

# Our milestones:

1. Why do ATLAS and CMS look like they look like today?
- 2. Tools used in the analyses**
3. Some highlights, and comments
4. What next?

# Our Master Equation

Event rates (absolute, relative, differential)  
Stat vs syst errors, backgrounds from data or MC?  
Resolution, Energy Scale, Signal Significance

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon L}$$

Proton-Proton Luminosity  
uncertainty a few %;  
eliminated in ratios

**Experimental issues** : Triggers, reconstruction, isolation cuts, low- $p_T$  jets (jet veto)  
acceptance, efficiency determination

**Theoretical issues** :  $p_T$  distributions at N<sup>n</sup>LO + resummation;  
differential calculations for detectable acceptance.

$$\sigma_{\text{theo}} = PDF(x_1, x_2, Q^2) \otimes \hat{\sigma}_{\text{hard}}$$

constrain, define uncertainties

HO calculations,  
implemented in MC?

Goal : test SM (in)consistency :  $\sigma_{\text{exp}} \pm \Delta_{\text{exp}} \stackrel{?}{=} \sigma_{\text{SM}} \pm \Delta_{\text{th}}$

# Efficiencies and Acceptance

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\varepsilon L}$$

# Efficiencies and acceptances

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon L}$$

Identification eff.

Number of “reconstructed” objects/events,  
which have passed the ID criteria

Number of “all reconstructed” objects/events

Trigger eff.

Number of “detectable” objects/events,  
which have been triggered on

Number of “all detectable” objects/events

$$= \epsilon_{\text{ID}} \cdot \epsilon_{\text{RECO}} \cdot \epsilon_{\text{TRIG}} \cdot A$$

Acceptance

Reconstruction eff.

Number of “detectable, triggered” objects/events,  
which have been reconstructed

Number of “detectable/triggered” objects/events

Number of “detectable” objects/events

Number of “all produced” objects/events

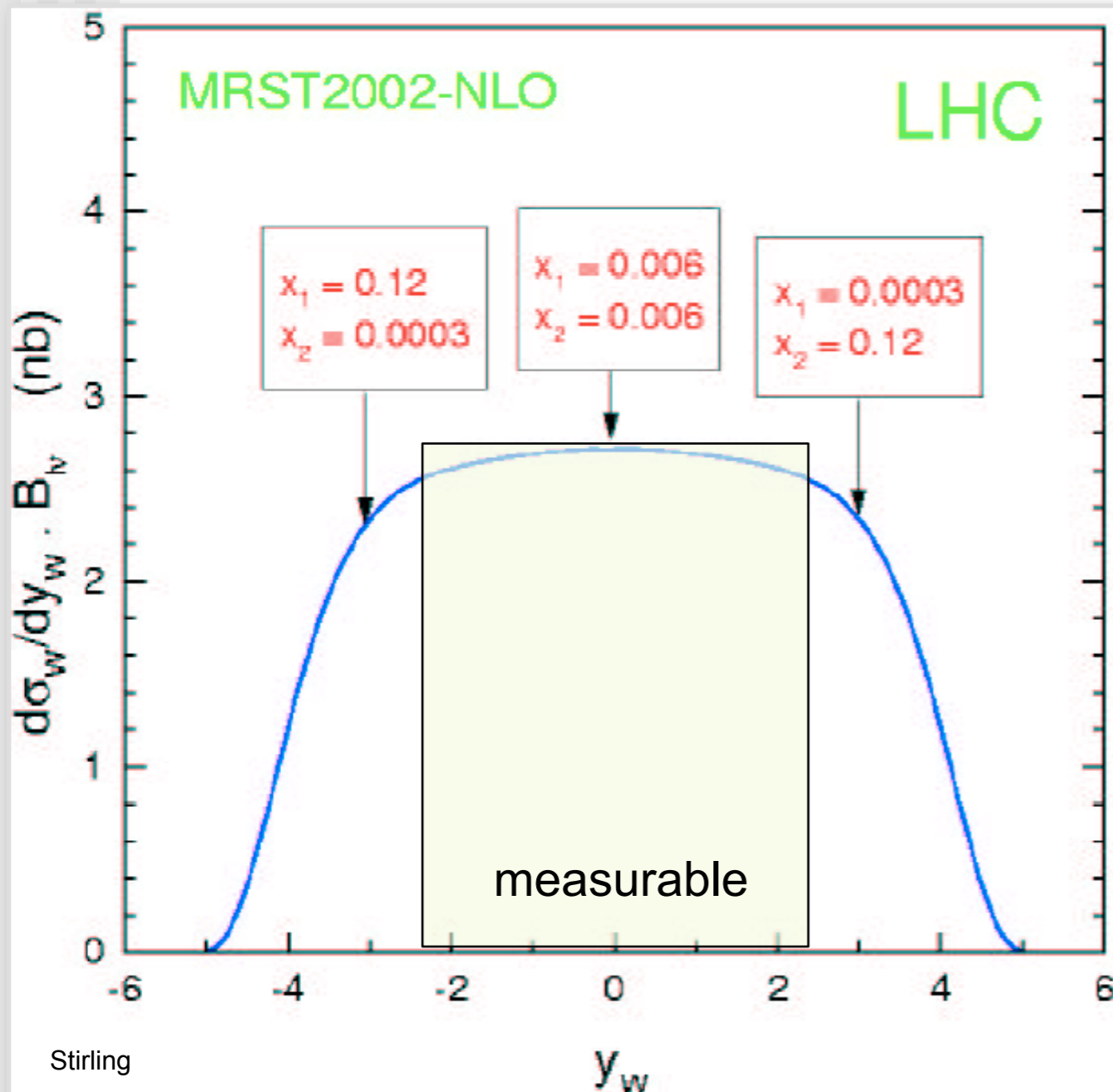
**example, from MC:**

N(muons) with  $p_T > 10$  GeV and  $\eta < 2$

N(all generated muons)

# Issue of acceptance...

## Example: W or Z production



It is a convolution of the acceptances for the leptons

Do we really want to correct for acceptance?

### Pros:

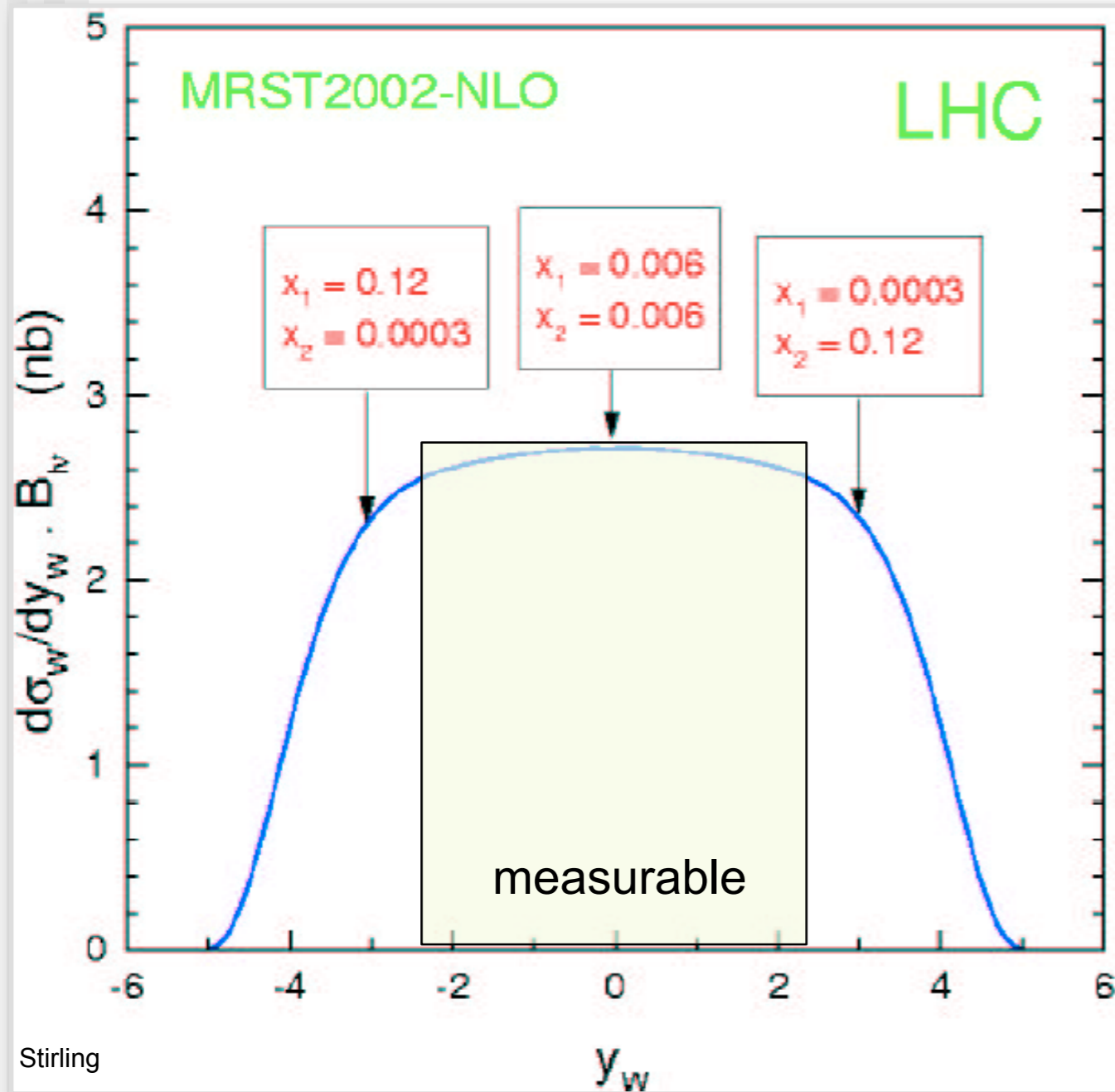
- The cross section measurement can be directly compared to other (corrected) measurements, from other exps
- The measurement can be compared to theory predictions which cannot be obtained for arbitrary acceptance

### Cons:

- The measurement becomes model dependent
- we introduce a systematic error, eg. because of uncertain extrapolation to full acceptance

# Issue of acceptance...

## Example: W or Z production



if we want **precise measurement** of Luminosity or some parameter in hard-interaction cross section, it is essential to have HO calc.

restricted to measurable acceptance

⇒ avoid **extrapolation errors**

eg. from extrapolation to large  $y_W$  where uncertainties from pdfs are large!

... fortunately, nowadays more and more fully differential calculations are available.... thus it becomes possible to calculate “EXACTLY” what is measured.... and such kind of so-called “fiducial cross sections” have indeed been published.

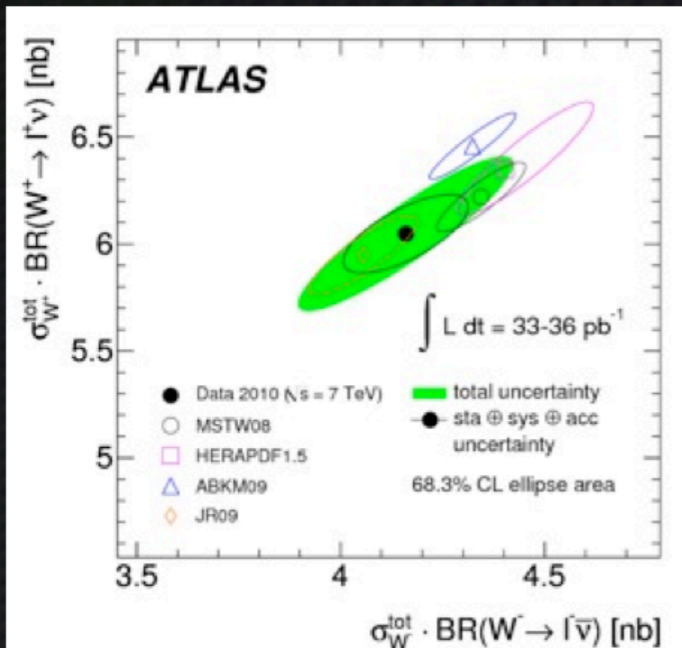
# Fiducial cross sections...

## Fiducial W and Z Cross Sections

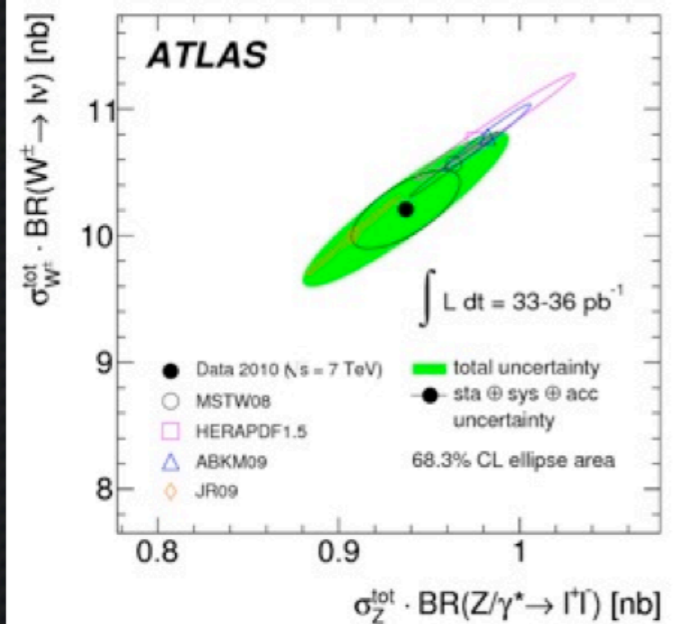
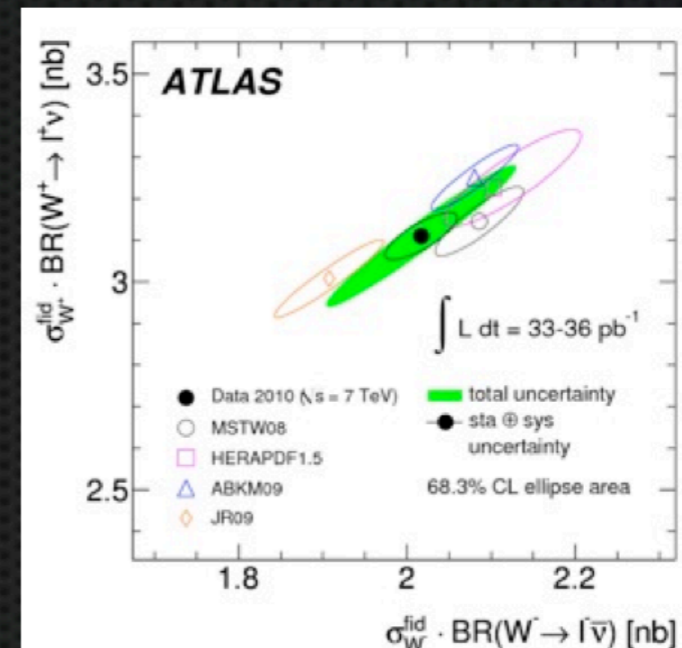
Phys. Rev. D85 (2012) 072004

$\sigma_{\text{Total}}$

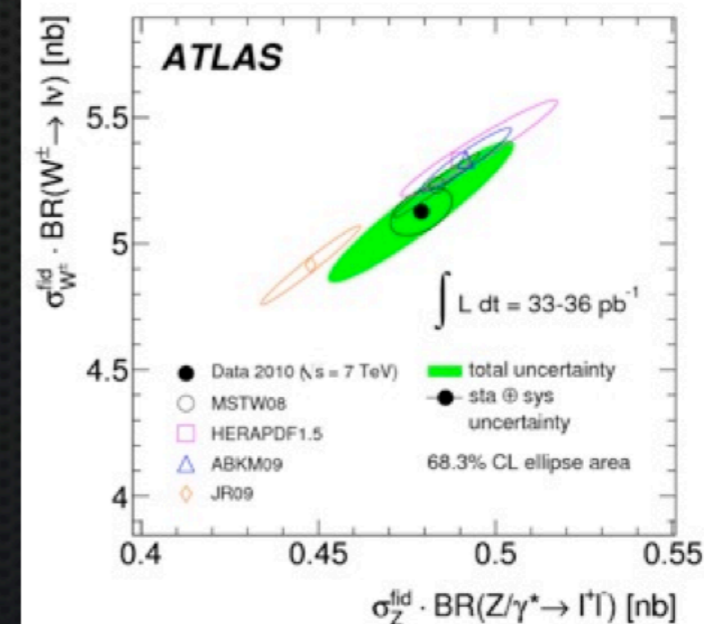
$\sigma_{\text{Fiducial}}$



**$W^+$  versus  $W^-$**



**$W^\pm$  versus Z**



**Luminosity 3.4%**

**Some differentiation between PDF sets already observed**

**JR09 seems to be the most discrepant**

Particle Physics in the LHC Era -- Zurich 2013 -- Joao Guimaraes

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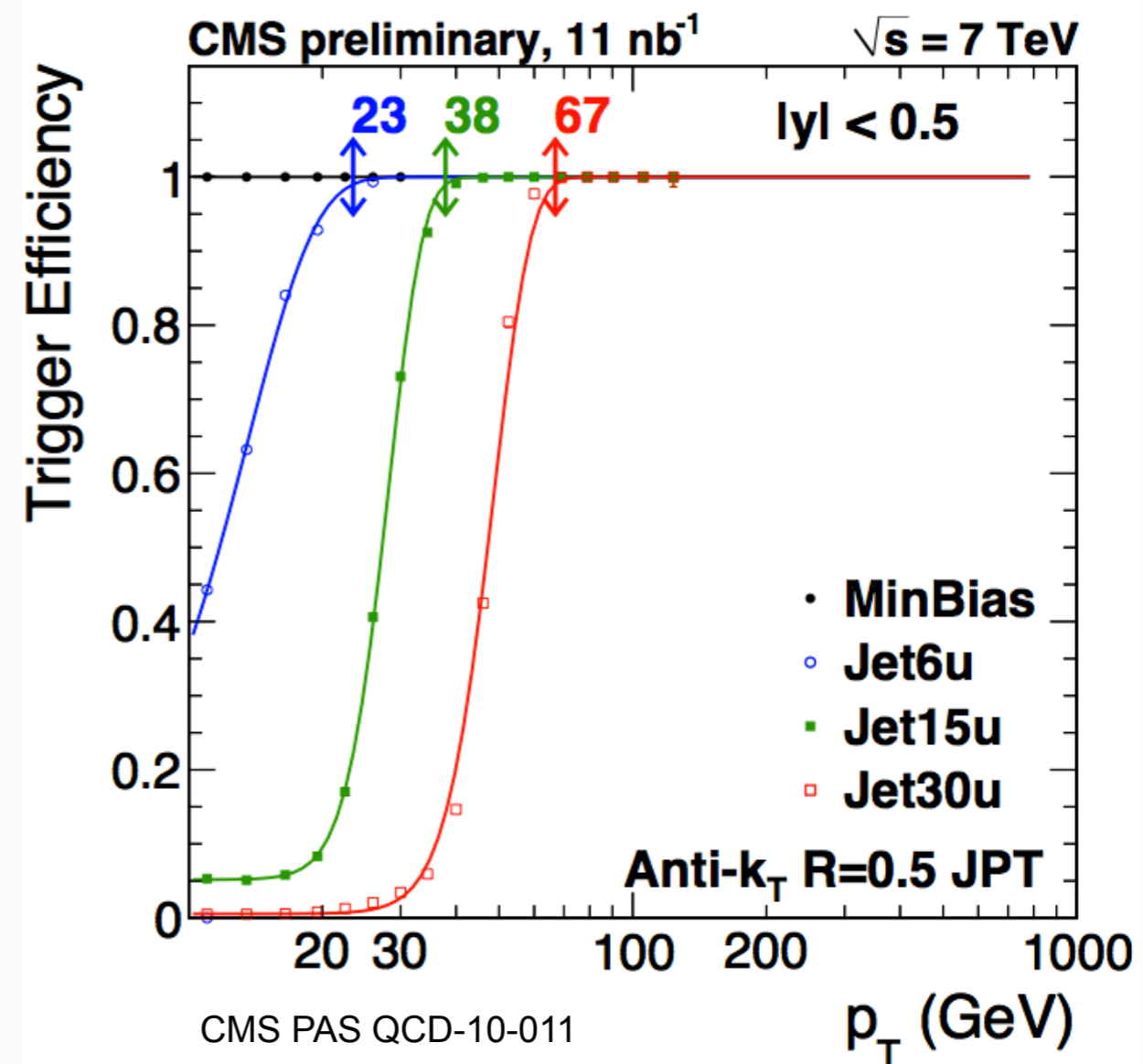
# Trigger efficiencies

