

Two texture zeros in neutrino mass matrices

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Lepton mass terms (we assume **Majorana neutrinos**):

$$\begin{aligned}\mathcal{L} &= -\bar{\ell}'_R \mathcal{M}_\ell \ell'_L + \frac{1}{2} \nu'^T_L C^{-1} \mathcal{M}_\nu \nu'_L + \text{H.c.} = \\ &= -\bar{\ell}'_R \hat{\mathcal{M}}_\ell \ell_L + \frac{1}{2} \nu'^T_L C^{-1} \hat{\mathcal{M}}_\nu \nu_L + \text{H.c.}\end{aligned}$$

Diagonalization by a biunitary transformation:

$$\begin{aligned}\mathcal{M}_\ell &= U_R^\ell \hat{\mathcal{M}}_\ell U_L^{\ell\dagger}, \quad \mathcal{M}_\nu = U_L^{\nu*} \hat{\mathcal{M}}_\nu U_L^{\nu\dagger} \\ &\Rightarrow U_L^{\nu\dagger} \nu'_L = \nu_L, \quad U_L^{\ell\dagger} \ell'_L = \ell_L\end{aligned}$$

Charged current interaction:

$$\mathcal{L}_{\text{CC}} = -\frac{g}{\sqrt{2}} W_\mu^- \bar{\ell}'_L \gamma^\mu \nu'_L + \text{H.c.} = -\frac{g}{\sqrt{2}} W_\mu^- \bar{\ell}_L \gamma^\mu \underbrace{U_L^{\ell\dagger} U_L^\nu}_{U_{\text{PMNS}}} \nu_L + \text{H.c.}$$

Parameters of the lepton mixing matrix

U_{PMNS} can be parameterized by **three mixing angles**, the CP-violating phase δ and two physical Majorana phases.

| parameter | best fit | 2σ | HPS |
|---------------------|---------------------------|-----------------|---------------|
| $\sin^2\theta_{12}$ | $0.320^{+0.015}_{-0.017}$ | $0.29 - 0.35$ | $\frac{1}{3}$ |
| $\sin^2\theta_{23}$ | $0.49^{+0.08}_{-0.05}$ | $0.41 - 0.62$ | $\frac{1}{2}$ |
| | $0.53^{+0.05}_{-0.07}$ | $0.42 - 0.62$ | $\frac{1}{2}$ |
| $\sin^2\theta_{13}$ | $0.026^{+0.003}_{-0.004}$ | $0.019 - 0.033$ | 0 |
| | $0.027^{+0.003}_{-0.004}$ | $0.020 - 0.034$ | 0 |

M. Tórtola, J.W.F. Valle and D. Vanegas (2012)¹

¹arXiv: 1205.4018 (17 May 2012)

Symmetries in the lepton sector

For a long time U_{PMNS} seemed to be compatible with the **Harrison-Perkins-Scott** (HPS) mixing matrix

$$U_{\text{HPS}} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix}.$$

→ Idea: **symmetries in the lepton sector**

Recent data from oscillation experiments (T2K, MINOS, Double Chooz, Daya Bay, RENO)

$$\Rightarrow \theta_{13} > 0 \quad \text{at} > 6\sigma!$$

$$\text{Global fit}^2 : \quad \sin^2\theta_{13} = 0.026_{-0.004}^{+0.003} \quad (0.027_{-0.004}^{+0.003}).$$

Although $\theta_{13} > 0$, the idea of symmetries in the lepton sector is still interesting.

²M. Tórtola, J.W.F. Valle and D. Vanegas (2012), arXiv: 1205.4018

A simple way to restrict the lepton mass matrices (and therefore also U_{PMNS}) is to assume so-called **texture zeros** in the mass matrices:

$$\mathcal{M}_{ij} = 0 \quad \text{for some } i, j.$$

In the following we will assume:

- The charged lepton mass matrix \mathcal{M}_ℓ is diagonal.
- Neutrinos are Majorana particles ($\rightarrow \mathcal{M}_\nu$ symmetric).
- Texture zeros in the neutrino mass matrix \mathcal{M}_ν .

Within this framework: More than two texture zeros in $\mathcal{M}_\nu \Rightarrow$ incompatible with experimental data.

\Rightarrow (At most) two texture zeros.

