

Generic cosmological solution without singularity from bifurcation analysis

Franciszek Humieja¹, Marek Szydłowski^{1,2}

¹Astronomical Observatory of the Jagiellonian University

²Mark Kac Complex Systems Research Centre, Jagiellonian University

59th Cracow School of Theoretical Physics

Zakopane, June 19, 2019

arXiv:1901.06578

Scalar field approach:

Scalar field approach:

- ▶ Uses the **scalar field** ϕ in order to describe **dark energy and inflaton field**.

Scalar field approach:

- ▶ Uses the **scalar field** ϕ in order to describe **dark energy and inflaton field**.
- ▶ Provides description of the acceleration of the Universe with the **energy density of ϕ dependent on time**.

Scalar field approach:

- ▶ Uses the **scalar field** ϕ in order to describe **dark energy and inflaton field**.
- ▶ Provides description of the acceleration of the Universe with the **energy density of ϕ dependent on time**.
- ▶ In general, **requires implementation of the potential** $U(\phi)$, affecting the scalar field ϕ .

Purpose of the research

Purpose of the research

- ▶ An investigation of the cosmological model with **scalar field non-minimally coupled to the gravity**, describing both **inflation** and **late-time acceleration**.

Purpose of the research

- ▶ An investigation of the cosmological model with **scalar field non-minimally coupled to the gravity**, describing both **inflation** and **late-time acceleration**.
- ▶ A meticulous analysis of changes of the dynamics of the cosmological equations **under a variation of the parameters**.

Purpose of the research

- ▶ An investigation of the cosmological model with **scalar field non-minimally coupled to the gravity**, describing both **inflation** and **late-time acceleration**.
- ▶ A meticulous analysis of changes of the dynamics of the cosmological equations **under a variation of the parameters**.
- ▶ An extraction of **scenarios with an emergence of a non-singular evolution of the universe** starting from de Sitter state (inflation) to another de Sitter state (late-time acceleration after inclusion of matter).

Purpose of the research

- ▶ An investigation of the cosmological model with **scalar field non-minimally coupled to the gravity**, describing both **inflation** and **late-time acceleration**.
- ▶ A meticulous analysis of changes of the dynamics of the cosmological equations **under a variation of the parameters**.
- ▶ An extraction of **scenarios with an emergence of a non-singular evolution of the universe** starting from de Sitter state (inflation) to another de Sitter state (late-time acceleration after inclusion of matter).
- ▶ A study of bifurcations within given scenarios—how many **qualitatively different behaviours** occur?

Non-minimal coupling model

Non-minimal coupling model

- The action

$$S = S_g + S_\phi,$$

where

$$S_g = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} R \quad (\text{Einstein-Hilbert action}),$$

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{R-\phi \text{ coupling}} \right],$$

with

Non-minimal coupling model

- ▶ The action

$$S = S_g + S_\phi,$$

where

$$S_g = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} R \quad (\text{Einstein-Hilbert action}),$$

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{R-\phi \text{ coupling}} \right],$$

with

- ▶ $\varepsilon = \pm 1$ —**canonical** or **phantom** scalar field, respectively,

Non-minimal coupling model

- ▶ The action

$$S = S_g + S_\phi,$$

where

$$S_g = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} R \quad (\text{Einstein-Hilbert action}),$$

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{R-\phi \text{ coupling}} \right],$$

with

- ▶ $\varepsilon = \pm 1$ —**canonical** or **phantom** scalar field, respectively,
- ▶ ξ —**coupling parameter**, $\xi \neq 0$ —non-minimal coupling,

Non-minimal coupling model

- ▶ The action

$$S = S_g + S_\phi,$$

where

$$S_g = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} R \quad (\text{Einstein-Hilbert action}),$$

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{R\text{-}\phi \text{ coupling}} \right],$$

with

- ▶ $\varepsilon = \pm 1$ —**canonical** or **phantom** scalar field, respectively,
- ▶ ξ —**coupling parameter**, $\xi \neq 0$ —non-minimal coupling,
- ▶ $\kappa^2 = 8\pi G$, $c = 1$ and the signature is $(-+++)$.

