

# Rope Hadronisation

54th Cracow School of Theoretical Physics, Zakopane, 2014-06-13

Christian Bierlich  
bierlich@thep.lu.se

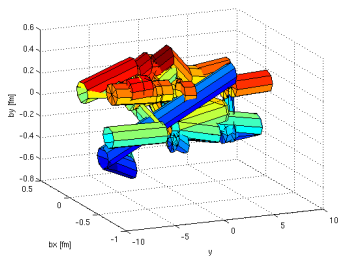
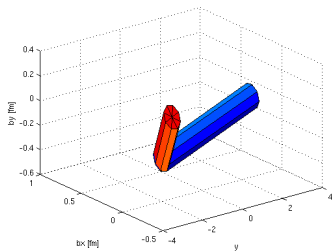
With Gösta Gustafson, Leif Lönnblad and Andrey Tarasov  
Lund University, Dept. Astronomy and Theoretical Physics

Support by the MCnet, Marie Curie Initial Training Network is gratefully acknowledged



# Introduction

- The *string model* is the primary hadronization workhorse for the Lund family of Monte Carlo EG.
- The string model is developed for the very clean LEP environment...
- ...while  $pp$  min. bias or Heavy Ion is more messy
- Corrections based on *coherence effects* or *rope hadronisation*.
- Part of the DIPSY/Ariadne 5 Monte Carlo EG.

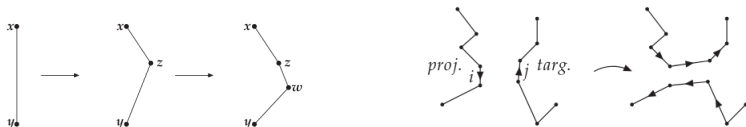


# The DIPSY event generator [Flensburg et al. arXiv:1103.4321 [hep-ph], 2011]

- A Monte Carlo implementation of the Mueller Dipole Cascade (with corrections) [see talks this morning or G. Gustafson, proceedings 2011].
- Builds up virtual Fock states of proton, colliding dipoles interact via gluon exchange.
- A dipole  $(\vec{x}, \vec{y})$  can emit a gluon at position  $\vec{z}$  with probability  $(P)$  per unit rapidity  $(Y)$ ; dipoles  $i$  and  $j$  interact with probability  $2f_{ij}$ :

$$\frac{dP}{dY} = \frac{\bar{\alpha}}{2\pi} d^2\vec{z} \frac{(\vec{x} - \vec{y})^2}{(\vec{x} - \vec{z})^2 (\vec{z} - \vec{y})^2}$$

$$f_{ij} = \frac{\alpha_s^2}{8} \left[ \log \left( \frac{(\vec{x}_i - \vec{y}_j)^2 (\vec{y}_i - \vec{x}_j)^2}{(\vec{x}_i - \vec{x}_j)^2 (\vec{y}_i - \vec{y}_j)^2} \right) \right]^2$$



# Monte Carlo event generation

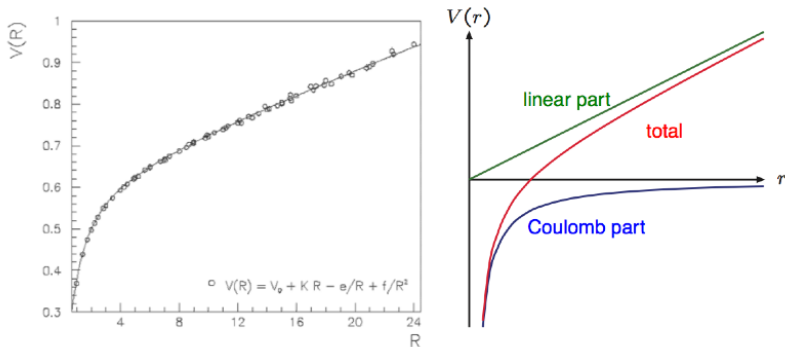
- Many corrections beyond Muellers model – see reference on previous page.
- Rough sketch of MC generation of an event:
  - 1 Generate projective and target cascades;
  - 2 Determine which dipoles interact;
  - 3 Reabsorption of non-interacting chains (dipole loops);
  - 4 Final state radiation (ARIADNE 5);
  - 5 Hadronisation;
- The rest of this talk will concentrate on the last part!