P.H. Frampton, S.L. Glashow, D. Marfatia (2002)³:

In the basis where \mathcal{M}_ℓ is diagonal, there are **seven cases of two texture zeros** that are compatible with the experimental data.

$$A_1 = \begin{pmatrix} 0 & 0 & \times \\ 0 & \times & \times \\ \times & \times & \times \end{pmatrix}, \quad A_2 = \begin{pmatrix} 0 & \times & 0 \\ \times & \times & \times \\ 0 & \times & \times \end{pmatrix}, \quad B_1 = \begin{pmatrix} \times & \times & 0 \\ \times & 0 & \times \\ 0 & \times & \times \end{pmatrix},$$
$$B_2 = \begin{pmatrix} \times & 0 & \times \\ 0 & \times & \times \\ \times & \times & 0 \end{pmatrix}, \quad B_3 = \begin{pmatrix} \times & 0 & \times \\ 0 & 0 & \times \\ \times & \times & \times \end{pmatrix}, \quad B_4 = \begin{pmatrix} \times & \times & 0 \\ \times & \times & \times \\ 0 & \times & 0 \end{pmatrix},$$
$$C = \begin{pmatrix} \times & \times & \times \\ \times & 0 & \times \\ \times & \times & 0 \end{pmatrix}.$$

³Phys.Lett. B536 (2002) 79-82 [hep-ph/0201008]

C.I. Low (2005)⁴:

The only extremal mixing angle which can be enforced by texture zeros is $\theta_{13} = 0$.

⇒ Problems:

- $\theta_{13} = 0$ is ruled out by experiment.
- $\theta_{23} = 45^\circ$ cannot be enforced by texture zeros.

However, even if texture zeros do not restrict the mixing angles to certain values:

Texture zeros imply **relations among the observables** ($m_i, \Delta m_{ij}^2, \sin^2\theta_{ij}, \delta, \sigma_i$).

⁴Phys. Rev. D71 (2005) 073007 [hep-ph/0501251].

Near maximal atmospheric mixing from texture zeros

Texture zeros cannot enforce $\theta_{23} = 45^\circ$

→ Can $\theta_{23} \approx 45^\circ$ be achieved with an additional assumption?

W. Grimus, POL (2011)⁵:

- 1 Texture zeros of type B_3 or B_4
- 2 Quasi-degenerate neutrino mass spectrum: $m_i \gg \sqrt{|\Delta m_{ij}^2|}$

$$\Rightarrow \theta_{23} \approx 45^\circ.$$

Note: If one assumes a quasi-degenerate neutrino mass spectrum one has to take into account the [cosmological bounds](#): PDG (2010):

$$\sum_i m_i < 1 \text{ eV} \Rightarrow \text{quasi-degenerate spectrum not excluded yet.}$$

⁵Phys. Lett. B700 (2011) 356-361 [arXiv:1104.4340]

Near maximal atmospheric mixing from texture zeros

$$B_3 : \mathcal{M}_\nu \sim \begin{pmatrix} \times & 0 & \times \\ 0 & 0 & \times \\ \times & \times & \times \end{pmatrix}, \quad B_4 : \mathcal{M}_\nu \sim \begin{pmatrix} \times & \times & 0 \\ \times & \times & \times \\ 0 & \times & 0 \end{pmatrix}$$

\mathcal{M}_ℓ diagonal $\Rightarrow U_{\text{PMNS}} = U_L^\nu =: U$.

$$\mathcal{M}_\nu = U^* \text{diag}(m_1, m_2, m_3) U^\dagger$$

Two texture zeros in $\mathcal{M}_\nu \Rightarrow$ two equations of the form

$$(\mathcal{M}_\nu)_{ij} = \sum_k m_k U_{ik}^* U_{jk} = 0.$$

Input: $m_1, m_2, m_3, \theta_{12}, \theta_{13} \Rightarrow$ predictions: $\theta_{23}, \delta, \sigma_{1,2}$

\rightarrow One obtains a cubic equation for $\lambda = \tan^2 \theta_{23}$.

The limit of a quasi-degenerate spectrum

In the limit of a quasi-degenerate neutrino mass spectrum the cubic equation becomes

$$\lambda^3 + \lambda^2 - \lambda - 1 = 0$$

$$\Rightarrow \lambda = \tan^2 \theta_{23} = 1 \Rightarrow \sin^2 \theta_{23} = \frac{1}{2} \Rightarrow \theta_{23} = 45^\circ.$$

Limit independent of the values of θ_{12} and θ_{13} !

