

Kerr Killing tensor and its conserved charge at \mathcal{I}^+

Aruna Kesavan

The Pennsylvania State University

July 04, 2013

Work done in collaboration with Abhay Ashtekar.

- Physical problem
- Killing vector fields and conserved quantities
- Asymptotic Kerr Killing tensor and conserved quantity
- Summary and Outlook

- **Physical system:** Initial conditions that evolve under Einstein's equations and settle down finally to a Kerr blackhole.

- **Physical system:** Initial conditions that evolve under Einstein's equations and settle down finally to a Kerr blackhole.
- **Problem:** Associate conserved charge with Kerr Killing tensor.

- **Killing vectors** K^a : Isometries of spacetime.
 - $\mathcal{L}_K g_{ab} = 0$ or $\nabla_{(a} K_{b)} = 0$

- **Killing vectors** K^a : Isometries of spacetime.
 - $\mathcal{L}_K g_{ab} = 0$ or $\nabla_{(a} K_{b)} = 0$
 - $t^a = (\partial_t)^a$, $\phi_3^a = (\partial_\phi)^a$ correspond to stationarity and axi-symmetry in Kerr.

- **Killing vectors K^a :** Isometries of spacetime.
 - $\mathcal{L}_K g_{ab} = 0$ or $\nabla_{(a} K_{b)} = 0$
 - $t^a = (\partial_t)^a$, $\phi_3^a = (\partial_\phi)^a$ correspond to stationarity and axi-symmetry in Kerr.
- **Conserved charges:**
 - **Komar mass for t^a :** $Q_t = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c t^d dS^{ab}$
 S is a 2 surface enclosing the blackhole.

- **Killing vectors** K^a : Isometries of spacetime.

- $\mathcal{L}_K g_{ab} = 0$ or $\nabla_{(a} K_{b)} = 0$
- $t^a = (\partial_t)^a$, $\phi_3^a = (\partial_\phi)^a$ correspond to stationarity and axi-symmetry in Kerr.

- **Conserved charges:**

- **Komar mass for t^a :** $Q_t = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c t^d dS^{ab}$
 S is a 2 surface enclosing the blackhole.
- Works in the Kerr region, but not in the dynamical region.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .
- τ^a exists on \mathcal{I}^+ even when t^a does not exist as a spacetime Killing vector field.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .
- τ^a exists on \mathcal{I}^+ even when t^a does not exist as a spacetime Killing vector field.
- In general, Bondi energy: $Q_{\tau, dyn} = -\frac{1}{8\pi} \int_S C_{abcd} l^a \tau^b l^c \tau^d +$
Radiation part (Ashtekar, Streubel 1981)

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .
- τ^a exists on \mathcal{I}^+ even when t^a does not exist as a spacetime Killing vector field.
- In general, Bondi energy: $Q_{\tau, dyn} = -\frac{1}{8\pi} \int_S C_{abcd} l^a \tau^b l^c \tau^d +$
Radiation part (Ashtekar, Streubel 1981)
 C_{abcd} : Weyl tensor, l^a null vector satisfying $l_a \tau^a = -1$.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .
- τ^a exists on \mathcal{I}^+ even when t^a does not exist as a spacetime Killing vector field.
- In general, Bondi energy: $Q_{\tau, dyn} = -\frac{1}{8\pi} \int_S C_{abcd} l^a \tau^b l^c \tau^d +$
Radiation part (Ashtekar, Streubel 1981)
 C_{abcd} : Weyl tensor, l^a null vector satisfying $l_a \tau^a = -1$.
- ADM mass – Bondi mass = Radiation energy flux.

Charges at null infinity

- Go to future null infinity, \mathcal{I}^+ , consider generator τ^a , which in the limit is same as t^a .
- Energy at null infinity: $Q_\tau = -\frac{1}{8\pi} \int_S \epsilon_{abcd} \nabla^c \tau^d dS^{ab}$.
- Same as Komar mass in the Kerr region.
- **CAN** be defined in the dynamical region!
- τ^a is an asymptotic symmetry vector field (BMS generator), can be rigidly transported on all of \mathcal{I}^+ .
- τ^a exists on \mathcal{I}^+ even when t^a does not exist as a spacetime Killing vector field.
- In general, Bondi energy: $Q_{\tau, \text{dyn}} = -\frac{1}{8\pi} \int_S C_{abcd} l^a \tau^b l^c \tau^d +$
Radiation part (Ashtekar, Streubel 1981)
 C_{abcd} : Weyl tensor, l^a null vector satisfying $l_a \tau^a = -1$.
- ADM mass – Bondi mass = Radiation energy flux.
- Similarly ∂_ϕ gives Angular momentum.

- **Killing tensor:** Generalisation of Killing vector field.

