

# *NLO corrections to DDVCS, DVCS and TCS processes*

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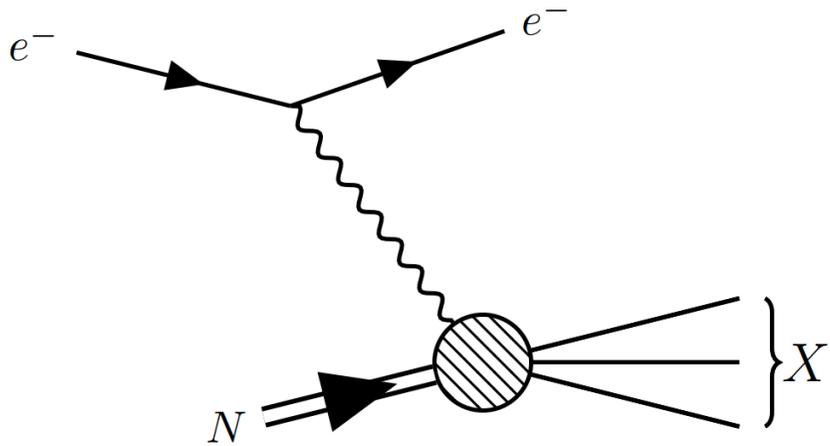
Seminar at 61 Cracow School for Theoretical Physics: Electron-Ion Collider Physics

September 22, 2021

# Why NLO corrections?

- ▶ Precise determination of GPDs (and so nucleon tomography). Also, gluon GPDs appear at NLO!!
- ▶ Universality testing in processes that at Born level look alike: DVCS & TCS
  - “On timelike and spacelike hard exclusive reactions”, D. Mueller, B. Pire, L. Szymanowski, J. Wagner - DOI: [10.1103/PhysRevD.86.031502](https://doi.org/10.1103/PhysRevD.86.031502)
  - “Data-driven study of timelike Compton scattering”, O. Grocholsky, H. Moutarde, B. Pire, P. Sznajder, J. Wagner - DOI: [10.1140/epjc/s10052-020-7700-9](https://doi.org/10.1140/epjc/s10052-020-7700-9)
- ▶ Apparent divergences of the form  $\ln 0$
- ▶ Study of feasibility of DDVCS at EIC
- ▶ Framework for analysis of JLab 12 GeV data

