

Many-body localization due to random interactions

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- 1 Many-body localization in optical lattices
- 2 Bosons in an optical lattice with disorder
- 3 MBL due to random interactions
 - Ergodicity breaking
 - Spectral statistics
 - Mobility edge
- 4 Random on-site potential

MBL in optical lattice

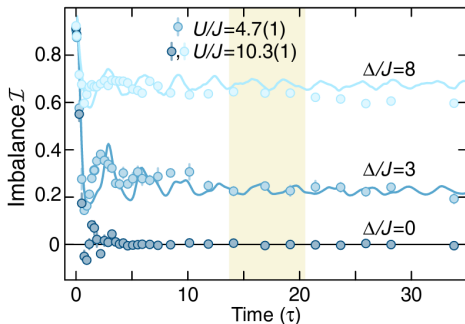
D. Basko, et. al., Ann. Phys. (NY) 321, 1126 (2006) – theoretical prediction of MBL

M. Schreiber et. al., Science 349, 842 (2015):

Interacting fermions in quasi-random optical lattice:

$$H = -J \sum_{i,\sigma} \left(c_{i,\sigma}^\dagger c_{i+1,\sigma} + h.c. \right) + \Delta \sum_{i,\sigma} \cos(2\pi\beta i + \phi) c_{i,\sigma}^\dagger c_{i,\sigma} + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

Time evolution of initial charge-density wave state:



- non-zero stationary imbalance – ergodicity breaking
- lack of thermalization – ETH is false
- the system is many-body localized
- fully localized single-particle spectrum

Bosons in optical lattice

Consider the Bose–Hubard model

$$H = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{U}{2} \sum_i \hat{n}_i(\hat{n}_i - 1) + \mu \sum_i \hat{n}_i$$

- Optical speckle field \rightarrow random on-site potential

$$H = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{U}{2} \sum_i \hat{n}_i(\hat{n}_i - 1) + \sum_i \mu_i \hat{n}_i$$

$$\mu_i \in [-W/2, W/2]$$

- Atom chip providing spatially random magnetic field \rightarrow random interactions
H. Gimpel et. al., Phys. Rev. Lett. **95**, 170401 (2005)

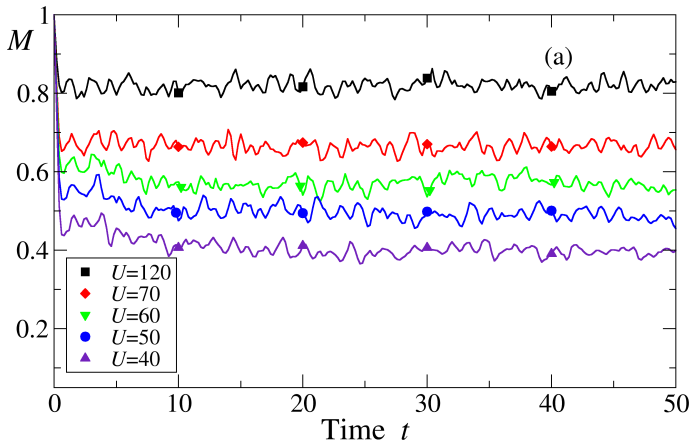
$$H = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \sum_i \frac{U_i}{2} \hat{n}_i(\hat{n}_i - 1) + \sum_i \mu \hat{n}_i$$