# Rest of this talk

- 1 Ordinary string hadronisation
- 2 Rope model
- 3 Results for  $pp$
- 4 Prospects
- 5 Outlook and conclusion

# String Hadronisation [Sjöstrand et al. hep-ph/0603175, 2006]

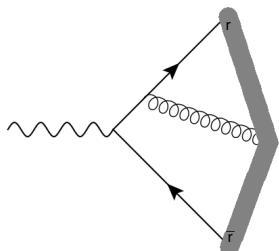
- Linear confinement potential  $V(r) \approx \kappa r$ ,  $\kappa \approx 1$  GeV/fm.
- Valid for large distances – for small distances perturbation theory should be valid.



- Realized in a 1+1 dimensional string with tension  $\kappa$ .

## String Hadronisation II [Andersson et al. Z. Phys. C20 317, 1983]

- Repeated *breaking* with  $\mathcal{P} \propto \exp\left(-\frac{\pi m_{\perp}^2}{\kappa}\right)$  gives hadrons.
- Left-right symmetry in the breaking gives  $f(z) \propto z^{-1}(1-z)^a \exp\left(\frac{-bm_{\perp}}{z}\right)$ .



- $a$  and  $b$  related to total multiplicity.
- Flavours determined by relative probabilities:

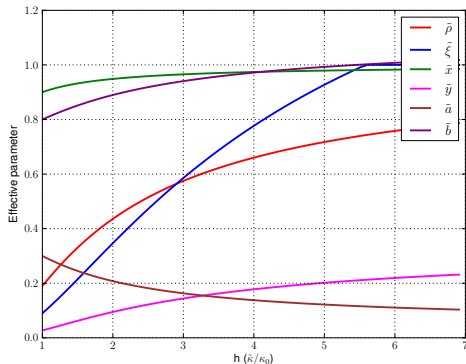
$$\rho = \frac{\mathcal{P}_{\text{strange}}}{\mathcal{P}_{\text{u or d}}}, \xi = \frac{\mathcal{P}_{\text{diquark}}}{\mathcal{P}_{\text{quark}}}$$

$$x = \frac{\mathcal{P}_{\text{strange diquark}}}{\mathcal{P}_{\text{diquark}}}, y = \frac{\mathcal{P}_{\text{spin 1 diquark}}}{\mathcal{P}_{\text{spin 0 diquark}}}$$

- Notice that probabilities are related to  $\kappa$  via tunneling equation.

# Change of string tension

- All parameters related through string tension.
- Used to produce a set of *effective parameters*.



- Let  $\kappa \mapsto \tilde{\kappa} = h\kappa$ .
- All effective parameters calculable from the two governing equations.
- Parameters  $\rho$  and  $\xi$  are very sensitive to change in  $\kappa$ .
- Can this be connected to overlapping strings?

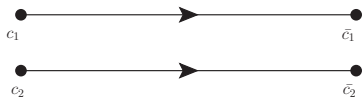


# The Rope Model at a glance

- The *Rope Model*:
  - ① Let charges act coherently to form a rope.
  - ② The rope is an SU(3) multiplet, with  $\kappa \propto C_2$ .
  - ③ All ropes described by two quantum numbers  $\{p, q\}$ .
  - ④ Rope breaks in steps, one strand at a time.
- Must still map colour configuration to  $\{p, q\}$  for individual strands.
- Hadronisation of strands by Pythia 8, using effective parameters.
- *Input*: Colour connected output from DIPSY/Ariadne 5 shower.
- *Output*: Individual strands with an associated set of effective parameters.

## Rope model – example 1

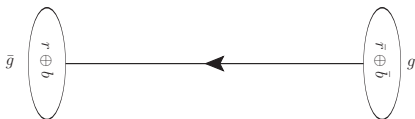
- The simplest example: Two  $q\bar{q}$  pairs act coherently, colour flow in same direction:



Case (a),  $c_1 = c_2$ :

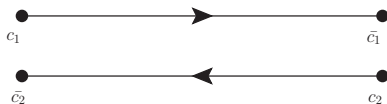


Case (b),  $c_1 \neq c_2$ :



## Rope model – example 2

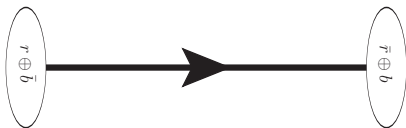
- Next-to-simplest: Two  $q\bar{q}$  pairs act coherently, having oppositely directed colour flow:



Case (a),  $c_1 = c_2$  :



Case (b),  $c_1 \neq c_2$  :