# Refreshing Deep Inelastic Scattering (DIS)



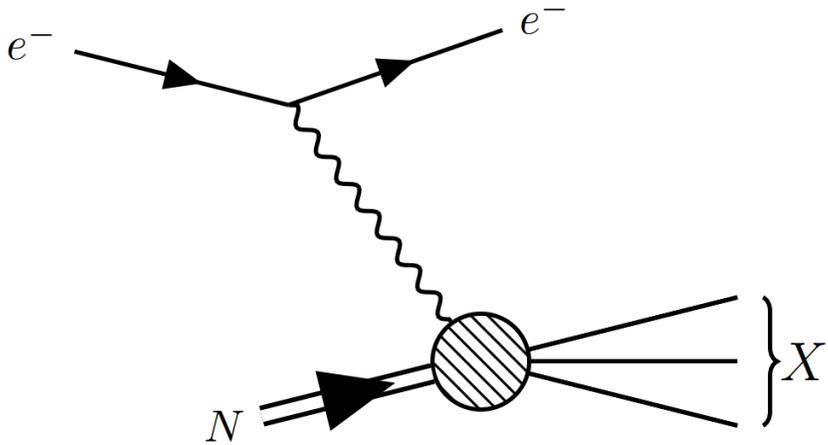
Factorization in cross-section

$$d\sigma_{\text{DIS}} \sim L^{\mu\nu} W_{\mu\nu}$$

pQED

Non-perturbative QCD

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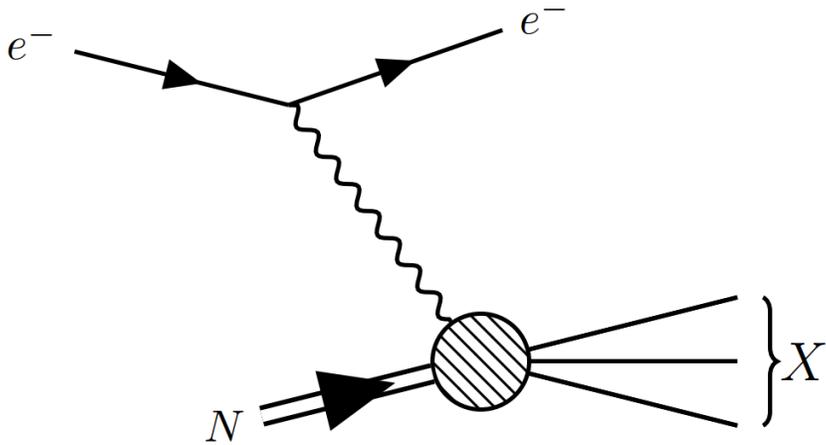
Optical theorem to DIS

$$\sum_X \int_X \left| \begin{array}{c} \text{Diagram: } e^- \text{ emits } \gamma^* \text{ which interacts with } N \text{ to produce } X \\ \hline \text{Diagram: } N \text{ emits } \gamma^* \text{ which interacts with } e^- \text{ to produce } X \end{array} \right|^2$$

=

$$2\text{Im} \left( \begin{array}{c} \text{Diagram: } e^- \text{ emits } \gamma^* \text{ which interacts with } N \text{ to produce } X \\ \hline \text{Diagram: } N \text{ emits } \gamma^* \text{ which interacts with } e^- \text{ to produce } X \end{array} \right)$$

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Optical theorem to DIS

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$$\sum_X \int_X |X\rangle\langle X| = 1 \rightarrow \text{DIS is inclusive process}$$

Cf. Riedl's talk for more details in kinematics and experimental results

# Parton Distribution Functions (PDFs)

- ▶ Extracted from the hadronic tensor  $W_{\mu\nu}$ , they determine the internal nucleon structure

$$\text{PDF}(x) = \frac{1}{2} \int \frac{dz^+}{2\pi} e^{ixp^- z^+} \langle P | \bar{q}_f(-z/2) \gamma^+ \mathcal{W}[-z/2, z/2] q_f(z/2) | P \rangle \Big|_{z_\perp = z^- = 0}$$

Wilson line definition

$$\mathcal{W}[z_1^+, z_2^+] = \mathbb{P} \exp \left[ ig \int_{z_2^+}^{z_1^+} d\eta^+ A^- \right]$$

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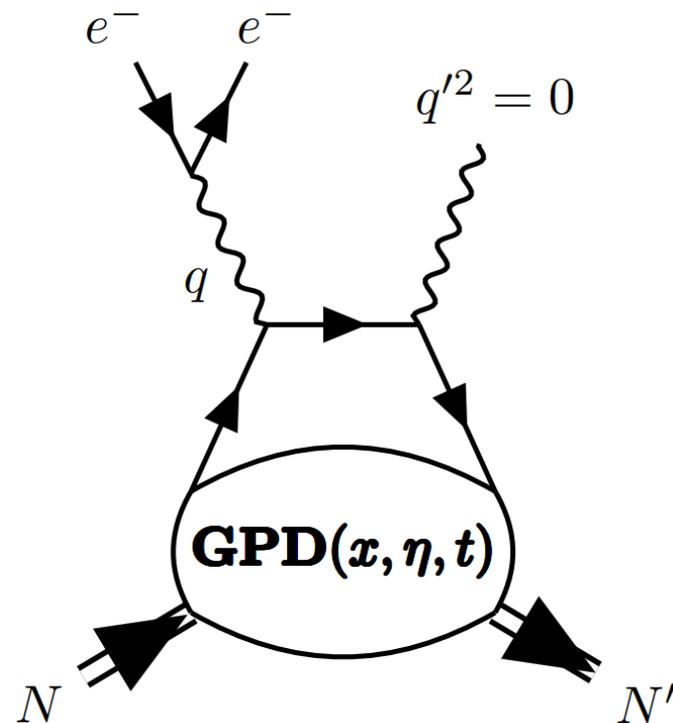
$$\gamma^+ = \frac{\gamma^0 + \gamma^3}{\sqrt{2}}$$