$$U_i \in [0, U], \text{ delocalized single-particle spectrum}$$

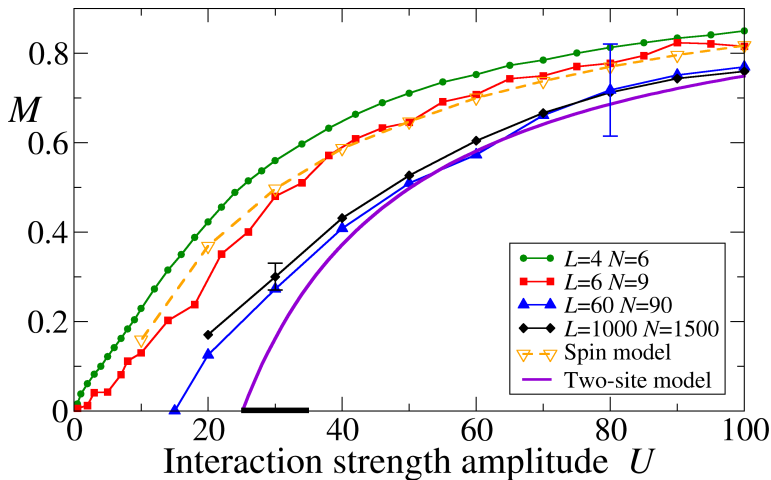
Magnetization

time evolution of the density wave state $|DW\rangle = |2121\dots\rangle$

define: $M = 3(N_e - N_o)/(N_e + N_o)$



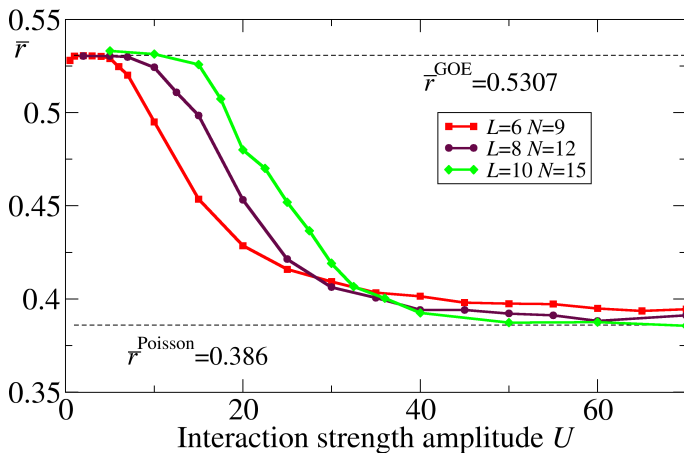
Stationary value of magnetization



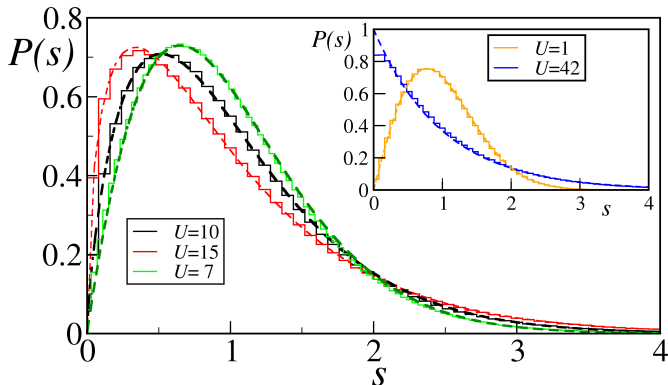
Spectral statistics - \bar{r}

Average ratio of adjacent energy gaps:

$$r_n = \min[\delta_n^E, \delta_{n-1}^E] / \max[\delta_n^E, \delta_{n-1}^E] \text{ with } \delta_n^E = E_n - E_{n-1}$$

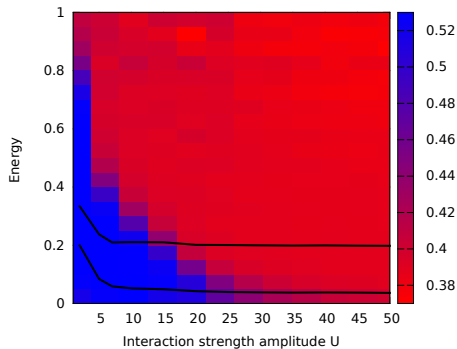
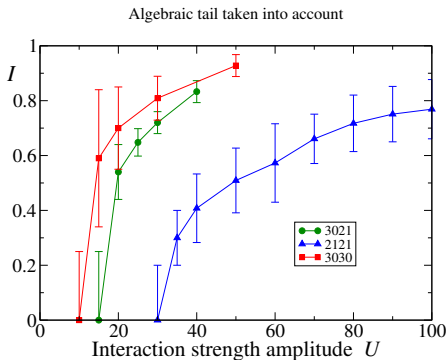


