## Glueball Searches with the STAR Detector at RHIC

## Włodek Guryn BNL (for the STAR Collaboration)

- Process of central production - physics program at STAR with forward protons
- The present: 2009 run preliminary results of Phase I program
- The near future: Phase I program continued
- The future: Phase II program, large data samples for Central Production
- Ultra peripheral AuAu collisions at STAR
- Other QCD topics of interest in Central Exclusive Production (Odderon, $\chi_{\mathrm{c}}$, glueball production and QCD scale invariance...)
- Summary


## Physics Processes I



In t-channel it is an exchange with quantum numbers of vacuum


Non Pert. QCD


PQCD picture

## Processes with Tagged Forward Protons

QCD color singlet exchange: $\mathrm{C}=+1$ (IP), $\mathrm{C}=-1(\mathrm{O})$
Discovery Physics

$p+p \rightarrow p+p$ elastic


$$
p+p \rightarrow p+X+p
$$

diffractive $X=$ particles, glueballs


## This Topic Has a Long History

No 1

CENTRAL DIFFRACTIVE PRODUCTION
By G. Biafkowski and J. Kalinowski
Institute of Theoretical Physics, Warsaw University*
(Received April 1, 1974; Revised version received June 1, 1974)
The topological cross sections and some characteristics of the multiplicity distribution for central diffractive production via double pomeron exchange are discussed.


Fig. 1 a) Single diffractive excitation, b) double diffractive excitation,c) central diffractive productic diagrams

## Summary of the Existing Data



Włodek Guryn
STAR

## Central Exclusive Production in Double Pomeron Exchange (DPE)



Method is complementary to:

- GLUEX experiment (2015)
- PANDA experiment (>2015)
- COMPASS experiment (taking data)
- BESIII

In the Double Pomeron Exchange (DPE) process each proton "emits" a Pomeron and the two Pomerons interact producing a massive system $\mathrm{M}_{\mathrm{X}}$
$\mathbf{M}_{\mathbf{X}}=\sqrt{\xi_{1} \xi_{2}}$ invariant mass


For each proton vertex one has
t four-momentum transfer
$\xi=\Delta \mathrm{p} / \mathrm{p}$

$$
\text { where } M_{X}=\pi^{+} \pi^{-}, \chi_{c}\left(\chi_{b}\right), q q(j e t s), H(\text { Higgs boson }), ~ g g(g l u e b a l l s)
$$

The massive system could form resonances. We expect that because of the constraints provided by the double Pomeron interaction, glueballs, hybrids, and other states coupling preferentially to gluons, will be produced with much reduced backgrounds compared to standard hadronic production processes.

## Glueball Spectrum

## Sparse spectrum!

New I=0 mesons starting with
$\mathbf{0}^{++} \quad$ 1.6 GeV
$\mathbf{0}^{-+}, 2^{++} \quad 2.3-2.5 \mathrm{GeV}$
No JPC -exotic glueballs until
$2^{+-}$at 4 GeV


## The Relativistic Heavy Ion Collider



RHIC is a QCD Laboratory:
Nucleus- Nucleus collisions (AuAu, CuCu...); Asym. Nucl. (dAu);
Polarized proton-proton; eRHIC - Future

## RHIC: the world's first polarized pp collider



- Spin orientation varies bunch by bunch
- Spin pattern changes from fill to fill
- Spin rotators provide choice of spin orientation
- "Billions" of spin reversals during a fill with little if any depolarization


## Implementation at RHIC - tag forward protons PP2PP Setup

Phys. Lett. B 579 (2004) 245-250, Phys. Lett. B 632 (2006) 167-172, Phys. Lett. B 647 (2007) 98-103 (Polish coauthors Chwastowski, Pawlik, Sandacz)


