

# Gravity and Non-Commutative Geometry

An alternative bridge in two dimensions

Sara Tavares

Zakopane, 30<sup>th</sup> June



# Outline

- 1 **Problem:** Geometric interpretation of spinors in gravity
- 2 **Strategy I:** Lattice gauge theory
- 3 **Strategy II:** Spinors and non-commutative geometry
- 4 **Epilogue:** Future directions

## Euclidean gravity in two dimensions

The gravity action can be written in terms of constrained BF theory:

$$S[A, B, \lambda] = \int_{\Sigma} \text{Tr}(BF) - \int_{\Sigma} \lambda (\text{Tr}(B^2) + 1),$$

invariant under the adjoint action of  $SO(2)$  with Lie algebra  $\mathfrak{so}(2)$ .

## Euclidean gravity in two dimensions

The gravity action can be written in terms of constrained BF theory:

$$S[A, B, \lambda] = \int_{\Sigma} \text{Tr}(BF) - \int_{\Sigma} \lambda (\text{Tr}(B^2) + 1),$$

invariant under the adjoint action of  $SO(2)$  with Lie algebra  $\mathfrak{so}(2)$ .

- $\Sigma$  is (typically) a compact surface
- $B$  is a scalar field;  $F$  the curvature of a connection  $A$
- $B$  and  $F$  are  $\mathfrak{so}(2)$ -valued
- $\text{Tr}: \mathfrak{so}(2) \rightarrow \mathbb{R}$  is a Killing form
- $\lambda$  is a volume form (Lagrange multiplier)

## Euclidean gravity in two dimensions

The gravity action can be written in terms of constrained BF theory:

$$S[A, B, \lambda] = \int_{\Sigma} \text{Tr}(BF) - \int_{\Sigma} \lambda (\text{Tr}(B^2) + 1),$$

invariant under the adjoint action of  $SO(2)$  with Lie algebra  $\mathfrak{so}(2)$ .

- $\Sigma$  is (typically) a compact surface
- $B$  is a scalar field;  $F$  the curvature of a connection  $A$
- $B$  and  $F$  are  $\mathfrak{so}(2)$ -valued
- $\text{Tr}: \mathfrak{so}(2) \rightarrow \mathbb{R}$  is a Killing form
- $\lambda$  is a volume form (Lagrange multiplier)

We are interested in the **quantum version**.

## Quantisation – ‘Path integral’ approach

The quantum theory is described by expectation values of operators:

$$\langle O \rangle_{\Sigma} = \int DADBD\lambda O[A, B, \lambda] e^{-iS[A, B, \lambda]}$$

**Topological** – all expectation values are **topological invariants**

# Purpose of this talk

## Problem

Can we describe operators encoding spinor information and find

$$\langle O \rangle_{\Sigma} = \int DADBD\lambda \int D\psi D\bar{\psi} O[A, B, \lambda, \psi, \bar{\psi}] e^{-iS[A, B, \lambda]}$$

**without** breaking topological invariance?

# Expectation values in BF theory

## Auxiliary problem

Calculate exactly the expectation values

$$\langle O \rangle_{\Sigma} = \int DADBD\lambda O[A, B, \lambda] e^{-i \int_{\Sigma} \text{Tr}(BF) + i \int_{\Sigma} \lambda (\text{Tr}(B^2) + 1)}$$

by giving them a **geometrical** interpretation.



# Expectation values in BF theory

## Auxiliary problem

Calculate exactly the expectation values

$$\langle O \rangle_{\Sigma} = \int DADB \int D\lambda O[A, B, \lambda] e^{i \int_{\Sigma} \lambda (\text{Tr}(B^2) + 1)} e^{-i \int_{\Sigma} \text{Tr}(BF)}$$

by giving them a **geometrical** interpretation.

# Expectation values in BF theory

## Auxiliary problem

Calculate exactly the expectation values

$$\langle O \rangle_{\Sigma} = \int DADBO[A, B] e^{-i \int_{\Sigma} \text{Tr}(BF)}$$

by giving them a **geometrical** interpretation.