Level spacing distribution



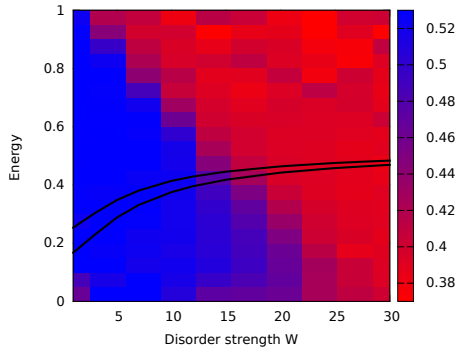
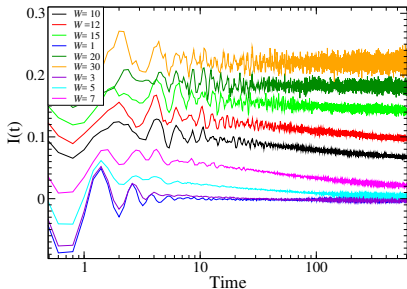
- $U = 1$: $P(s) \propto s e^{-\frac{\pi}{4}s^2}$ – GOE statistics
- $U = 7$: $P(s) \propto s^\beta e^{-C_2 s^{2-\gamma}}$ – M. Serbyn, J.E. Moore, PRB 93, 041424 (2016)
- $U = 10, 15$: $P(s) \propto s^\beta e^{-(\beta+1)s}$ – Semi-Poisson
- $U = 42$: $P(s) \propto e^{-s}$ – Poisson statistics

Localization transition depends strongly on the energy density

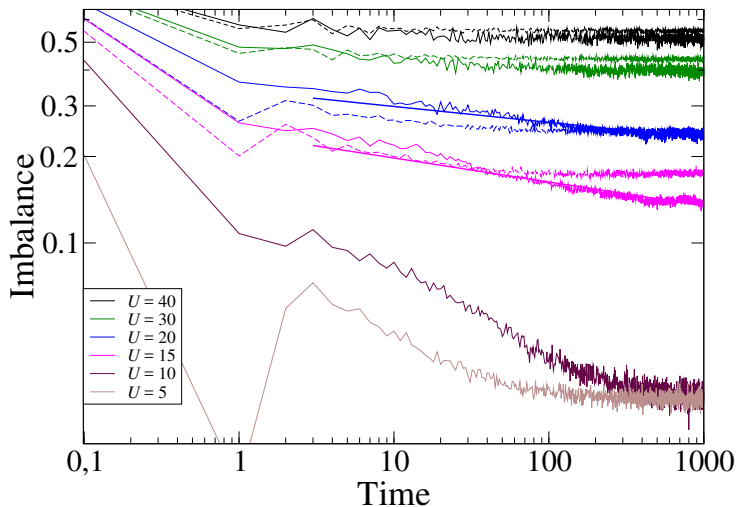


Random chemical potential

$$H = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i \hat{n}_i$$



$L = 8 N = 12$ vs $L = 6 N = 9$

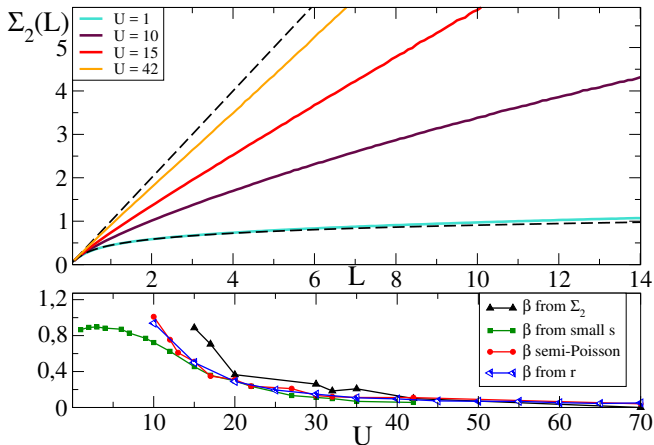


- random interactions - MBL occurs solely due to the interactions in the system with extended single-particle states
- the models are experimentally accessible
- bosons allow for much larger density variations than fermions
- more details about random interactions case in PRA **95**, 021601(R), 2017

