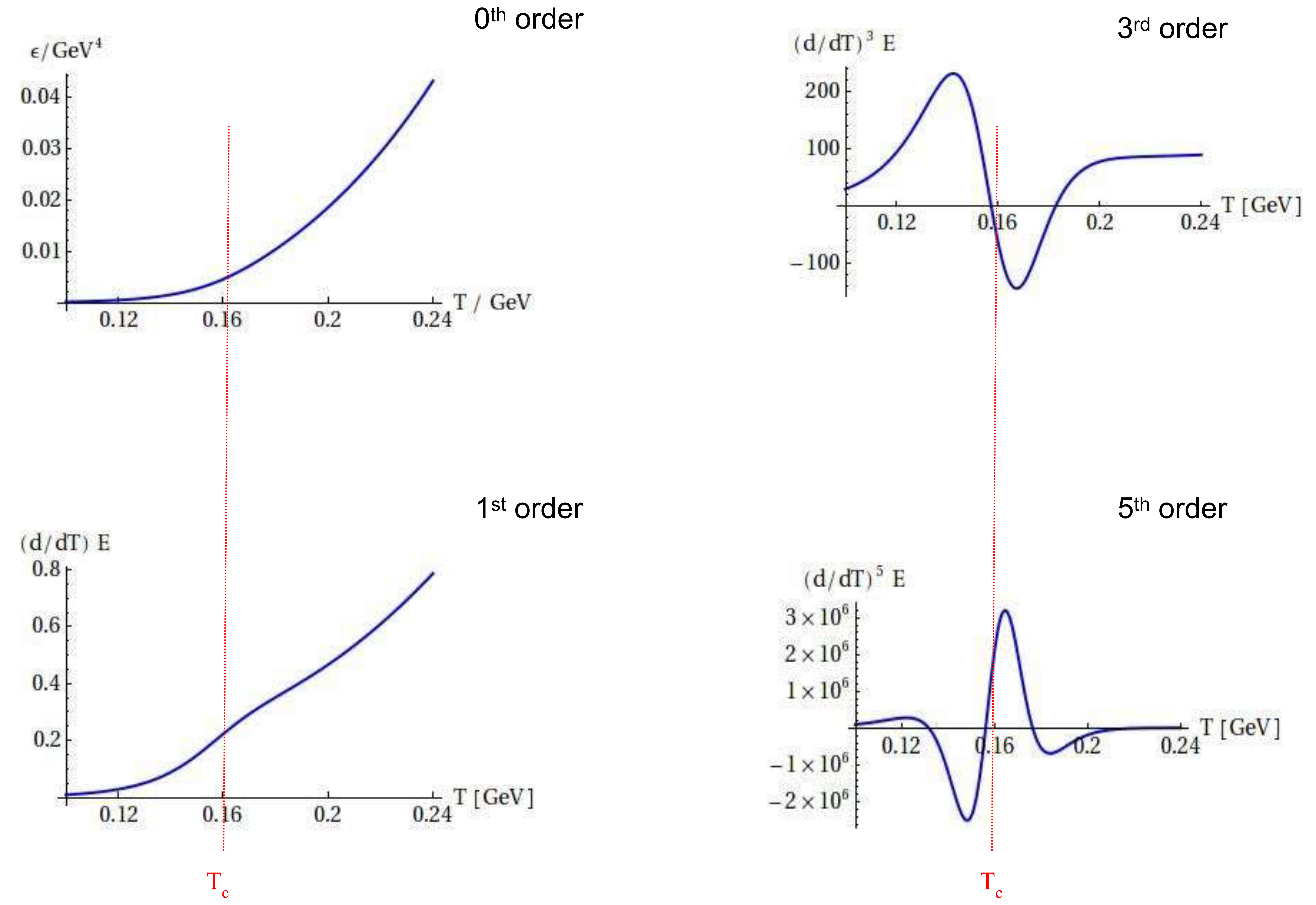
What we always see....



What it really means....

“T_c” ~ 155 MeV

Derivatives



How to measure derivatives

$$Z = tr\, e^{-\hat{E}/T + \mu/T \hat{N}_B}$$

$$\langle E \rangle = \frac{1}{Z} tr\, \hat{E}\, e^{-\hat{E}/T + \mu/T \hat{N}_B} = -\frac{\partial}{\partial 1/T} \ln(Z)$$

$$\langle (\delta E)^2 \rangle = \langle E^2 \rangle - \langle E \rangle^2 = \left(-\frac{\partial}{\partial 1/T} \right)^2 \ln(Z) = \left(-\frac{\partial}{\partial 1/T} \right) \langle E \rangle$$

$$\langle (\delta E)^n \rangle = \left(-\frac{\partial}{\partial 1/T} \right)^{n-1} \langle E \rangle$$

Cumulants of Energy measure the temperature derivatives of the EOS

Cumulants of Baryon number measure the chem. pot. derivatives of the EOS

Cumulants of (Baryon) Number

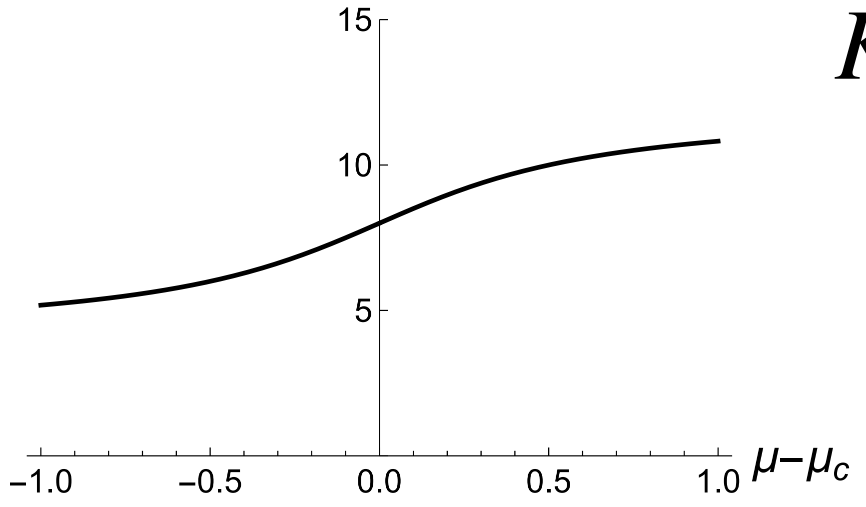
$$K_n = \frac{\partial^n}{\partial(\mu/T)^n} \ln Z = \frac{\partial^{n-1}}{\partial(\mu/T)^{n-1}} \langle N \rangle$$

$$K_1 = \langle N \rangle \,, \;\; K_2 = \langle N - \langle N \rangle \rangle^2 \,, \;\; K_3 = \langle N - \langle N \rangle \rangle^3$$

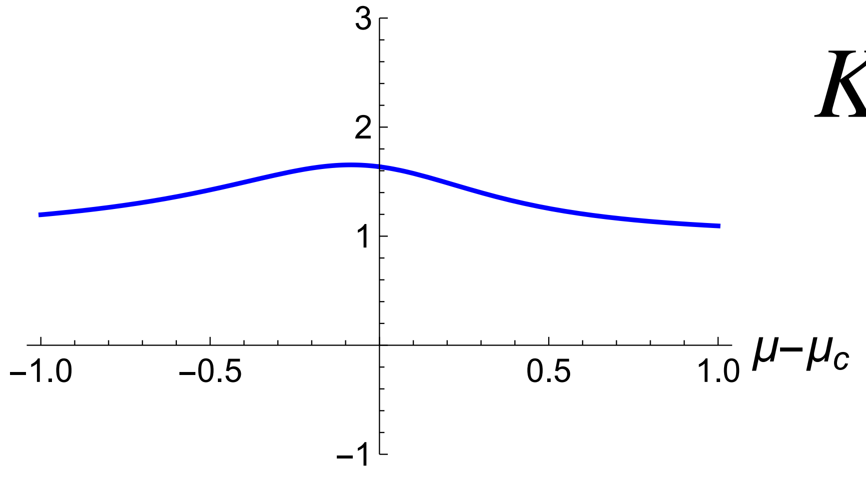
Cumulants scale with volume (extensive): $K_n \sim V$

Volume not well controlled in heavy ion collisions

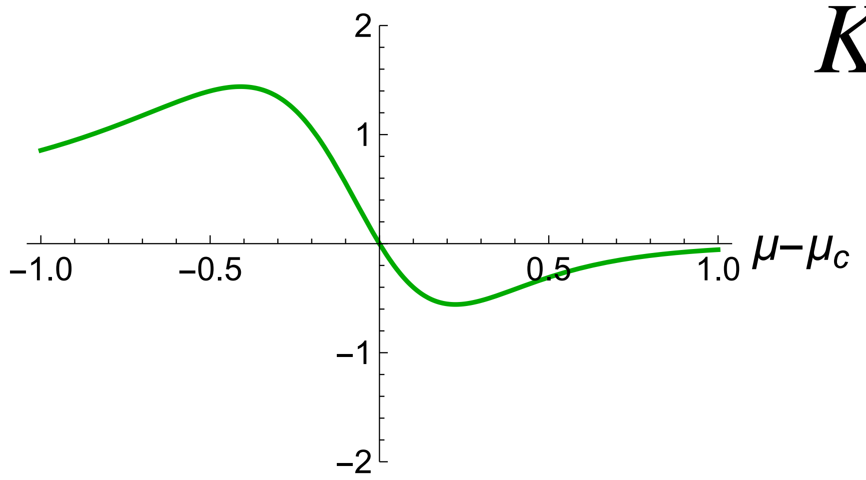
Cumulant Ratios: $\frac{K_2}{\langle N \rangle}, \;\; \frac{K_3}{K_2}, \;\; \frac{K_4}{K_2}$



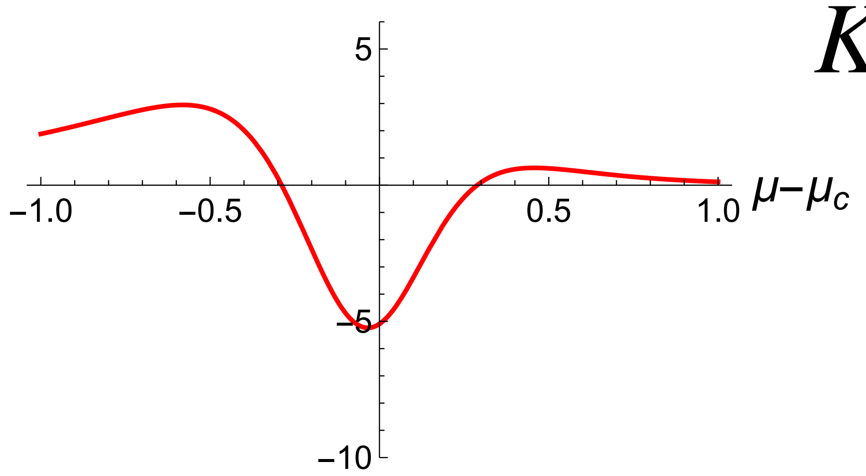
$K_1 = \langle N \rangle$



$K_2 = \frac{\partial}{\partial(\mu/T)} K_1$

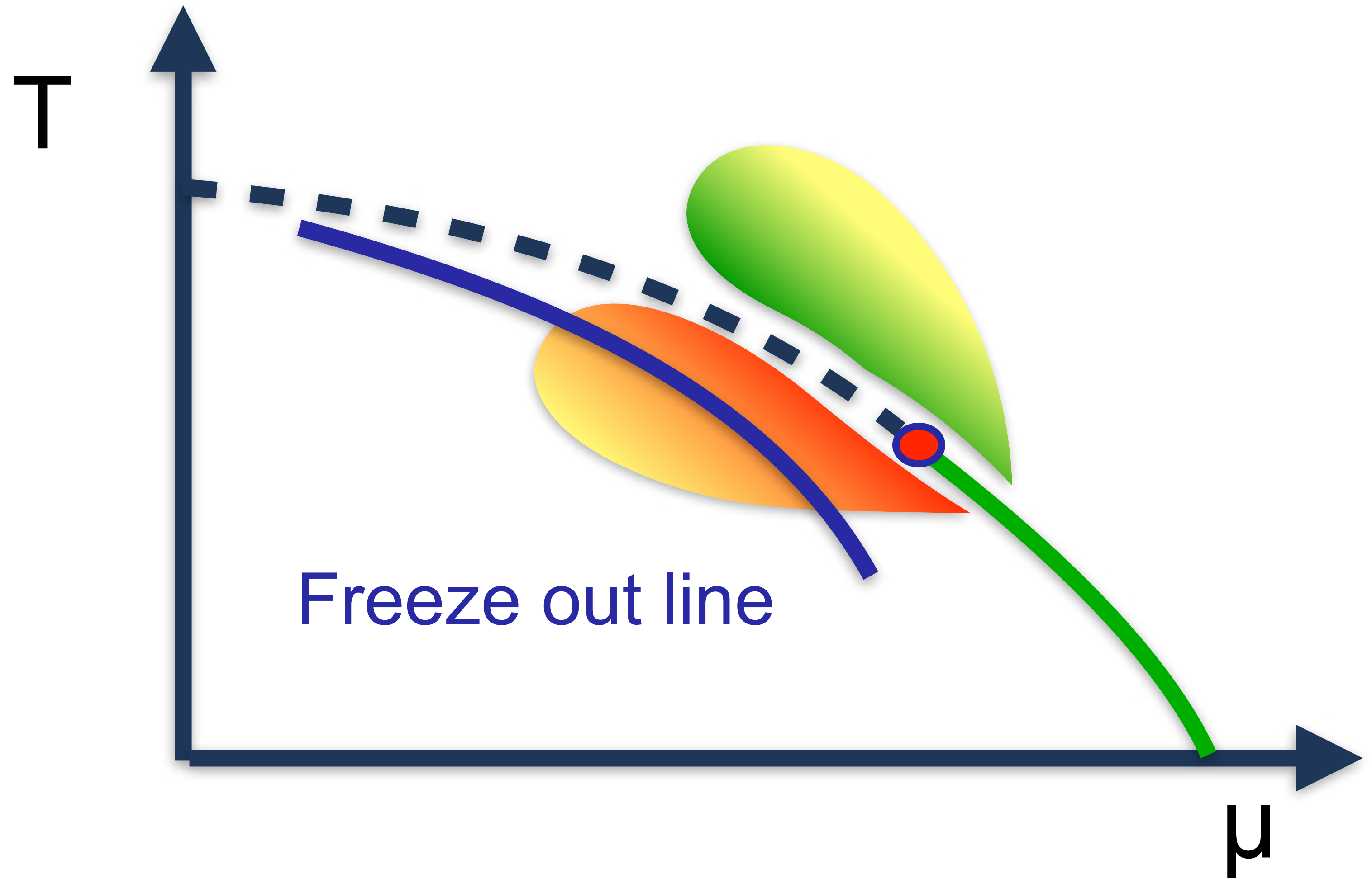


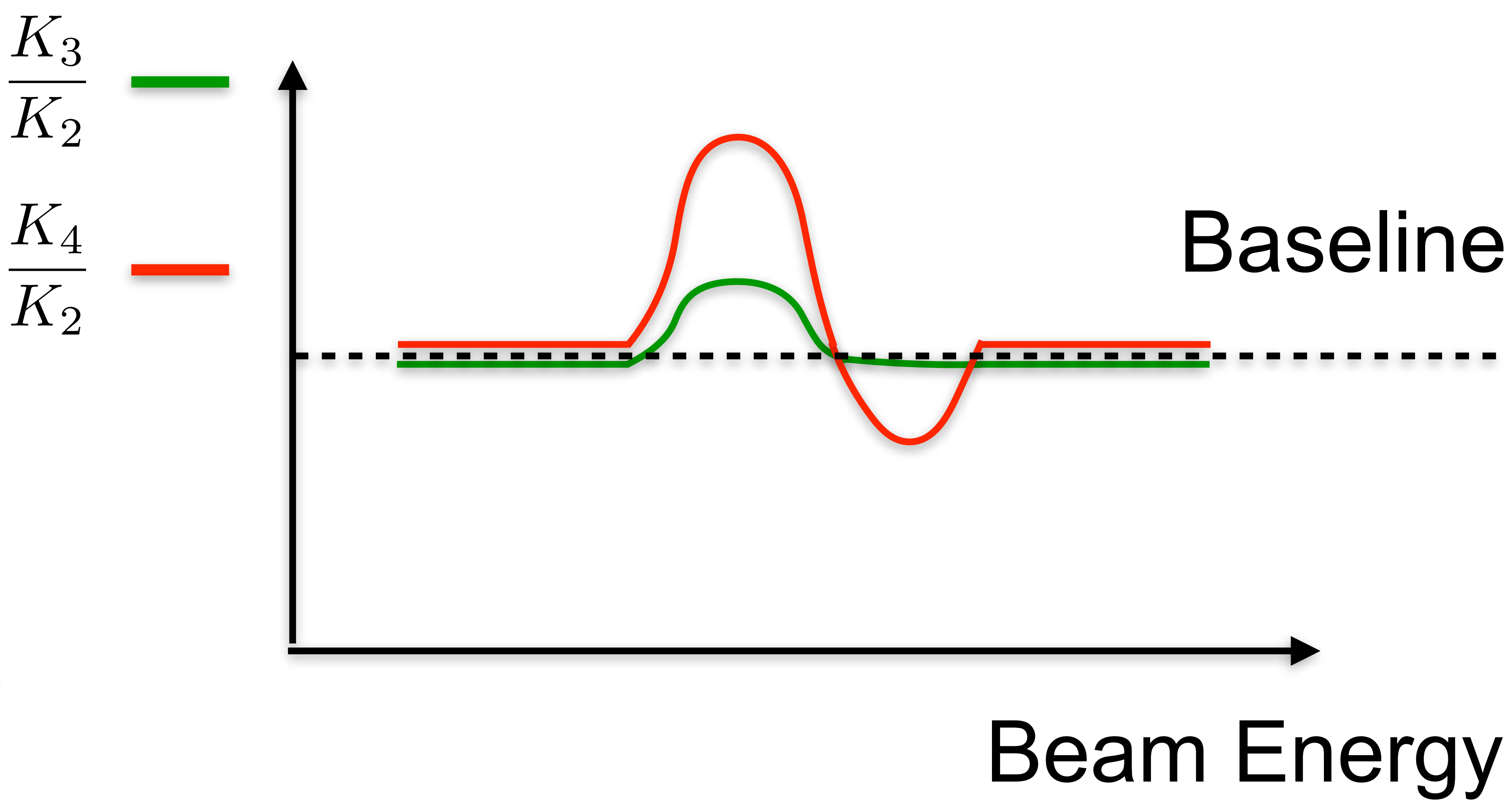
$K_3 = \frac{\partial}{\partial(\mu/T)} K_2$

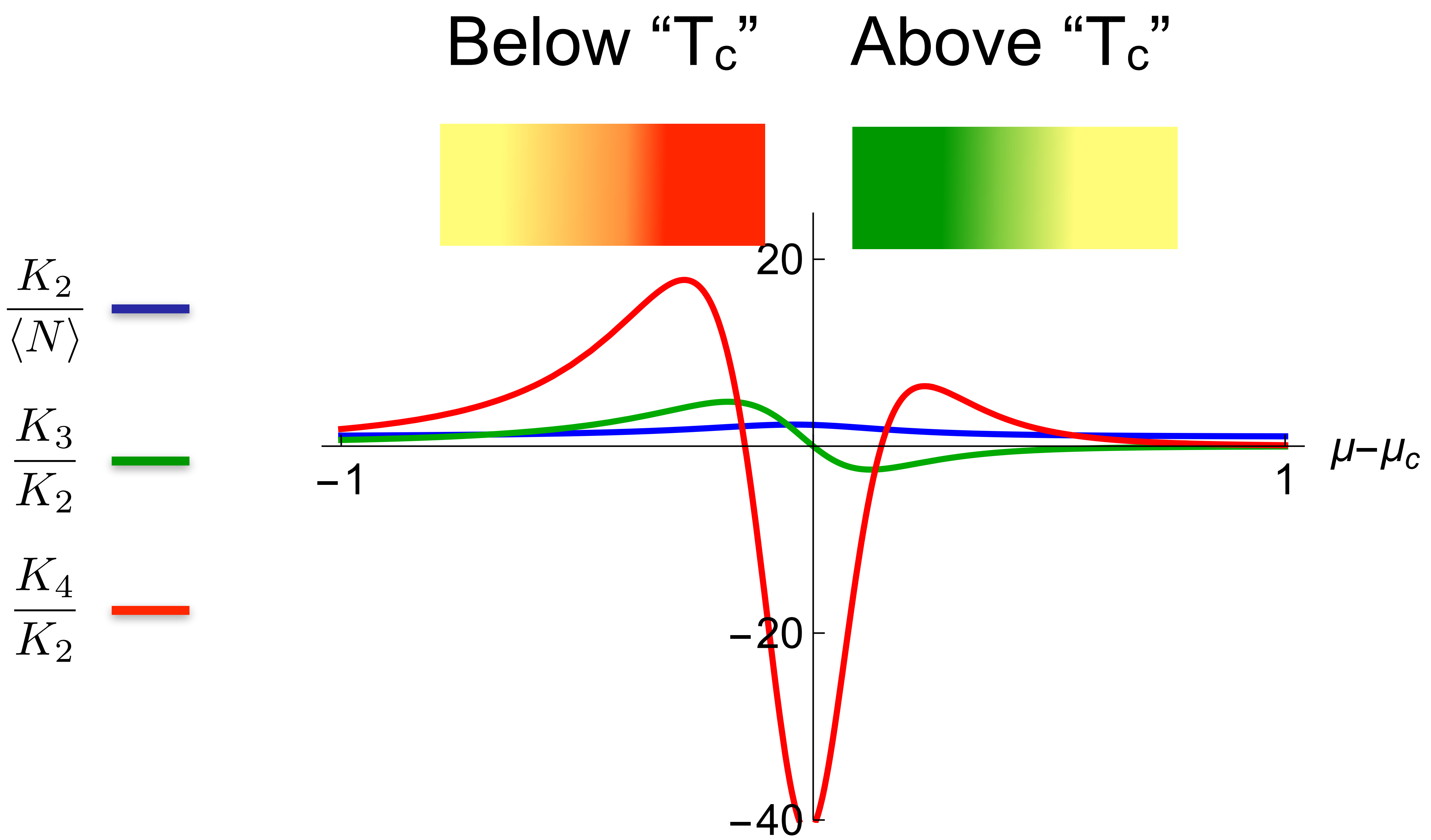


$K_4 = \frac{\partial}{\partial(\mu/T)} K_3$

What to expect?



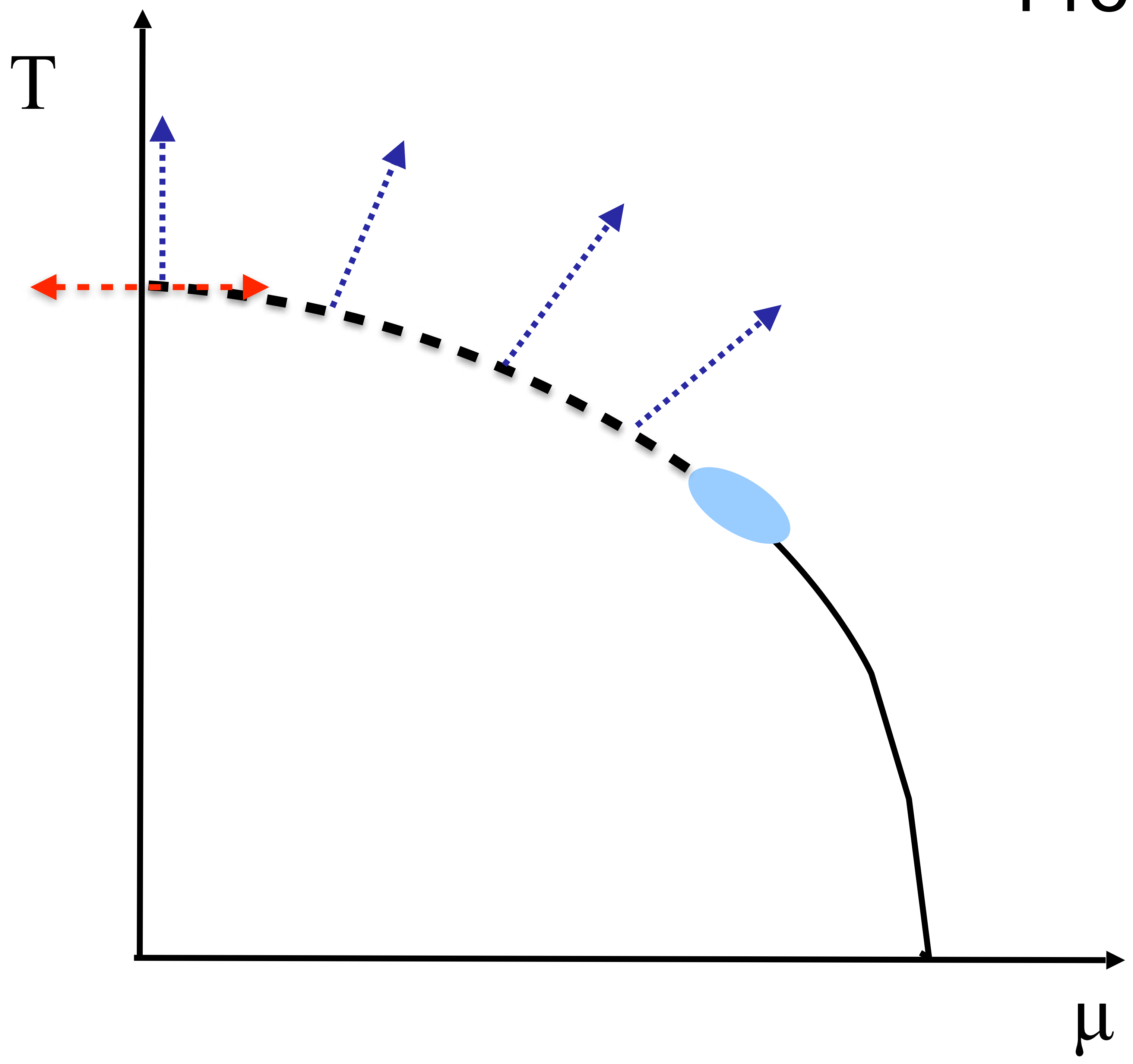




Close to μ=0

Free energy: $F = F(r), \quad r = \sqrt{T^2 + a\mu^2}$

a ~ curvature of critical line

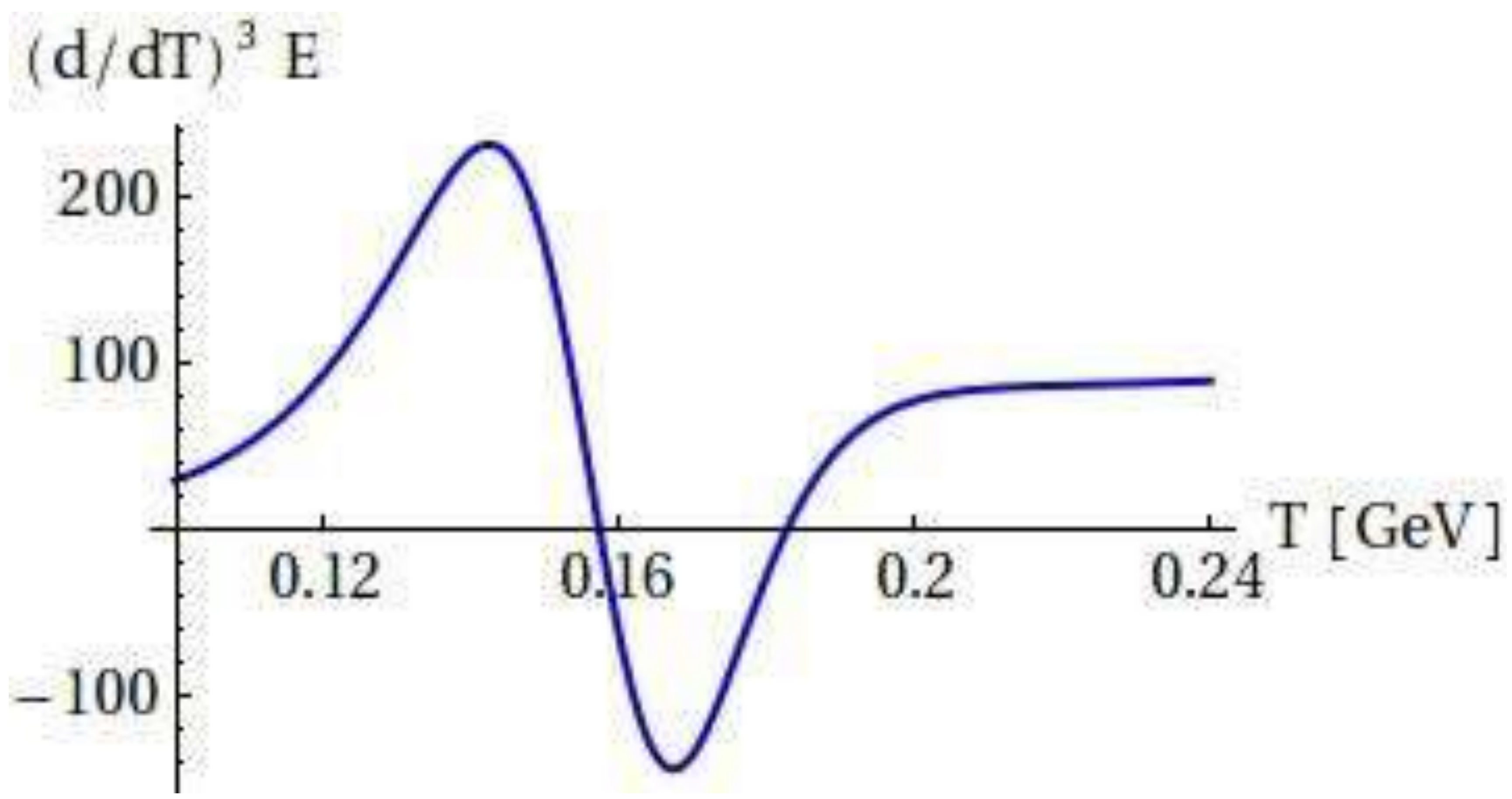


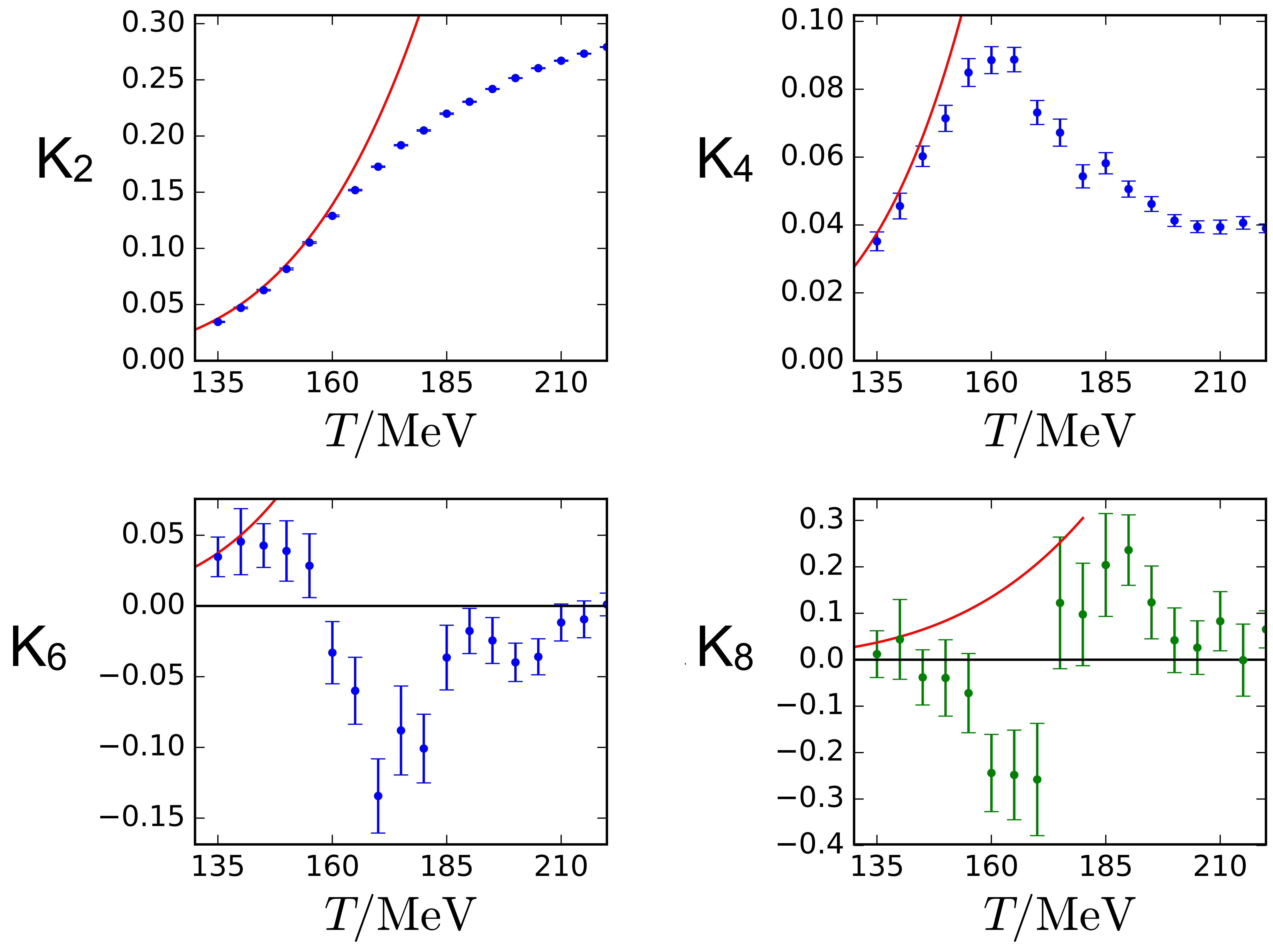
$$\frac{\partial^2}{\partial \mu^2} F(T,\mu)|_{\mu=0} = \frac{a}{T}\frac{\partial}{\partial T} F(T,\mu=0) \sim \langle E \rangle$$

