A story of chaos and purity

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Outline

Chaos

- Diagnosing chaos with out-of-time-order (OTO) correlators.
- Chaos at large central charge.
- OTO in $SU(N)_k$ WZW models.
- OTO for rational CFT, and fractional statistics.

Purity

- Purity for excited states.
- Purity scrambling for large-c $SU(N)_k$.

Based on: 1602.06542 with P. Caputa and T. Numasawa

Diagnosing chaos with OTO

An indicator of the buterfly effect

$$\left\langle \left[\mathcal{O}_{j}(0),\mathcal{O}_{i}(t)\right]^{2}\right\rangle _{\beta}\overset{?}{=}0$$

Larkin, Ovchinnikov '69, Kitaev '14 Maldacena, Shenker, Roberts, Stanford '15

Consider late times behaviour

$$egin{aligned} &-\left<\left[\mathcal{O}_{j},\mathcal{O}_{i}(t)
ight]^{2}
ight
angle_{eta} = \left<\mathcal{O}_{j}\mathcal{O}_{i}(t)\mathcal{O}_{i}(t)\mathcal{O}_{j}
ight
angle_{eta} + \left<\mathcal{O}_{i}(t)\mathcal{O}_{j}\mathcal{O}_{j}\mathcal{O}_{i}(t)
ight
angle_{eta} \ &-\left<\mathcal{O}_{j}\mathcal{O}_{i}(t)\mathcal{O}_{j}\mathcal{O}_{i}(t)
ight
angle_{eta} - \left<\mathcal{O}_{i}(t)\mathcal{O}_{j}\mathcal{O}_{i}(t)\mathcal{O}_{j}
ight
angle_{eta} \end{aligned}$$

In general, first line $\sim \langle \mathcal{O}_j \mathcal{O}_j
angle_eta \langle \mathcal{O}_i(t) \mathcal{O}_i(t)
angle_eta$

If the second line vanishes the system is chaotic

The correlators in the second line are out-of-time-order (OTO)

OTO in CFT_2

Chaos indicator
$$C_{ij}^{\beta}(t) \equiv \frac{\langle \mathcal{O}_i(t)\mathcal{O}_j\mathcal{O}_i(t)\mathcal{O}_j \rangle_{\beta}}{\langle \mathcal{O}_i\mathcal{O}_i \rangle_{\beta} \langle \mathcal{O}_j\mathcal{O}_j \rangle_{\beta}} = f(z, \bar{z})$$

Cross-ratio z



OTOs $\langle \mathcal{O}_i(t) \mathcal{O}_j \mathcal{O}_j \mathcal{O}_i(t) \rangle_{\beta}$ go through the branch cut

Large-*c* and chaos

We can write (holography)

$$f(z,ar{z})
ightarrow \mathcal{F}(z)ar{\mathcal{F}}(ar{z})$$

For example, in the (heavy-heavy \rightarrow light-light) regime

$$\mathcal{F}(z) \sim \left(rac{z(1-z)^{-6h_w/c}}{1-(1-z)^{1-12h_w/c}}
ight)^{2h_v}$$

Fitzpatrick, Kaplan, Walters '14

At late times the chaos indicator will vanish

$$C_{ij}^eta(t)\sim \left(1+rac{2\pi h_w}{\epsilon_{12}^*\epsilon_{34}}\exp\left(rac{2\pi}{eta}(t-t_*)
ight)
ight)^{-2h_v}$$

Roberts, Stanford '15

Fast scrambling time $t_* = \frac{\beta}{2\pi} \log(c)$

Large-c and chaos

We can write (holography)

$$f(z,ar{z}) o \mathcal{F}(z)ar{\mathcal{F}}(ar{z})$$

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Fitzpatrick, Kaplan, Walters '14

Holographic theories display the butterfly effect

$$\mathcal{C}_{ij}^eta(t)\sim \left(1+rac{2\pi h_w}{\epsilon_{12}^*\epsilon_{34}}\exp\left(rac{2\pi}{eta}(t-t_*)
ight)
ight)^{-2h_v}$$

Roberts, Stanford '15

Matches bulk analysis (Shenker, Stanford '14)

WZW theories

CFT's that besides Virasoro they have an affine Lie algebra $\hat{\mathfrak{g}}_k$

$$\left[j_m^a, j_n^b\right] = i \sum_c f^{abc} j_{m+n}^c + k \, m \, \delta^{ab} \, \delta_{m+n}$$

Virasoro and $\hat{\mathfrak{g}}_k$ must be compatible

$$c = rac{k \dim \mathfrak{g}}{k + C_{\mathfrak{g}}}$$
 dual Coexeter number $C_{\mathfrak{g}}$

We consider $SU(N)_k$ theories

$$\dim(\mathfrak{g}) = N^2 - 1 \qquad C_{\mathfrak{g}} = N \qquad c = \frac{k(N^2 - 1)}{k + N}$$

't Hooft large-N

$$N, k o \infty$$
 $\lambda_s = rac{N}{k}$ fixed

OTO in $SU(N)_k$ WZW

Affine primaries (N - 1 rows, k columns)



Late time OTO fundamental (\Box)

$$C_{gg}^{eta}(t) o q^{rac{1}{N}+rac{1}{2}} rac{\left(q^{-rac{N+2}{2}}+[N-1]
ight)}{[N]}$$

Introducing the q-numbers

$$[x] = \frac{q^{x/2} - q^{-x/2}}{q^{1/2} - q^{-1/2}} \qquad q = e^{-\frac{2\pi i}{N+k}}$$