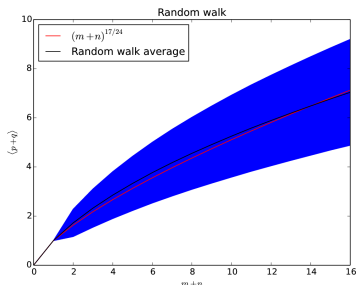
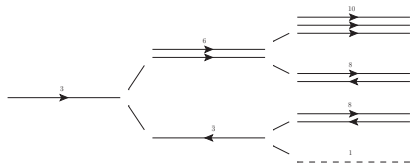
## Generalization – The Ropewalk

- Generalizable to arbitrarily many triplets and anti triplets  $\{m, n\}$ .
- The procedure is iterative, adding one (anti)-triplet at a time.
- This is similar to a random walk in colour space Biro et al., Nucl. Phys. B 245, 449-468, 1984.
- Below: A forming machine at the end of a *ropewalk* (reeperbahn).



# The Ropewalk II

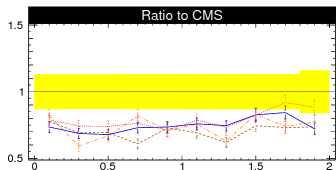
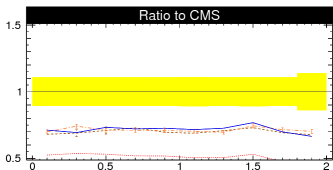
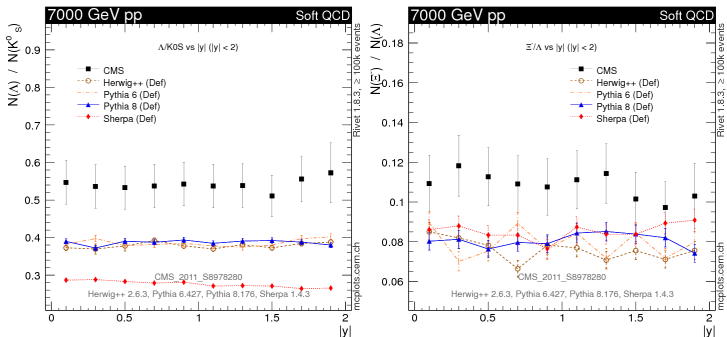
- Random walk carried out to form multiplets, using recursion relations.



- Rope tension equivalent to 2nd Casimir operator [Ambjørn, J et al. Nucl. Phys. B 240, 533, 1984] and [Bali, G. S, hep-lat/0006022, 2000].
- Averaged over  $N = p + q$  partial breakups:  $\tilde{\kappa}/\kappa = \frac{1}{4}(N + 3 - \frac{pq}{N})$

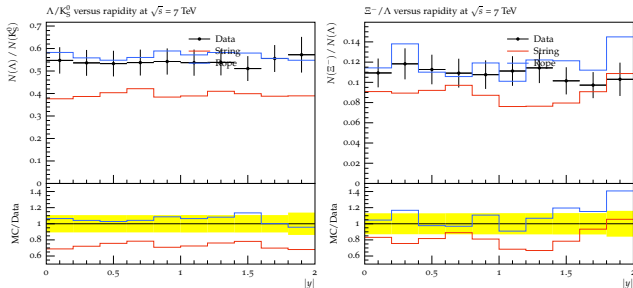
## Comparison to data (CMS)

- From MCplots, comparison to CMS data [CMS, arXiv:1102.4282 [hep-ex], 2012]:



# Comparison to data (CMS)

- Vanilla DIPSY/ARIADNE 5 (string hadronisation) suffers from similar problem.
- Including rope effects improves the picture.



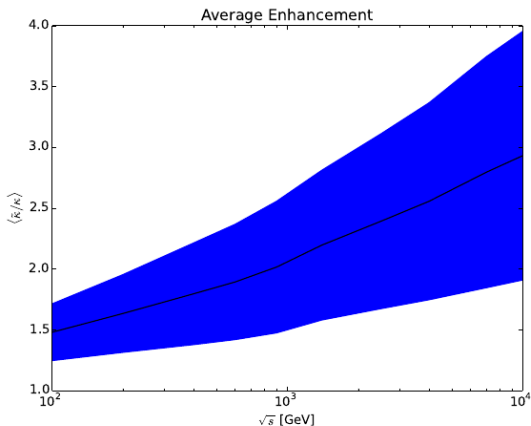
# Two prospects

- Two future prospects:
- Higher energy (trivial)
- Heavy ions



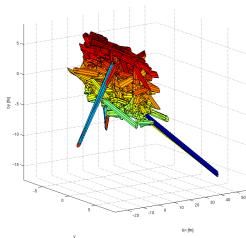
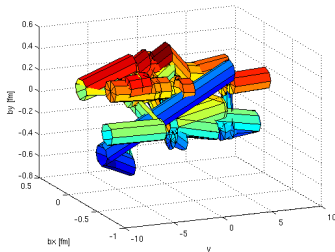
# Increased energy

- Increased energy results in more strings.
- With more overlap comes a larger effect!