Needs higher order cumulants (derivatives) at μ ~ 0

Cumulants at small μ

- Baryon number cumulants can be calculated in Lattice QCD
  - possible test of chiral criticality ? Friman et al, '11





Wuppertal/Budapest 1805.04445

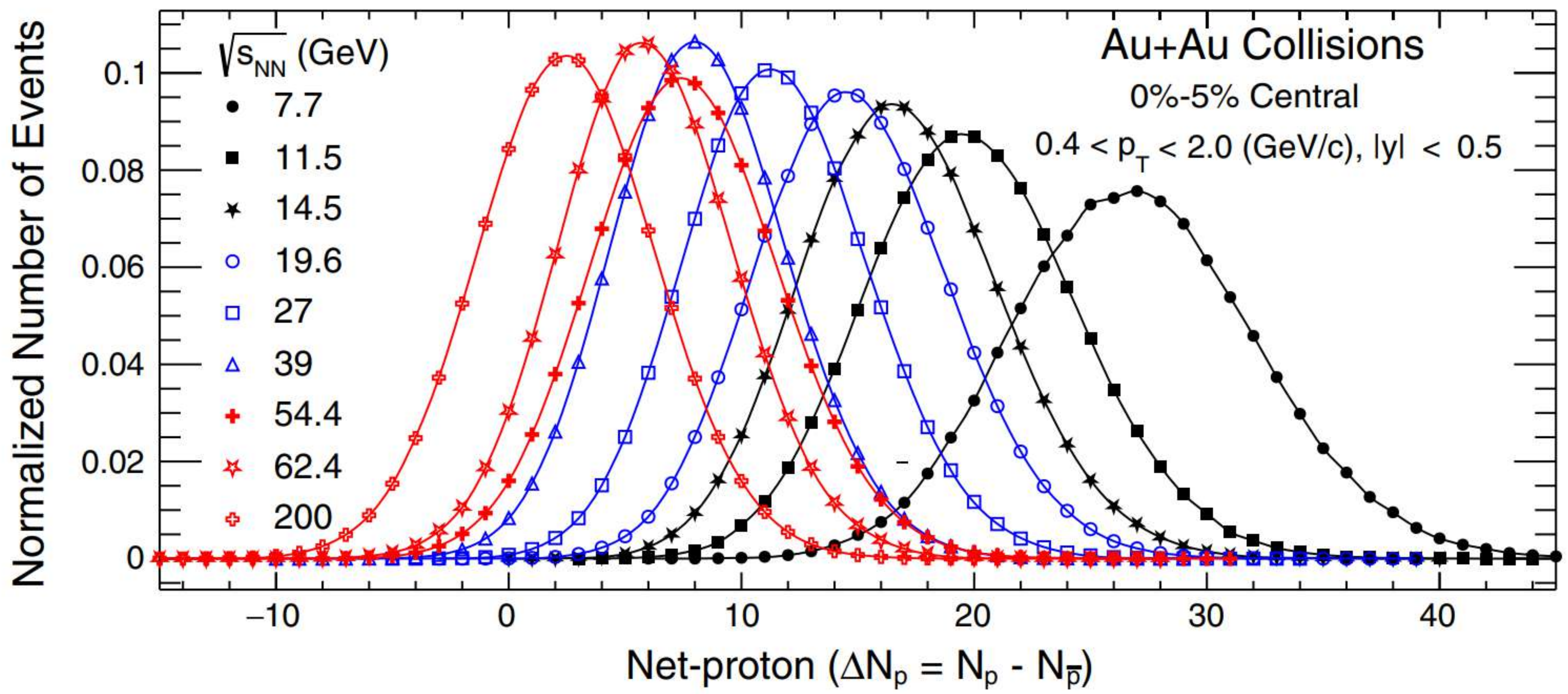
Measuring cumulants (derivatives)

$$K_2 = \langle N - \langle N \rangle \rangle^2 = \sum_N P(N)(N - \langle N \rangle)^2$$

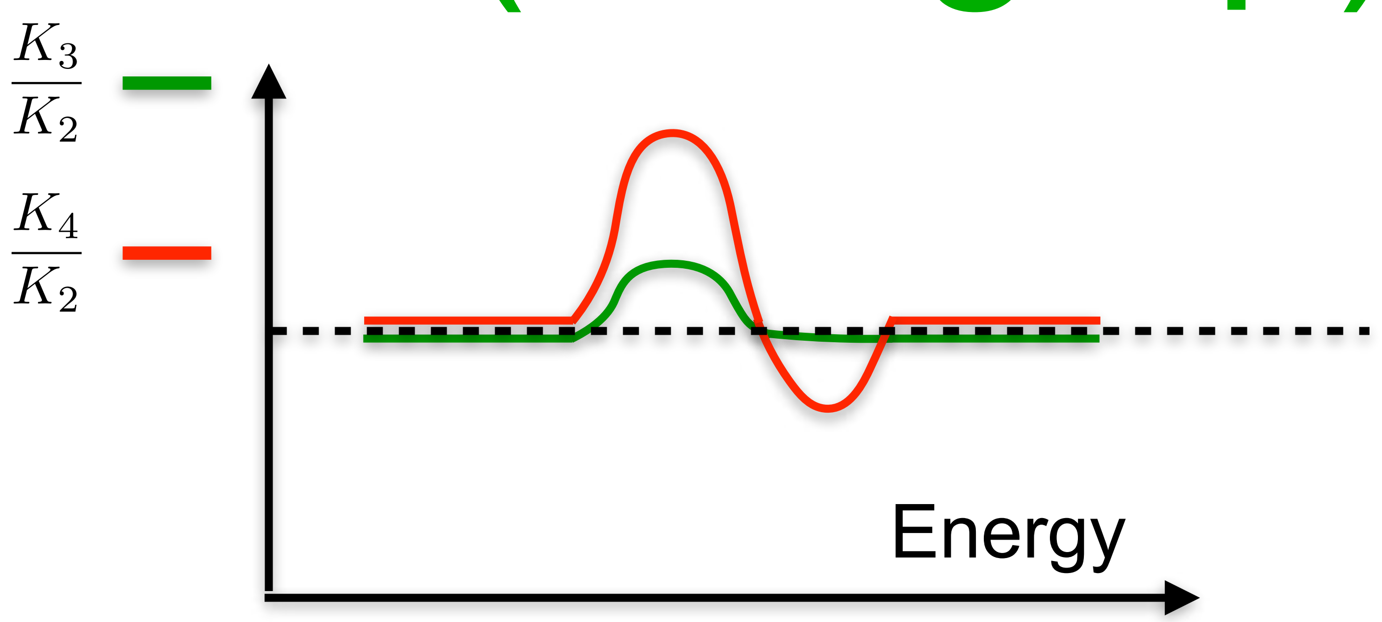
$$K_3 = \langle N - \langle N \rangle \rangle^3 = \sum_N P(N)(N - \langle N \rangle)^3$$

$$P(N) = \frac{N_{events}(N)}{N_{events}(total)}$$

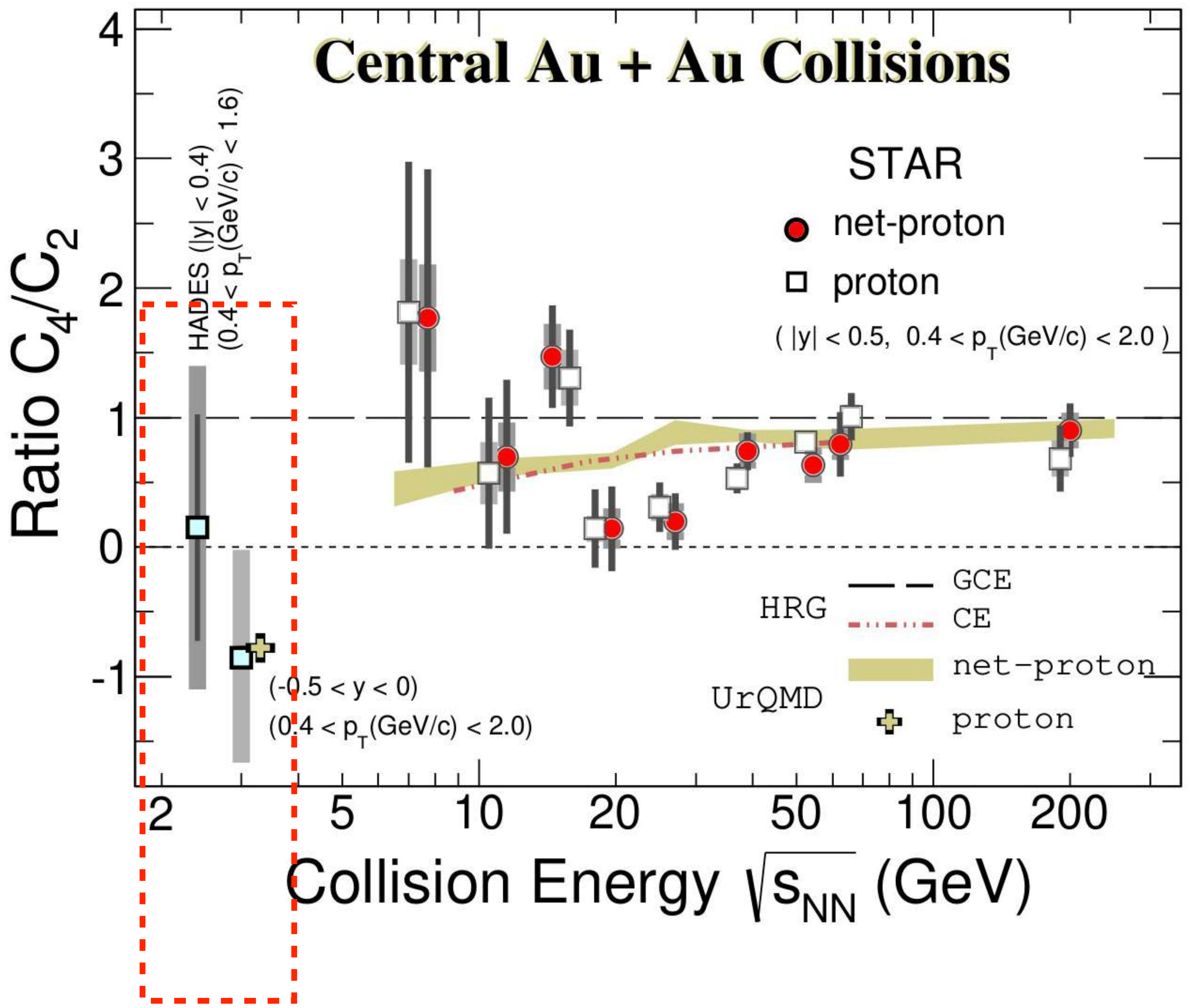
STAR Collaboration, Phys. Rev. Lett. 126, 092301 (2021)



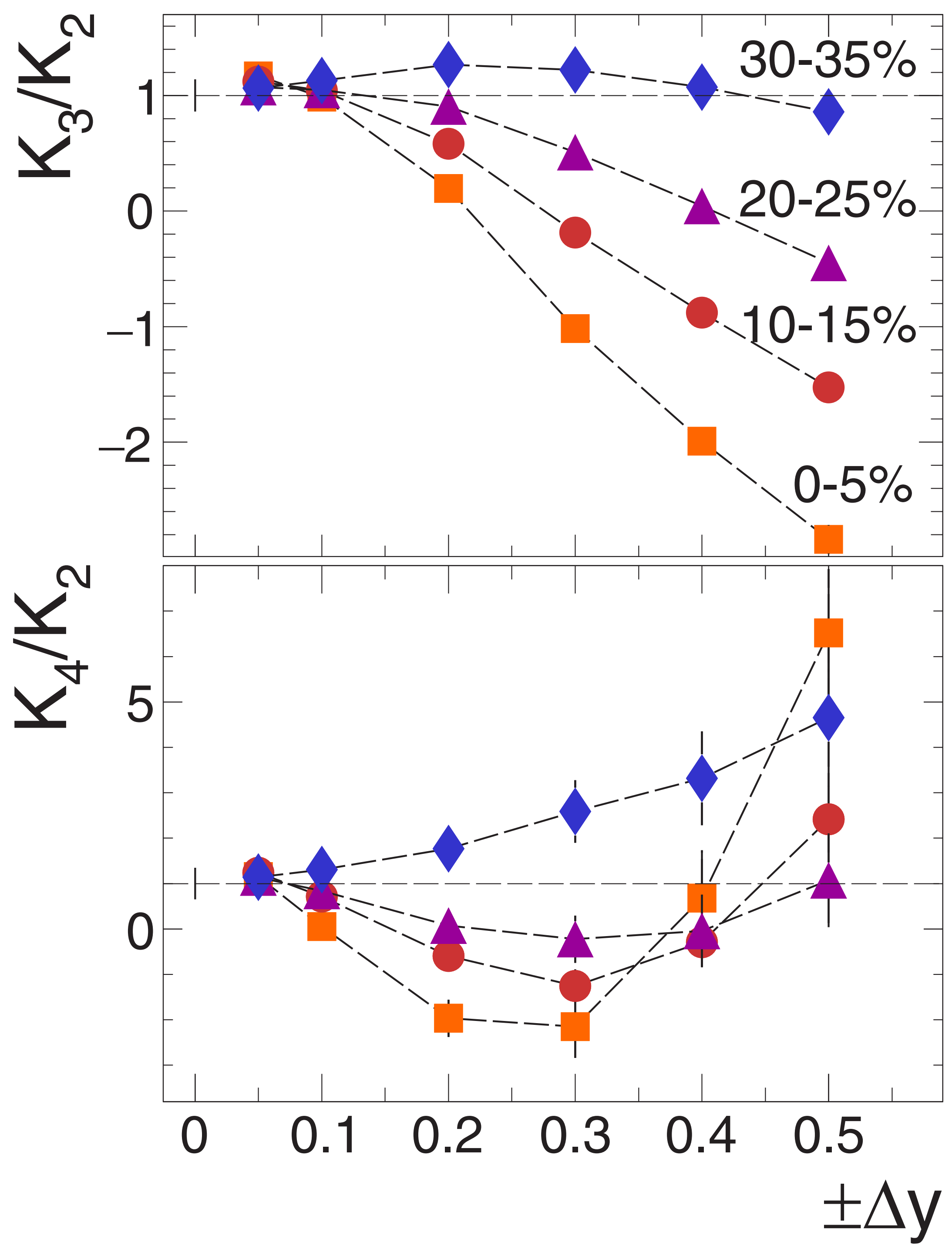
Data (at high μ)



STAR: arXiv:2112.00240



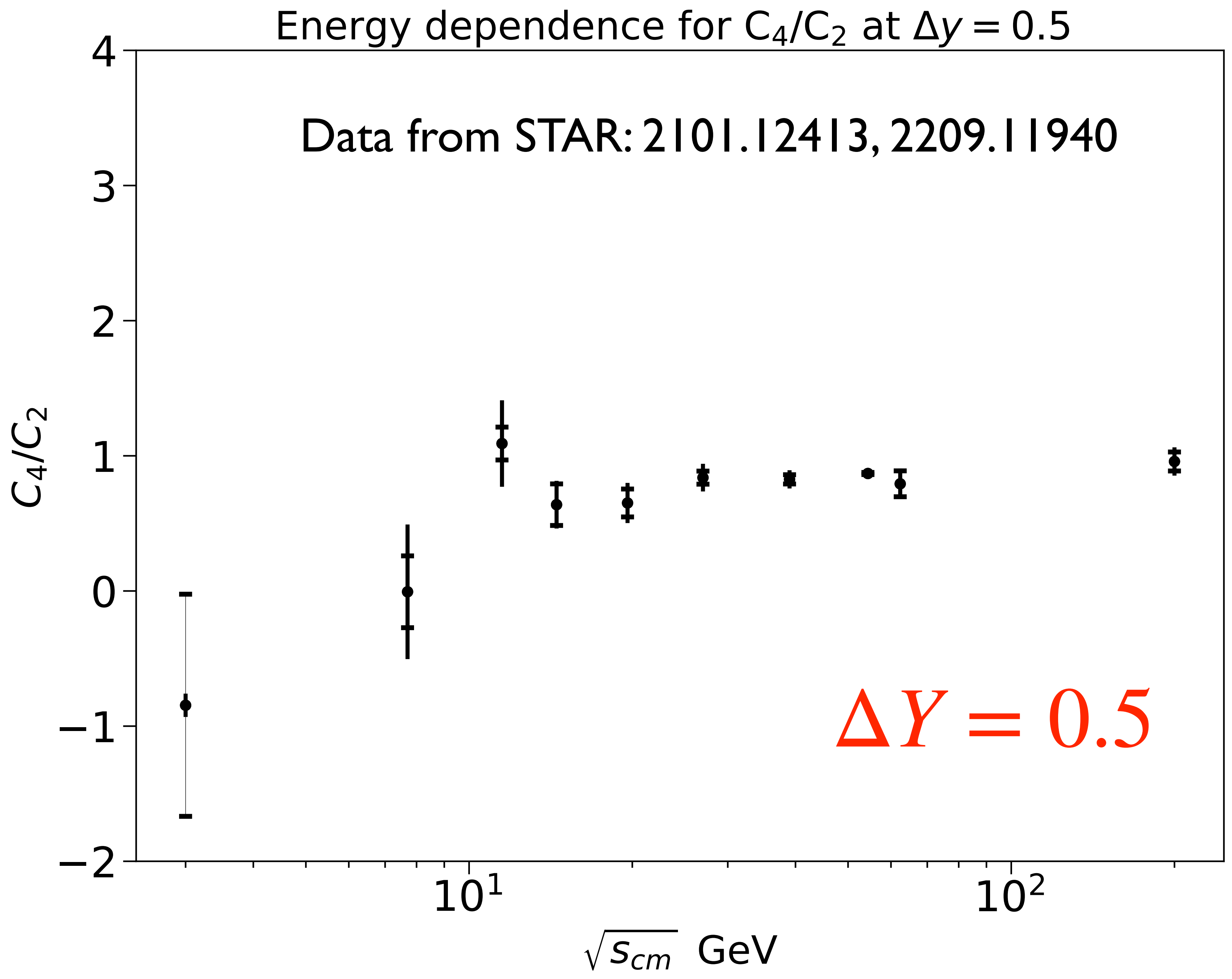
Different acceptance!



HADES: arXiv:2002.08701

$\sqrt{s} = 2.4$ GeV

Fair comparison



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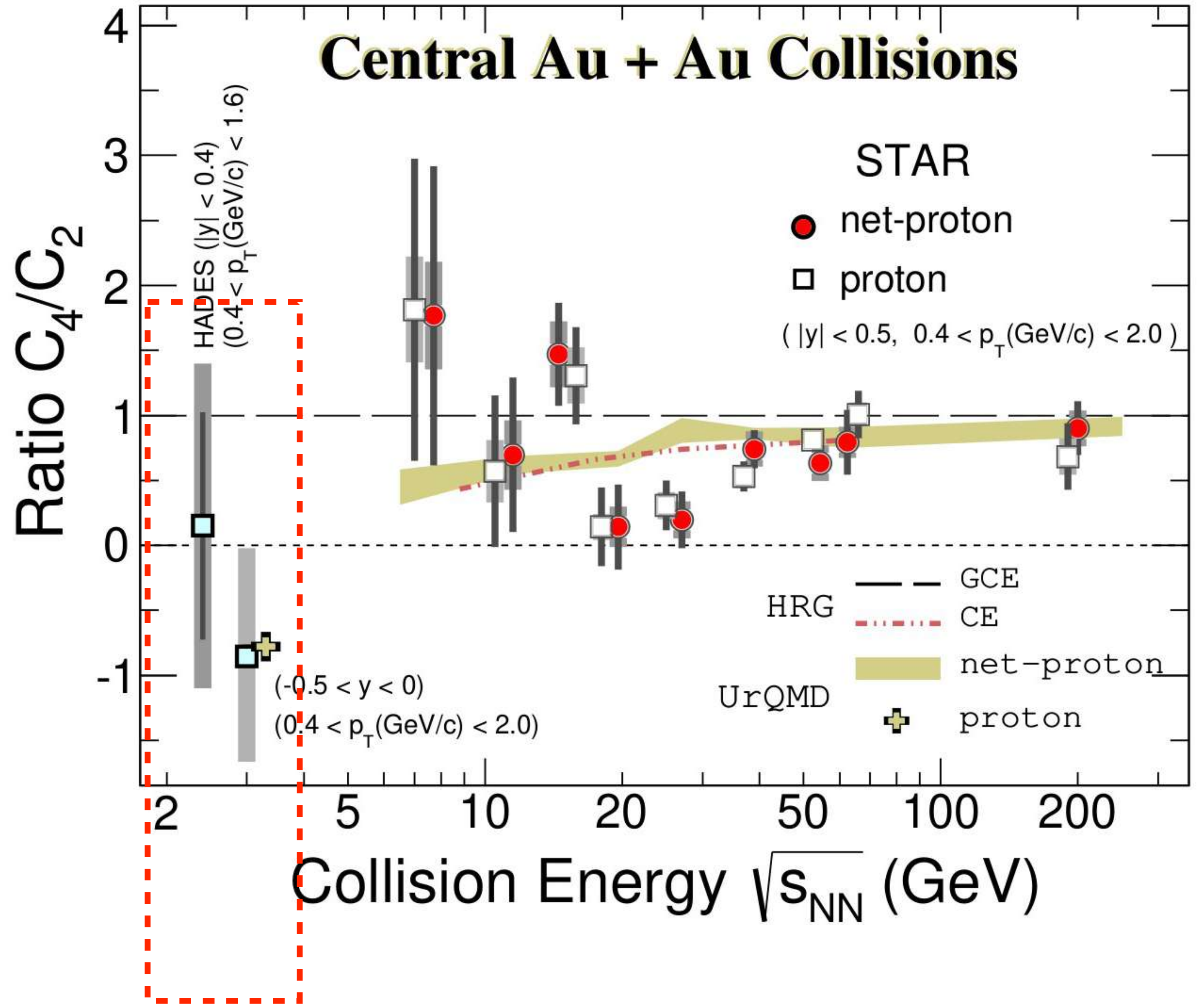
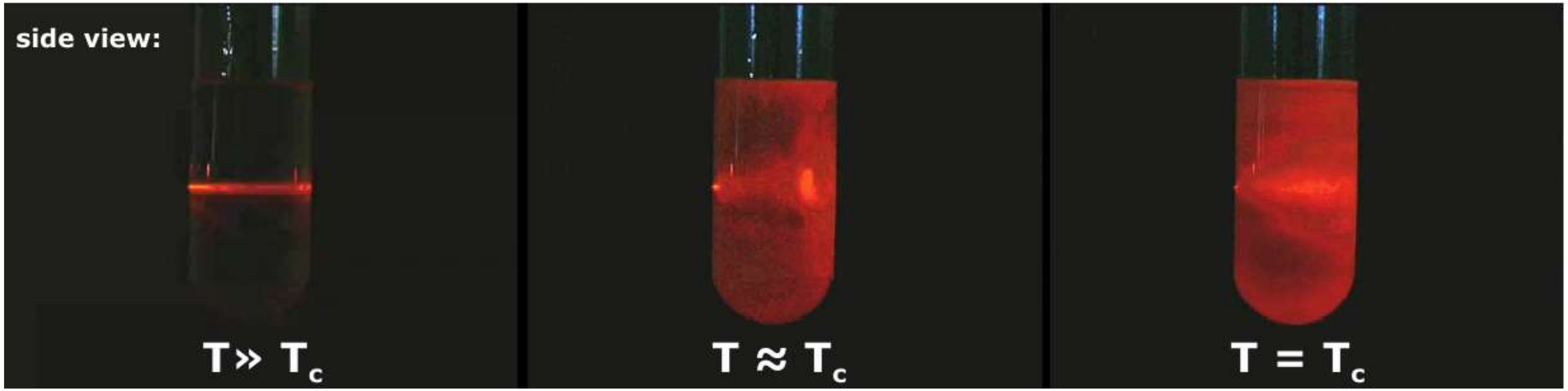


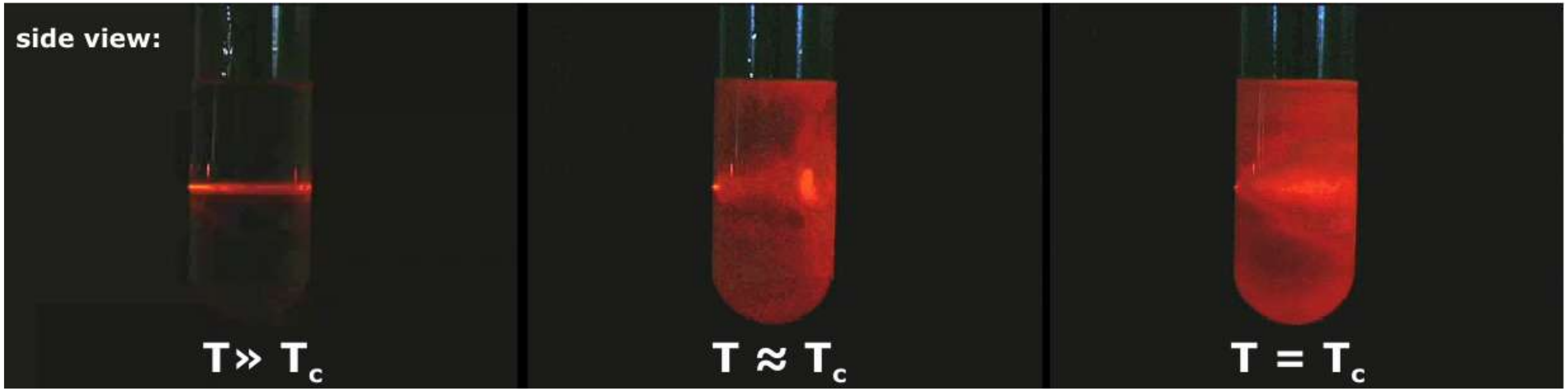
Different acceptance!

Macroscopic Systems



Macroscopic Systems

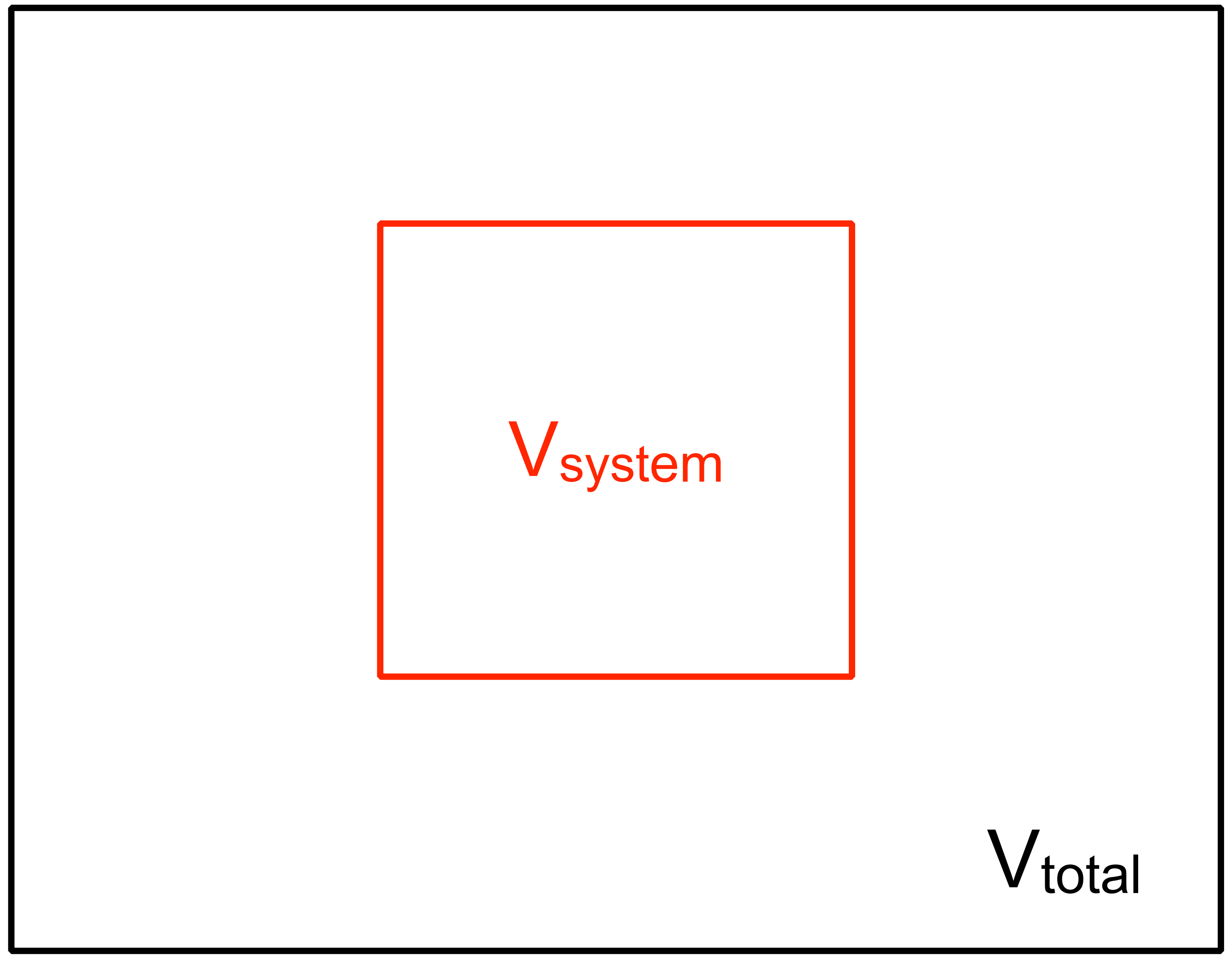
$$\langle N^2 \rangle - \langle N \rangle^2 \sim \langle N \rangle \sim 10^{23}$$



Compare Data with Lattice QCD and other field theoretical models

- Lattice cannot calculate hadron abundances
- Cumulants are well defined quantities
  - Compare cumulants !?
    - Detector fluctuates (efficiency etc…)
    - Volume is not fixed in experiment
      - Possible solution (Rustamov et al, 2211.14849)
    - Baryon number conservation (new development: Vovchenko et al, arXiv 2003.13905, arXiv:2007.03850)
    - Lattice uses grand canonical ensemble
    - Experiment measures protons not all baryons

Grand canonical ensemble

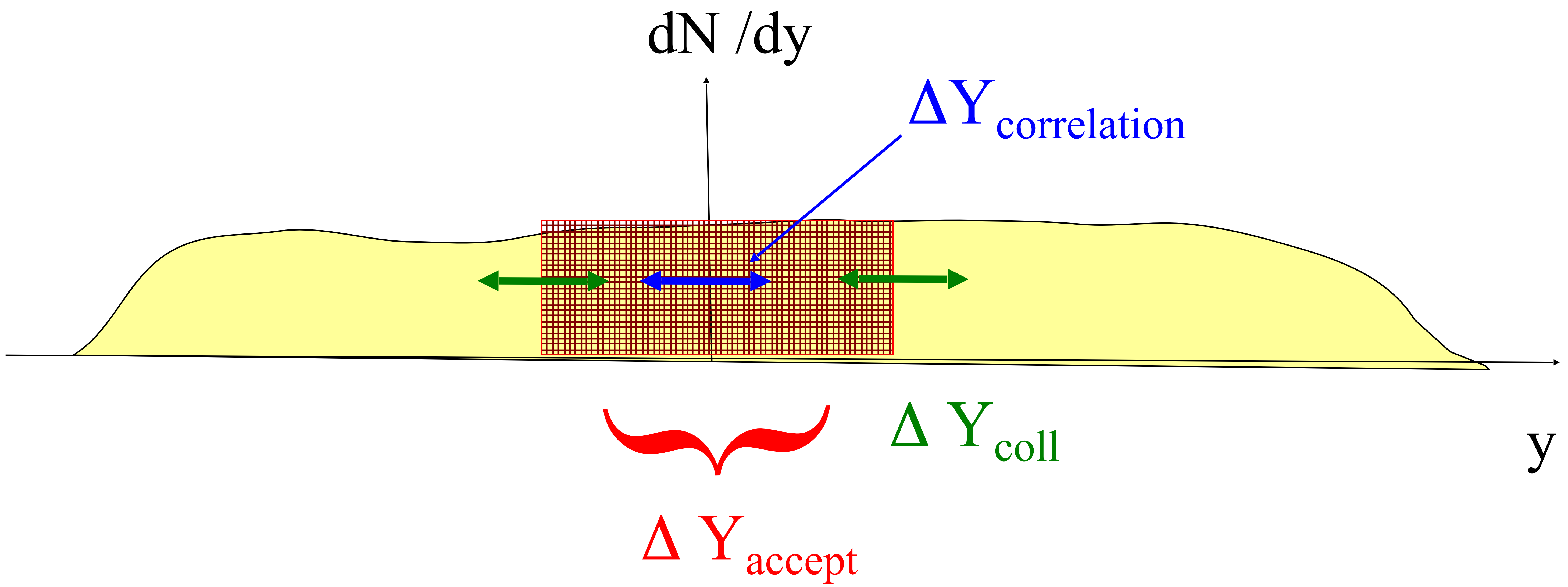


$$V_{total} \to \infty$$
$$V_{system} \to \infty$$

$$\frac{V_{system}}{V_{total}} \to 0$$

In coordinate space!!!!

