

Pomeron-Graviton duality and resummation at high energies

Anna Staśto

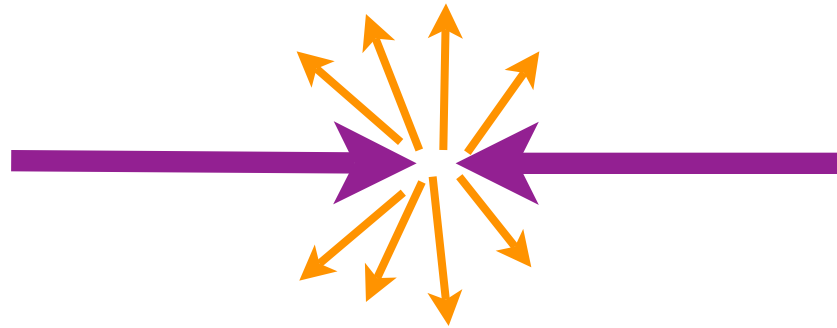
Penn State University, University Park, PA, USA

and

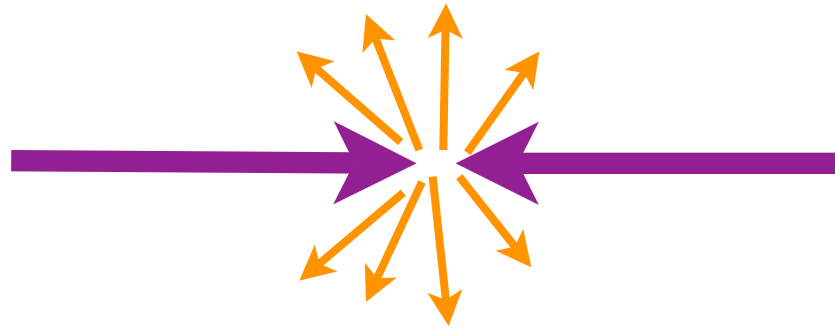
Institute of Nuclear Physics, Kraków, Poland

Cracow School of Theoretical Physics, XLVII Course, 2007

High - energy scattering of hadrons

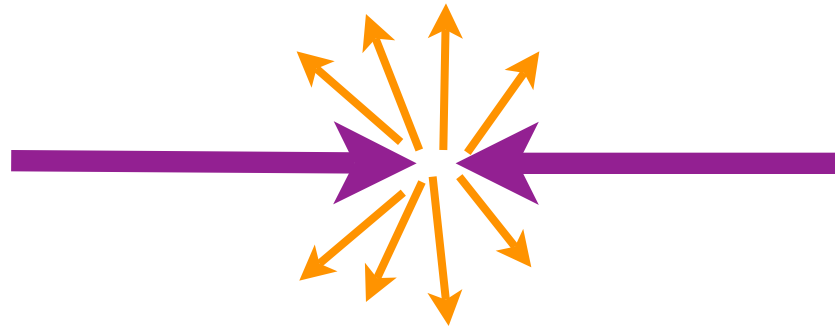


High - energy scattering of hadrons



What is the dependence of the cross section at high energies?

High - energy scattering of hadrons



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Gauge theory

weak gauge coupling

Pomeron:
collective state of gluons

High - energy scattering of hadrons



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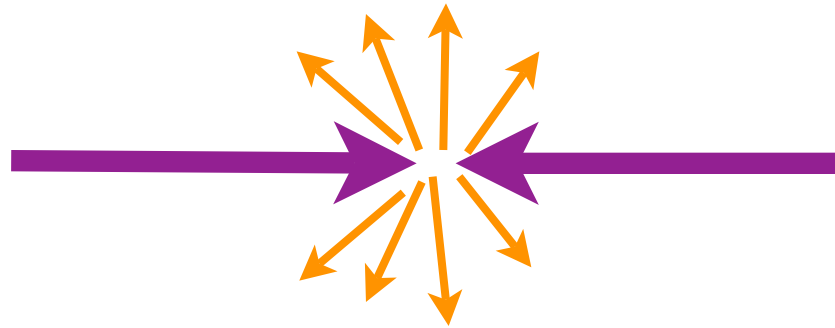
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collective state of gluons

String theory

strong gauge coupling

5-dimensional graviton
in anti - de Sitter space

High - energy scattering of hadrons



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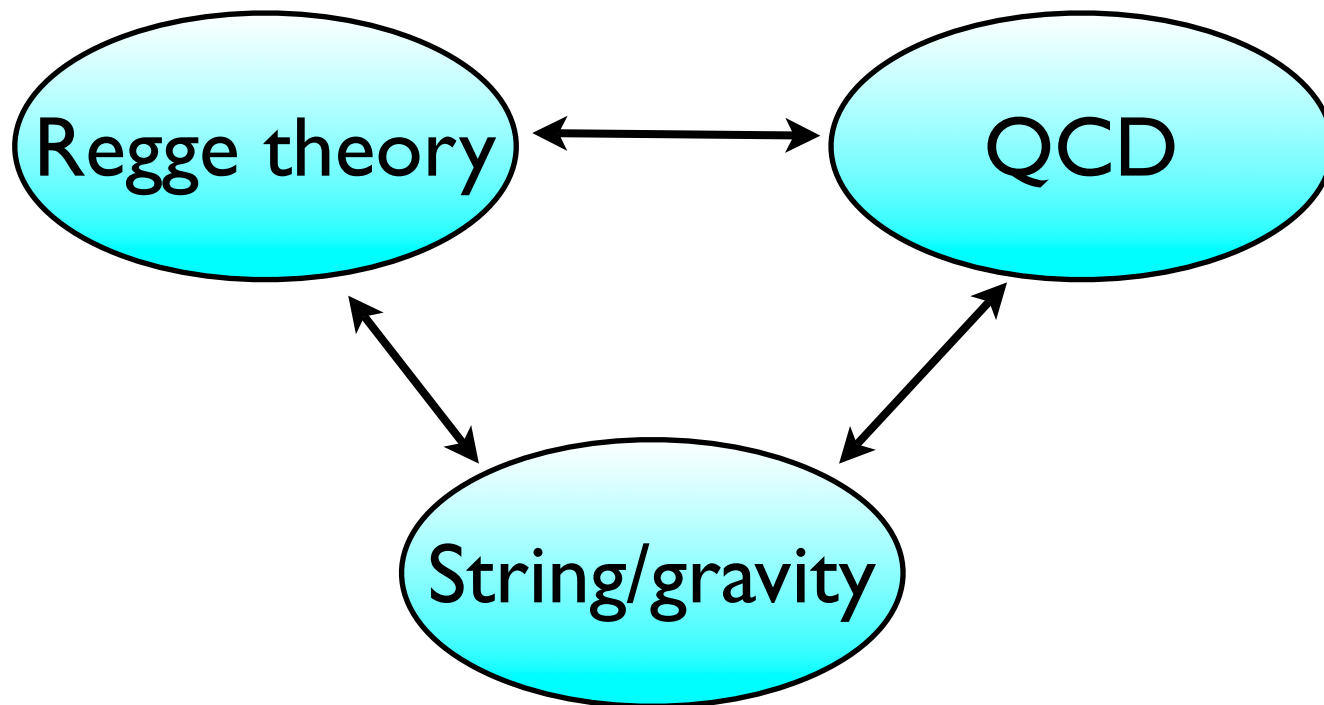
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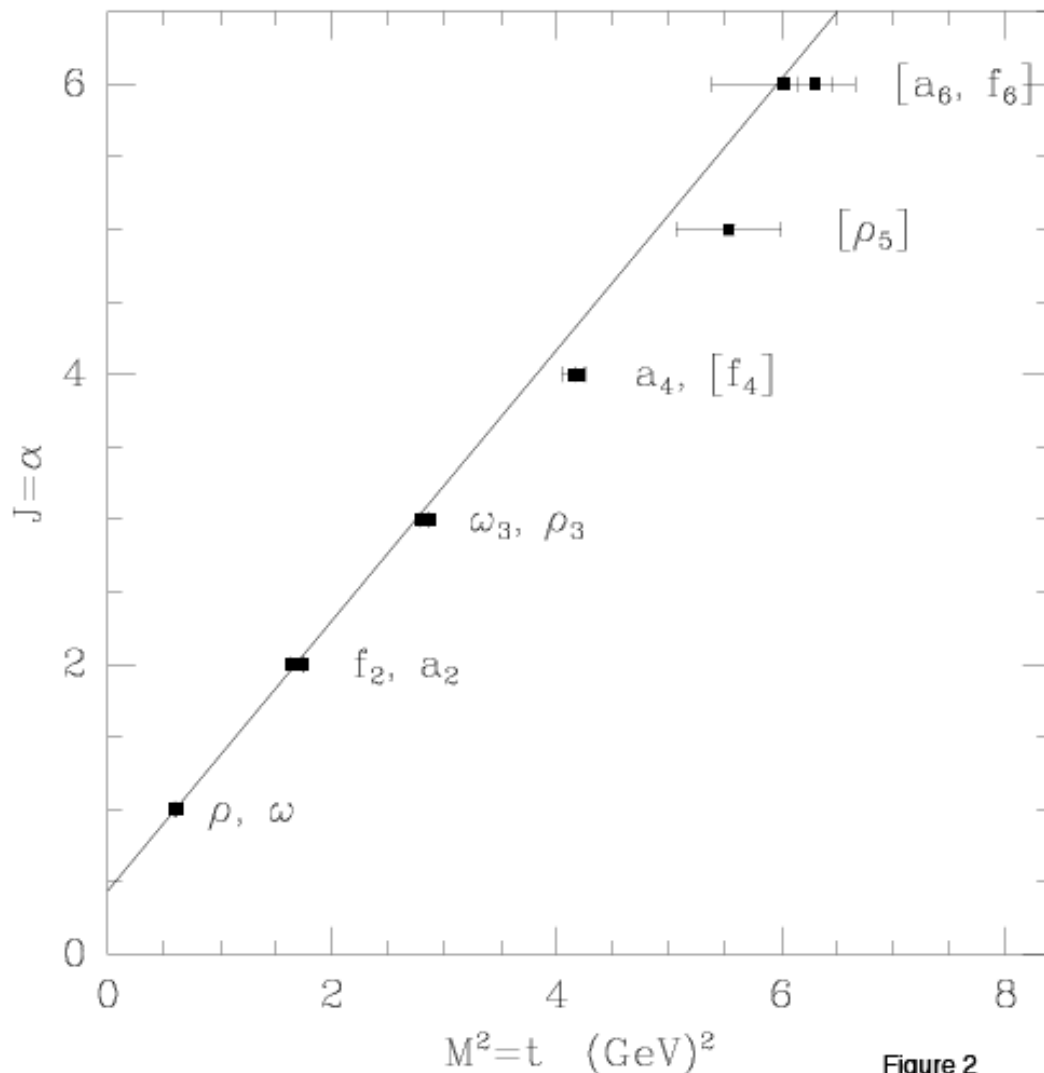
Resummation