Non-minimal coupling model

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{\text{R-}\phi \text{ coupling}} \right],$$

Non-minimal coupling model

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{R-\phi \text{ coupling}} \right],$$

- ▶ The introduction of NMC is forced upon us in many situations of physical and cosmological interest

Non-minimal coupling model

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\varepsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{2U(\phi)}_{\text{potential energy}} + \underbrace{\varepsilon \xi R \phi^2}_{\text{R-}\phi \text{ coupling}} \right],$$

- ▶ The introduction of NMC is forced upon us in many situations of physical and cosmological interest
- ▶ NMC arises at the quantum level due to quantum corrections and is required to the renormalizability of the scalar field theory

Non-minimal coupling model

$$S_\phi = -\frac{1}{2} \int d^4x \sqrt{-g} \left[\underbrace{\epsilon \nabla^\alpha \phi \nabla_\alpha \phi}_{\text{kinetic energy}} + \underbrace{+2U(\phi)}_{\text{potential energy}} + \underbrace{+\epsilon \xi R \phi^2}_{\text{R-}\phi \text{ coupling}} \right],$$

- ▶ The introduction of NMC is forced upon us in many situations of physical and cosmological interest
- ▶ NMC arises at the quantum level due to quantum corrections and is required to the renormalizability of the scalar field theory
- ▶ The value of ξ is not arbitrary but is determined by the underlying physics (allowed ranges for ξ)

Field equations

- Variation of the action w.r.t. $g^{\mu\nu}$ and ϕ produces Einstein and Klein-Gordon equations, respectively:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa^2 T_{\mu\nu}^{(\phi)},$$

$$\square\phi - \xi R\phi - \varepsilon U_{,\phi} = 0,$$

where the energy-momentum tensor

$$\begin{aligned} T_{\mu\nu}^{(\phi)} = & \varepsilon \nabla_{\mu}\phi \nabla_{\nu}\phi - \frac{1}{2}\varepsilon g_{\mu\nu} \nabla^{\alpha}\phi \nabla_{\alpha}\phi - g_{\mu\nu} U(\phi) + \\ & + \varepsilon \xi \phi^2 \left(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R \right) + \varepsilon \xi (g_{\mu\nu} \square\phi^2 - \nabla_{\mu}\nabla_{\nu}\phi^2), \end{aligned}$$

and

$$U_{,\phi} := dU/d\phi$$

Assumptions

Assumptions

- ▶ Spatially flat ($k = 0$) universe with Friedmann-Lemaître-Robertson-Walker (FLRW) symmetry

$$ds^2 = -dt^2 + a^2(t) (dx^2 + dy^2 + dz^2),$$

Assumptions

- ▶ Spatially flat ($k = 0$) universe with Friedmann-Lemaître-Robertson-Walker (FLRW) symmetry

$$ds^2 = -dt^2 + a^2(t) (dx^2 + dy^2 + dz^2),$$

- ▶ Linear barotropic equation of state between energy density ρ_ϕ and pressure p_ϕ

$$p_\phi = w_\phi \rho_\phi.$$

Assumptions

- ▶ Spatially flat ($k = 0$) universe with Friedmann-Lemaître-Robertson-Walker (FLRW) symmetry

$$ds^2 = -dt^2 + a^2(t) (dx^2 + dy^2 + dz^2),$$

- ▶ Linear barotropic equation of state between energy density ρ_ϕ and pressure p_ϕ

$$p_\phi = w_\phi \rho_\phi.$$

- ▶ We use the **Ratra-Peebles potential**

$$U(\phi) = \frac{M^{n+4}}{\phi^n},$$

where n is a dimensionless parameter and $M > 0$ is a dimensional constant.