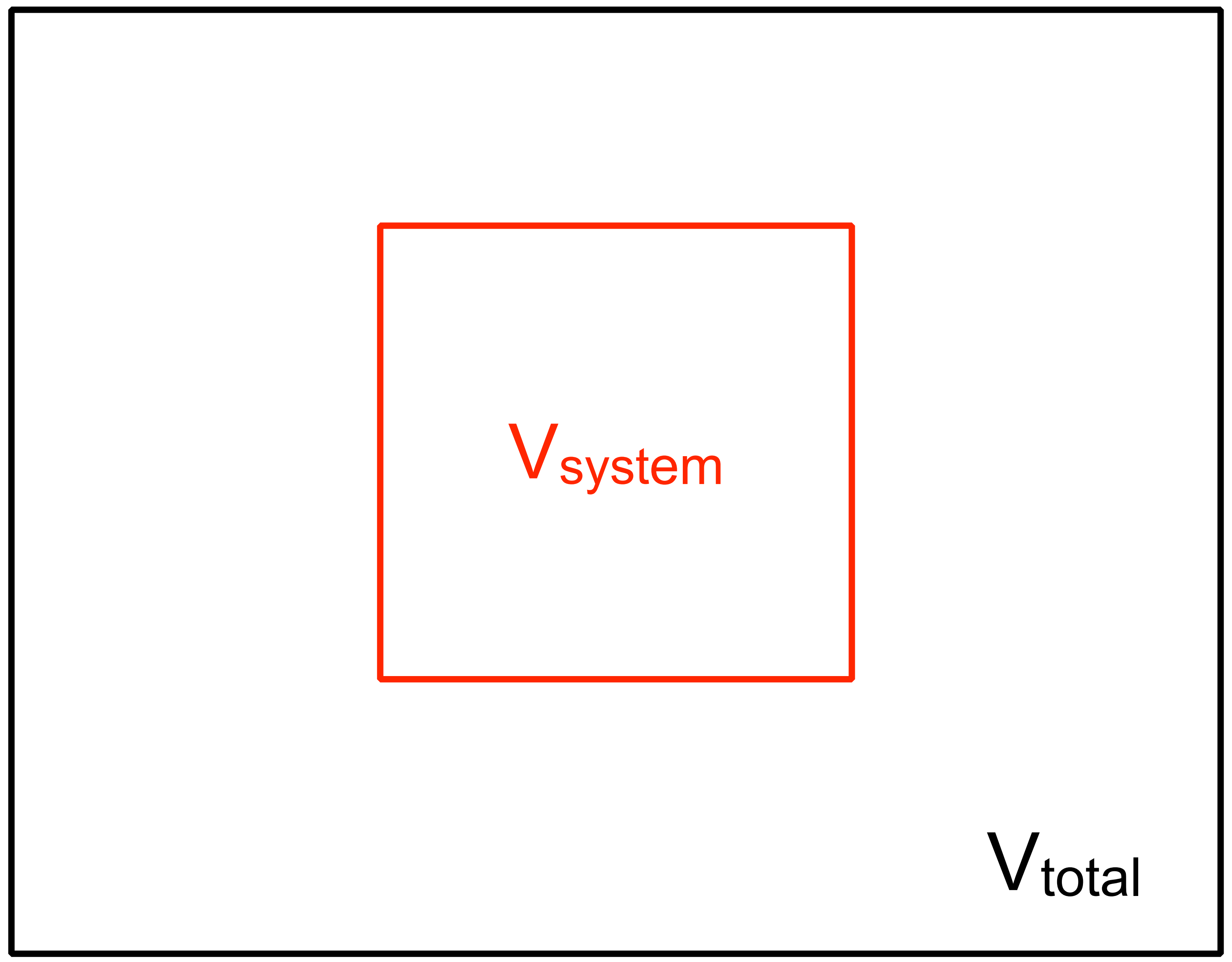
How to make a grand-canonical ensemble in experiment



Conditions for "charge" fluctuations:

- $\Delta Y_{\text{corrrelation}} << \Delta Y_{\text{accept}}$ **(catch the physics)**
- $\Delta Y_{\text{total}} >> \Delta Y_{\text{accept}} >> \Delta Y_{\text{coll}}$ **(keep the physics and minimize charge conservation effect)**

Grand canonical ensemble



$$V_{total} \rightarrow \infty$$

$$V_{system} \rightarrow \infty$$

$$\frac{V_{system}}{V_{total}} \rightarrow 0$$

In coordinate space!!!!

Lattice:

$V_{total} \rightarrow \infty$

grand-canonical ensemble

Coordinate space

Experiment:

V_{total} finite!

$V_{system} \ll V_{total}$ (hopefully)

effect of global charge conservation

Momentum Space

Subensemble acceptance method
(SAM)

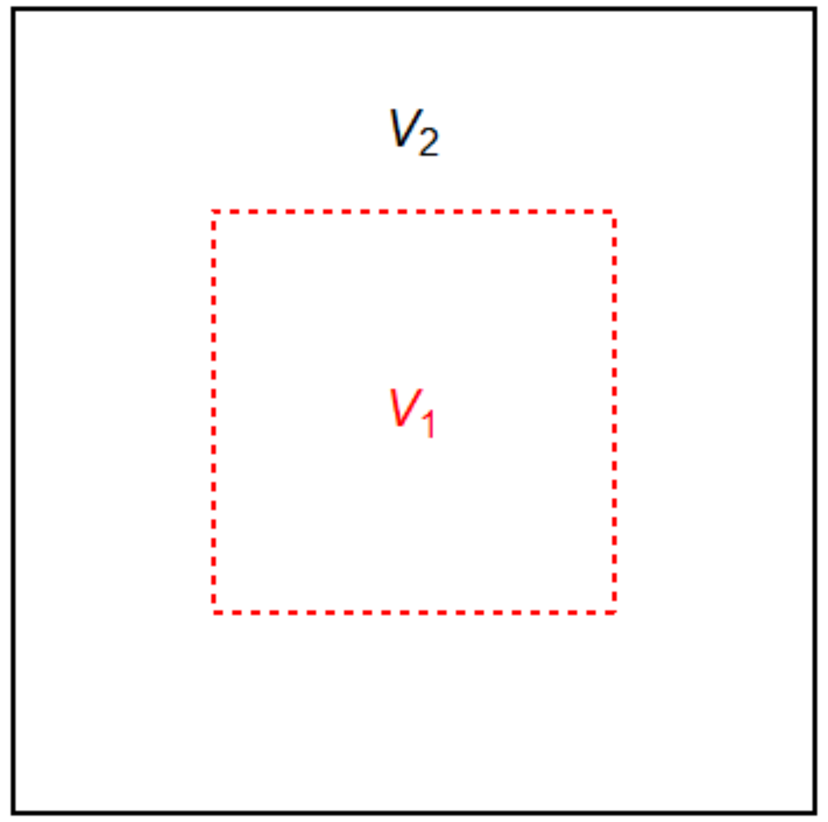
Partition a thermal system with a globally conserved charge *B (canonical ensemble)* into two subsystems which can exchange the charge

$$V = V_1 + V_2$$

Assume thermodynamic limit:

$$V, V_1, V_2 \to \infty; \quad \frac{V_1}{V} = \alpha = const; \quad \frac{V_2}{V} = (1-\alpha) = const;$$

$$V_1, V_2 \gg \xi^3 \qquad \xi = correlation\, length$$



The canonical partition function then reads:

$$Z^{ce}(T, V, B) = \sum_{B_1} Z^{ce}(T, V_1, B_1) Z^{ce}(T, V - V_1, B - B_1)$$

The probability to have charge B_1 *in* V_1 is:

$$P(B_1) \sim Z^{ce}(T, \alpha V, B_1) Z^{ce}(T, (1-\alpha)V, B - B_1), \qquad \alpha \equiv V_1/V$$

Subensemble acceptance method
(SAM)

In the thermodynamic limit, $V \to \infty$, Z^{ce} expressed through free energy *density*

$$Z^{\text{ce}}(T,V,B) \overset{V\to\infty}{\simeq} \exp\left[-\frac{V}{T}f(T,\rho_B)\right]$$

Cumulant generating function for B_1:

$$G_{B_1}(t) \equiv \ln\langle e^{t\,B_1}\rangle = \ln\left\{\sum_{B_1} e^{t\,B_1}\,\exp\left[-\frac{\alpha V}{T}f(T,\rho_{B_1})\right]\exp\left[-\frac{\beta V}{T}f(T,\rho_{B_2})\right]\right\} + \tilde{C}$$

Cumulants of B_1:

$$\kappa_n[B_1] = \left.\frac{\partial^n G_{B_1}(t)}{\partial t^n}\right|_{t=0} \equiv \tilde{\kappa}_n[B_1(t)]|_{t=0}$$

or

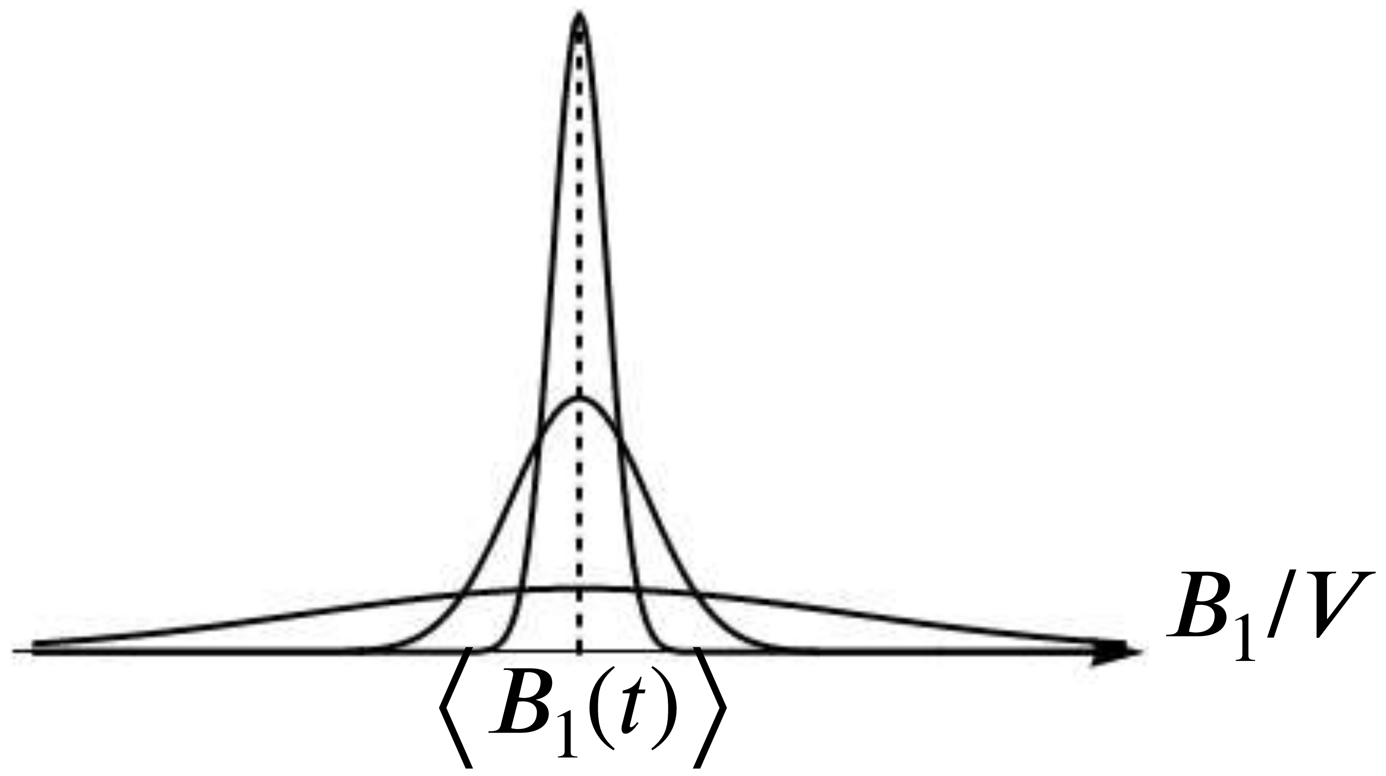
$$\kappa_n[B_1] = \left.\frac{\partial^{n-1}\tilde{\kappa}_1[B_1(t)]}{\partial t^{n-1}}\right|_{t=0}$$

All κ_n can be calculated by determining the *t*-dependent first cumulant $\tilde{\kappa}_1[B_1(t)]$

Making the connection…

$$\tilde{\kappa}_1[B_1(t)] = \frac{\sum_{B_1} B_1 \tilde{P}(B_1;t)}{\sum_{B_1} \tilde{P}(B_1;t)} \equiv \langle B_1(t) \rangle \qquad \text{with} \qquad \tilde{P}(B_1;t) = \exp\left\{ tB_1 - V\,\frac{\alpha f(T,\rho_{B_1}) + \beta f(T,\rho_{B_2})}{T} \right\}.$$

Thermodynamic limit: $\tilde{P}(B_1;t)$ *highly peaked* at $\langle B_1(t) \rangle$



$\langle B_1(t) \rangle$ is a solution to equation $d\tilde{P}/\mathrm{d}B_1 = 0$:

$$t = \hat{\mu}_B[T,\rho_{B_1}(t)] - \hat{\mu}_B[T,\rho_{B_2}(t)] \qquad \text{with} \qquad \hat{\mu}_B \equiv \mu_B/T, \qquad \mu_B(T,\rho_B) = \partial f(T,\rho_B)/\partial \rho_B$$

t = 0:

$$\rho_{B_1} = \rho_{B_2} = B/V,\; B_1 = \alpha B,$$

i.e. conserved charge uniformly distributed between the two subsystems

Second order cumulant

Differentiate condition for maximum of $\widetilde{P}(B_1;t)$,

$$t = \hat{\mu}_B[T, \rho_{B_1}(t)] - \hat{\mu}_B[T, \rho_{B_2}(t)] \qquad (*)$$

$$\frac{\partial(*)}{\partial t}: \qquad 1 = \left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B1}}\right)_T \left(\frac{\partial \rho_{B1}}{\partial \langle B_1 \rangle}\right)_V \frac{\partial \langle B_1 \rangle}{\partial t} - \left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B2}}\right)_T \left(\frac{\partial \rho_{B2}}{\partial \langle B_2 \rangle}\right)_V \frac{\partial \langle B_2 \rangle}{\partial \langle B_1 \rangle} \frac{\partial \langle B_1 \rangle}{\partial t}$$

$$\left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B1,2}}\right)_T \equiv \left[\chi_2^B\left(T, \rho_{B_{1,2}}\right) T^3\right]^{-1}, \qquad \rho_{B_1} \equiv \frac{\langle B_1 \rangle}{\alpha V}, \qquad \rho_{B_2} \equiv \frac{\langle B_2 \rangle}{(1-\alpha)V}, \qquad \langle B_2 \rangle = B - \langle B_1 \rangle, \qquad \frac{\partial \langle B_1 \rangle}{\partial t} \equiv \tilde{\kappa}_2[B_1(t)]$$

Solve the equation for $\widetilde{\kappa}_2$:

$$\tilde{\kappa}_2[B_1(t)] = \frac{V\, T^3}{[\alpha\, \chi_2^B(T, \rho_{B_1})]^{-1} + [(1-\alpha)\, \chi_2^B(T, \rho_{B_2})]^{-1}}$$

t = 0: $$\kappa_2[B_1] = \alpha\,(1-\alpha)\, V\, T^3\, \chi_2^B$$

Higher-order cumulants: iteratively differentiate $\widetilde{\kappa}_2$ w.r.t. t

Full result up to sixth order

$$\kappa_1[B_1] = \alpha\, VT^3\, \chi_1^B$$

$$\beta = 1 - \alpha$$

$$\kappa_2[B_1] = \alpha\, VT^3\, \beta\, \chi_2^B$$

$$\kappa_3[B_1] = \alpha\, VT^3\, \beta\, (1-2\alpha)\, \chi_3^B$$

$$\kappa_4[B_1] = \alpha\, VT^3\, \beta\, \left[\chi_4^B \; - 3\alpha\beta \frac{(\chi_3^B)^2 + \chi_2^B\, \chi_4^B}{\chi_2^B}\right]$$

$$\kappa_5[B_1] = \alpha\, VT^3\, \beta\, (1-2\alpha) \left\{[1-2\beta\alpha]\chi_5^B \; - 10\alpha\beta \frac{\chi_3^B \chi_4^B}{\chi_2^B}\right\}$$

$$\kappa_6[B_1] = \alpha\, VT^3\, \beta\, [1 - 5\alpha\beta(1-\alpha\beta)]\, \chi_6^B + 5\, VT^3\, \alpha^2\, \beta^2 \left\{9\alpha\beta \frac{(\chi_3^B)^2\, \chi_4^B}{(\chi_2^B)^2} - 3\alpha\beta \frac{(\chi_3^B)^4}{(\chi_2^B)^3} \right.$$
$$\left. -2(1-2\alpha)^2 \frac{(\chi_4^B)^2}{\chi_2^B} - 3[1-3\beta\alpha] \frac{\chi_3^B\, \chi_5^B}{\chi_2^B} \right\}$$

$$\chi_n^B = \frac{\partial^n (p/T^4)}{\partial(\mu_B/T)^n}$$

– grand-canonical susceptibilities e.g from Lattice QCD!!

Details: Vovchenko, et al. arXiv:2003.13905

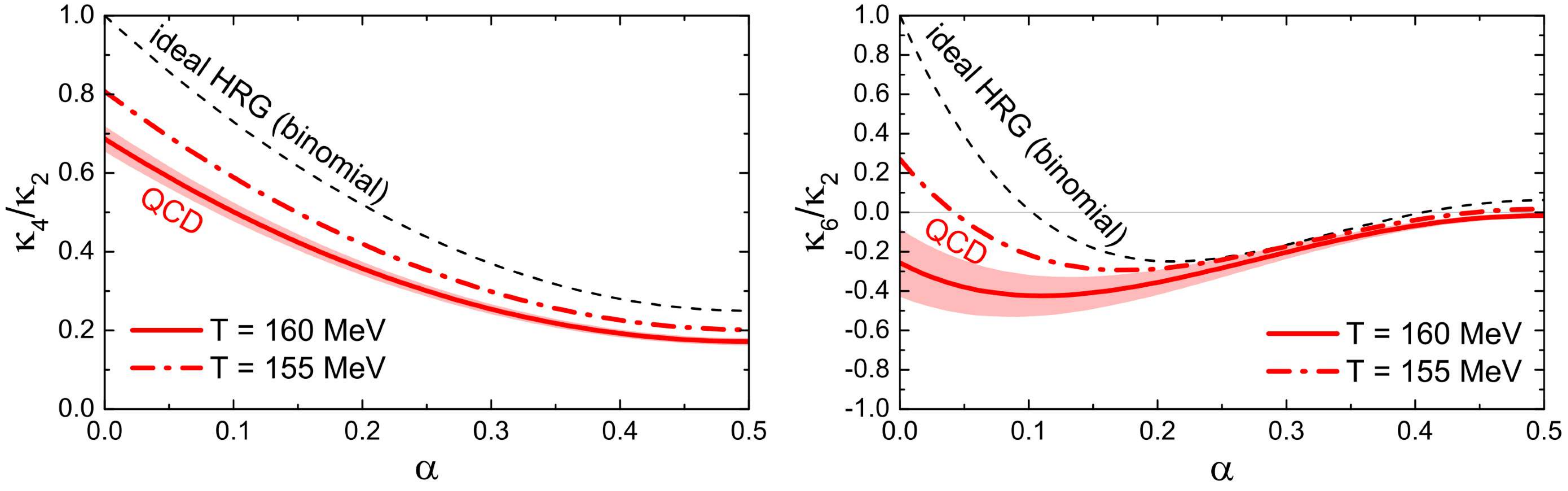
Global charge conservation

Solved for ANY equation of state (including QCD)

V. Vovchenko et al, arXiv 2003.13905, arXiv:2007.03850

$$\left(\frac{\kappa_4}{\kappa_2}\right)_{LHC} = (1-3\alpha\beta)\,\frac{\chi_4^B}{\chi_2^B}$$

$$\left(\frac{\kappa_6}{\kappa_2}\right)_{LHC} = [1-5\alpha\beta(1-\alpha\beta)]\,\frac{\chi_6^B}{\chi_2^B} - 10\alpha(1-2\alpha)^2\beta\left(\frac{\chi_4^B}{\chi_2^B}\right)^2$$



$\alpha =$ fraction of measured baryons

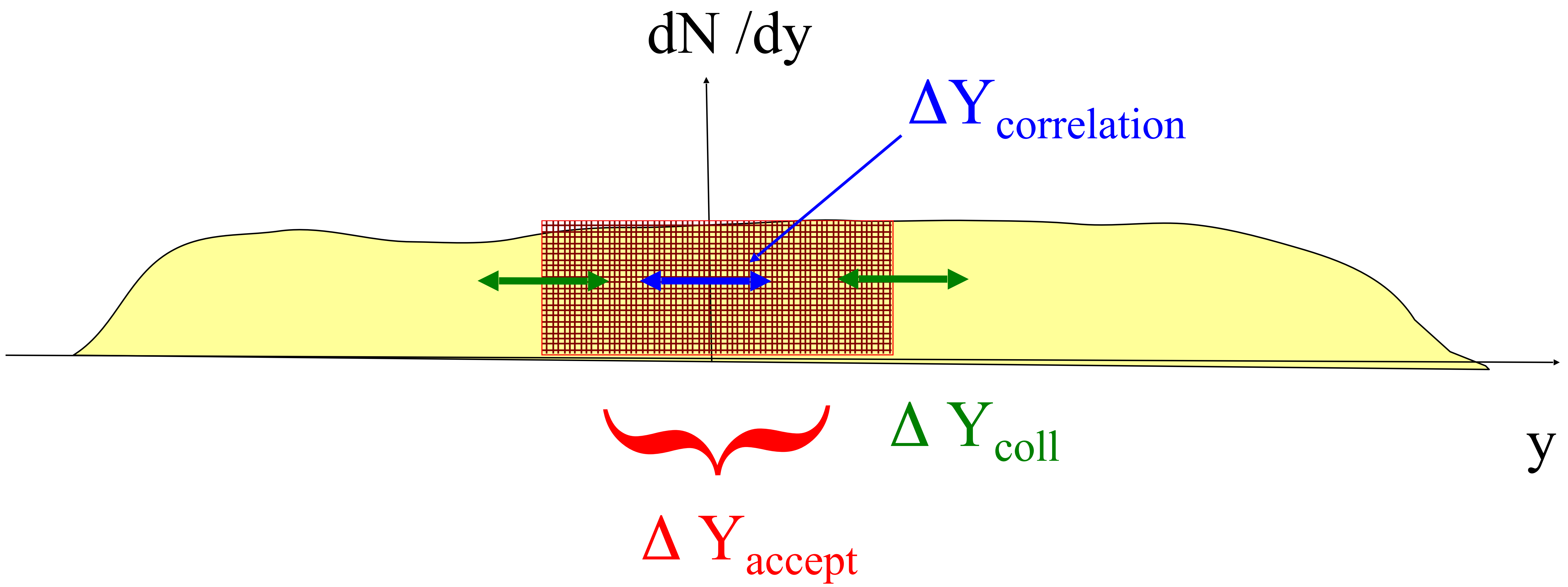
$\beta = 1 - \alpha$

Lattice data for χ_4^B/χ_2^B and χ_6^B/χ_2^B from Borsanyi et al., 1805.04445

Alternative derivation:
M. Barey, and A. Bzdak 2205.05497, 2210.15394

For ideal gas:
Bleicher et al: hep-ph/0006201
Bzdak et al: 1203.4529
Braun-Munzinger et al, 1807.08927

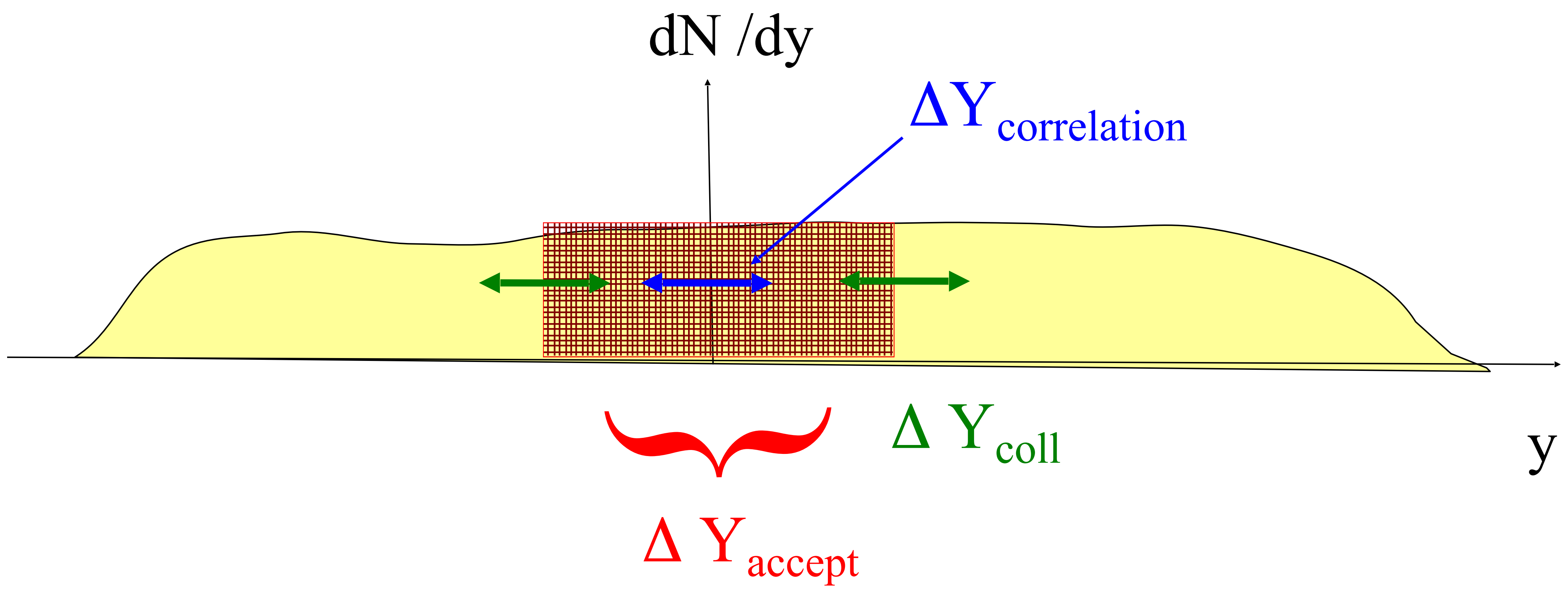
How to make a grand-canonical ensemble in experiment



Conditions for "charge" fluctuations:

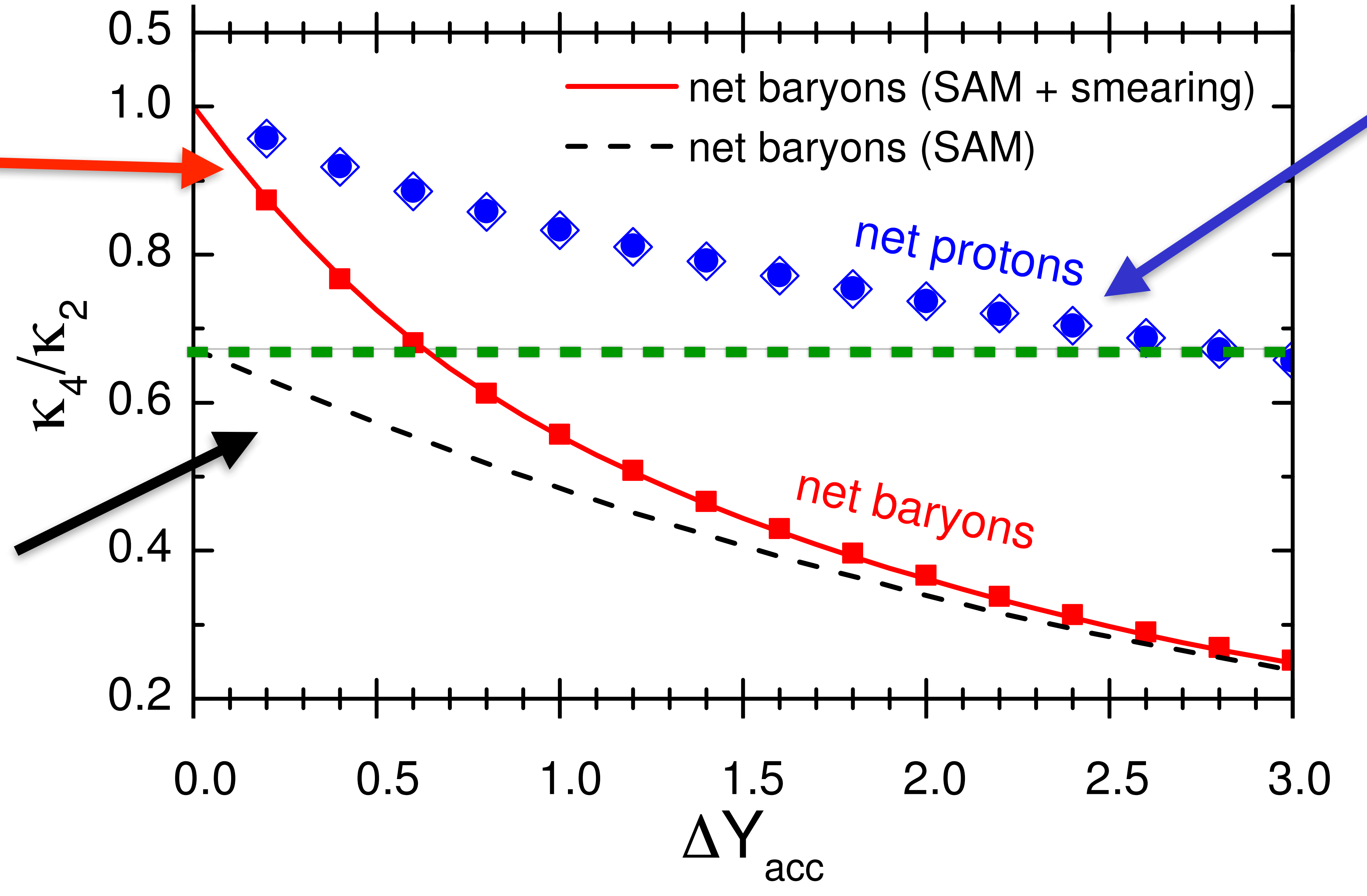
- $\Delta Y_{corrrelation} << \Delta Y_{accept}$ **(catch the physics)**
- $\Delta Y_{total} >> \Delta Y_{accept} >> \Delta Y_{coll}$ **(keep the physics and minimize charge conservation effect)**

Cumulant ratios in "experiment" (μ_B~0)