## Principle of the Measurement



- (Elastically) scattered protons have very small scattering angle $\theta^{*}$, hence beam transport magnets determine trajectory scattered protons
- The optimal position for the detectors is where scattered protons are well separated from beam protons
- Need Roman Pot to measure scattered protons close to the beam without breaking accelerator vacuum
Beam transport equations relate measured position at the detector to scattering angle

$$
\left(\begin{array}{l}
x_{D} \\
\Theta_{D}^{x} \\
y_{D} \\
\Theta_{D}^{y}
\end{array}\right)=\left(\begin{array}{llll}
a_{11} & L_{e f f}^{x} & a_{13} & a_{14} \\
a_{21} & a_{22} & a_{23} & a_{24} \\
a_{31} & a_{32} & a_{33} & L_{e f f}^{y} \\
a_{41} & a_{42} & a_{43} & a_{44}
\end{array}\right)\left(\begin{array}{l}
x_{0} \\
\Theta_{x}^{*} \\
y_{0} \\
\Theta_{y}^{*}
\end{array}\right) \quad \begin{aligned}
& \mathrm{x}_{0}, \mathrm{y}_{0}: \text { Position at Interaction Point } \\
& \Theta^{*} \Theta^{*} \Theta_{\mathrm{y}}^{*}: \text { Scattering Angle at IP } \\
& \mathrm{x}_{\mathrm{D}}, \mathrm{y}_{\mathrm{D}}: \text { Position at Detector } \\
& \Theta_{\mathrm{D}}, \Theta_{\mathrm{D}} \text { : Angle at Detector }
\end{aligned}
$$

## Reconstruction of the Momentum Loss $\xi$

For inelastic forward protons there is an additional momentum loss term

1. Need to measure vector at the detection point, hence two RPs are needed on each side of STAR.
2. For a proton, which scatters with $\Theta$ and $\xi$ we have:

$$
\begin{aligned}
& x_{1}=a_{1} x_{0}+L_{1} \Theta_{x}+\eta_{1} \xi ; \quad \text { detection point } 1 \\
& x_{2}=a_{2} x_{0}+L_{2} \Theta_{x}+\eta_{2} \xi ; \quad \text { detection point } 2 \\
& \binom{\Theta_{x}}{\xi}=\frac{1}{D e t}\left(\begin{array}{c}
\eta_{2} ; \\
-\eta_{1} ; \\
-L_{1} \\
\xi
\end{array}\right)\binom{x_{1}-a_{1} x_{0}}{x_{2}-a_{2} x_{0}}
\end{aligned}
$$

## Current STAR detector in cross section



Large acceptance detector running since 2000

- High resolution tracking device: TPC in $-1<\eta<1,-\pi<\varphi<\pi$
- Forward rapidity gap veto
- FTPC: $2.5<|\eta|<4.2, B B C: 3.8<|\eta|<5.2$


## Great Charged Particle ID in the STAR TPC

- High resolution tracking device: TPC in $-1<\eta<1,-\pi<\varphi<\pi$
- Excellent particle identification capability: TPC dE/dx, ToF
Particle Identification at STAR


Reconstructed hadrons: $K_{s}, \boldsymbol{\phi}, \Lambda, \equiv$, and $\Omega$ in Au+Au collisions at $\sqrt{ } \mathbf{s}_{\mathrm{NN}}=39 \mathrm{GeV}$

dE/dx vs. rigidity compared with theoretical expectations


Particle identification with new barrel Time-of-Flight system.

## Implementation: STAR + pp2pp

1. Need detectors to measure forward protons: $\mathbf{t}$ - four-momentum transfer, $\xi=\Delta \mathrm{p} / \mathrm{p}, \mathrm{M}_{\mathrm{X}}$ invariant mass and;
2. Detector with good acceptance and particle ID to measure central system


- Phase I, present- low-t coverage
- Phase II, future- higher-t coverage, large data samples


## Kinematic "filter" ( $\mathrm{dp}_{\mathrm{T}}$ ) for " gg "

(F. Close et al./W102)

PLB 397339 (1997)

- Coupling of the exchange particles to the final state mesons for gluon exchange (small $\mathrm{dp}_{\mathrm{T}}$ ) and quark exchange (large $\mathrm{dp}_{\mathrm{T}}$ )
- Spin-dependence of the coupling can be studied at RHIC

As predicted by Regge theory the diffractive cross section at RHIC is dominated by the Pomeron (gluonic) exchange, :

$$
\begin{aligned}
& \sigma_{R R} \sim \mathrm{~s}^{-2} \\
& \sigma_{R P} \sim \mathrm{~s}^{-1} \\
& \sigma_{\mathrm{PP}} \sim \text { const. or } \mathrm{s}^{\alpha} \text { where } \alpha \sim(0.1)
\end{aligned}
$$



