Phases of one-dimensional Bose-Hubbard model in optical lattices with dipole-dipole interactions

K. Biedroń¹ J. Zakrzewski^{1,2}

¹Instytut Fizyki im. M. Smoluchowskiego Uniwersytet Jagielloński

²Mark Kac Complex Systems Research Center Jagiellonian University

Cracow School of Theoretical Physics, LVII Course, 2017 Entanglement and Dynamics

ヘロト ヘワト ヘビト ヘビト

Outline



Motivation

- Bose-Hubbard models
- EBH in optical lattices theory
- 2 Generalized BHM
 - The model
 - Parameter values
- Phases in one dimension
 - Phase diagram for $\rho = 1$
 - Phase diagram for fixed V/U

(七日)) (日日)

Bose-Hubbard models EBH in optical lattices - theory

Bose-Hubbard Model

Basic model of bosons in an optical lattice:

$$H = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) \quad (+ \sum_i \epsilon_i b_i^{\dagger} b_i)$$

Here, U is contact interaction between two bosons at the same lattice site, which we can parametrize using effective scattering length (a_s)

Two phases in limiting cases (for ho = 1)

- $J \gg U$: superfluid (SF), delocalized wavefunctions
- $J \ll U$: Mott insulator (MI), localized wavefunctions

Experimentally realized in an optical lattice in 2002¹:



Absorption images at time of flight at 15ms, pictures a-h correspond to V_0/E_B : 0, 3, 7, 10, 13, 14, 16 and 20

¹Markus Greiner et al. "Quantum phase transition from a superfluid to a Mott insulator in a gas of ultracold atoms". In: nature 415.6867 (2002), pp. 39–44.

K. Biedroń, J. Zakrzewski

Phases of one-dimensional Bose-Hubbard model in optical lattice

æ

Bose-Hubbard models EBH in optical lattices - theory

Bose-Hubbard Model

Basic model of bosons in an optical lattice:

$$H = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) \quad (+ \sum_i \epsilon_i b_i^{\dagger} b_i)$$

Here, U is contact interaction between two bosons at the same lattice site, which we can parametrize using effective scattering length (a_s)

Two phases in limiting cases (for $\rho = 1$)

- $J \gg U$: superfluid (SF), delocalized wavefunctions
- J « U: Mott insulator (MI), localized wavefunctions

Experimentally realized in an optical lattice in 2002¹:



Absorption images at time of flight at 15ms, pictures a-h correspond to V_0/E_B : 0, 3, 7, 10, 13, 14, 16 and 20

¹Markus Greiner et al. "Quantum phase transition from a superfluid to a Mott insulator in a gas of ultracold atoms". In: nature 415.6867 (2002), pp. 39–44.

K. Biedroń, J. Zakrzewski

Phases of one-dimensional Bose-Hubbard model in optical lattice

3

Bose-Hubbard models EBH in optical lattices - theory

Bose-Hubbard Model

Basic model of bosons in an optical lattice:

$$H = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) \quad (+ \sum_i \epsilon_i b_i^{\dagger} b_i)$$

Here, U is contact interaction between two bosons at the same lattice site, which we can parametrize using effective scattering length (a_s)

Two phases in limiting cases (for $\rho = 1$)

- $J \gg U$: superfluid (SF), delocalized wavefunctions
- J « U: Mott insulator (MI), localized wavefunctions

Experimentally realized in an optical lattice in 20021:



Absorption images at time of flight at 15ms, pictures a-h correspond to V_0/E_R : 0, 3, 7, 10, 13, 14, 16 and 20

ъ

¹Markus Greiner et al. "Quantum phase transition from a superfluid to a Mott insulator in a gas of ultracold atoms". In: nature 415.6867 (2002), pp. 39–44.

K. Biedroń, J. Zakrzewski Phases of one-dimensional Bose-Hubbard model in optical lattice

Bose-Hubbard models EBH in optical lattices - theory

Extended Bose-Hubbard Model

Addition of nearest-neighbour tunneling to previous hamiltonian:

$$H = -t \sum_{\langle i,j \rangle} b_i^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j$$

V is mostly due to dipole-dipole interaction of bosonic particles:

$$V_{dip}(r - r') = \frac{C_{dd}}{4\pi} \frac{1 - 3\cos^2{\theta}}{|r - r'|^3}$$

For particles having magnetic dipole moment $C_{dd} = \mu_0 \mu^2$, and for particles having electric dipole moment $C_{dd} = d^2 / \varepsilon_0$.

Additional phases with this extension to BHM

density wave (DW) phase: $U \gg V$, $U \gg t$: "checkerboard" density distribution

Haldane insulator: finite non-local string correlations: $C_{str}(r) = \langle \delta n_j e^{i\pi \sum_{j < k < j+r} \delta n_k} \delta n_{j+r} \rangle$.