Lattice +
baryon conservation
+ thermal smearing

Lattice +
baryon conservation



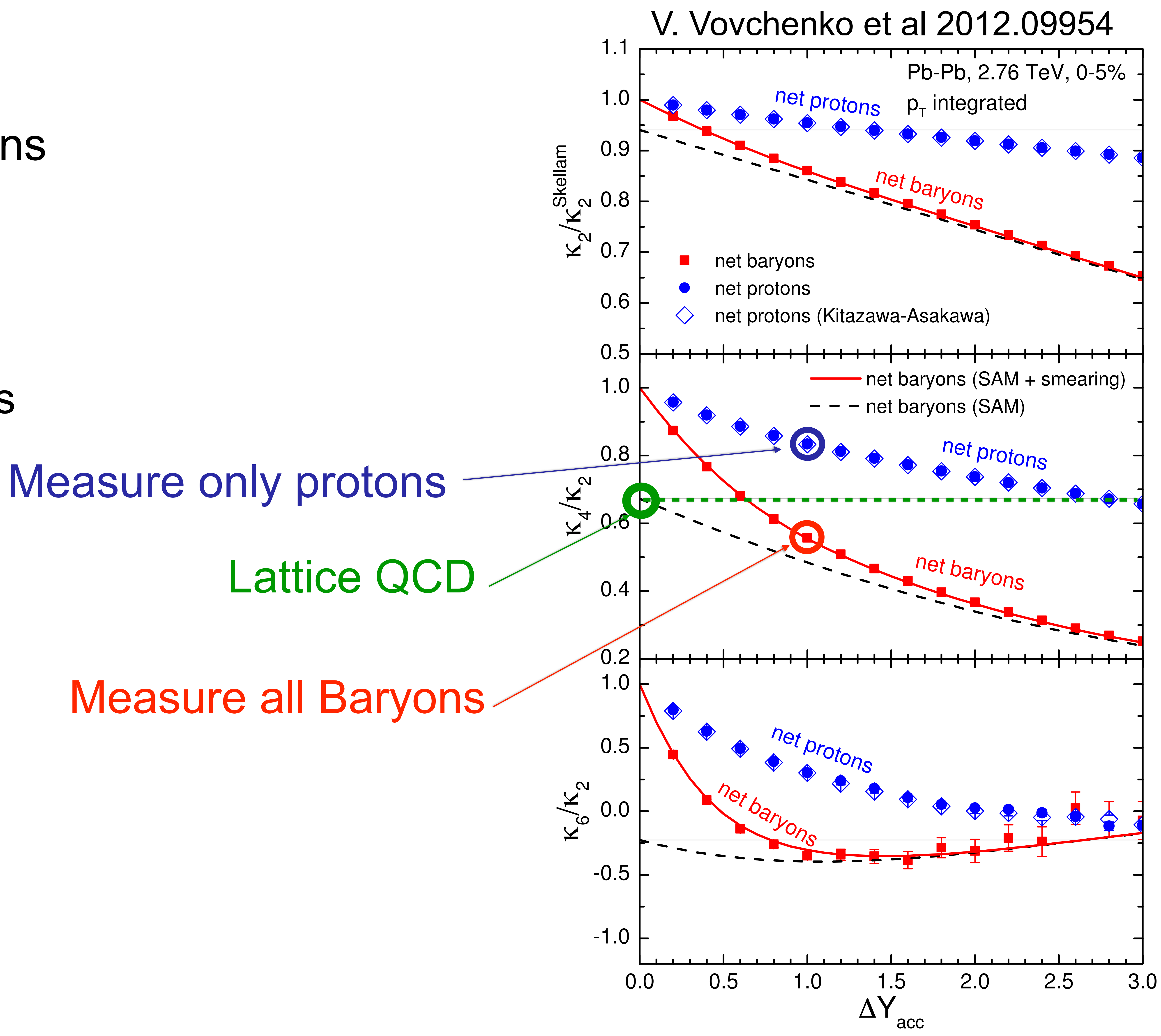
What is REALLY measured

net baryons
net protons
net protons (Kitazawa-Asakawa)

Lattice result

Protons vs Baryons

- Proton are subset of all baryons
  - dilutes the signal
  - need to do binomial unfolding
    - Kitazawa, Asakawa PRC '12
  - Otherwise Apples vs. Oranges



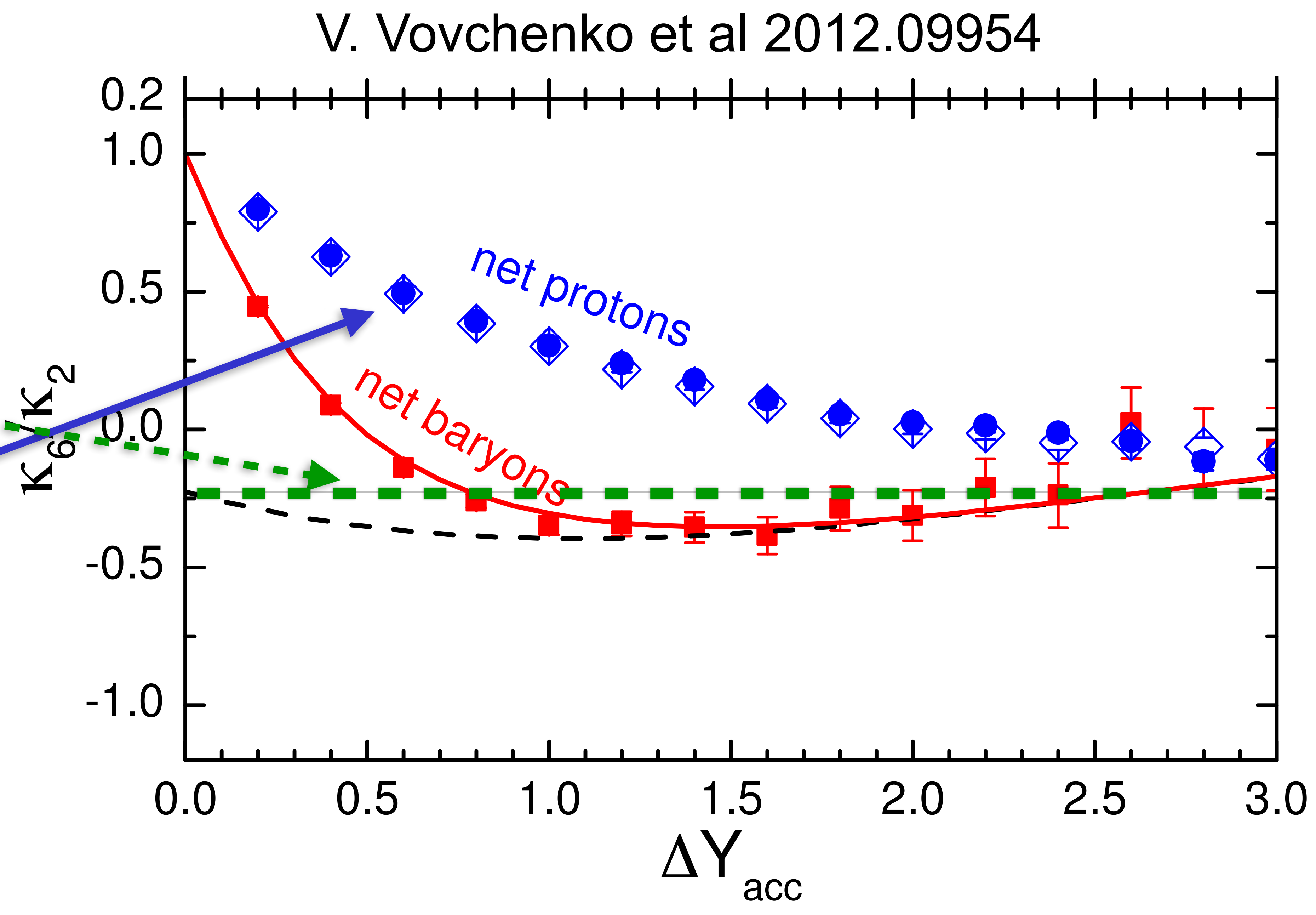
Cumulant ratios in "experiment" (μ_B~0)

NOTE:

$\frac{\kappa_6}{\kappa_2} < 0$ for Baryons (Lattice prediction)

results in

$\frac{\kappa_6}{\kappa_2} > 0$ for PROTONS



$$\frac{\kappa_6}{\kappa_2}$$

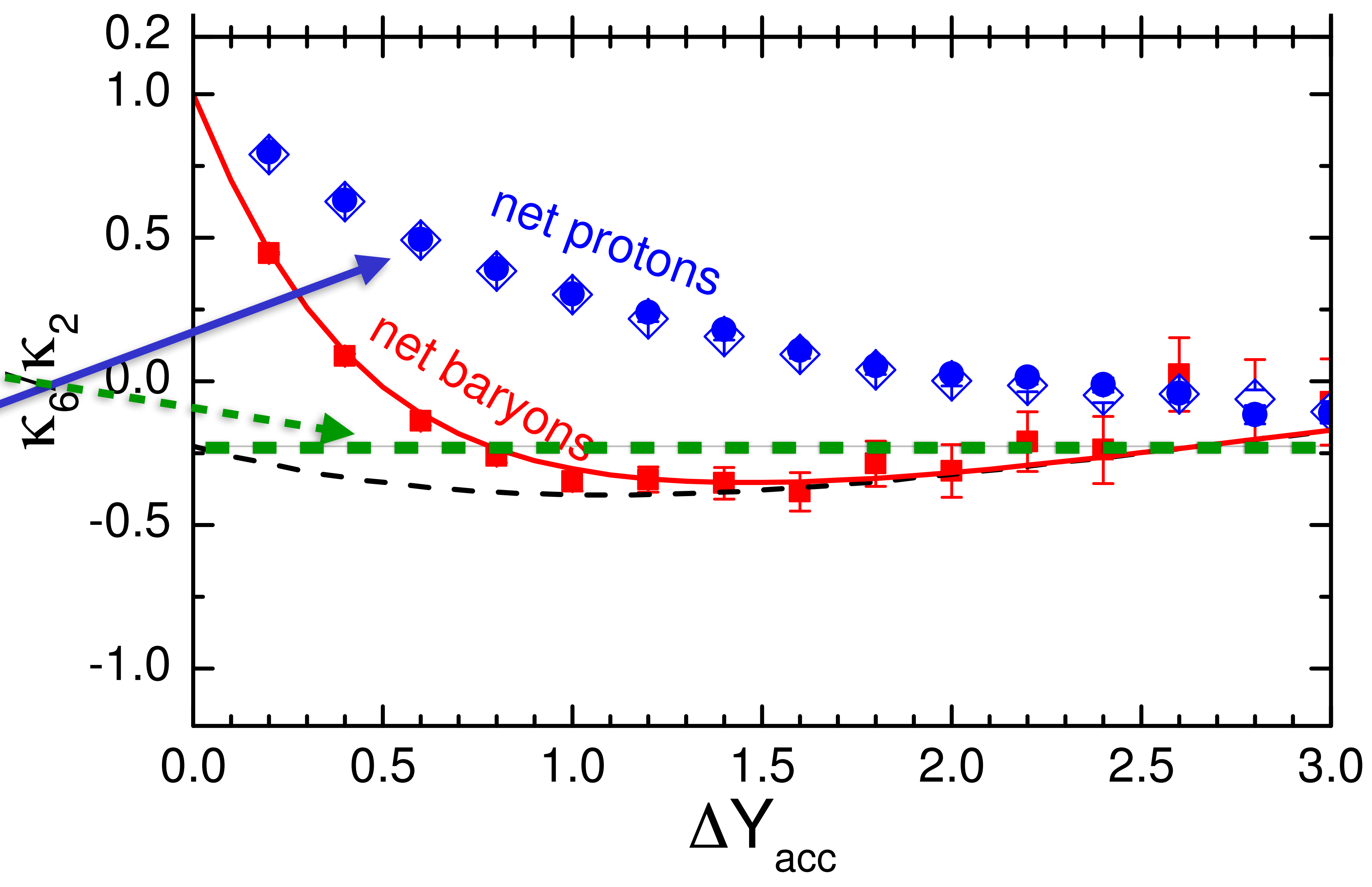
Cumulant ratios in "experiment" (μ_B~0)

NOTE:

$\frac{\kappa_6}{\kappa_2} < 0$ for Baryons (Lattice prediction)

results in

$\frac{\kappa_6}{\kappa_2} > 0$ for PROTONS



$\frac{\kappa_6}{\kappa_2}$

STAR finds at 200 GeV: $\frac{\kappa_6}{\kappa_2} \simeq -5 +/- 4$???? !!!!

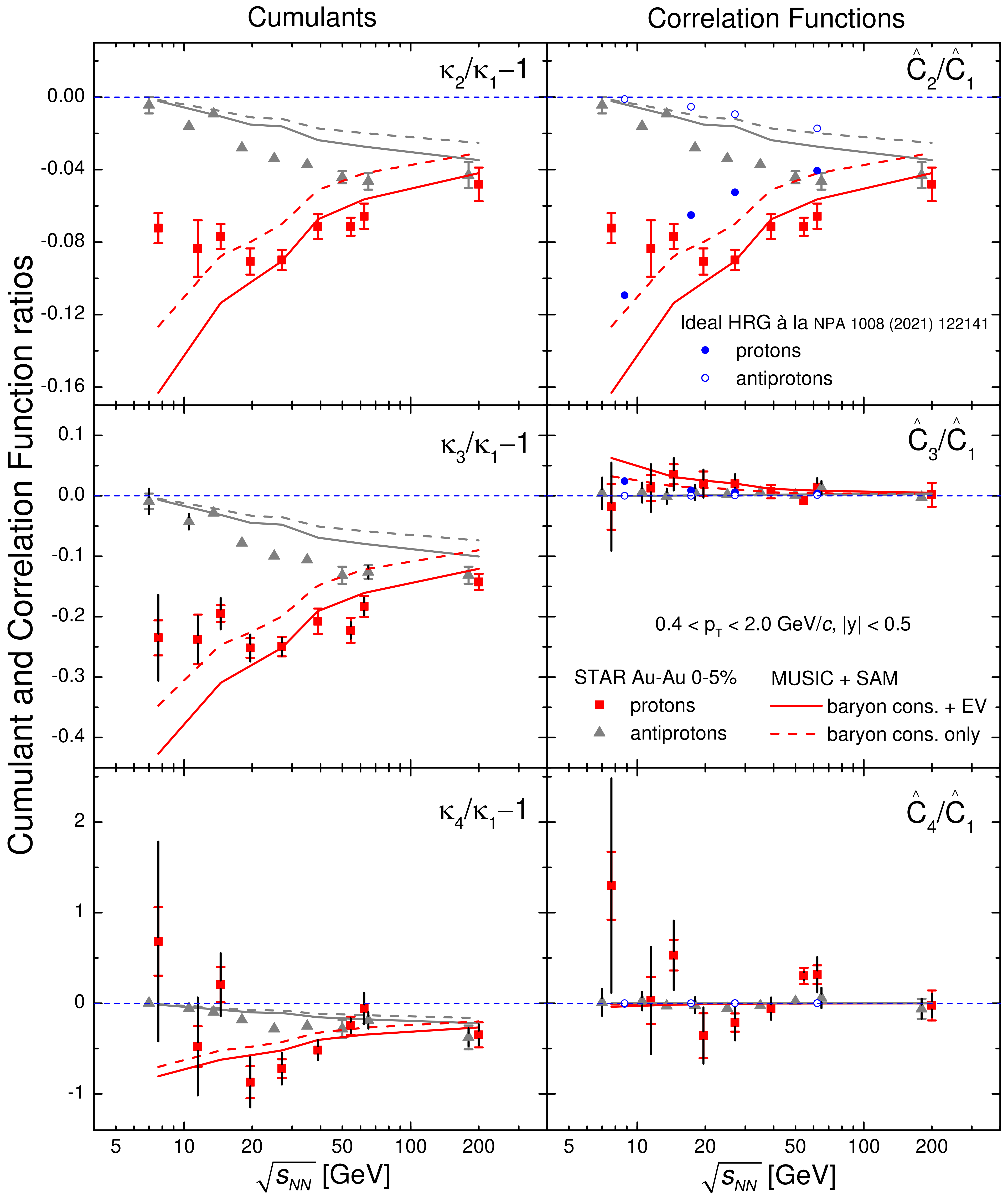
PRL130, 082301 (2023)

Comparison with data from Beam energy scan

Vovchenko et al, 2107.00163

- Viscous hydro
- EOS tuned to LQCD
- Correct for global charge conservation
- Protons NOT baryons

- Baseline!
  No critical point or phase transition



Comparison with data from Beam energy scan

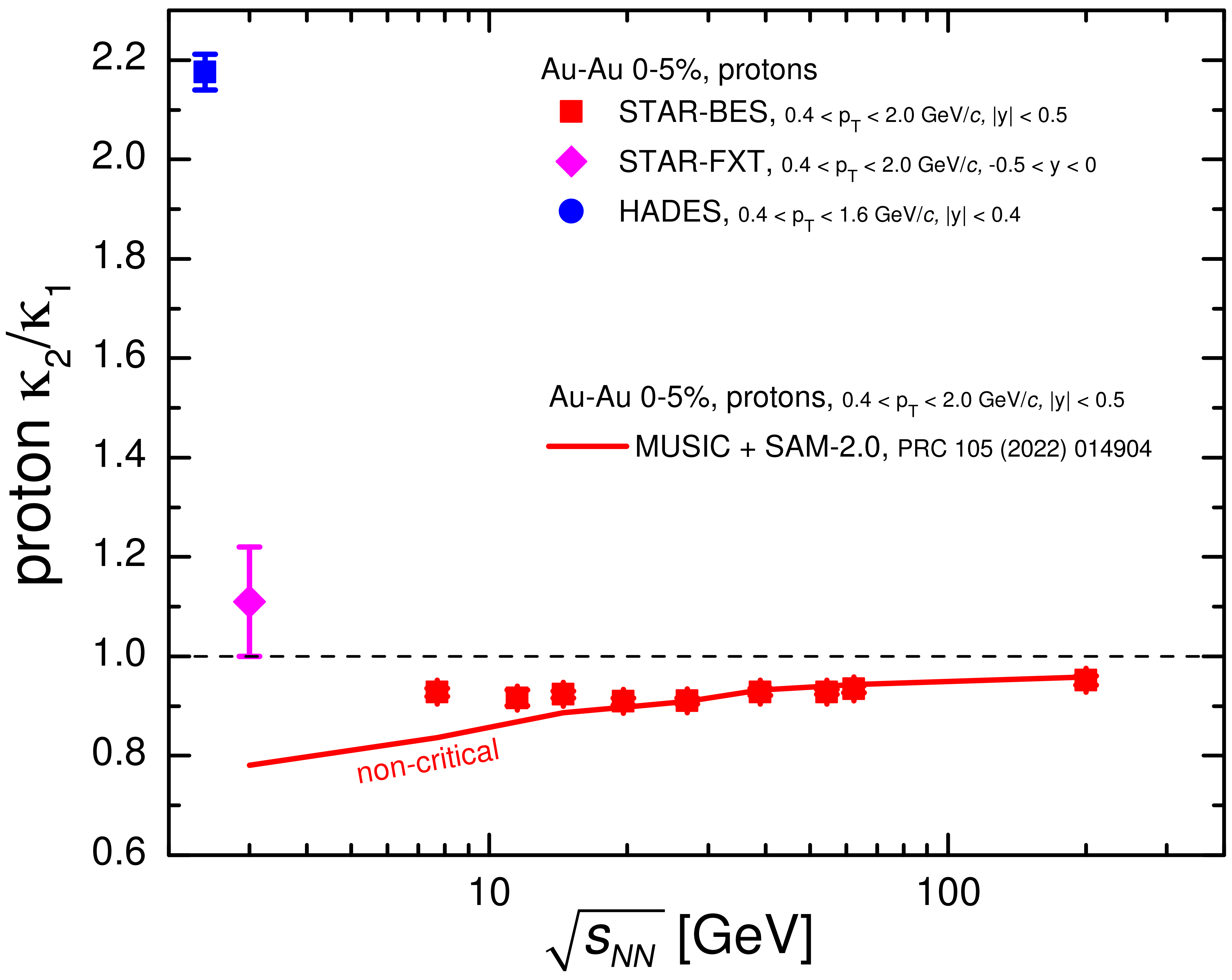


Figure courtesy of V. Vovchenko

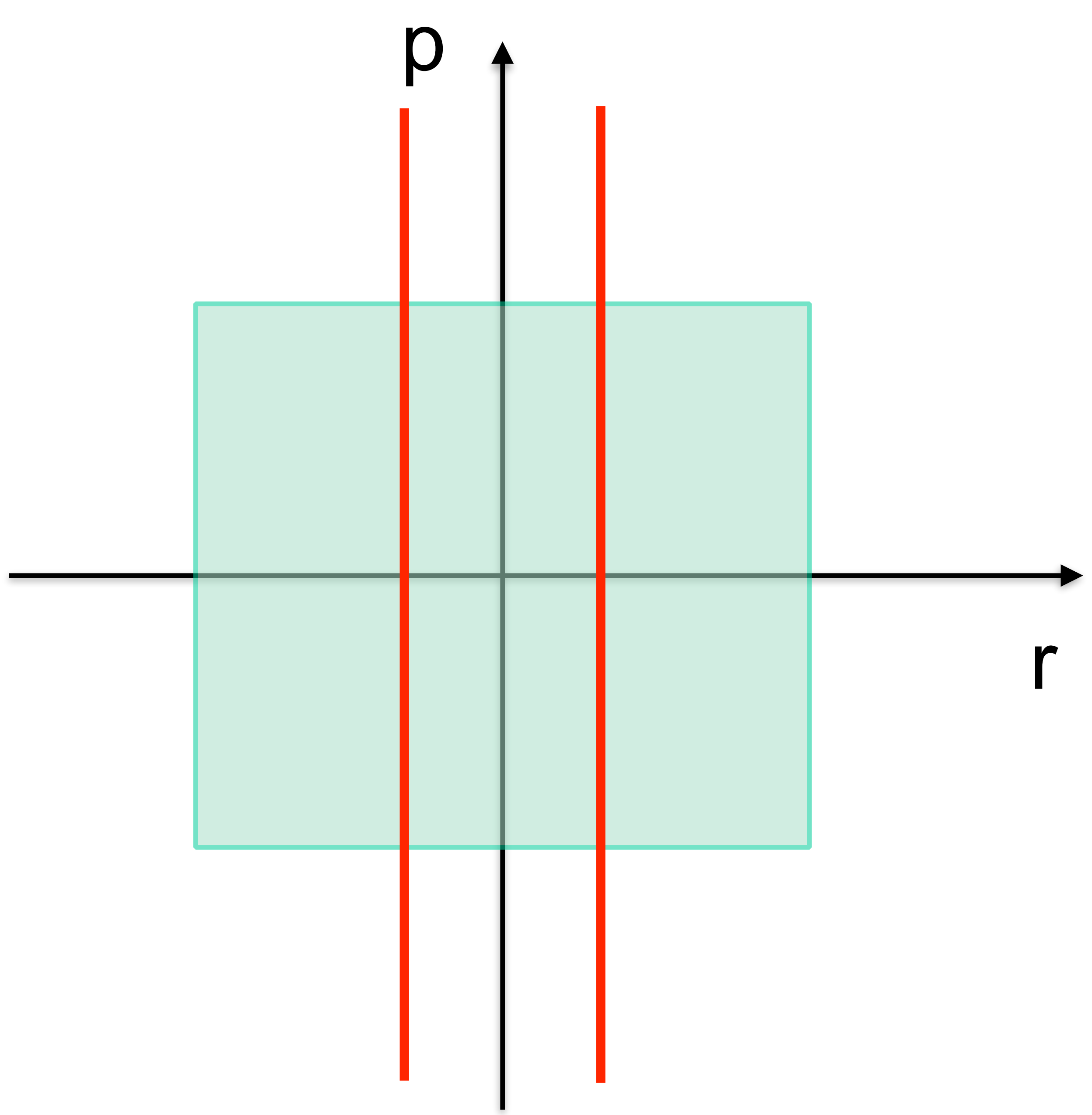
Let's understand the second order cumulants first!

Compare Data with Lattice QCD and other field theoretical models

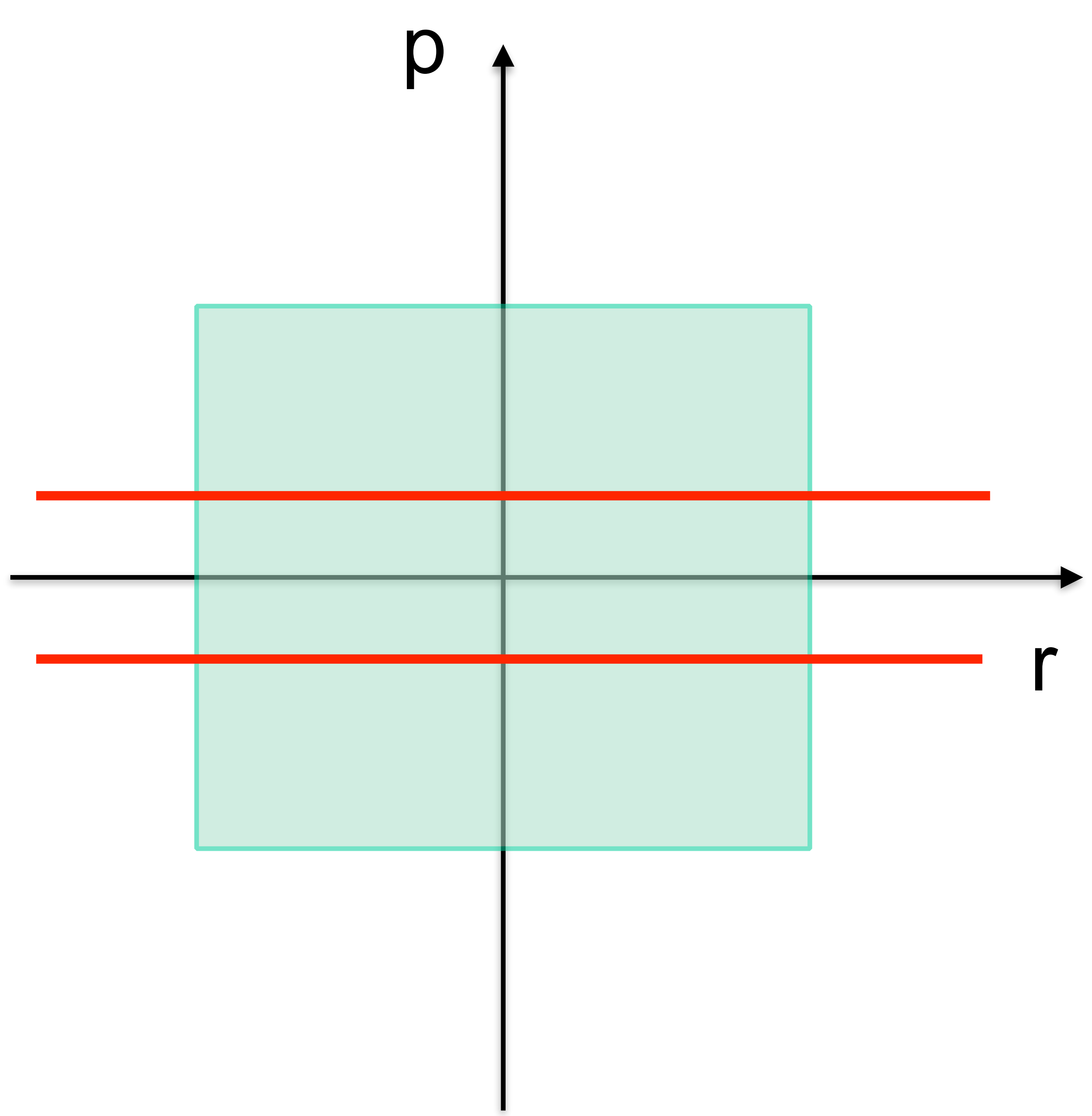
- Lattice cannot calculate hadron abundances
- Cumulants are well defined quantities
- Compare cumulants !?
  - Detector fluctuates (efficiency etc…)
  - Experiment measures protons not all baryons
  - Volume is not fixed in experiment
    - Possible solution (Rustamov et al, 2211.14849)
  - Baryon number conservation
    - Lattice uses grand canonical ensemble
    - Experiment measures protons not all baryons
  - Experiment cuts **momentum** space, Theory cuts **configuration** space

Coordinate vs momentum space cuts

Small Spatial Volume



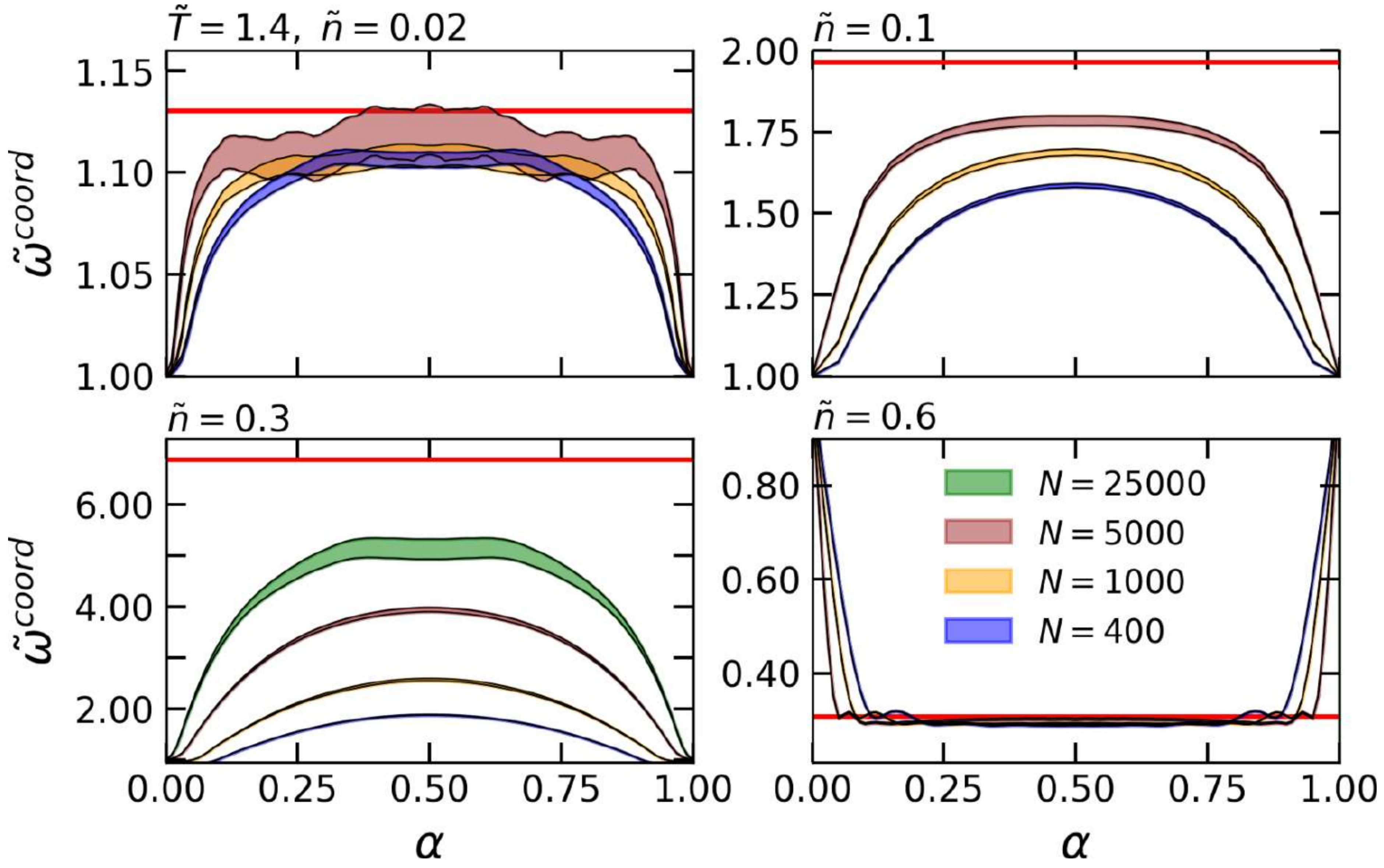
Limited Acceptance in p



classical Molecular dynamics: $H = \frac{p^2}{2m} + V(x_1 - x_j)$

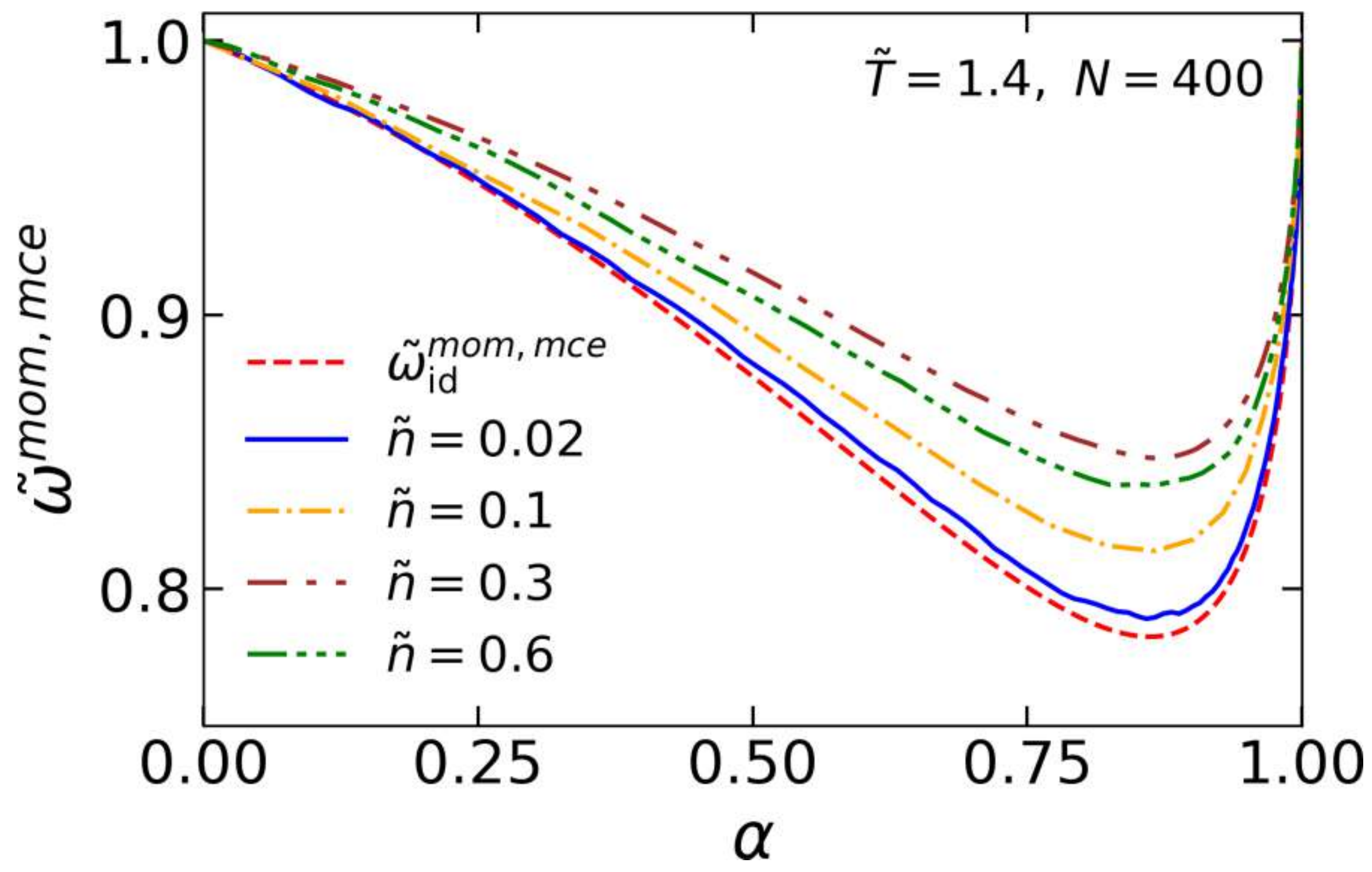
partition function $Z = \int \prod_i (dp_i dx_i) e^{[} = \beta H = \int \prod_i dp_i e^{\frac{p^2}{2m}} \int \prod_i dx_i e^{V(x_i - x_j)} = Z_p \times Z_x$ factorizes!

