Lambda hyperon spin polarization evolution using spin hydro

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Related papers: arXiv:2103.02592, arXiv:2011.14907, arXiv:1901.09655

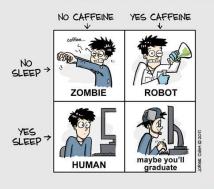
23 Sep, 2021

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Supported by IFJPAN and the NCN Grants No. 2016/23/B/ST2/00717 and No. 2018/30/E/ST2/00432

GRAD SCHOOL ENERGY LEVELS

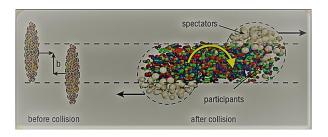


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Heavy-ion collisions:

- Non-central relativistic heavy-ion collisions create global rotation of matter, which may induce spin polarization.
- Emerging particles are expected to be globally polarized with their spins on average pointing along the systems angular momentum.

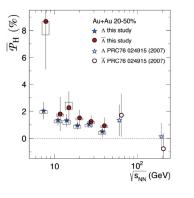
nucl-th/0410079, nucl-th/0410089, arXiv:0708.0035.



Source: CERN Courier

Global polarization:

First positive measurements of global spin polarization of Λ hyperons by STAR





$$\begin{array}{c} \text{thermal approach} \longrightarrow P_{\Lambda} \approx \frac{1}{2} \frac{\omega}{T} + \frac{\mu_{\Lambda} B}{T} P_{\Lambda} \approx \frac{1}{2} \frac{\omega}{T} - \frac{\mu_{\Lambda} B}{T} \\ \text{Becattini, F., Karpenko, I., Lisa, M., Upsal, I., Voloshin, S., PRC 95, 054902 (2017)} \end{array}$$

...the hottest, least viscous – and now, most vortical – fluid produced in the laboratory ... $\omega = \left(P_{\Lambda} + P_{\overline{\Lambda}}\right) k_{B}T/\hbar \sim 0.6 - 2.7 \times 10^{22} \text{s}^{-1}$ L. Adamczyk et al. (STAR) (2017). Nature 548 (2017) 62-65

Even larger than...

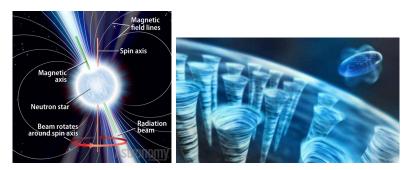


Figure: PSR J1748-2446ad (716 s^{-1}) & Nanodroplets of superfluid helium (10⁷ s^{-1}).

10.1126/science.1123430, Science 345, 906-909 (2014)

MCU



Global polarization:

 Good agreement between experiment and models based on local thermodynamic equilibrium of spin degrees of freedom.

0711.1253, 1304.4427, 1303.3431, 1501.04468, 1610.02506, 1610.04717, 1605.04024, 1703.03770, etc...

But...

Longitudinal polarization:

 Good agreement between experiment and models based on local thermodynamic equilibrium of spin degrees of freedom.

0711.1253, 1304.4427, 1303.3431, 1501.04468, 1610.02506, 1610.04717, 1605.04024, 1703.03770, etc...

But...

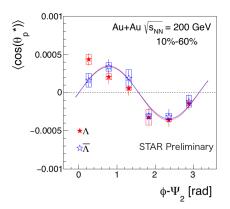


Figure: Longitudinal polarization of $\Lambda - \bar{\Lambda}$ (1905.11917)

Bigger picture:

 This study will help us to know the formation and characteristics of the QGP, a state of matter believed to exist at sufficiently high energy densities.

 Detecting and understanding the QGP allows us to understand better the universe in the moments after the Big Bang.

Our approach:

 Include spin degrees of freedom into the ideal standard hydrodynamics to form spin hydrodynamics formalism.

•
$$J^{\mu,\alpha\beta}(x) = x^{\alpha} T^{\mu\beta}(x) - x^{\beta} T^{\mu\alpha}(x) + S^{\mu,\alpha\beta}(x)$$

- And, conservation of total angular momentum, $\partial_{\lambda}J^{\lambda,\mu\nu}(x)=0$ gives $\partial_{\lambda}S^{\lambda,\mu\nu}(x)=T^{\nu\mu}(x)-T^{\mu\nu}(x)$
- For symmetric energy-momentum tensor, $T_{\rm GLW}^{\nu\mu}(x)=T_{\rm GLW}^{\mu\nu}(x)$, we have $\partial_{\lambda}S_{\rm GLW}^{\lambda,\mu\nu}(x)=0$
- Hence conservation of the angular momentum implies the conservation of its spin part in the de Groot-van Leeuwen-van Weert (GLW) formulation.