Regge theory

Regge trajectories

- Linear dependence of the spin J on the square mass for mesons.



$$\alpha(t) = \alpha(0) + \alpha' t$$

$$\alpha(0)_\rho = 0.45 \quad \text{intercept}$$

$$\alpha'_\rho = 0.93 \text{ GeV}^{-2} \quad \text{slope}$$

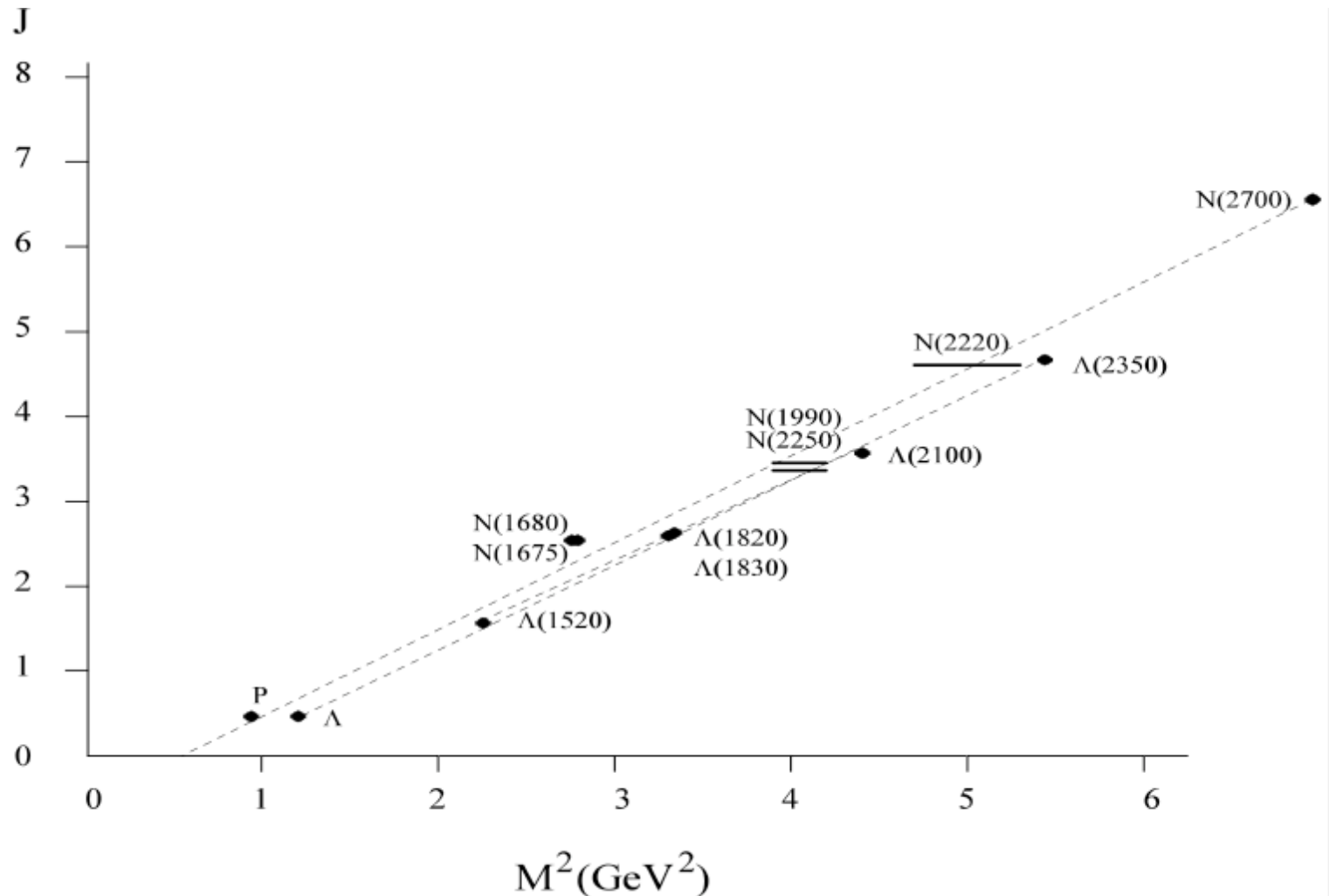
$$t = M^2 > 0$$

for bound states

Figure 2

Regge trajectories

Similar regularity for baryons



S-matrix and Regge theory

S-matrix and Regge theory

General assumption about the S matrix:

- Lorentz invariance $\mathcal{A}(s, t)$
- Unitarity $S^\dagger S = S S^\dagger = 1$
- Analyticity
- Crossing $\mathcal{A}_{ab \rightarrow cd}(t, s, u) = \mathcal{A}_{a\bar{c} \rightarrow \bar{b}d}(s, t, u)$

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Decomposition into partial waves

$$\mathcal{A}(s, t) = \frac{1}{2i} \oint_C dl (2l + 1) \frac{a(l, t)}{\sin \pi l} P(l, 1 + 2s/t)$$

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Partial wave
amplitude

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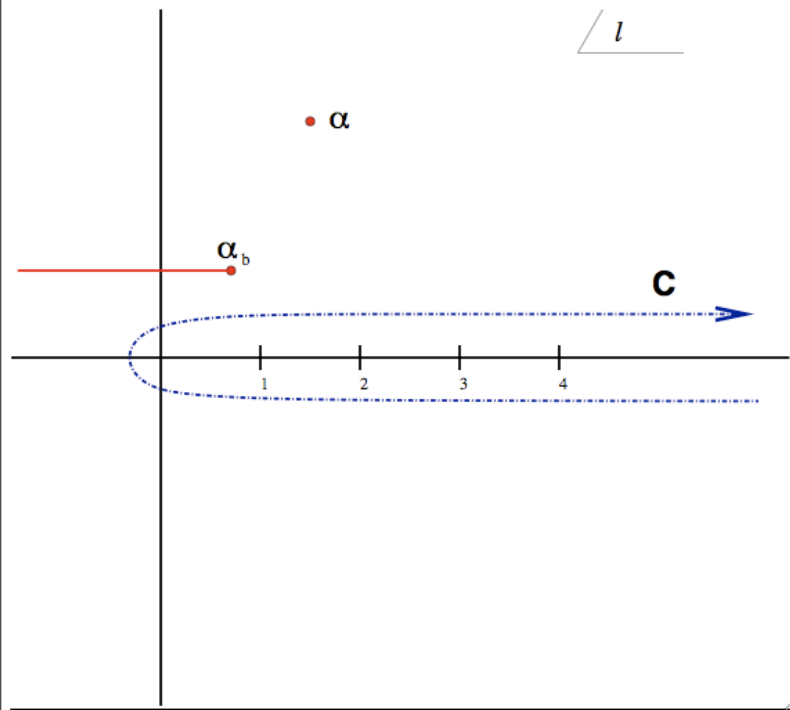
Angular momentum

Partial wave
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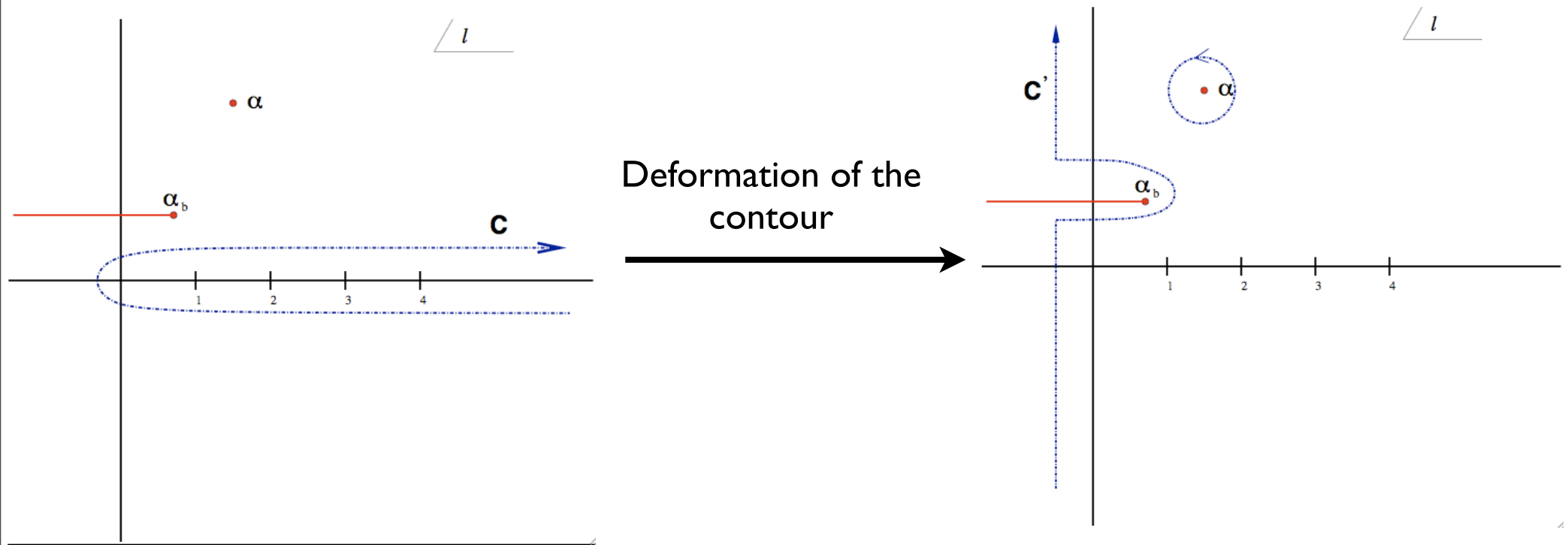
Legendre
polynomial