# Heavy Ion (very preliminary!)

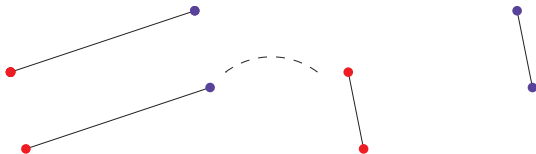
- Heavy ion events (AuAu at 130 GeV, RHIC) are way more stringy!



- It is absolutely necessary to deal with coherence effects.

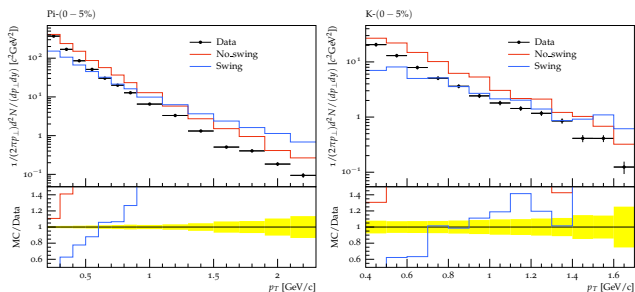
# The Colour Swing

- We deal with the simple case  $3 \otimes \bar{3} \rightarrow 1$  using a colour swing mechanism (also in  $pp$ ).
- Let nearby dipoles »swing« if compatible.
- Far reaching consequences for kinematics!



# Comparison to data (PHENIX) [PHENIX, 2002, arXiv:nucl-ex/0112006]

- Hard tail from the swing, rope effect not visible.
- Even more pronounced for higher  $p_{\perp}$  (eg. ALICE PbPb).



# Outlook and conclusions

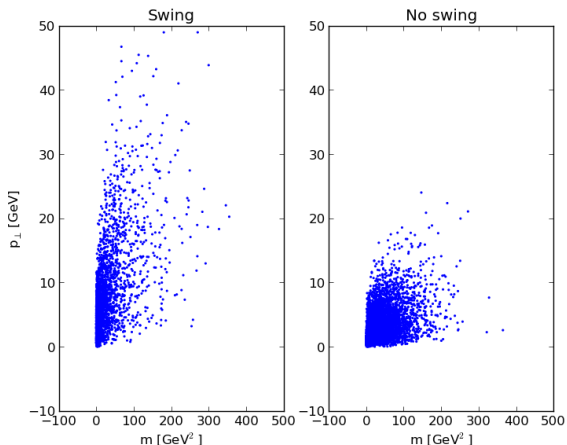
- The string model needs corrections in dense environments.
- Spatially overlapping strings enhance string tension.
- Implemented in context of DIPSY/Ariadne 5.
- Ropes formed by colour charges acting coherently gets strange/baryonic content right in  $pp$ .
- Effects should increase for higher LHC energies/SLHC.
- Promise for Heavy Ion events, colour reconnection by swing is a work in progress.

*The End*

*Bonus slides*

## Effect of swing

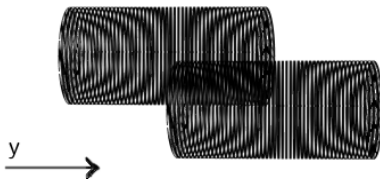
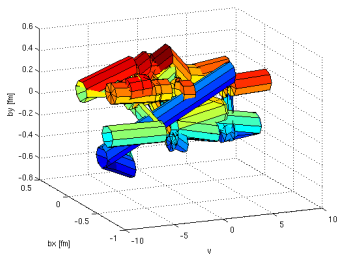
- Swing transforms low  $p_{\perp}$  particles to high  $p_{\perp}$  particles.
- Seems to be right for  $pp$ , but overshooting the effect in HI collisions.