Assumptions

- ▶ Spatially flat ($k = 0$) universe with Friedmann-Lemaître-Robertson-Walker (FLRW) symmetry

$$ds^2 = -dt^2 + a^2(t) (dx^2 + dy^2 + dz^2),$$

- ▶ Linear barotropic equation of state between energy density ρ_ϕ and pressure p_ϕ

$$p_\phi = w_\phi \rho_\phi.$$

- ▶ We use the **Ratra-Peebles potential**

$$U(\phi) = \frac{M^{n+4}}{\phi^n},$$

where n is a dimensionless parameter and $M > 0$ is a dimensional constant.

- ▶ Well known form of potential in quintessence models

Assumptions

- ▶ Spatially flat ($k = 0$) universe with Friedmann-Lemaître-Robertson-Walker (FLRW) symmetry

$$ds^2 = -dt^2 + a^2(t) (dx^2 + dy^2 + dz^2),$$

- ▶ Linear barotropic equation of state between energy density ρ_ϕ and pressure p_ϕ

$$p_\phi = w_\phi \rho_\phi.$$

- ▶ We use the **Ratra-Peebles potential**

$$U(\phi) = \frac{M^{n+4}}{\phi^n},$$

where n is a dimensionless parameter and $M > 0$ is a dimensional constant.

- ▶ Well known form of potential in quintessence models
- ▶ Allows for a solution of the fine tuning of initial conditions

Cosmological equations

Cosmological equations

- ▶ Applying these assumptions, we obtain cosmological equations

Cosmological equations

- ▶ Applying these assumptions, we obtain cosmological equations
 - ▶ Friedmann equation

$$\frac{3}{\kappa^2} H^2 = \rho_\phi = \frac{1}{2} \varepsilon \dot{\phi}^2 + U + 3\varepsilon \xi H^2 \phi^2 + 6\varepsilon \xi H \phi \dot{\phi},$$

Cosmological equations

- ▶ Applying these assumptions, we obtain cosmological equations
 - ▶ Friedmann equation

$$\frac{3}{\kappa^2} H^2 = \rho_\phi = \frac{1}{2} \varepsilon \dot{\phi}^2 + U + 3\varepsilon \xi H^2 \phi^2 + 6\varepsilon \xi H \phi \dot{\phi},$$

- ▶ acceleration equation

$$-\frac{1}{\kappa^2} (2\dot{H} + 3H^2) = p_\phi = \frac{\frac{1}{2} \varepsilon \dot{\phi}^2 (1 - 4\xi) + 6\varepsilon \xi^2 \phi^2 H^2 + 2\varepsilon \xi H \phi \dot{\phi} - U + 2\xi \phi U_{,\phi}}{1 - \varepsilon \kappa^2 \xi \phi^2 (1 - 6\xi)},$$

Cosmological equations

- ▶ Applying these assumptions, we obtain cosmological equations
 - ▶ Friedmann equation

$$\frac{3}{\kappa^2} H^2 = \rho_\phi = \frac{1}{2} \varepsilon \dot{\phi}^2 + U + 3\varepsilon \xi H^2 \phi^2 + 6\varepsilon \xi H \phi \dot{\phi},$$

- ▶ acceleration equation

$$-\frac{1}{\kappa^2} (2\dot{H} + 3H^2) = p_\phi = \frac{\frac{1}{2} \varepsilon \dot{\phi}^2 (1 - 4\xi) + 6\varepsilon \xi^2 \phi^2 H^2 + 2\varepsilon \xi H \phi \dot{\phi} - U + 2\xi \phi U_{,\phi}}{1 - \varepsilon \kappa^2 \xi \phi^2 (1 - 6\xi)},$$

- ▶ Klein-Gordon equation

$$\ddot{\phi} + 3H\dot{\phi} + 6\xi\phi \left(\dot{H} + 2H^2 \right) + \varepsilon U_{,\phi} = 0,$$

Cosmological equations

- ▶ Applying these assumptions, we obtain cosmological equations
 - ▶ Friedmann equation

$$\frac{3}{\kappa^2} H^2 = \rho_\phi = \frac{1}{2} \varepsilon \dot{\phi}^2 + U + 3\varepsilon \xi H^2 \phi^2 + 6\varepsilon \xi H \phi \dot{\phi},$$