**PDF is 1D**

$x$  = longitudinal parton momentum fraction

# Improving PDF's 1D picture

- ▶ In the late '90s, Ji, Müller and Radyushkin introduce Generalized Parton Distributions (GPDs) through Deeply Virtual Compton Scattering (DVCS) process
- ▶ The point now is to study the conversion of a virtual photon into a real one



Cf. Ji's talk for theory; and Ent and Newman's talks for phenomenology

# DVCS = exclusive process = factorization in amplitude

Sketch of DVCS amplitude

$$\begin{aligned}\mathcal{A}_{\text{DVCS}} &\sim \int_{-1}^1 dx \frac{1}{x - \eta + i0} \text{GPD}(x, \eta, t) + \dots \\ &= \text{PV} \left( \frac{1}{x - \eta} \right) (\text{GPD}(x, \eta, t)) - \int_{-1}^1 dx i\pi \delta(x - \eta) \text{GPD}(x, \eta, t) + \dots\end{aligned}$$

So we can measure GPDs at  $x = \eta$  only, i.e., we can access  $\text{GPD}(\eta, \eta, t)$

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Dispersion relation  $\Rightarrow$  real part can be computed in terms of imaginary:

$$\Re \mathcal{A} = \text{PV} \int_{-1}^1 dx \frac{\Im \mathcal{A}}{x - \eta} + D(t)$$

Cf. Riedl's talk for info. about experimental measurement of Re and Im parts

# GPD definition: 3D distribution

$$\text{GPD}(x, \eta, t) = \frac{1}{2} \int \frac{dz^+}{2\pi} e^{ixp^- z^+} \langle P' | \bar{q}_f(-z/2) \gamma^+ \mathcal{W}[-z/2, z/2] q_f(z/2) | P \rangle \Big|_{z_\perp = z^- = 0}$$

Measure the difference  
between  $P$  and  $P'$

$$\eta = -\frac{(q - q')(q + q')}{(P + P')(q + q')}$$

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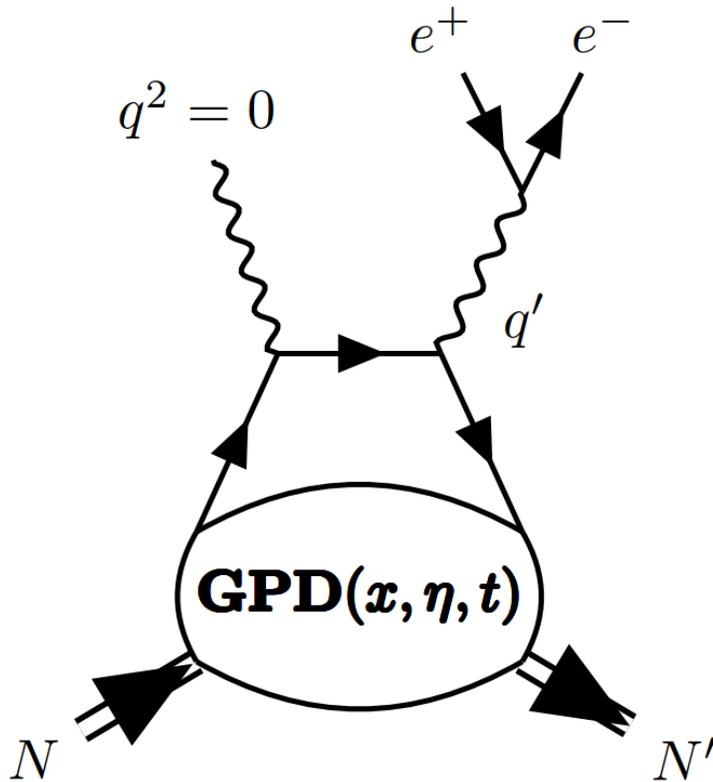
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Nucleon tomography via Fourier  
transform in the plane transverse to  
proton motion

# Other 2 “golden channels”: TCS & DDVCS

but experimentally more challenging than DVCS

- ▶ TCS or timelike Compton scattering
- ▶ Counterpart of DVCS
- ▶ A real photon transforms into a virtual one (lepton photo-production)