## Expectation values in BF theory

### Auxiliary problem

Calculate exactly the expectation values

$$\langle O \rangle_{\Sigma} = \int DADB O[A, B] e^{-i \int_{\Sigma} \text{Tr}(BF)}$$

by giving them a **geometrical** interpretation.

**Starting point:** calculating the vacuum expectation value

$$\langle 1 \rangle_{\Sigma} = \int DADB e^{-i \int_{\Sigma} \text{Tr}(BF)} = \int DA \delta(F)$$

What is the importance of flat connections?

## Quantisation – ‘Canonical’ approach

$$\begin{array}{ccc} T^*(\mathcal{A}) & \xrightarrow{\text{quantise}} & L^2(\mathcal{A}) \\ \downarrow \text{constraint} & & \downarrow \text{constraint} \\ T^*(\mathcal{A}_0/\mathcal{G}) & \xrightarrow{\text{quantise}} & L^2(\mathcal{A}_0/\mathcal{G}) \end{array}$$

- $\mathcal{A}$  is the space of connections  $A$  (configuration space)
- $\mathcal{A}_0/\mathcal{G}$  – the **moduli space of flat connections**
  - ‘Moduli space’ refers to the invariance under the action of  $\mathcal{G}$
  - ‘Flat connection’ is the choice of ‘gauge’,  $F = 0$

## The vacuum expectation value

$$\langle 1 \rangle_{\Sigma} = \int DA \delta(F) = \text{vol}(\mathcal{A}_0/\mathcal{G})$$

**Breakthrough of rigorous path integral approaches:**

Using the **geometry of  $\Sigma$**  to probe the **geometry of  $\mathcal{A}_0/\mathcal{G}$**

## The vacuum expectation value

$$\langle 1 \rangle_{\Sigma} = \int DA \delta(F) = \text{vol}(\mathcal{A}_0/\mathcal{G})$$

### Breakthrough of rigorous path integral approaches:

Using the **geometry of  $\Sigma$**  to probe the **geometry of  $\mathcal{A}_0/\mathcal{G}$**

#### Key points

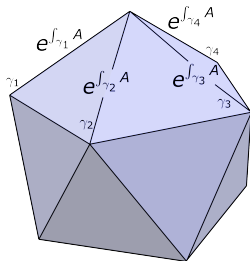
- Find some 'good' functions  $f(A)$  – parallel transport variables
- Investigate them along some 'regions' of  $\Sigma$  – edges of a triangulation
- Impose gauge invariance –  $\mathcal{G} = SO(2)$  – and gauge fixing

# The vacuum expectation value

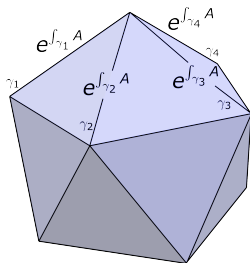
$$\langle 1 \rangle_{\Sigma} = \int DA \delta(F) = \text{vol}(\mathcal{A}_0/\mathcal{G})$$

**Breakthrough of rigorous path integral approaches:**

Using the **geometry of  $\Sigma$**  to probe the **geometry of  $\mathcal{A}_0/\mathcal{G}$**



# The vacuum expectation value



$$e^{\int_{\gamma} A} = g \in SO(2)$$

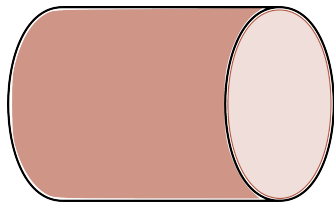
Basis for  $L^2(\mathcal{A})$ :  $\rho^i(g)^a_b$

This basis is also dense in  $\mathbb{R}[SO(2)]$ , the group algebra.

