

Out-of-Equilibrium Energy Flow and Entanglement Entropy in AdS/CFT

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E.M. PoS EPS-HEP2015 (2015) 366, J. Erdmenger et al, in
progress (2016).

Issues

- 1 Motivation
- 2 Energy flow in strongly correlated systems
 - A universal regime of thermal transport
 - Holographic Model
 - Linearized solution
- 3 Information Flow
 - Entanglement Entropy
 - Mutual Information
 - Results in $d = 2$

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Far from equilibrium dynamics

- Compute real time dynamics.
- Transition to hydrodynamics at late times.
- In strongly correlated systems:
 - Organizing principles out of equilibrium?
 - Universality classes?

Examples:

- Thermalization of quark-gluon plasma.
- Quenches in condensed matter systems.
- Fluctuations in the fractional Hall effect.

- **Much simpler class of problems** → (non)equilibrium steady state.
- **Concrete setup:** Two heat reservoirs at different temperatures are put in contact at time $t = 0$.

$$\text{Energy density} \equiv \varepsilon(x, t = 0) = \frac{c\pi}{12} [T_L^2 \Theta(-x) + T_R^2 \Theta(x)]$$

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A universal regime of thermal transport

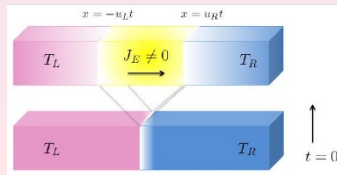
[Bernard, Doyon '12], [Bhaseen, Doyon, Lucas, Schalm '13]

- Two systems initially at equilibrium.
- Independently thermalized.
- Conservation equations and tracelessness:

$$\partial_x \langle T^{xx} \rangle = -\partial_t \langle T^{tx} \rangle = 0, \quad \langle T^{xx} \rangle = \langle T^{tt} \rangle$$

→ "shockwaves" emanating from interface:
$$\begin{cases} \langle T^{tx} \rangle = F(x-t) - F(x+t) \\ \langle T^{tt} \rangle = F(x-t) + F(x+t) \end{cases}$$

- In $d+1$ dim: $\langle T^{\mu\nu} \rangle = a_d T^{d+1} (\eta^{\mu\nu} + (d+1)u^\mu u^\nu)$



Hydrodynamics

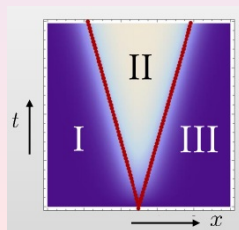
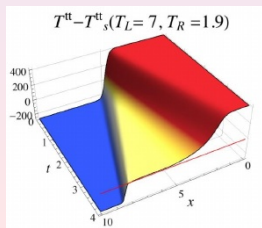
[Bernard, Doyon '12], [Bhaseen, Doyon, Lucas, Schalm '13]

Prediction for the **steady-state**:

- **Lorentz-boosted energy density** with $u^\mu = (\cosh \theta, \sinh \theta, \vec{0})$ and temperature

$$T_L = T \cdot e^\theta, \quad T_R = T \cdot e^{-\theta} \quad \equiv \text{Doppler shift}$$

- **Heat current**: $J_E \equiv \langle T^{tx} \rangle_{st} = a_d \frac{T_L^{d+1} - T_R^{d+1}}{u_L + u_R}$ with $a_d \sim \frac{L^d}{G_N}$.
- **Temperature**: $T_{st} = \sqrt{T_L T_R}$
- **Shockwave velocities**: $c_s^2 = \frac{1}{d} = u_L u_R$.



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Gravitational Dual

- Energy transport: Lorentz-boosted thermal distribution
→ Gravity dual: **Boosted black brane**

$$S = \frac{1}{16\pi G} \int d^{d+1}x \sqrt{-g} \{R + 2\Lambda\}.$$

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= \frac{L^2}{z^2} \left[\frac{dz^2}{f(z)} - f(z) (\cosh \theta dt - \sinh \theta dx)^2 + (\cosh \theta dx - \sinh \theta dt)^2 + dx_\perp^2 \right] \end{aligned}$$

where $f(z) = 1 - \left(\frac{z}{z_h}\right)^d$ and $z_h = \frac{d}{4\pi T}$.