イロト 不得 とくほ とくほとう

3

Bose-Hubbard models EBH in optical lattices - theory

Extended Bose-Hubbard Model

Addition of nearest-neighbour tunneling to previous hamiltonian:

$$H = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j$$

V is mostly due to dipole-dipole interaction of bosonic particles:

$$V_{dip}(\boldsymbol{r}-\boldsymbol{r}') = \frac{C_{dd}}{4\pi} \frac{1-3\cos^2\theta}{|\boldsymbol{r}-\boldsymbol{r}'|^3}$$

For particles having magnetic dipole moment $C_{dd} = \mu_0 \mu^2$, and for particles having electric dipole moment $C_{dd} = d^2 / \varepsilon_0$.

Additional phases with this extension to BHM

density wave (DW) phase: $U \gg V$, $U \gg t$: "checkerboard" density distribution

Haldane insulator: finite non-local string correlations: $C_{str}(r) = \langle \delta n_j e^{i\pi \sum_{j < k < j+r} \delta n_k} \delta n_{j+r} \rangle$.

イロト 不得 とくほと くほとう

∃ <2 <</p>

Bose-Hubbard models EBH in optical lattices - theory

Extended Bose-Hubbard Model

Addition of nearest-neighbour tunneling to previous hamiltonian:

$$H = -t \sum_{\langle i,j \rangle} b_i^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j$$

V is mostly due to dipole-dipole interaction of bosonic particles:

$$V_{dip}(\boldsymbol{r}-\boldsymbol{r}') = \frac{C_{dd}}{4\pi} \frac{1-3\cos^2\theta}{|\boldsymbol{r}-\boldsymbol{r}'|^3}$$

For particles having magnetic dipole moment $C_{dd} = \mu_0 \mu^2$, and for particles having electric dipole moment $C_{dd} = d^2 / \varepsilon_0$.

Additional phases with this extension to BHM

• density wave (DW) phase: $U \gg V$, $U \gg t$: "checkerboard" density distribution

• Haldane insulator: finite non-local string correlations: $C_{str}(r) = \langle \delta n_j e^{i\pi \sum_{j < k < j+r} \delta n_k} \delta n_{j+r} \rangle$.

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ● □ ● ● ● ●

Bose-Hubbard models EBH in optical lattices - theory

EBH phase diagram



The phase diagram of EBH for filling ρ = 1 has been calculated using density matrix renormalization group (DMRG) ².

²Davide Rossini and Rosario Fazio. "Phase diagram of the extended Bose-Hubbard model". In: New Journal of Physics 14.6 (2012), p. 065012

Bose-Hubbard models EBH in optical lattices - theory

EBH phase diagram



The phase diagram of EBH for filling ρ = 1 has been calculated using density matrix renormalization group (DMRG) ².

Different phases for V/U = 3/4 and varying filling (here measured by chemical potential μ) - calculated using DMRG (black lines) and Quantum Monte Carlo (symbols)³.

²Davide Rossini and Rosario Fazio. "Phase diagram of the extended Bose–Hubbard model". In: New Journal of Physics 14.6 (2012), p. 065012

³GG Batrouni et al. "Competing phases, phase separation, and coexistence in the extended one-dimensional bosonic Hubbard model". In: *Physical Review B* 90.20 (2014), p. 205123

The model Parameter values

The model

We can extend our model even more by including next highest terms in our hamiltonian:⁴

$$H_{GBH} = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j + V_{nnn} \sum_{\langle i,j \rangle} n_i n_j - T \sum_{\langle i,j \rangle} b_i^{\dagger} (n_i + n_j) b_j + \frac{P}{2} b_i^{\dagger 2} b_j^{\dagger 2} b_j^{$$

Additional terms:

- V_{nnn} next-nearest neighbour interaction
- T density dependent tunneling
- P pair tunneling

Describing states of the bosons using Wannier functions ($w_i(r)$ describing state at *i*-th site), we can express above terms as:

$$t = -\int dr w_j^*(r) \left[\frac{-\hbar^2 \nabla^2}{2m} \right] w_j(r) \tag{1}$$

$$U_{ijkl} = g \int d^3 r_1 d^3 r_2 w_i^*(r_1) w_j^*(r_2) U(r_1 - r_2) w_k(r_1) w_l(r_1)$$
(2)

$$U = U_{1111}, V = U_{1212}, V_{nnn} = U_{1313}, T = -U_{1112}, P = U_{1122}$$
 (3)

The model Parameter values

The model

We can extend our model even more by including next highest terms in our hamiltonian:⁴

$$H_{GBH} = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j + V_{nnn} \sum_{\langle i,j \rangle} n_i n_j - T \sum_{\langle i,j \rangle} b_i^{\dagger} (n_i + n_j) b_j + \frac{P}{2} b_i^{\dagger 2} b_j^{\dagger 2} b_j^{$$

