On the stability and topology of the universe

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Introduction Questions Vlasov matter Model solutions, approximating fluids Causality

Introduction

The standard models of the universe

- satisfy the cosmological principle (i.e., they are spatially homogeneous and isotropic),
- are spatially flat,
- have matter content consisting of ordinary matter, dark matter and dark energy.

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Introduction

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Current model of the universe



Current model of the universe: NASA/WMAP Science Team.

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Questions

Since the universe is not exactly spatially homogeneous and isotropic, it is natural to ask: are the standard models future stable?

Since the assumption of spatial homogeneity and isotropy yields strong restrictions on the allowed topologies, it is also of interest to ask: what are the restrictions on the global topology of the universe imposed by the constraint that what we observe seems to be close to a standard model?

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Matter models

Perfect fluids: Matter described by energy density ρ and pressure \mathfrak{p} . *Dust:* $\mathfrak{p} = 0$. *Radiation:* $\mathfrak{p} = \rho/3$.

Vlasov matter: collection of particles, where

- the particles all have unit mass,
- collisions are neglected,
- the particles follow geodesics,
- collection described statistically by a distribution function.

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Stress energy tensor, Vlasov matter

In the Vlasov setting, the relevant mathematical structures are

- the mass shell P; the future directed unit timelike vectors in (M,g),
- the distribution function $f: P \to [0, \infty)$,
- the stress energy tensor

$$T_{\alpha\beta}|_{\xi} = \int_{P_{\xi}} f p_{\alpha} p_{\beta} \mu_{P_{\xi}},$$

the Vlasov equation

$$\mathcal{L}f=0.$$

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The Einstein-Vlasov system

The Einstein-Vlasov system consists of the equations

$$G + \Lambda g = T,$$

$$\mathcal{L}f = 0$$

for g and f. Note that the second equation corresponds to the requirement that f be constant along timelike geodesics.

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Standard models

Spatial homogeneity, isotropy and flatness imply that the metric takes the form

$$g=-dt^2+a^2(t)\bar{g}.$$

The stress energy tensor of the Vlasov matter then takes perfect fluid form, and the energy density and pressure are given by

$$egin{array}{rll}
ho_{ ext{Vl}}(t) &=& \int_{\mathbb{R}^3} ar{\mathsf{f}}\left(\mathsf{a}(t)ar{q}
ight) (1+|ar{q}|^2)^{1/2} dar{q}, \ \mathfrak{p}_{ ext{Vl}}(t) &=& rac{1}{3} \int_{\mathbb{R}^3} ar{\mathsf{f}}\left(\mathsf{a}(t)ar{q}
ight) rac{|ar{q}|^2}{(1+|ar{q}|^2)^{1/2}} dar{q}, \end{array}$$

where $\overline{f},$ a function on $\mathbb{R}^3,$ is the initial datum for the distribution function.

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Approximating fluids



Figure: An illustration of an initial datum for the distribution function which is appropriate when approximating a standard model.

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Minkowski space; non-silent causality

Let $\gamma(t) = (t, 0, 0, 0)$. Then γ is an observer in Minkowski space. How much of the t = 0 hypersurface does γ see?



Figure: The causal past of $\gamma(t)$ intersected with the causal future of the t = 0 hypersurface for t = 1/2, t = 1 and t = 2.

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de Sitter space; silent causality

Consider the metric

 $g = -dt^2 + e^{2t}\bar{g}.$



Figure: The causal past of $\gamma(t)$ intersected with the causal future of the t = 0 hypersurface for t = 1/2 and for all t.

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Induced initial data

Let (M, g, f) be a solution and Σ be a spacelike hypersurface in (M, g). Then the **initial data induced on** Σ consist of

- the Riemannian metric induced on Σ by g, say \overline{g} ,
- the second fundamental form induced on Σ by g, say \bar{k} ,
- the induced distribution function $\overline{f} : T\Sigma \to [0,\infty)$.