- $g_{\mu\nu}$ is a solution of e.o.m. as long as $z_h \neq z_h(t, x)$ and $\theta \neq \theta(t, x)$.
- (t, x) -dependent solutions → **linearization** [E.Megías '15].

$$z_h(t, x) = z_h^{(0)} + \epsilon z_h^{(1)}(t, x) + \dots, \quad \theta(t, x) = \epsilon \theta^{(1)}(t, x) + \dots.$$

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Linearization

$$z_h(t, x) = z_h^{(0)} + \epsilon z_h^{(1)}(t, x) + \dots, \quad \theta(t, x) = \epsilon \theta^{(1)}(t, x) + \dots$$

$$g_{\mu\nu} = g_{\mu\nu}^{(0)} + \epsilon \begin{pmatrix} \delta g_{tt}(t, x, z) & \delta g_{tx}(t, x, z) & 0 & 0 \\ \delta g_{tx}(t, x, z) & \delta g_{xx}(t, x, z) & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \delta g_{zz}(t, x, z) \end{pmatrix}$$

e.o.m. $\rightarrow \begin{cases} 0 = 2\partial_t^2 z_h^{(1)}(t, x) - \partial_x^2 z_h^{(1)}(t, x) \\ 0 = -z_h^{(0)} \partial_x \theta^{(1)}(t, x) + 2\partial_t z_h^{(1)}(t, x) \end{cases}$

Solution: Assuming some initial profile $T_{ini}(x) \equiv T(t=0, x)$ $\rightarrow$

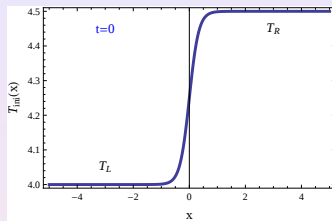
$$T(t, x) = \frac{3}{4\pi z_h(t, x)} = \frac{1}{2} \left[T_{ini}\left(x + \frac{t}{\sqrt{2}}\right) + T_{ini}\left(x - \frac{t}{\sqrt{2}}\right) \right]$$

$$\theta(t, x) = -\epsilon \frac{2\sqrt{2}}{3} \pi z_h^{(0)} \left[T_{ini}\left(x + \frac{t}{\sqrt{2}}\right) - T_{ini}\left(x - \frac{t}{\sqrt{2}}\right) \right]$$

with $z_h^{(0)} = \frac{3}{2\pi(T_L + T_R)}$. Note $c_s = \frac{1}{\sqrt{2}}$.

Linearization

- Let us consider the initial profile:



$$T_{\text{ini}}(x) = \frac{T_R + T_L}{2} + \frac{T_R - T_L}{2} \tanh(\alpha x)$$

- Then one gets the t -dependent solution:

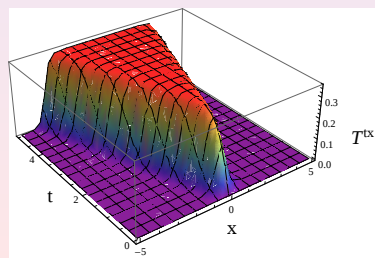
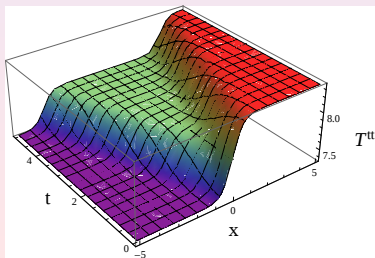
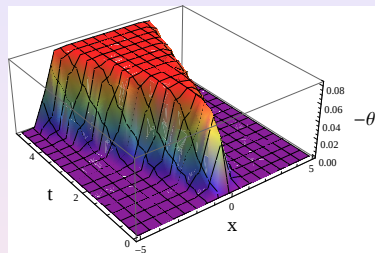
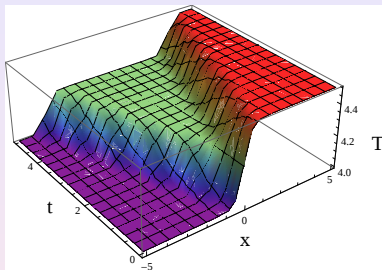
$$\langle T^{tt}(t, x) \rangle = \frac{(d-1)}{8G} \frac{1}{z_{h(0)}^{d-1}} \left[-\frac{(d-1)}{2\pi z_{h(0)}} + \epsilon (T_{\text{ini}}(x + c_s t) + T_{\text{ini}}(x - c_s t)) \right],$$

$$\langle T^{tx}(t, x) \rangle = -\epsilon \frac{1}{c_s} \frac{1}{8G} \frac{1}{z_{h(0)}^{d-1}} [T_{\text{ini}}(x + c_s t) - T_{\text{ini}}(x - c_s t)],$$

→ $\mathcal{O}(\epsilon)$ are corrections in $\frac{|T_R - T_L|}{T_R + T_L} \ll 1$ → Small gradients in $T(t, x)$.