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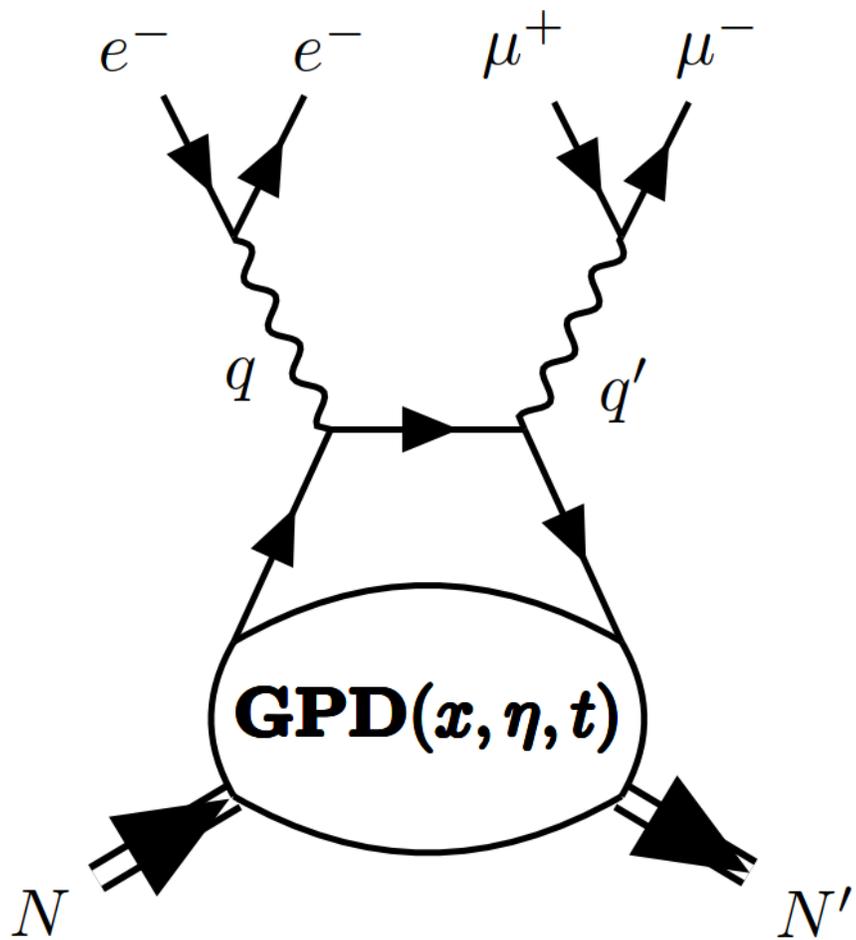
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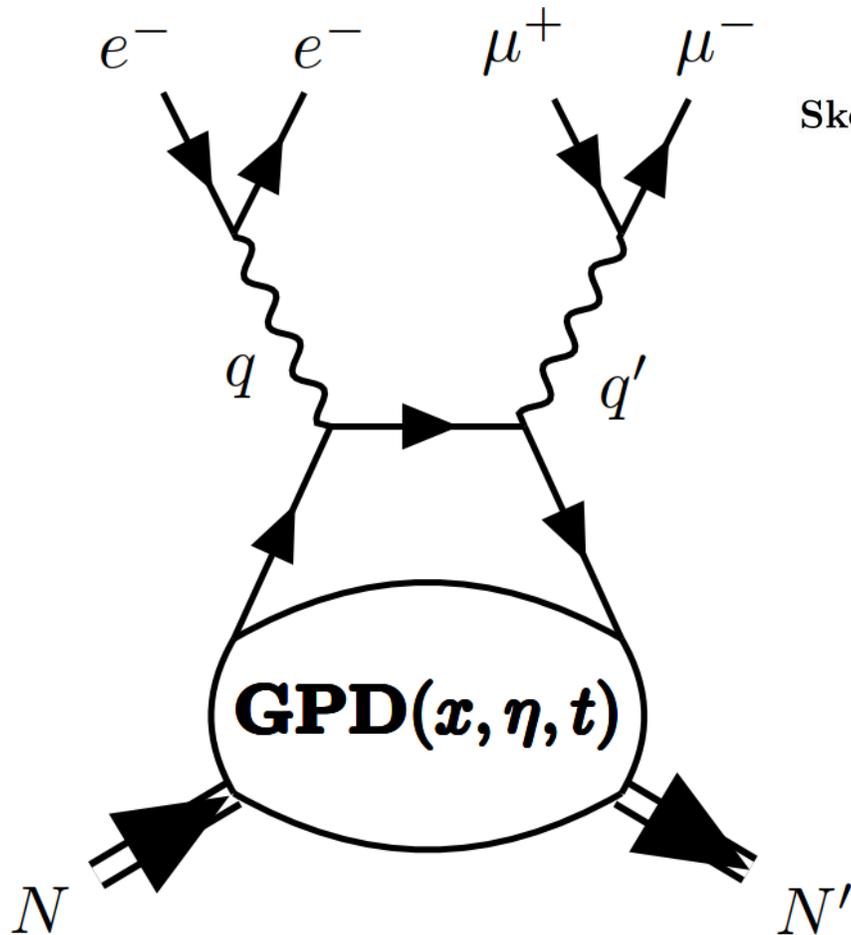
**Recent paper (from this August) on the 1st measurement of TCS by CLAS collaboration:**