- **Usual recipe:** try to have a “more inclusive” trigger, where you “know” that it is “100% efficient”, and calculate rate w.r.t. this one
- Example: trigger rate for a Jet Trigger with  $E_T > 15$  GeV:

$$\epsilon_{\text{TRIG}} = \frac{N(\text{Jet15 Trigger AND MinBias Trigger})}{N(\text{MinBias Trigger})}$$

- **Minimum Bias Trigger:** a minimal set of selection criteria are applied, eg. a few hits in the beam scintillation counters
- compare, eg. to **Zero Bias Trigger**
- Then, the efficiency of a higher Jet ET trigger, eg. 30 GeV, can be found from:

$$\epsilon_{\text{TRIG}} = \frac{N(\text{Jet30 Trigger AND Jet15 Trigger})}{N(\text{Jet15 Trigger})}$$



- Typically, apply selection cuts only above a  $p_T$  where your trigger is >99% efficient!

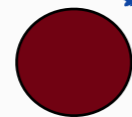
# The Tag & Probe Method

## Useful to measure efficiencies from data

- trigger eff, reconstruction eff., identification eff.
- eg. **single muon trigger eff.**: what is the fraction of reconstructed muons, which would also have been triggered on?
- eg. **electron ID eff**: what is the fraction of reconstructed electron candidates, which also pass a tight isolation criterium?

$$\epsilon_{ID} = \frac{N(\text{ Probes which pass further criteria})}{N(\text{ all tags})}$$

**Probe Object:** “loosely” selected:  
now apply further criteria



**Tag Criterion:** eg. di-lepton system close to invariant mass of Z or J/Psi; or a very pure W candidate: one isolated lepton, large MET, no further activity in the event, transverse mass > X

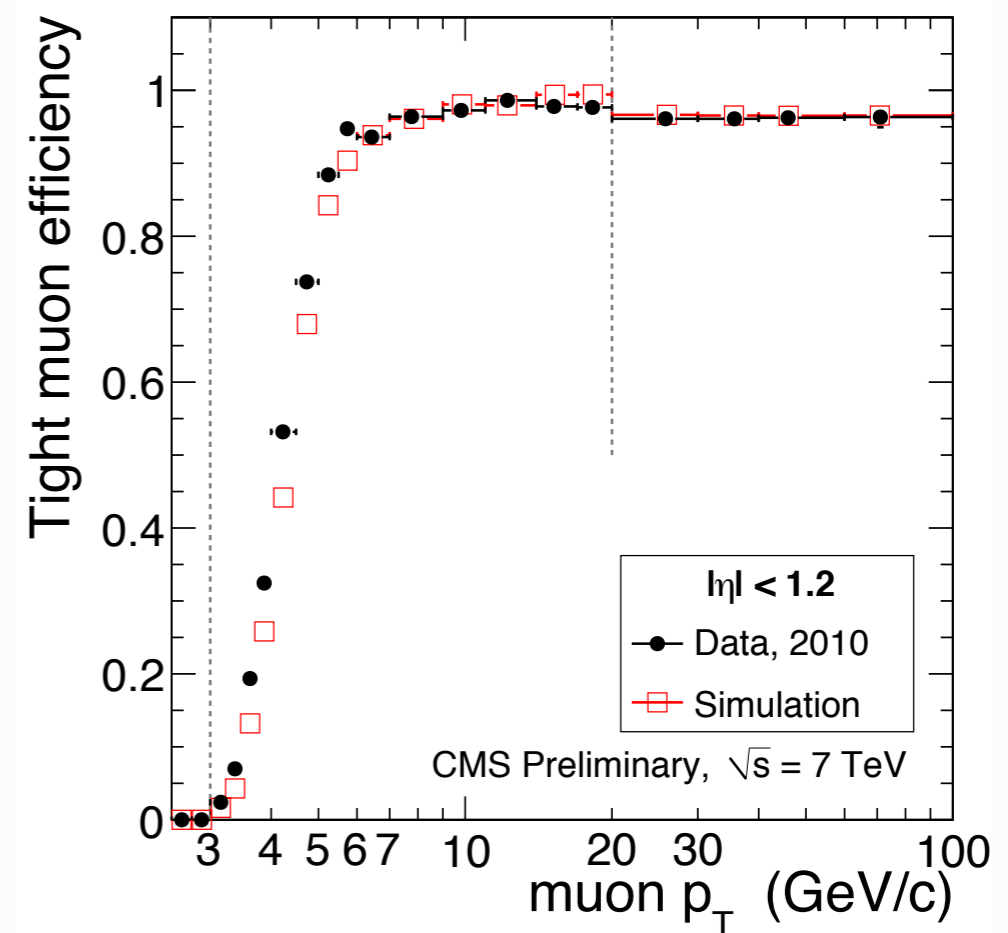
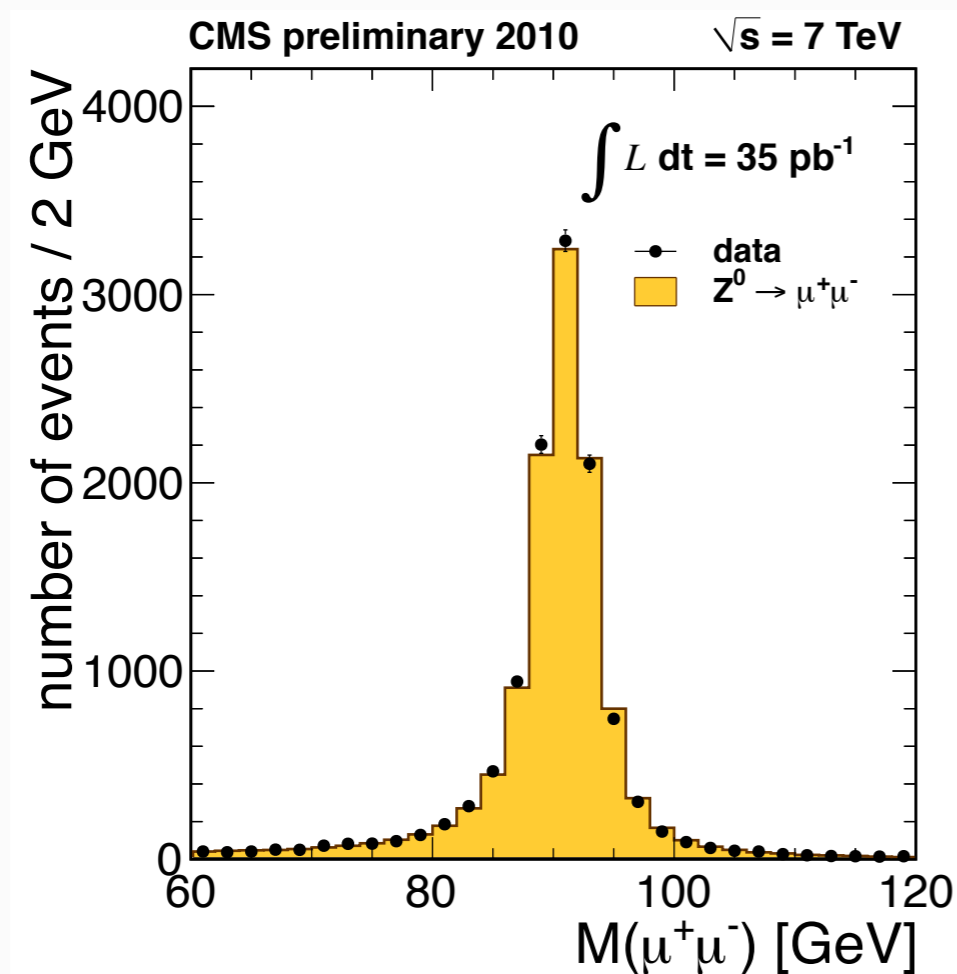


**Tag Object:** “tight” selection applied:  
defined the tag together with the additional criterium above

# Tag & Probe....

## Careful:

- make sure no background left, or subtract it
- make sure no correlations introduced
- Apply same method in data and MC.  
In MC: compare to “True Eff.” and if necessary apply (hopefully small) additional correction factors, if some bias is observed