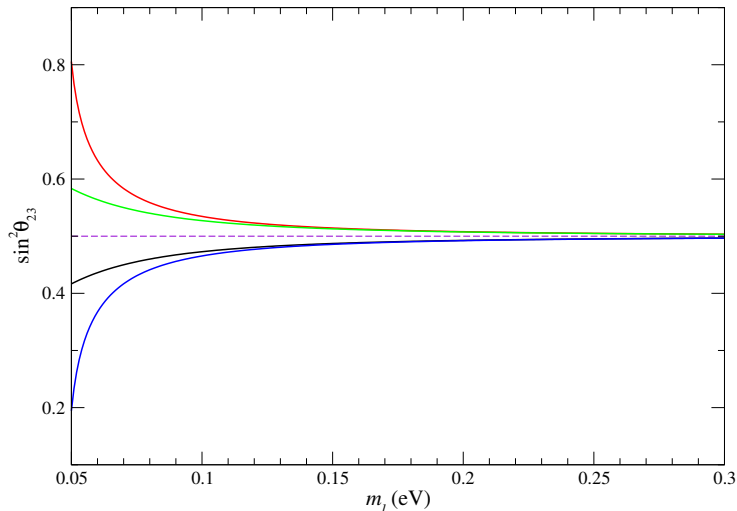
Approximate solution of the exact equation:

$$\sin^2 \theta_{23} \simeq \frac{1}{2} \mp \frac{1}{8} \frac{\Delta m_{31}^2}{m_1^2} (1 + \sin^2 \theta_{13})$$

The limit of a quasi-degenerate spectrum

(Δm_{21}^2 , Δm_{31}^2 , $\sin^2\theta_{12}$ and $\sin^2\theta_{13}$ fixed to best fit values.)

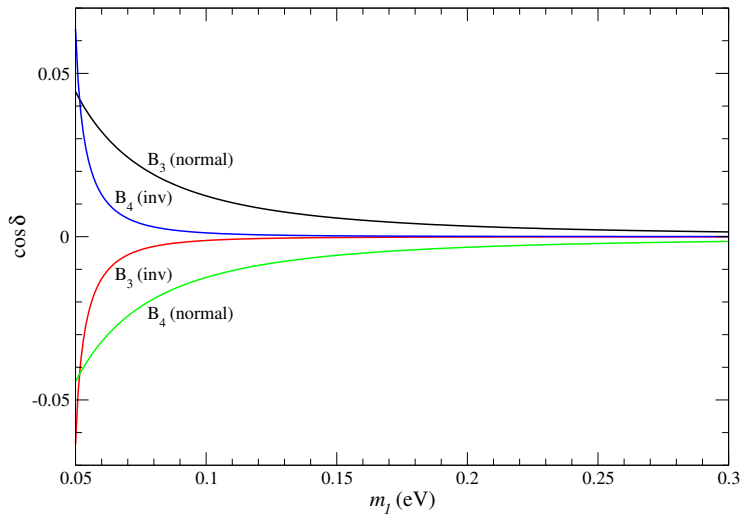
B_3 (inverted) B_4 (normal) B_3 (normal) B_4 (inverted)



The limit of a quasi-degenerate spectrum

In the limit of a quasi-degenerate mass spectrum:

$$\tan\theta_{12} \sin\theta_{13} \cos\delta \rightarrow 0.$$



Two texture zeros and the relation between θ_{13} and δ

- $\sin^2\theta_{13}$ well known by now.
- CP-phase δ almost not restricted at a significant level.

⇒ Interesting question: Correlation between θ_{13} and δ in models with two texture zeros?

POL, S. Morisi, E. Peinado (2011)⁶:

Two texture zeros A_1, \dots, C (\mathcal{M}_ℓ diagonal)

⇒ two equations of the form

$$(\mathcal{M}_\nu)_{ij} = \sum_k m_k U_{ik}^* U_{jk} = 0.$$

Leads (in the seven viable cases) to an at most cubic equation for $\cos \delta$ as a function of the ratio $\Delta m_{21}^2 / \Delta m_{31}^2$ and the three mixing angles.