$$L^2(\mathcal{A}) \simeq \mathbb{R}[SO(2)]$$



## The vacuum expectation value



Loop variables:  $\chi^i(g) = \text{Tr} \rho^i(g)$

Invariance under  $SO(2)$ :  $\chi^i(g) = \chi^i(hgh^{-1})$

$$L^2(\mathcal{A}/SO(2)) \simeq Z(\mathbb{R}[SO(2)])$$

## The vacuum expectation value

$$\langle 1 \rangle_{\Sigma} = \int DA \delta(F) = \text{vol}(\mathcal{A}_0/\mathcal{G})$$

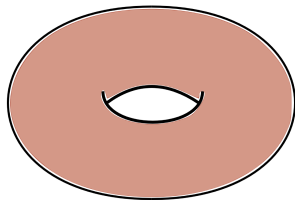
### What do we conclude?

**Geometrical** properties have been **translated into algebraic** ones

- The Hilbert spaces have extra structure – they are also algebras
- Triangulation independence  $\Leftrightarrow$  axioms of the algebra  $\mathbb{R}[SO(2)]$

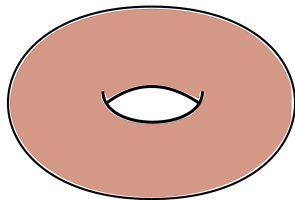
# Ansatz

$$\langle \mathbf{1} \rangle_{\Sigma}$$

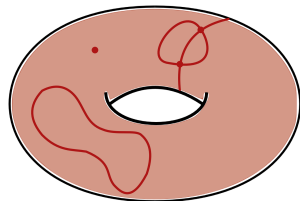


# Ansatz

$$\langle \mathbf{1} \rangle_{\Sigma}$$

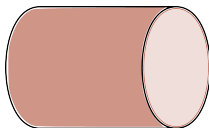


$$\langle \mathbf{0} \rangle_{\Sigma}$$



# Ansatz

The linear map

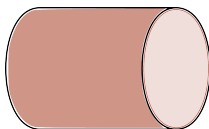


$$L^2(\mathcal{A}/SO(2)) \rightarrow L^2(\mathcal{A}/SO(2))$$

is an **operator in the Hilbert space**.

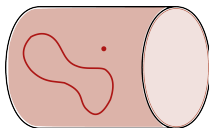
# Ansatz

The linear map



$$L^2(\mathcal{A}/SO(2)) \rightarrow L^2(\mathcal{A}/SO(2))$$

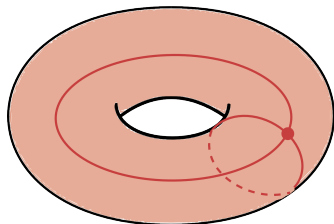
is an **operator in the Hilbert space**. It follows



$$L^2(\mathcal{A}/SO(2)) \rightarrow L^2(\mathcal{A}/SO(2))$$

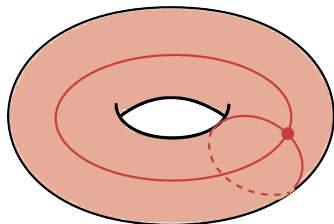
**must also be one.**

## What is the information encoded in the defects?



- **Defect lines:** vector spaces with inner product.
- The vector space carries left and right actions from  $\mathbb{R}[SO(2)]$
- **Vertices:** linear maps

# What is the information encoded in the defects?



## Some results:

- Expectation values constructed from defects are **homotopy invariant**
- Operator implementing **gravity constraint** has been constructed
- Simply connected cover data can be recovered



# What about the spinors?

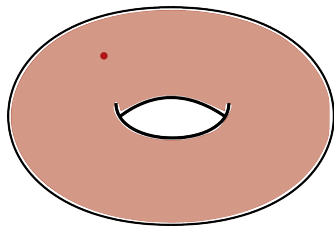
## Problem

Can we describe operators encoding spinor information and find

$$\langle O \rangle_{\Sigma} = \int DADBD\lambda \int D\psi D\bar{\psi} O[A, B, \lambda, \psi, \bar{\psi}] e^{-iS[A, B, \lambda]}$$