**“First-time measurement of Timelike Compton Scattering”, P. Chatagnon et al., arXiv:2108.11746 [hep-ex]**

- ▶ DDVCS or double DVCS
- ▶ 2 virtual photons: spacelike (incoming) and timelike (outgoing)



- ▶ DDVCS or double DVCS
- ▶ 2 virtual photons: spacelike (incoming) and timelike (outgoing)
- ▶ **Allows to measure GPDs outside  $x = \eta$**



Sketch of DDVCS amplitude

$$\begin{aligned} \mathcal{A}_{\text{DDVCS}} &\sim \int_{-1}^1 dx \frac{1}{x - \xi + i0} \text{GPD}(x, \eta, t) + \dots \\ &= \text{PV} \left( \frac{1}{x - \xi} \right) (\text{GPD}(x, \eta, t)) - \int_{-1}^1 dx i\pi \delta(x - \xi) \text{GPD}(x, \eta, t) + \dots \end{aligned}$$

So now we can access  $\text{GPD}(\xi, \eta, t)$

Related to Orginos' talk: GPDs extraction is not ill-defined problem thanks to DDVCS. LQCD will render GPDs for sure, but DDVCS makes the problem well-defined and experimentally solvable

# Details in DDVCS

Here,  $\xi$  is the *generalized* Bjorken variable,

$$\xi = \frac{-\bar{q}^2}{2p\bar{q}}, \quad \bar{q} = \frac{q + q'}{2}, \quad p = \frac{P + P'}{2}$$

$q, q'$  are the 4-momenta of incoming & outgoing photon

$P, P'$  are the 4-momentum of incoming & outgoing proton

**Experimentally, DDVCS is very demanding: x-sec smaller than DVCS' → EIC will have enough luminosity to accurate measurements**

# Renormalization

- ▶ Need to renormalize both GPDs and hard part. Amplitude:

$$\mathcal{A}^{\mu\nu} = g_T^{\mu\nu} \int_{-1}^1 dx \left[ \sum_q^{n_F} \tilde{T}^q(x) \tilde{F}^q(x) + \tilde{T}^g(x) \tilde{F}^g(x) \right]$$

Bare GPD

Bare hard-part coefficients

- ▶ We can work in DDVCS and use

$$\text{DDVCS}|_{\xi=\eta} \rightarrow \text{DVCS}$$

$$\text{DDVCS}|_{\xi=-\eta} \rightarrow \text{TCS}$$

# GPDs at NLO

$$\begin{aligned}\tilde{F}^q(x) = & F^q(x) - \left( \frac{1}{\epsilon} + \frac{1}{2} \ln \frac{e^\gamma \mu_F^2}{4\pi\mu_R^2} \right) K^{qq}(x, x') \otimes F^q(x') \\ & - \left( \frac{1}{\epsilon} + \frac{1}{2} \ln \frac{e^\gamma \mu_F^2}{4\pi\mu_R^2} \right) K^{qg}(x, x') \otimes F^g(x')\end{aligned}$$