Additional terms:

- V_{nnn} next-nearest neighbour interaction
 - T density dependent tunneling
 - P pair tunneling

Describing states of the bosons using Wannier functions ($w_i(r)$ describing state at *i*-th site), we can express above terms as:

$$t = -\int dr w_i^*(r) \left[\frac{-h^2 \nabla^2}{2m} \right] w_j(r) \tag{1}$$

$$U_{ijkl} = g \int d^3 r_1 d^3 r_2 w_i^*(r_1) w_j^*(r_2) U(r_1 - r_2) w_k(r_1) w_l(r_1)$$
(2)

$$U = U_{1111}, V = U_{1212}, V_{nnn} = U_{1313}, T = -U_{1112}, P = U_{1122}$$
 (3)

The model Parameter values

The model

We can extend our model even more by including next highest terms in our hamiltonian:⁴

$$H_{GBH} = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j + V_{nnn} \sum_{\langle i,j \rangle} n_i n_j - T \sum_{\langle i,j \rangle} b_i^{\dagger} (n_i + n_j) b_j + \frac{P}{2} b_i^{\dagger 2} b_j^{\dagger 2} b_j^{$$

Additional terms:

- V_{nnn} next-nearest neighbour interaction
- T density dependent tunneling
 - P pair tunneling

Describing states of the bosons using Wannier functions ($w_i(r)$ describing state at *i*-th site), we can express above terms as:

$$t = -\int dr w_j^*(r) \left[\frac{-\hbar^2 \nabla^2}{2m} \right] w_j(r) \tag{1}$$

$$U_{ijkl} = g \int d^3 r_1 d^3 r_2 w_i^*(r_1) w_j^*(r_2) U(r_1 - r_2) w_k(r_1) w_l(r_1)$$
(2)

 $U = U_{1111}, V = U_{1212}, V_{nnn} = U_{1313}, T = -U_{1112}, P = U_{1122}$ (3)

The model Parameter values

The model

We can extend our model even more by including next highest terms in our hamiltonian:⁴

$$H_{GBH} = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j + V_{nnn} \sum_{\langle i,j \rangle} n_i n_j - T \sum_{\langle i,j \rangle} b_i^{\dagger} (n_i + n_j) b_j + \frac{P}{2} b_i^{\dagger 2} b_j^{\dagger 2} b_j^{$$

Additional terms:

- V_{nnn} next-nearest neighbour interaction
- T density dependent tunneling
- P pair tunneling

Describing states of the bosons using Wannier functions ($w_i(r)$ describing state at *i*-th site), we can express above terms as:

$$t = -\int dr w_i^*(r) \left[\frac{-h^2 \nabla^2}{2m} \right] w_j(r) \tag{1}$$

$$U_{ijkl} = g \int d^3 r_1 d^3 r_2 w_i^*(r_1) w_j^*(r_2) U(r_1 - r_2) w_k(r_1) w_l(r_1)$$
(2)

 $U = U_{1111}, V = U_{1212}, V_{nnn} = U_{1313}, T = -U_{1112}, P = U_{1122}$ (3)

The model Parameter values

The model

We can extend our model even more by including next highest terms in our hamiltonian:⁴

$$H_{GBH} = -t \sum_{\langle i,j \rangle} b_j^{\dagger} b_j + \frac{U}{2} \sum_i n_i (n_i - 1) + V \sum_{\langle i,j \rangle} n_i n_j + V_{nnn} \sum_{\langle i,j \rangle} n_i n_j - T \sum_{\langle i,j \rangle} b_i^{\dagger} (n_i + n_j) b_j + \frac{P}{2} b_i^{\dagger 2} b_j^{\dagger 2} b_j^{$$

Additional terms:

- V_{nnn} next-nearest neighbour interaction
- T density dependent tunneling
- P pair tunneling

Describing states of the bosons using Wannier functions ($w_i(r)$ describing state at *i*-th site), we can express above terms as:

$$t = -\int d\mathbf{r} w_i^*(\mathbf{r}) \left[\frac{-\hbar^2 \nabla^2}{2m} \right] w_j(\mathbf{r})$$
⁽¹⁾

$$U_{ijkl} = g \int d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 w_i^*(\mathbf{r}_1) w_j^*(\mathbf{r}_2) U(\mathbf{r}_1 - \mathbf{r}_2) w_k(\mathbf{r}_1) w_l(\mathbf{r}_1)$$
(2)

$$U = U_{1111}, \quad V = U_{1212}, \quad V_{nnn} = U_{1313}, \quad T = -U_{1112}, \quad P = U_{1122} \tag{3}$$