## WA102 $\mathrm{f}_{0}(1500) \pi^{+} \pi^{-} \pi^{+} \pi^{-}$

$\sigma\left(f_{1}\right)=7 \mu$ barn

We are sensitive to this level of cross section
$\sigma\left(\mathrm{f}_{0}\right)=3 \mu \mathrm{barn}$

| $\mathrm{f}_{1}(1285)$ |  |
| :---: | :---: |
|  | F.E.Close and A.Kirk, PLB397, 333 (1997). $\mathrm{dP}_{\mathrm{T}}>0.5 \mathrm{GeV} / \mathrm{c}$ |
|  | $\begin{aligned} & 0.2 \mathrm{GeV}<\mathrm{dP}_{\mathrm{T}}<0.5 \mathrm{GeV} / \mathrm{c} \\ & -\mathrm{f}_{0}(1500) \end{aligned}$ |
|  | $\mathrm{dP}_{\mathrm{T}}<0.2 \mathrm{GeV} / \mathrm{c}$ |

Figure 3: The $4 \pi$ mass spectra (i) With $d P_{T}>0.5 \mathrm{GeV}$ exhibiting a clear $f_{1}(1285)$; (ii) $0.2<d P_{T}<0.5 \mathrm{GeV}$ (iii) $d P_{T}<0.2 \mathrm{GeV}$ where the $f_{1}(1285)$ has disappeared while the $f_{0}(1500)$ is seen more clearly.

## Central Exclusive Production Process in DPE

$$
p_{1} p_{2} \rightarrow p_{1} \cdot M_{x} p_{2^{\prime}}
$$

- Exclusive process with "small" momentum transfer:

$$
-\mathrm{t}_{1}\left(\mathrm{p}_{1} \rightarrow \mathrm{p}_{1}^{\prime}\right) \text { and }-\mathrm{t}_{2}\left(\mathrm{p}_{2} \rightarrow \mathrm{p}_{2}^{\prime}\right)
$$

- $M_{X}$ is centrally produced, nearly at rest, through DPE process
- In pQCD, Pomeron is considered to be made of two gluons: natural place to look for gluon bound state
- $\mathrm{M}_{\mathrm{x}}\left(\sim 1-3 \mathrm{GeV} / \mathrm{c}^{2}\right) \rightarrow \pi^{+} \pi^{-}, \pi^{+} \Pi^{-} \pi^{+} \pi^{-}, \mathrm{K}^{+} \mathrm{K}^{-}, \ldots$
- Lattice calculation: Lightest glueball $\mathrm{M}\left(0^{++}\right)=1.5-1.7 \mathrm{GeV} / \mathrm{c}^{2}$ (PRD73 2006)
- Search for glueball (gg) candidates in $M_{x}$
- Candidates with conventional quantum numbers: need to be studied in a wide kinematical range


## Phase I: 2009 data



- Data taken with RP and ToF multiplicity triggers for the central process
- Tracks reconstructed off line in TPC
- Require two reconstructed tracks in opposite direction in the RPs
- Work in progress for identifying exclusive DPE events: rapidity gaps, PID, pтbalance, missing-mass...


## Run 2009 Candidate Central Production Event

Event Information run: 10183036 Events seen: 25 Event \#127


Triggers:

## Phase I: First Look at Central Production

(Vector Meson Production)
Exclusive production of vectors meson in pp and pbarp collisions
A.Cisek, W. Schäfer and A. Szczurek, Meson 2010


The $\rho^{0}$ cross section $\sim$ few hundred nb



## Phase I: First Look at Central Production

(non exclusive channels)


- Use RP and ToF multiplicity trigger online
- Reconstruct TPC tracks off line
- Confirm two tracks in the Roman Pots
- Reconstruct $\mathrm{M}_{\mathrm{x}}$ using TPC tracks assuming all pions




## Phase I Elastic Scattering: First high-statistics measurement of $A_{N}$ at $\mathrm{HE}(\sqrt{ } \mathrm{s}=200 \mathrm{GeV})$

$$
A_{N}(t)=\frac{\sigma^{\wedge}(t)-\sigma^{\downarrow}(t)}{\sigma^{\wedge}(t)+\sigma^{\downarrow}(t)}=C_{1} \phi_{f l i p}^{e n^{*}} \phi_{\text {non-flip }}^{\text {had }}+C_{2} \phi_{f i p}^{\text {had }} \phi_{\text {non-flip }}^{e m}
$$