Similar formula for gluons

Kernels  $K^{qq, qg}$  can be read from

M. Diehl, Phys. Rept. **388** (2003) 41;

A. V. Belitsky and A. V. Radyushkin, Phys. Rept. **418**, 1 (2005)

# Corrections and $\ln 0$

- ▶ When calculating hard part you find expressions that can be written down as

$$\tilde{T}^q \sim \text{prefactors} \times \ln \left( \frac{\bar{Q}^2}{\xi \mu_R^2} (\xi - x) - i0 \right) + \dots \quad \bar{Q}^2 = \frac{Q^2 - Q'^2}{2}$$

If  $Q^2 \sim Q'^2$ , then  $\bar{Q}^2 \rightarrow 0$  and  $\ln$  is divergent

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Complete expressions in

B. Pire, L. Szymanowski, and J. Wagner. NLO corrections to timelike, spacelike and double deeply virtual Compton scattering. *Phys. Rev. D*, 83:034009, 2011.

$$\frac{\bar{Q}^2}{\xi} = \frac{Q^2}{x_B} - \hat{Q}^2 \xrightarrow{Q^2 \rightarrow Q'^2} 2Pq - Q^2 = s - P^2 = s - M^2 > 0$$

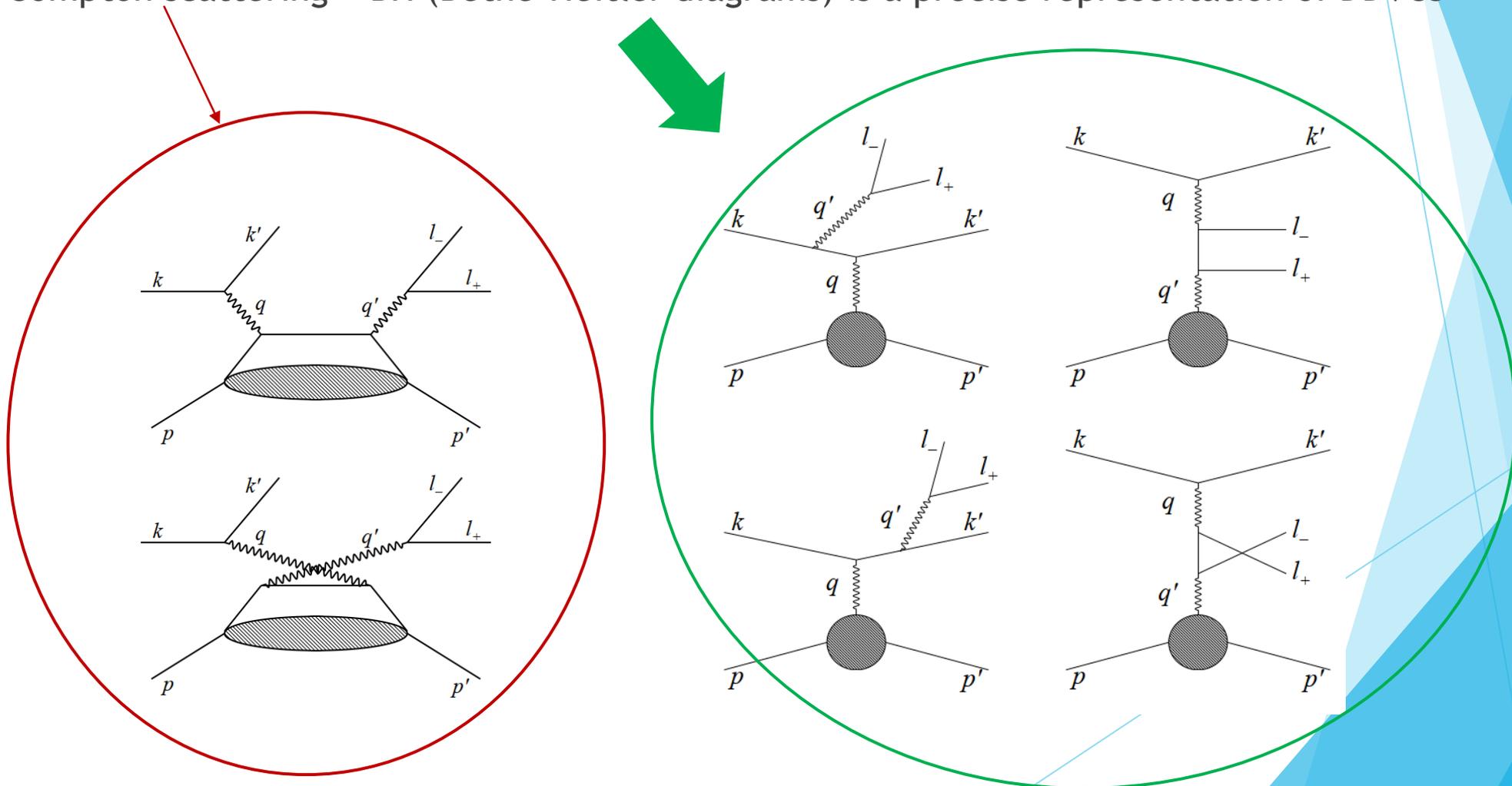
$$\Rightarrow \frac{\bar{Q}^2}{\xi} \not\rightarrow 0$$

