Spplications of On-Shell Physics

Jacob L. Bourjaily

Cracow School of Theoretical Physics LVI Course, 2016 A Panorama of Holography



Friday, 27th May Cracow School of Theoretical Physics, Zakopane Part III: Applications of On-Shell Physics: Generalized Unitarity (Redux)

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Organization and Outline

- Spiritus Movens: the Discovery of On-Shell Physics
 - Using Generalized Unitarity to Compute One-Loop Amplitudes
- 2 Revisiting Generalized Unitarity: Improving the One-Loop Toolbox
 - Finite Scalar Box Integrals and their Infrared-Divergent Limits
 - Maximally Preserving Dual-Conformal Invariance of Divergences
- **Output** Upgrading Unitarity at One-Loop: the *Chiral* Box Expansion
 - Chiral Boxes Expansion for One-Loop Integrands
 - Making Manifest the Finiteness of All Finite Observables
- Generalizing Unitarity for Two-Loop Amplitudes & Integrands
 - The Two-Loop Chiral Integrand Expansion
 - Novel Contributions at Two-Loops and Transcendentality
- 5 The Ongoing Revolution in Our Understanding of Quantum Field Theory

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Using Generalized Unitarity to Compute One-Loop Amplitudes

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Advantages:

- each standardized, scalar integral need only be computed once
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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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The Scalar "Four-Mass Box" Integral

The four-mass box integral is a manifestly finite, symmetric function of two dual-conformally invariant cross ratios, denoted u and v.

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$$(a,b) \equiv (x_a - x_b)^2 = (p_a + p_{a+1} + \dots + p_{b-1})^2$$

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$$I_{a,b,c,d} \equiv d \xrightarrow{a} b = -\int_{\ell \in \mathbb{R}^{3,1}} d^4 \ell \frac{(a,c)(b,d)\Delta}{(\ell,a)(\ell,b)(\ell,c)(\ell,d)},$$

$$\Delta \equiv \sqrt{(1-u-v)^2 - 4uv} c \qquad u \equiv \frac{(a,b)(c,d)}{(a,c)(b,d)}, \quad v \equiv \frac{(b,c)(a,d)}{(a,c)(b,d)},$$

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$$I_{a,b,c,d} \equiv \underbrace{I_{a,b,c,d}}_{I = a,b,c,d} = -\int_{\ell \in \mathbb{R}^{3,1}} d^4 \ell \frac{(a,c)(b,d)\Delta}{(\ell,a)(\ell,b)(\ell,c)(\ell,d)},$$
$$\equiv \sqrt{(1-u-v)^2 - 4uv} \cdot C \qquad u \equiv \frac{(a,b)(c,d)}{(a,c)(b,d)}, \quad v \equiv \frac{(b,c)(a,d)}{(a,c)(b,d)},$$
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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

The Scalar "Four-Mass Box" Integral and its Divergences

When any corner becomes massless, the integral becomes infrared divergent

Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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When any corner becomes massless, the integral becomes **infrared divergent** *e.g.*, if we send $(a, b) \rightarrow O(\epsilon)$, then $u \rightarrow O(\epsilon)$, causing a divergence:

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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$$I = \frac{(u,v)}{(a,c)(b,d)} = \operatorname{Li}_{\ell}(a) + \operatorname{Li}_{\ell}(b) = \operatorname{Li}_{\ell}(a) + \operatorname{Li}_{$$

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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The Scalar "Four-Mass Box" Integral and its Divergences

When any corner becomes massless, the integral becomes infrared divergent *e.g.*, if we send $(a, b) \rightarrow \mathcal{O}(\epsilon)$, then $u \rightarrow \mathcal{O}(\epsilon)$, causing a divergence:

$$I_{a,b,c,d} \equiv \underbrace{d}_{\ell \in \mathbb{R}^{3,1}} b \equiv -\int_{\ell \in \mathbb{R}^{3,1}} d^{4}\ell \frac{(a,c)(b,d)\Delta}{(\ell,a)(\ell,b)(\ell,c)(\ell,d)},$$
$$\Delta \equiv \sqrt{(1-u-v)^{2}-4uv} \cdot C \qquad u \equiv \frac{(a,b)(c,d)}{(a,c)(b,d)}, \quad v \equiv \frac{(b,c)(a,d)}{(a,c)(b,d)},$$
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$$L_{a,b,c,d}(u,v) = Li_2(\alpha) + Li_2(\beta) - Li_2(1) + \frac{1}{2}\log(u)\log(v) - \log(\alpha)\log(\beta)$$