- Statistical errors + systematic $t$-scale uncertainty (10\%) in the fit
- Higher- $t$ reach planned from the upcoming $\sqrt{ } s=500 \mathrm{GeV}$ (and with Phase II set-up) at RHIC


## Phase II - Simulation Performance Plots

We assume the DPE cross section $140 \mu$ barn, and branching ratios as measured at the ISR
A. Breakstone et al., Z. Phys. C42, (1989) 387

| Reaction | Number <br> of events | Cross section <br> $[\mu \mathrm{b}]$ |
| :--- | :---: | :---: |
| (1) $p p \rightarrow p p\left(\pi^{+} \pi^{-}\right)$ | 16400 | $79.0 \pm 13.0$ |
| (2) $p p \rightarrow p p\left(2 \pi^{+} 2 \pi^{-}\right)$ | 5800 | $46.0 \pm 10.0$ |
| (3) $p p \rightarrow p p\left(3 \pi^{+} 3 \pi^{-}\right)$ | 1900 | $32.0 \pm 9.0$ |
| (4) $p p \rightarrow p p\left(K^{+} K^{-}\right)$ | 560 | $6.5 \pm 1.7$ |
| (5) $p p \rightarrow p p\left(K^{+} K^{-} \pi^{+} \pi^{-}\right)$ | 150 | $10.0 \pm 3.3$ |
| (6) $p p \rightarrow p p(p \bar{p})$ | 120 | $0.8 \pm 0.17$ |
| (7) $p p \rightarrow p p\left(p \bar{p} \pi^{+} \pi^{-}\right)$ | 65 | $1.3 \pm 0.36$ |

## Phase II - Simulation Performance Plots

- Mass $M_{X}$ calculated from the proton kinematics
- Use phase space to determine the decay of mass $M_{x}$ in a particular channel
- Use STAR TPC acceptance to make sure that all decay products are measured.
- High- $\mathrm{M}_{\mathrm{x}}$ reconstruction is limited by PID ( $\pi / \mathrm{K}$ separation up to $\sim 1.6 \mathrm{GeV} / \mathrm{c}$ )




## Acceptance and expected yields in $\mathrm{M}_{\mathrm{x}}$



Mass Acceptance


Event yields for 20 week run at 500 GeV

$$
\begin{array}{lc}
\pi^{+} \pi^{-} \pi^{+} \pi^{-}-2.7 \times 10^{6} \text { events } \\
\pi^{+} \pi^{-}- & 10.4 \times 10^{6} \text { events } \\
K^{+} K^{-}- & 0.8 \times 10^{6} \text { events }
\end{array}
$$

## Other QCD Processes: $\chi_{c}$ Production

 arXiv:1005.0695 L.A. Harland-Lang, V.A. Khoze, M.G. Ryskin, W.J. Stirling

RHIC

| $\sqrt{s}(\mathrm{TeV})$ | 0.5 | 1.96 | 7 | 10 | 14 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\frac{\mathrm{~d} \sigma}{\mathrm{~d} y_{\chi_{c}}(p p \rightarrow p p(J / \psi+\gamma))}$ | 0.57 | 0.73 | 0.89 | 0.94 | 1.0 |
| $\frac{\mathrm{~d} \sigma\left(1^{+}\right)}{\mathrm{d} \sigma\left(+^{+}\right)}$ | 0.59 | 0.61 | 0.69 | 0.69 | 0.71 |
| $\frac{\mathrm{~d} \sigma\left(2^{+}\right)}{\mathrm{d} \sigma\left(0^{+}\right)}$ | 0.21 | 0.22 | 0.23 | 0.23 | 0.23 |

Table 3: Differential cross section (in nb) at rapidity $y_{\chi}=0$ for central exclusive $\chi_{c J}$ production via the $\chi_{c J} \rightarrow J / \psi \gamma$ decay chain, summed over the $J=0,1,2$ contributions, at RHIC, Tevatron and LHC energies, and calculated using GRV94HO partons, as explained in the text.

$$
\operatorname{BR}\left(\chi_{c}=>J / \psi+\gamma\right)=1.14 \pm 0.08 \%
$$

CDF: $\sigma\left(\chi_{\mathrm{c}}\right)=76 \pm 10 \pm 10 \mathrm{nb}$ PRL 102, 242001 (2009)