Complex angular momentum plane

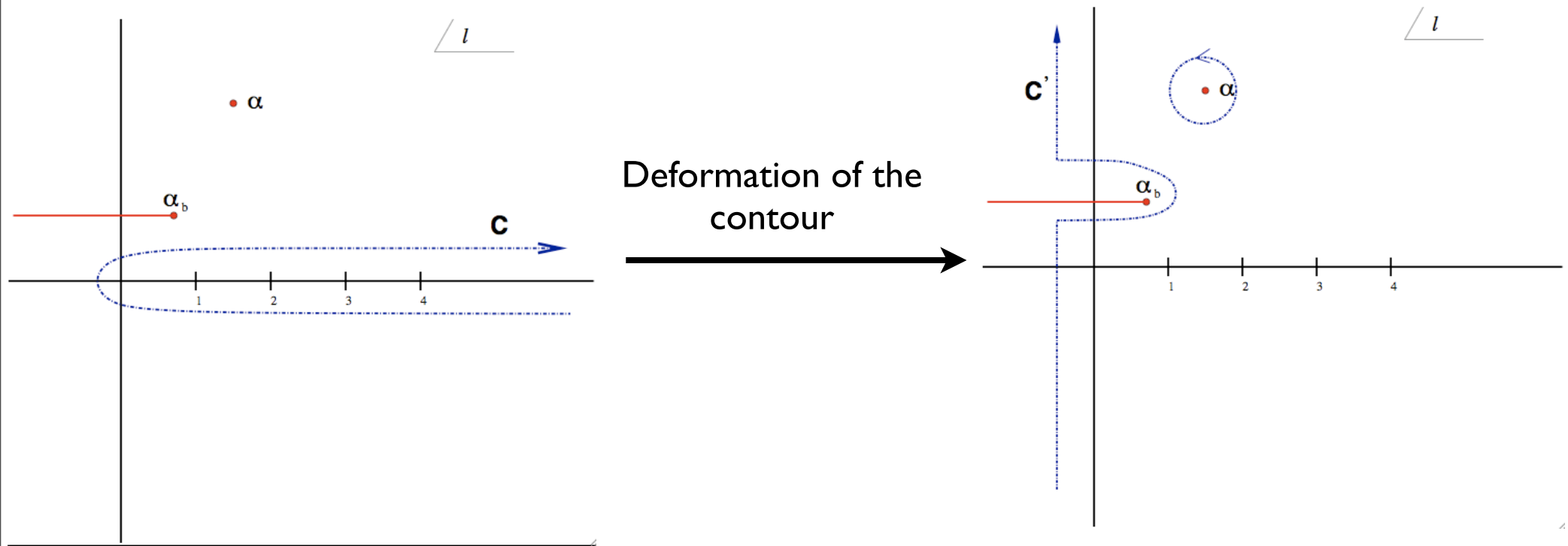
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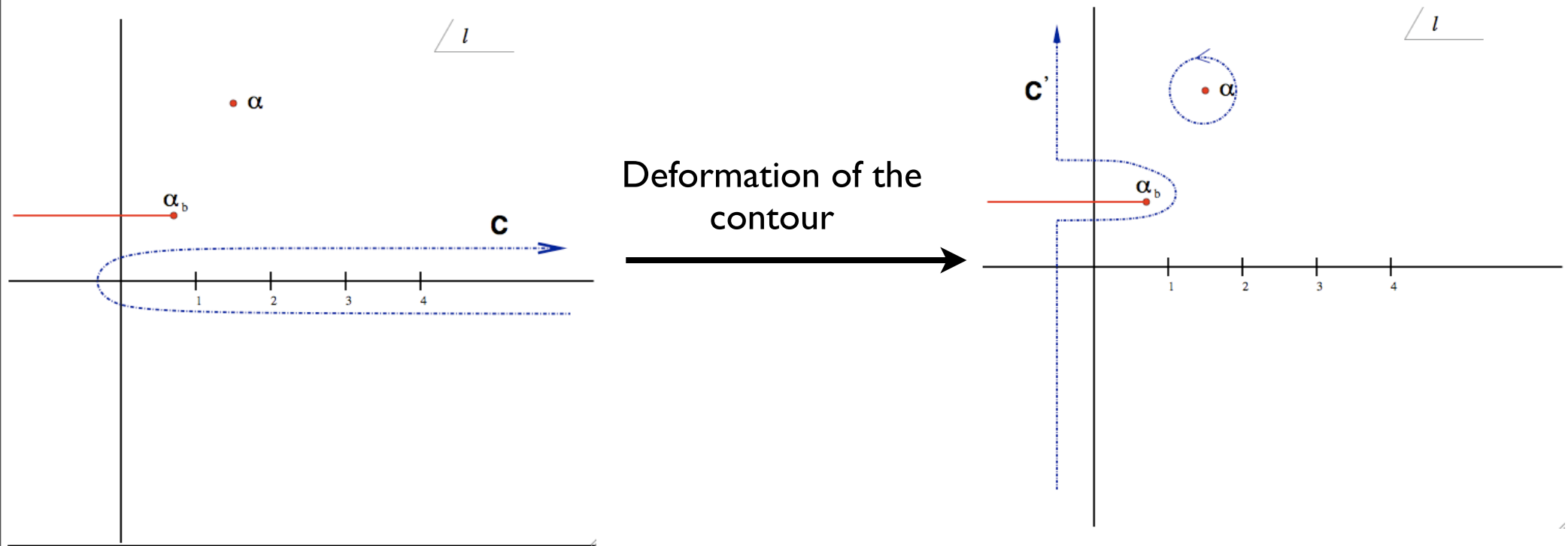
Complex angular momentum plane



In the Regge limit:

$$s \gg |t| \quad (s \rightarrow \infty, t \text{ fixed})$$

Complex angular momentum plane

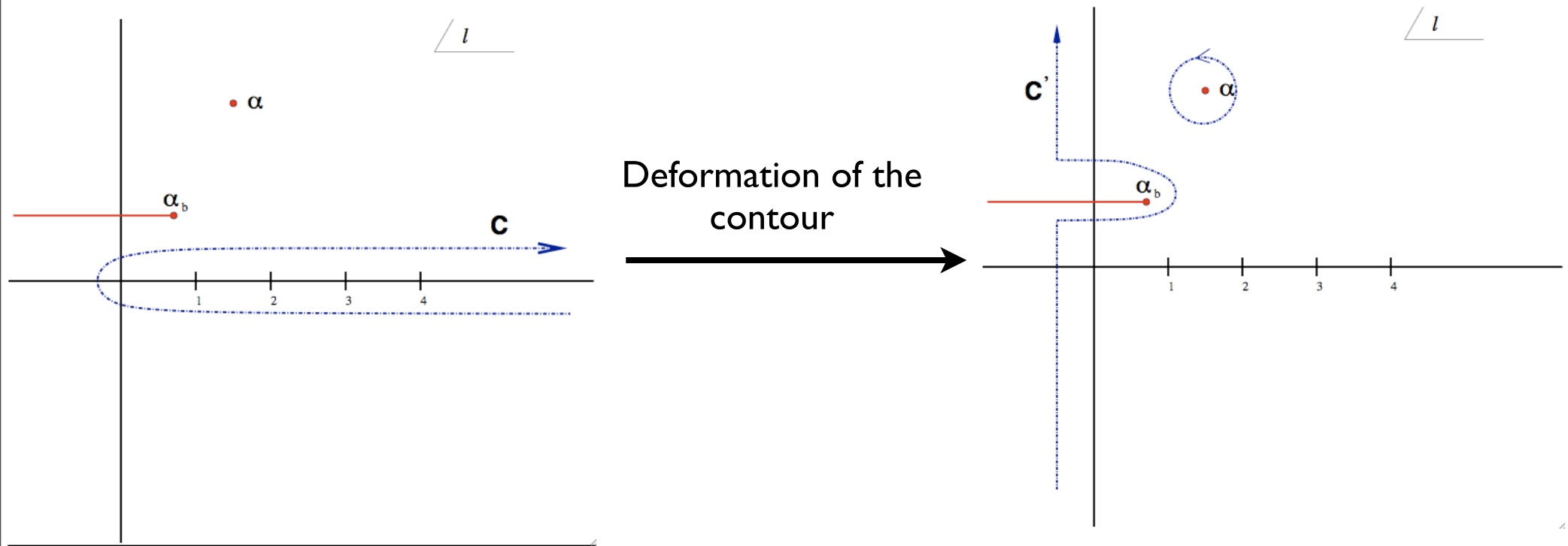


In the Regge limit:

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$$A(s, t) \rightarrow \frac{\eta + e^{-i\pi\alpha(t)}}{2} \beta(t) s^{\alpha(t)}$$

Complex angular momentum plane



In the Regge limit: $s \gg |t|$ ($s \rightarrow \infty, t$ fixed)

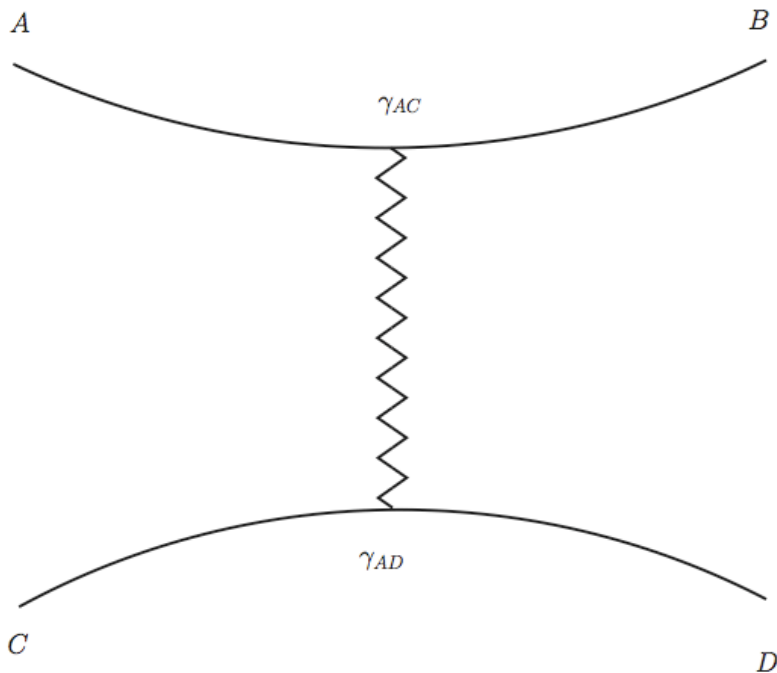
$$A(s, t) \rightarrow \frac{\eta + e^{-i\pi\alpha(t)}}{2} \beta(t) s^{\alpha(t)}$$