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

The Scalar "Four-Mass Box" Integral and its Divergences

When any corner becomes massless, the integral becomes **infrared divergent** *e.g.*, if we send $(a, b) \rightarrow O(\epsilon)$, then $u \rightarrow O(\epsilon)$, causing a divergence:

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$$I_{a,b,c,d} \equiv \begin{array}{c} & \overset{a}{\longrightarrow} b^{-1} \\ & \overset{b}{\longleftarrow} b \end{array} \equiv -\int_{\ell \in \mathbb{R}^{3,1}} d^{4}\ell \ \frac{(a,c)(b,d)\Delta}{(\ell,a)(\ell,b)(\ell,c)(\ell,d)}, \\ & \overset{(1-\nu)}{\longrightarrow} + \mathcal{O}(\epsilon) \end{array}$$

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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Finite Scalar Box Integrals and their Infrared-Divergent Limits Maximally Preserving Dual-Conformal Invariance of Divergences

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A Dual-Conformal Regularization of Infrared Divergences

In order to regulate the infrared divergences of the box integrals, we render **all** external legs off-shell by displacing the coordinates according to:

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Under this shift, all cross-ratios are displaced proportional to cross-ratios!



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Under this shift, all cross-ratios are displaced proportional to cross-ratios!

e.g., when a = b - 1



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A Dual-Conformal Regularization of Infrared Divergences

In order to regulate the infrared divergences of the box integrals, we render **all** external legs off-shell by displacing the coordinates according to:

$$x_a \to \hat{x}_a \equiv x_a + \epsilon (x_{a+1} - x_a) \frac{(a-2,a)}{(a-2,a+1)}$$

Under this shift, all cross-ratios are displaced proportional to cross-ratios!

e.g., when a = b - 1



Friday, 27th May Cracow School of Theoretical Physics, Zakopane Part III: Applications of On-Shell Physics: Generalized Unitarity (Redux)

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Under this shift, all cross-ratios are displaced proportional to cross-ratios!

e.g., when
$$a = b - 1$$
, we have: $(a, \widehat{b}) \mapsto \epsilon(a, b+1) \frac{(b-2, b)}{(b-2, b+1)} + \mathcal{O}(\epsilon^2)$



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Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \,\mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

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Consider for example the 'MHV' amplitude

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Consider for example the 'MHV' amplitude (k=2)

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Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$

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$$f_{a,a+1,c,c+1}^{1} = \begin{array}{c} c+1 & Q_{1} \\ Q_{1} & Q_{2} \\ \vdots & \vdots \\ c &$$

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A 'Box'-Expansion for One-Loop Integrands

A *Chiral* 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(2),1} \stackrel{\scriptscriptstyle G}{=} \sum_{a,c} \mathcal{I}_{a,a+1,c,c+1}^1 f_{a,a+1,c,c+1}^1$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are:



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A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are: d^a

$$f_{a,a+1,c,c+1}^{1} = \underbrace{c+1}_{c} \underbrace{Q_{1}}_{c} a+1$$

$$c+1 \underbrace{Q_{1}}_{c} a+1 \Leftrightarrow \int d^{4}\ell \underbrace{(a,c)(a,a+1)-(a,c+1)(c,a+1)}_{(\ell,a)(\ell,a+1)(\ell,c)(\ell,c+1)} \underbrace{(\ell,Q_{2})(X,Q_{1})}_{(\ell,X)(Q_{2},Q_{1})}$$

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This ansatz matches the correct integrand on all co-dimension four residues

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

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This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*.

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

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A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}^2}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**!