1705.00587, 1712.07676, 1806.02616, 1811.04409, S. R. De Groot *et. al.*, Relativistic Kinetic Theory: Principles and Applications (1980).

Steps of spin hydrodynamic framework:

- Solving the standard perfect-fluid hydrodynamic equations without spin.
- Determination of the spin evolution in the hydrodynamic background.

- Determination of the Pauli-Lubański (PL) vector on the freeze-out hypersurface.
- Calculation of the spin polarization of particles in their rest frame.
 The spin polarization obtained is a function of the three-momenta of particles and can be directly compared with the experiment.

1901 09655

Conservation laws:

- $d_{\alpha} N^{\alpha} \equiv d_{\alpha} (\mathcal{N} U^{\alpha}) = 0 \rightarrow \text{Conservation of net baryon number}$ where $\mathcal{N} = 4 \sinh(\mu/T) \mathcal{N}_{(0)}$.
- $d_{\alpha} T^{\alpha\beta} \equiv d_{\alpha} \left[(\mathcal{E} + \mathcal{P}) U^{\alpha} U^{\beta} \mathcal{P} g^{\alpha\beta} \right] = 0 \rightarrow \text{Conservation of EMT}$ where $\mathcal{E} = 4 \cosh(\mu/T) \mathcal{E}_{(0)}$ and $\mathcal{P} = 4 \cosh(\mu/T) \mathcal{P}_{(0)}$.

These laws provide closed system of 5 eqns. for 5 funcs: μ , T, and three independent components of U^{μ} which need to be solved to get the perfect-fluid background evolution.

•
$$d_{\alpha}S^{\alpha,\beta\gamma} \equiv d_{\alpha} \Big[\mathcal{A}_1 U^{\alpha} \omega^{\beta\gamma} + \mathcal{A}_2 U^{\alpha} U^{[\beta} \kappa^{\gamma]} + \mathcal{A}_3 (U^{[\beta} \omega^{\gamma]\alpha} + g^{\alpha[\beta} \kappa^{\gamma]}) \Big] = 0$$

Conservation of spin

where
$$\mathcal{A}_1 = \cosh(\mu/T) \left(n_{(0)} - \mathcal{B}_{(0)} \right)$$
, $\mathcal{A}_2 = \cosh(\mu/T) \left(\mathcal{A}_{(0)} - 3\mathcal{B}_{(0)} \right)$, $\mathcal{A}_3 = \cosh(\mu/T) \mathcal{B}_{(0)}$ with, $\mathcal{B}_{(0)} = -\frac{2}{(m/T)^2} (\mathcal{E}_{(0)} + \mathcal{P}_{(0)})/T$ and $\mathcal{A}_{(0)} = -3\mathcal{B}_{(0)} + 2\mathcal{N}_{(0)}$. Here, $\omega^{\beta\gamma}$ is known as spin polarization tensor.

1811.04409

Spin polarization tensor:

 $\omega_{\mu\nu}$ is an anti-symmetric tensor of rank 2 and can be parameterized by the four-vectors κ^μ and ω^μ ,

$$\omega_{\mu\nu} = \kappa_{\mu} U_{\nu} - \kappa_{\nu} U_{\mu} + \epsilon_{\mu\nu\alpha\beta} U^{\alpha} \omega^{\beta},$$

where $\kappa^{\alpha} = \frac{C_{\kappa X} X^{\alpha} + C_{\kappa Y} Y^{\alpha} + C_{\kappa Z} Z^{\alpha}}{U, X, Y \text{ and } Z \text{ form a 4-vector basis satisfying the following normalization conditions: } U \cdot U = 1, \qquad X \cdot X = Y \cdot Y = Z \cdot Z = -1.$

$$\omega_{\alpha\beta} = \begin{bmatrix} 0 & e^1 & e^2 & e^3 \\ -e^1 & 0 & -b^3 & b^2 \\ -e^2 & b^3 & 0 & -b^1 \\ -e^3 & -b^2 & b^1 & 0 \end{bmatrix}$$

$$\begin{split} C_{\kappa X} &= e^1 \cosh(\vartheta + \eta) - b^2 \sinh(\vartheta + \eta), \\ C_{\omega Y} &= b^2 \cosh(\vartheta + \eta) - e^1 \sinh(\vartheta + \eta), \\ C_{\kappa Y} &= e^2 \cosh(\vartheta + \eta) + b^1 \sinh(\vartheta + \eta), \\ C_{\omega X} &= b^1 \cosh(\vartheta + \eta) + e^2 \sinh(\vartheta + \eta), \\ C_{\kappa Z} &= e^3, \quad C_{\omega Z} &= b^3. \end{split}$$