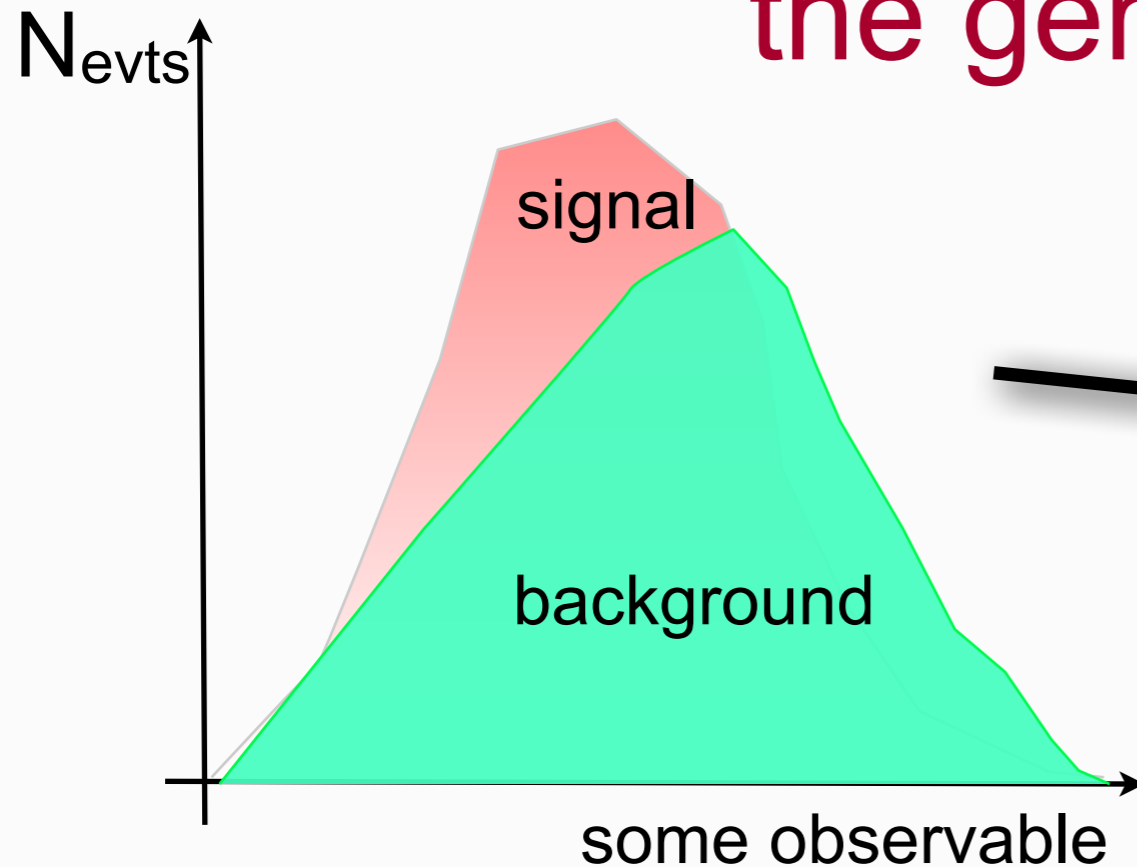


see eg. <https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsMUO>

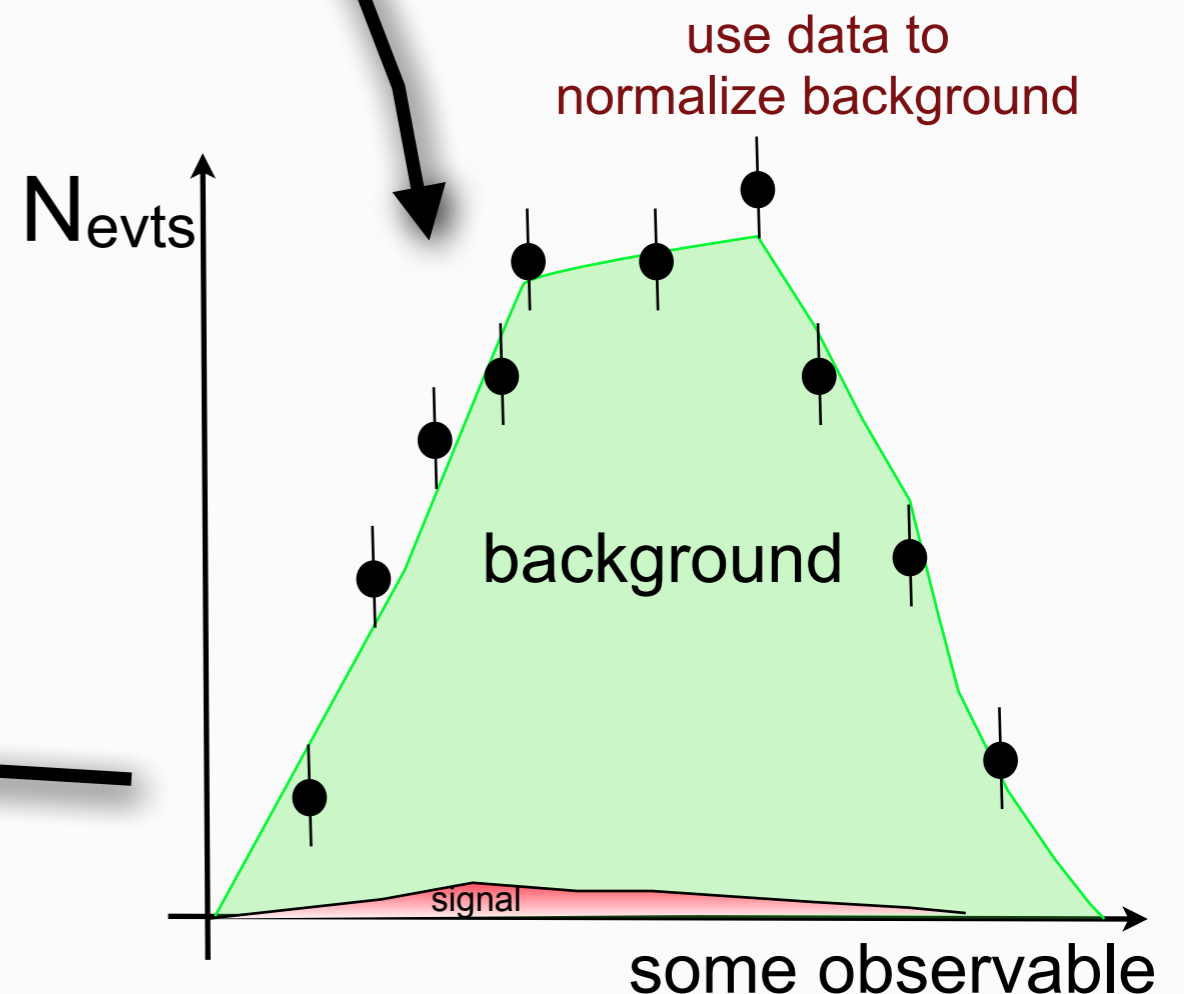
# Backgrounds

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\varepsilon L}$$

## the general idea



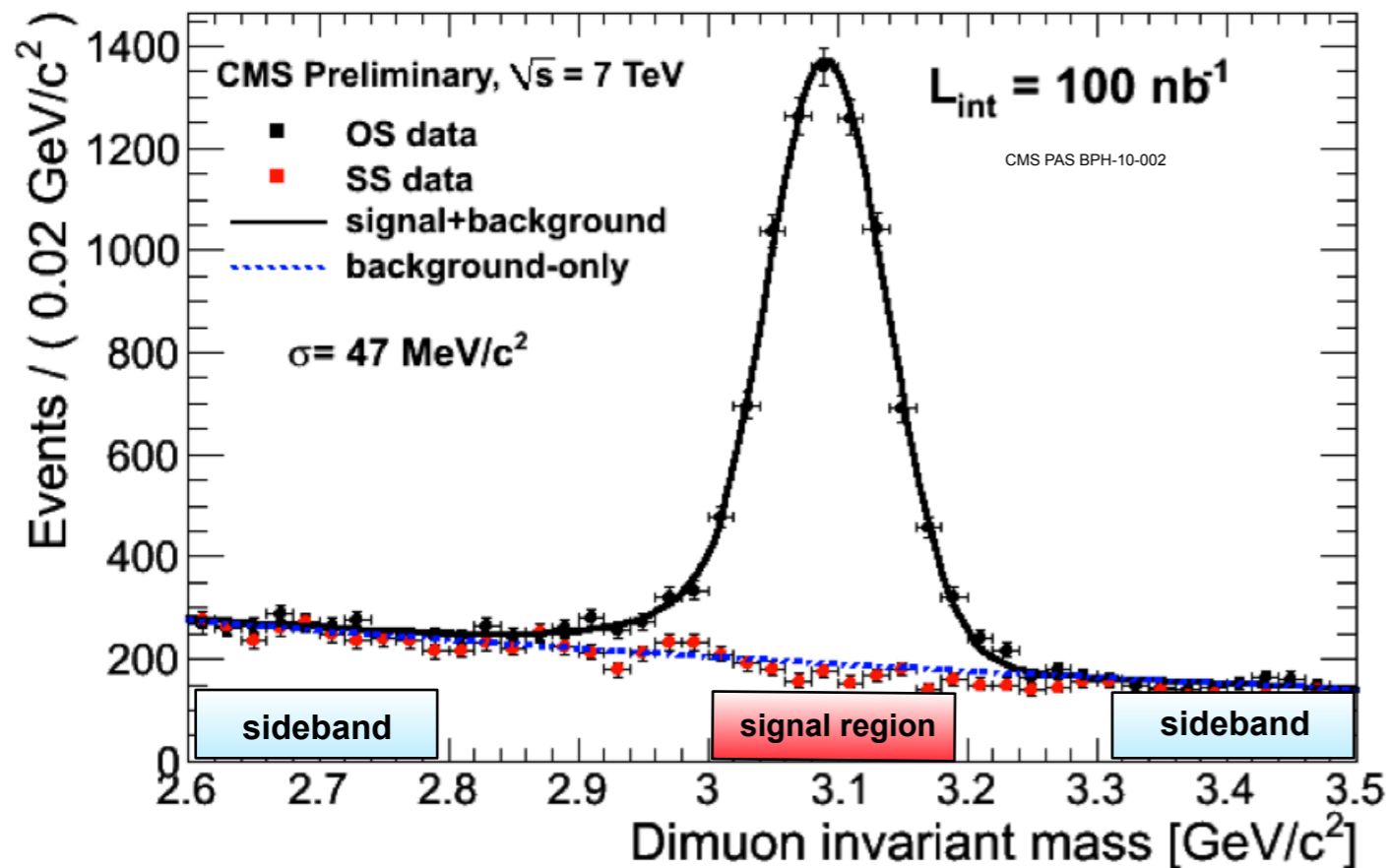
invert cuts :  
from signal enhancement to  
background enhancement



going back:  
use theory to compute  
change in background  
when inverting cuts, i.e.  
a ratio (**see later**);  
or use some well-  
motivated extrapolation,  
data only (**see next  
slide**)

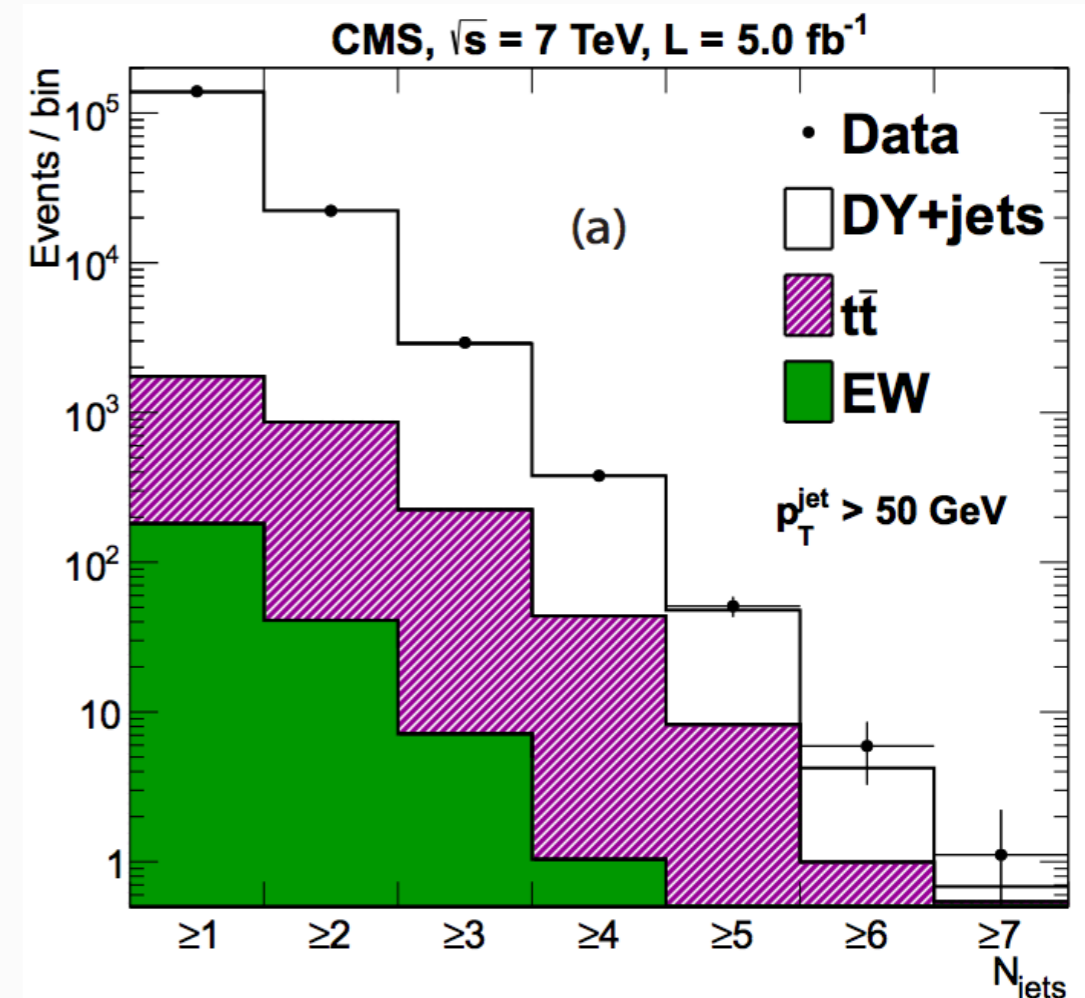
# Examples of (fully) data-driven bkg predictions

the trivial case: sidebands of a mass peak



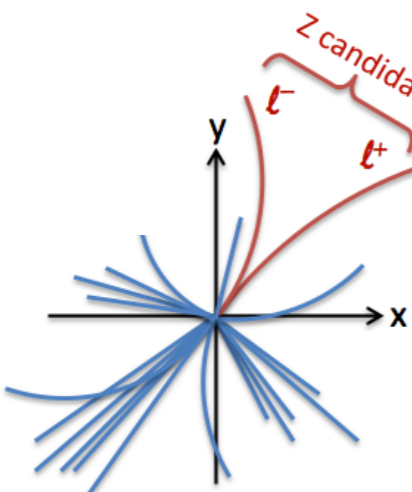
$t\bar{t}$  bckg to Z+jets:

take an e-mu dilepton sample with the same kinematic selection as for the same-flavour (ee, mumu) dilepton sample. This is almost pure in  $t\bar{t}$ .

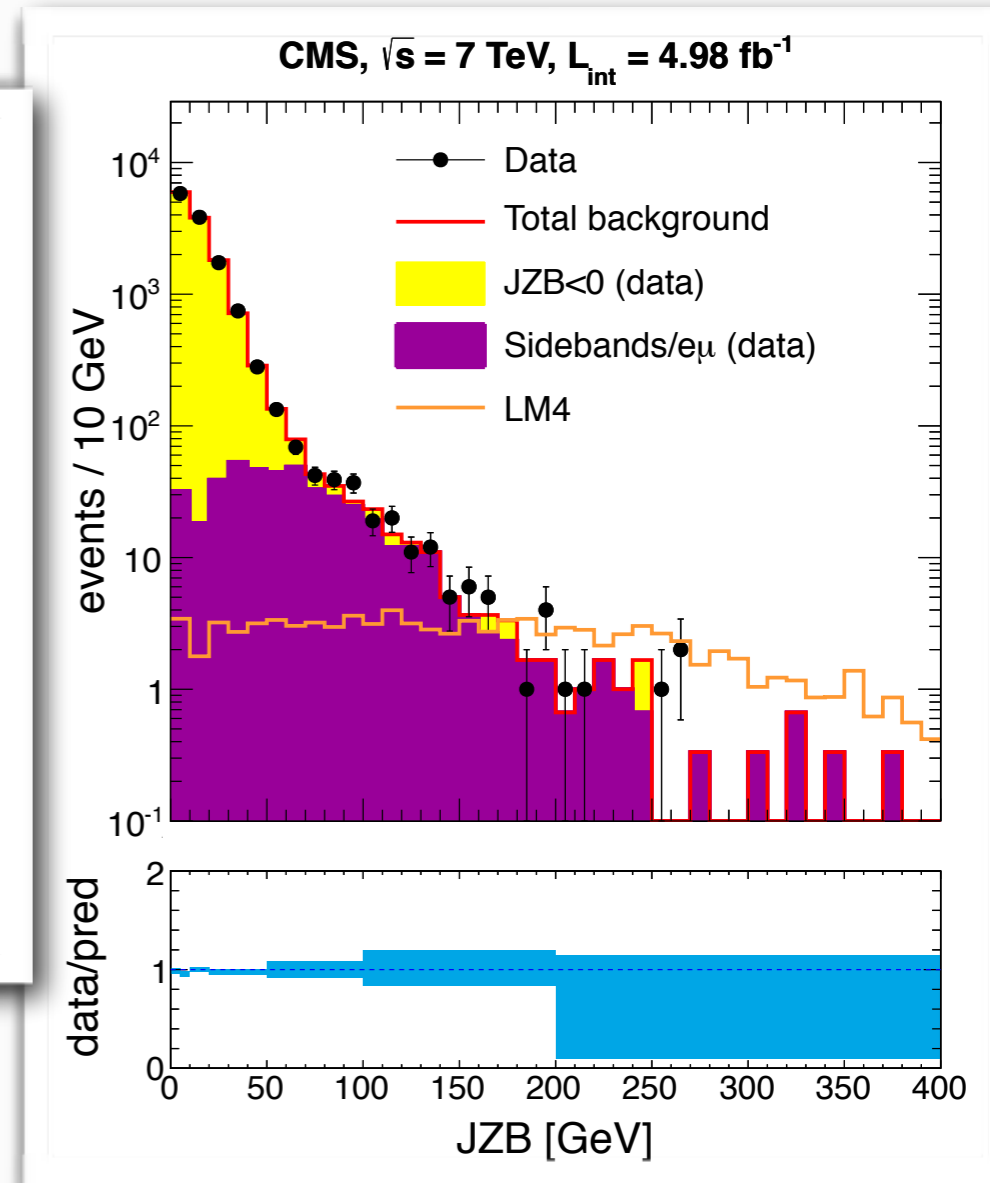
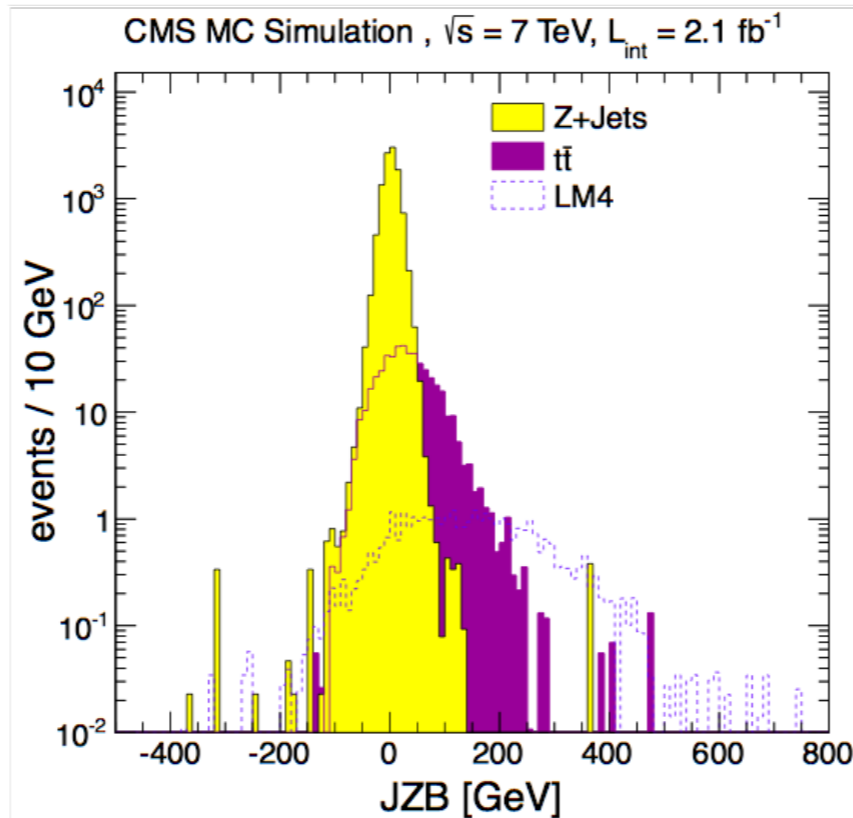


# Examples of (fully) data-driven bkg predictions

Search for topologies with jets, MET, Z :



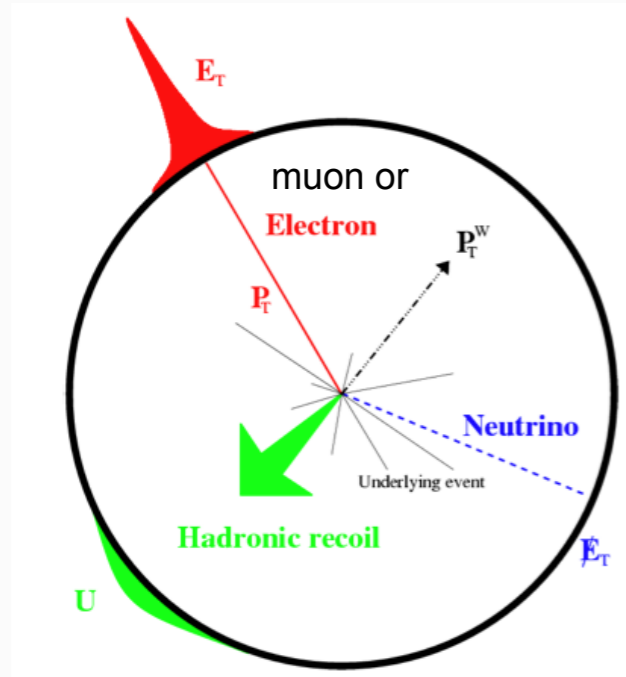
$$JZB \equiv \left| \sum_{\text{recoil}} \vec{p}_T \right| - \left| \vec{p}_T(\ell\ell) \right|$$



- MET in Z+jets events: fake, from jet mis-measurements
- Z+jets bkg on positive JZB: from negative JZB part
- top backg : use opposite-flavour events

arXiv:1204.3774

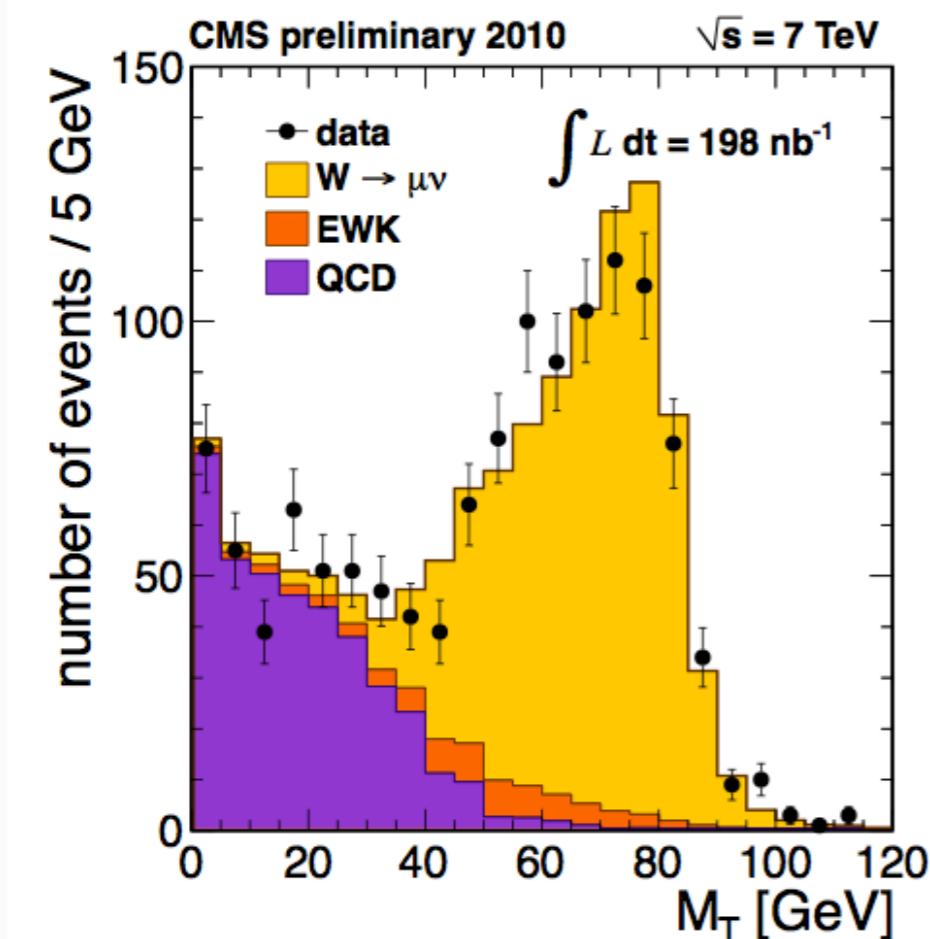
# Often encountered: W selection



W: decay to charged leptons

- high- $p_T$
- isolated
- $E_{T,miss}$  (from neutrino)

$$\text{transverse mass: } M_T = \sqrt{2p_T(\mu)E_T(1 - \cos(\Delta\phi_{\mu,E_T}))}$$