Actually,

$$\frac{\bar{Q}^2}{\xi} = \frac{\hat{Q}^2}{\eta} > 0 \Rightarrow \text{no divergent ln}$$

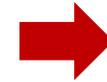
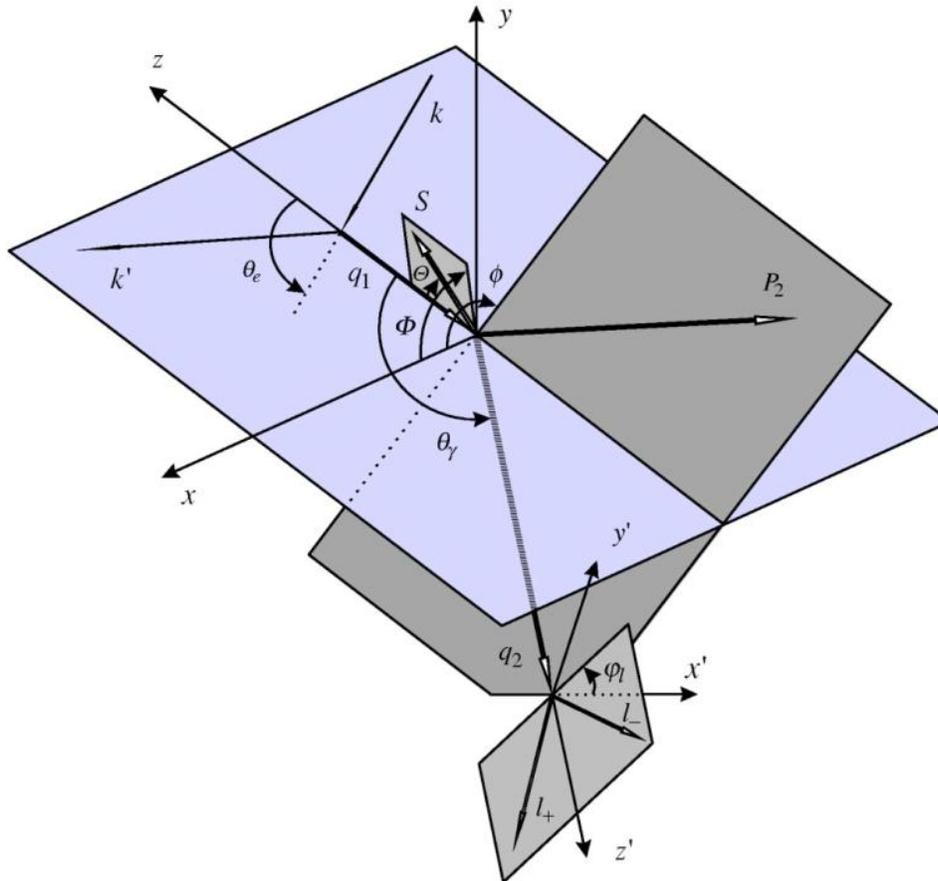
# Not only Compton scattering counts...

