

# Higher Conformal Couplings from Weyl Gauge Theory



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LIII School, 29.06.2013, Zakopane

# Outline

- Motivation
- Weyl gauge theory
- General method
- Applications
- Conclusions

# Motivation

- On flat spacetime:  
no dimensionfull couplings  $\Rightarrow$  scale invariance  
Symmetry may be enhanced to full conformal (Iorio, O'Raiheartaigh, Sachs, Wiesendanger)
- On curved spacetime:  
Conformal (Weyl) rescalings, coupling to spacetime curvature

$$L = (\partial \phi)^2 + \frac{d-2}{4(d-1)} R \phi^2$$

- How to obtain this non-minimal couplings preserving conformal covariance?

# Motivation

Various methods have been used:

- BRST covariant algebra for Weyl-gravity (Boulanger)
  - Unfolded Dynamics Approach (Shaynkman, Tipunin, Vasiliev)
  - Tractor calculus (Gover, Shaukat, Waldron)
  - Ambient space techniques (Fefferman, Graham)
  - Ricci gauging (Iorio, O'Raiheartaigh, Sachs, Wiesendanger)
  - Others (Manvelyan, Mkrtchyan, Erdmenger)
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- In this talk method based on decoupling of Weyl gauge fields (related to Weyl geometry)

# Weyl gauge theory

- Gauge transformation of lengths  $ds'^2 = \Omega^2 ds^2$
- Gauge transformations on fields  $g'_{\mu\nu} = \Omega^2 g_{\mu\nu}$   $\Phi' = \Omega^w \Phi$
- If  $\Omega = \Omega(x)$  then need to introduce abelian Weyl gauge potential  $b_\mu$
- Weyl curvature (field strength)  $B_{\mu\nu} = \frac{1}{2}(\partial_\mu b_\nu - \partial_\nu b_\mu)$
- Coordinates and derivatives do not transform  $x'^\mu = x^\mu$   $\partial'_\mu = \partial_\mu$
- However Riemannian connection is not Weyl covariant!

$$\Gamma'^{\mu}_{LC\nu\rho} \neq \Gamma^{\mu}_{LC\nu\rho}$$

# Weyl gauge theory

- Unique, symmetric, torsionfree linear connection  $\Rightarrow$  Weyl connection

$$\tilde{\Gamma}_{\nu\rho}^{\mu} = \Gamma_{\text{LC}\nu\rho}^{\mu} - \delta_{\nu}^{\mu} b_{\rho} - \delta_{\rho}^{\mu} b_{\nu} + g_{\nu\rho} b^{\mu}$$

- 1st consequence: Weyl connection is Weyl invariant  $\tilde{\Gamma}'^{\mu}_{\nu\rho} = \tilde{\Gamma}^{\mu}_{\nu\rho}$

- 2nd consequence: cov. derivative  $\tilde{\nabla}_{\mu}$  with  $\tilde{\Gamma}^{\mu}_{\nu\rho}$  preserve Weyl symmetry

- Weyl covariant derivatives

on scalar fields:  $D_{\mu} \phi = (\partial_{\mu} - w b_{\mu}) \phi$

on tensor fields:  $D_{\mu} T = (\tilde{\nabla}_{\mu} - w b_{\mu}) T$

- Transformation law of cov. Weyl derivatives of any tensor  $T$  (indices suppressed)

$$(D_{\mu} T)' = \Omega^w D_{\mu} T \quad T' = \Omega^w T$$

# Weyl curvatures

- Riemann-Weyl curvature tensor  $[D_\mu, D_\nu]v^\rho = r_{\mu\nu}{}^\rho{}_\sigma v^\sigma$
- Ricci-Weyl tensor and scalar  $r_{\nu\sigma} = r_{\mu\nu}{}^\mu{}_\sigma$   $r = g^{\nu\sigma} r_{\nu\sigma}$
- Riemann-Weyl and Ricci-Weyl tensors are conf. inv.,
- Ricci-Weyl scalar is conf. cov. with  $w=-2$
- Trivial Weyl gauge field  $B_{\mu\nu} = 0$   
 $\Rightarrow$  Weyl curvatures in terms of ordinary spacetime curvatures and