CMS PAS EWK-10-002

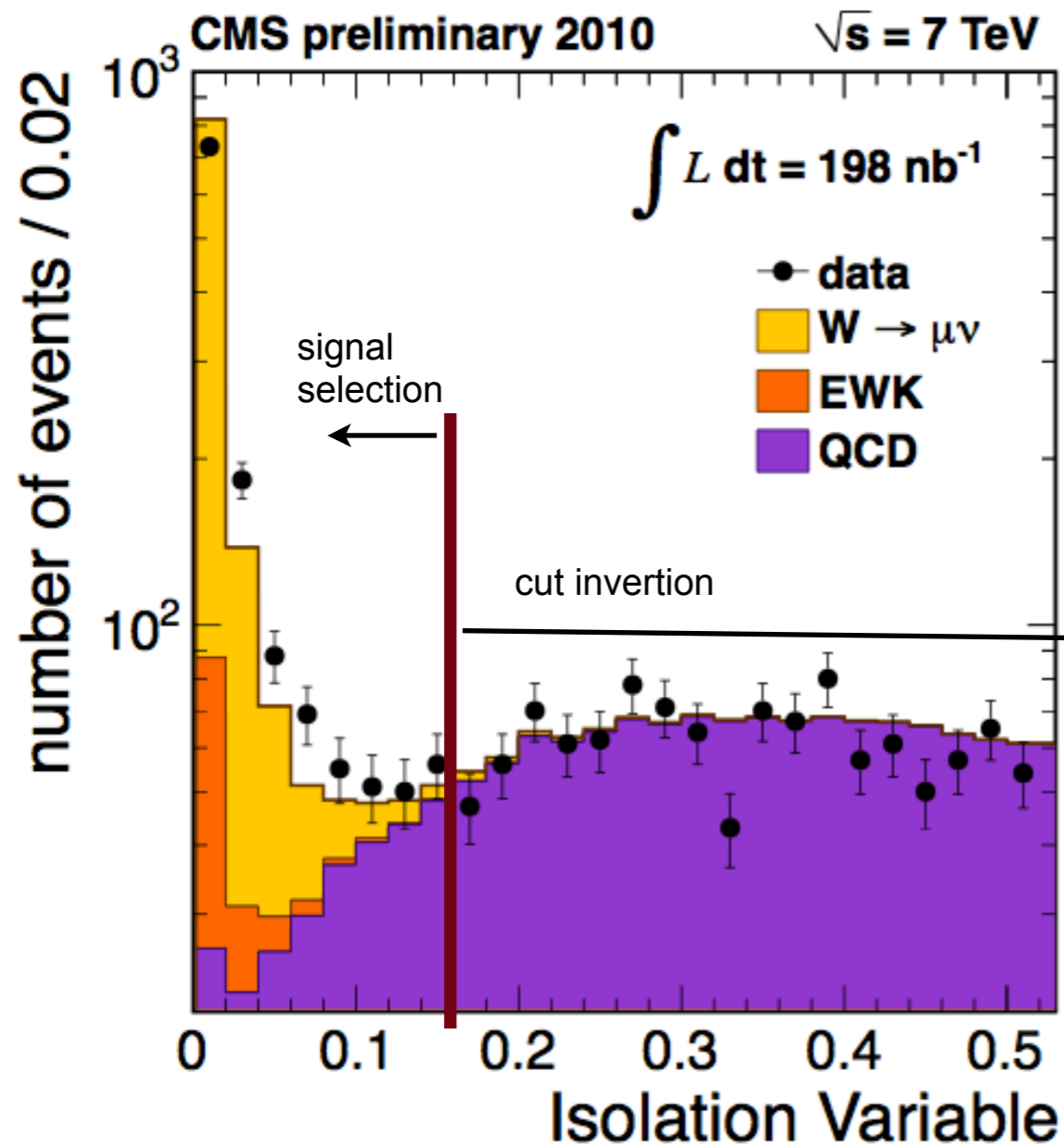
after cut on important selection variable,  
the relative isolation:

$$I_{\text{comb}}^{\text{rel}} = \left\{ \sum (p_T(\text{tracks}) + E_T(\text{em}) + E_T(\text{had})) \right\} / p_T(\mu)$$

in cone  $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} < 0.3$   
around the muon

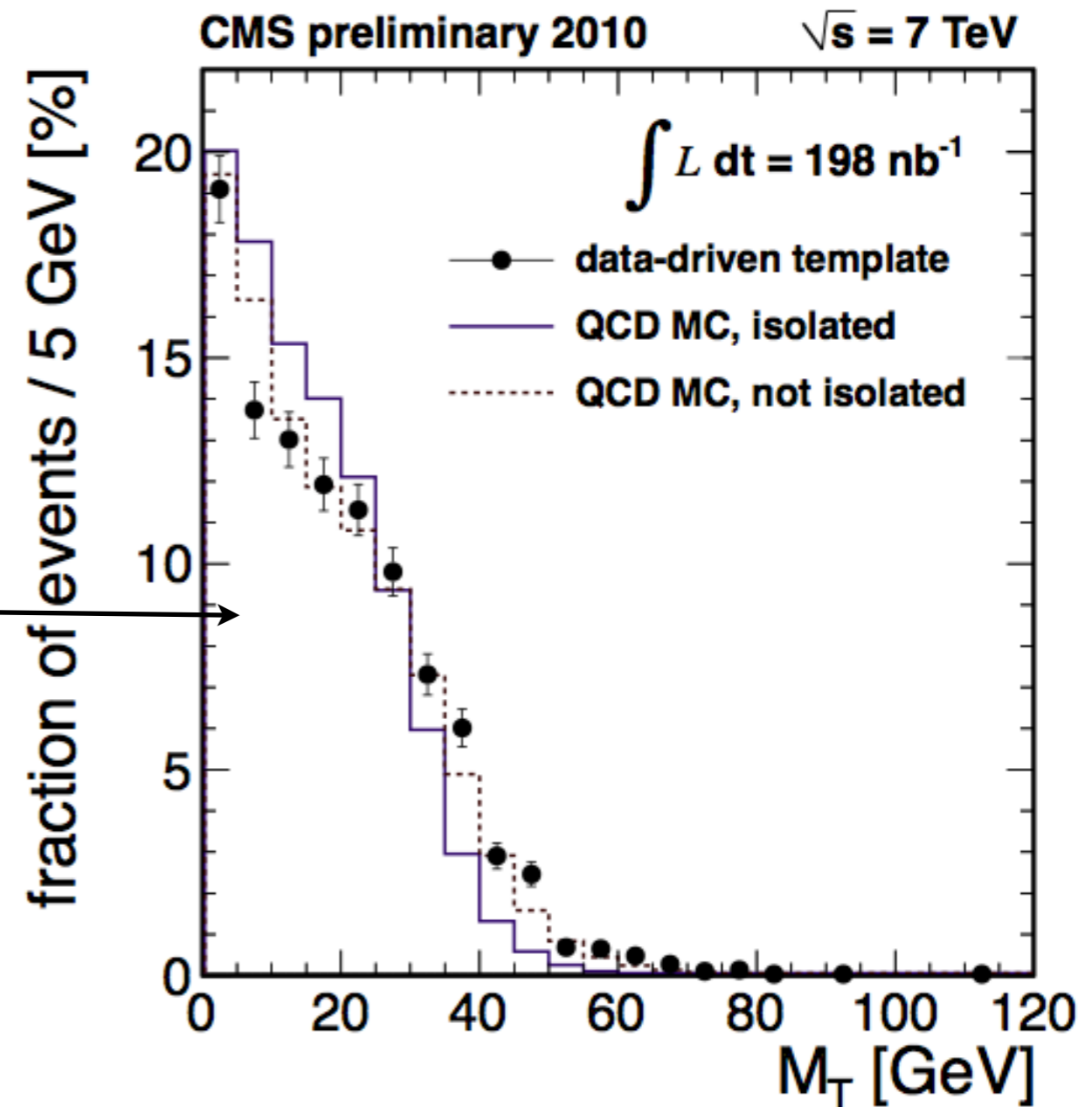
# W selection: cont.ed

CMS PAS EWK-10-002



$p_T > 20 \text{ GeV}/c$  in  $|\eta| < 2.1$

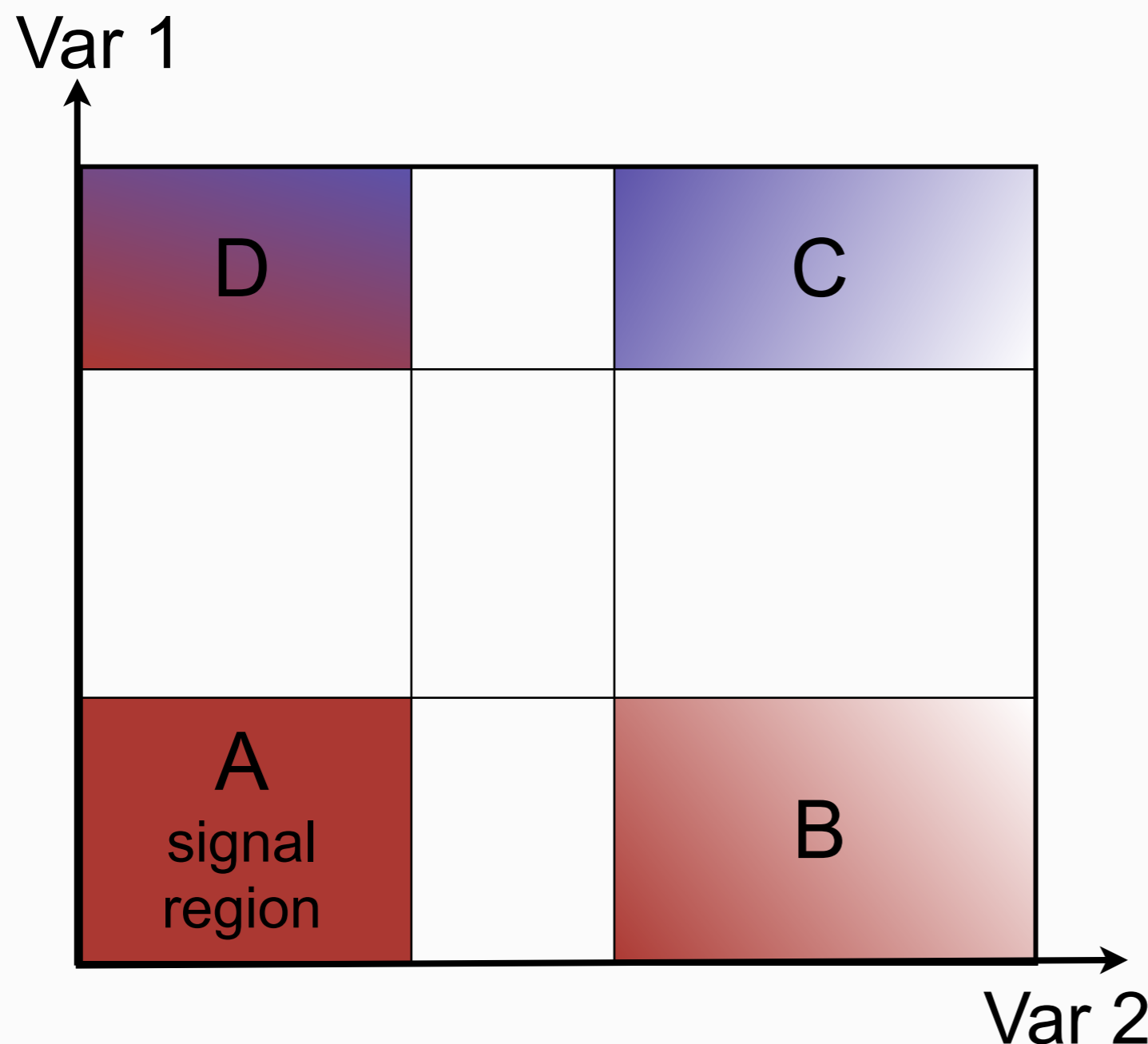
QCD bkg: mostly b-decays



take this shape for fit  
to  $M_T$  distribution

# The “ABCD” method

- find two variables, which characterize the events of interest
- A=signal region, B,C,D: background regions
- hypothesis of **un-correlated** variables:  
background shape in AD sector is the same as in BC sector



If this hypothesis is true,  
and **no** signal contamination  
in B,C,D:

estimate for background in  
signal region is

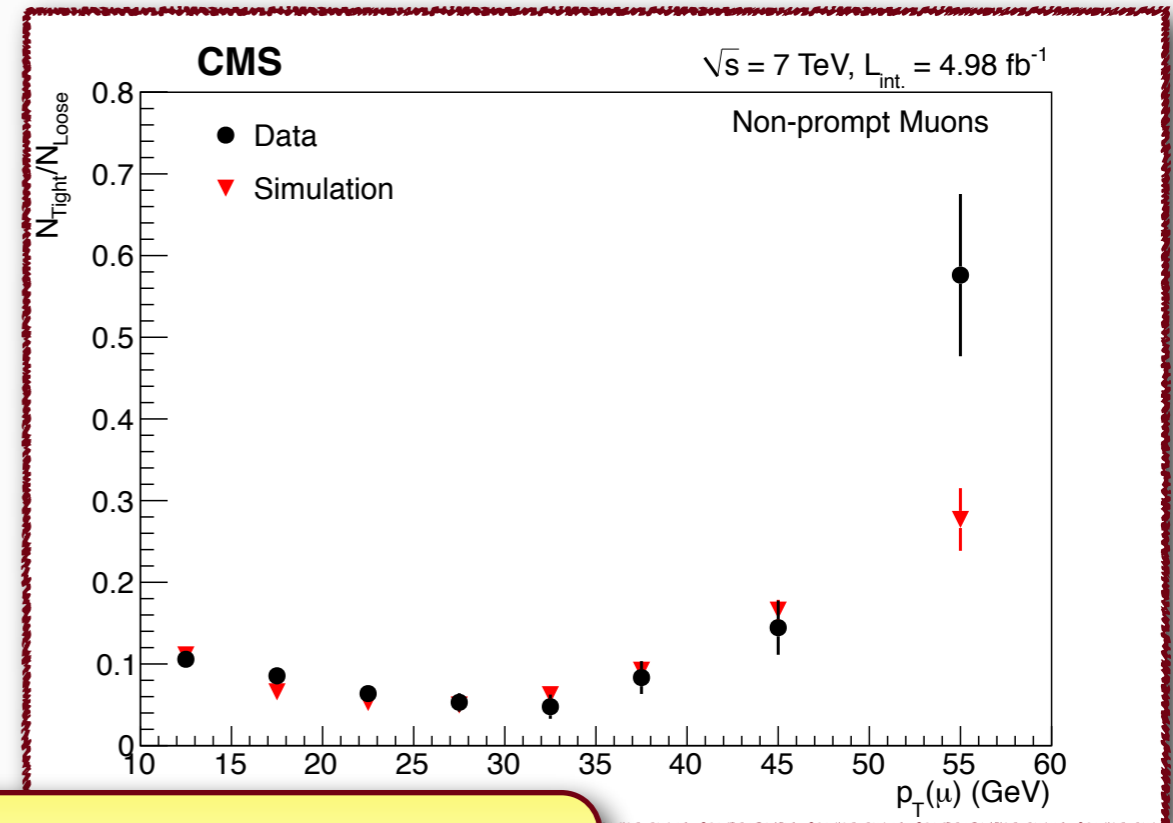
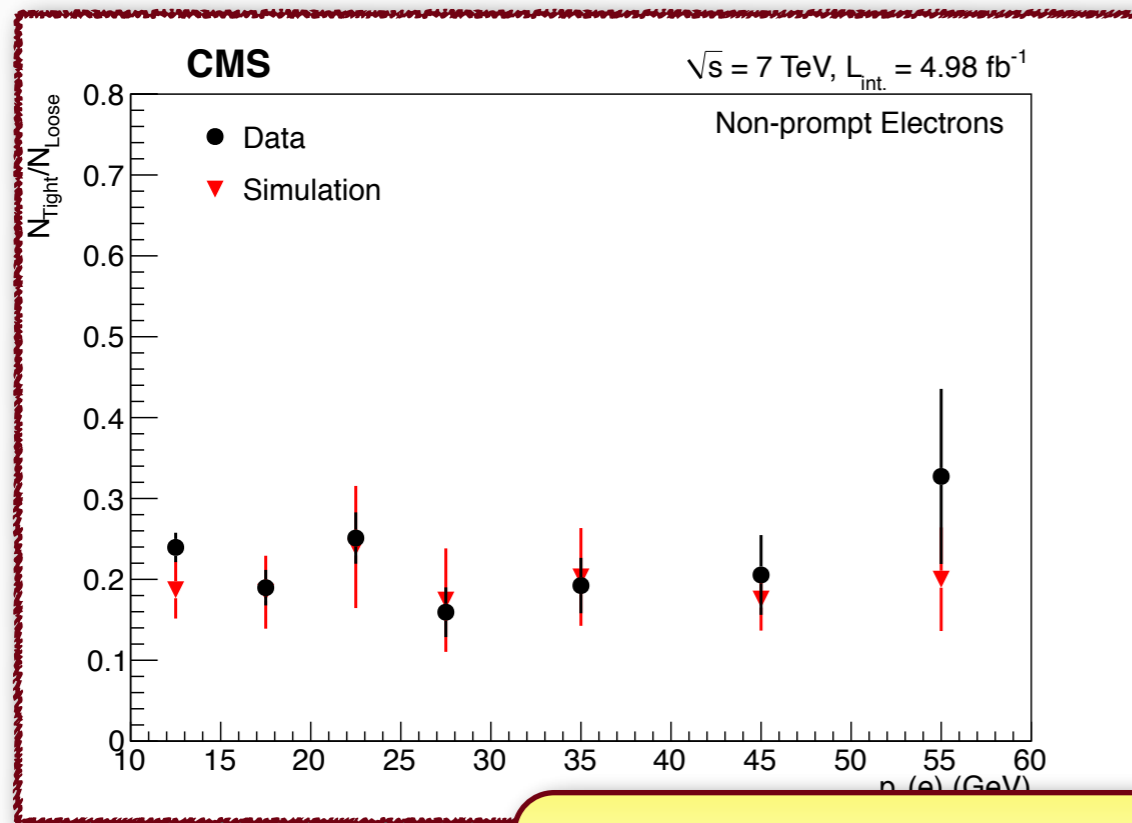
$$\frac{N(A)}{N(D)} = \frac{N(B)}{N(C)} \Rightarrow N(A)$$

from the counted number  
of events  $N(B,C,D)$  in the  
background regions.