Correlations live in coordinate space



Cut in coordinate space
Integrate over all Momenta

Fluctuations close to expectation from grand canonical
Correlations clearly visible



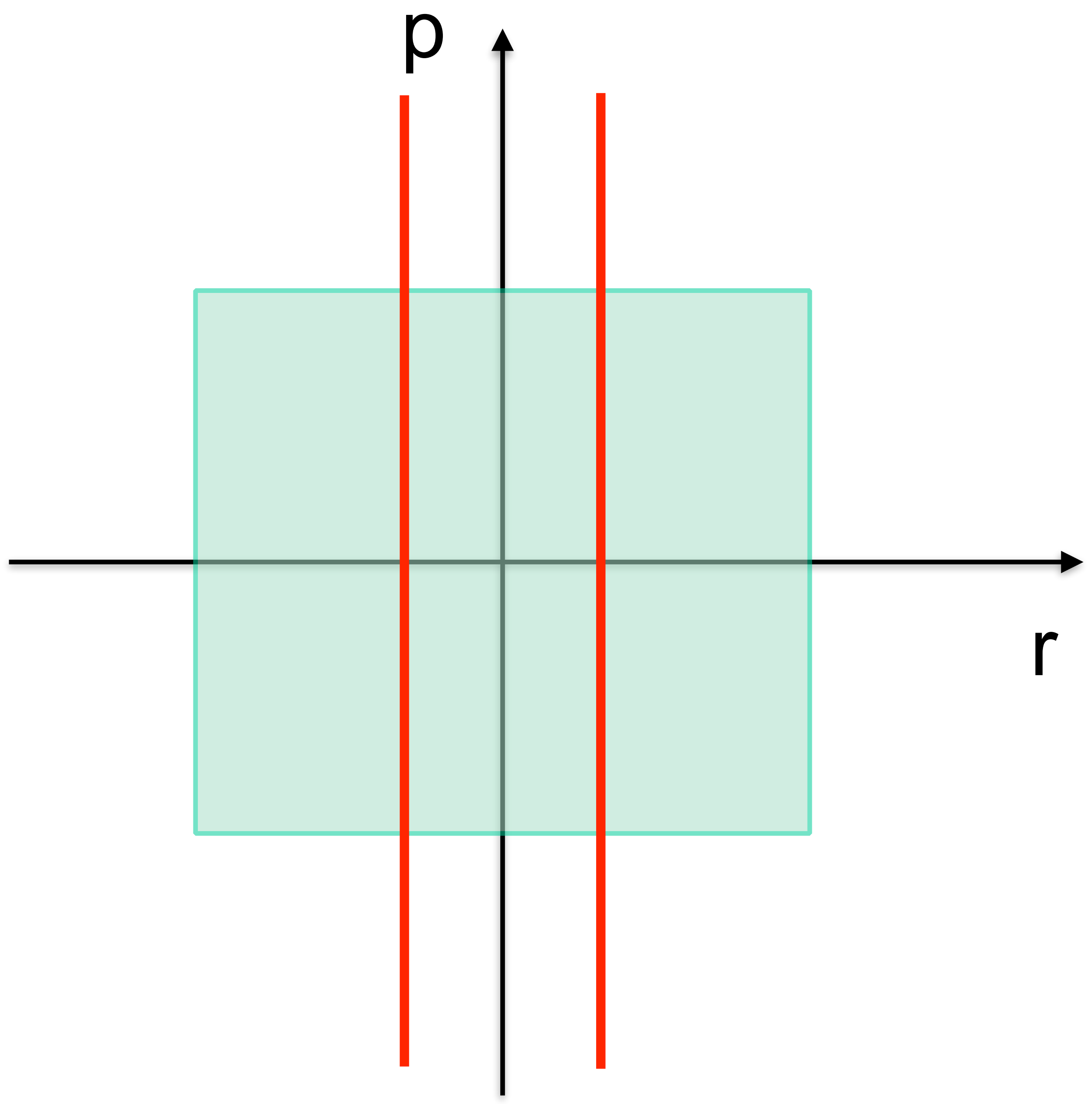
Cut in momentum space
Integrate over all Space

Fluctuations close to non-interacting gas
NO correlations or criticality visible

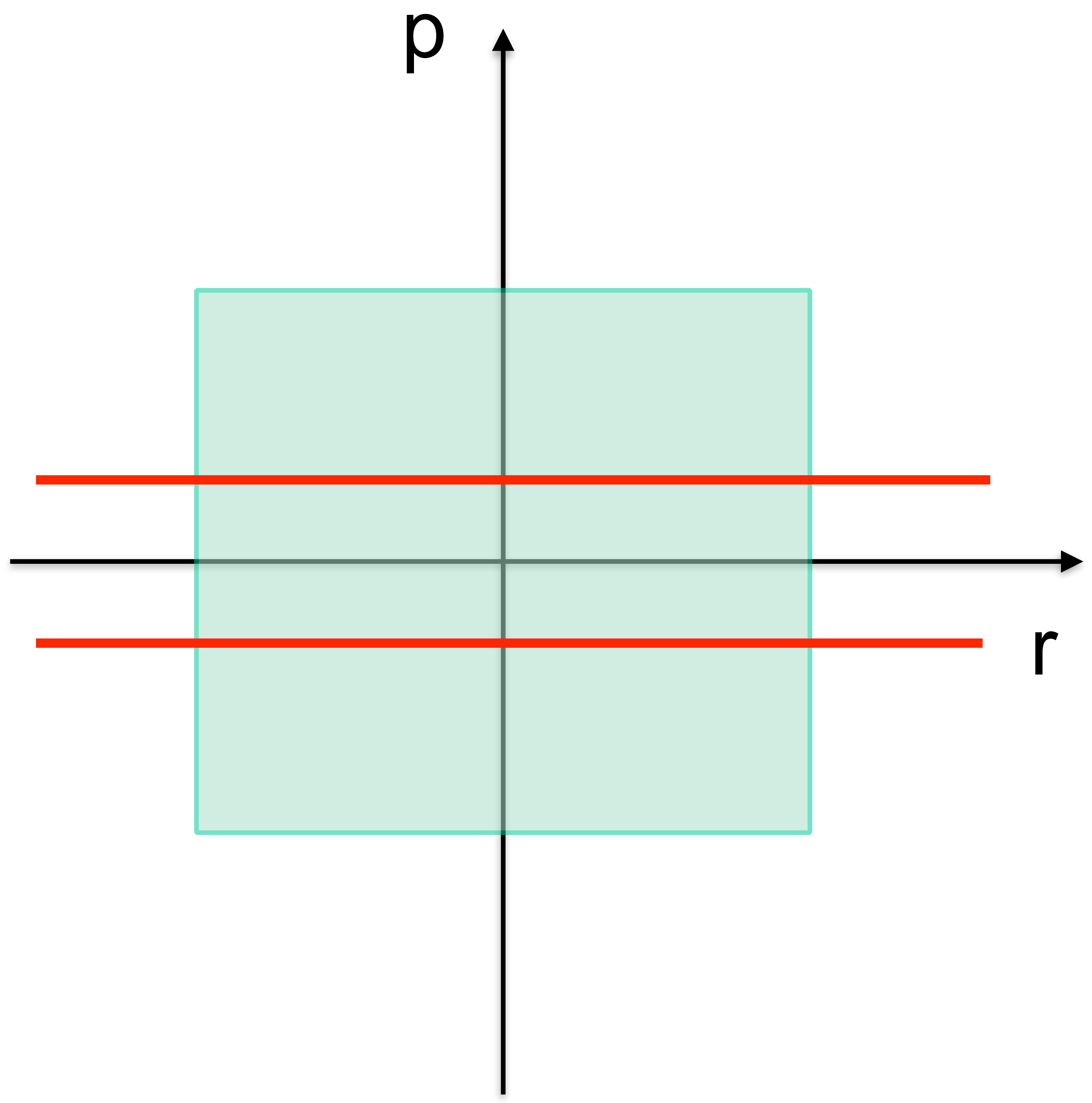
Need Space-Momentum correlations → Flow!

Coordinate vs momentum space cuts

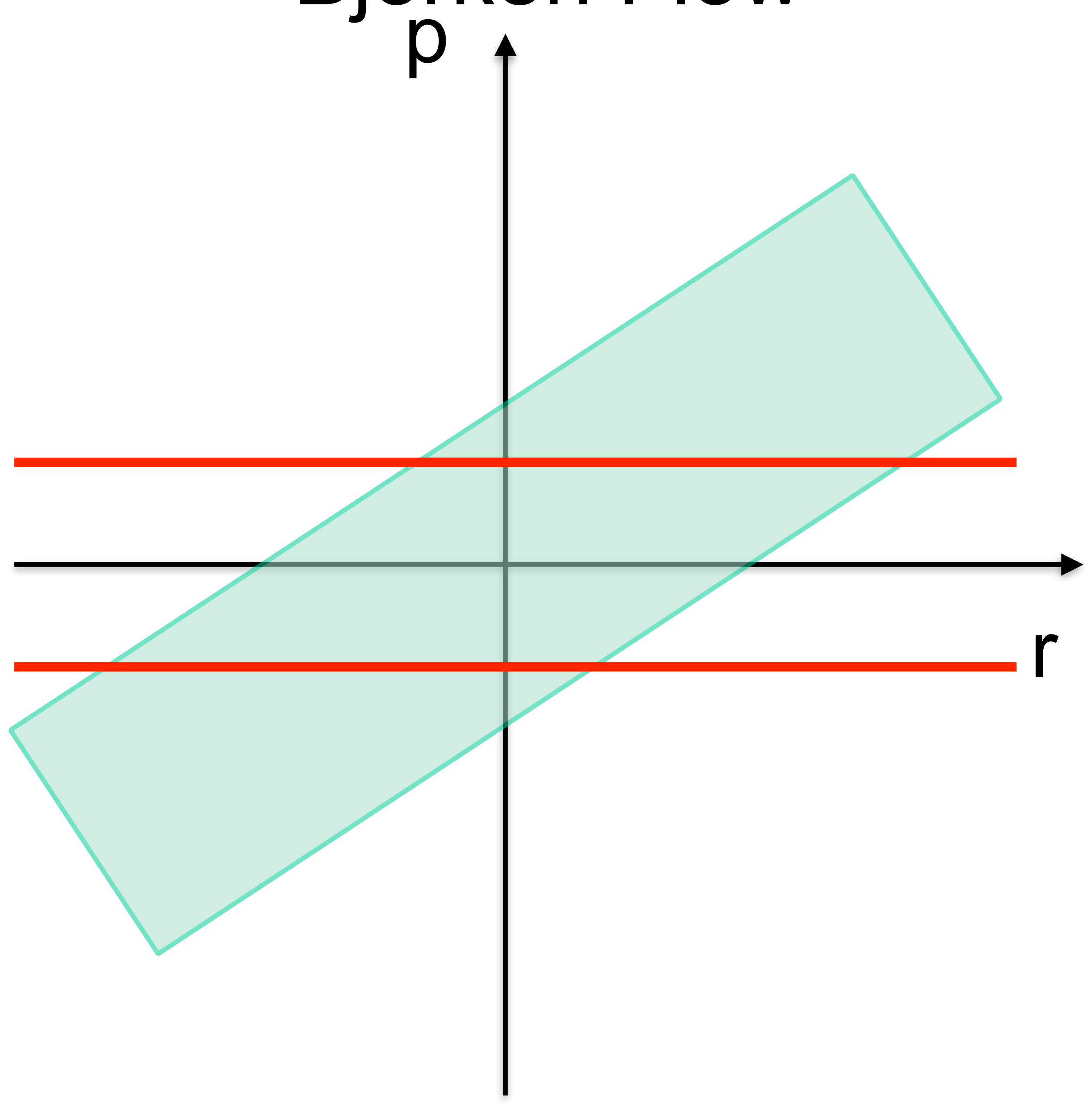
Small Spatial Volume



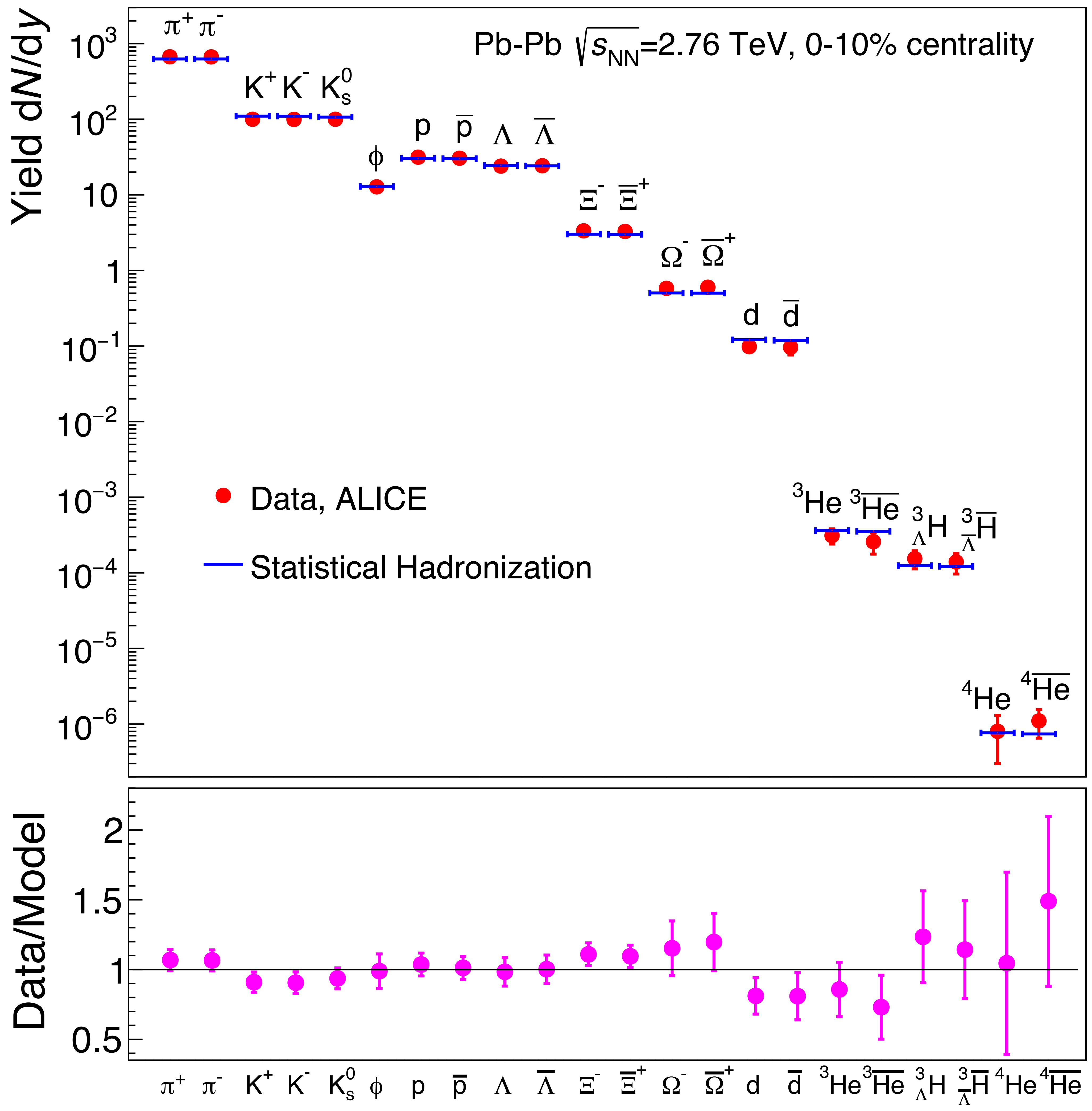
Limited Acceptance in p
No Flow



Limited Acceptance in p
Bjorken Flow



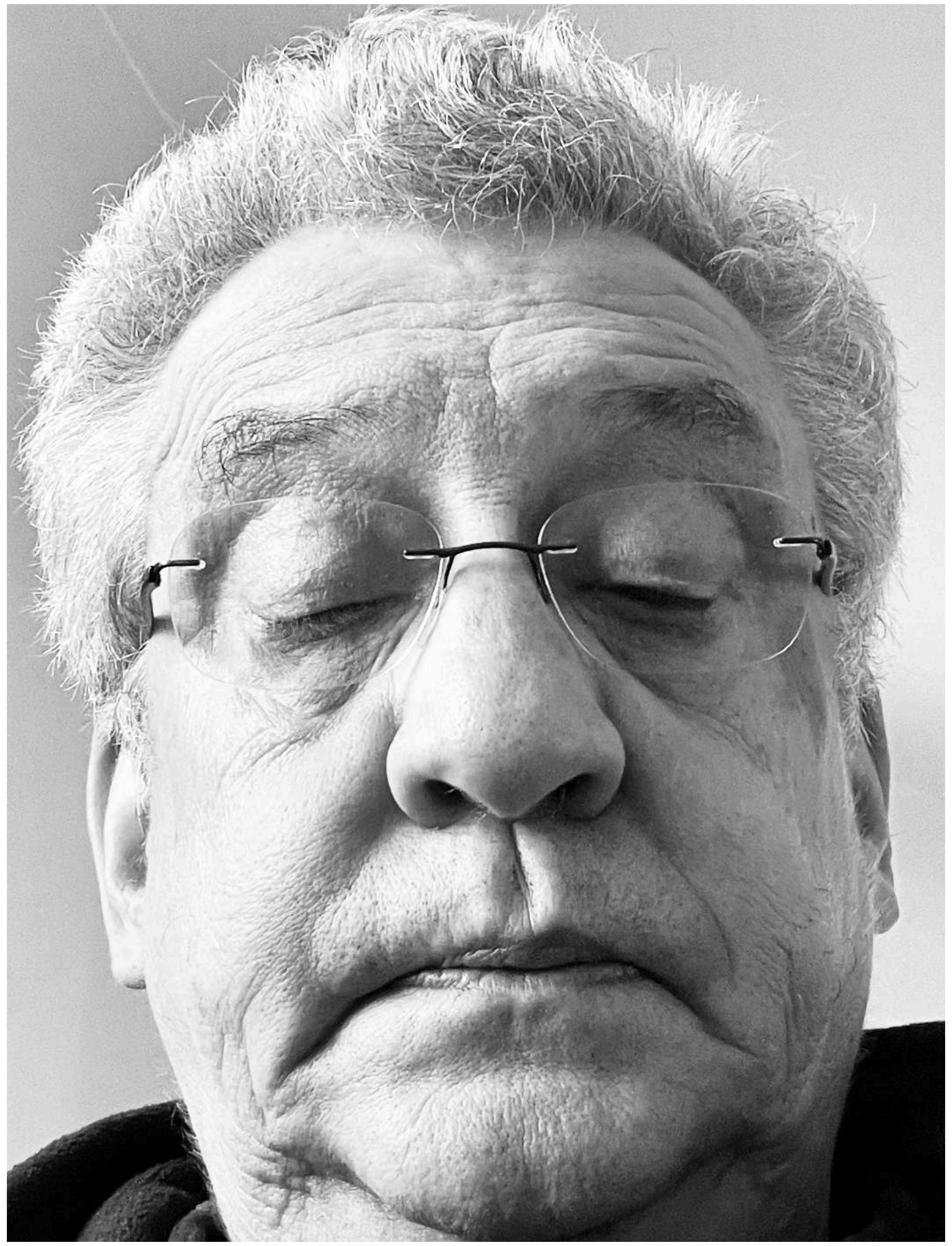
Proton annihilation in the hadronic phase ?



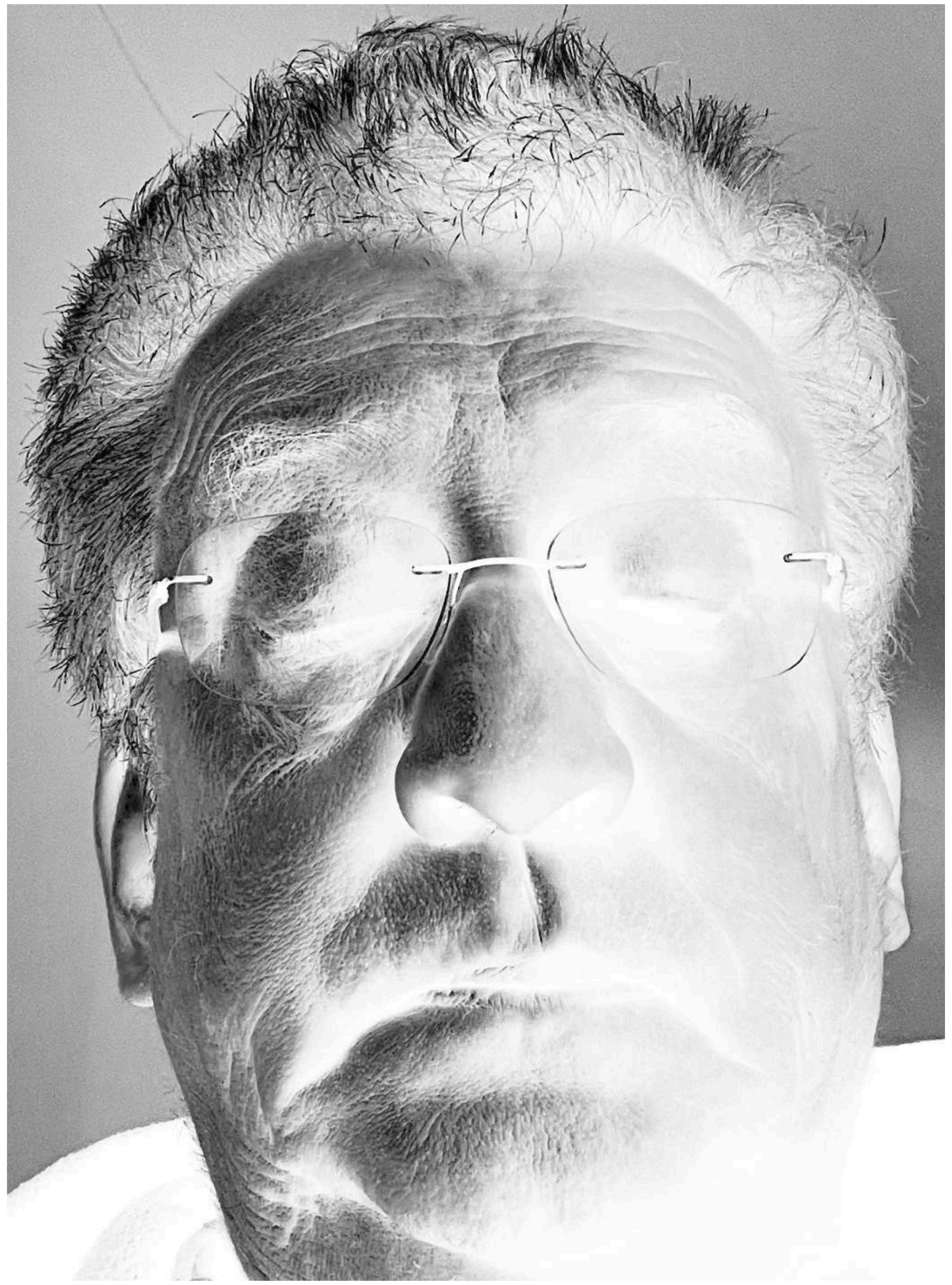
Andronic et al, 2101.05747

- Thermal Model with phase shift corrections:
  - No room for annihilation in hadronic phase

Me



$\overline{\text{Me}}$



Why the discussion?

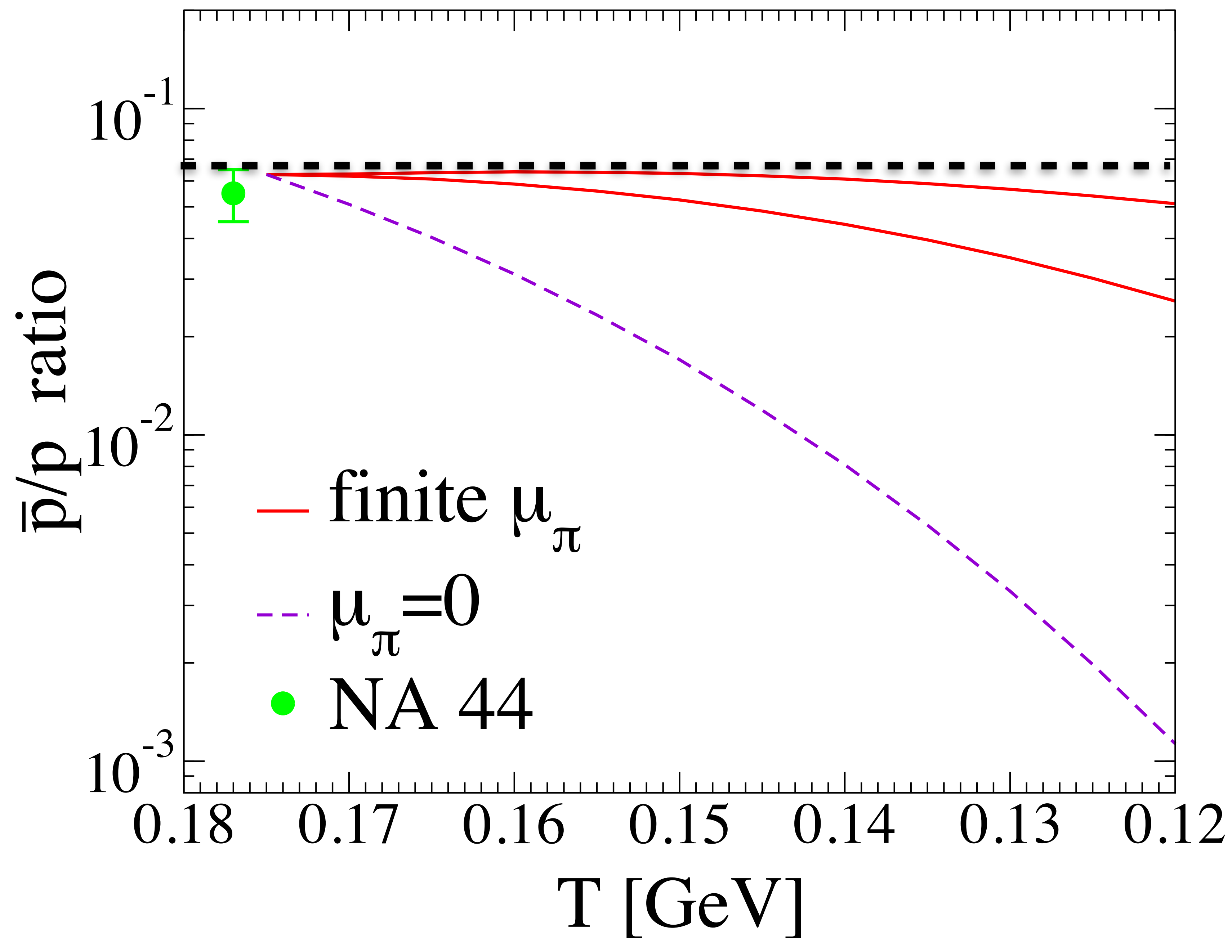
- Lifetime of hadronic phase is short
- pion number effectively conserved

  - $4\pi \Leftrightarrow 2\pi$ suppressed (chiral symmetry)
- $\Rightarrow$ finite μ_π
- increased re-generation of anti-protons

  - $5\pi \Leftrightarrow p + \bar{p}$
- Most transport calculations violate
  detailed balance
  exceptions:
  E. Seifert, W. Cassing, PRC 97 (2018) 024913,
  O. Garcia-Montero et al, Phys. Rev. C 105 (2022) 064906

.