## Other QCD Processes: $\chi_{c}$ Production

## Szczurek and collaborators

1. $\chi_{c}(0+)$ production:
R. S. Pasechnik, A. Szczurek and O. V. Teryaev, Phys. Rev. D 78, 014007 (2008) [arXiv:0709.0857 [hep-ph]];
2. $\chi_{c}(1+)$ production:
R. S. Pasechnik, A. Szczurek and O. V. Teryaev, Phys. Lett. B 680, 62 (2009) [arXiv:0901.4187 [hep-ph]];
3. $\chi_{c}(2+)$ production:
R. S. Pasechnik, A. Szczurek and O. V. Teryaev, Phys. Rev. D 81, 034024 (2010) [arXiv:0912.4251 [hep-ph]].

TABLE I: Integrated over full phase space cross sections (in nb) for the central exclusive $\chi_{c}\left(0^{+}, 1^{+}, 2^{+}\right)$ production at RHIC energy $W=200 \mathrm{GeV}$. Absorption effects and NLO QCD corrections to the $g g \rightarrow \chi_{c}$ vertex are included here. Gap survival factor for all $\chi_{c}$ states is taken here to be equal 0.1 . Branching ratio to the channel of interest should be included in addition.

| $\chi_{c}$ | without absorption | with absorption |
| :---: | :---: | :---: |
| $\chi_{c}\left(0^{+}\right)$ | 45 | 4.5 |
| $\chi_{c}\left(1^{+}\right)$ | 2 | 0.2 |
| $\chi_{c}\left(2^{+}\right)$ | 2 | 0.2 |

## Odderon Contribution to $\mathrm{J} / \psi$ production

A. Bzdak et. al Phys. Rev.D 75.094023, also S. Klein, et. al PRL 92, 142003 (2004) for $\gamma \mathrm{IP}$
$E\left(s, m_{V}\right)=\left(x_{0} \sqrt{s} / m_{V}\right)^{2 \lambda}$.

| $\sigma^{\text {corr }} / d y$ | $J / \psi$ |  | $\Upsilon$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | odderon | photon | odderon | photon |
| Tevatron | $0.3-1.3-5 \mathrm{nb}$ | $0.8-5-9 \mathrm{nb}$ | $0.7-4-15 \mathrm{pb}$ | $0.8-5-9 \mathrm{pb}$ |
| LHC | $0.3-0.9-4 \mathrm{nb}$ | $2.4-15-27 \mathrm{nb}$ | $1.7-5-21 \mathrm{pb}$ | $5-31-55 \mathrm{pb}$ |

Use scaling from HERA to calculate cross sections from Tevatron to RHIC

$$
\begin{gathered}
\mathrm{J} / \psi(500 / 2000)^{0.4} \sim 0.4 \text { and for } \mathrm{Y}(500 / 2000)^{0.7 \sim 5} \\
\mathrm{~d} \sigma^{\text {corr } / \mathrm{dy}}=>\mathrm{J} / \psi \sim 0.5 \mathrm{nb} \text { and } \mathrm{Y} \sim 0.8 \mathrm{pb}
\end{gathered}
$$

## The Glueball Filter in Central Production and Broken Scale Invariance

## John Ellis and Dmitri Kharzeev arXiv:hep-ph/9811222



Abstract
We propose a possible explanation of the kinematical dependence of the central production of the scalar glueball candidate observed recently by the WA91 and WA102 Collaborations, and discussed by Close and Kirk, in the context of the broken scale invariance of QCD. The dependences of glueball production on the transverse momenta and azimuthal angles of the final-state protons may be related to the structure of the trace anomaly in QCD.

## The $\phi$ Correlations and Broken Scale Invariance

John Ellis and Dmitri Kharzeev arXiv:hep-ph/9811222
The azimuthal angle correlation between outgoing protons are an alternative filter to identify glueball states

$\phi$
Figure 2: Distribution in the azimuthal angle $\phi$ (in radians) between the final-state protons in double-diffractive production of the $f_{0}(980)$ (left panel) and $f_{0}(1500)$ (right panel) states, calculated for the center-of-mass energy of the WA91 and WA102 experiments, and with $t=-0.5 \mathrm{GeV}^{2}$ for both final-state protons.