⁶ Nucl. Phys. B857 (2012) 411-423 [arXiv:1109.3393]

Two texture zeros and the relation between θ_{13} and δ

\Rightarrow Vary $\Delta m_{21}^2/\Delta m_{31}^2$, $\sin^2\theta_{12}$ and $\sin^2\theta_{23}$ in their $n\sigma$ -ranges.

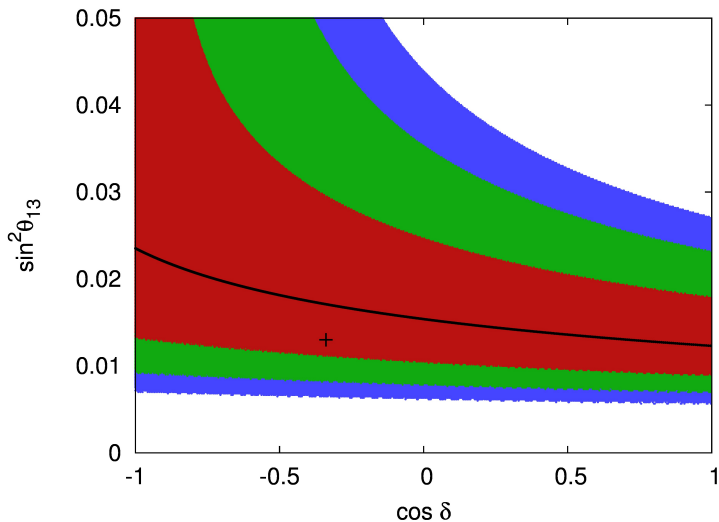
\Rightarrow One obtains plots of the allowed regions for $\sin^2\theta_{13}$ and $\cos\delta$.

Colors: black: best fit, red: 1σ , green: 2σ , blue: 3σ .

The best fit points⁷ for $\sin^2\theta_{13}$ and $\cos\delta$ are indicated by black crosses.

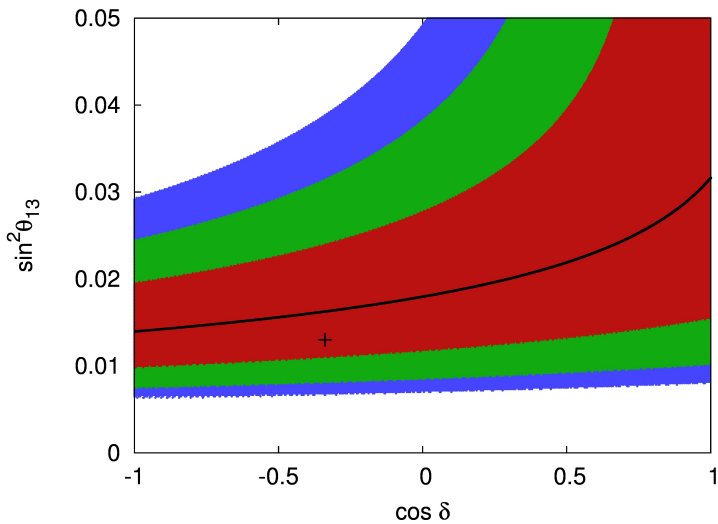
⁷T. Schwetz, M. Tórtola and J.W.F. Valle, New J. Phys. 13 (2011) 109401 [arXiv:1108.1376].

A_1 : only normal spectrum allowed:



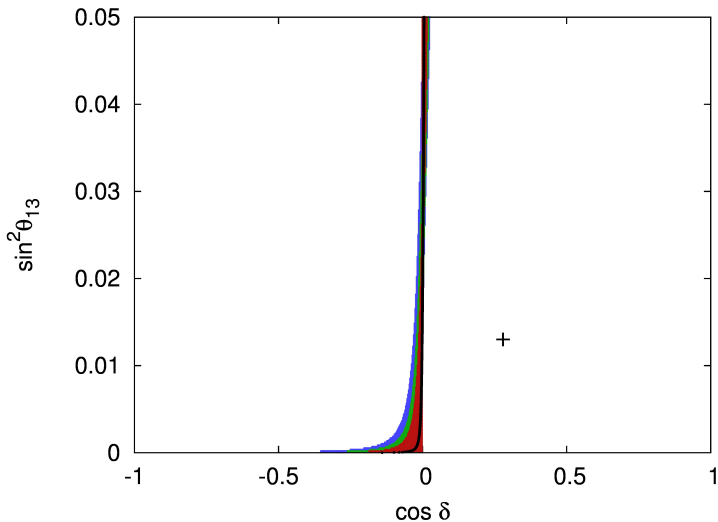
“Predicts” $\sin^2 \theta_{13} > 0$.

