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Finite size giant
magnons and
interactions

Olof Ohlsson Sax

Giant magnons

Finite size
corrections

Magnons from
finite gap

Magnons from
sine-Gordon

Summary

Finite size giant magnons and interactions

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Zakopane, June 20, 2008

J. A. Minahan and O. Ohlsson Sax, Nucl. Phys. B **801**, 97 (2008)

arXiv:0801.2064 [hep-th].



Giant magnons

- Consider a string moving in $\mathbb{R} \times S^2 \subset \text{AdS}_5 \times S^5$.

Hofman–Maldacena limit

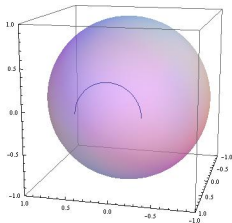
$$\begin{array}{ll} E, J \rightarrow \infty & p = \text{fixed} \\ E - J = \text{fixed} & \lambda = \text{fixed} \end{array}$$

- We can now relax the level matching condition $\sum p \in 2\pi\mathbb{Z}$ and consider single magnon excitations.



Giant magnons

- Use conformal gauge: $\frac{dJ}{dx} = \text{const} \rightarrow$ the worldsheet becomes infinitely long.
- Fundamental excitations are local on worldsheet, but macroscopic in space-time.
- Dispersion relation: $E - J = \sqrt{1 + \frac{\lambda}{\pi^2} \sin^2 \frac{p}{2}} \approx \frac{\sqrt{\lambda}}{\pi} \sin \frac{p}{2}$
- Compare with spin-chain magnon.





Finite size corrections

- The leading order corrections for finite J to the single giant magnon dispersion relation were calculated by Arutyunov, Frolov and Zamaklar, as a soliton solution in uniform gauge.

$$\Delta(E - J) = -\frac{4}{e^2} \frac{\sqrt{\lambda}}{\pi} \sin^3 \frac{p}{2} e^{-\mathcal{R}},$$

$$\mathcal{R} = 2 \frac{J}{E - J} + ap \cot \frac{p}{2}$$

- The result is dependent on the gauge parameter a .



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- The result is dependent on the gauge parameter a .
- We can get a gauge independent result (corresponding to $a = 0$) by considering a giant magnon with momentum $p = 2\pi m/M$ on an orbifold $\mathbb{R} \times S^2/\mathbb{Z}_M$. [Astolfi, Forini, Grignani and Semenoff]



Multi-magnon states

For $J \rightarrow \infty$ we can construct multi-magnon states in two simple ways:

- 1 Put them on separate patches, separated by an infinite worldsheet.
- 2 Put several magnons on the same patch, in a scattering state.

As $t \rightarrow \infty$ the interacting magnons become infinitely separated. The two types of multi-magnon states thus share the same energy spectrum.



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For finite J there is always a finite separation between the magnons, so any multi-magnon state is interacting.



Finite gap equations

- Encodes the spectrum of a classical string solution in terms of a differential on a Riemann surface.
- The differential is specified as a density along the square root **cuts** and log cuts (**condensates**).
- The density is constant on the condensates, but is given as solutions of an integral equation on the cuts:

$$\oint \frac{\rho(x)dx}{z-x} = f_i(E, J) \quad z \in \mathcal{C}_i$$

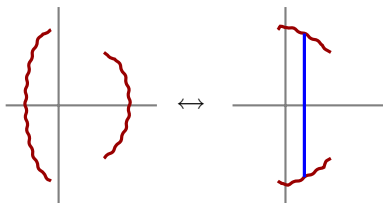


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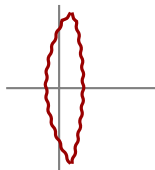
- We can re-cut the surface by reconnecting the branch points in another way, but we may need to add some extra condensates.





Finite gap equations

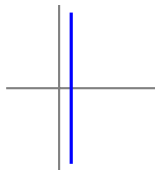
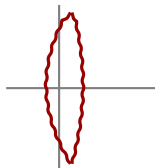
- A giant magnon is described by a two-cut solution in the singular limit where the end points of the cuts merge. [Vicedo]





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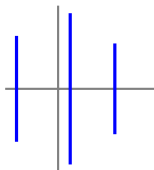
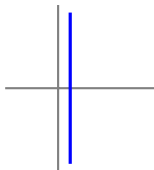
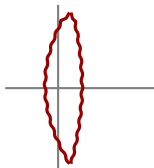
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- It can also be described by a single condensate. [Minahan, Tirziu, Tseytlin]





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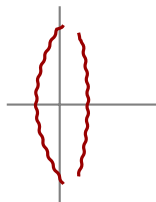
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- It can also be described by a single condensate. [Minahan, Tirziu, Tseytlin]
- A multi-magnon state is given by a set of condensates.