- ▶ acceleration equation

$$-\frac{1}{\kappa^2} (2\dot{H} + 3H^2) = p_\phi = \frac{\frac{1}{2} \varepsilon \dot{\phi}^2 (1 - 4\xi) + 6\varepsilon \xi^2 \phi^2 H^2 + 2\varepsilon \xi H \phi \dot{\phi} - U + 2\xi \phi U_{,\phi}}{1 - \varepsilon \kappa^2 \xi \phi^2 (1 - 6\xi)},$$

- ▶ Klein-Gordon equation

$$\ddot{\phi} + 3H\dot{\phi} + 6\xi\phi \left(\dot{H} + 2H^2 \right) + \varepsilon U_{,\phi} = 0,$$

- ▶ where $U_{,\phi} := dU/d\phi$.

Dynamical equations

Dynamical equations

- ▶ Let us introduce dimensionless phase space variables

$$u = \frac{\dot{\phi}}{H\phi}, \quad v = \frac{\sqrt{6}}{\kappa} \frac{1}{\phi}.$$

Dynamical equations

- ▶ Let us introduce dimensionless phase space variables

$$u = \frac{\dot{\phi}}{H\phi}, \quad v = \frac{\sqrt{6}}{\kappa} \frac{1}{\phi}.$$

- ▶ In new variables, cosmological equations produce the **dynamical system**

$$\begin{aligned} u' = & \left[-\frac{1}{2}u^2(2+n) - \frac{3}{2}u(1+4\xi n) + \frac{1}{2}\epsilon n v^2 - 3\xi(1+n) \right] \left[\frac{1}{3}\epsilon v^2 - 2\xi(1-6\xi) \right]^2 + \\ & + (6\xi + u) \left[\frac{1}{3}\epsilon v^2 - 2\xi(1-6\xi) \right] \cdot \\ & \cdot \left(u^2 [1 - \xi(2-n)] + 4\xi u(2 + 3\xi n) - \frac{1}{2}\epsilon v^2(1 + 2\xi n) + 3\xi [1 + 2\xi(1+n)] \right), \\ v' = & -uv \left[\frac{1}{3}\epsilon v^2 - 2\xi(1-6\xi) \right]^2, \end{aligned}$$

where $f' = \frac{df}{d \ln a} = H^{-1}\dot{f}$.

Investigation of the dynamics of the system

Investigation of the dynamics of the system

- ▶ We start with finding **equilibria**, i.e. points, in which $u' = 0$ and $v' = 0$ (no evolution).

Investigation of the dynamics of the system

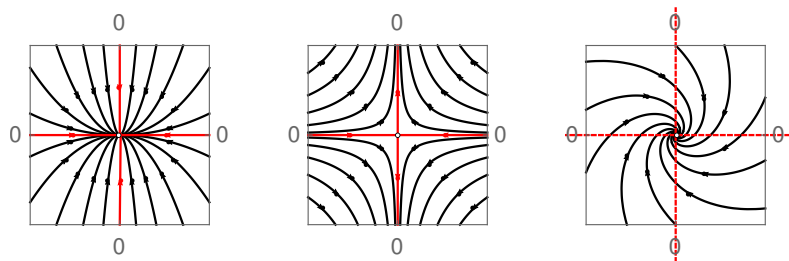
- ▶ We start with finding **equilibria**, i.e. points, in which $u' = 0$ and $v' = 0$ (no evolution).
- ▶ There are **six equilibria** ($A-F$) in finite space.

Investigation of the dynamics of the system

- ▶ We start with finding **equilibria**, i.e. points, in which $u' = 0$ and $v' = 0$ (no evolution).
- ▶ There are **six equilibria** ($A-F$) in finite space.
- ▶ Coordinates and stability features of these points **depend on model parameters ε , ξ and n** .

Investigation of the dynamics of the system

- ▶ We start with finding **equilibria**, i.e. points, in which $u' = 0$ and $v' = 0$ (no evolution).
- ▶ There are **six equilibria** ($A-F$) in finite space.
- ▶ Coordinates and stability features of these points **depend on model parameters** ε , ξ and n .