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a}K_{bc)} = 0$$

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a}K_{bc)} = 0$$

Trivial Example: Sum of products of KVF's.

$K^{ab} = \phi_1^a\phi_1^b + \phi_2^a\phi_2^b + \phi_3^a\phi_3^b$ in any spherically symmetric spacetime gives J^2 , square of the total angular momentum.

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a}K_{bc)} = 0$$

Trivial Example: Sum of products of KVF's.

$K^{ab} = \phi_1^a\phi_1^b + \phi_2^a\phi_2^b + \phi_3^a\phi_3^b$ in any spherically symmetric spacetime gives J^2 , square of the total angular momentum.

- Kerr has a **non-trivial** Killing tensor.

$$K^{ab} = 2\rho^2\tilde{l}^{(a}\tilde{n}^{b)} + r^2\tilde{g}^{ab};$$

\tilde{l}^a, \tilde{n}^a principal null directions with $\partial_t, \partial_\phi$ and ∂_r components.

Kerr Killing tensor

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a}K_{bc)} = 0$$

Trivial Example: Sum of products of KVF's.

$K^{ab} = \phi_1^a\phi_1^b + \phi_2^a\phi_2^b + \phi_3^a\phi_3^b$ in any spherically symmetric spacetime gives J^2 , square of the total angular momentum.

- Kerr has a **non-trivial** Killing tensor.

$$K^{ab} = 2\rho^2\tilde{l}^{(a}\tilde{n}^{b)} + r^2\tilde{g}^{ab};$$

\tilde{l}^a, \tilde{n}^a principal null directions with $\partial_t, \partial_\phi$ and ∂_r components.

- NOT trivially obtained from Killing vectors ∂_t and ∂_ϕ .

Kerr Killing tensor

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a} K_{bc)} = 0$$

Trivial Example: Sum of products of KVF's.

$K^{ab} = \phi_1^a \phi_1^b + \phi_2^a \phi_2^b + \phi_3^a \phi_3^b$ in any spherically symmetric spacetime gives J^2 , square of the total angular momentum.

- Kerr has a **non-trivial** Killing tensor.

$$K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab};$$

\tilde{l}^a, \tilde{n}^a principal null directions with $\partial_t, \partial_\phi$ and ∂_r components.

- NOT trivially obtained from Killing vectors ∂_t and ∂_ϕ .

- **Problem:**

- For **trivial** KT, say, $K_{ab} = \sum_{i,j} v_i^{(a} v_j^{b)}$, Conserved charge

$$Q_K = \sum_{i,j} Q_{v_i} Q_{v_j}$$

- **Killing tensor:** Generalisation of Killing vector field.

$$\nabla_{(a} K_{bc)} = 0$$

Trivial Example: Sum of products of KVF's.

$K^{ab} = \phi_1^a \phi_1^b + \phi_2^a \phi_2^b + \phi_3^a \phi_3^b$ in any spherically symmetric spacetime gives J^2 , square of the total angular momentum.

- Kerr has a **non-trivial** Killing tensor.

$$K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab};$$

\tilde{l}^a, \tilde{n}^a principal null directions with $\partial_t, \partial_\phi$ and ∂_r components.

- NOT trivially obtained from Killing vectors ∂_t and ∂_ϕ .

- **Problem:**

- For **trivial** KT, say, $K_{ab} = \sum_{i,j} v_i^{(a} v_j^{b)}$, Conserved charge $Q_K = \sum_{i,j} Q_{v_i} Q_{v_j}$
- For non-trivial Killing tensors, NO general prescription!

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

$$K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - (a \cos\theta \tau^{(a})(a \cos\theta \tau^{b)})$$

where ϕ_1, ϕ_2, ϕ_3 are rotational fields, τ^a is the generator of null infinity, a : angular momentum parameter of Kerr BH, θ : polar coordinate.

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

$$K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - (a \cos\theta \tau^{(a})(a \cos\theta \tau^{b)})$$

where ϕ_1, ϕ_2, ϕ_3 are rotational fields, τ^a is the generator of null infinity, a : angular momentum parameter of Kerr BH, θ : polar coordinate.

- Product of generators of **BMS** group, the *asymptotic* symmetry group!

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

$$K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - (a \cos\theta \tau^{(a})(a \cos\theta \tau^{b)})$$

where ϕ_1, ϕ_2, ϕ_3 are rotational fields, τ^a is the generator of null infinity, a : angular momentum parameter of Kerr BH, θ : polar coordinate.

- Product of generators of **BMS** group, the *asymptotic* symmetry group!
 - Preserves universal structure of asymptotic flatness at null infinity. (BMS, Ashtekar etc.)