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

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This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

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$$\Leftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

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$$\mathcal{I}_{\rm div}^a \equiv : \qquad \Longleftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

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and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

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A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

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A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} = \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv : \qquad \Leftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

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A 'Box'-Expansion for One-Loop Integrands



Because the divergences are universal, the ratio function is manifestly finite!

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Because the divergences are universal, the ratio function is manifestly finite!

$$\mathcal{R}_n^{(k),1} \equiv \mathcal{A}_n^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_n^{(2),1}$$

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A 'Box'-Expansion for One-Loop Integrands



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$$\begin{aligned} \mathcal{R}_n^{(k),1} &\equiv \mathcal{A}_n^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_n^{(2),1} \\ &= \mathcal{A}_{n,\text{fin}}^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_{n,\text{fin}}^{(2),1} \end{aligned}$$

Manifesting the Exponentiation of Divergences to All Orders

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders.

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$$\mathcal{R}_{n}^{(k),\ell} = \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} - \sum_{q=1}^{\ell} \mathcal{R}_{n}^{(k),\ell-q} \mathcal{A}_{n,\mathrm{fin}}^{(2),q}$$

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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

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"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) + \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

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 $\mathcal{I}_L(X) \otimes \mathcal{I}_R(X)$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

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 $\mathcal{I}_L(\mathbf{X}) \bigotimes \mathcal{I}_R(\mathbf{X}) \equiv$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right) + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

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$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) + \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

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"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) + \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{n,\mathrm{div}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$
The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathrm{H}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{(k),0} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

 $\frac{(b-1,b+1)(b,X)}{(\ell_1,b-1)(\ell_1,b)(\ell_1,b+1)(\ell_1,X)} \otimes \frac{(X,a)(a-1,a+1)}{(X,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$

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Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{(k),\mathrm{div}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

$$\frac{(b-1,b+1)(b,a)(a-1,a+1)}{(\ell_1,b-1)(\ell_1,b)(\ell_1,b+1)(\ell_1,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$$

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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

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"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{(k),\mathrm{div}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

$$\frac{(b-1,b+1)(b,a)(a-1,a+1)}{(\ell_1,b-1)(\ell_1,b)(\ell_1,b+1)(\ell_1,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$$

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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

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$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{b} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{b} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

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Constructing Local Integrands for Two-Loop Amplitudes

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"Merging" One-Loop, Chiral (X-dependent) Integrands

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$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)$$

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Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

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$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{b}$$

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Constructing Local Integrands for Two-Loop Amplitudes

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"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

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Constructing Local Integrands for Two-Loop Amplitudes

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"Merging" One-Loop, Chiral (X-dependent) Integrands

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$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{a} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(k),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{($$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Constructing Local Integrands for Two-Loop Amplitudes

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"Merging" One-Loop, Chiral (X-dependent) Integrands

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$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{a} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(k),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{($$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Finite Integrand Contributions to Two-Loop Amplitudes

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



2. Finite Penta-Boxes:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



2. <u>Finite</u> Penta-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



2. Finite Penta-Boxes:



Friday, 27th May

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

3. Finite Double-Boxes:



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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

3. Finite Double-Boxes:



2. <u>Finite</u> Penta-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

3. <u>Finite</u> Double-Boxes:



2. <u>Finite</u> Penta-Boxes:



Friday, 27th May

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Finite Integrand Contributions to Two-Loop Amplitudes

3. <u>Finite</u> Double-Boxes:



2. <u>Finite</u> Penta-Boxes:



Friday, 27th May

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Novel Contributions Required

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Novel Contributions Required

4. "Shifted" Double-Boxes:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'.

