

Yangian Symmetry of $N=4$ SYM

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Joint work with
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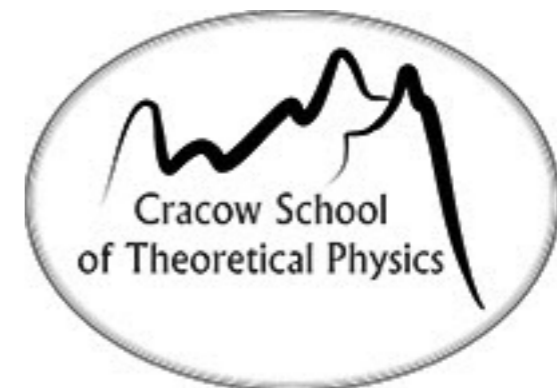


Table of contents

1. Introduction and Motivation*
2. Yangian Symmetry
3. Correlation Functions in Pictures
4. Conclusions & Outlook

*Shortened due to overlap with Jan Plefka's lecture

$N=4$ SYM

Maximally supersymmetric gauge theory in 4D

1 gauge field A_μ , 4 spinors Ψ , 4 conjugate spinors $\bar{\Psi}$,
6 scalars ϕ_m

$psu(2,2|4)$ - symmetry algebra of the action

Integrability in $N=4$ SYM

Spectra of local operators

Wilson loops

Scattering amplitudes

Increased
symmetry algebra:
Yangian

Yangian Algebra

Infinitely many levels of generators.

Level 0 - original $\mathfrak{psu}(2,2|4)$:

$$[J_a, J_b] = f_{ab}^c J_c$$

Level 1 - first higher level Yangian generators:

$$[J_a, \hat{J}_b] = f_{ab}^c \hat{J}_c$$

Coproduct (action on multiparticle states)

$$\Delta(J_a) = id \otimes J_a + J_a \otimes id$$

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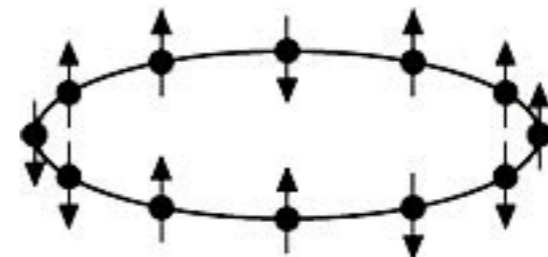
Coproduct (action on multiparticle states)

$$\Delta(\hat{J}_a) = id \otimes \hat{J}_a + \hat{J}_a \otimes id + f_a^{bc} J_b \otimes J_c$$

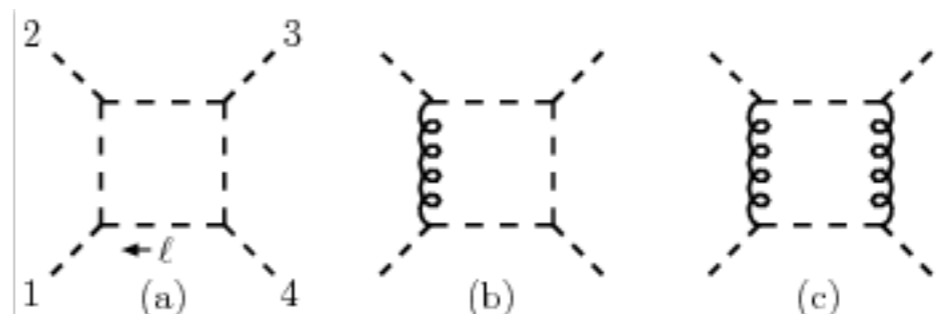
Motivation (for advanced)