Amplitude dominated by the Regge pole with largest

$\text{Re } \alpha(t)$

Reggeon exchange

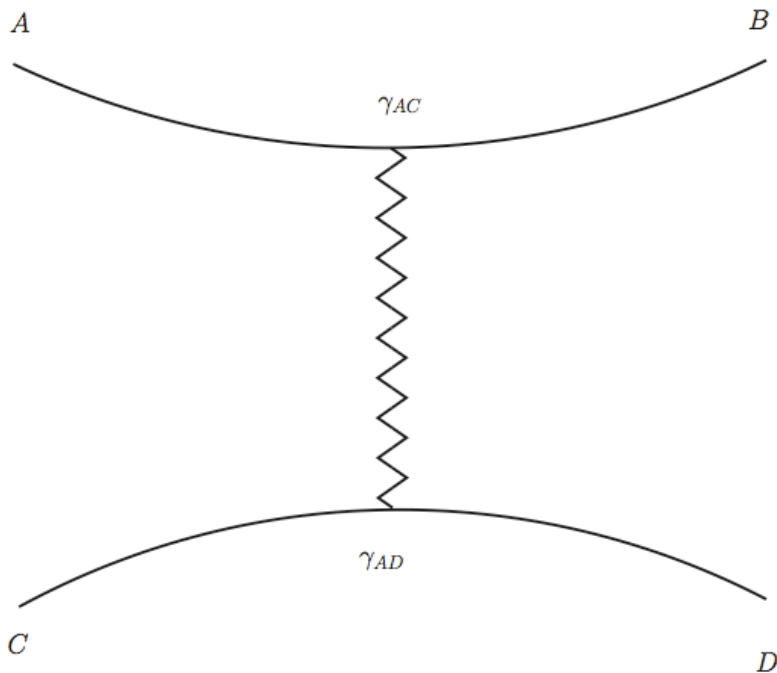
Reggeon exchange



Factorization of the couplings
and the Reggeon exchange

$$\mathcal{A}(s, t) \rightarrow \frac{\eta + e^{-i\pi\alpha(t)}}{2 \sin \pi\alpha(t)} \frac{\gamma_{ac}(t)\gamma_{bd}(t)}{\Gamma(\alpha(t))} s^{\alpha(t)}$$

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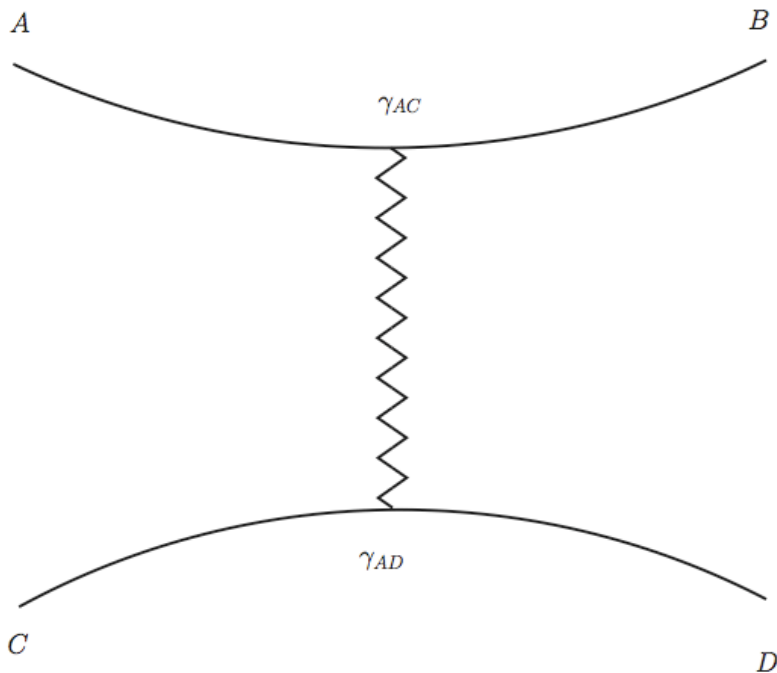


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Signature

Reggeon exchange



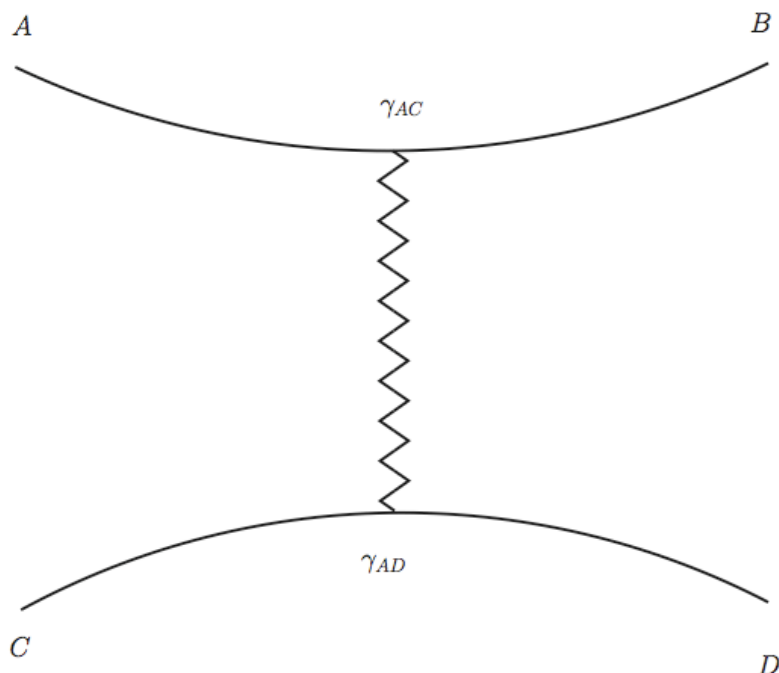
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Signature

Couplings

The energy behavior of the amplitude is determined
by the exchange of the quasi-particle: **Reggeon**