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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Novel Contributions Required

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It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

 $\mathcal{A}_{10}^{(5)}\big(\varphi_{12},\varphi_{12},\varphi_{12},\varphi_{23},\varphi_{23},\varphi_{34},\varphi_{34},\varphi_{34},\varphi_{41},\varphi_{41}\big)$

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 $\mathcal{A}_{10}^{(5)}(\varphi_{12},\varphi_{12},\varphi_{12},\varphi_{23},\varphi_{23},\varphi_{34},\varphi_{34},\varphi_{34},\varphi_{41},\varphi_{41})$

 $\propto (\widetilde{\eta}_1^1 \widetilde{\eta}_1^2) (\widetilde{\eta}_2^1 \widetilde{\eta}_2^2) (\widetilde{\eta}_3^1 \widetilde{\eta}_3^2) (\widetilde{\eta}_4^2 \widetilde{\eta}_4^3) (\widetilde{\eta}_5^2 \widetilde{\eta}_5^3) (\widetilde{\eta}_6^3 \widetilde{\eta}_6^4) (\widetilde{\eta}_7^3 \widetilde{\eta}_7^4) (\widetilde{\eta}_8^3 \widetilde{\eta}_8^4) (\widetilde{\eta}_9^4 \widetilde{\eta}_9^1) (\widetilde{\eta}_{10}^4 \widetilde{\eta}_{10}^1)$

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The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

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H. The analytic expression of the remainder function

In this appendix we present the full analytic expression of the remainder function. The result is also resultable in electronic form from wev. arXiv.org. Using the notation introduced in Eqs. (3.23) and (5.7), the full expression reads,

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The Two-Loop Hexagon Wilson Loop in $\mathcal{N} = 4$ SYM