Bifurcation theory

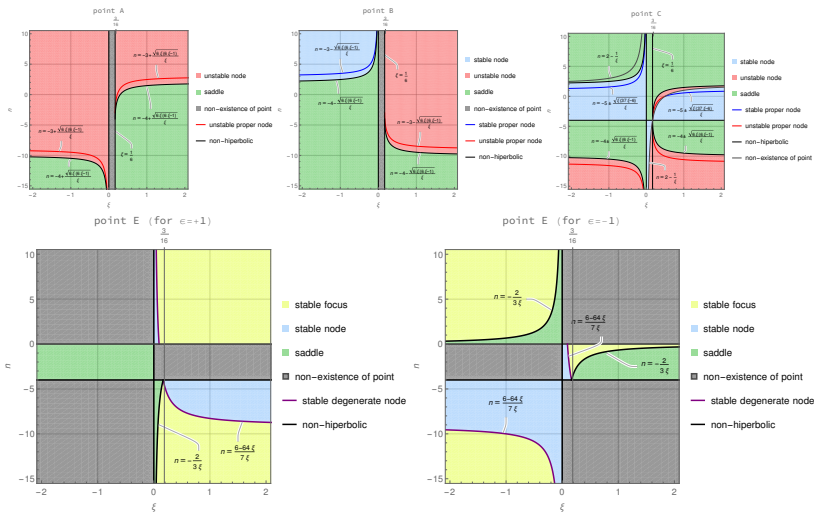
- ▶ As said before: Stability features of equilibria depend on parameters of the model...

Bifurcation theory

- ▶ As said before: Stability features of equilibria depend on parameters of the model...
- ▶ ...this leads to the methods of bifurcation theory.

- ▶ As said before: Stability features of equilibria depend on parameters of the model...
- ▶ ...this leads to the methods of bifurcation theory.
- ▶ **Definition:** The appearance of topologically nonequivalent phase portrait under variation of parameters is called a **bifurcation**.

Bifurcation diagrams of local stability in equilibria



What is the de Sitter universe?

What is the de Sitter universe?

- ▶ In the **de Sitter universe** the dynamics is dominated by the cosmological constant Λ , so the matter component (both baryonic and dark) is neglected.

What is the de Sitter universe?

- ▶ In the **de Sitter universe** the dynamics is dominated by the cosmological constant Λ , so the matter component (both baryonic and dark) is neglected.
- ▶ Pressure and energy density satisfy

$$p_{dS} = -\rho_{dS} = -\frac{\Lambda}{\kappa^2} \implies w_{dS} = \frac{p_{dS}}{\rho_{dS}} = -1.$$

What is the de Sitter universe?

- ▶ In the **de Sitter universe** the dynamics is dominated by the cosmological constant Λ , so the matter component (both baryonic and dark) is neglected.
- ▶ Pressure and energy density satisfy

$$p_{dS} = -\rho_{dS} = -\frac{\Lambda}{\kappa^2} \implies w_{dS} = \frac{p_{dS}}{\rho_{dS}} = -1.$$

- ▶ For the spatially flat universe ($k = 0$) the scale factor a depends on time as

$$a(t) \propto e^{\pm\sqrt{\frac{\Lambda}{3}}t}.$$

What is the de Sitter universe?

- ▶ In the **de Sitter universe** the dynamics is dominated by the cosmological constant Λ , so the matter component (both baryonic and dark) is neglected.
- ▶ Pressure and energy density satisfy

$$p_{dS} = -\rho_{dS} = -\frac{\Lambda}{\kappa^2} \implies w_{dS} = \frac{p_{dS}}{\rho_{dS}} = -1.$$

- ▶ For the spatially flat universe ($k = 0$) the scale factor a depends on time as

$$a(t) \propto e^{\pm\sqrt{\frac{\Lambda}{3}}t}.$$

- ▶ For '+' sign it is **de Sitter state**, while for '-' it is **anti de Sitter state**.

What is the de Sitter universe?