The model Parameter values

Parameter values

For one-dimensional lattice (lattice potentials $V_y = V_z = 50E_R$, $V_x = V_1$):



Total tunnelings and interactions are the sum of contact and dipole-dipole ones, $t = t_c + t_{dd}$, $U = U_c + U_{dd}$

イロト イポト イヨト イヨト

3

Parameter values

Engineering U/t and V/t

3 parameters we can adjust

- V₁: here I used range from 2E_B to 20E_B
- as: scattering length of contact interactions adjustable by Feshbach resonance
- $d = \frac{C_{dd}}{4\pi E_{D}}$, can be changed by adjusting the orientation of magnetic dipoles

(ロ) (同) (目) (日) (日) (の)

The model Parameter values

Engineering U/t and V/t

3 parameters we can adjust

- V₁: here I used range from 2E_R to 20E_R
- as: scattering length of contact interactions adjustable by Feshbach resonance
- $d = \frac{C_{dd}}{4\pi E_{D}}$, can be changed by adjusting the orientation of magnetic dipoles

(ロ) (同) (目) (日) (日) (の)

The model Parameter values

Engineering U/t and V/t

3 parameters we can adjust

- V₁: here I used range from 2E_R to 20E_R
- as: scattering length of contact interactions adjustable by Feshbach resonance
- $d = \frac{C_{dd}}{4\pi E_p}$, can be changed by adjusting the orientation of magnetic dipoles

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ ● □ ● ○ ○ ○

The model Parameter values

Engineering U/t and V/t

3 parameters we can adjust

- V₁: here I used range from 2E_R to 20E_R
- as: scattering length of contact interactions adjustable by Feshbach resonance
- $d = \frac{C_{dd}}{4\pi E_P}$, can be changed by adjusting the orientation of magnetic dipoles

For aforementioned lattice depths and set value of d, there is a specific range of V/t values we can get by modifying V_1 and a_s , which is marked here as a shaded region:



K. Biedroń, J. Zakrzewski

Phases of one-dimensional Bose-Hubbard model in optical lattice

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



★ E > ★ E

э

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

SF - MI transition:

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

MI - HI transition:

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

HI - DW transition:

Phase diagram for $\rho = 1$

For d = 0.25, $V2 = 50E_R$ and varying V_1 and a_s phase diagram was calculated with open boundary conditions DMRG for sizes up to 400 sites (using itensor library):



イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:



ヘロト ヘワト ヘビト ヘビト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:





Figure: Simple EBH phase diagram.

イロト イポト イヨト イヨト

э

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:



Figure: $\rho = 7/4$

イロト イポト イヨト イヨト

ъ

SF - SS (supersolid) and SF - DW transitions:

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:

Correlations in DW regime:



イロト イポト イヨト イヨト

э

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:



Correlations in SS regime:





Figure: Log-log plot of correlations, $\rho = 7/4$, t/U = 0.15

イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Phase diagram for fixed V/U

Here V/U = 0.75, $V_{nnn} = V/8$, T and P are neglected. For various lattice fillings (from $\rho = 1$ to $\rho = 2$) phase diagram was calculated with open boundary conditions DMRG with 192 sites:

Correlations in SF regime:



Figure: Log-log plot of correlations, $\rho = 7/4$, t/U = 0.25

ヘロト ヘワト ヘビト ヘビト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Summary

- Bose-Hubbard model can be generalized with additional terms that can be calculated straightforwardly using Wannier functions
- Including those additional terms do not qualitatively change phase diagram, proving HI phase robustness
- Including next-nearest neighbour interactions influences SS regime extent for higher lattice fillings and creates additional DW patterns

イロト イポト イヨト イヨト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Summary

- Bose-Hubbard model can be generalized with additional terms that can be calculated straightforwardly using Wannier functions
- Including those additional terms do not qualitatively change phase diagram, proving HI phase robustness
- Including next-nearest neighbour interactions influences SS regime extent for higher lattice fillings and creates additional DW patterns

イロト イポト イヨト イヨト

æ

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Summary

- Bose-Hubbard model can be generalized with additional terms that can be calculated straightforwardly using Wannier functions
- Including those additional terms do not qualitatively change phase diagram, proving HI phase robustness
- Including next-nearest neighbour interactions influences SS regime extent for higher lattice fillings and creates additional DW patterns

ヘロト ヘワト ヘビト ヘビト

Phase diagram for $\rho = 1$ Phase diagram for fixed V/U

Acknowledgements

Thank you for your attention!

This work was realized under National Science Centre (Poland) Project No. DEC-2012/04/A/ST2/00088 and was supported in part by PL-Grid Infrastructure.

イロト イポト イヨト イヨト