Pomeron

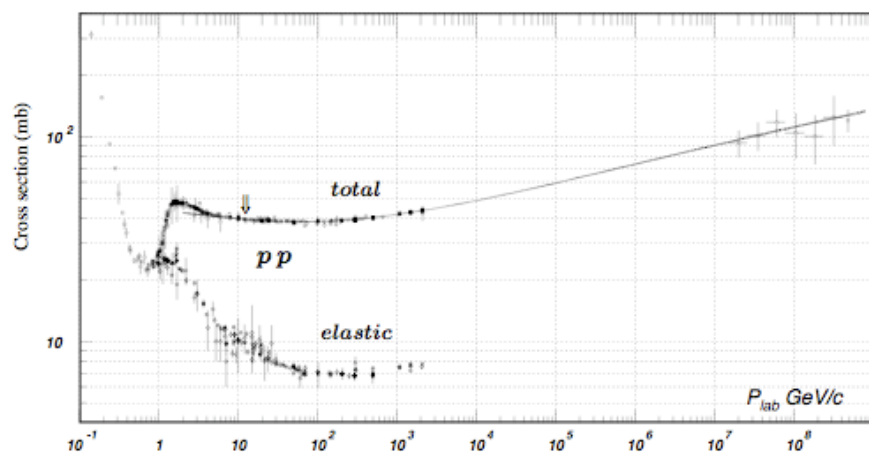
Pomeron

Vacuum exchange dominates cross sections at high energies

Pomeron

Vacuum exchange dominates cross sections at high energies

$$\sigma_{\text{TOT}}(p\bar{p}) \sim \sigma_{\text{TOT}}(pp)$$



vs GeV

1.9 2 10 10² 10³ 10⁴

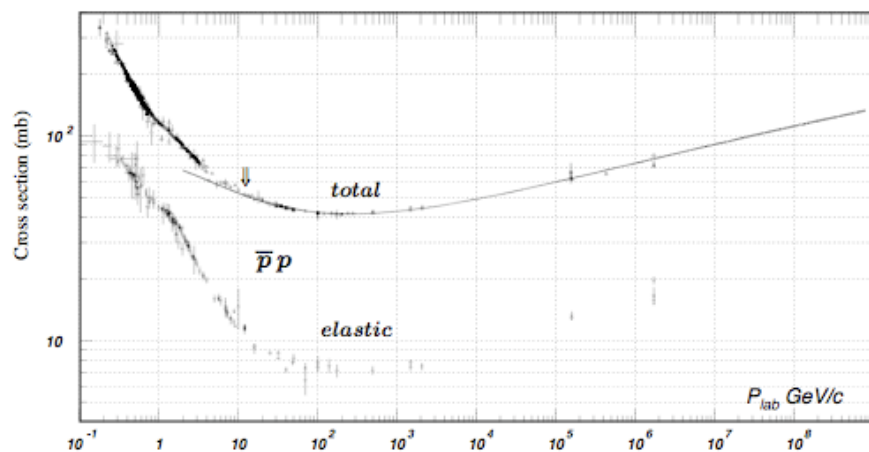


Figure 40.11: Total and elastic cross sections for pp and $\bar{p}p$ collisions as a function of laboratory beam momentum and total center-of-mass energy. Corresponding computer-readable data files may be found at <http://pdg.lbl.gov/xsect/contents.html>. (Courtesy of the COMPAS group, IHEP, Protvino, August 2005.)

Pomeron

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Vacuum exchange

$$\alpha(0)_P \geq 1$$

experimentally: $\alpha_P(0) \simeq 1.08, \sigma_{\text{TOT}} \sim s^{(\alpha(0)_P - 1)}$

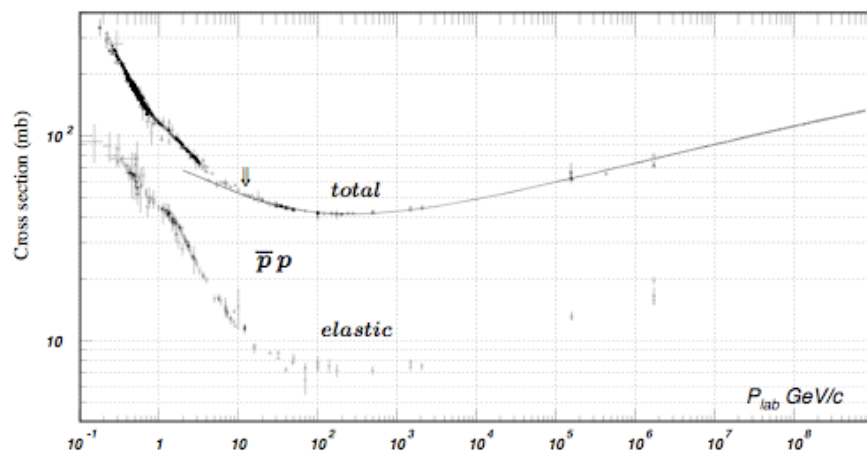
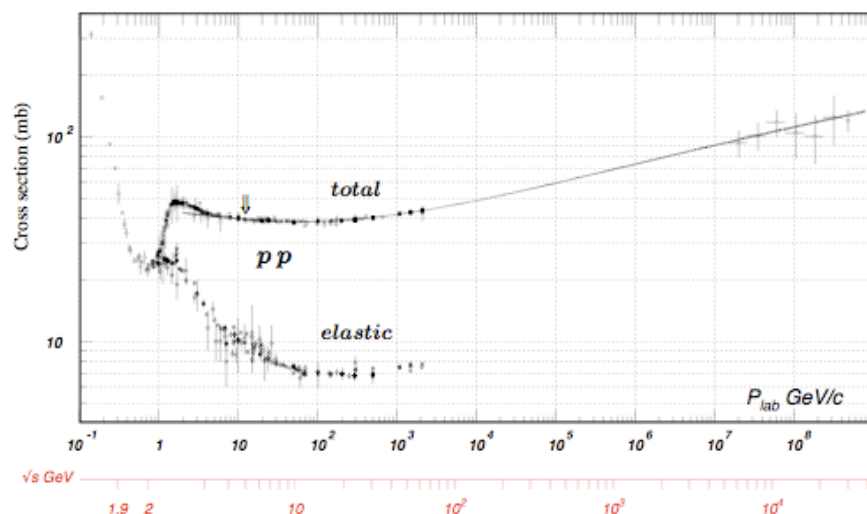
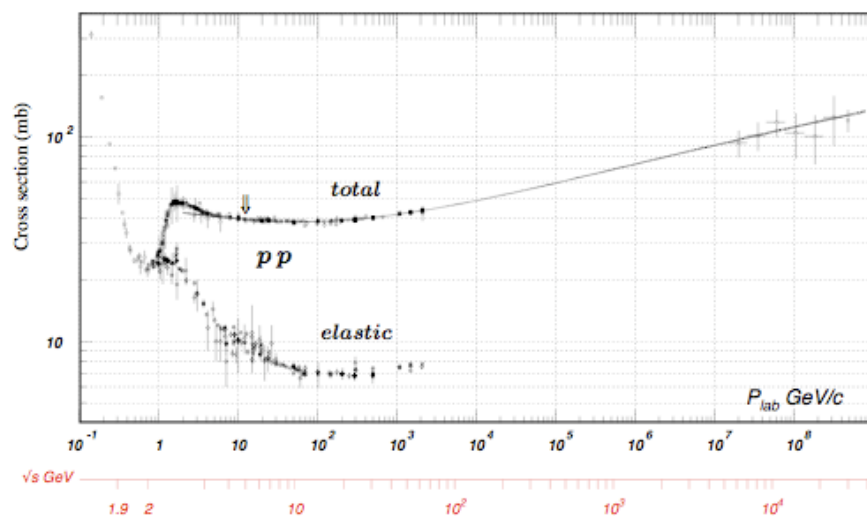


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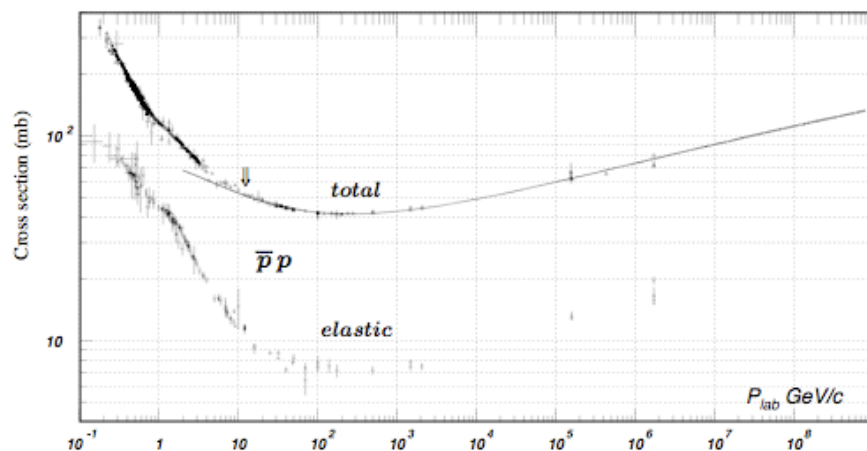