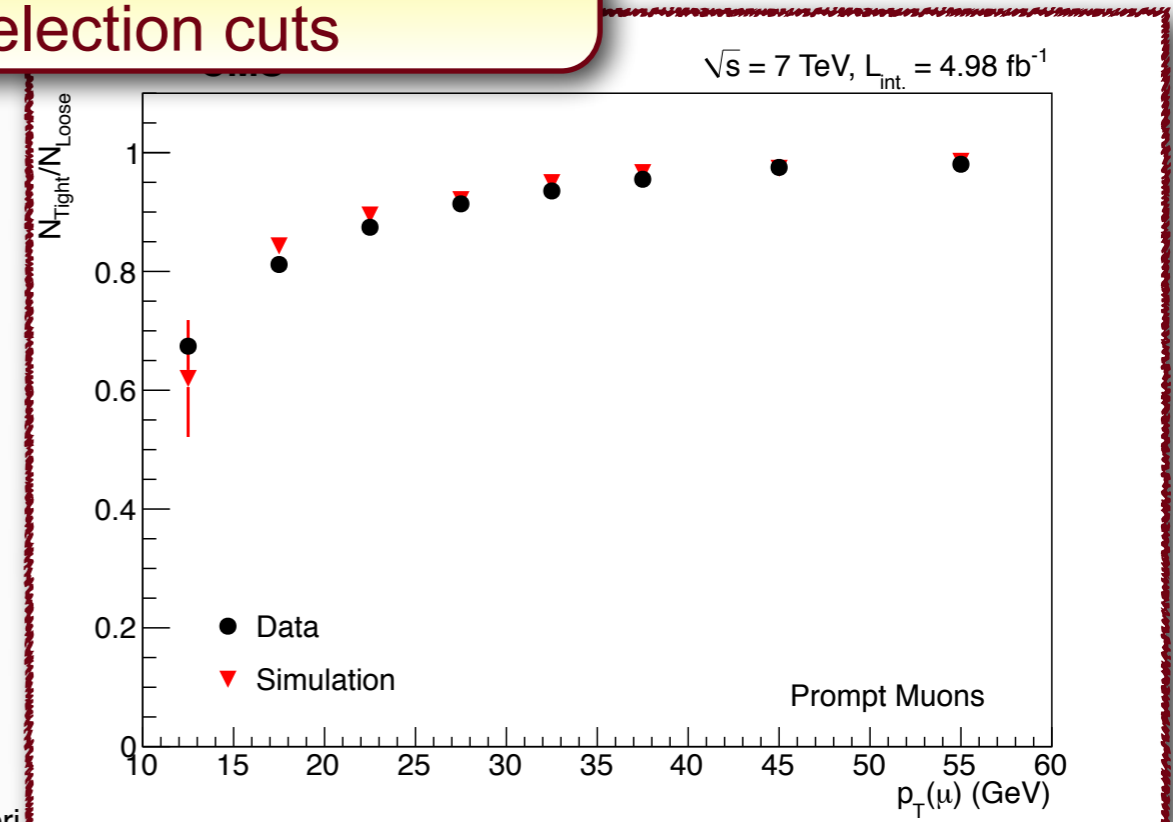
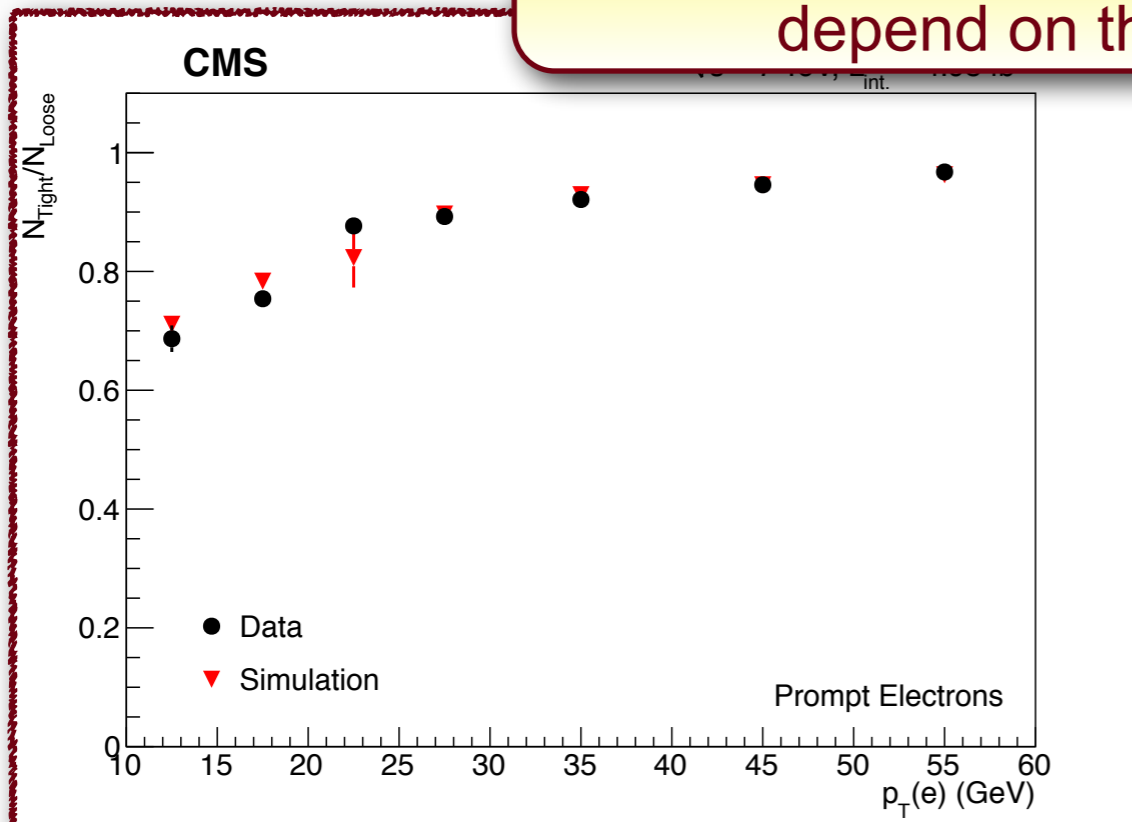
- What is probability that eg. a jet is mis-identified as an “isolated” lepton?
  - Important to know for leptonic analyses, especially in case of search for rare “multi-lepton” signatures, or e.g., same-sign dilepton SUSY search
  - even if tight isolation requirements are applied, the probability of faking is not zero, and a small number, multiplied with the huge cross section of multi-jet production, can still lead to a sizeable background
  - difficult (impossible?) to trust the simulation on this faking probability, rather try to get it from data?
- “Standard Method”:
  - “Fakeable Object method”, or “Tight-To-Loose Ratio”
    - Idea** : define two selection steps, one with **LOOSE** criteria, and one with **TIGHT** criteria (eg. on isolation)
    - determine the “**fake ratio**”, or “probability for a jet to fake a lepton” from the ratio of tightly to loosely selected objects, in a control sample that should not have any prompt leptons (eg. multi-jet sample)
    - determine this number as function of basic kinematics ( $p_T$ , rapidity)
    - apply it to a MC background simulation, or at a preselection level, to determine this fake background on the final selection level

# Fake rates: example, tight-loose ratios

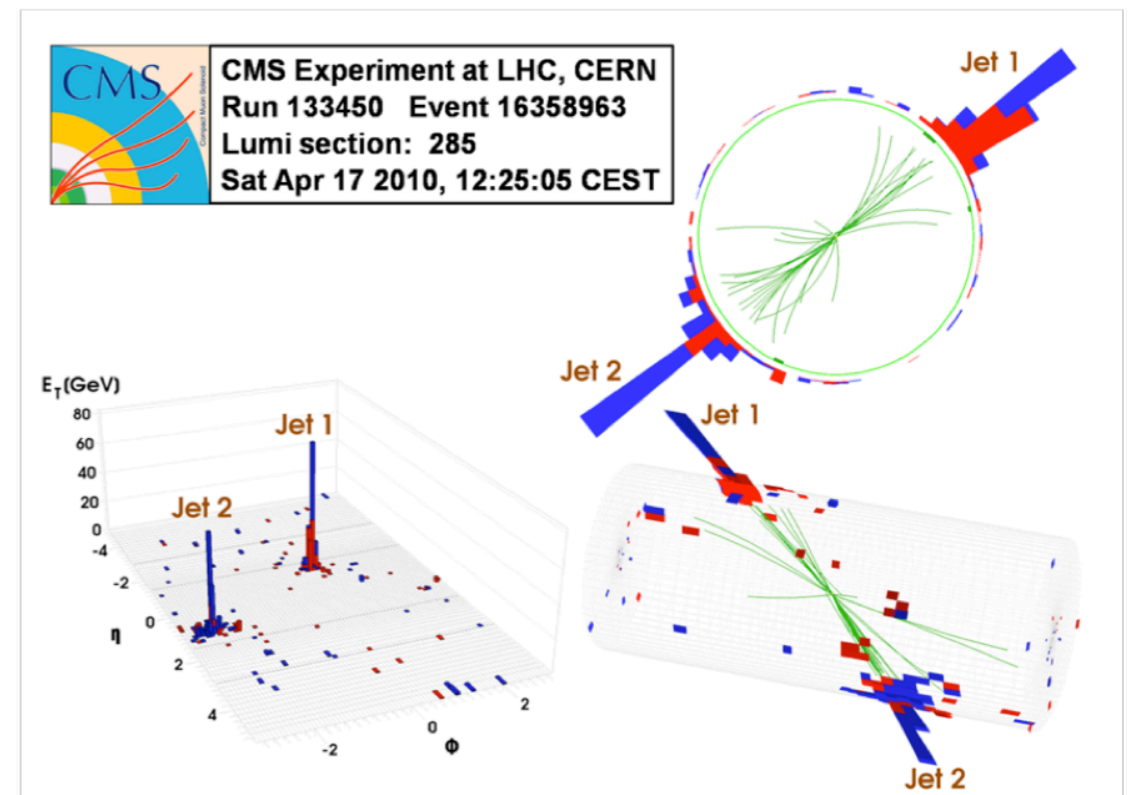
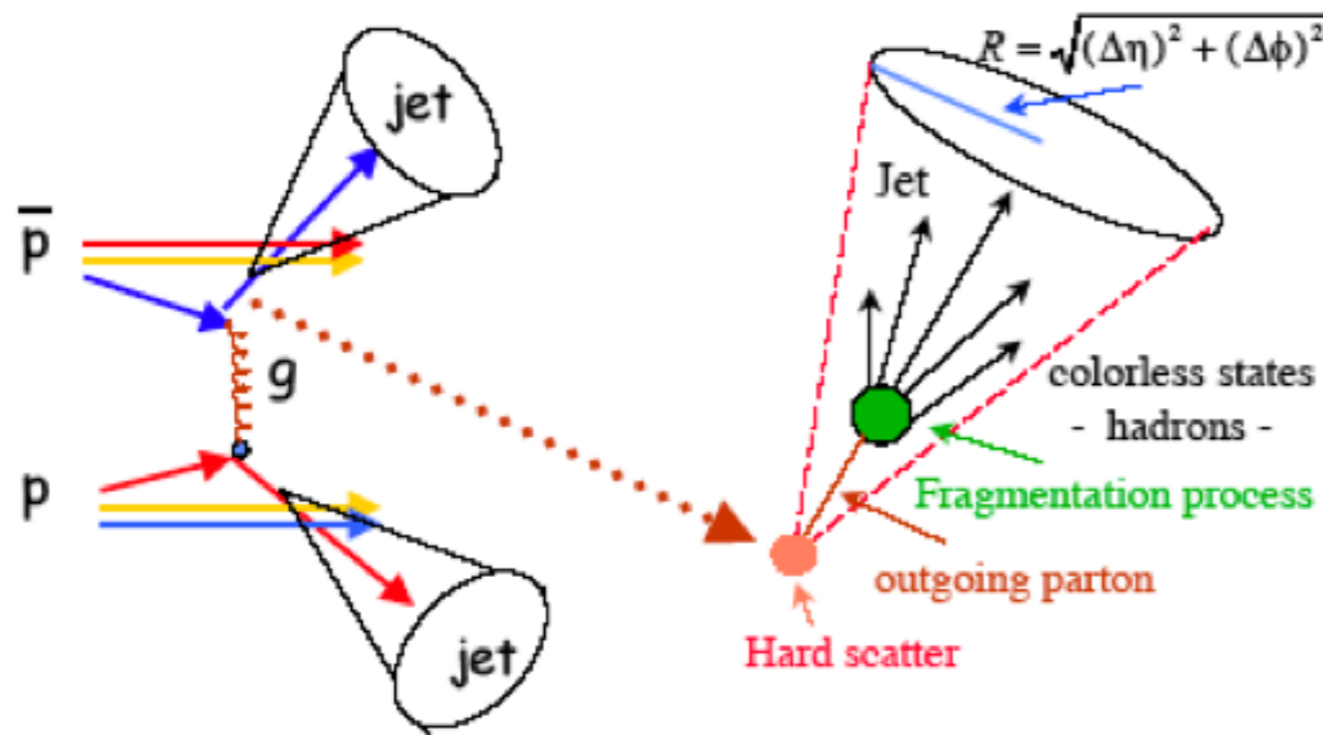
<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSUS110105fb>



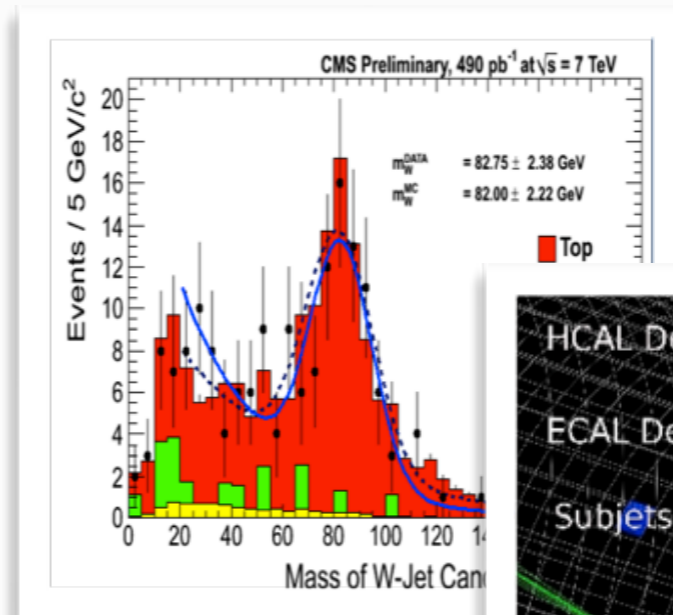
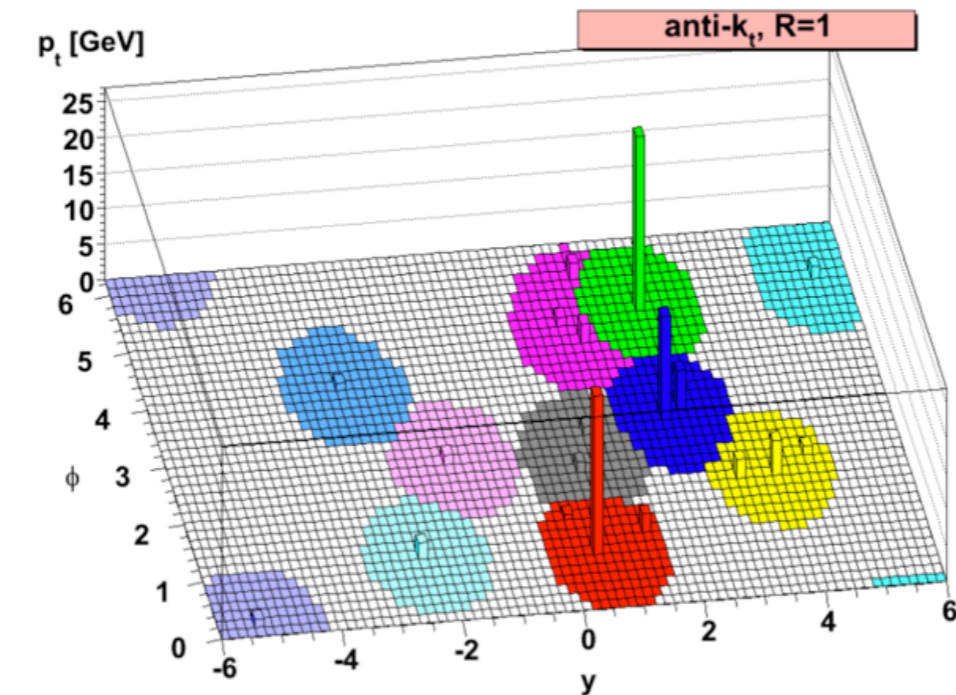
of course, the concrete numbers strongly  
depend on the selection cuts