Example: $T_{\text{steady state}} \equiv T(t \rightarrow +\infty, x) = \sqrt{T_L T_R} = \frac{T_L + T_R}{2} + \mathcal{O}(\epsilon^2)$

Linearization



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Entanglement Entropy

How does information get exchanged between the systems which are isolated at $t = 0$?

- Entanglement entropy: Holographic measure $\implies$ Generalization of Bekenstein-Hawking entropy formula

[Ryu, Takayanagi '06], [Hubeny, Rangamani, Takayanagi '07]

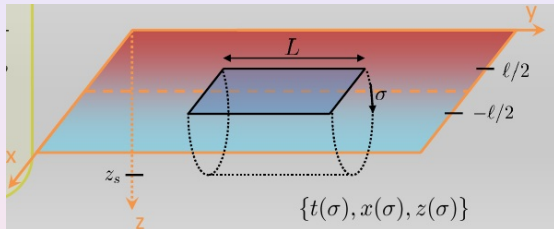
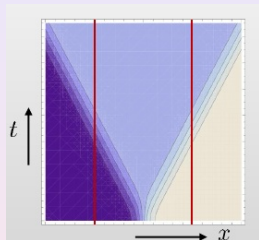
$$S_A = \frac{\text{Area}(\gamma_A)}{4G_N^{(d+2)}}; \quad \text{Area with boundary } \partial A$$

- Regularization: $\text{Area}(\gamma_A^{\text{div}}) \sim \frac{1}{z_{uv}}$.
- Spatial induced metric on surface: $ds^2 = h_{\sigma\sigma} d\sigma^2 + g_{yy} dy^2$.
- Action: $S = L \int ds \sqrt{\frac{\partial x^a}{\partial \sigma} \frac{\partial x^b}{\partial \sigma} \tilde{g}_{ab}}$
- **Minimal surface** $\rightarrow$ **Geodesic equations** for metric $\tilde{g}_{ab} = g_{yy} g_{ab}$:

$$\frac{d^2 x^a}{ds^2} + \tilde{\Gamma}_{bc}^a \frac{dx^b}{ds} \frac{dx^c}{ds} = 0, \quad a = t, x, z$$

$$s \equiv \text{affine parameter} \rightarrow \frac{\partial x^a}{\partial s} \frac{\partial x^b}{\partial s} \tilde{g}_{ab} = 1.$$

Entanglement Entropy



$$S_A^{reg} = \frac{1}{4G_N^{d+2}} \left(\text{Area}(\gamma_A) - \text{Area}(\gamma_A^{div}) \right)$$

Entanglement Entropy

$$S = L \int_{-s_{UV}}^{s_{UV}} ds = 2L \cdot s_{UV}$$

$$S^{reg} = S - S^{div} = 2L \cdot \left(s_{UV} - \frac{1}{z(s_{UV})} \right)$$

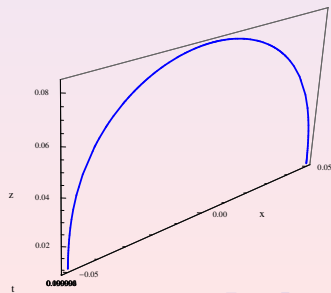
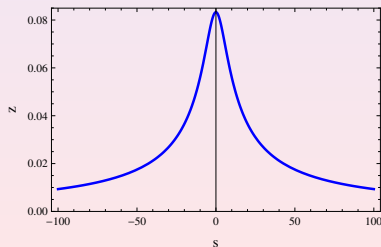
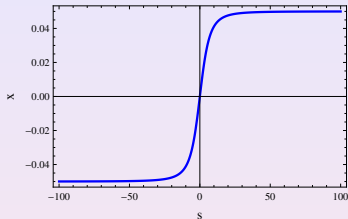
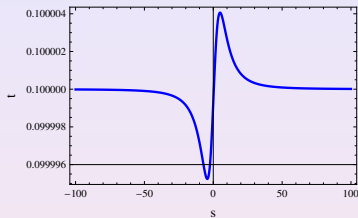
- **Geodesic equations** are 3 coupled equations of 2nd order $\implies$ solutions:

$$t = t(s), \quad x = x(s), \quad z = z(s)$$

$\rightarrow$ 6 boundary conditions:

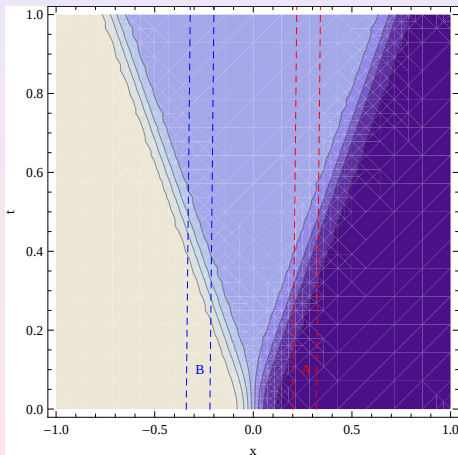
$$\begin{cases} t(s_{UV}) = t(-s_{UV}) = t_0 \\ x(s_{UV}) = -x(-s_{UV}) = \ell/2 \\ z(s_{UV}) = z(-s_{UV}) = z_{UV} \end{cases}$$

Geodesics

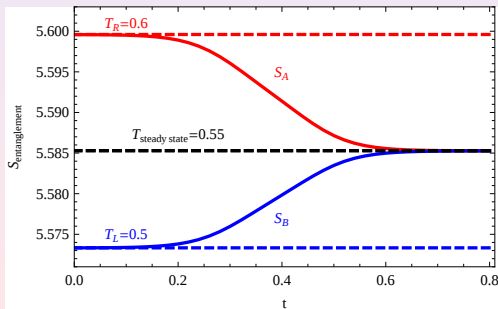
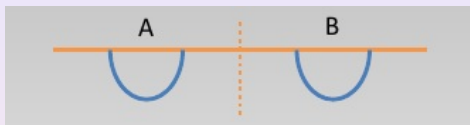


Entanglement Entropy

Contour plot of $T(t, x)$ with $T_L = 0.5$ and $T_R = 0.6$.



Entanglement Entropy



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Mutual Information

What do we learn about A by looking at B?

→ Mutual information: $I(A, B) = S(A) + S(B) - S(A \cup B)$



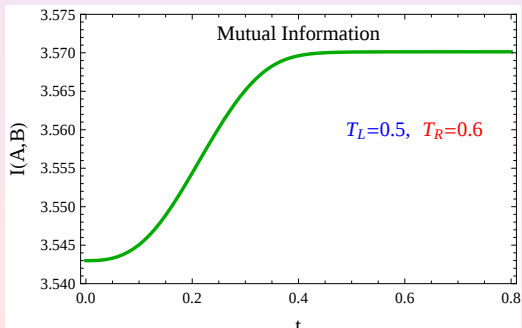
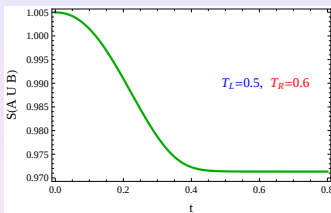
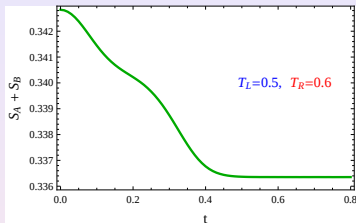
Basic features:

- Shockwaves transport information.
- When does system A find out about the existence of system B?
- $I(A, B) = \text{finite}$ → no need to regularize!!!

We find:

$$\partial_t I(A, B) \geq 0$$

Mutual Information



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Results in $d = 2$

- Exact analytical solution [Bhaseen et al. '13]

$$ds^2 = \frac{L^2}{z^2} \left(dz^2 + g_{\mu\nu} dx^\mu dx^\nu \right),$$

$$g_{tt} = - \left(1 - \frac{z^2}{L^2} (f_R(x-t) + f_L(x+t)) \right)^2 + \left(\frac{z^2}{L^2} (f_R(x-t) - f_L(x+t)) \right)^2$$

$$g_{tx} = -2 \frac{z^2}{L^2} (f_R(x-t) - f_L(x+t))$$

$$g_{xx} = \left(1 + \frac{z^2}{L^2} (f_R(x-t) + f_L(x+t)) \right)^2 - \left(\frac{z^2}{L^2} (f_R(x-t) - f_L(x+t)) \right)^2$$

- Energy-momentum tensor

$$\langle T^{tt} \rangle = \langle T^{xx} \rangle = \frac{c}{6\pi^2 L} (f_R(x-t) + f_L(x+t))$$

$$\langle T^{tx} \rangle = \frac{c}{6\pi^2 L} (f_R(x-t) - f_L(x+t))$$

Results in $d = 2$

- Entanglement entropy in a **boosted BH** in $(\Delta x, \Delta t = 0) \iff$
Entanglement entropy in a **non-boosted BH** in $(\Delta \tilde{x}, \Delta \tilde{t})$

$$S_{\text{boost}}(T_L, T_R, \ell) = 2 \log \left(\frac{1}{\pi \sqrt{T_L T_R}} \sinh^{\frac{1}{2}}(\pi \ell T_L) \sinh^{\frac{1}{2}}(\pi \ell T_R) \right)$$

Valid in the steady-state regime $t > |x|$.