Rapp, Shuryak, PRL 86 (2001) 2980;

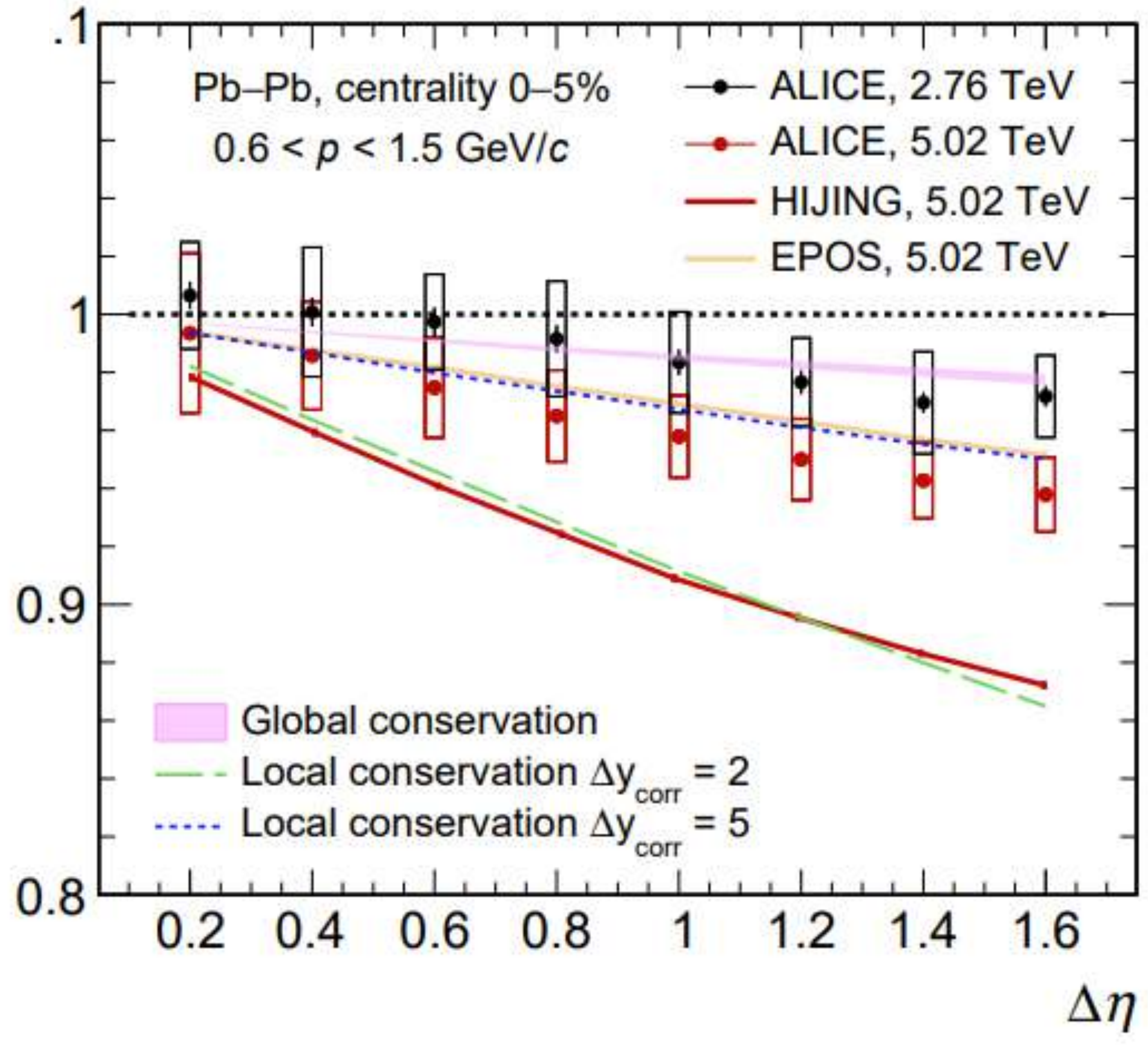


Need additional data to settle this issue

"Local" or "Global" Charge conservations

ALICE Coll., arXiv:2206.03343

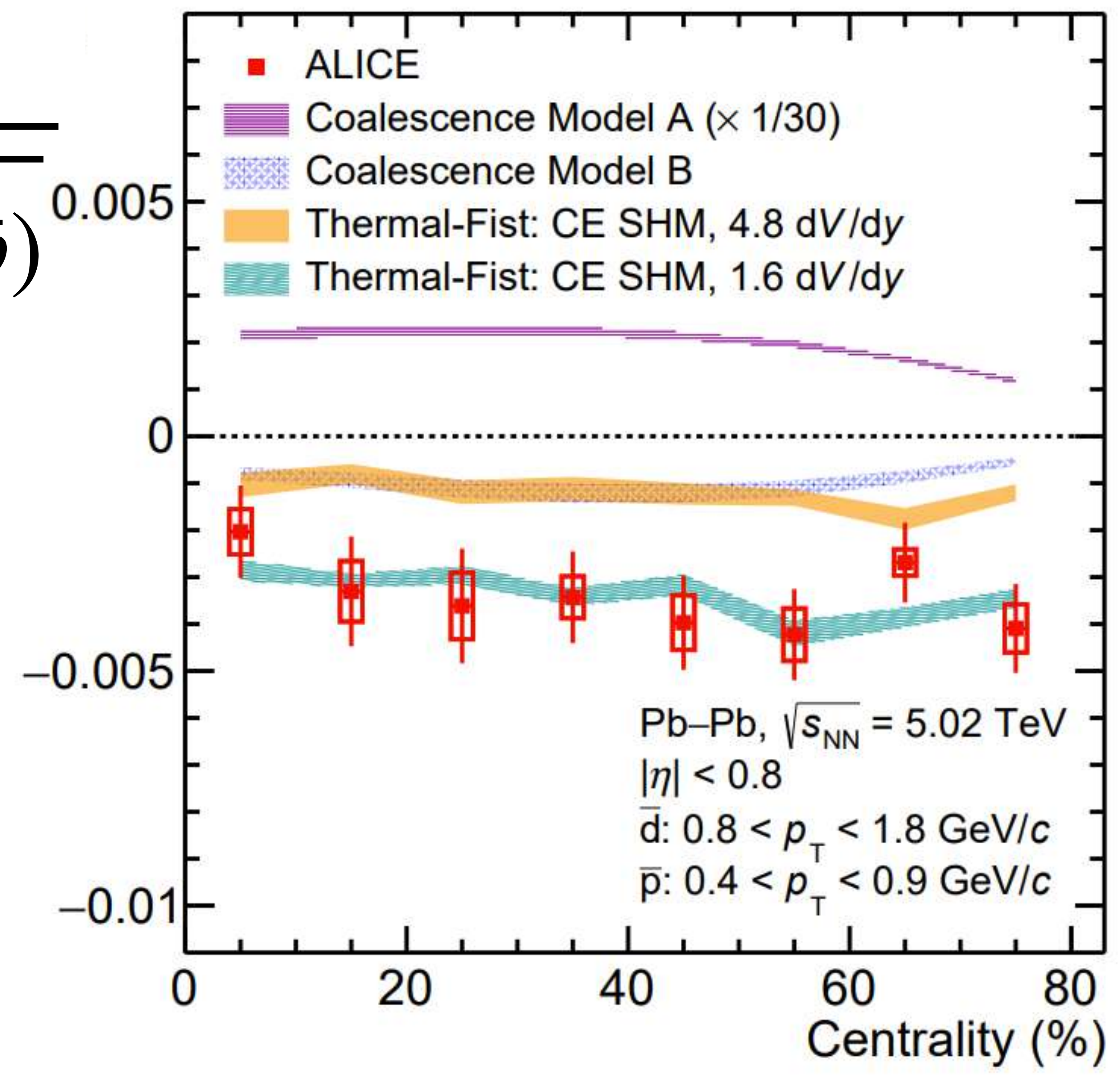
$$\frac{\kappa_2(p-\bar{p})}{\langle p+\bar{p}\rangle}$$



No annihilation

ALICE Coll., arXiv:2204.10166

$$\frac{\mathrm{cov}(\bar{\mathrm{d}},\bar{\mathrm{p}})}{\sqrt{\kappa_2(\bar{d})\kappa_2(\bar{p})}}$$



No annihilation

"wants" **long** range charge correlation

"wants" **short** range charge correlations

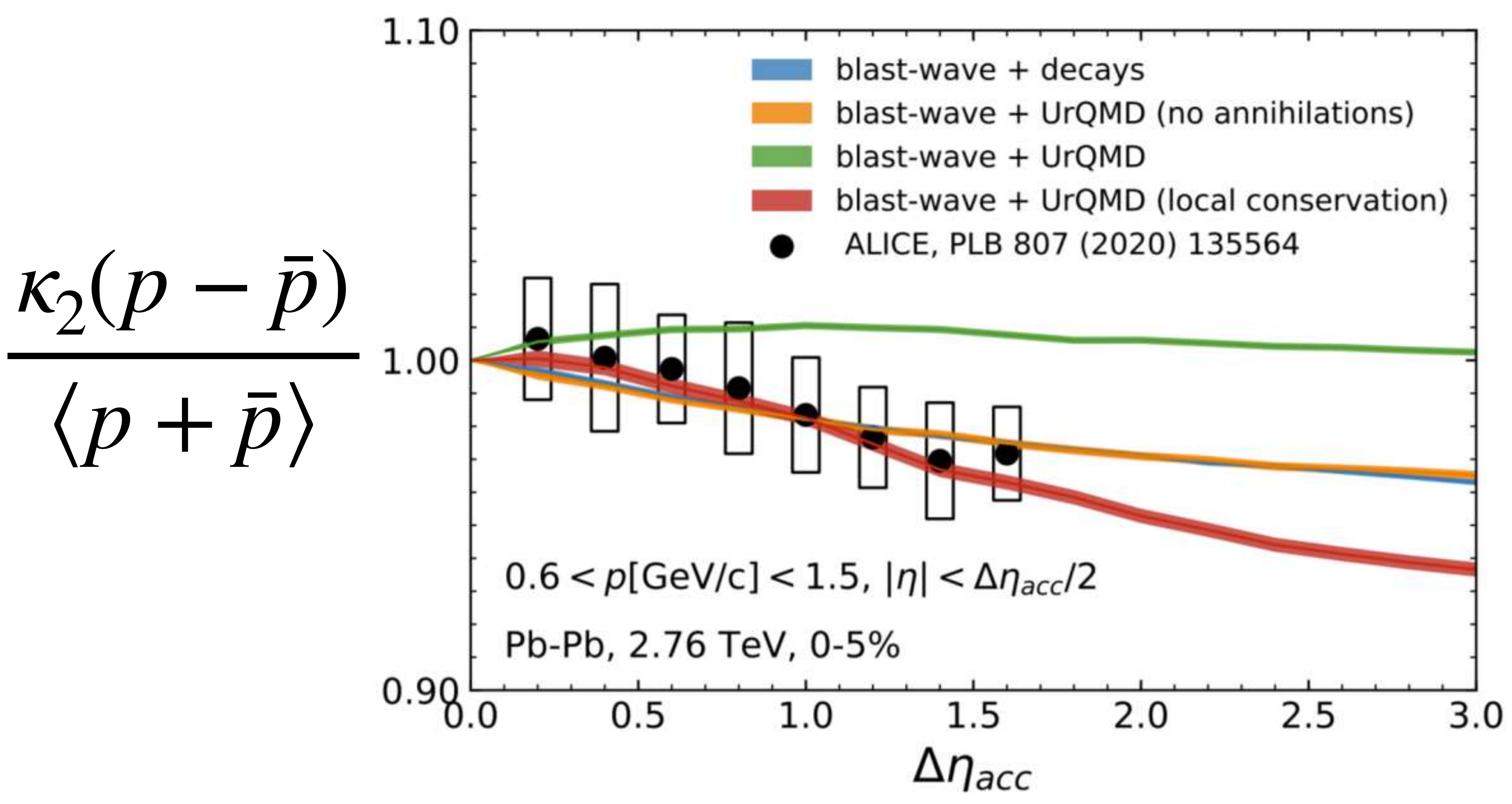
May resolve the tension between proton fluctuations that seem to prefer "global" baryon conservation vs light $\bar{d}$ - $\bar{p}$ correlations that prefer more "local" baryon conservation

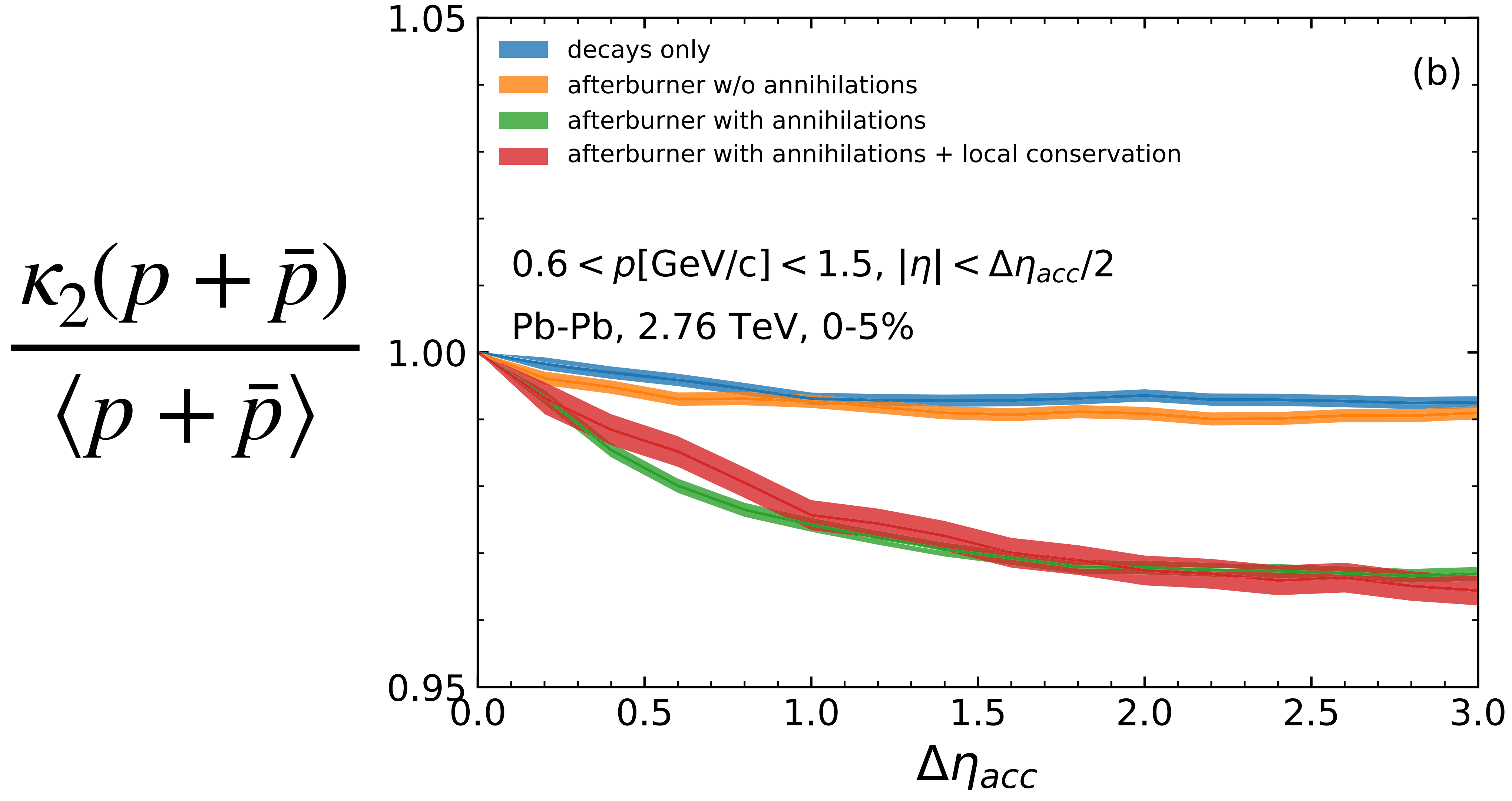
Baryon annihilation and fluctuations

Savchuk et al., PLB 827, 136983 (2022)

- $\kappa_2(p-\bar{p})$:
  - **Not (really)** affected by annihilation
  - affected by baryon number conservation
- $\kappa_2(p+\bar{p})$:
  - affected by annihilation
  - **NOT** affected by baryon number conservation

N.B.:
In UrQMD annihilation has NO detailed balance
$\rightarrow$*No reaction* $5\pi \rightarrow p+\bar{p}$
$\rightarrow$ *maximum effect*





Measure $\kappa_2(p-\bar{p})$ AND $\kappa_2(p+\bar{p})$ to constrain both amount of annihilation AND baryon correlation length

Summary

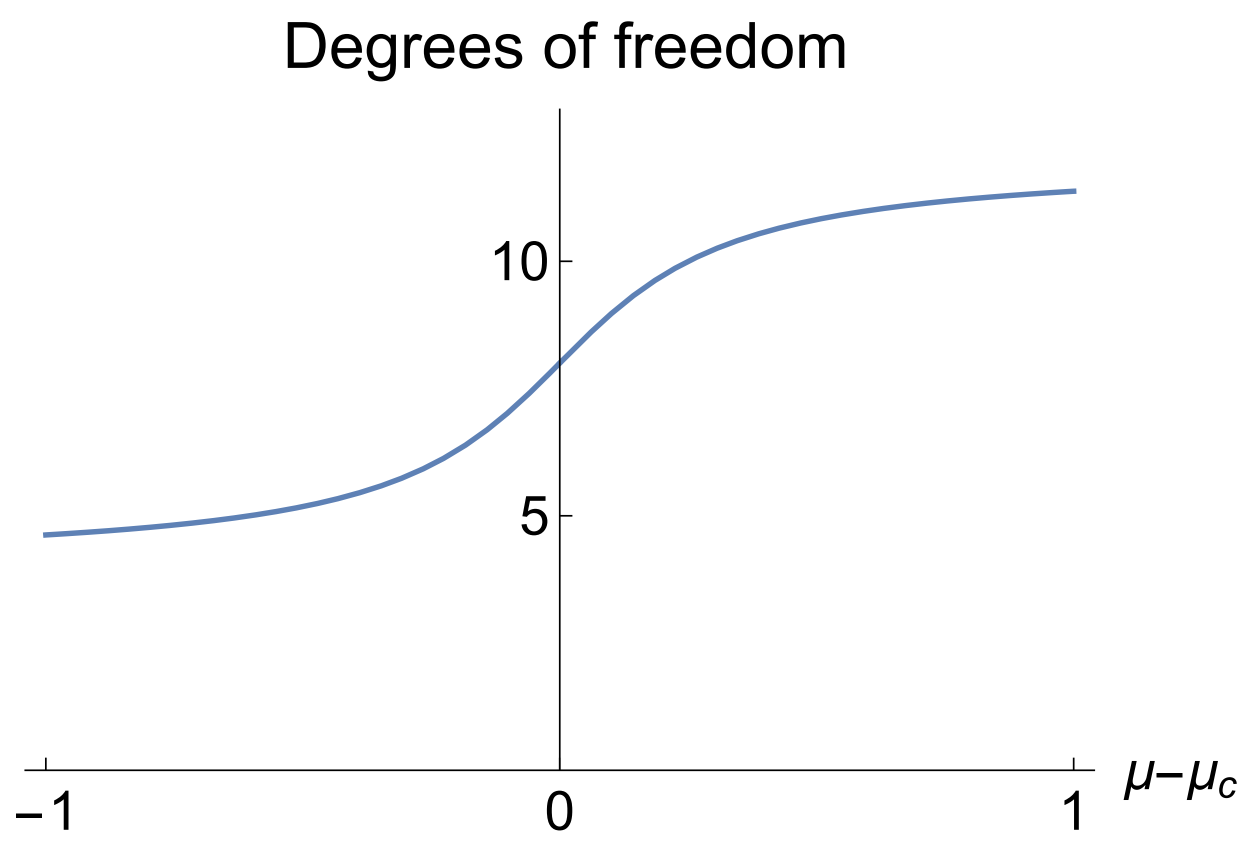
- Fluctuations measure derivatives of the Free Energy
  - They are a powerful tool to explore QCD phase diagram and other stuff
    - critical point
    - nuclear liquid gas transition
    - remnants of chiral criticality at $\mu \sim 0$
- Quantitative interpretation of measurements requires care:
  - Global (local) charge conservation
  - Protons vs baryons
  - momentum vs coordinate space
- Fluctuations also useful for other stuff: For example constrain baryon annihilation

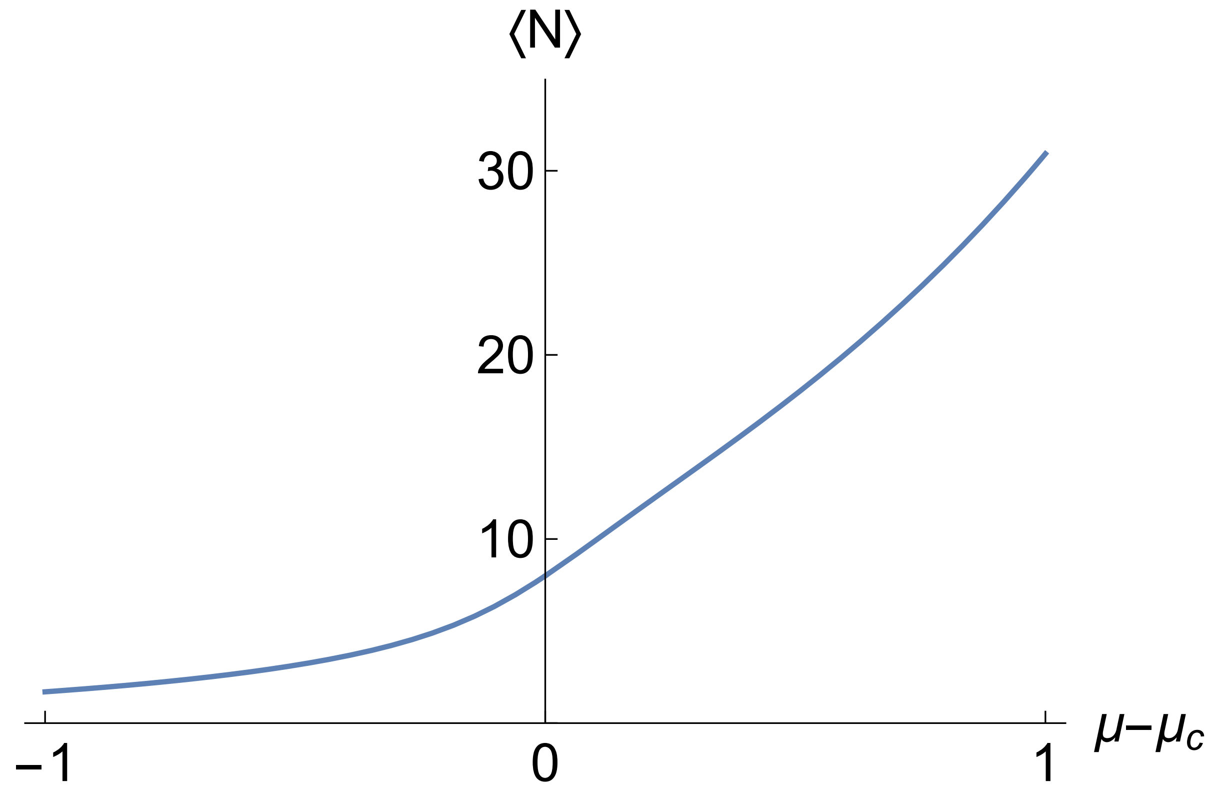
Thank You

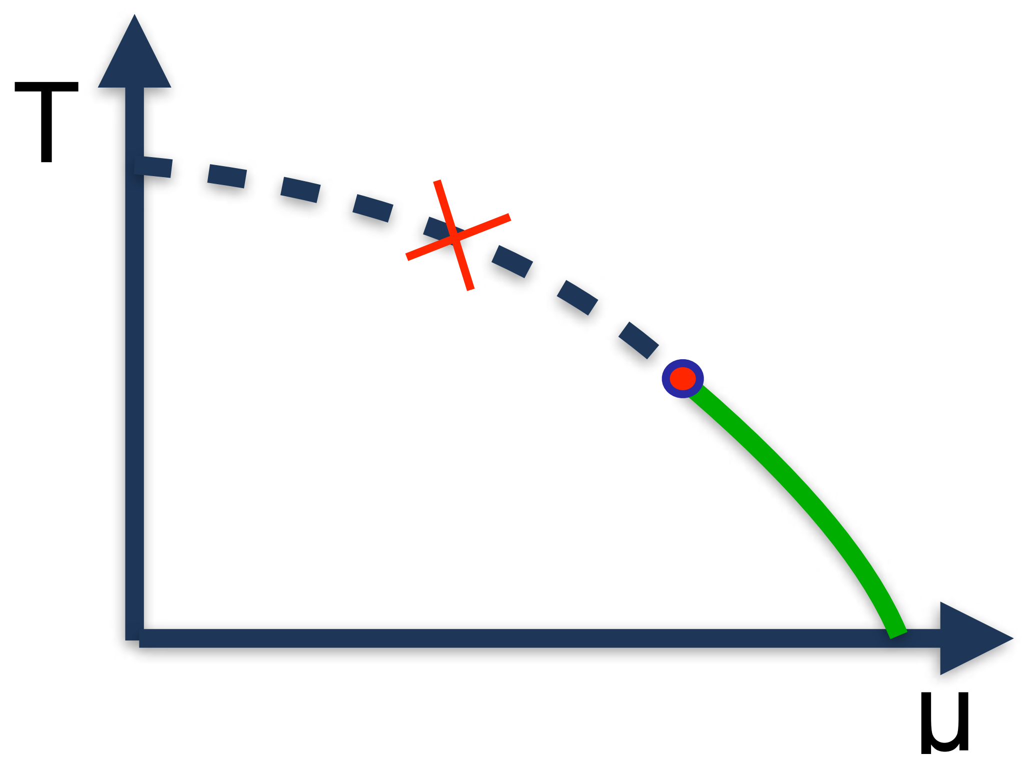
Simple model

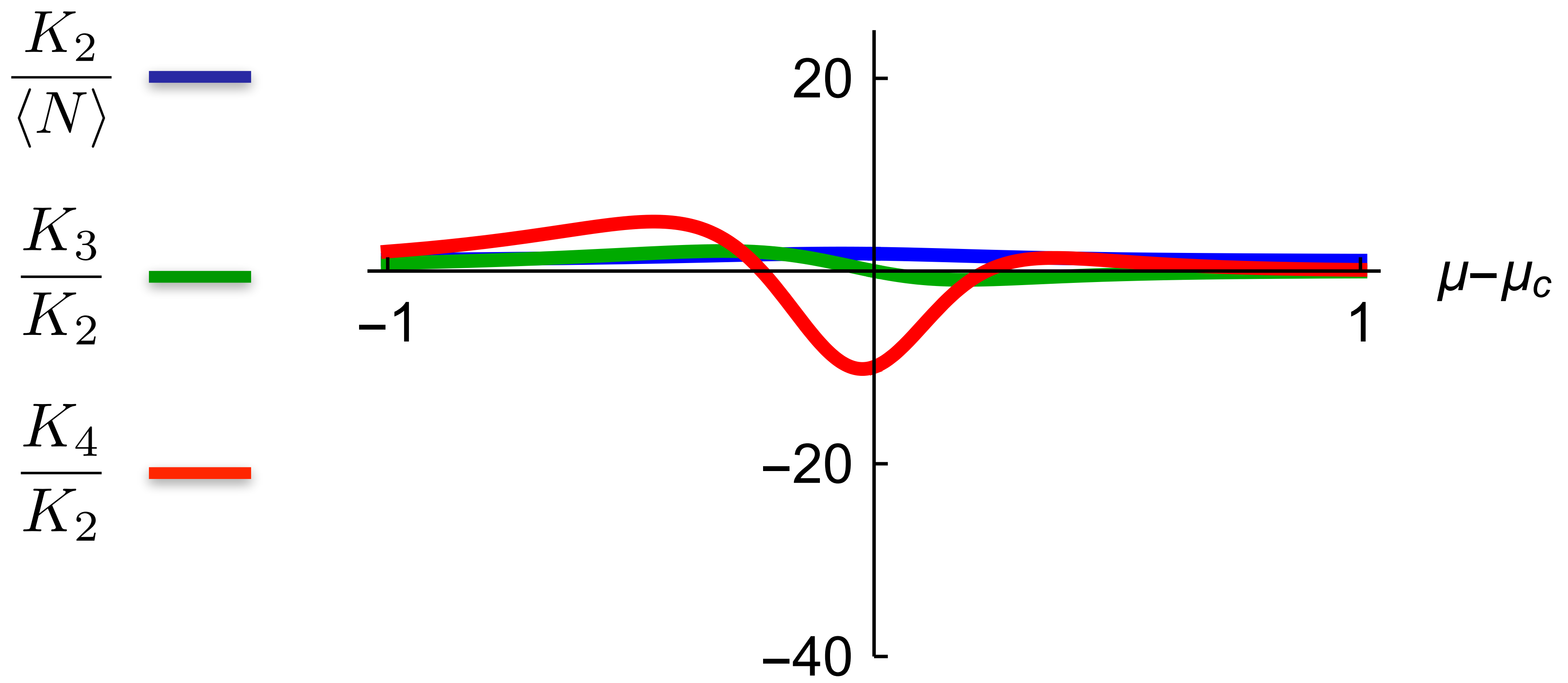
Change degrees of freedom at phase transition

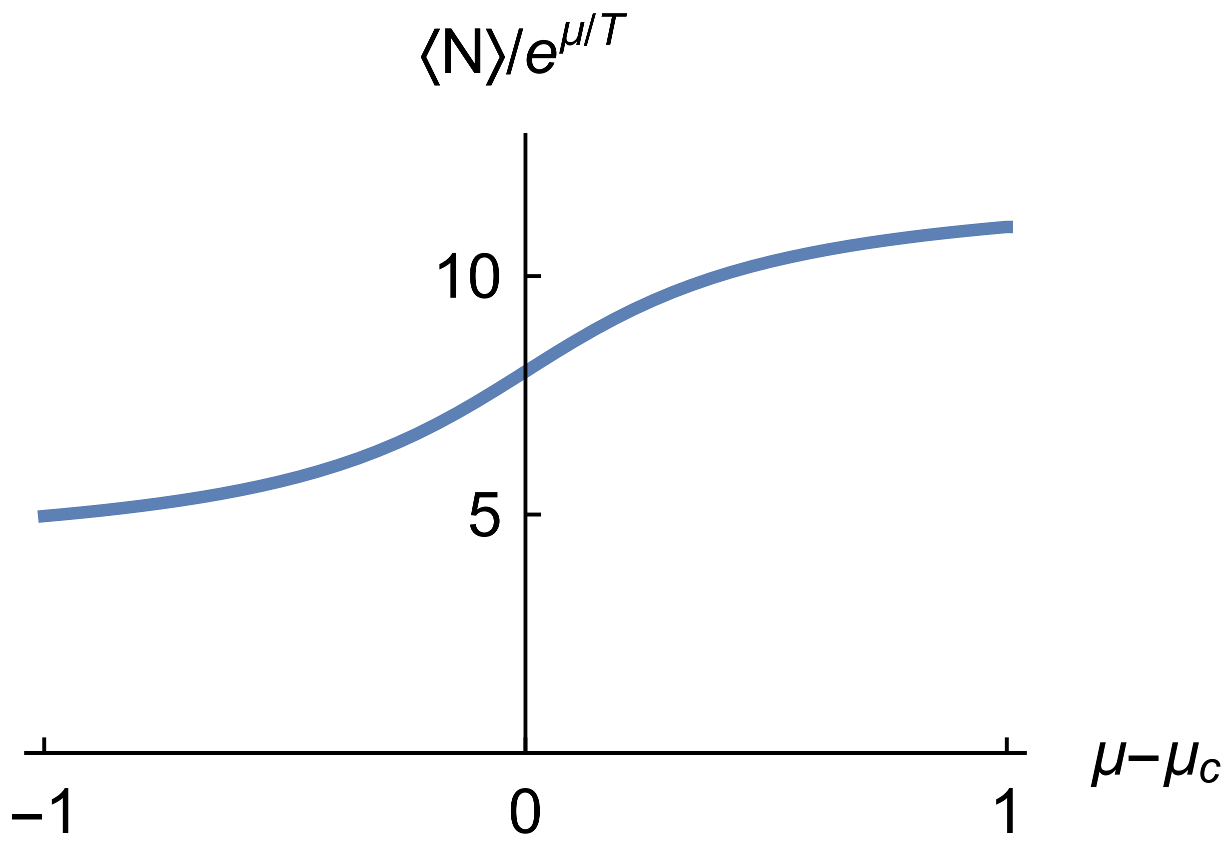
$$\langle N \rangle = dof(\mu)\, e^{\mu/T} \int d^3p e^{-E/T}$$