## $\mathrm{J} / \psi$ Production in DPE and Photoproduction

1. In many cases $J / \psi$ is produced.
2. The experimental challenge is to distinguish between production channels

Are $\phi$ correlations and $\mathrm{dP}_{\mathrm{T}}$ distributions of the forward protons different for different reactions?

Need $d^{2} \sigma / d t d \xi$ for each process

## Ultra Peripheral Collisions (UPC) of Heavy Ions



Phys.Rept.458:1-171,2008 (arXiv:0706.3356)

- Ultra Peripheral Collisions: nuclei miss each other ( $\gamma \gamma$ and $\gamma \mathrm{P}$ interactions)
- Requires: $\boldsymbol{b}>\boldsymbol{R}_{\boldsymbol{A}}+\boldsymbol{R}_{\boldsymbol{B}}$
-Weizsacker-Williams approach: a field of almost real photons


## Particle Production in UPC



- Meson spectroscopy: $\rho, \omega, \varphi, \rho^{\prime}=$ mix of $\rho(1450)$ and $\rho(1700)$;
- Transition from soft physics $(\rho, \omega, \varphi)$ to pQCD $(J / \Psi, Y)$;
- $\mathrm{e}^{+} \mathrm{e}^{-}$pair production;
- Fundamental tests of Quantum Mechanics: interference between non overlapping particles


## Signature and Triggering

- Signatures:
- Coherent production dominates
- Low transverse momentum ( $\mathrm{p}_{\mathrm{T}} \leq 2 \mathrm{~h} / \mathrm{RA} \approx 60$ MeV )
- Low multiplicity events with vertex
- Events with nuclear breakup accompanied by forward neutrons
- Trigger:
- Minimum bias
- Low multiplicity
- Neutrons in both ZDCs
- Topology ToF and TPC:
- Low multiplicity events
- Coincidence of North and South
- Top and Bottom veto cosmics



## The $\rho^{0}$ photoproduction cross section

- Goncalves \& Machado (EPJ C29,2003)
- QCD color dipole approach
- Nuclear effects and parton saturation phenomena
- Frankfurt, Strikman \& Zhalov (PRC67 034901,2003)
- Generalized vector dominance (VDM)
- QCD - Gribov-Glauber approach
- Klein \& Nystrand (PR C60 014903, 1999)
- VDM
- Classical mechanical approach for scattering

Klein \& Nystrand model agrees well with the data, $|y|<1$ does not allow to discriminate based on shape



50th Cracow School
Zakopane 2010

Red total cross section Blue cross section with mutual excitation.
Simulation based on Klein \& Nystrand

Coherent and incoherent form factors Double exponential fit function $\sigma($ incoh $) / \sigma($ coh $) \sim 0.29 \pm 0.03$


Włodek Guryn

## Photoproduction of $\pi^{+} \pi^{-} \pi^{+} \pi^{-}$

(arXiv:0912.0604) Phys.Rev.C81:044901,2010

- Expected to be largely through a radially excited
- Could be $\rho(1450)$ and/or $\rho(1700)$
- Studies of the substructure showed low mass pion pairs accompanied by $\rho(770)$
$\sigma_{\text {coh }}\left(\pi^{+} \pi^{-} \pi^{+} \pi^{-}\right) / \sigma_{\text {coh }}(\rho[770])=13.4 \pm 0.8 \%$



Peak at low $\mathrm{p}_{\mathrm{T}}$ is due to the coherent production


STAR is $^{37}$

## Summary

1. A new rich diffractive physics program with tagged forward protons in polarized proton-proton scattering at RHIC, has been launched and its significant expansion has been proposed.
2. It will search for new physics, including glueballs, Odderon and sphalerons (particle correlations in DPE process).
3. It will search for diffractive production of light and massive systems in double Pomeron exchange process. Possible Pomeron - Odderon interaction $=>\mathrm{J} / \psi$ production, C -odd glueball.
4. Other QCD processes will be studied a) breaking of the scale invariance as a glueball filter; and b) $\chi_{c}$ central production.
5. Not discussed here - systematic study of the spin dependence of elastic scattering, of unpolarized quantities $\sigma_{\text {tot }}, \rho, d \sigma / d t$ in unexplored ranges of $t$ and $\sqrt{ } s$.
6. STAR has had a reach program with many common topics in heavy ion UPC collisions.
RHIC is an exciting, and complementary to other hadron colliders, place to do diffractive physics