- Entanglement entropy in a stepwise BH

$$ds^2 = ds_L^2 \Theta(-x) + ds_R^2 \Theta(x)$$

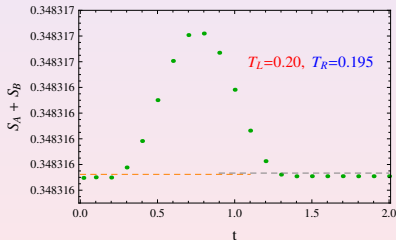
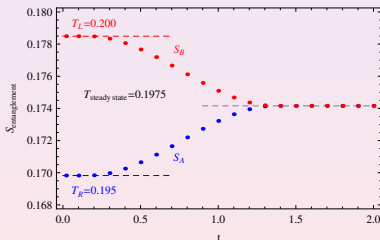
$$\begin{aligned} S(T_L, x_L; T_R, x_R) &= \\ &= 2 \log \left[\frac{T_L \cosh(\pi T_L x_L) \sinh(\pi T_R x_R) - T_R \sinh(\pi T_L x_L) \cosh(\pi T_R x_R)}{\pi T_L T_R} \right] \end{aligned}$$

Valid in the regime $t = 0$.

Results in $d = 2$

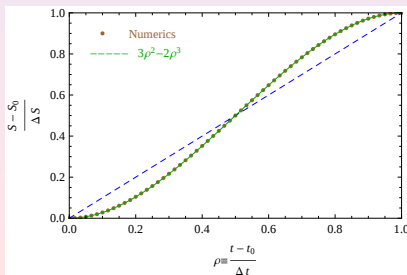
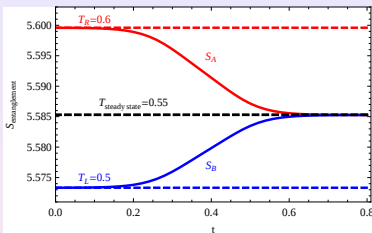
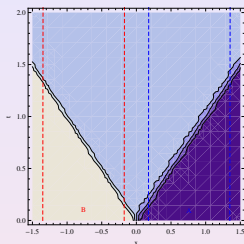
- 'Conservation' of entanglement entropy ($t = 0 \leftrightarrow t = \infty$):

$$S(T_L, \ell; t = 0) + S(T_R, \ell; t = 0) = 2S_{\text{boost}}(T_L, T_R, \ell; t = \infty)$$



$\Rightarrow S_A(t) + S_B(t)$ not exactly conserved at intermediate times.

Universal t-evolution of entanglement entropy ($d = 2$)



$$\frac{S(t) - S(t=0)}{S(t=\infty) - S(t=0)} = f(\rho) = 3\rho^2 - 2\rho^3$$

with $\rho \equiv (t - t_0) / \Delta t = (t - x_L) / \ell$

Good for $|T_L - T_R| < T_L + T_R$

Conclusions and future directions

- We have studied a **holographic model** for time-dep energy flow.
- **Linearization for $d \geq 3$** $\rightarrow$ analytical solution valid for $|T_L - T_R| \ll T_L + T_R$ (*linear response regime*).
- **Exact solution for $d = 2$** [Bhaseen et al. '13].
- We have computed the time evolution of **Entanglement Entropies** $\rightarrow$ **Mutual information grows with time**.
- **Universal time evolution of entanglement entropy** in the linear response regime ($d = 2$).

What next?

- Complete analysis beyond linear response regime for $d \geq 3$: $0 < T_L/T_R < 1$ $\rightarrow$ **Full numerical solution of Einstein e.o.m.** [Chesler, Yaffe '09], [Amado, Yarom '15], [Ecker, Grumiller, Stricker '15], [Erdmenger, Fernandez, Flory, EM, Straub, Witkowski '16].
- **Other possible solutions?** [Chang, Karch, Yarom '14].
- **Higher dimensions:** What changes?

Thank You!