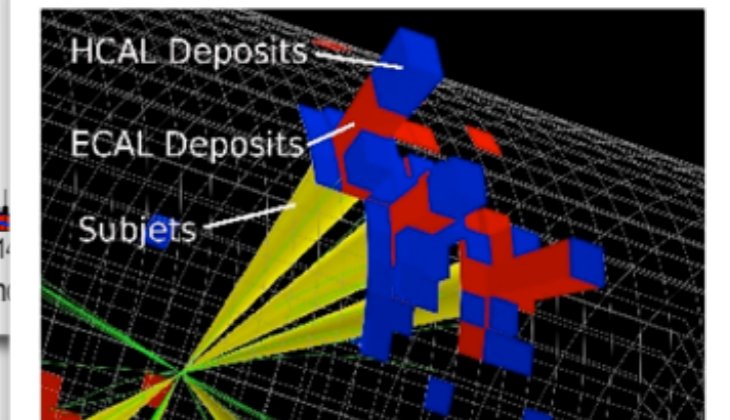
# Jet Reconstruction and Pile-Up subtraction



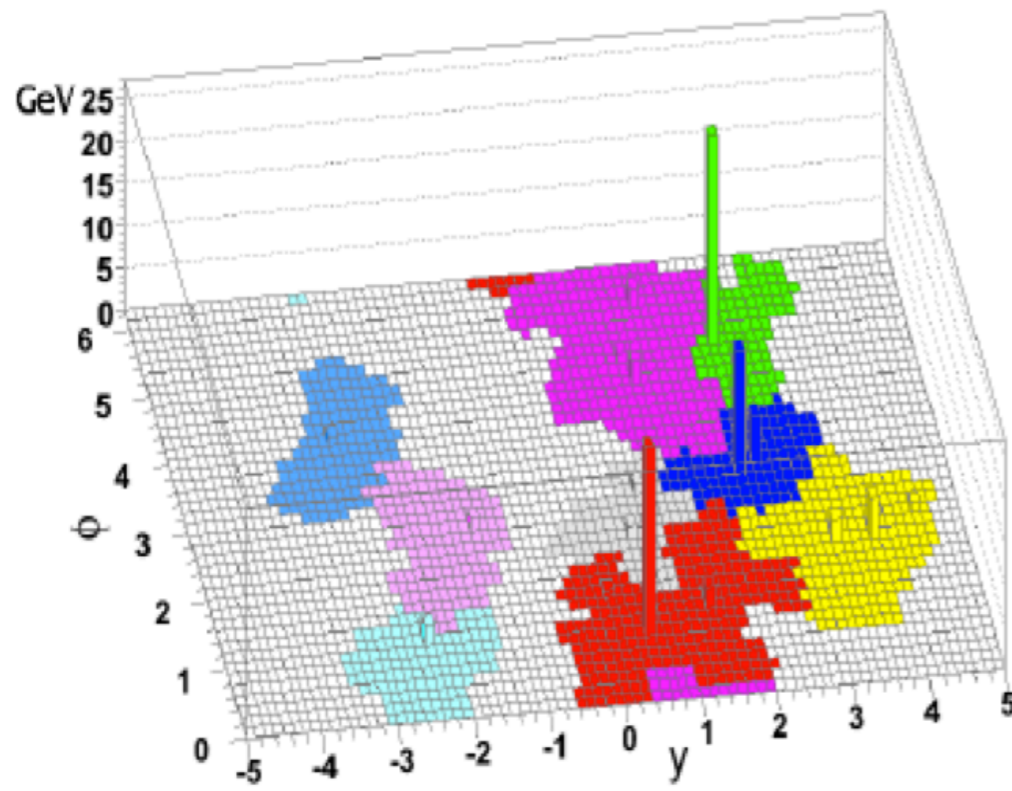
# Essential tools!



boosted tops,  
top tagging,  
subjects



G. Salam, "Towards Jetography"



A lot of the experimental success is due to fantastic tools, developed and proposed prior to the LHC startup, such as

- new jet algorithms (anti-KT)
- PU subtraction algorithms
- jet substructure scrutiny

They had a great impact; all proposed, mostly by theorists, only few years ago, but adopted quickly and became standards!

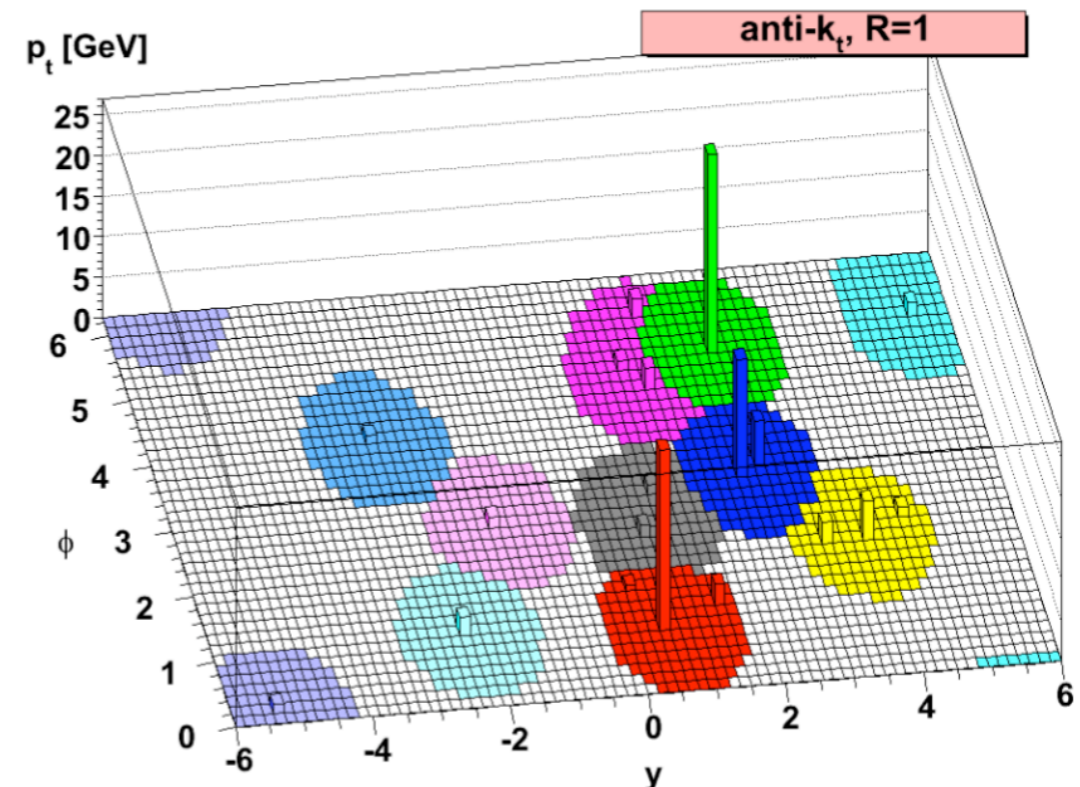
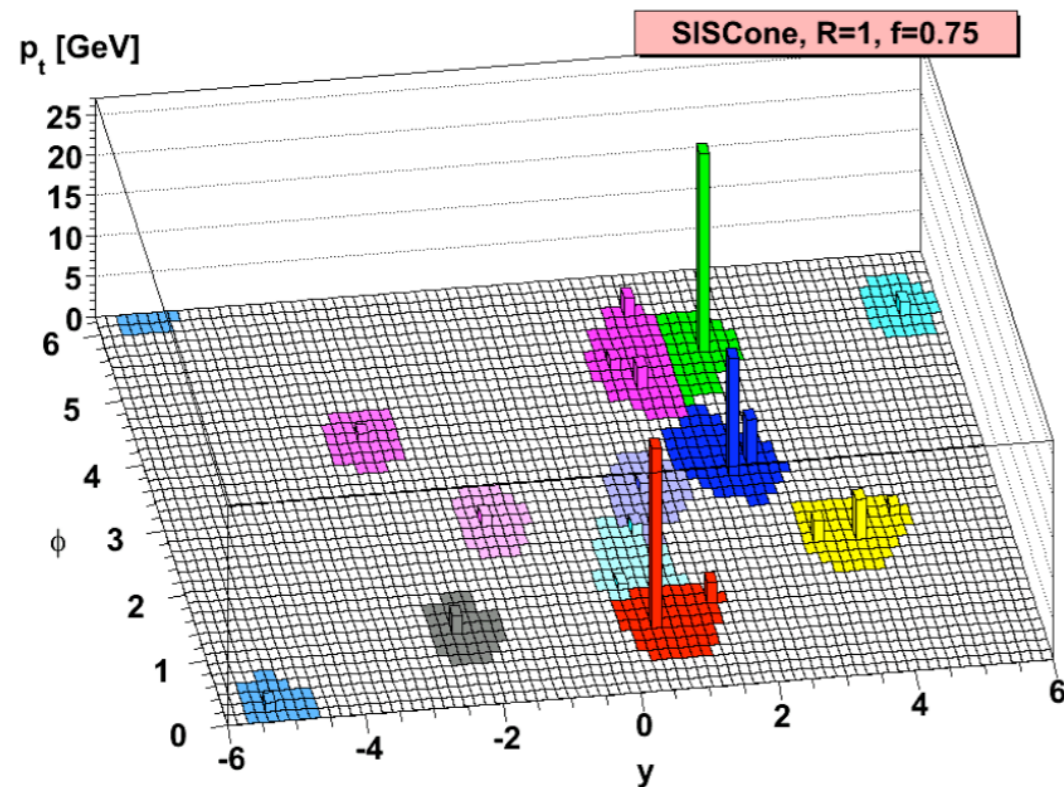
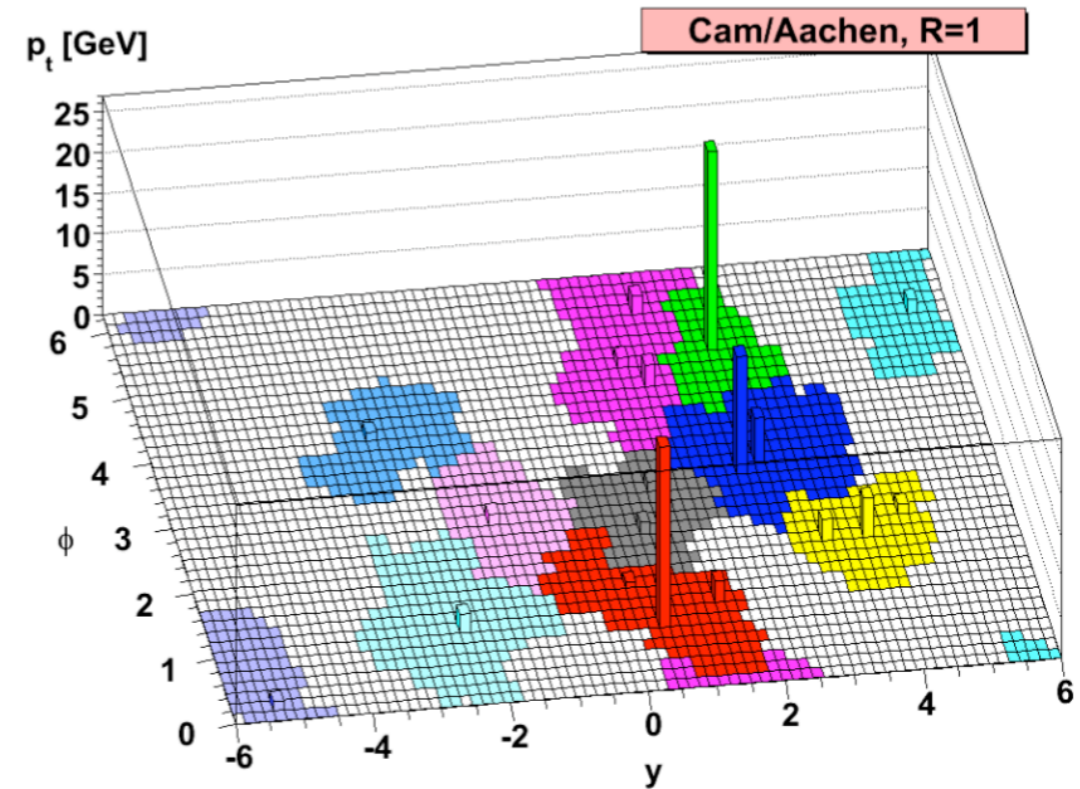
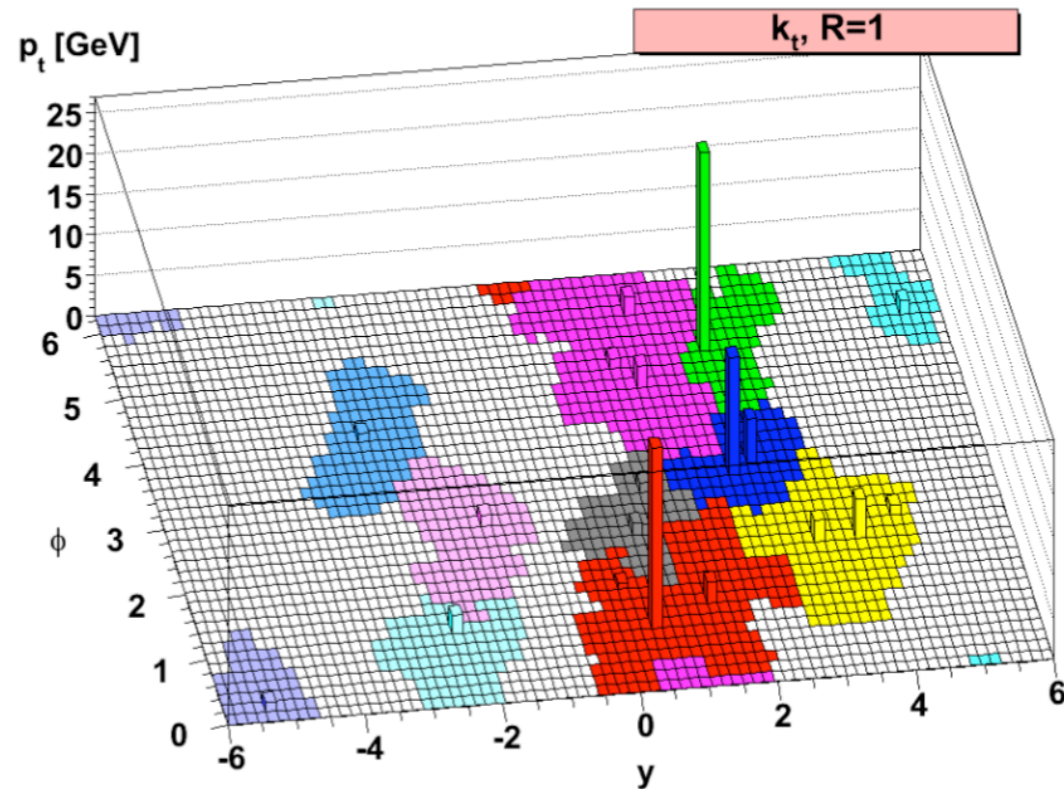
Thus, for the theorists among you: do not only think about super-difficult, non-planar three-loop master integrals, but: **also think about clever observables, clever tools, in order to get out the most from our data!**

- Once upon a time, there were
  - cone** algorithms (iterative cone, midpoint cone, ...)
  - known not to be infrared and collinear (IRC) safe (i.e., if implemented in a pert. QCD calculation, they would give non-sense and/or infinities at a certain order of pert. theory)
  - there was also the **IRC-safe kT-algorithm** (sequential recombination), but considered to be too slow for application at hadron colliders
- Then came Cacciari, Salam, Soyez
  - first, they developed an IRC-safe cone algorithm (**SisCone**)
  - then they found a clever way to tremendously speed up the kT-algorithm
  - then they generalized the clustering metric, kT became a special case, and another special case appeared and was adopted by the experiments:  
**the anti-kT algorithm**
  - Then: the fact that the kT (or similarly the anti-kT) algorithm is so fast, allowed them to introduce the jet-area method, which then led to a widely-used pile-up (PU) subtraction method! see later.....