# Calculation of $\{m, n\}$

- Finding the  $\{m, n\}$  strings producing a rope is not an entirely trivial matter.
- Must deal with partial overlaps.
- Crude approximation: Draw cylinders around each string, calculate volume overlap of cylinders:

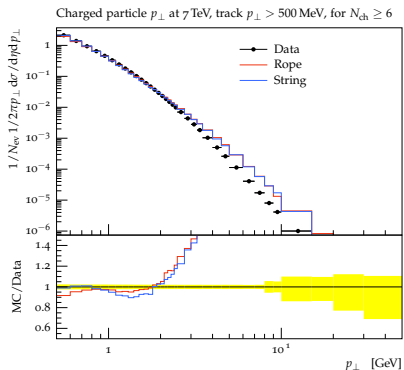
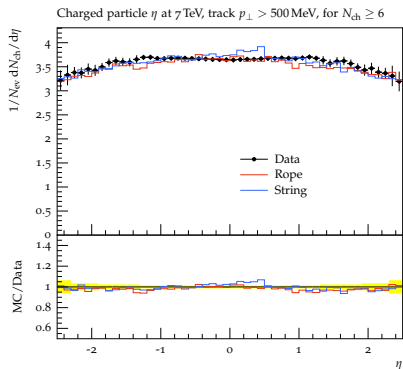
$$\langle N \rangle_{\text{total}} = \langle N \rangle_{\text{self}} + \sum_{i \neq \text{self}} \frac{V_{o,i}}{V_i} \langle N \rangle_i$$





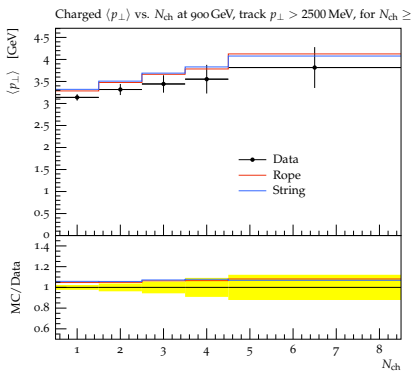
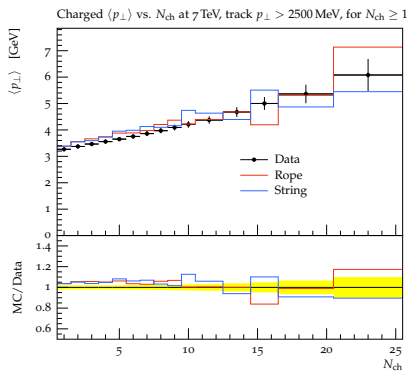
# Effect on total multiplicity [ATLAS, Phys. Lett. B688, 21-42, 2010]

- Must not destroy total multiplicity, which is so carefully tuned...



# Effect on $\langle p_{\perp} \rangle$ [ATLAS, Phys. Lett. B688, 21-42, 2010]

- The mean  $p_{\perp}$  is also an interesting observable.
- We see that it is not destroyed by neither swing nor ropes.



## Recursion relations etc.

- All multiplets are given by pairs of fundamental quantum numbers  $\{p, q\}$ .
- We can choose another set of fundamental quantum numbers, eg.:

$$N = \frac{1}{2}(p+1)(q+1)(p+q+2),$$

$$C_2(\{p, q\})/C_2(\{1, 0\}) = \frac{1}{4}(p^2 + q^2 + pq + 3(p+q))$$

- Related to Young tableaux as:  $q = \lambda_2$  and  $p+q = \lambda_1$  so eg:

$$\{2, 3\} = \begin{array}{|c|c|c|c|c|} \hline \square & \square & \square & \square & \square \\ \hline \square & \square & & & \\ \hline \end{array}$$

- Following the usual rules we have:

$$\{1, 0\} \otimes \{1, 0\} = \square \otimes \square = \square\square \oplus \begin{array}{|c|} \hline \square \\ \hline \end{array} = \{2, 0\} \oplus \{0, 1\},$$

$$\{1, 0\} \otimes \{0, 1\} = \square \otimes \begin{array}{|c|} \hline \square \\ \hline \end{array} = \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \oplus I = \{1, 1\} \oplus \{0, 0\}.$$

# Recursion relations cont'd

- Directly generalizable:

$$\{p, q\} \otimes \{1, 0\} = \underbrace{\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}}_q \dots \underbrace{\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}}_p \otimes \square =$$

$$\{p, q-1\} \oplus \{p-1, q+1\} \oplus \{p+1, q\}.$$

- and:

$$\{p, q\} \otimes \{0, 1\} = \underbrace{\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}}_q \dots \underbrace{\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}}_p \otimes \begin{array}{|c|} \hline \square \\ \hline \end{array} =$$

$$\{p-1, q\} \oplus \{p, q+1\} \oplus \{p+1, q-1\}.$$