Weyl gauge potential  $b_\mu$

$$r_{\mu\nu\rho\sigma} = R_{\mu\nu\rho\sigma} + g_{\mu\rho}(\nabla_\nu b_\sigma + b_\nu b_\sigma) - g_{\mu\sigma}(\nabla_\nu b_\rho + b_\nu b_\rho) - g_{\nu\rho}(\nabla_\mu b_\sigma + b_\mu b_\sigma) + g_{\nu\sigma}(\nabla_\mu b_\rho + b_\mu b_\rho) - (g_{\mu\rho}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\rho})b^2$$

$$r_{\mu\nu} = R_{\mu\nu} + (d-2)b_\mu b_\nu + (d-2)\nabla_\mu b_\nu - (d-2)b^2 g_{\mu\nu} + (\nabla \cdot b) g_{\mu\nu}$$

$$r = R + 2(d-1)(\nabla \cdot b) - (d-1)(d-2)b^2$$

# General method

- 1) Start with generally covariant expression  $X$
- 2) Weyl covariantize all derivatives in your expression  $\nabla_\mu \rightarrow D_\mu$
- 3) Impose condition of vanishing Weyl gauge field strength  $B_{\mu\nu} = 0$
- 4) Add all possible Weyl covariant terms  $T_i$  with the same weight  $w$  constructed out of Weyl curvatures, Weyl derivatives and original fields in  $X$   
 $X + a_i T_i \quad w(T_i) = w(X)$   
2 Weyl derivatives  $D \Rightarrow$  one Weyl curvature
- 5) Expand resulting sum in Weyl gauge potentials –  $b_\mu$  fields
- 6) Search for coefficients  $a_i$  and  $w$ , s.t. Weyl potentials decouple and terms with  $b_\mu$  vanish
- 7) If a solution exists, then  $X$  conformally coupled to spacetime is given by  $X + a_i T_i$



# General method

- Linear algebra problem
- Applicable to any expression w/o dimensionful constants
- Works for fields with any spin and in any spacetime dimension
- Found conformal weights  $w$  agree with canonical dimensions of fields as assigned on flat spacetime
- Shows obvious links with more general gravitation-dilaton-conformal theory (Weyl gauge theory of gravitation)

# Applications

- Scalar field with 4 derivatives  $L = \phi \square^2 \phi$   $\square^2 \phi = 0$
- Old result from  $d=4$ , when  $w(\phi)=0$

$$\square^2 \phi + \frac{1}{3} (\nabla_\mu R) (\nabla^\mu \phi) + 2 \text{Ric}^{\mu\nu} \nabla_\mu \nabla_\nu \phi - \frac{2}{3} R \square \phi = 0$$

- In general dimension 3 new terms must be added

$$r^2 \phi \quad \text{ric}^2 \phi \quad (\square r) \phi \quad \Delta_{(4)} \phi = 0$$

- General result

$$\Delta_{(4)} \phi = \square^2 \phi + A R^2 \phi + B \text{Ric}^2 \phi + C (\square R) \phi + D (\nabla_\mu R) (\nabla^\mu \phi) + E \text{Ric}^{\mu\nu} \nabla_\mu \nabla_\nu \phi + F R \square \phi$$

- Dimension-dependent coefficients

$$A = \frac{(d-4)(d^3 - 4d^2 + 16d - 16)}{(4(d-1)(d-2))^2} \quad B = \frac{4-d}{(d-2)^2} \quad C = \frac{4-d}{4(d-1)} \quad D = \frac{6-d}{2(d-1)}$$

$$E = \frac{4}{d-2} \quad F = -\frac{(d-2)^2 + 4}{2(d-1)(d-2)} \quad w = \frac{4-d}{2}$$

# Applications 2

- Abelian gauge field in  $d=6$  with Minkowski spacetime action  $F \nabla^2 F$   
 Gauge fields have no conformal weights  $w(A_\mu) = w(F_{\mu\nu}) = 0$   
 However ordinary cov. derivatives are not conf. cov.

$$D_\mu A_\rho = \tilde{\nabla}_\mu A_\rho \neq \nabla_\mu A_\rho$$

- Conformally covariant lagrangian exists

Terms to combine  $F^{\rho\sigma} \square F_{\rho\sigma}$   $F^{\rho\sigma} \nabla_\rho \nabla_\mu F^\mu_\sigma$

- In our method we add  $r F^{\rho\sigma} F_{\rho\sigma}$   $\text{ric}_{\rho\mu} F^{\rho\sigma} F^\mu_\sigma$

- Final conf. coupled lagrangian with weight  $w=-6$  on curved spacetime

$$F^{\rho\sigma} \left( \square - \frac{R}{10} \right) F_{\rho\sigma} - 2 F^{\rho\sigma} \left( \nabla_\rho \nabla_\mu + \frac{R_{\rho\mu}}{2} \right) F^\mu_\sigma$$

# Summary

- Conformal couplings to curved spacetimes
- Weyl gravitation as a gauge theory with Weyl rescalings
- Weyl curvatures and Weyl geometry
- Decoupling of Weyl gauge potential as an useful method for finding conformally coupled lagrangians and equations
- Applications to old and new results

Thank you  
for attention!