Similar for A_2 :



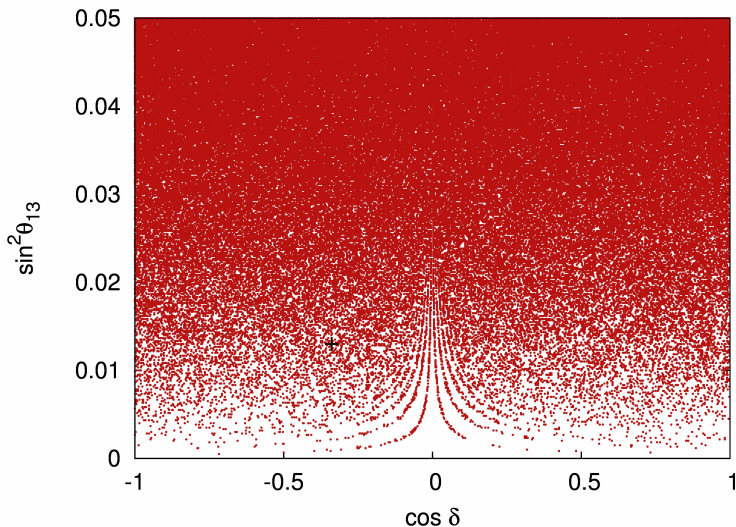
“Predicts” $\sin^2 \theta_{13} > 0$.

Cases B_1 to B_4 all predict $\cos \delta \approx 0$.

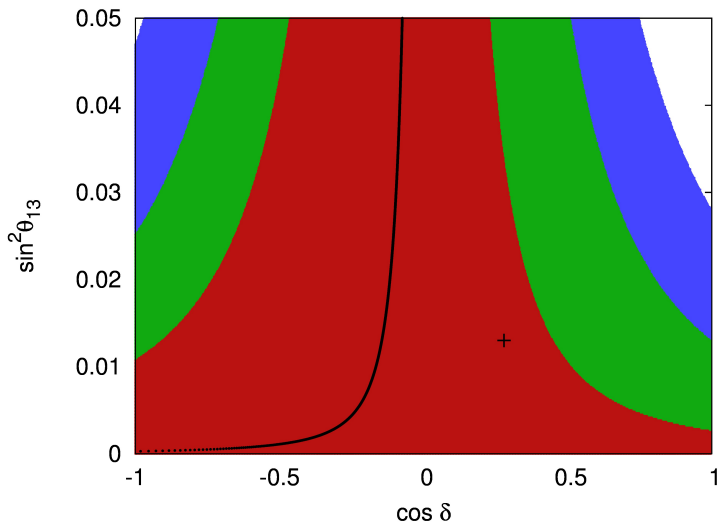


As an example: B_1 (inverted spectrum)

Case C (normal): No correlation between $\sin^2\theta_{13}$ and $\cos\delta$.



Case C (inverted):



Conclusions

- We have found two instances in the framework of **two texture zeros** in \mathcal{M}_ν (\mathcal{M}_ℓ diagonal) which lead to **maximal atmospheric neutrino mixing** in the limit of a **quasi-degenerate** neutrino mass spectrum.
- In this scenario the limit

$$\sin^2\theta_{23} \rightarrow \frac{1}{2}$$

is **independent** of the values of θ_{13} and θ_{12} !

- $\theta_{23} \approx 45^\circ$ and $\cos\delta \approx 0$ achievable by the use of an **Abelian symmetry** (texture zeros). \rightarrow Alternative road to explain near maximal atmospheric neutrino mixing.
- Seven cases of two texture zeros in \mathcal{M}_ν are compatible with the experimental data.
- Two of these cases are compatible with the experimental data on $\sin^2\theta_{12}$, $\sin^2\theta_{23}$ and Δm_{ij}^2 **only if $\sin^2\theta_{13} > 0$.**