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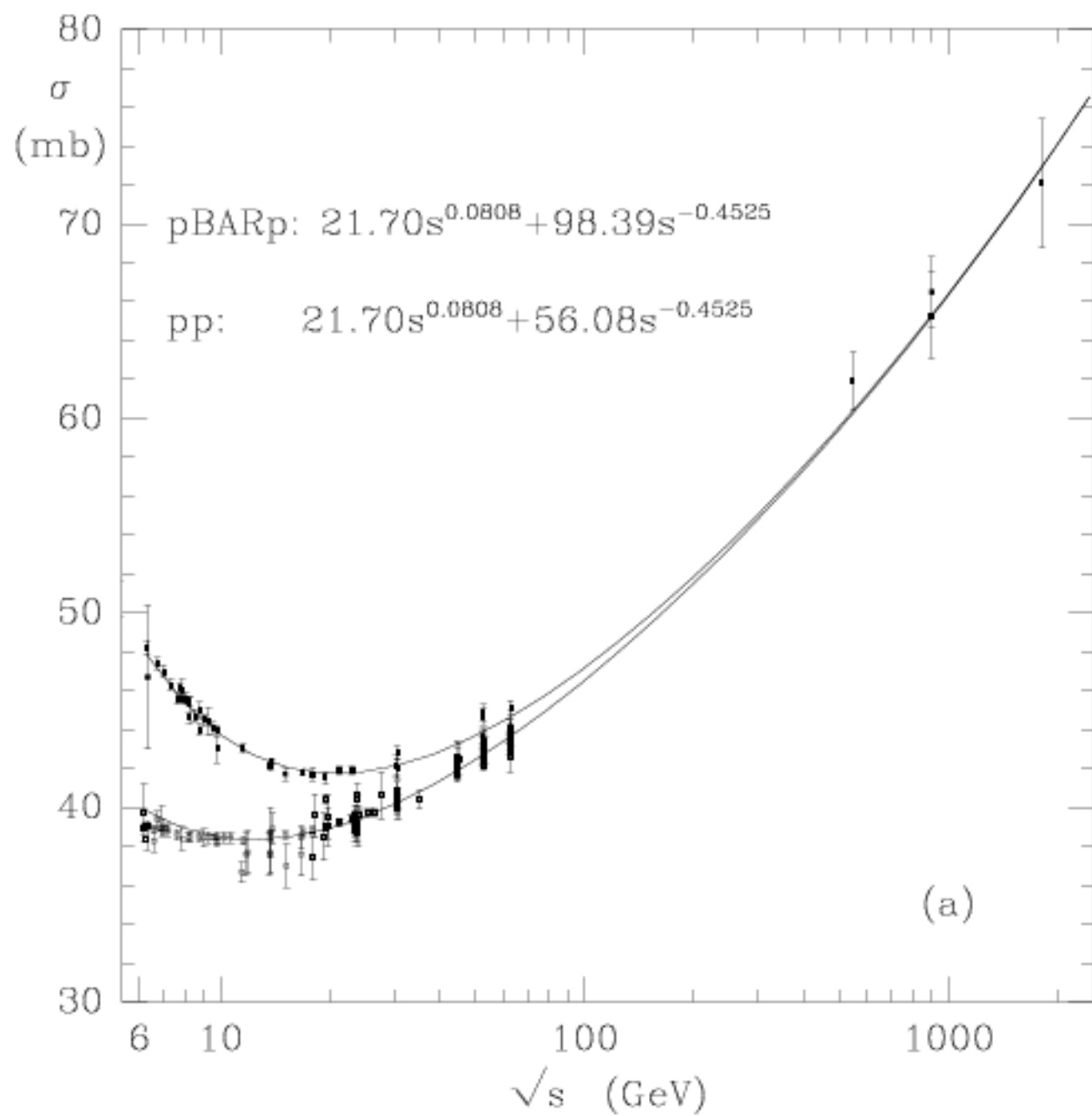
Non-vacuum exchange