Here

$$\bar{f} = f \circ \operatorname{proj}_{\Sigma}^{-1},$$

where $\mathrm{proj}_{\Sigma}: \mathcal{P}_{\Sigma} \to \mathcal{T}\Sigma$ represents projection orthogonal to the normal.

Induced initial data Norms

Projection to the tangent space



Figure: In suitable coordinates, proj_{Σ} corresponds to the ordinary projection to the p^1p^2 -plane in the figure.

Induced initial data Norms

Function spaces

If Σ is a compact manifold, $\overline{\mathfrak{D}}^{\infty}_{\mu}(T\Sigma)$ denotes the space of smooth functions $f: T\Sigma \to \mathbb{R}$ such that

$$\|\bar{f}\|_{H^{I}_{\mathrm{V1},\mu}} = \left(\sum_{i=1}^{j} \sum_{|\alpha|+|\beta| \leq I} \int_{\bar{\mathbf{x}}_{i}(U_{i}) \times \mathbb{R}^{n}} \langle \bar{\varrho} \rangle^{2\mu+2|\beta|} \bar{\chi}_{i}(\bar{\xi}) (\partial_{\bar{\xi}}^{\alpha} \partial_{\bar{\varrho}}^{\beta} \bar{\mathbf{f}}_{\bar{\mathbf{x}}_{i}})^{2} (\bar{\xi}, \bar{\varrho}) d\bar{\xi} d\bar{\varrho} \right)^{1/2}$$

is finite for every $l \ge 0$, where

$$\langle \bar{\varrho} \rangle = (1 + |\bar{\varrho}|^2)^{1/2}.$$

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Previous results, exponential expansion Future stability Topology

Previous results, exponential expansion

- Stability of de Sitter space in 3 + 1-dimensions, etc., Helmut Friedrich, '86, '91.
- Stability of even dimensional de Sitter spaces, Michael Anderson '05.
- Stability in the non-linear scalar field case, H.R. '08.
- Einstein-Euler with a positive cosmological constant, Igor Rodnianski and Jared Speck, *preprint* '09.

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Previous results, exponential expansion Future stability Topology

Future stability, assumptions

Let $(G, \bar{g}_{\rm bg}, \bar{k}_{\rm bg}, \bar{f}_{\rm bg})$ be initial data corresponding to an expanding spatially homogeneous solution to the Einstein–Vlasov system. Assume, that

- the solution is neither of Bianchi type IX nor of Kantowski Sachs type,
- there is a cocompact subgroup Γ of the isometry group of the initial data.

Let $\Sigma = G/\Gamma$ be the compact quotient.

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Future stability, conclusions

Then there is an $\epsilon > 0$ such that if $(\Sigma, \bar{g}, \bar{k}, \bar{f})$ are initial data satisfying

$$\|\bar{g}-\bar{g}_{\mathrm{bg}}\|_{H^5}+\|\bar{k}-\bar{k}_{\mathrm{bg}}\|_{H^4}+\|\bar{f}-\bar{f}_{\mathrm{bg}}\|_{H^4_{\mathrm{Vl},\mu}}\leq\epsilon,$$

then the maximal Cauchy development of $(\Sigma, \overline{g}, \overline{k}, \overline{f})$ is future causally geodesically complete. Moreover, the solution is asymptotically de Sitter like.

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What is the shape of the universe?



Previous results, exponential expansion Future stability Topology

Question, topology

Assume that

- the observational data indicate that, to our past, the universe is well approximated by one of the standard models,
- ► interpreting the data in this model, we only have information concerning the universe for t ≥ t₀,
- there is a big bang,
- analogous statements apply to all observers in the universe (with the same t₀).

The question is then: what conclusions are we allowed to draw concerning the global spatial topology of the universe?

Previous results, exponential expansion Future stability Topology

Ingredients

Assume we are given

- a standard model, characterised by an existence interval *I*, a scale factor *a* etc.,
- ► a t₀ ∈ *I*, which represents the time to the future of which we wish the approximation to be valid,
- An *l* ∈ N, specifying the norm with respect to which we measure proximity to the standard model,
- an $\epsilon > 0$, characterising the size of the distance,
- a closed 3-manifold Σ .