- ▶ In the **de Sitter universe** the dynamics is dominated by the cosmological constant Λ , so the matter component (both baryonic and dark) is neglected.
- ▶ Pressure and energy density satisfy

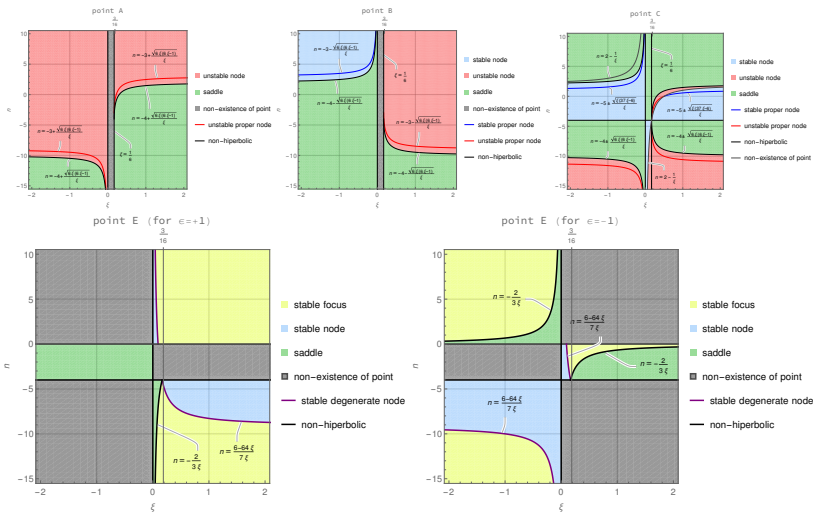
$$p_{dS} = -\rho_{dS} = -\frac{\Lambda}{\kappa^2} \implies w_{dS} = \frac{p_{dS}}{\rho_{dS}} = -1.$$

- ▶ For the spatially flat universe ($k = 0$) the scale factor a depends on time as

$$a(t) \propto e^{\pm\sqrt{\frac{\Lambda}{3}}t}.$$

- ▶ For '+' sign it is **de Sitter state**, while for '-' it is **anti de Sitter state**.
- ▶ **De Sitter state describes dynamics of cosmic inflation.**

Bifurcation diagrams of local stability in equilibria



Cases of de Sitter–de Sitter evolution of the universe

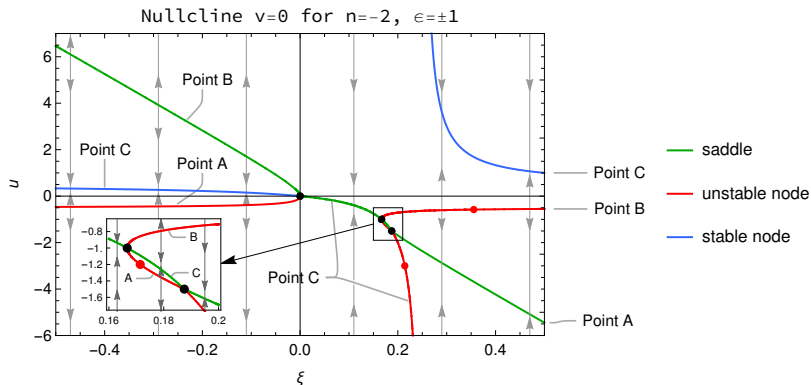
- ▶ Looking at bifurcations diagrams, we could extract ranges of parameters for which evolution of the universe **starts in de-Sitter state (inflation) and finishes in another de Sitter state (late time acceleration after the inclusion of matter)**.

Table: Sets of parameters for which the universe undergoes the evolution starting from the de Sitter state and finishing in the de Sitter state.

