Dark Matter and Colliders

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1 Introduction



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2 Producing WIMPs



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- **3 Producing Mediators**



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- 4 Summary

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Counter–examples: axions; Gravitinos; FIMPs; dark atoms; primordial black holes; keV neutrinos: not covered in this talk. Note: Proves that LHC does *not* "recreate conditions of the early universe"!

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• Indicates weak-scale $\chi\chi$ annihilation cross section: $\langle \sigma(\chi\chi \to any)v \rangle \simeq (2 \text{ to } 4.5) \cdot 10^{-26} \text{cm}^3 \text{s}^{-1}$

WIMPs and Early Universe

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Determining $\sigma(\chi\chi \to SM)$ allows probe of very early Universe, once χ has been established to be "the" DM particle! e.g. MD, Iminniyaz, Kakizaki, arXiv:0704.1590

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- Thermal WIMP: Only know total $\chi\chi \to SM$ cross section; contribution of specific final states $(e^+e^-, u\bar{u} + d\bar{d})$ not known
- $\Omega_{\chi}h^2$ determined from $\sigma(\chi\chi \to SM)$ near threshold $(T_F \simeq m_{\chi}/20 \Longrightarrow s \simeq 4m_{\chi}^2)$. At colliders need ≥ 3 body final state to get signature (e.g. $e^+e^- \to \chi\chi\gamma, \ q\bar{q} \to \chi\chi g$) \Longrightarrow typically need $\sigma(\chi\chi \to SM)$ at $s \gtrsim 6$ to $10m_{\chi}^2$!

"Model-independent" approach

Goodman et al., arXiv:1005.1286 and 1008.1783; Bai, Fox, Harnik, arXiv:1005.3797; Wang, Li, Shao, Zhang, arXiv:1107.2048; Fox, Harnek, Kopp, Tsai, arXiv:1103.0240 Parameterize χ interaction with relevant SM fermion through dim–6 operator; e.g. for hadron colliders:

 $\mathcal{L}_{\text{eff}} = G_{\chi,q} \, \bar{\chi} \Gamma_{\chi} \chi \, \bar{q} \Gamma_{q} q$

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$$\begin{split} \chi \text{ Majorana} &\Longrightarrow \Gamma_{\chi} \in \{1, \gamma_5, \gamma_{\mu} \gamma_5\} \\ \Gamma_q \in \{1, \gamma_5, \gamma_{\mu}, \gamma_{\mu} \gamma_5\} \\ \text{ If } \Gamma_{\chi}, \Gamma_q \in \{1, \gamma_5\} : \ G_{\chi,q} = m_q/(2M_*^3) \text{ (chirality violating!),} \\ \text{ else } \Gamma_{\chi} = 1/(2M_*^2) \text{ Rajamaran, Shepherd, Tait, Wijango, arXiv:1108.1196.} \end{split}$$

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LHC can only hope to be sensitive if χ couples to quarks at tree level. Two options:

s-channel \bar{q} g_{q} g_{q} g_{q} $G_{\chi,q} = \frac{g_{\chi}g_{q}}{m_{M}^{2}}$ χ g_{M} χ M $G_{\chi,q} = \frac{g_{\chi}g_{q}}{m_{M}^{2}}$ q χ q χ

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Case $\Gamma_{\chi} = \Gamma_q = \gamma_5$ also has poor LHC reach. Gives velocity–dependent interaction for $\chi p \rightarrow \chi p \Longrightarrow$ very poor reach in direct detection as well

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Bound on $\Lambda^2 \equiv 1/G_{\chi,q}$ depends on ratio $g_{\chi}/g_q!$
Monojet analysis is not model-independent!



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For *s*-channel: bound on 4q contact interaction stronger than bound from monojet searches, unless $g_{\chi} \gg g_q!$

Bounds on Λ

S. Belwal, MD, J.S. Kim, in preparation



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For perturbative couplings: search for off-shell mediator $\rightarrow \chi \chi$ is hopeless! E.g. SUSY: signal for $\tilde{\chi} \tilde{\chi} j$ is *much* smaller than $Z (\rightarrow \nu \bar{\nu}) j$ background, even for 100 GeV higgsino-like $\tilde{\chi}$ (Baer, Mustafayev, Tata, arXiv:1401.1162)

DM and Light (Gauge) Bosons

(At least) 3 kinds of WIMP models require light ($m \le$ few GeV) (gauge) bosons U:

● <u>MeV DM</u>: Suggested as explanation of 511 keV line (⇒ slow e^+) excess from central region of our galaxy (Boehm et al., astro-ph/0309686). Should have $m_{\chi} \leq 10$ MeV (γ constraints)

 $\Rightarrow m_{\chi} \leq m_U \leq 200 \text{ MeV to mediate } \chi\chi \rightarrow e^+e^-$; fixes $g_{U\chi\chi}g_{Ue^+e^-}/m_U^2!$ (Unless $2m_{\chi} \simeq m_U$.)

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• PAMELA/FermiLAT inspired TeV DM: Needs light boson for Sommerfeld enhancement (e.g. Arkani-Hamed et al., arXiv:0810.0713(4)) ($\chi\chi \rightarrow UU \rightarrow 4l$ is also somewhat less constrained by γ spectrum than $\chi\chi \rightarrow 2l$.)

• DAMA/CoGeNT inspired few GeV DM: Needs light mediator to achieve sufficiently large $\sigma_{\chi p}$. (2 different mediators for isospin violation to evade bounds: Cline, Frey, arXiv:1108.1391)

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 $U\chi\chi$ coupling may well be large.

Signatures of light gauge bosons

 $\frac{\text{If } m_U > 2m_{\chi}:}{\text{tag, e.g. } e^+e^- \to \gamma U \to \gamma + \text{ nothing.}} \text{ Is invisible } \Rightarrow \text{ need extra}$

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- Instrumental backgrounds (not from e^+e^- annihilation) seem large

Sensitivity at B-factories (100 fb⁻¹)



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Also, KLOE-2 performed search, mostly for $\phi \rightarrow U\eta$: no signal. arXiv:1107.2531

A1 results



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- HLS theorem, relation to superstrings: don't single out weak scale.

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- In simplest, *R*-parity invariant scenario: lightest superparticle LSP is stable: satisfies one condition for DM candidate!

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- but DM-allowed regions of parameter space do exist even in constrained models!
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- Note: DM-allowed region of $(m_0, m_{1/2})$ plane of cMSSM depends on $A_0, \tan \beta!$

Strongly interacting sparticles have biggest cross sections; may decay via long decay "cascades". Example:

$$gg \to \tilde{g}\tilde{g} \to (\tilde{b}_1\bar{b}) (\tilde{u}_L\bar{u}) \to (\tilde{\chi}_2^0 b\bar{b}) (\tilde{\chi}_1^+ d\bar{u}) \to (\tilde{\chi}_1^0 e^+ e^- b\bar{b}) (\tilde{\chi}_1^0 c\bar{s} d\bar{u})$$

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 $\mathcal{O}(100)$ searches have been performed, but no signal has been found.

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- In pMSSM10: $m_{\tilde{\chi}^0_1} \simeq 50~{\rm GeV~still~ok!}$ de Vries et al., arXiv:1504.03260

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- Most interesting to me: Predict $\Omega_{\chi}h^2$, compare with observation: Constrain very early universe!

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- Scenarios with new light gauge bosons with suppressed couplings to SM fermions are now being probed at low-E colliders, fixed-target expts.
- Absence of missing E_T signal at LHC is disappointing, but plenty of parameter space in reasonably well motivated WIMP models left to explore