**without** breaking topological invariance?

# The role of non-commutative geometry



## Non-commutative geometry

\* algebra  $\mathcal{A}$

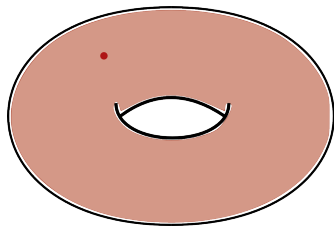
Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

# The role of non-commutative geometry



## Classical case

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

## Non-commutative geometry

$\star$  algebra  $\mathcal{A}$

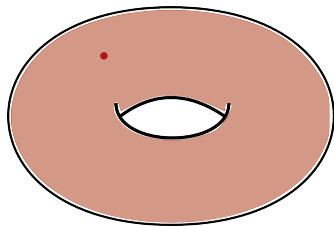
Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

# The role of non-commutative geometry



**Classical case**

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

**Non-commutative geometry**

$\star$  algebra  $\mathcal{A}$

Hilbert space  $\mathcal{H}$

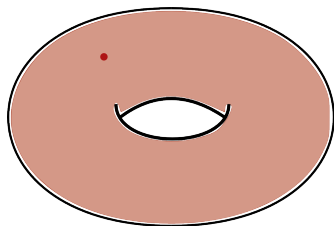
$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

**Defect formalism**

# The role of non-commutative geometry



**Classical case**

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

**Non-commutative geometry**

$\star$  algebra  $\mathcal{A}$

Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

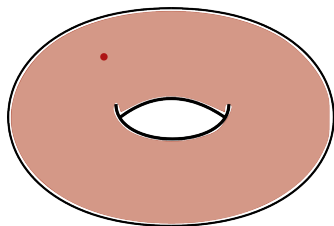
$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

**Defect formalism**

$\mathbb{R}[SO(2)]$

# The role of non-commutative geometry



**Classical case**

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

**Non-commutative geometry**

$\star$  algebra  $\mathcal{A}$

Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

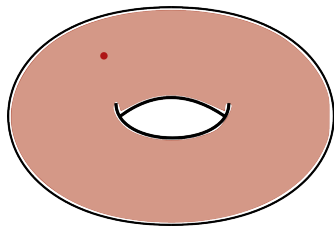
$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

**Defect formalism**

$\mathbb{R}[SO(2)]$

Vector space  $V$

# The role of non-commutative geometry



## Classical case

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

## Non-commutative geometry

$\star$  algebra  $\mathcal{A}$

Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

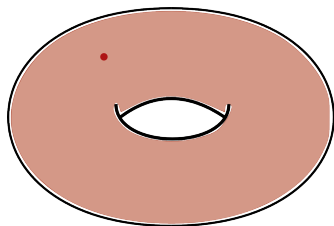
## Defect formalism

$\mathbb{R}[SO(2)]$

Vector space  $V$

Linear map  $V \rightarrow V$

# The role of non-commutative geometry



## Classical case

$C^\infty(\Sigma)$

Spinor fields on  $\Sigma$

Dirac operator

Charge conjugation

$\gamma^5$

## Non-commutative geometry

$\star$  algebra  $\mathcal{A}$

Hilbert space  $\mathcal{H}$

$D: \mathcal{H} \rightarrow \mathcal{H}$

$J: \mathcal{H} \rightarrow \mathcal{H}, J^2 = \pm 1$

$\Gamma: \mathcal{H} \rightarrow \mathcal{H}, \Gamma^2 = \pm 1$

## Defect formalism

$\mathbb{R}[SO(2)]$

Vector space  $V$

Linear map  $V \rightarrow V$

Specific  $V$

Specific  $V$



## Work in progress:

- Find all defect vector spaces  $V$  satisfying the NC axioms
- Use non-orientable case to suppress false candidates
- Relate expectation values with classical fermionic fields

# Thank you!

Sara Tavares  
sara.orian@gmail.com