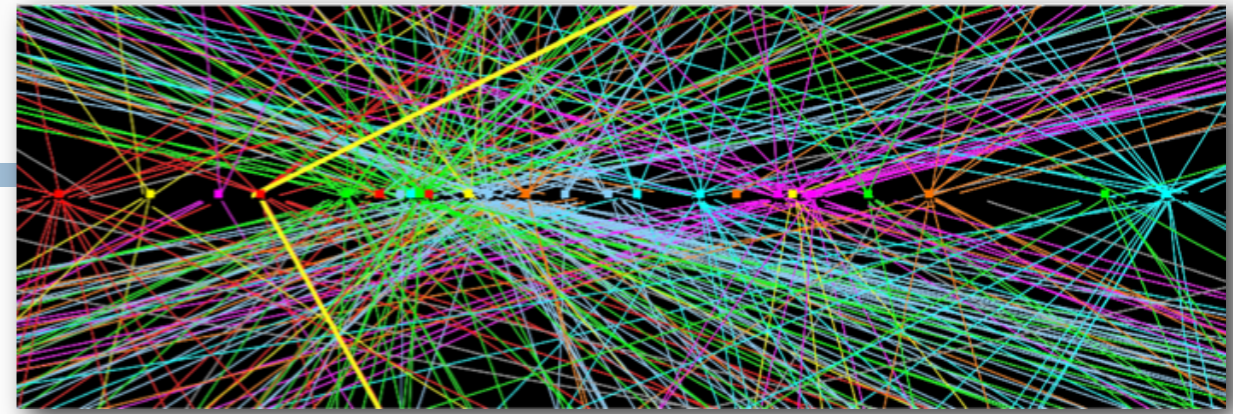
In the past few years, a suite of IRC safe algorithm  
has become widely used

$k_t$	<p>SR</p> $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ <p>hierarchical in rel <math>p_t</math></p>	<p>Catani et al '91 Ellis, Soper '93</p>
Cambridge/ Aachen	<p>SR</p> $d_{ij} = \Delta R_{ij}^2 / R^2$ <p>hierarchical in angle</p>	<p>Dokshitzer et al '97 Wengler, Wobish '98</p>
anti- $k_t$	<p>SR</p> $d_{ij} = \min(k_{ti}^{-2}, k_{tj}^{-2}) \Delta R_{ij}^2 / R^2$ <p>gives perfectly conical hard jets</p>	<p>MC, Salam, Soyez '08 (Delsart, Loch)</p>
SISCone	<p>Seedless iterative cone with split-merge gives 'economical' jets</p>	<p>Salam, Soyez '07</p>

# Jet areas



# Pile-up subtraction



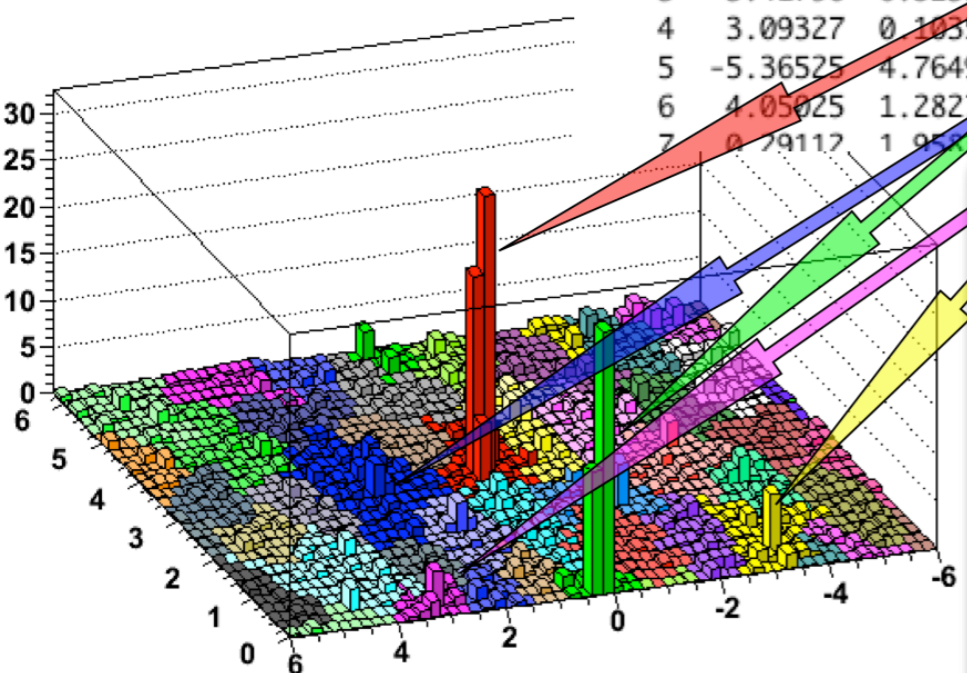
- **A bit of history....**
- Once upon a time, we thought
  - well, if we use a **cone** algorithm, then we get jets with nice cones, where we know the area. Then we measure the average PU energy falling into such an area, subtract it, and we are done
- Then, Cacciari et al. “took us our cone algorithms away”. They gave us a fast kT algorithm, but which produces terribly shaped jet areas! What to do now?
- But: since the kT algorithm is so fast, they had the idea:
  - throw into an event a very large number of almost-zero-energy particles (“**ghosts**”)
  - run the kT algorithm, find the jets (most of them soft PU jets), and take the **median** of the jet pT over the “**active (or catchment) area**”
  - this is basically the ratio of all ghosts over those ghosts clustered into a particular jet
  - taking the median removes bias from possible appearance of real high-pt jets
- Now there was a way to measure, event by event (!), the average PU energy (or rather pT- ) density, which then can be subtracted from the jets, as well as from areas/cones around other objects, such as isolated leptons or photons!!

# Pile-up subtraction

A concrete example:  
a 50 GeV di-jet event at the  
LHC with pile-up  
(10 min-bias events added)

iev 0 (irepeat 24): number of particles = 1428  
strategy used = NlnN  
number of particles = 9051  
Total area: 76.0265  
Expected area: 76.0265

ijet	eta	phi	Pt	area	+-	err
0	0.15050	3.24498	69.970	2.625	+-	0.020
1	0.18579	0.13150	59.133	1.896	+-	0.020
2	2.33840	3.23960	31.976	4.749	+-	0.028
3	-3.41796	0.52394	26.595	3.084	+-	0.021
4	3.09327	0.10350	20.072	2.688	+-	0.023
5	-5.36525	4.76491	19.584	2.780	+-	0.012
6	4.05025	1.28279	15.361	3.592	+-	0.028
7	0.29112	1.95745	14.566	2.114	+-	0.018



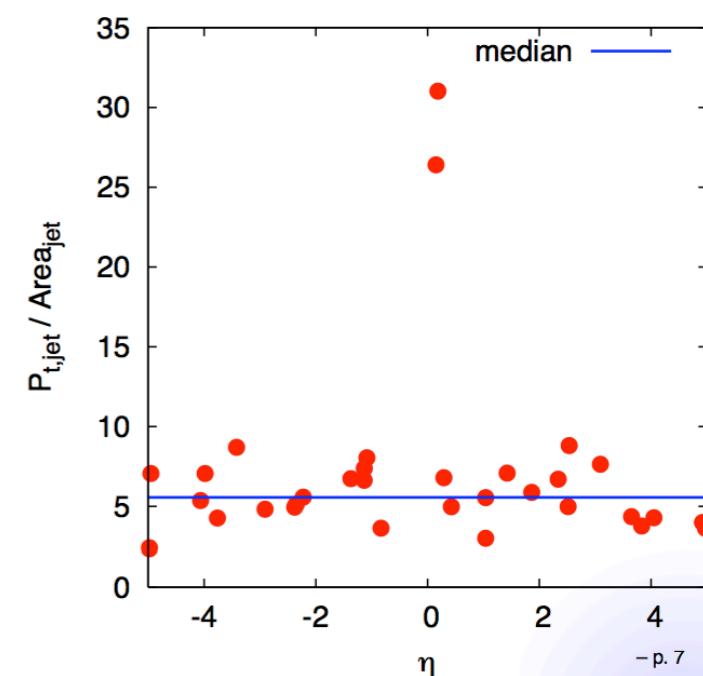
slides from M. Cacciari and G. Soyez

[M.Cacciari, G.P. Salam, GS, 2007-08]

$$p_{t,\text{jet}}^{(\text{sub})} = p_{t,\text{jet}} - \rho_{\text{bkg}} A_{\text{jet}}$$

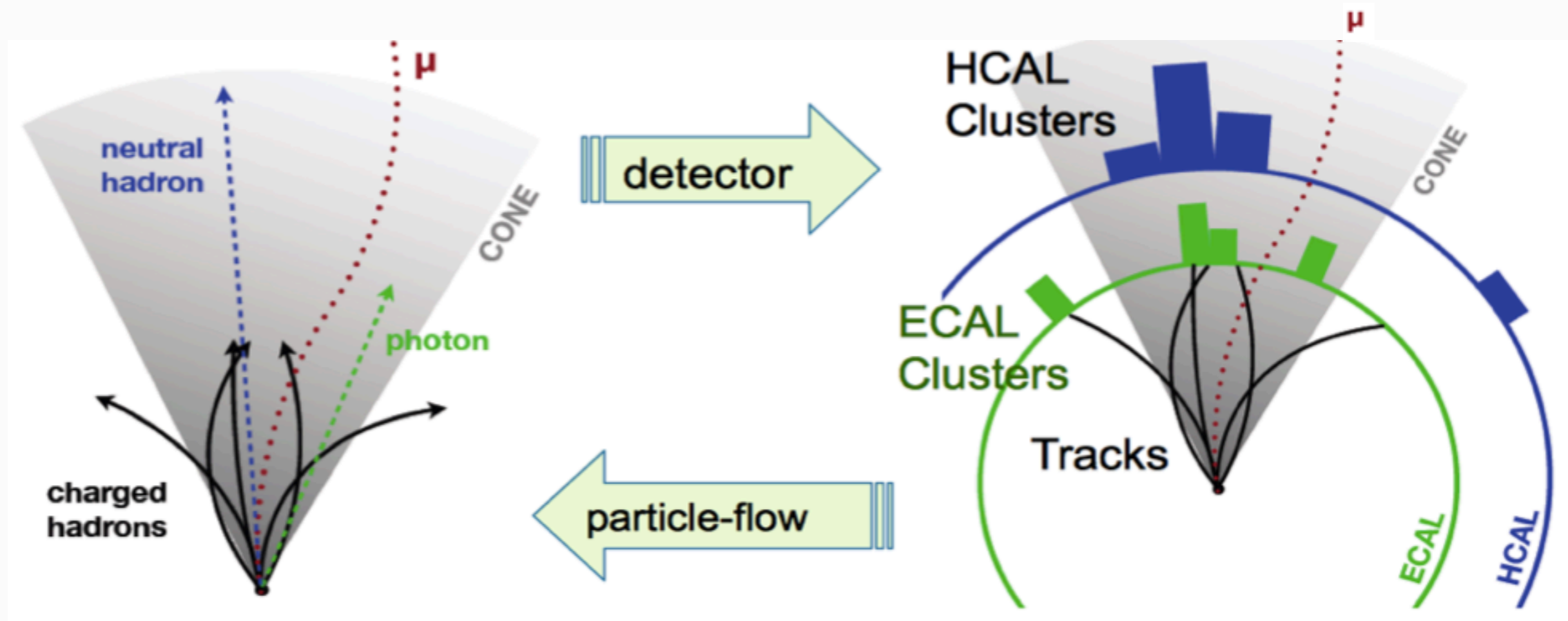
- jet area: available with jet clustering
- $\rho_{\text{bkg}}$ , the background  $p_t$  density per unit area
  - break the event in patches of similar size  
e.g. cluster with  $k_t$
  - Estimate  $\rho_{\text{bkg}}$  using

$$\rho_{\text{bkg}} = \text{median}_{j \in \text{patches}} \left\{ \frac{p_{t,j}}{A_j} \right\}$$