Vittorio Del Duca

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Friday, 27th May

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Vladimir A. Smirnov

Nuclear Physics Institute of Moscow State University Moscow 119992, Russia E-mail: smirnov@theory.sinp.msu.ru $H(0; u_3) H(0, 0, 1; (u_1 + u_3)) - \frac{1}{2} H(0; u_1) H(0, 0, 1; \frac{u_2}{2})$ $\frac{1}{2}H(0; u_3) H\left(0, 0, 1; \frac{u_2 + u_3 - 1}{2}\right) - H(0; u_2) H(0, 0, 1; (u_2 + u_3)) H(0; u_3) H(0, 0, 1; (u_2 + u_3)) - \frac{1}{2} H(0; u_2) H(0, 1, 0; u_1) \frac{1}{2}H(0; u_3)H(0, 1, 0; u_2) =$ $\frac{1}{2}H(0; u_1)H(0, 1, 0; u_3) + \frac{1}{4}H(0; u_2)H(0, 1, 1; \frac{u_1 + u_2}{2})$ $\frac{1}{2}H(0; u_1)H(0, 1, 1; \frac{u_1 + u_2 - 1}{2}) + \frac{1}{2}H(0; u_1)H(0, 1, 1; \frac{u_1 + u_3 - 1}{2})$ $\frac{1}{2}H(0; u_2)H(0, 1, 1; \frac{u_1 + u_1 - 1}{2}) - \frac{1}{4}H(0; u_1)H(0, 1, 1; \frac{u_2 + v_1}{2})$ $\frac{1}{H}(0; u_2)H(0, 1, 1; \frac{u_2 + u_3 - 1}{2}) + \frac{1}{H}(0; u_2)H(1, 0, 0; u_1) - \frac{1}{H}(0; u_2)H(1, 0; u_1) - \frac{1}{H}(1, 0; u_1)H(1, 0; u_2)H(1, 0; u_1) - \frac{1}{H}(1, 0; u_1)H(1, 0; u_1)H(1, 0; u_1) - \frac{1}{H}(1, 0; u_1)H(1, 0; u_1)H(1, 0; u_1)H(1, 0; u_1)H(1, 0; u_1)H(1$ $\frac{1}{2}H(0; u_1)H(1, 0, 0; u_2) + \frac{1}{2}H(0; u_3)H(1, 0, 0; u_2) + \frac{1}{2}H(0; u_1)H(1, 0, 0; u_3) \frac{1}{2}H(0; u_2)H(1, 0, 0; u_3) - \frac{1}{2}H(0; u_3)H(1, 0, 1; \frac{u_3}{2})$ $\left(\frac{1}{2}\right) - \frac{1}{2}H(0; u_1)H(1, 0, 1; \frac{u_2}{2})$ $-7H(0, 0, 0, 0; u_2) - 7H(0, 0, 0, 0; u_3) + \frac{\delta}{2}H(0, 0, 0, 1; \frac{1}{2})$ $3H(0, 0, 0, 1; (u_1 + u_2)) + \frac{3}{2}H(0, 0, 0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}) + 3H(0, 0, 0, 1; (u_1 + u_2)) +$ $\frac{3}{n}H\left(0, 0, 0, 1; \frac{u_2 + u_3 - 1}{n}\right) + 3H\left(0, 0, 0, 1; (u_2 + u_3)\right) + \frac{9}{r}H\left(0, 0, 1, 0; u_1\right) +$ $(1, 0; u_2) + \frac{9}{2}H(0, 0, 1, 0; u_3) - \frac{1}{2}H(0, 1, 0, 0; u_1) - \frac{1}{2}H(0, 1, 0, 0; u_2) - \frac{1}{2}H(0, 1, 0, 0; u_2) - \frac{1}{2}H(0, 1, 0, 0; u_3) - \frac{1}{2}H(0, 0; 0; u_3) - \frac{1}{2}H(0, 0; 0; u_3) - \frac{1}{2}H(0, 0; 0; u_3) - \frac{1}{2}H(0, 0;$ $\frac{4}{2}H(0, 1, 0, 0; u_3) + \frac{4}{2}H(0, 1, 0, 1; \frac{u_1 + u_2 - 1}{m - 1}) + \frac{1}{2}H(0, 1, 0, 1; \frac{u_1 + u_3 - 1}{m - 1}) +$ $+u_2-1$ $-\frac{1}{2}H(0, 1, 1, 1; \frac{u_1+u_3-1}{2})$ $(1, 1, 1; \frac{u_2 + u_3 - 1}{w_3 - 1}) + H(1, 0, 0, 1; \frac{u_1 + u_2 - 1}{w_3 - 1}) + H(1, 0, 0, 1; \frac{u_1 + u_3 - 1}{w_3 - 1})$ $\left(1, 0, 0, 1; \frac{u_{2} + u_{3} - 1}{u_{*} - 1}\right) + 2H(1, 0, 1, 0; u_{1}) + 2H(1, 0, 1, 0; u_{2}) + 2H(1, 0, 1, 0; u_{3}) +$ $\frac{1}{4}H\left(1, 1, 0, 1; \frac{u_1 + u_2 - 1}{u_2 - 1}\right) + \frac{1}{4}H\left(1, 1, 0, 1; \frac{u_1 + u_3 - 1}{u_1 - 1}\right) +$ $\frac{1}{4}H\left(1, 1, 0, 1; \frac{u_{2} + u_{3} - 1}{2}\right) + \frac{1}{2}H\left(1, 1, 1, 0; u_{1}\right) + \frac{1}{2}H\left(1, 1, 1, 0; u_{2}\right) + \frac{1}{2}H\left(1, 1, 1, 0; u_{3}\right) \frac{1}{2}x^{2}H(0; u_{3}) \mathcal{H}\left(1; \frac{1}{1}\right) - \frac{1}{2}x^{2}H(0; u_{1}) \mathcal{H}\left(1; \frac{1}{1}\right) - \frac{1}{2}x^{2}H(0; u_{2}) \mathcal{H}\left(1; \frac{1}{1}\right) +$ $\frac{1}{2}\pi^{2}H(0; u_{2}) \mathcal{H}\left(1; \frac{1}{1; \frac{1}{m_{max}}}\right) - \frac{1}{8}\pi^{2}H(0; u_{3}) \mathcal{H}\left(1; \frac{1}{m_{max}}\right) + \frac{1}{24}\pi^{2}H(0; u_{2}) \mathcal{H}\left(1; \frac{1}{m_{max}}\right) - \frac{1}{24}\pi$

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Not long ago, del Duca, Duhr, and Smirnov determined the 2-loop, 6-particle amplitude $\mathcal{A}_6^{(2),2}$ **analytically**—a truly heroic computation on par with Parke and Taylor's