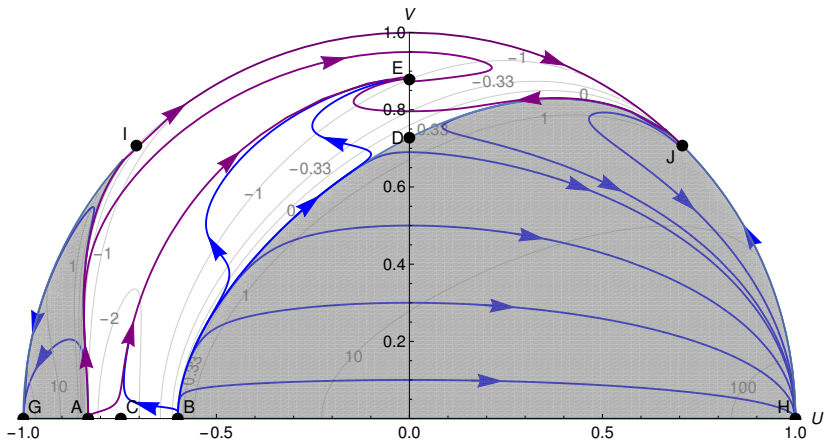
no.	ξ	n	ε	starting point	final point
1. Generic de Sitter–de Sitter evolution					
(a)	$\frac{3}{16}$	$(0, +\infty)$	+1	unstable node <i>A</i>	stable focus <i>E</i>
(b)	$\frac{3}{16}$	$(-2, 0)$	-1	unstable node <i>A</i>	stable focus <i>E</i>
(c)	$(\frac{3}{16}, \frac{1}{4})$	-2	-1	unstable node <i>C</i>	stable focus <i>E</i>
2. Non-generic de Sitter–de Sitter evolution					
(d)	$(-\infty, 0]$	-2	+1	saddle <i>E</i>	stable node <i>C</i>
(e)	$\frac{3}{16}$	$(-3\frac{5}{9}, -2)$	-1	saddle <i>A</i>	stable focus <i>E</i>
(f)	$\frac{3}{16}$	$[-4, -3\frac{5}{9})$	-1	saddle <i>A</i>	saddle <i>E</i>
(g)	$[0, \frac{3}{16})$	-2	-1	saddle <i>C</i>	stable node/focus <i>E</i>
(h)	$(\frac{1}{3}, +\infty)$	-2	-1	saddle <i>E</i>	stable node <i>C</i>

'Collisions' of equilibria for generic dS-dS scenarios

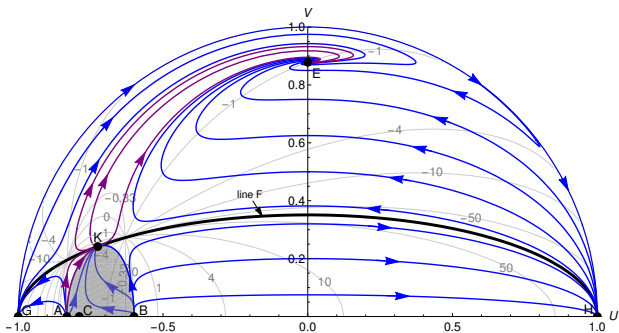
- ▶ Again, bifurcation diagrams (showing 'collisions' of equilibria) indicated that one phase portrait is fully representative for each of scenarios (a)–(c).



Phase portrait for scenario (a)

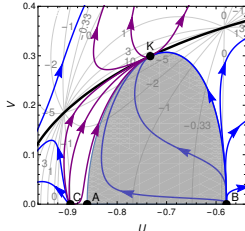
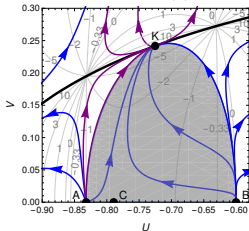


Phase portrait for scenarios (b) and (c)