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

$$K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a)})(\phi_3^{b)} + a\tau^{b)}) - (a \cos\theta \tau^{(a)})(a \cos\theta \tau^{b)})$$

where ϕ_1, ϕ_2, ϕ_3 are rotational fields, τ^a is the generator of null infinity, a : angular momentum parameter of Kerr BH, θ : polar coordinate.

- Product of generators of **BMS** group, the *asymptotic* symmetry group!
 - Preserves universal structure of asymptotic flatness at null infinity. (BMS, Ashtekar etc.) Larger than the Poincaré group.

Killing tensor at \mathcal{I}^+

- $K^{ab} = 2\rho^2 \tilde{l}^{(a} \tilde{n}^{b)} + r^2 \tilde{g}^{ab}$

This is **NOT** a trivial sum of products of Killing vector fields.

- However, at \mathcal{I}^+ ..

$$K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a)})(\phi_3^{b)} + a\tau^{b)}) - (a \cos\theta \tau^{(a)})(a \cos\theta \tau^{b)})$$

where ϕ_1, ϕ_2, ϕ_3 are rotational fields, τ^a is the generator of null infinity, a : angular momentum parameter of Kerr BH, θ : polar coordinate.

- Product of generators of **BMS** group, the *asymptotic* symmetry group!
 - Preserves universal structure of asymptotic flatness at null infinity. (BMS, Ashtekar etc.) Larger than the Poincaré group.
- Associate charge to Killing tensor as follows:
 $Q_K = \sum_i Q_{V_i} Q_{V_i}$ where $Q_{V_i} =$ Weyl part + Radiation part.

- $K^{ab} \rightarrow$
 $\phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - a^2 \cos^2 \theta \tau^{(a} \tau^{b)}$

Asymptotic Kerr Killing tensor

- $K^{ab} \rightarrow$
 $\phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - a^2 \cos^2 \theta \tau^{(a} \tau^{b)}$
- Naively, due to asymptotic flatness, may expect square of total angular momentum, but it is **NOT!**

Asymptotic Kerr Killing tensor

- $K^{ab} \rightarrow$
 $\phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - a^2 \cos^2 \theta \tau^{(a} \tau^{b)}$
- Naively, due to asymptotic flatness, may expect square of total angular momentum, but it is **NOT!**
- In addition to angular momentum, it contains information about z direction momentum!

Asymptotic Kerr Killing tensor

- $K^{ab} \rightarrow \phi_1^{(a} \phi_1^{b)} + \phi_2^{(a} \phi_2^{b)} + (\phi_3^{(a} + a\tau^{(a})(\phi_3^{b)} + a\tau^{b)}) - a^2 \cos^2 \theta \tau^{(a} \tau^{b)}$
- Naively, due to asymptotic flatness, may expect square of total angular momentum, but it is **NOT!**
- In addition to angular momentum, it contains information about z direction momentum!
- For Kerr, $Q_K = 4(Ma)^2$

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .
- Kerr Killing tensor decomposes into sum of product of asymptotic symmetry vector fields.

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .
- Kerr Killing tensor decomposes into sum of product of asymptotic symmetry vector fields.
- Can then compute charge of Killing tensor at various stages of blackhole formation!

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .
- Kerr Killing tensor decomposes into sum of product of asymptotic symmetry vector fields.
- Can then compute charge of Killing tensor at various stages of blackhole formation!
- Use this to write a conservation law.

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .
- Kerr Killing tensor decomposes into sum of product of asymptotic symmetry vector fields.
- Can then compute charge of Killing tensor at various stages of blackhole formation!
- Use this to write a conservation law.
- Conserved charge is **NOT** square of total angular momentum.

Summary

- Blackhole formed in asymptotic future provides asymptotic timelike Killing vector field τ^a and rotation Killing vector field which leaves τ^a invariant.
- Rigidity of asymptotic symmetry vector fields then determines the vector fields throughout \mathcal{I}^+ .
- Kerr Killing tensor decomposes into sum of product of asymptotic symmetry vector fields.
- Can then compute charge of Killing tensor at various stages of blackhole formation!
- Use this to write a conservation law.
- Conserved charge is **NOT** square of total angular momentum.
- It contains information about linear momentum: **blackhole kick?**

- Charge of Kerr Killing tensor at blackhole horizon.

- Charge of Kerr Killing tensor at blackhole horizon.
- Calculate rate of change of Carter's constant, solve for inspiralling geodesic in EMRI calculations.

THANK YOU!

- **Conserved charge:** For test particles along geodesics

$$Q_K = K^{a_1 \dots a_n} p_{a_1} \dots p_{a_n}$$

$$p_a \nabla^a Q_K = 0$$

- **Conserved charge:** For test particles along geodesics

$$Q_K = K^{a_1 \dots a_n} p_{a_1} \dots p_{a_n}$$

$$p_a \nabla^a Q_K = 0$$