Number variance

$\mathcal{N}(L)$ - number of energy levels in interval $(E, E + L)$,
 $\text{Var}[\mathcal{N}(L)]$ fulfills

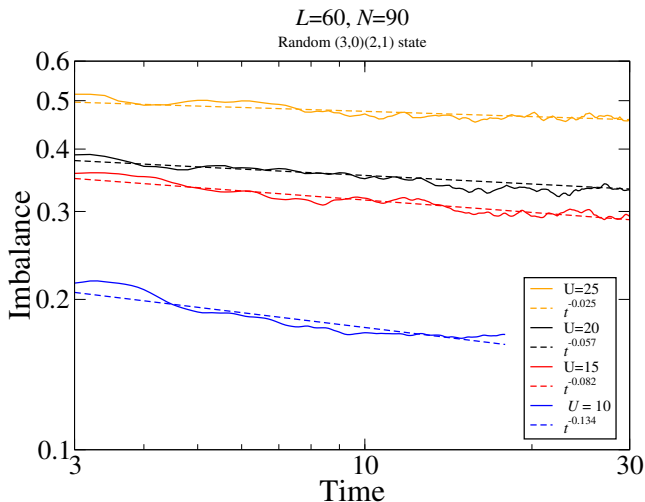
$$\Sigma_2(L) = L - 2 \int_0^L dr (L-r) Y_2(r)$$



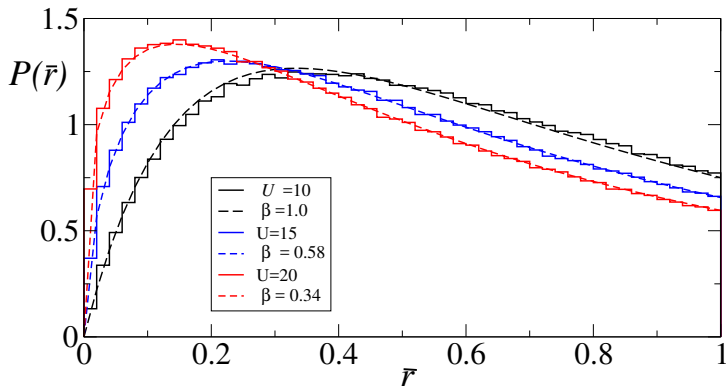
Algebraic decay of magnetization

Close to transition – Griffiths regime

$$M(t) \propto t^{-1/z} \quad (1)$$



\bar{r} distribution in the crossover region fitted with distribution for Semi-Poisson statistics from Y.Atas et. al., J. Phys. A: Math. Theor. 46, 355204 (2013)



Approximate models in MBL phase

- 1 low-energy subspace with $\hat{n}_i = 1, 2$; the XX Heisenberg spin chain in random magnetic field

$$H = J \sum_i (S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+) + \sum_i h_i S_i^z,$$

\Leftrightarrow Anderson localization non-interacting fermions, physically different – entanglement entropy saturates

- 2 two-site model: in basis $|21\rangle, |12\rangle$, the Hamiltonian reads

$$H = \begin{bmatrix} U_1 & -2J \\ -2J & U_2 \end{bmatrix}, \Rightarrow$$

diagonalize, average over time oscillations and disorder, consider transfer of $n_i = 2$ into both directions to get $M = 1 - 8\pi J/U$.

Stationary value of magnetization