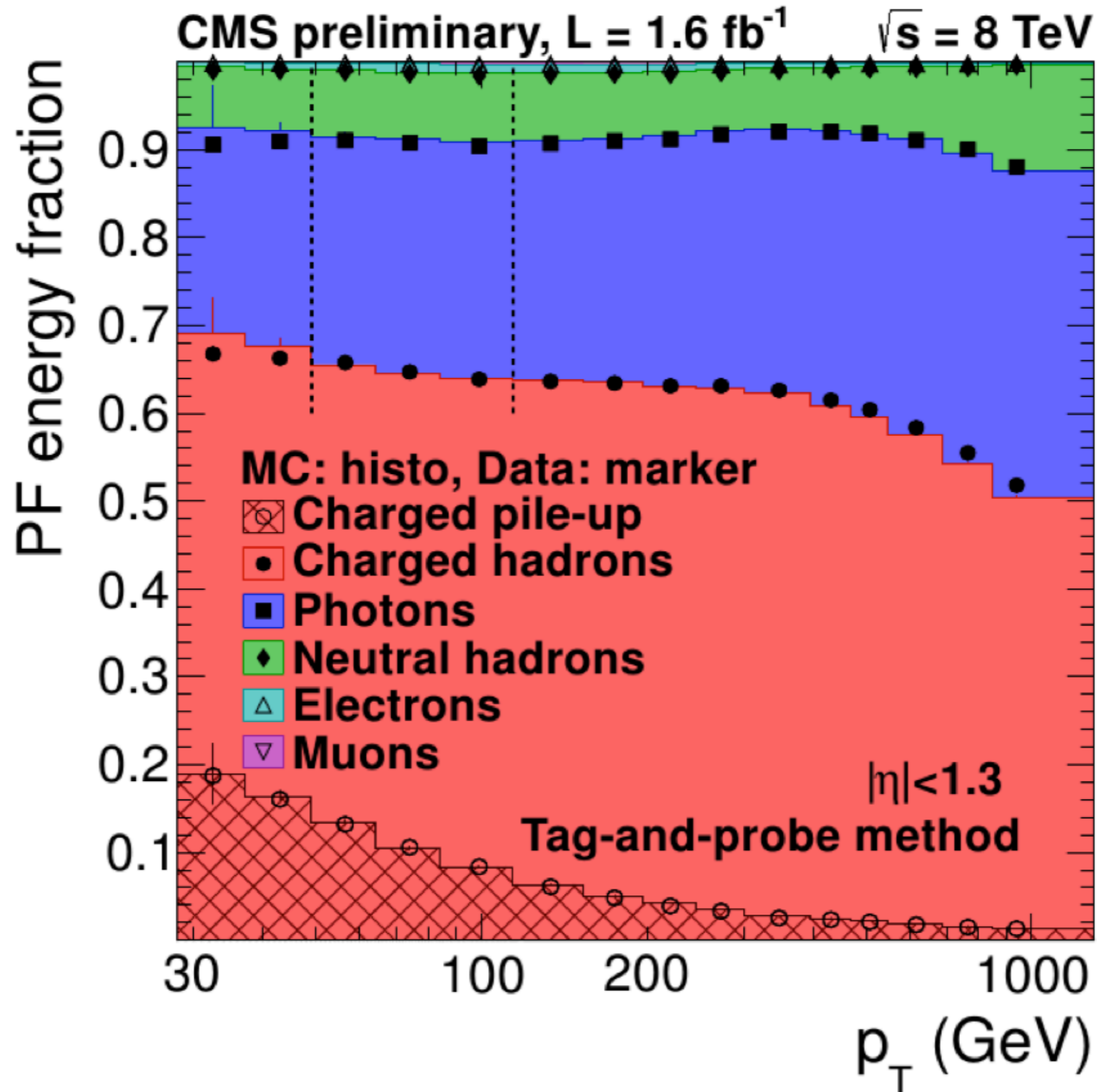
# Particle Flow and Consequences

# Use of global event description

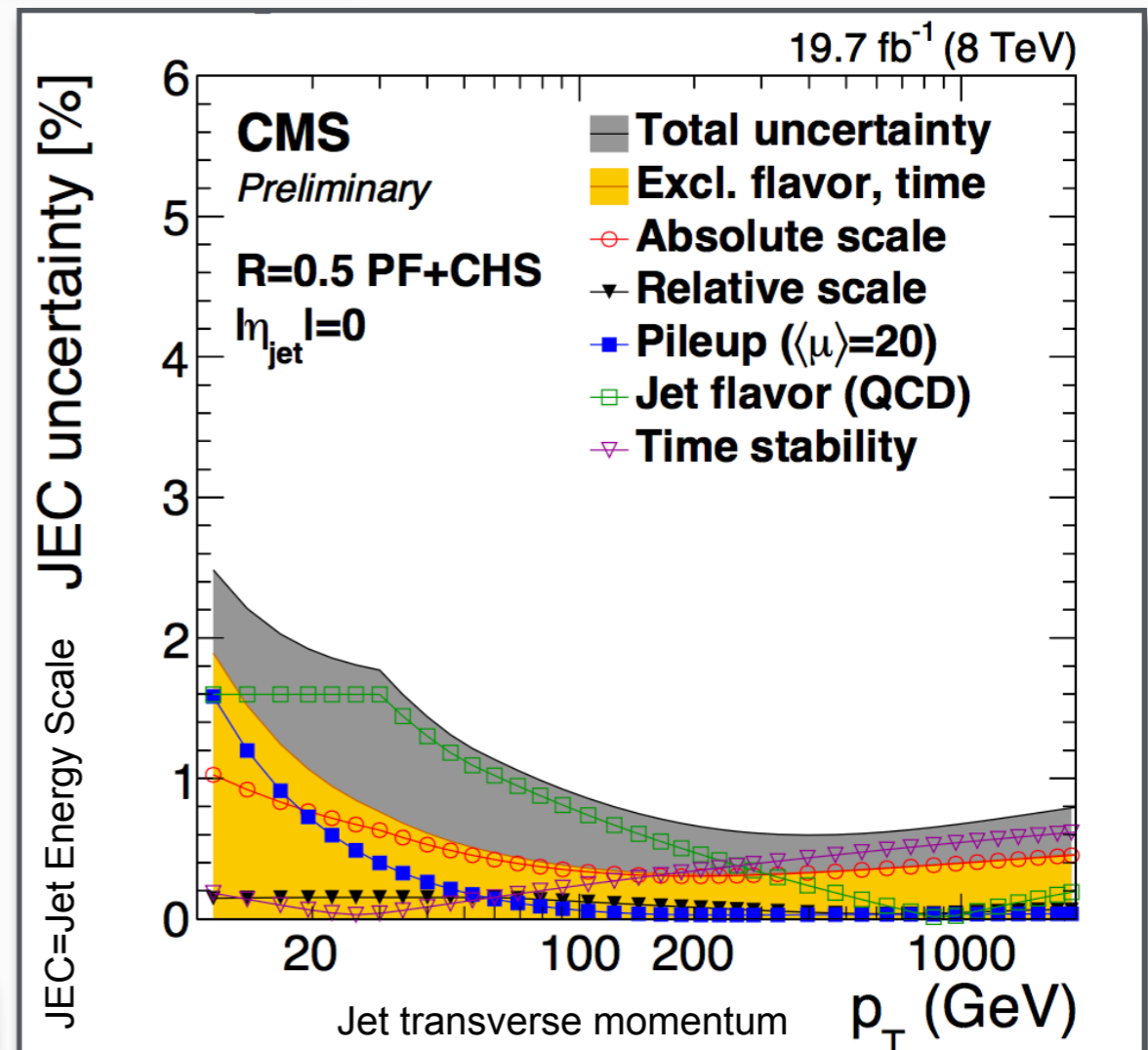
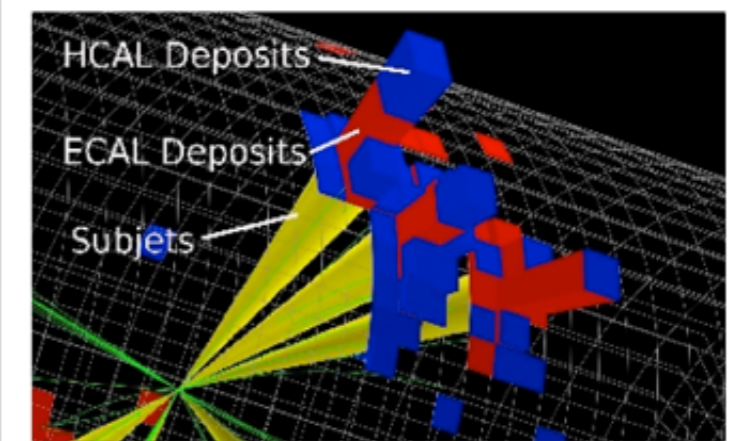
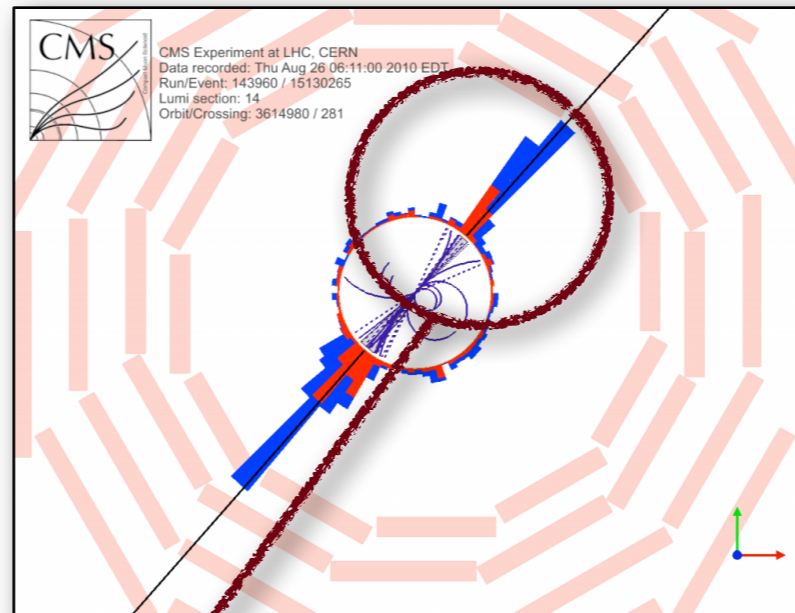
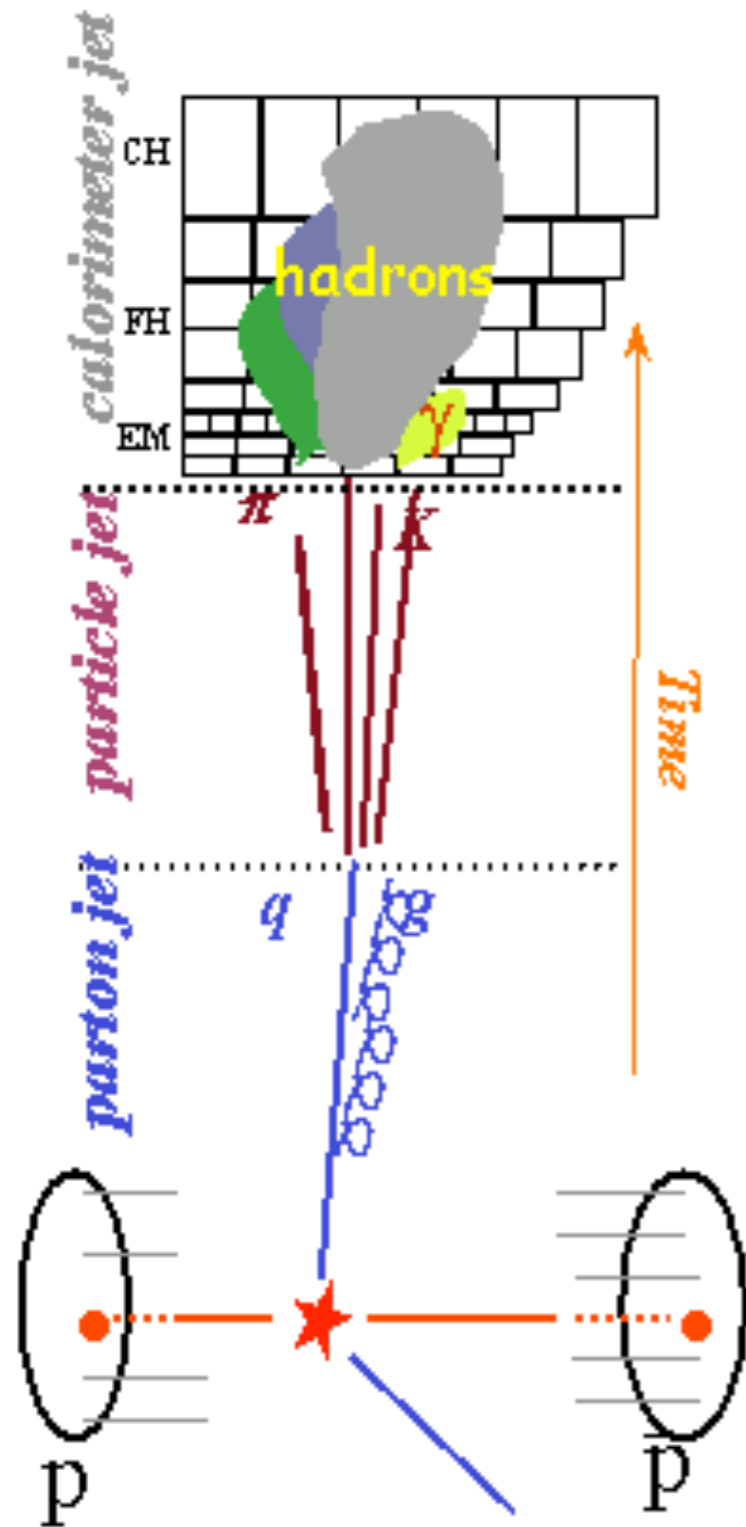


- Charged particles well separated in large tracker volume & 3.8T B field
- Excellent tracking, able to go down to very low momenta ( $\sim 100$  MeV)
- Granular electromagnetic calorimeter with excellent energy resolution
- In multi-jet events, only 10% of the energy goes to neutral (stable) hadrons ( $\sim 60\%$  charged,  $\sim 30\%$  neutral electromagnetic)
- Therefore: **Use a global event description :**
  - Optimal combination of information from all subdetectors
  - Returns a list of reconstructed particles (e,  $\mu$ , photons, charged and neutral hadrons)
  - Used as building blocks for jets, taus, missing transverse energy, isolation and PU particle ID

# The Pflow jet composition



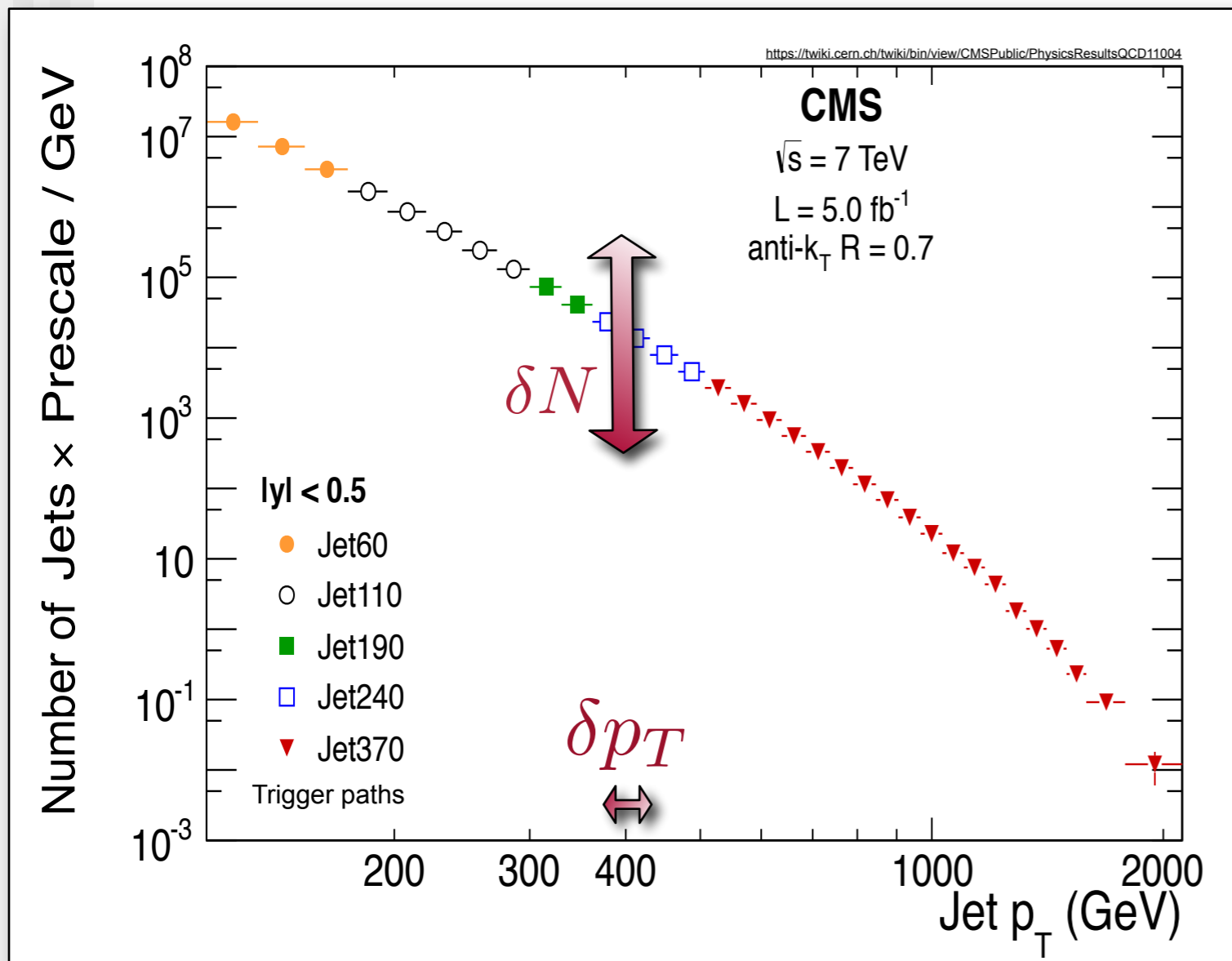
# Impact on Jet Calibration



JES uncertainty  $< 2\%$  for most of the  $p_T$  range, JER about 10%

# Jet energy calibration

- Question : how well do we know the **calibration** of the variable on the x-axis, eg. jet energy?
- A general problem for a very steeply falling spectrum!
- It makes a **big** difference if the jet energy scale uncertainty is 1%, 2% or 5%



$$\frac{d\sigma}{dp_T} \approx \text{const} \cdot p_T^{-n}$$

$n$  large, eg.  $n \approx 6$



relative uncertainties

$$\frac{\delta N}{N} \approx 6 \cdot \frac{\delta p_T}{p_T}$$

**so beware:**

eg. an uncertainty of **5%** on absolute energy scale (calibration)

→ an uncertainty of **30%** (!) on the measured cross section

# The Luminosity

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\varepsilon L}$$

# Two possible approaches

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon L}$$

invert the  
problem

$$L = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon \sigma_{\text{theo}}}$$

**Needs a very precisely  
calculable process,  
eg. W and Z production,  
as well as low exp. uncertainties**

measure the  
luminosity from first  
principles

$$L = \frac{N_1 N_2 \nu_{\text{orb}} N_b}{2\pi \sigma_{\text{eff}}(x) \sigma_{\text{eff}}(y)} \equiv \frac{\dot{N}}{\sigma}$$

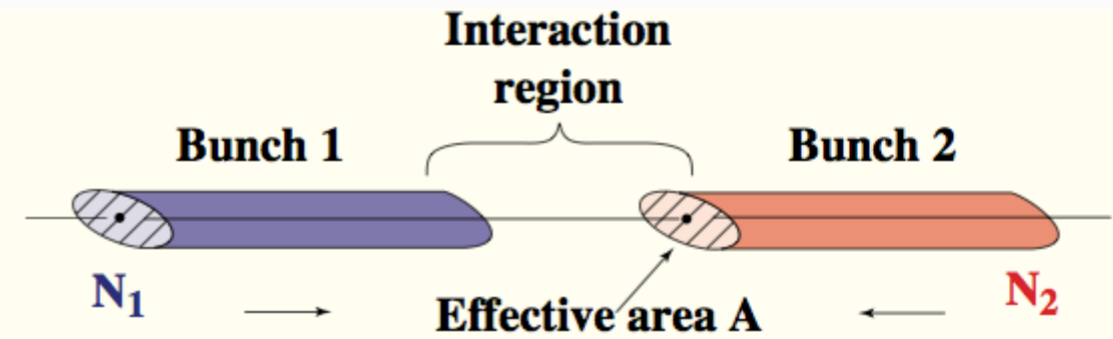
**Have to measure:**

- beam currents
- effective beam size --> Van der Meer scan !
- then, after absolute calibration:

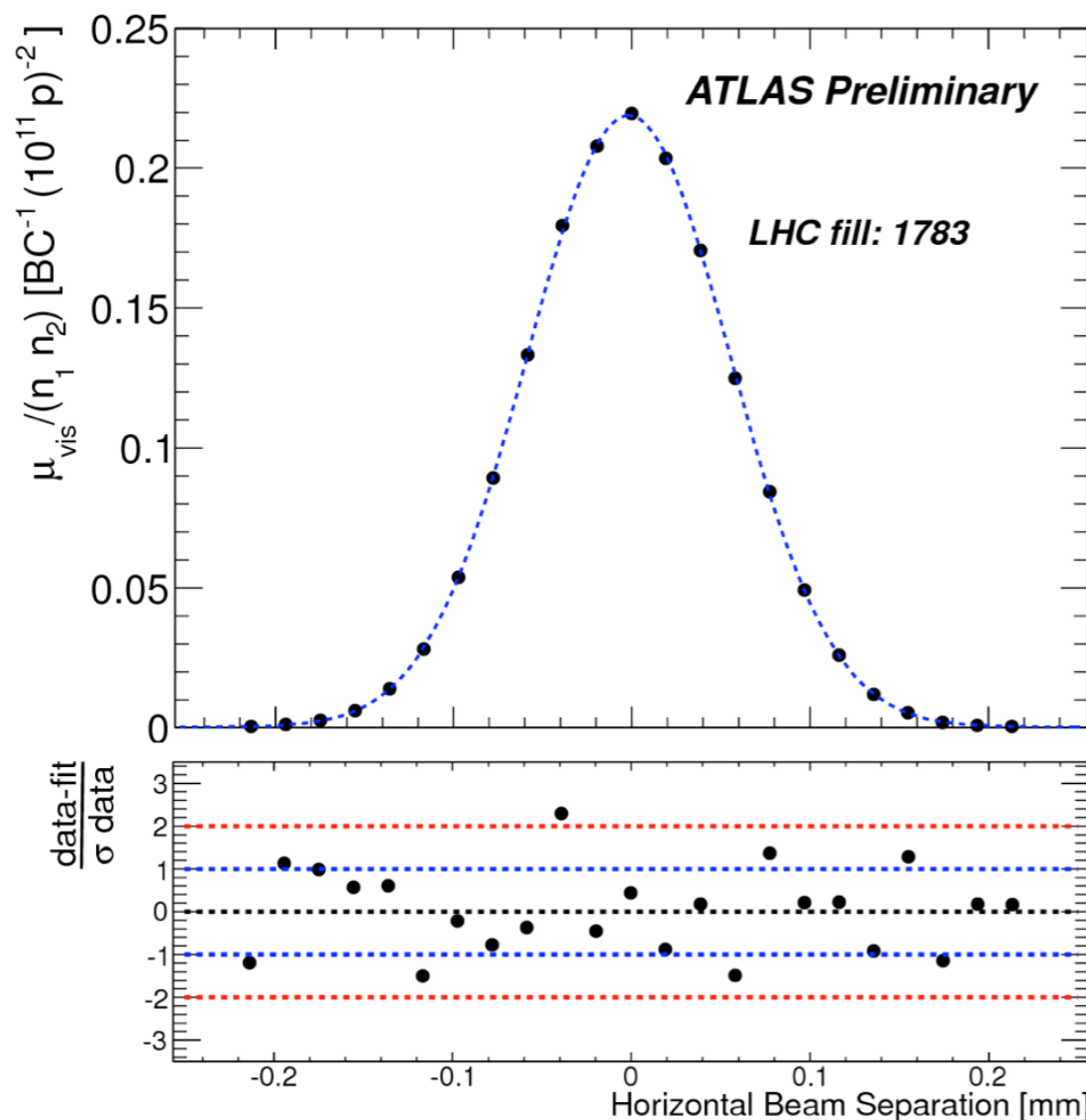
take stable process to measure evolution in  
time, eg. number of counts in forward calorimeters

# Van der Meer scans

- Move the beams relative to each other and monitor the rate of some basic process, eg. MinBias triggers



for Gaussian bunches with rms sizes  $\sigma_x \sigma_y$   $A = 4 \pi \sigma_x \sigma_y$



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Uncertainty Source	$\delta\mathcal{L}/\mathcal{L}$
<i>vdM</i> Scan Calibration	3.4%
Afterglow Correction	0.2%
Long-term consistency	1.0%
$\mu$ Dependence	1.0%
Total	3.7%