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 $\frac{1}{\pi}\pi^{2}H(0; u_{1}) \mathcal{H}\left(1; \frac{1}{1}\right) + \frac{1}{\pi}\pi^{2}H(0; u_{3}) \mathcal{H}\left(1; \frac{1}{1}\right) + \frac{1}{\pi}\pi^{2}H(0; u_{1}) \mathcal{H}\left(1; \frac{1}{1}\right)$ $\frac{1}{8}\pi^2 H(0; u_2) \mathcal{H}\left(1; \frac{1}{u_{uax}}\right) + \frac{1}{24}\pi^2 H(0; u_1) \mathcal{H}\left(1; \frac{1}{u_{uax}}\right) - \frac{1}{24}\pi^2 H(0; u_2) \mathcal{H}\left(1; \frac{1}{$ $\frac{1}{4}H(0, u_2) H(0; u_3) \mathcal{H}\left(0, 1; \frac{1}{w_{100}}\right) - \frac{1}{4}H(1, 0; u_2) \mathcal{H}\left(0, 1; \frac{1}{w_{100}}\right) + \frac{1}{24}\pi^2 \mathcal{H}\left(0, 1; \frac{1}{w_{100}}\right) +$ $\frac{1}{24}\pi^2 \mathcal{H}\left(0, 1; \frac{1}{u_{00}}\right) - \frac{1}{4}\mathcal{H}\left(0; u_1\right) \mathcal{H}\left(0; u_3\right) \mathcal{H}\left(0, 1; \frac{1}{u_{00}}\right) - \frac{1}{4}\mathcal{H}\left(1, 0; u_3\right) \mathcal{H}\left(1, 0; u_3\right) \mathcal{H}$ $\frac{1}{4}H(0, u_1)H(0; u_2)H(0, 1; \frac{1}{u_{max}}) - \frac{1}{4}H(1, 0; u_1)H(0, 1; \frac{1}{u_{max}}) + \frac{1}{24}\pi^2 H(0, 1; \frac{1}{u_{max}})$ $\frac{1}{4}H(0; u_2) H(0; u_3) \mathcal{H}(0; 1; \frac{u_{312}}{v_{111}}) + \frac{1}{4}H(0, 0; u_2) \mathcal{H}(0, 1; \frac{u_{312}}{v_{111}}) +$ $\frac{1}{4}H(0, 0; u_3) \mathcal{H}\left(0, 1; \frac{1}{a-c}\right) + \frac{1}{a}e^{2}\mathcal{H}\left(0, 1; \frac{1}{a-c}\right) - \frac{1}{4}H(0; u_2) H(0; u_3) \mathcal{H}\left(0, 1; \frac{1}{0; u_2}\right) +$ $\frac{1}{4}H(0, 0; u_2) \mathcal{H}\left(0, 1; \frac{1}{u_{u_1}}\right) + \frac{1}{4}H(0, 0; u_3) \mathcal{H}\left(0, 1; \frac{1}{u_{u_2}}\right) + \frac{1}{6}\pi^2 \mathcal{H}\left(0, 1; \frac{1}{u_{u_1}}\right) - \frac{1}{6}\pi^2 \mathcal{H}\left($ $\frac{1}{4}H(0; u_1)H(0; u_3)\mathcal{H}(0, 1; \frac{1}{u_{max}}) + \frac{1}{4}H(0, 0; u_1)\mathcal{H}(0, 1; \frac{1}{u_{max}}) +$ $\frac{1}{4}H(0, 0; u_3) \mathcal{H}\left(0, 1; \frac{1}{v_{013}}\right) + \frac{1}{6}\pi^2 \mathcal{H}\left(0, 1; \frac{1}{v_{013}}\right) - \frac{1}{4}H(0; u_1) H(0; u_3) \mathcal{H}\left(0, 1; \frac{1}{v_{013}}\right) +$ $\frac{1}{4}H(0, 0; u_1) \mathcal{H}(0, 1; \frac{1}{v_{vu}}) + \frac{1}{4}H(0, 0; u_3) \mathcal{H}(0, 1; \frac{1}{v_{vu}}) + \frac{1}{6}\pi^2 \mathcal{H}(0, 1; \frac{1}{v_{vu}}) \frac{1}{4}H(0, u_1)H(0; u_2)\mathcal{H}(0, 1; \frac{1}{v_{11}v}) + \frac{1}{4}H(0, 0; u_1)\mathcal{H}(0, 1; \frac{1}{v_{11}v}) +$ $\frac{1}{4}H(0, 0; u_2) \mathcal{H}\left(0, 1; \frac{1}{v_{111}}\right) + \frac{1}{6}\pi^2 \mathcal{H}\left(0, 1; \frac{1}{v_{112}}\right) - \frac{1}{4}H(0; u_1) H(0; u_2) \mathcal{H}\left(0, 1; \frac{1}{v_{111}}\right) +$ $\frac{1}{4}H(0, 0; u_1) \mathcal{H}(0, 