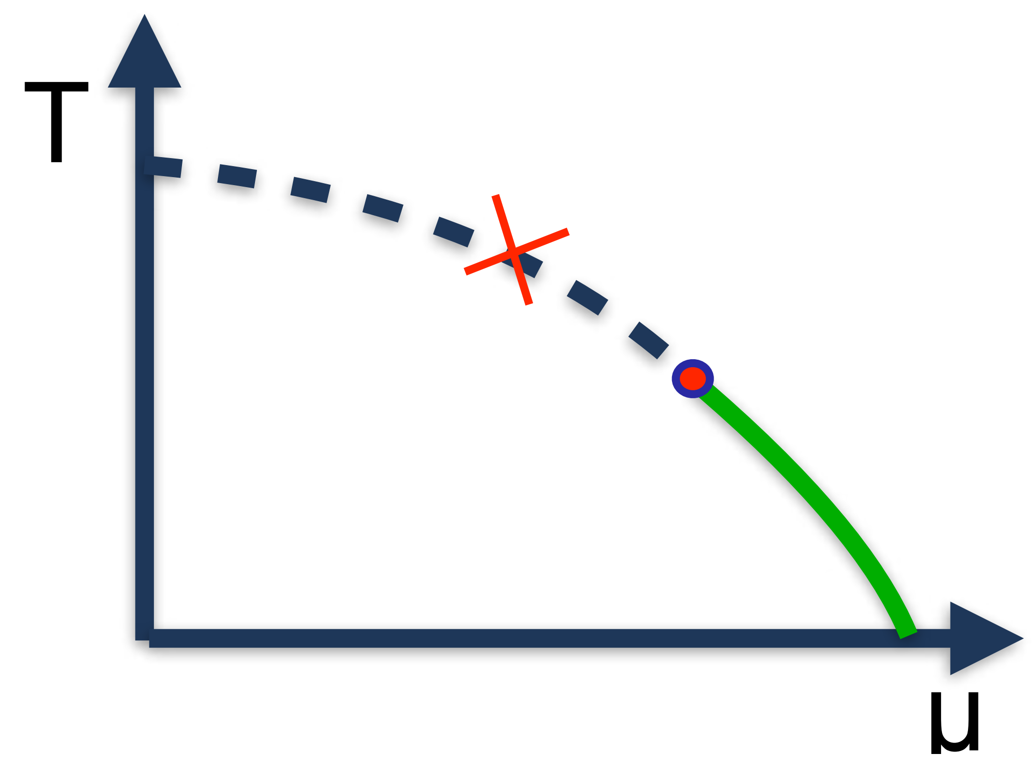


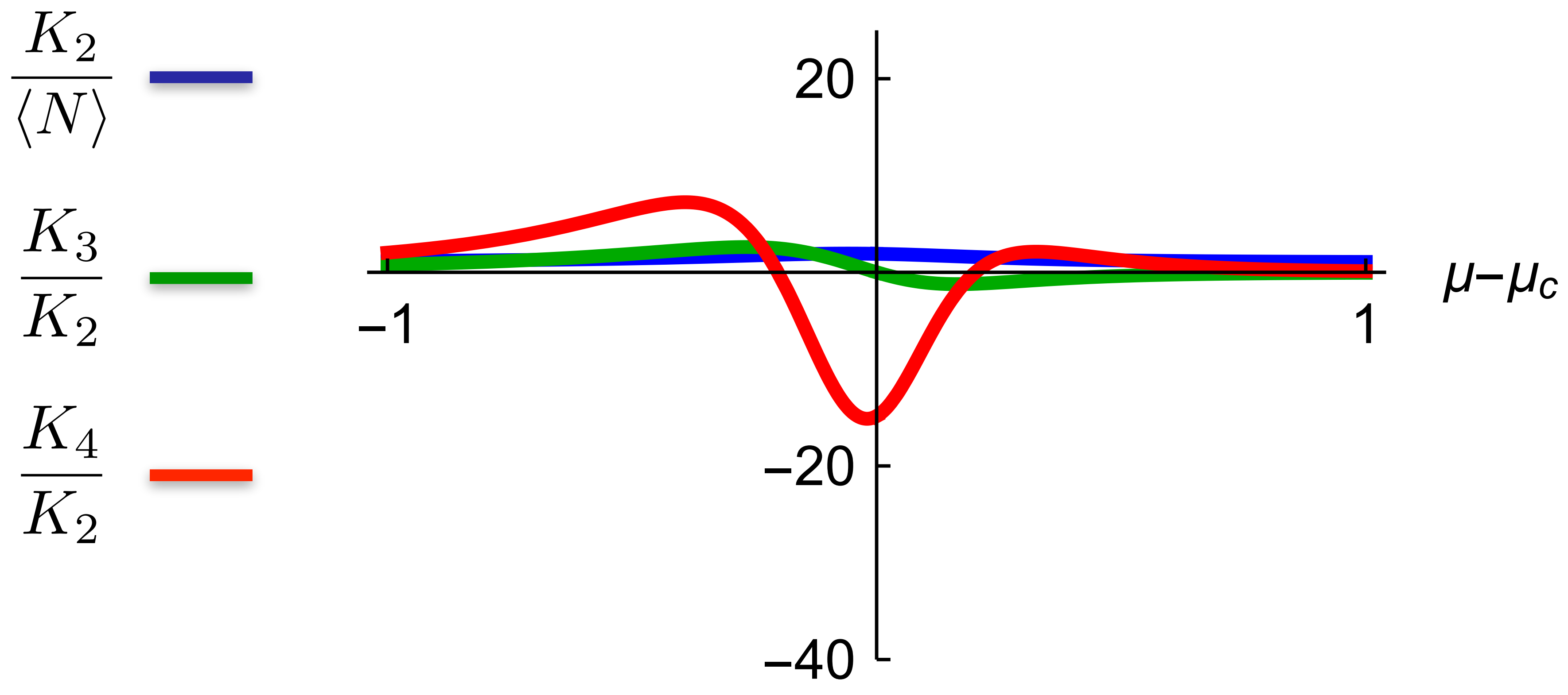


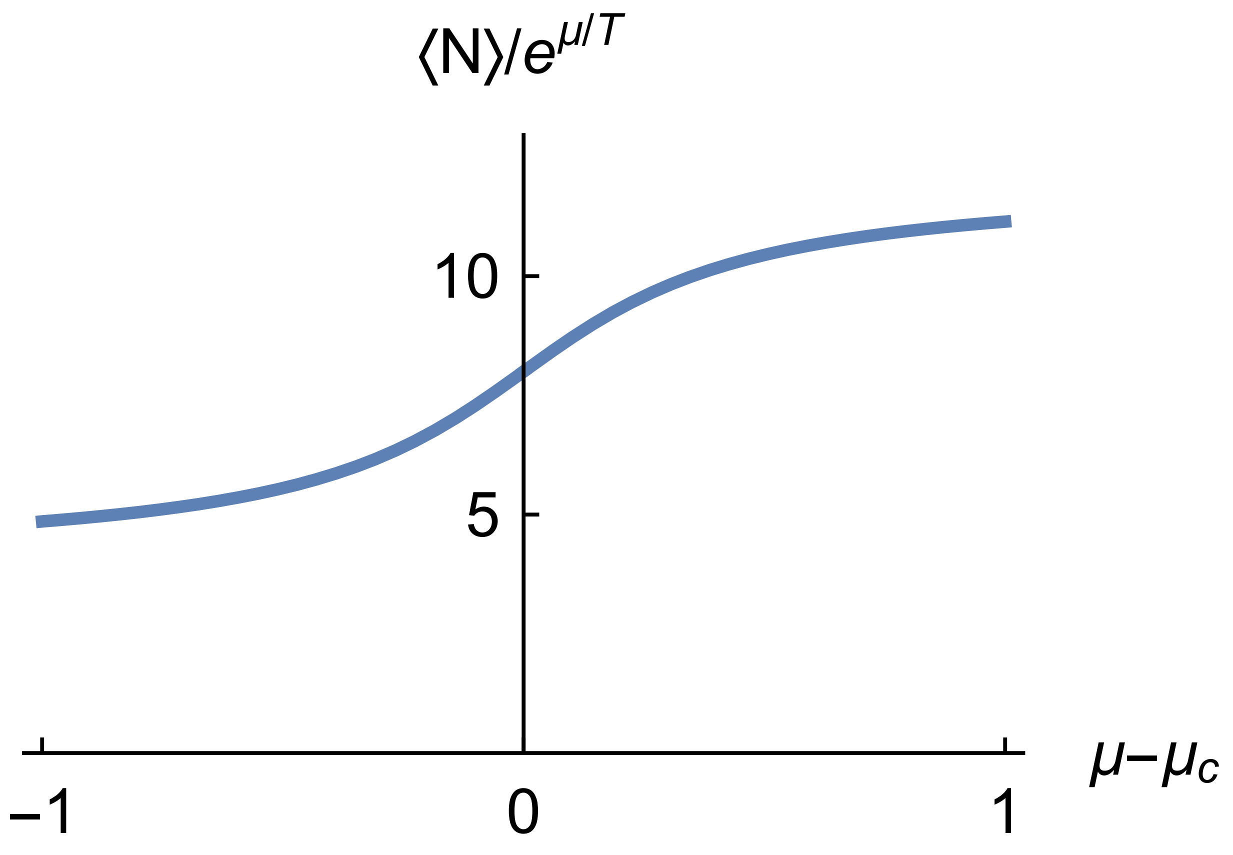

T
μ


$\frac{K_2}{\langle N \rangle}$
$\frac{K_3}{K_2}$
$\frac{K_4}{K_2}$
20
−20
−40
−1
1
μ−μc


⟨N⟩/e^{μ/T}
10
5
−1
0
1
μ−μc


T
μ


$\frac{K_2}{\langle N \rangle}$
$\frac{K_3}{K_2}$
$\frac{K_4}{K_2}$
20
−20
−40
−1
1
μ−μc


⟨N⟩/e^{μ/T}
10
5
−1
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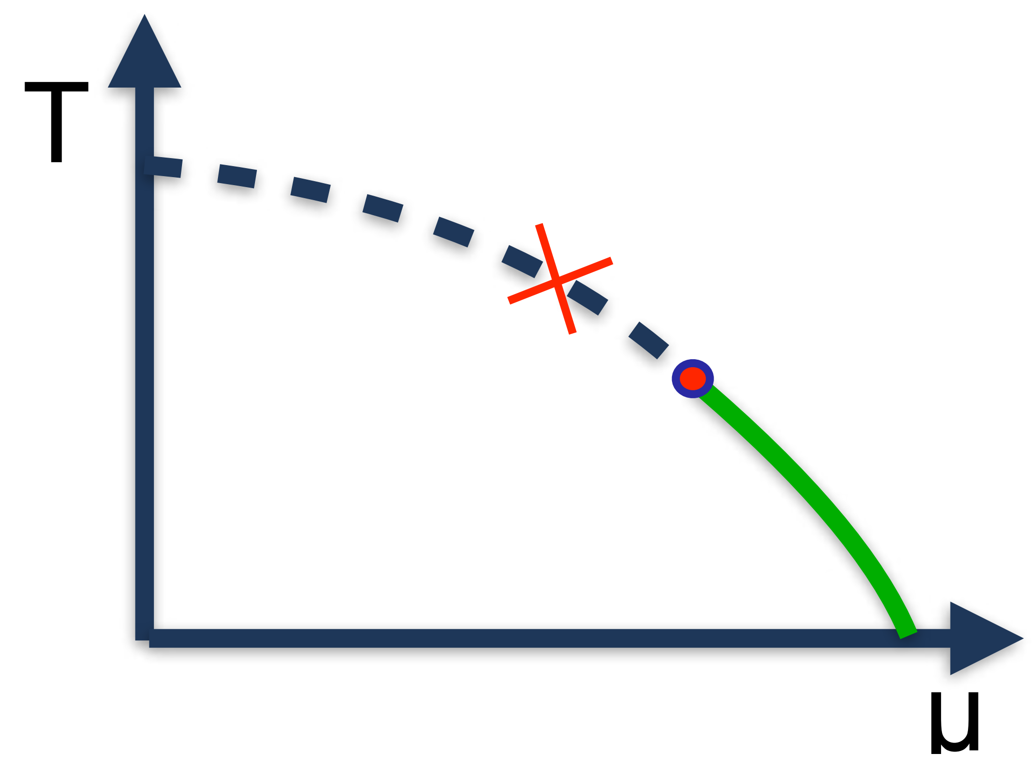
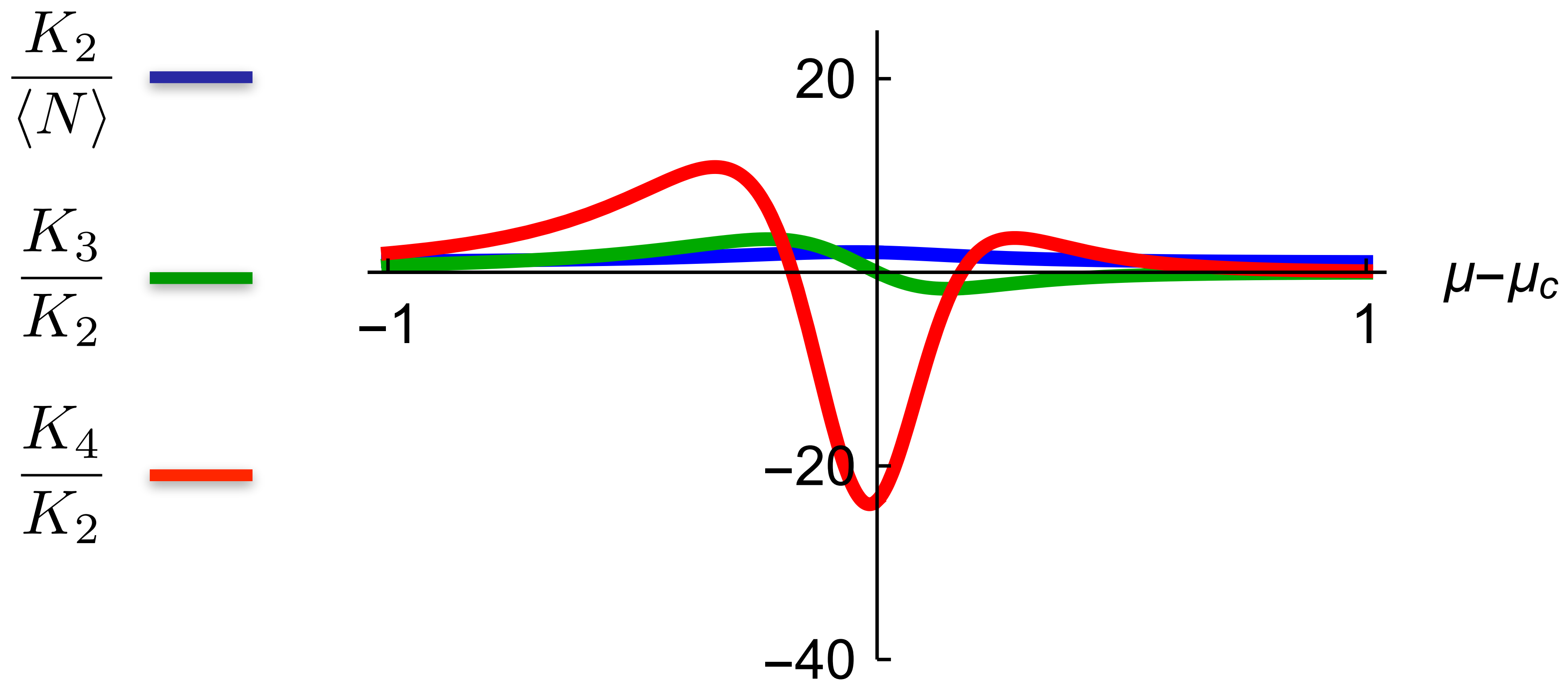
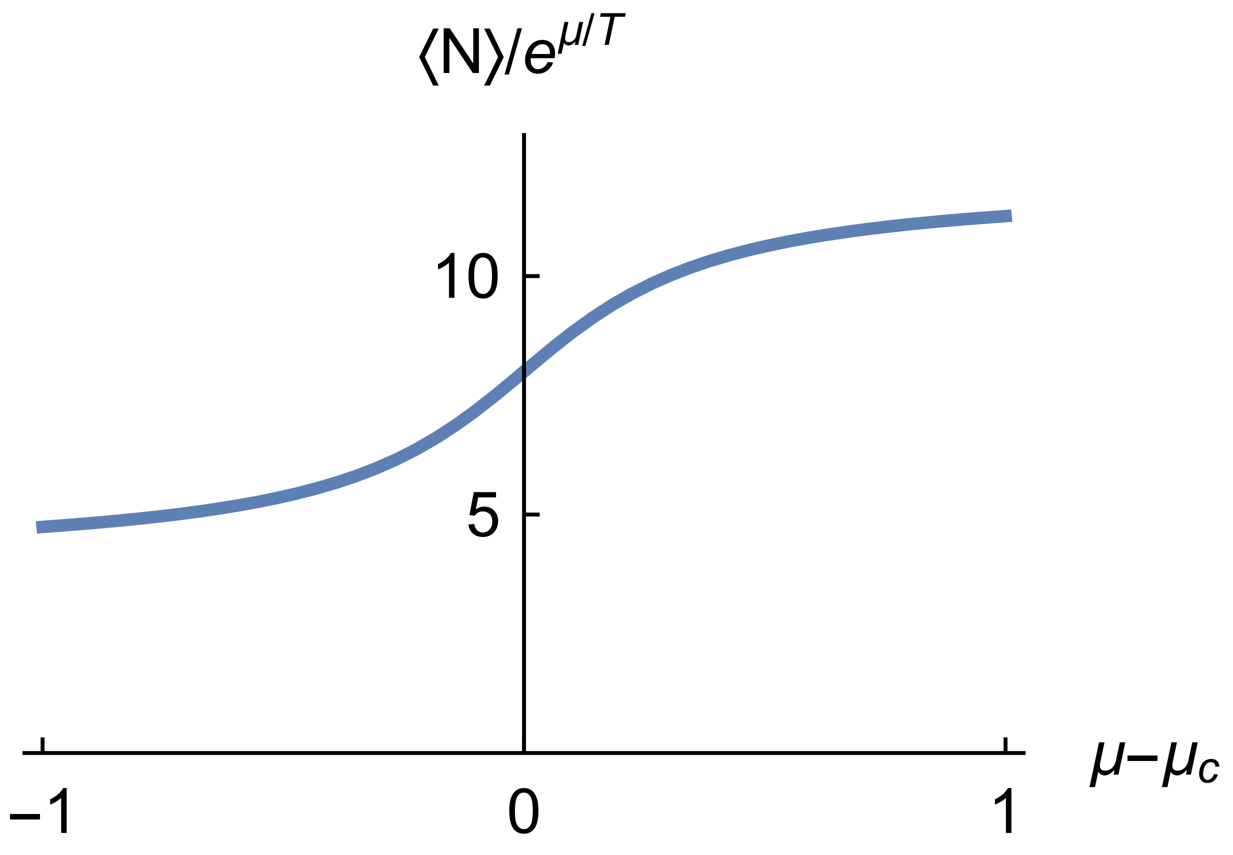
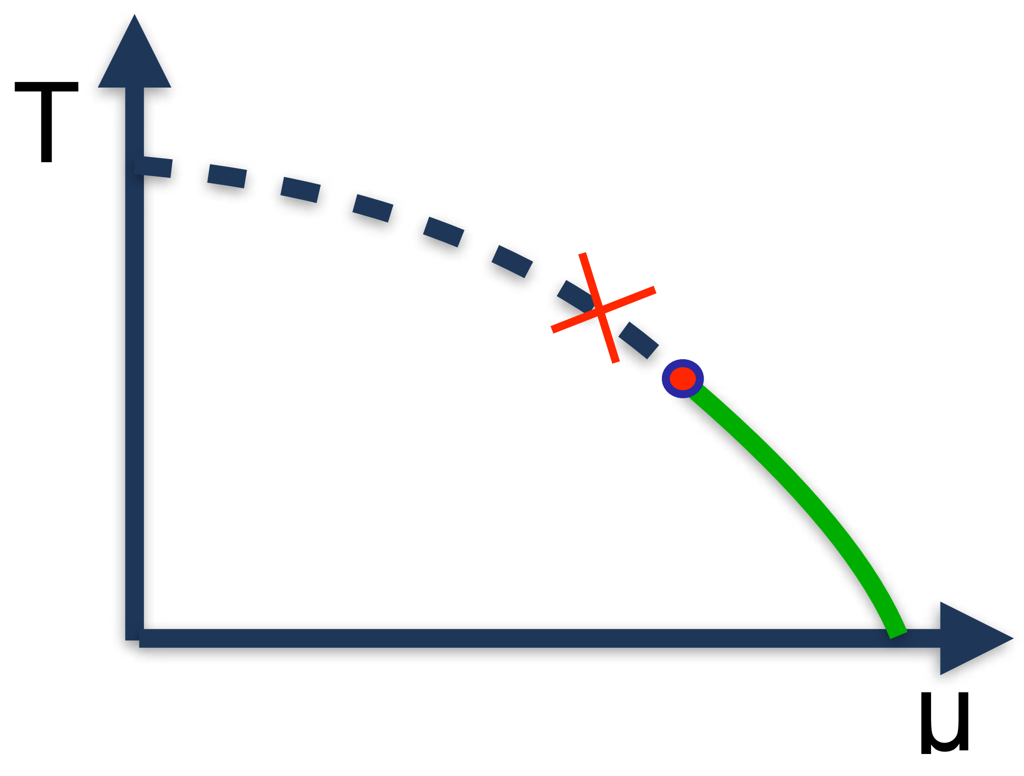
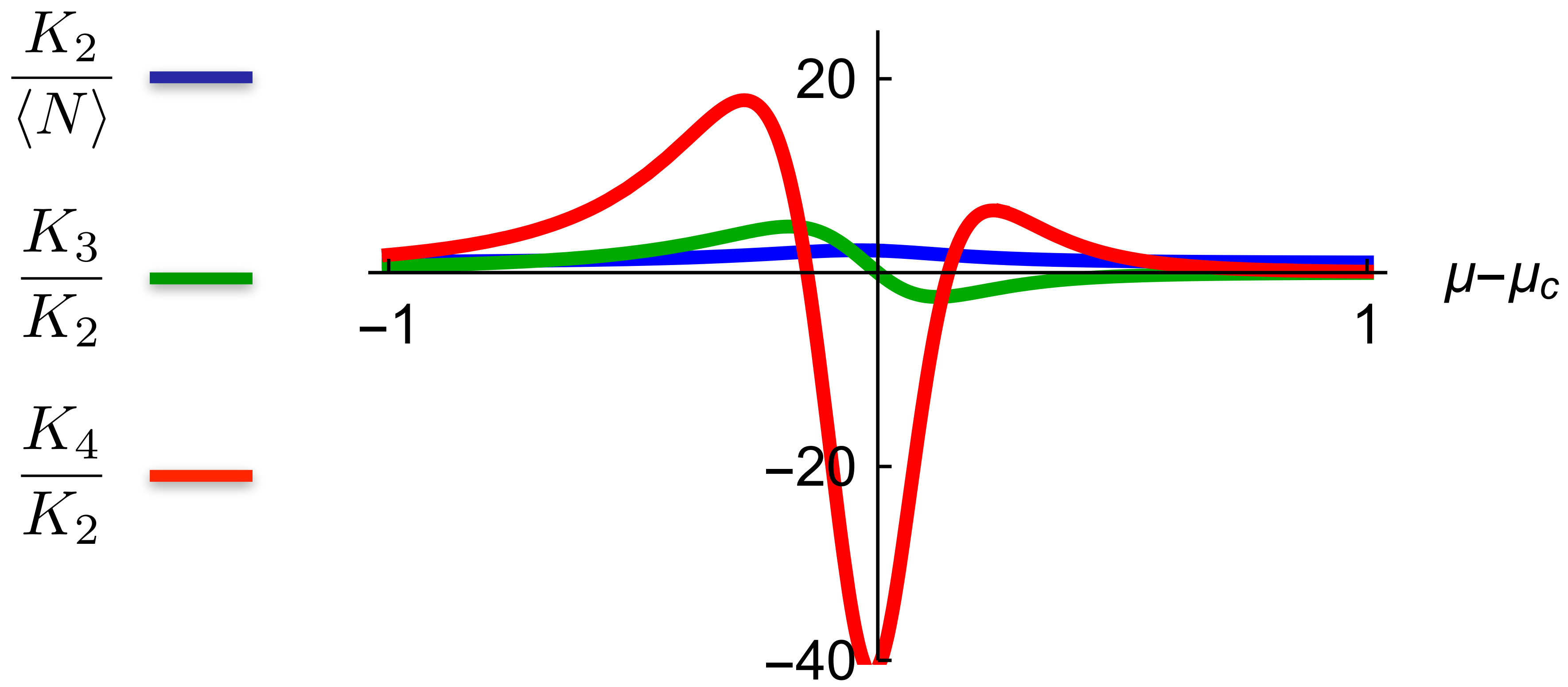
μ−μc

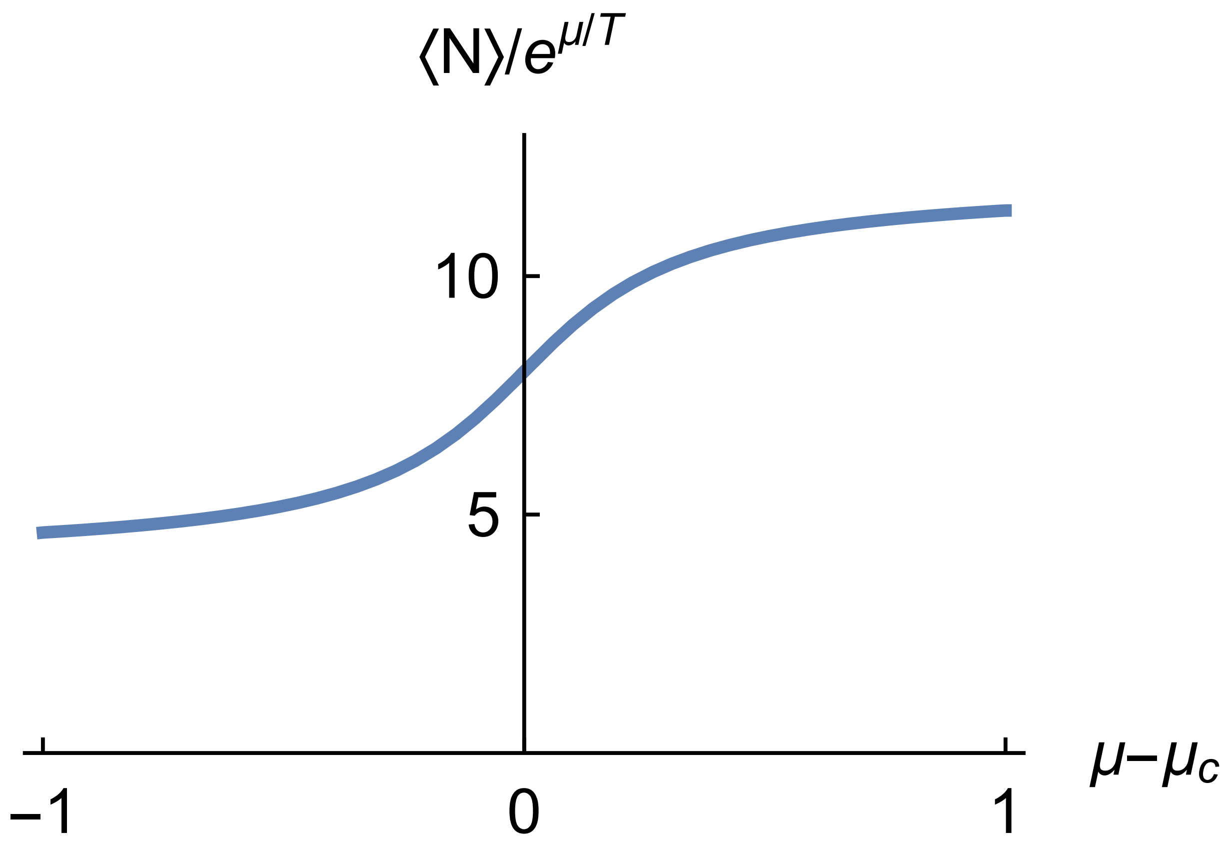

T
μ


$\frac{K_2}{\langle N \rangle}$
$\frac{K_3}{K_2}$
$\frac{K_4}{K_2}$
20
−20
−40
−1
1
μ−μc

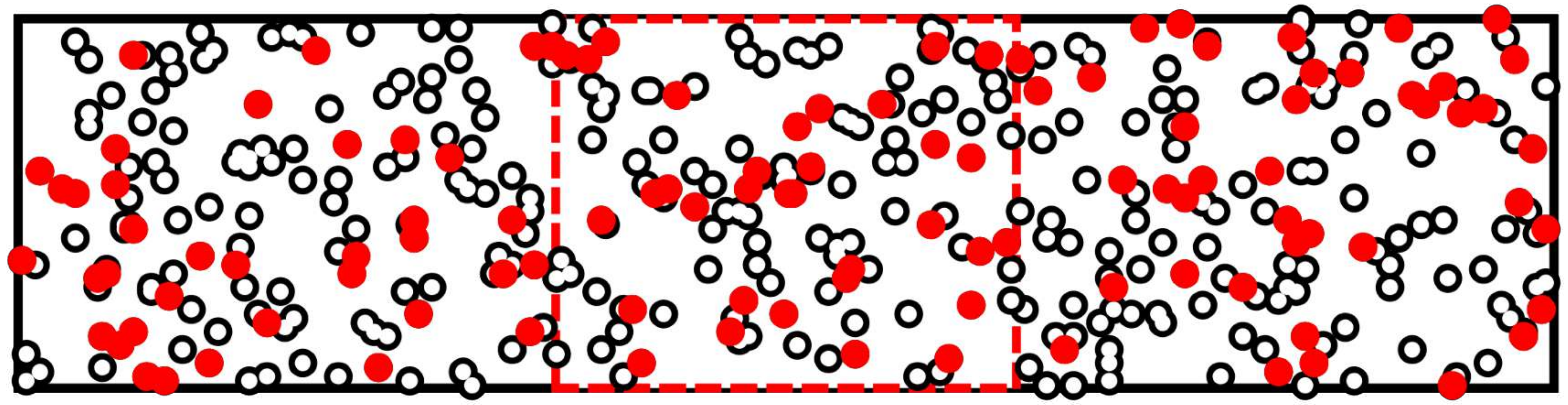

⟨N⟩/e^{μ/T}
10
5
−1
0
1
μ−μc

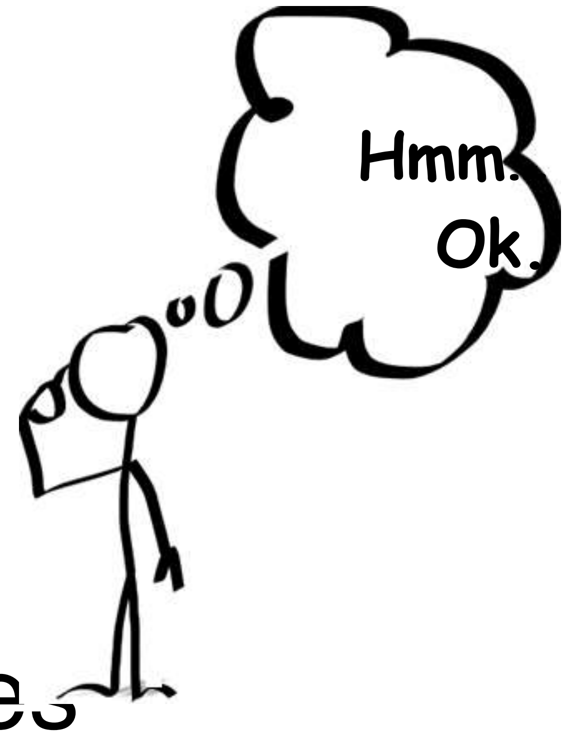

T
μ


$\frac{K_2}{\langle N \rangle}$
$\frac{K_3}{K_2}$
$\frac{K_4}{K_2}$
20
−20
−40
−1
1
μ−μc


⟨N⟩/e^μ/T
10
5
−1
0
1
μ−μc

Binomial acceptance vs actual acceptance

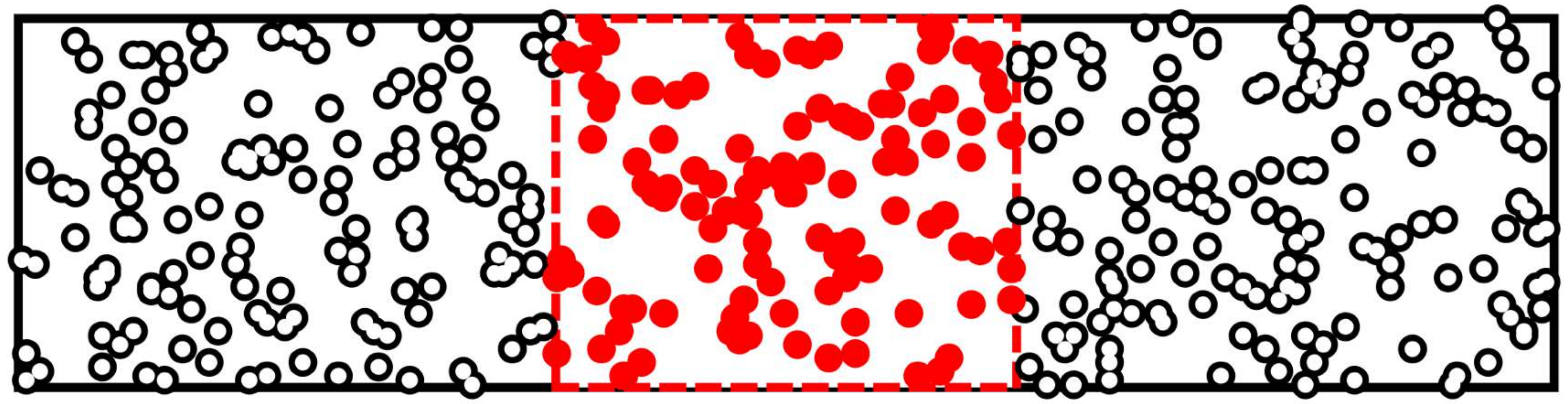




Binomial acceptance: accept each particle (charge) with probability α independently from all other particles

The binomial acceptance will not provide the correct result (except for a gas of uncorrelated particles)