We know how Yangian generators act on

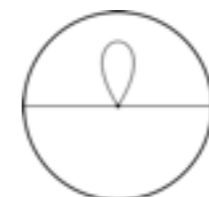
spin chains



scattering amplitudes



Wilson loops



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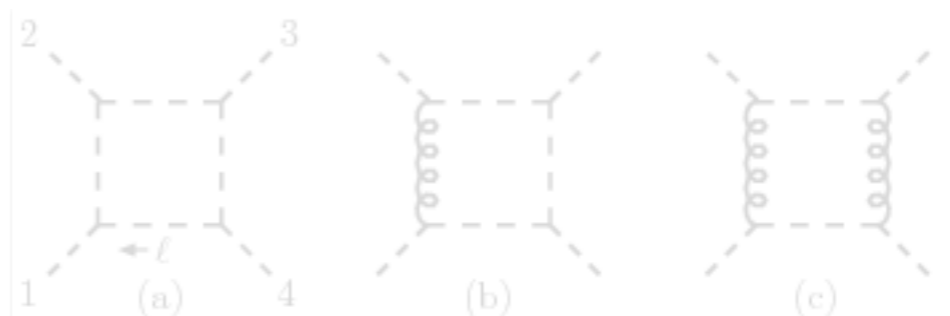
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See Jan Plefka's lecture for details!
(also mentioned in Jacob Bourjaily's talk)



Yangian Action

on Action

$$S = \int d^4x \text{Tr}(\dots)$$

Due to the trace, the action is cyclic:

$$\text{Tr}(ABC) = \text{Tr}(CAB)$$

However:

$$\hat{J}(ABC) \neq \hat{J}(CAB)$$

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Incompatible!

Symmetry of the EoMs

(they are no longer cyclic)

Simplest level-1 generator:

$$\Delta(\hat{P}^\mu) = D \wedge P^\mu + L^{\mu\nu} \wedge P_\nu - \frac{i}{4} \sigma^{\mu\alpha\dot{\alpha}} \bar{Q}_{\dot{\alpha}}^a \wedge Q_{a\alpha}$$

Easiest equation of motion:

$$\sigma^{\rho\dot{\alpha}\beta} [D_\rho, \Psi_{\beta a}] - ig\epsilon^{\dot{\alpha}\dot{\beta}} \sigma_{ab}^m [\Phi_m, \bar{\Psi}_{\dot{\beta}}^b] = 0$$

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etc...

Symmetry of the EoMs

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Introducing the single field action:

$$\hat{P}^\mu(D^\rho) = \frac{1}{2}g^2\eta^{\mu\rho}\{\Phi^m, \Phi_m\}$$

$$\hat{P}^\mu(\Psi_{a\alpha}) = \frac{ig}{2}\sigma_{\alpha\dot{\beta}}^\mu\epsilon^{\dot{\beta}\dot{\gamma}}\sigma_{ab}^m\{\bar{\Psi}_{\dot{\gamma}}^b, \Phi_m\}$$

we obtain schematically:

$$\hat{P}(\text{Dirac}) = 0$$

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we obtain schematically:

$$\hat{P}(Dirac) = \{ \Phi, \overline{Dirac} \}$$

Similarly for all the other equations of motion!

Off-shell?

For level-0 symmetry generators we have:

$$0 = JS = \int (JZ_a) \frac{\delta S}{\delta Z_a}$$

Varying again w. r. t. to a field:

Vanishes on-shell \rightarrow

$$J \frac{\delta S}{\delta Z_c} + \frac{\delta(JZ_a)}{\delta Z_c} \frac{\delta S}{\delta Z_a} = 0$$

Off-shell equality regarding behavior of eoms under J.

Off-shell!

Generalize to level-1 generators:

$$\hat{j}^k \frac{\delta S}{\delta Z_c} = - \frac{\delta(\hat{J}^k Z_a)}{\delta Z_c} \frac{\delta S}{\delta Z_a} + f^k_{mn} \left((J^m Z_a) \frac{\delta}{\delta Z_a} \wedge \frac{\delta S}{\delta Z_b} \frac{\delta}{\delta Z_c} \right) (J^n Z_b)$$

Vanishes on shell

Holds off-shell!

Off-shell!

Generalize to level-1 generators:

$$\hat{J}^k \frac{\delta S}{\delta \Phi_c} = \frac{\delta(\hat{J}^k \Phi_a)}{\delta \Phi_c} \frac{\delta S}{\delta \Phi_a} + \epsilon^{kmn} \left((J^m \Phi_a) \frac{\delta}{\delta \Phi_a} \wedge \frac{\delta S}{\delta \Phi_b} \frac{\delta}{\delta \Phi_c} \right) (J^n \Phi_b)$$

Variables on shell

Holds off-shell!