Key point

At large-
$$c \qquad C^{eta}_{gg}(t) = 1$$

Caputa, Numasawa, AVO '16

Late time OTO in RCFT

In a rational conformal field theory (RCFT)

$$f(z,ar{z}) = \sum_{p} \mathcal{F}^{ii}_{jj}(p|z)ar{\mathcal{F}}^{ii}_{jj}(p|ar{z}).$$

We pick the monodromy



$$C_{ij}^{eta}(t)
ightarrow rac{1}{d_i d_j} rac{S_{ij}^*}{S_{00}}$$

Quantum dimension (d_i) (S_{ii}) Modular S-matrix

Caputa, Numasawa, AVO '16, Gu, Qi '16

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What is the quantum dimension?

Modular S-matrix

$$\chi_i(\tau) = \operatorname{Tr}_{\mathcal{H}_i}\left(e^{i\pi\tau(L_0 - c/24)}\right) \qquad \chi_i\left(-1/\tau\right) = \sum_j S_{ij}\chi_j\left(\tau\right)$$

Quantum dimension of \mathcal{O}_i

$$d_i = \lim_{\tau \to 0} \frac{\chi_i(\tau)}{\chi_0(\tau)} \qquad \qquad d_i = \frac{S_{0i}}{S_{00}}$$

Dijkgraaf, Verlinde '88

Example: Ising model (1, σ , ϵ)

$$S_{ij} = \begin{pmatrix} 1 & \sigma & \epsilon \\ 1/2 & 1/\sqrt{2} & 1/2 \\ 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/2 & -1/\sqrt{2} & 1/2 \end{pmatrix} \begin{pmatrix} 1 & & |C_{\sigma\sigma}^{\beta}| = 0 \\ \sigma & & |C_{\epsilon\sigma}^{\beta}| = 1 \\ \epsilon & & |C_{\epsilon\sigma}^{\beta}| = 1 \\ \epsilon & & |C_{\epsilon\epsilon}^{\beta}| = 1 \end{pmatrix}$$



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3d TQFT and anyon monodromy

Anyon worldlines



Hopf link



Caputa, Numasawa, AVO '16 Gu, Qi '16

This quantity can be measured in anyon interferometry experiments

Bonderson, Shtengel, Slingerland '06

Connections between the buterfly effect and fractional statistics? OTO for *Fractional Quantum Hall (FQH) fluids:* Ising anyons, Fibonacci anyons, quantum magnets

Locally excited states

Physical setup:



- Start with the ground state of a CFT₂.
- At t = 0 insert an operator $\mathcal{O}(x, t)$ at x = -I.
- Smear the operator with ϵ .

Purity $\Delta S_A^{(n)}(t)$ corresponds to a four-point function.

Cross-ratios
$$z = z(I, L, t, \epsilon)$$

Purity and quantum dimension

For a Rational CFT

$$\Delta S^{(2)}_{\mathcal{A}}(t) \simeq \left\{egin{array}{cc} 0 & t \notin [l,L+l]\ \log(d_{\mathcal{O}}) & t \in [l,L+l] \end{array}
ight.$$



He, Numasawa, Takayanagi, Watanabe '14; Caputa, AVO '15

Purity and quantum dimension

Quasiparticle interpretation



He, Numasawa, Takayanagi, Watanabe '14; Caputa, AVO '15

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Purity in $SU(N)_k$ WZW

Some quantum dimensions



Removing the affine cut-off

$$\lim_{k\to\infty} [x] = x$$

Quantum dimension \longrightarrow irrep dimension

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The jump in purities $\Delta S_A^{(2)}$ will diverge in the large-*c* limit.

Purity in $SU(N)_k$ WZW, large-c

The jump in purity comes from the divergence

$$\mathcal{G}(z,\bar{z})
ightarrow d_{\mathcal{O}}^{-1} \left(\left(1-z
ight) ar{z}
ight)^{-2h} + \ldots$$

as $(z, \overline{z}) \rightarrow (1, 0)$.

If we take $c
ightarrow \infty$ we find

$$\Delta S_A^{(2)}(t) \simeq 2h \log\left(rac{2t}{\epsilon}
ight) - \log(2)$$

This matches holographic computations.

Punchline

No quasiparticles \rightarrow Memory loss \rightarrow Entanglement scrambling

Caputa, Numasawa, AVO '16

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Conclusions

- OTO correlators serve to characterize the buterfly effect.
- Late OTO in RCFT are captured by the (measurable) monodromy scalar.
- There are non-chaotic channels even at large-c.
- ► The purity grows logarithmically (scrambles) at large-c.

Puzzles

- What is the precise relationship between chaos, holography and scrambling?
- Fractional statistics and the butterfly effect?
- Other settings: Higher-spin, coset models, GCFT, LCFT ...
- Could OTO and purity be explored experimentally?
- Interesting models with signatures of chaos: Spin chains, SYK, FKV, etc

... and many other interesting questions?

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Thank you for your attention!!

Entanglement scrambling

After introducing a quench



The mutual information

$$I(A,B) = S_A + S_B - S_{A \cup B}$$

Displays a long term memory effect

Calabrese, Cardy '05

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Entanglement scrambling

Mutual information spike



Spike fades in holographic theories

Breakdown of quasiparticles \longleftrightarrow Entanglement scrambling

Asplund, Bernamonti, Galli, Hartman '15

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