What we really need is



Cumulants of (baryon) number distribution

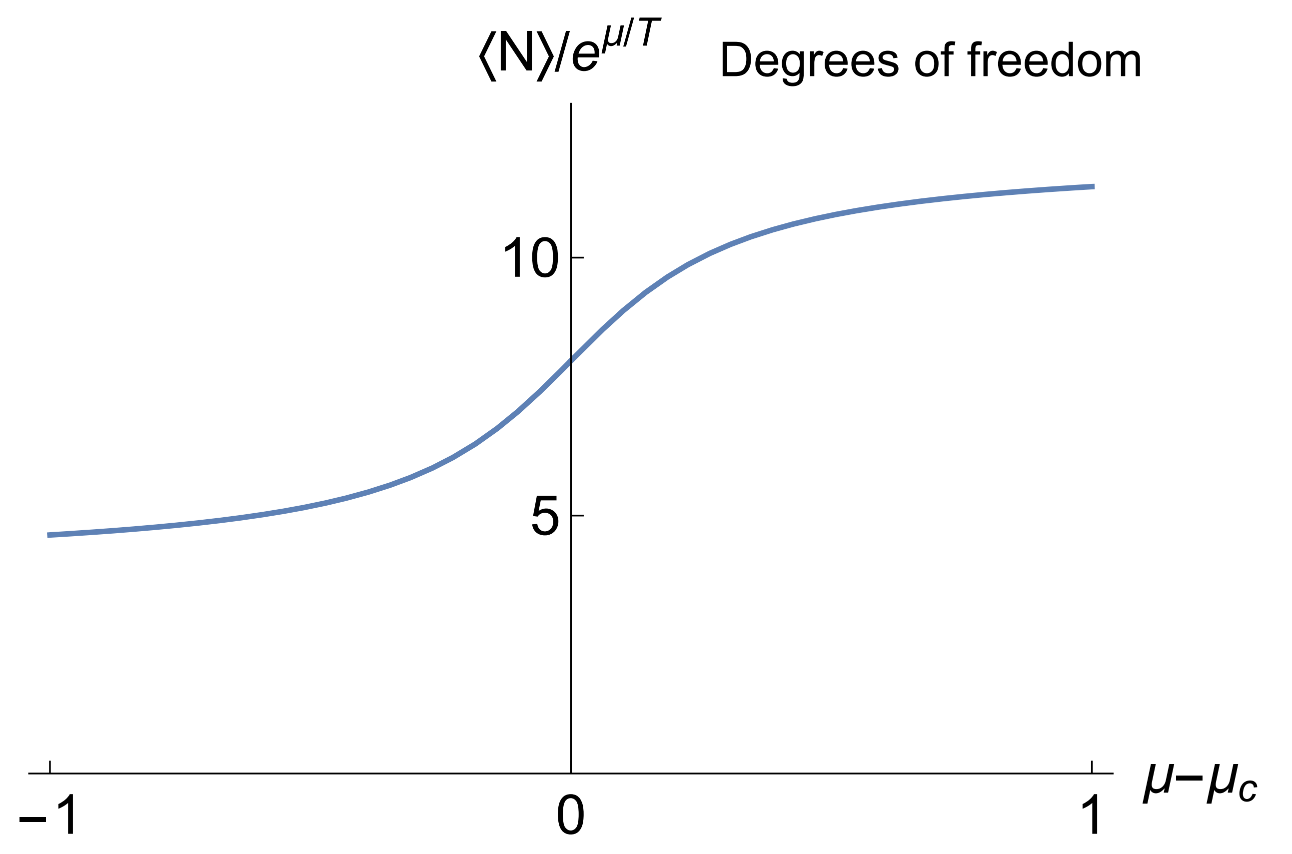
$$K_n = \frac{\partial^n}{\partial(\mu/T)^n} \ln Z = \frac{\partial^{n-1}}{\partial(\mu/T)^{n-1}} \langle N \rangle$$

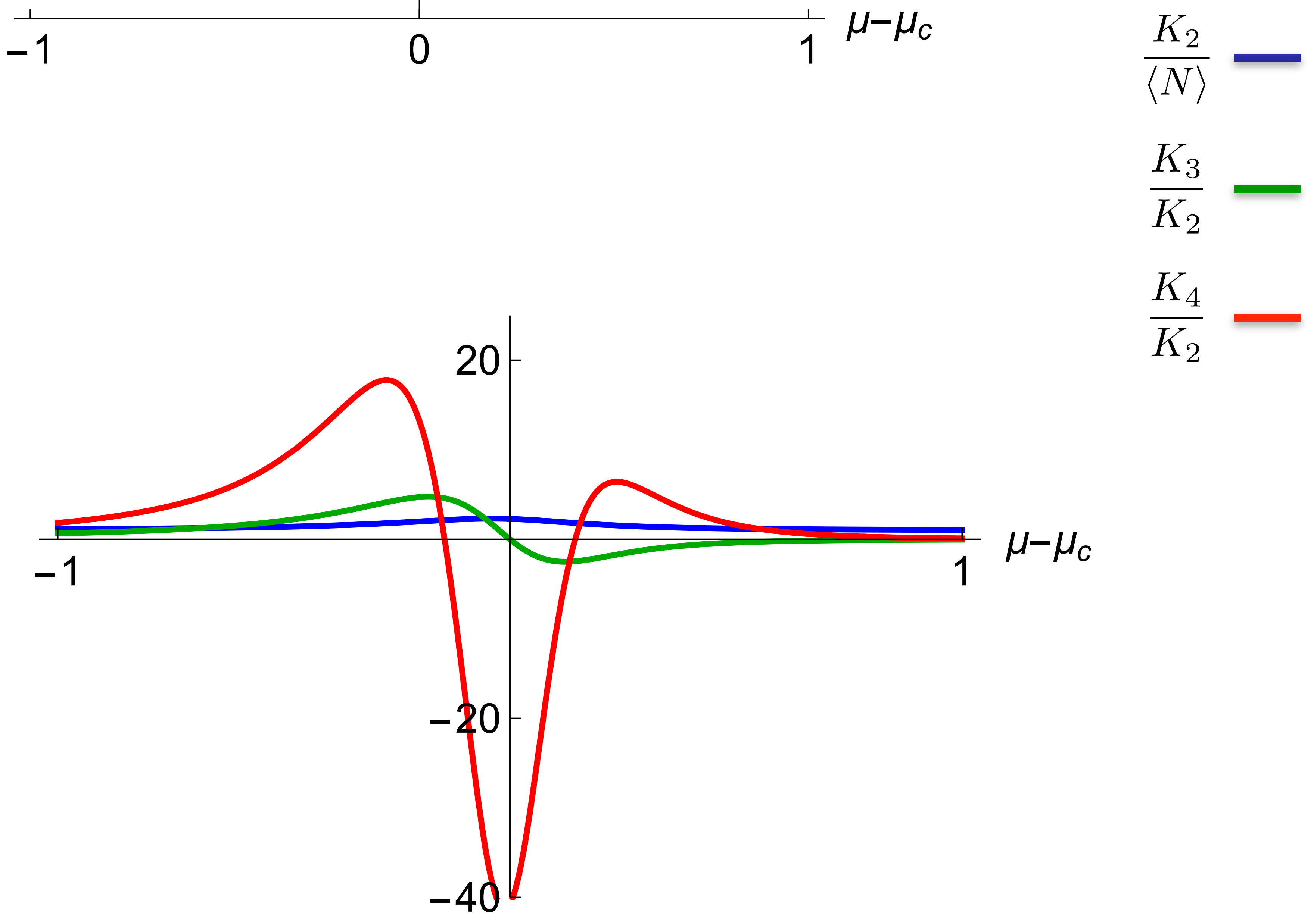
$$K_1 = \langle N \rangle\,, \quad K_2 = \langle N - \langle N \rangle \rangle^2\,, \quad K_3 = \langle N - \langle N \rangle \rangle^3$$

Cumulants scale with volume (extensive): $K_n \sim V$

Volume not well controlled in heavy ion collisions

Cumulant Ratios: $$\frac{K_2}{\langle N \rangle},\ \frac{K_3}{K_2},\ \frac{K_4}{K_2}$$



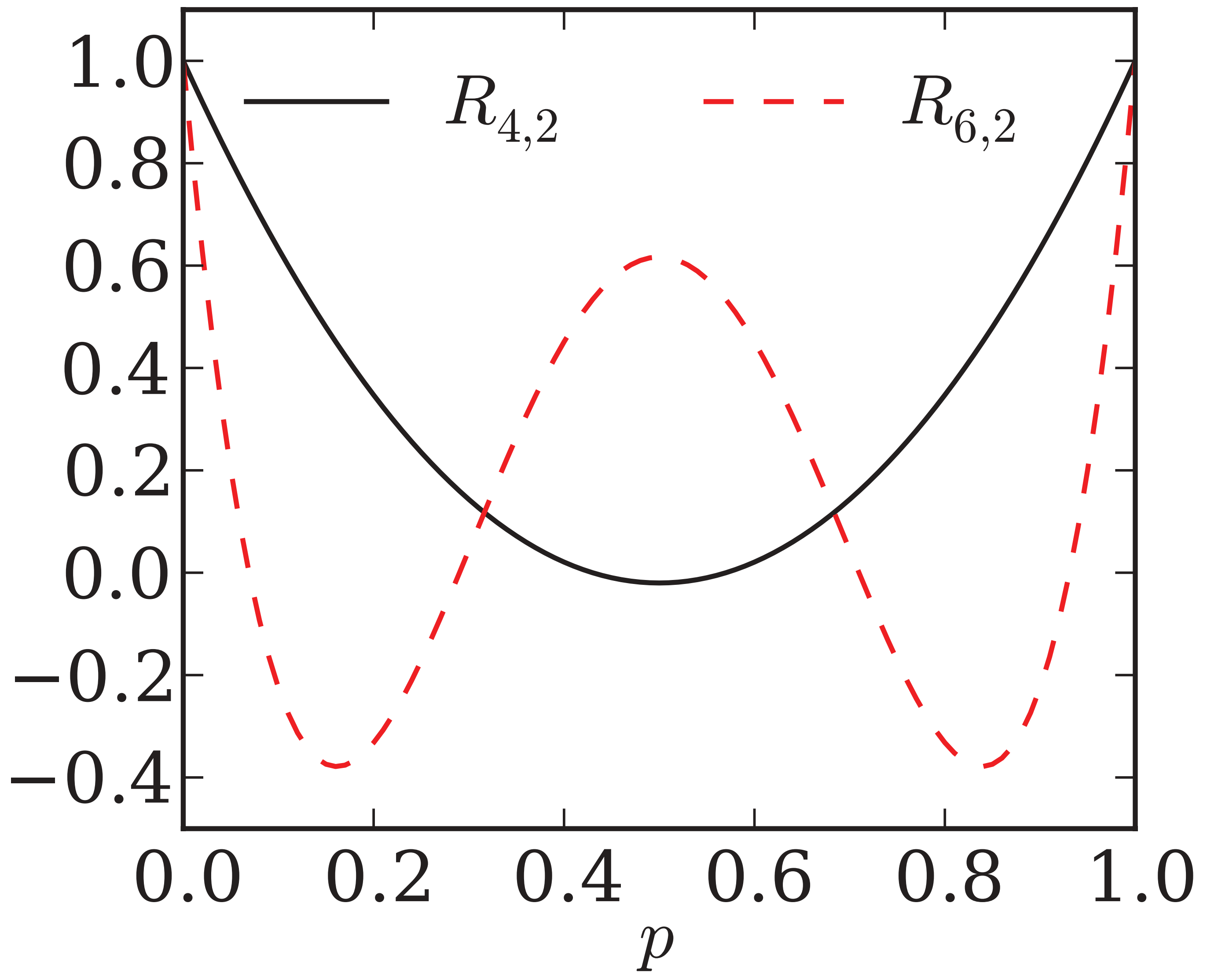


Baryon number cumulants measure derivatives of the EOS w.r.t chemical potential

Baryon number conservation and lattice susceptibilities

- Charges (baryon number, strangeness, electric charge) are conserved globally in HI collisions
- Lattice (and most other calculations) work in the grand canonical ensemble: charges may fluctuate
- Effect of charge conservation have been calculated in the **ideal gas/HRG** limit. NON-neglibile corrections especially for higher order cumulants
  (Bzdak et al 2013, Rustamov et al. 2017,…)
- Wouldn't it be nice to know what the effect of charge conservation on **real QCD** (aka lattice) susceptibilities is?

This can actually be done!



Bzdak et al, 2013

V. Vovchenko, O. Savchuk, R. Poberezhnyuk, M. Gorenstein, V.K., arXiv 2003.13905,

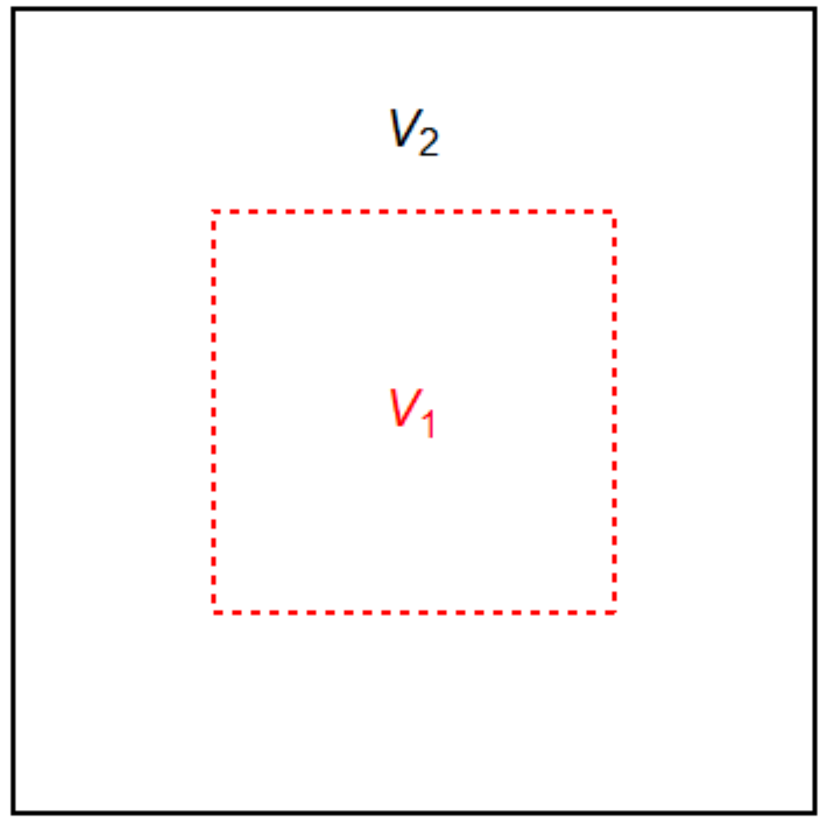
V. Vovchenko, R. Poberezhnyuk, V.K., arXiv:2007.03850

Subensemble acceptance method
(SAM)

Partition a thermal system with a globally conserved charge *B (canonical ensemble)* into two subsystems which can exchange the charge

$$V = V_1 + V_2$$

Assume thermodynamic limit:



$$V, V_1, V_2 \to \infty; \quad \frac{V_1}{V} = \alpha = const; \quad \frac{V_2}{V} = (1-\alpha) = const;$$

$$V_1, V_2 \gg \xi^3 \qquad \xi = correlation\, length$$

The canonical partition function then reads:

$$Z^{ce}(T,V,B) = \sum_{B_1} Z^{ce}(T,V_1,B_1) Z^{ce}(T,V-V_1,B-B_1)$$

The probability to have charge B_1 *in* V_1 is:

$$P(B_1) \sim Z^{ce}(T,\alpha V,B_1) Z^{ce}(T,(1-\alpha)V,B-B_1), \qquad \alpha \equiv V_1/V$$

Subensemble acceptance method
(SAM)

In the thermodynamic limit, $V \to \infty$, Z^{ce} expressed through free energy *density*

$$Z^{\mathrm{ce}}(T,V,B) \overset{V\to\infty}{\simeq} \exp\left[-\frac{V}{T} f(T,\rho_B)\right]$$

Cumulant generating function for B_1:

$$G_{B_1}(t) \equiv \ln\langle e^{t\,B_1}\rangle = \ln\left\{\sum_{B_1} e^{t\,B_1} \exp\left[-\frac{\alpha V}{T} f(T,\rho_{B_1})\right] \exp\left[-\frac{\beta V}{T} f(T,\rho_{B_2})\right]\right\} + \tilde{C}$$

Cumulants of B_1:

$$\kappa_n[B_1] = \left.\frac{\partial^n G_{B_1}(t)}{\partial t^n}\right|_{t=0} \equiv \tilde{\kappa}_n[B_1(t)]|_{t=0}$$

or

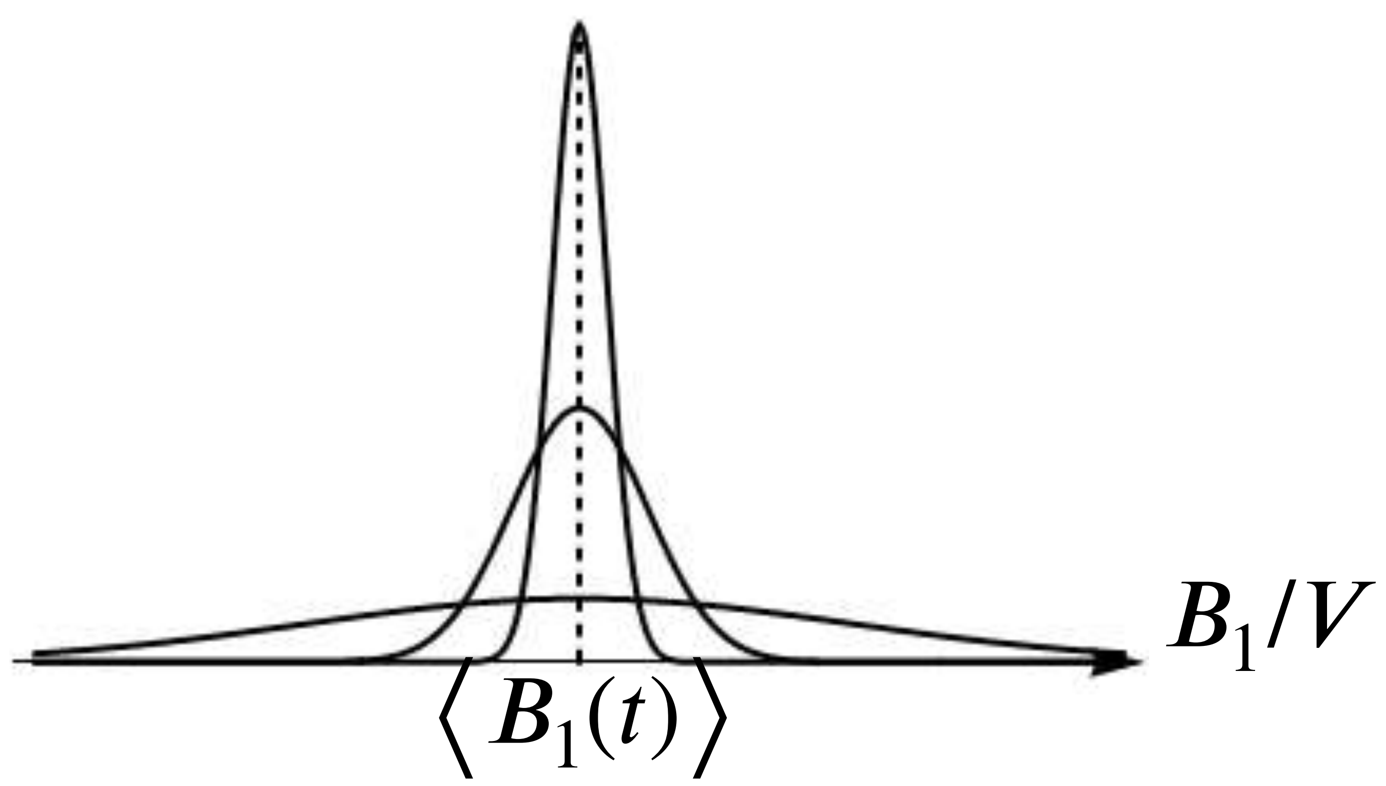
$$\kappa_n[B_1] = \left.\frac{\partial^{n-1}\tilde{\kappa}_1[B_1(t)]}{\partial t^{n-1}}\right|_{t=0}$$

All κ_n can be calculated by determining the *t*-dependent first cumulant $\tilde{\kappa}_1[B_1(t)]$

Making the connection…

$$\tilde{\kappa}_1[B_1(t)] = \frac{\sum_{B_1} B_1 \tilde{P}(B_1;t)}{\sum_{B_1} \tilde{P}(B_1;t)} \equiv \langle B_1(t) \rangle \qquad \text{with} \qquad \tilde{P}(B_1;t) = \exp\left\{ tB_1 - V\,\frac{\alpha f(T,\rho_{B_1}) + \beta f(T,\rho_{B_2})}{T} \right\}.$$

Thermodynamic limit: $\tilde{P}(B_1;t)$ *highly peaked* at $\langle B_1(t) \rangle$



$\langle B_1(t) \rangle$ is a solution to equation $d\tilde{P}/\mathrm{d}B_1 = 0$:

$$t = \hat{\mu}_B[T,\rho_{B_1}(t)] - \hat{\mu}_B[T,\rho_{B_2}(t)] \qquad \text{with} \qquad \hat{\mu}_B \equiv \mu_B/T, \qquad \mu_B(T,\rho_B) = \partial f(T,\rho_B)/\partial\rho_B$$

t = 0:

$$\rho_{B_1} = \rho_{B_2} = B/V,\ B_1 = \alpha B,$$

i.e. conserved charge uniformly distributed between the two subsystems

Second order cumulant

Differentiate condition for maximum of $\widetilde{P}(B_1;t)$,

$$t = \hat{\mu}_B[T,\rho_{B_1}(t)] - \hat{\mu}_B[T,\rho_{B_2}(t)] \qquad (*)$$

$$\frac{\partial(*)}{\partial t}: \qquad 1 = \left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B1}}\right)_T \left(\frac{\partial \rho_{B1}}{\partial \langle B_1 \rangle}\right)_V \frac{\partial \langle B_1 \rangle}{\partial t} - \left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B2}}\right)_T \left(\frac{\partial \rho_{B2}}{\partial \langle B_2 \rangle}\right)_V \frac{\partial \langle B_2 \rangle}{\partial \langle B_1 \rangle} \frac{\partial \langle B_1 \rangle}{\partial t}$$

$$\left(\frac{\partial \hat{\mu}_B}{\partial \rho_{B1,2}}\right)_T \equiv \left[\chi_2^B\left(T,\rho_{B_{1,2}}\right)\, T^3\right]^{-1}, \quad \rho_{B_1} \equiv \frac{\langle B_1 \rangle}{\alpha V}, \quad \rho_{B_2} \equiv \frac{\langle B_2 \rangle}{(1-\alpha)V}, \quad \langle B_2 \rangle = B - \langle B_1 \rangle, \quad \frac{\partial \langle B_1 \rangle}{\partial t} \equiv \tilde{\kappa}_2[B_1(t)]$$

Solve the equation for $\tilde{\kappa}_2$:

$$\tilde{\kappa}_2[B_1(t)] = \frac{V\,T^3}{[\alpha\,\chi_2^B(T,\rho_{B_1})]^{-1} + [(1-\alpha)\,\chi_2^B(T,\rho_{B_2})]^{-1}}$$

t = 0: $$\kappa_2[B_1] = \alpha\,(1-\alpha)\,V\,T^3\,\chi_2^B$$

Higher-order cumulants: iteratively differentiate $\tilde{\kappa}_2$ w.r.t. t

Full result up to sixth order

$$\beta = 1 - \alpha$$

$$\kappa_1[B_1] = \alpha\, VT^3\, \chi_1^B$$

$$\kappa_2[B_1] = \alpha\, VT^3\, \beta\, \chi_2^B$$

$$\kappa_3[B_1] = \alpha\, VT^3\, \beta\, (1-2\alpha)\, \chi_3^B$$

$$\kappa_4[B_1] = \alpha\, VT^3\, \beta\, \left[\chi_4^B \; - 3\alpha\beta \frac{(\chi_3^B)^2 + \chi_2^B\, \chi_4^B}{\chi_2^B}\right]$$

$$\kappa_5[B_1] = \alpha\, VT^3\, \beta\, (1-2\alpha) \left\{[1-2\beta\alpha]\chi_5^B \; - 10\alpha\beta \frac{\chi_3^B \chi_4^B}{\chi_2^B}\right\}$$

$$\kappa_6[B_1] = \alpha\, VT^3\, \beta\, [1 - 5\alpha\beta(1-\alpha\beta)]\, \chi_6^B + 5\, VT^3\, \alpha^2\, \beta^2 \left\{9\alpha\beta \frac{(\chi_3^B)^2\, \chi_4^B}{(\chi_2^B)^2} - 3\alpha\beta \frac{(\chi_3^B)^4}{(\chi_2^B)^3} \right.$$
$$\left. -2(1-2\alpha)^2 \frac{(\chi_4^B)^2}{\chi_2^B} - 3[1-3\beta\alpha] \frac{\chi_3^B\, \chi_5^B}{\chi_2^B}\right\}$$

$$\chi_n^B = \frac{\partial^n (p/T^4)}{\partial(\mu_B/T)^n}$$ – grand-canonical susceptibilities e.g from Lattice QCD!!

Details: Vovchenko, et al. arXiv:2003.13905

Cumulant ratios

Some common cumulant ratios:

scaled variance $$\frac{\kappa_2[B_1]}{\kappa_1[B_1]} = (1-\alpha)\,\frac{\chi_2^B}{\chi_1^B},$$

skewness $$\frac{\kappa_3[B_1]}{\kappa_2[B_1]} = (1-2\alpha)\,\frac{\chi_3^B}{\chi_2^B},$$

kurtosis $$\frac{\kappa_4[B_1]}{\kappa_2[B_1]} = (1-3\alpha\beta)\,\frac{\chi_4^B}{\chi_2^B} - 3\alpha\beta\left(\frac{\chi_3^B}{\chi_2^B}\right)^2.$$

- Global conservation (α) and equation of state (χ_n^B) effects factorize in cumulants up to the 3rd order, starting from κ_4 not anymore
- $\alpha \to 0$: Grand canonical limit
- $\alpha \to 1$: canonical limit
- $\chi_{2n} = \langle N \rangle + \langle \bar{N} \rangle$; $\chi_{2n+1} = \langle N \rangle - \langle \bar{N} \rangle$: recover known results for ideal gas

New data @ 5.02 TeV

ALICE Collaboration, Phys. Rev. C 101 (2020) 044907

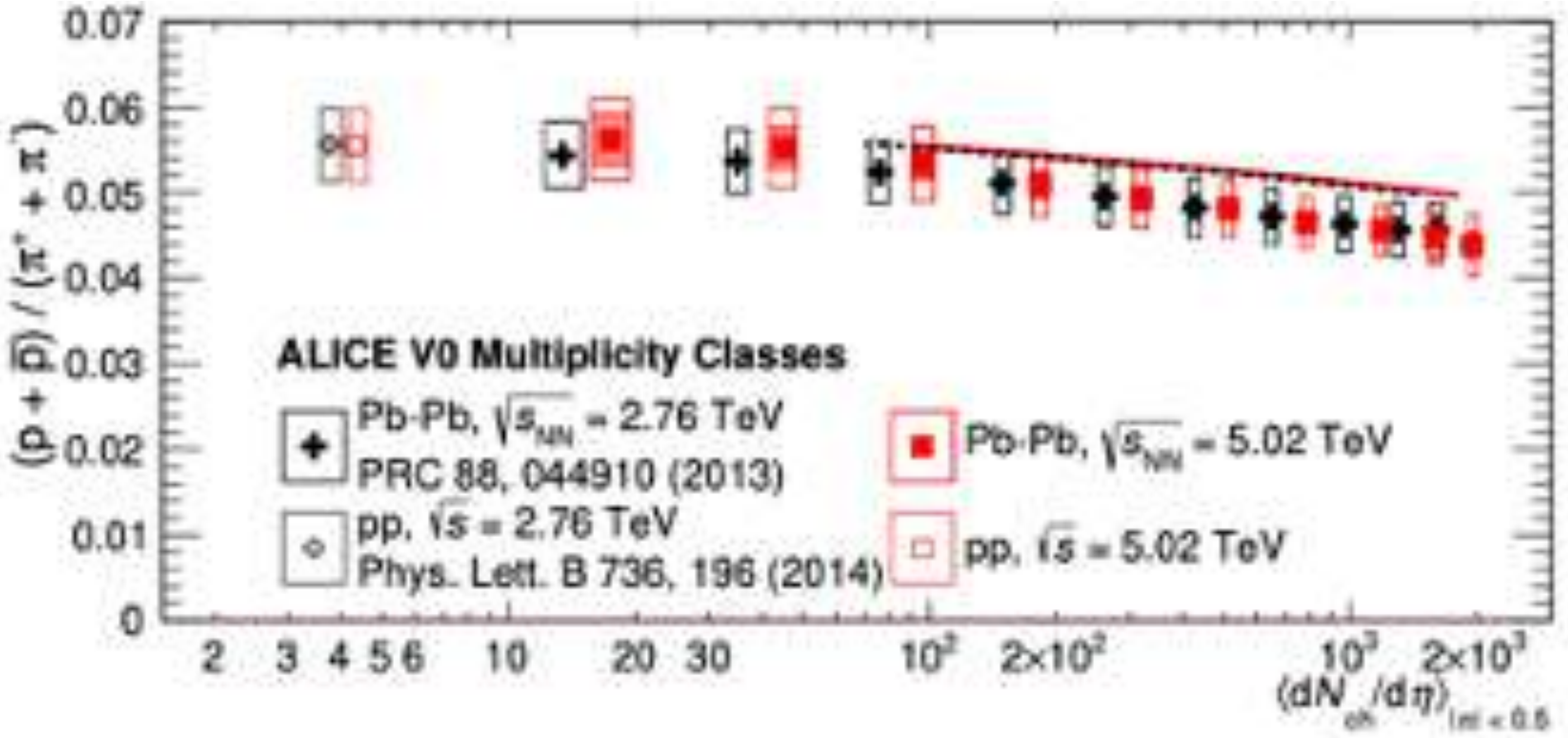
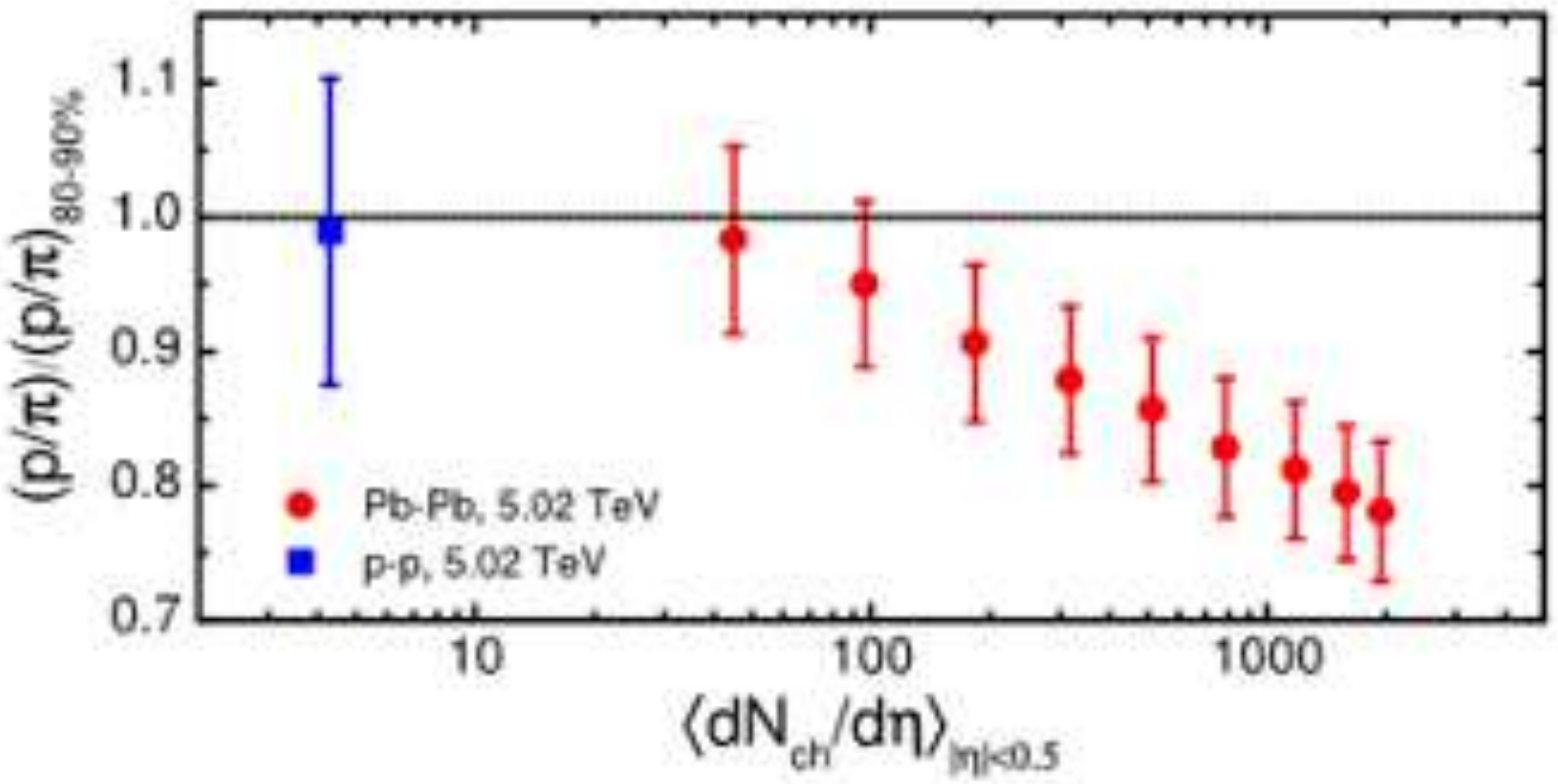


Figure 7: Transverse momentum integrated K/π (top) and p/π (bottom) ratios as a function of $\langle dN_{ch}/d\eta \rangle$ in Pb – Pb collisions at $\sqrt{s_{NN}}$ = 5.02 TeV, compared to Pb – Pb at 2.76 TeV [14]. The values in pp collisions at $\sqrt{s}$ = 5.02 and 2.76 TeV are also shown. The empty boxes show the total systematic uncertainty; the shaded boxes indicate the contribution uncorrelated across centrality bins (not estimated in Pb – Pb collisions at $\sqrt{s_{NN}}$ = 2.76 TeV).

Short hadronic phase

Long hadronic phase



- Evidence for suppression of p/π ration in central collisions (~20%, >4σ level)
- Due to hadronic phase?

For analysis and discussion: See V.Vovlchenko and V.K 2210.15641