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Construction

There is a solution (M, g, f) with the following properties:

- (M, g, f) is a maximal Cauchy development,
- (M, g) is future causally geodesically complete,
- ► there is a Cauchy hypersurface, say S
 , in (M,g), diffeomorphic to Σ,
- ▶ given an observer \(\gamma\) in (M,g), there is a neighbourhood, say U, of

 $J^-(\gamma)\cap J^+(\bar{S})$

such that the solution in U is ϵ -close to the standard model in a solid cylinder of the form $[t_0, \infty) \times \overline{B}_R(0)$,

- all timelike geodesics in (M, g) are past incomplete,
- the solution is stable with these properties.

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Questions Expanding direction The big bang Hierarchy

Questions

- Strong cosmic censorship; do initial data generically yield an inextendible maximal Cauchy development?
- Curvature blow up; does the curvature blow up in the incomplete directions of causal geodesics (in the maximal Cauchy development)?
- Stability; given a certain behaviour at the singularity or in the expanding direction, is it stable under perturbations?

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Asymptotic aspects

- Silence; do horizons form in the direction of the singularity? Do observers loose the ability to communicate in the expanding direction?
- Isotropy; Does the solution isotropize or not?
- Convergence; Is the behaviour convergent or oscillatory?

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Expanding vacuum solutions; a non-silent situation

Let Σ be a closed hyperbolic manifold and \bar{g}_H be a (suitable) hyperbolic metric on Σ . Then

$$g = -dt^2 + t^2 \bar{g}_H$$

is a solution to Einstein's vacuum equations.

This solution is future stable; Andersson and Moncrief '04.

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Hyperbolic dominance/isotropization

Due to geometrization, it is possible to **cut up** a closed 3-manifold and to endow the pieces with preferred geometries.

There is a general conjecture relating the asymptotic behaviour of vacuum solutions to such a division; due to *Moncrief, Fischer and Anderson.*

Rough conjecture: all the volume is in the hyperbolic pieces asymptotically. The solution isotropizes on the hyperbolic pieces.

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Geometric decomposition



Figure: A schematic depiction of a geometric decomposition; H_1 , H_2 and H_3 represent the hyperbolic pieces and S the Seifert manifold pieces.

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Support

- Andersson and Moncrief '05,
- Choquet-Bruhat and Moncrief '01,
- ▶ ...,
- Anderson '01,
- Reiris '10,
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Horizons and CMB

Misner '69:

... if the 3°K background radiation were last scattered at a redshift z = 7, then the radiation coming to us from two directions in the sky separated by more than about 30° was last scattered by regions of plasma whose prior histories had no causal relationship. [...] Robertson-Walker models therefore give no insight into why the observed microwave radiation from widely different angles in the sky has very precisely ($\leq 0.2\%$) the same temperature.

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Mixmaster and BKL

Misner: Mixmaster solutions (Bianchi IX) do not form horizons.

BKL: In generic collapse horizons form and the local behaviour is well described by Bianchi IX solutions (which are oscillatory).

Caveat: Convergent behaviour for scalar fields, stiff fluids and in higher dimensions.

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Support

Oscillatory:

- Belinskii, Khalatnikov, Lifschitz (BKL), '70, '82,
- Misner '69,
- Chitre '72,
- Damour, Henneaux, Nicolai '03,
- Heinzle, Uggla, Rohr '09,
- ► ...

Non-oscillatory:

- Andersson, Rendall '01.
- Damour, Henneaux, Rendall, Weaver '02.

Questions Expanding direction The big bang **Hierarchy**

Hierarchy

- Silent, convergent and isotropic,
- Silent, convergent and anisotropic,
- Non-silent, partially convergent and partially isotropic,
- Silent (?), oscillatory and anisotropic.

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Questions Expanding direction The big bang **Hierarchy**

Thank you!

Hans Ringström On the stability and topology of the universe

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