Doesn't hold for non-integrable field theories!

Brief recap

- Yangian algebra appears in many observables in $N=4$ SYM.
- At the first level, the Yangian generators act bilocally (here - in color space).
- Cannot act with them on the action, but they are an on-shell symmetry of equations of motion.
- We put forward an off-shell equality we believe equivalent to the invariance of the action.

Correlation Functions

(tree level)

If $JS = 0$ and J leaves the \mathcal{P} I measure invariant, the correlation functions are also invariant

(Slavnov-Taylor / Ward identity):

$$J \langle Z_1(x_1) Z_2(x_2) \dots Z_n(x_n) \rangle = \sum_{i=1}^N \langle Z_1(x_1) \dots (J Z_i(x_i)) \dots Z_n(x_n) \rangle = 0$$

Ars Gratia Artis*

Introduce graphical notation:

$$S = \begin{array}{|c} \hline \\ \hline \end{array} + g \begin{array}{|c|c} \hline \\ \hline \end{array} + g^2 \begin{array}{|c|c|c} \hline \\ \hline \end{array} + \dots$$

$$JZ = \begin{array}{|c} \hline \square \\ \hline \end{array} + g \begin{array}{|c|c} \hline \square \\ \hline \end{array} + \dots$$

$$\overset{v}{Z} = \begin{array}{|c} \bullet \\ \hline \\ \hline \end{array} + g \begin{array}{|c|c} \bullet \\ \hline \\ \hline \end{array} + \dots$$

Example:

Propagator Invariance

Level-0 generator \mathcal{J} acting on a propagator:

$$\mathcal{J} \left[\begin{array}{c} | \\ | \\ \hline | \\ | \end{array} \right] = \left[\begin{array}{c} | \\ \square \\ | \\ \hline | \\ | \end{array} \right] + \left[\begin{array}{c} | \\ | \\ \hline \square \\ | \end{array} \right]$$

Level-0 magic formula $\mathcal{J} \frac{\delta S}{\delta Z_c} + \frac{\delta(\mathcal{J} Z_a)}{\delta Z_c} \frac{\delta S}{\delta Z_a} = 0$ at 0-th order:

$$\left[\begin{array}{c} \bullet \\ | \\ \square \\ | \\ \hline | \\ | \end{array} \right] + \left[\begin{array}{c} \bullet \\ | \\ | \\ \hline \square \\ | \end{array} \right] = \left[\begin{array}{c} \bullet \\ | \\ \square \\ | \\ \hline | \\ | \end{array} \right] + \left[\begin{array}{c} \bullet \\ | \\ | \\ \hline \square \\ | \end{array} \right] = 0$$

Invariance under level-1

- Draw level-1 magic formula
- Draw action of Yangian generator on the correlation function
- Local action cancels easily - tricky part is the bilocal part of the coproduct
- Level-0 is a symmetry, we can move it around
- Dual coxeter number of $\mathfrak{psu}(2,2|4)$ vanishes:

$$f_{bc}^a f_d^{bc} = h \delta_b^a = 0$$

Final recap




- 🔔 Yangian of $\mathfrak{psu}(2,2|4)$ - correct symmetry algebra of planar $N=4$ SYM.
- 🔔 Invariance of equations of motion under level-1.
- 🔔 Cannot act on the action, but exists a way to circumvent it.
- 🔔 Gauge-fixing (BRST) - compatible.
- 🔔 Implies invariance of the correlation functions.

Final recap

- Yangian of $\mathfrak{psu}(2,2|4)$ - correct symmetry algebra of planar $N=4$ SYM.
- Invariance of equations of motion under level-1.
- Cannot act on the action, but exists a way to circumvent it.
- Gauge-fixing (BRST) - compatible.
- Implies invariance of the correlation functions.

Holds for ABJM too!

To investigate

-  Anomalies
-  Algebraic relations for Yangian generators
-  Proper Ward identities

Thanks for your attention!