1; \frac{1}{u_1}) + \frac{1}{4}H(0, 0; u_2) \mathcal{H}(0, 1; \frac{1}{u_1}) + \frac{1}{6}\pi^2 \mathcal{H}(0, 1; \frac{1}{u_1}) \frac{1}{n}H(0; u_2) H(0; u_3) H(1; \frac{1}{1}) + \frac{1}{n}H(0, 0; u_2) H(1; \frac{1}{1}) +$ $\frac{1}{2}H(0, 0; u_3) \mathcal{H}\left(1, 1; \frac{1}{m_{10}}\right) + \frac{1}{24}x^2 \mathcal{H}\left(1, 1; \frac{1}{m_{10}}\right) - \frac{1}{24}x^2 \mathcal{H}\left(1, 1; \frac{1}{m_{10}}\right) - \frac{1}{24}x^2 \mathcal{H}\left(1, 1; \frac{1}{m_{10}}\right) - \frac{1}{24}x^2 \mathcal{H}\left(1, 1; \frac{1}{m_{10}}\right)$ $\frac{1}{24}x^{2}\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right) - \frac{1}{2}\mathcal{H}\left(0; u_{1}\right)\mathcal{H}\left(0; u_{3}\right)\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right) + \frac{1}{2}\mathcal{H}\left(0, 0; u_{1}\right)\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right) + \frac{1}{2}\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right)\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right) + \frac{1}{2}\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right)\mathcal{H}\left(1, 1; \frac{1}{m_{11}}\right)$ $\frac{1}{2}H(0, 0; u_3) \mathcal{H}(1, 1; \frac{1}{m_{u_1}}) + \frac{11}{24}x^2 \mathcal{H}(1, 1; \frac{1}{m_{u_1}}) - \frac{1}{2}H(0; u_1) H(0; u_2) \mathcal{H}(1, 1; \frac{1}{m_{u_1}}) +$ $\frac{1}{2}H(0, 0; u_1) \mathcal{H}\left(1, 1; \frac{1}{u_{u_1}}\right) + \frac{1}{2}H(0, 0; u_2) \mathcal{H}\left(1, 1; \frac{1}{u_{u_2}}\right) + \frac{11}{24}\pi^2 \mathcal{H}\left(1, 1; \frac{1}{u_{u_2}}\right) \frac{1}{24}\pi^{2}\mathcal{H}\left(1, 1; \frac{1}{m_{W}}\right) + \frac{1}{2}H(0; u_{2})\mathcal{H}\left(0, 0, 1; \frac{1}{m_{W}}\right) + \frac{1}{2}H(0; u_{3})\mathcal{H}\left(0, 0, 1; \frac{1}{m_{W}}\right) +$ $\frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_3) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; u_1) H\begin{pmatrix} 0, 0, 1; \frac{1}{mu_1} \end{pmatrix} + \frac{1}{2}H(0; 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- dimensionally regulating thousands of separately divergent integrals
- final formula: 18 pages of so-called "Goncharov" polylogarithms

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Friday, 27th May

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Upon consulting with Goncharov about his polylogarithms, these 18 pages were found to simplify, [arXiv:1006.5703]:

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Friday, 27th May

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Vladimir A. Smirnov

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The Two-Loop Hexagon Wilson Loop in $\mathcal{N}=4$ SYM

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