Finite size corrections from finite gap equations

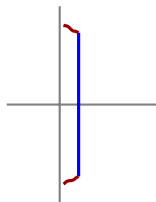
- A finite size giant magnon is described by a two cut solution for which the endpoints *almost* coincide.





Finite size corrections from finite gap equations

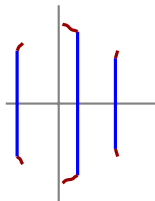
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- This is equivalent to a condensate with small cuts attached to the ends.





Finite size corrections from finite gap equations

- A finite size giant magnon is described by a two cut solution for which the endpoints *almost* coincide.
- This is equivalent to a condensate with small cuts attached to the ends.
- A multi-magnon state is given by a number of such condensate – cut configurations.





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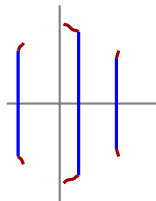
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Finite size corrections from finite gap equations

- A N magnon configuration is described by a genus $2N - 1$ Riemann surface.





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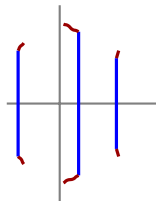
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- In general we would expect hyper-elliptic solutions.



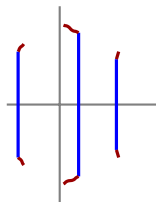


Finite size corrections from finite gap equations

- A N magnon configuration is described by a genus $2N - 1$ Riemann surface.
- In general we would expect hyper-elliptic solutions.
- The density for the cut \mathcal{C}_i is governed by

$$\int_{\mathcal{C}_i} \frac{\rho(x)}{z - x} = - \int_{\text{other}} \frac{\rho(x)}{z - x} + \dots,$$

where $z \in \mathcal{C}_i$ and the integral on the right is over all other cuts and condensates.





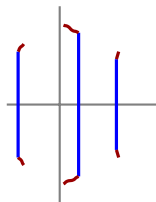
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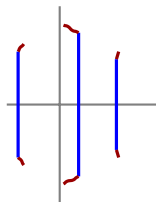
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where $z \in \mathcal{C}_i$ and the integral on the right is over all other cuts and condensates.

- Since the cuts are small, we can treat all integrands except for the closest condensate as constant in z .
- The problem is reduced to a one-cut problem which gives an algebraic solution.



Finite size corrections from finite gap equations

- For M magnons with momenta p_i , energy E_i

$$\Delta(E_i - J) = -\frac{4}{e^2} \frac{\sqrt{\lambda}}{\pi} \sin^3 \frac{p_i}{2} e^{-2 \frac{J}{E_i - J}} \\ \times \prod_{k \neq i}^m \frac{\sin^2 \frac{p_i + p_k}{4}}{\sin^2 \frac{p_i - p_k}{4}} e^{-2 \frac{E_k - J}{E_i - J}} .$$

This is singular as $p_i \rightarrow p_j$.



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This is singular as $p_i \rightarrow p_j$.

- For M magnons with $p = 2\pi m/M$

$$\Delta(E - J) = -\frac{4}{e^2} \frac{\sqrt{\lambda}}{\pi} \sin^3 \frac{p}{2} e^{-2 \frac{J/M}{E - J}}.$$

This agrees exactly with the result for one magnon on a \mathbb{Z}_M orbifold.



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Pohlmeyer reduction

- The classical equations of motion for the $O(3)$ sigma model – i.e. for the gauge fixed string – together with the Virasoro constraints are equivalent to the sine-Gordon equation.
- The giant magnon solution corresponds to the fundamental solitonic solution – the kink.



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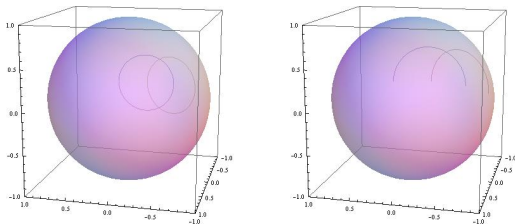
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Multi-magnon states from sine-Gordon

Interacting two-magnon states can be constructed from kink–kink and kink–anti-kink scattering solutions to the sine-Gordon equation.



Finite size solutions can be constructed by considering periodic generalizations of these solutions. [Klose and McLoughlin]



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Summary

- For infinite J , multi-magnon state can be either non-interacting or interacting.
- For finite J , multi-magnon states are always interacting.
- Explicit finite size interacting two-magnon states have been constructed from two-phase solutions to the sine-Gordon equation.
- The leading order finite size corrections for interacting states with *any* number of magnons have been computed using the finite gap equations.
- The finite gap computation can straightforwardly be extended to two-spin magnons on $\mathbb{R} \times S^3$.



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Outlook

- Can the sine-Gordon calculations be generalized to describe magnons on $\mathbb{R} \times S^3$? This would give a nice check of the finite gap results, but would require finding periodic solutions to the complex sine-Gordon model.
- Can the algebraic curve formalism be used to explicitly reconstruct the string solutions from the finite gap results?



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Thank you!