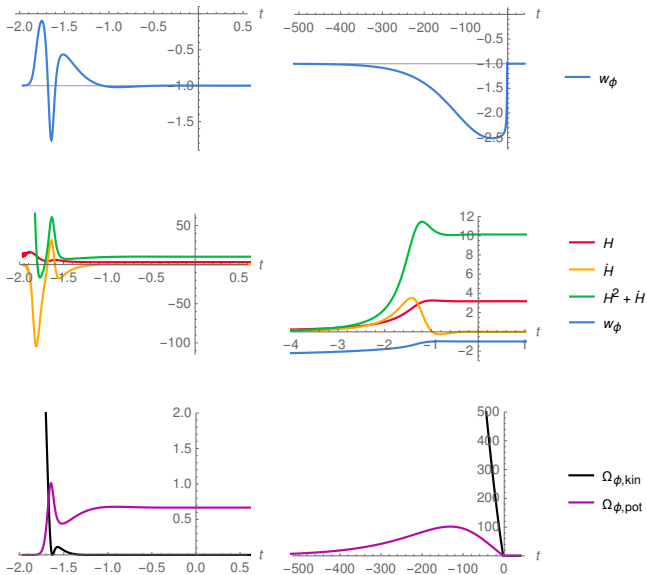


scenario (b)

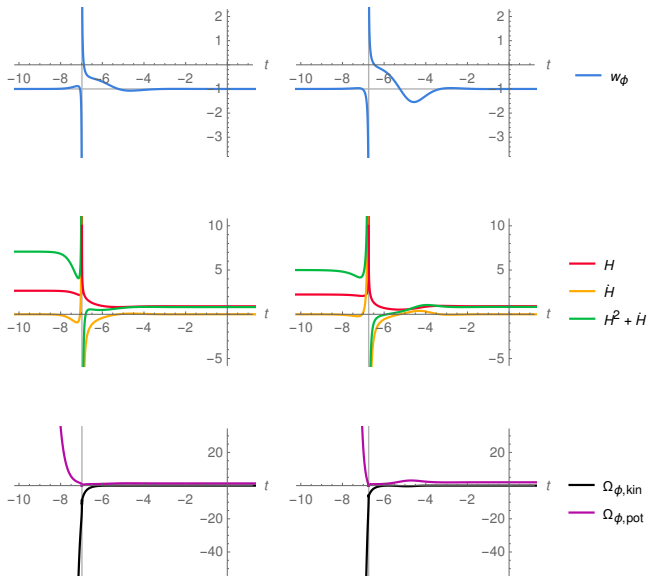
scenario (c)



Evolution of physical quantities for case (a)



Evolution of physical quantities for cases (b) and (c)



Conclusions

Conclusions

- ▶ We considered cosmology with scalar field coupled to the gravity; this approach **enabled us to include both the inflation and the late-time acceleration in a single model.**

Conclusions

- ▶ We considered cosmology with scalar field coupled to the gravity; this approach **enabled us to include both the inflation and the late-time acceleration in a single model.**
- ▶ The application of bifurcation theory allowed us to distinguish sets of parameters for which **the universe undergoes a generic evolution without presence of the initial singularity.**

Conclusions

- ▶ We considered cosmology with scalar field coupled to the gravity; this approach **enabled us to include both the inflation and the late-time acceleration in a single model.**
- ▶ The application of bifurcation theory allowed us to distinguish sets of parameters for which **the universe undergoes a generic evolution without presence of the initial singularity.**
- ▶ There occurred two types of non-singular initial states: **the de Sitter state and the static universe.**

Conclusions

- ▶ We considered cosmology with scalar field coupled to the gravity; this approach **enabled us to include both the inflation and the late-time acceleration in a single model.**
- ▶ The application of bifurcation theory allowed us to distinguish sets of parameters for which **the universe undergoes a generic evolution without presence of the initial singularity.**
- ▶ There occurred two types of non-singular initial states: **the de Sitter state and the static universe.**
- ▶ From the bifurcation analysis, we obtained pairs of the critical values of the parameters (ξ, n) which corresponded to bifurcation values.

Conclusions

- ▶ We considered cosmology with scalar field coupled to the gravity; this approach **enabled us to include both the inflation and the late-time acceleration in a single model.**
- ▶ The application of bifurcation theory allowed us to distinguish sets of parameters for which **the universe undergoes a generic evolution without presence of the initial singularity.**
- ▶ There occurred two types of non-singular initial states: **the de Sitter state** and **the static universe.**
- ▶ From the bifurcation analysis, we obtained pairs of the critical values of the parameters (ξ, n) which corresponded to bifurcation values.
- ▶ *THANK YOU FOR YOUR ATTENTION!*