Spin polarization (local and global):

$$\langle \pi_{\mu} \rangle_{p}({m p}) = rac{E_{p} \frac{d\Pi_{\mu}^{T}(p)}{d^{3}p}}{E_{p} \frac{dN(p)}{d^{3}p}} \quad {
m {\rightarrow Total Pauli Lubanski vector} \over
ightarrow Momentum density of all particles}} \sim {m p} \; dependent$$

$$\left\langle \boldsymbol{\pi}_{\boldsymbol{\mu}} \right\rangle = \frac{\int dP \left\langle \boldsymbol{\pi}_{\boldsymbol{\mu}} \right\rangle_{p} E_{p} \frac{d\mathcal{N}(p)}{d^{3}p}}{\int dP E_{p} \frac{d\mathcal{N}(p)}{d^{3}p}} = \frac{\int d^{3}p}{\int d^{3}p} \frac{\frac{d\Pi_{\boldsymbol{\mu}}^{'}(p)}{d^{3}p}}{\int d^{3}p} \sim \boldsymbol{p} \text{ integrated}$$

where,
$$E_p \frac{d\Pi_\mu^*(p)}{d^3p} = \frac{1}{(2\pi)^3 m} \int \cosh(\frac{\mu}{T}) \Delta \Sigma_\lambda p^\lambda \, e^{-\beta \cdot p} \left(\tilde{\omega}_{\beta\mu} \, \, p^\beta \right)^*$$

$$E_p rac{d\mathcal{N}(p)}{d^3p} = rac{4}{(2\pi)^3} \int \cosh(rac{\mu}{T}) \Delta \Sigma_\lambda p^\lambda \, e^{-eta \cdot p}$$

* meaning quantities calculated in particle rest frame.

1901.09655

Spin angular momentum components at the freeze-out:

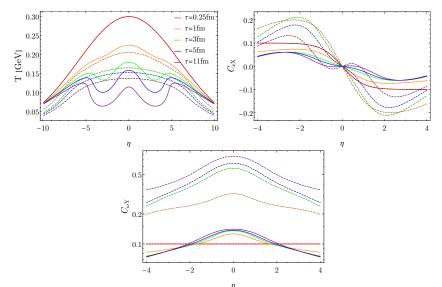
$$\begin{split} S_{\rm FO}^{\mu\nu} &= \int \Delta \Sigma_{\lambda} S_{\rm GLW}^{\lambda,\mu\nu} = \int dx dy \, \tau d\eta \, U_{\lambda}^{\rm B} S_{\rm GLW}^{\lambda,\mu\nu} \\ &= \pi R^2 \tau \int_{-\eta_{\rm FO}/2}^{+\eta_{\rm FO}/2} d\eta \, U_{\lambda}^{\rm B} \, S_{\rm GLW}^{\lambda,\mu\nu} \end{split}$$

$$S_{13}^{\rm FO} = \pi R^2 \tau \int_{-\eta_{\rm FO}/2}^{+\eta_{\rm FO}/2} d\eta \left[\mathcal{A}_1 \mathcal{C}_{\omega Y} \cosh(\eta) + \mathcal{A}_3 \mathcal{C}_{\kappa X} \sinh(\eta) \right].$$

$$\mathcal{C}_{\kappa X} = -\mathcal{C}_{\omega Y} \tanh(\vartheta + \eta).$$



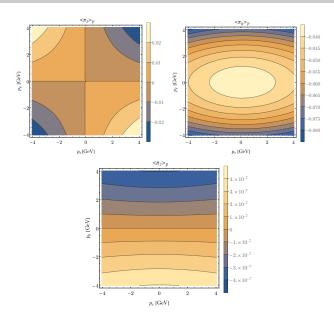
Perfect-fluid background and spin components evolution:



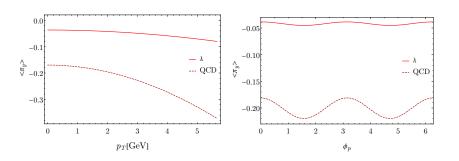
Solid lines: use of ideal gas EoS in both perfect fluid background and spin evolution. Dashed lines: use of lattice QCD EoS in

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Momentum dependence of polarization:



Global polarization:



Summary:

- Discussed relativistic hydrodynamics with spin based on the GLW formulation of energy-momentum and spin tensors.
- Showed how our formalism can be compared with the experiments.
- Obtained dynamics of spin polarization in the non-boost background.
- Incorporation of spin in full 3+1D hydro model required to address the problem of longitudinal polarization (which will be out pretty soon, stay tuned).

Thank you for your attention!

Acknowledgments

The authors wish to thank coffee, coff

AN HONEST ACKNOWLEDGMENT SECTION

JORGE CHAM @ 2015

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