- ▶ Compton scattering + BH (Bethe-Heitler diagrams) is a precise representation of DDVCS



# Amplitude structure

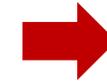
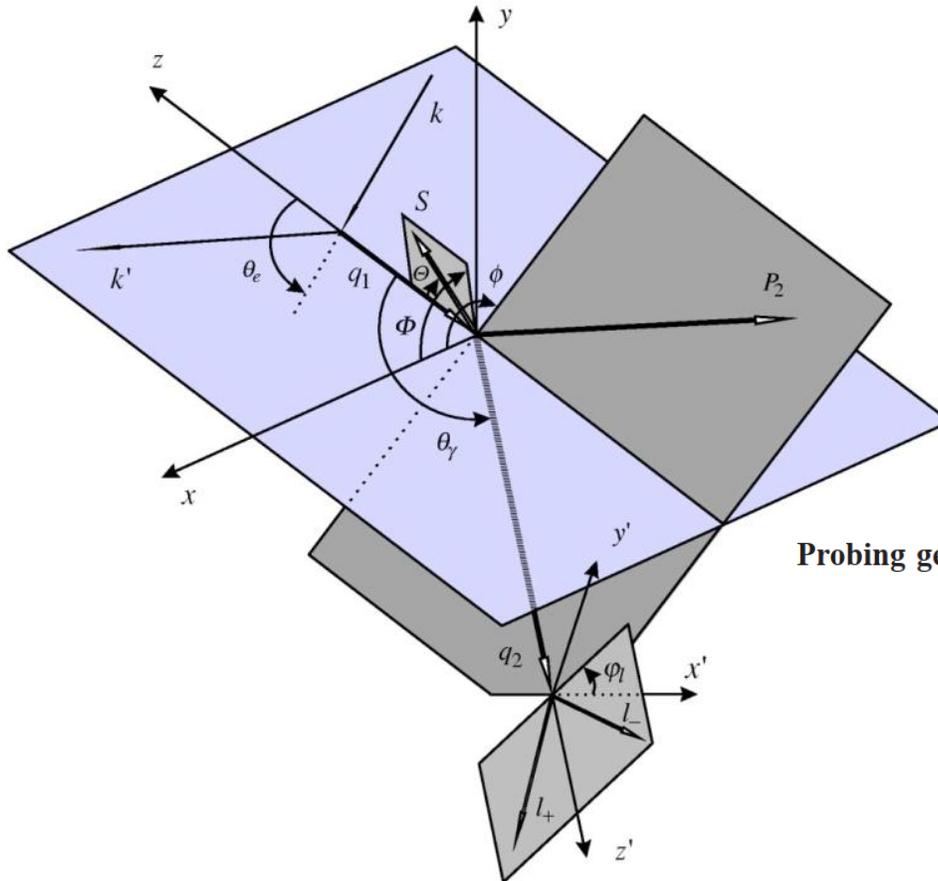
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- ▶ Set-up:



Amplitude can be written as a Fourier expansion in the angles.

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Amplitude can be written as a Fourier expansion in the angles.

LO expressions can be found in

PHYSICAL REVIEW D **68**, 116005 (2003)

Probing generalized parton distributions with electroproduction of lepton pairs off the nucleon

A. V. Belitsky     D. Müller

# Specialized software for this study

- ▶ PARTONS platform: open-source C++ program
  - ▶ Contains several GPD models
  - ▶ Leading twist... but higher twists will be included in near future
  - ▶ Can be used by theorists and experimentalists
  - ▶ Provides x-secs, Compton form factors, etc



PARTonic Tomography Of Nucleon Software

To download and for tutorials

<http://partons.cea.fr>

For detail description of architecture see:

[Eur. Phys. J. C78 \(2018\), 478](#)

# Summary

- ▶ NLO is needed to describe not only quark GPD but also gluon GPD, that appear at  $\mathcal{O}(\alpha_s)$  in hard part
- ▶ DDVCS is a remarkable process since it allows for GPD measurement outside  $x = \eta$ , at can be reduced to DVCS and TCS (measurements are reported for both)
- ▶ I have presented a framework for the study of DDVCS feasibility at EIC, and extraction of GPDs through JLab 12GeV data. All these, at NLO accuracy  $\rightarrow$  precise determination of GPDs and observables (asymmetries, xsecs, etc)