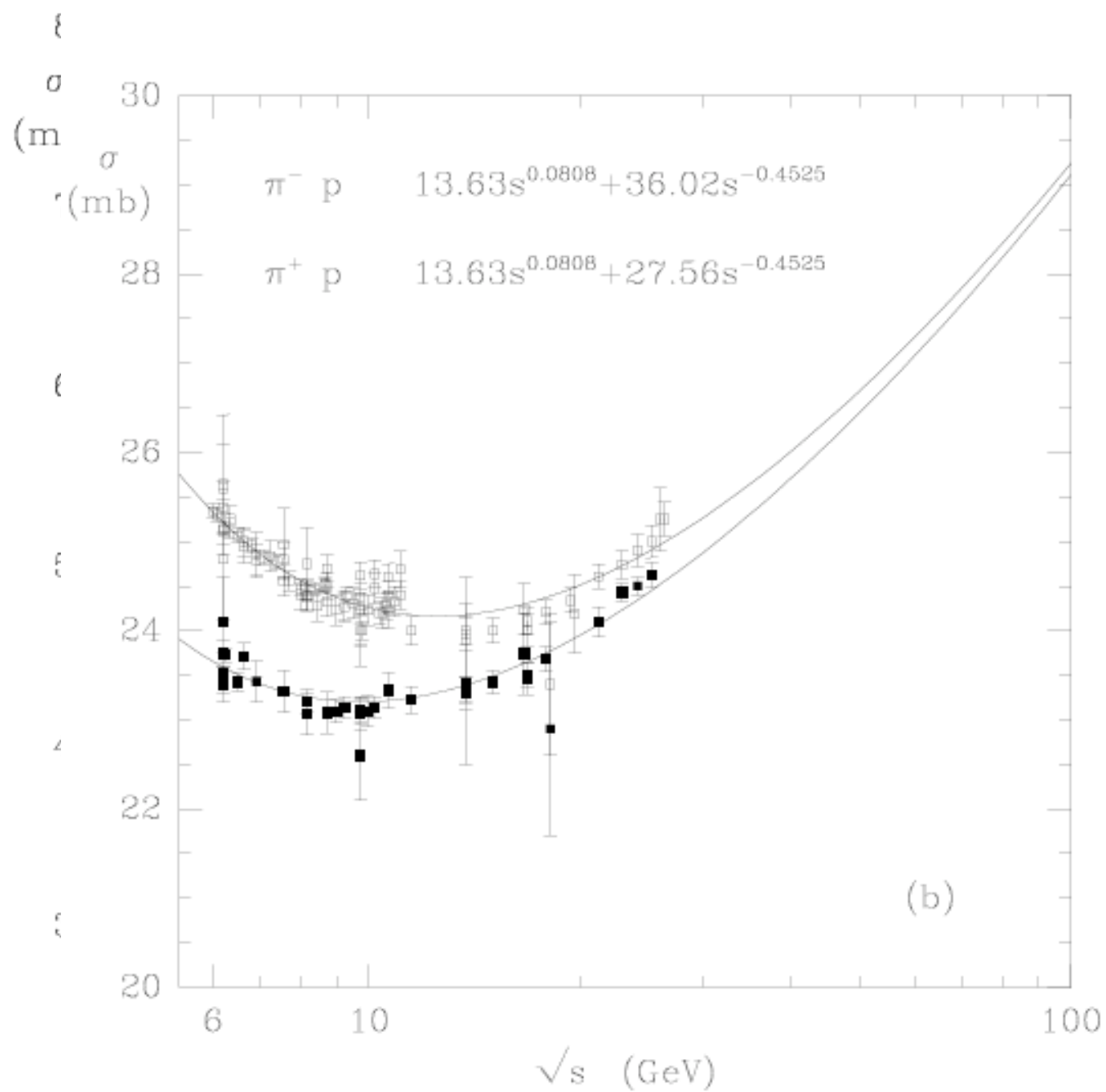
$$\alpha(0)_R < 1$$

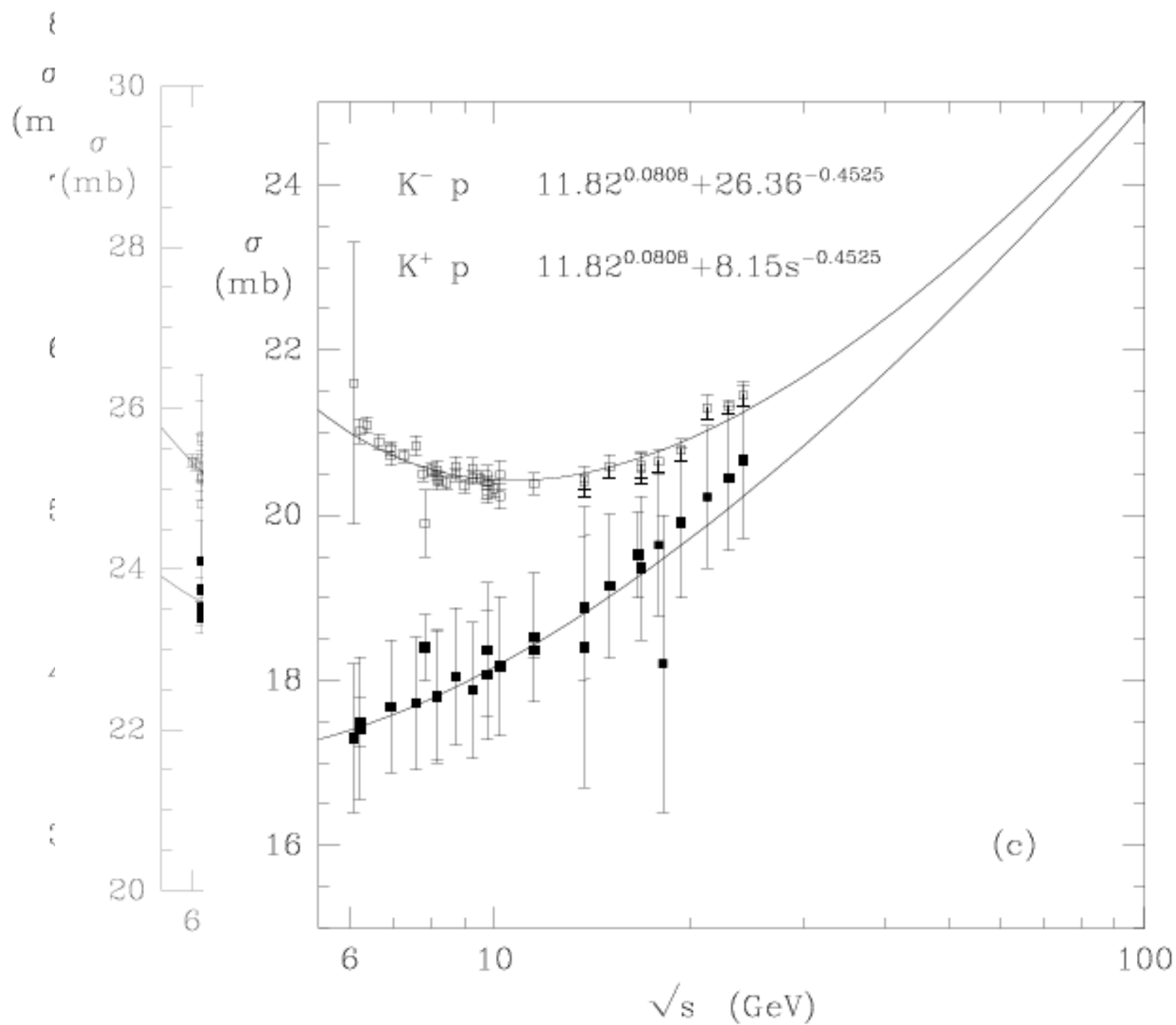
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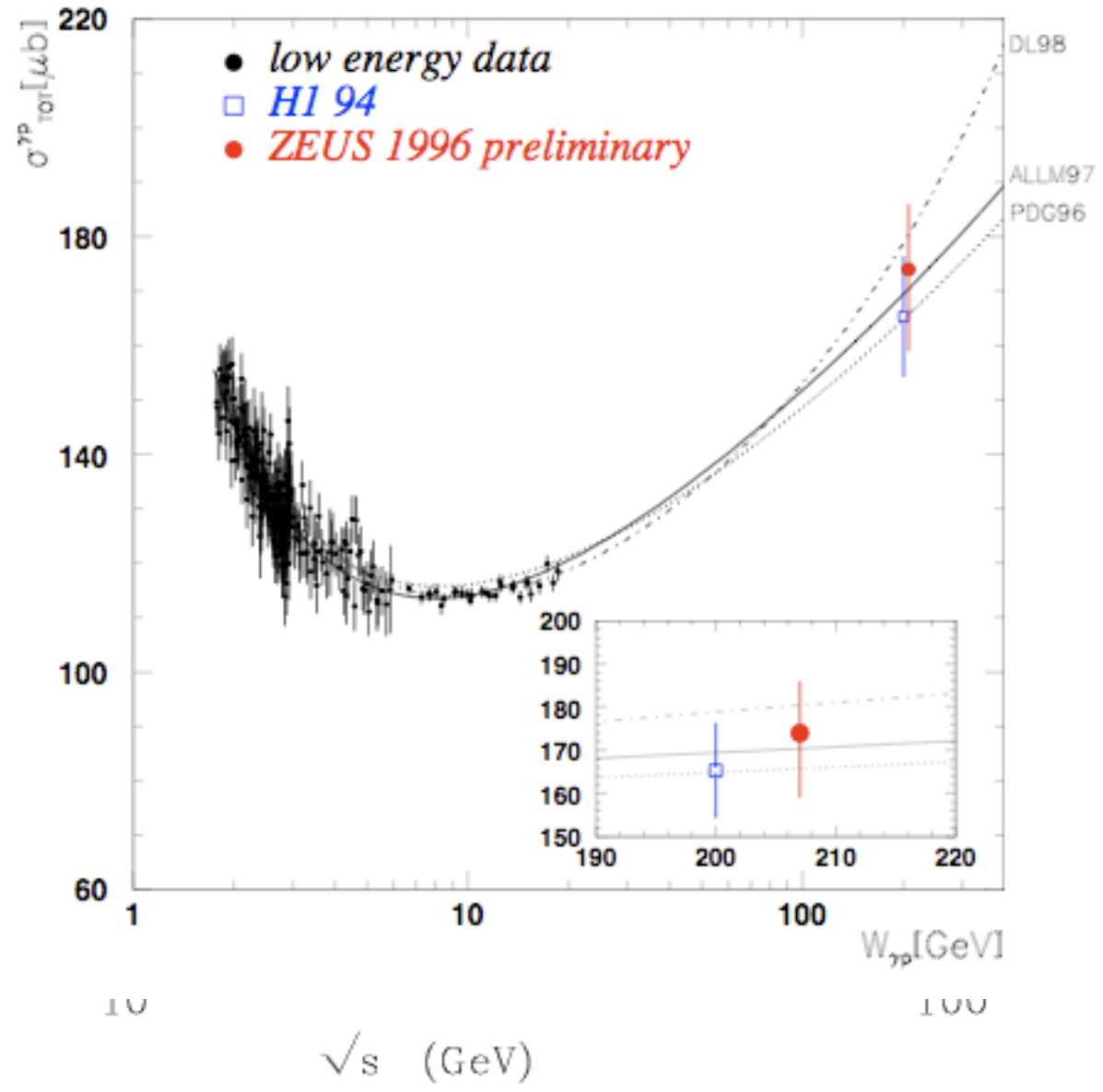
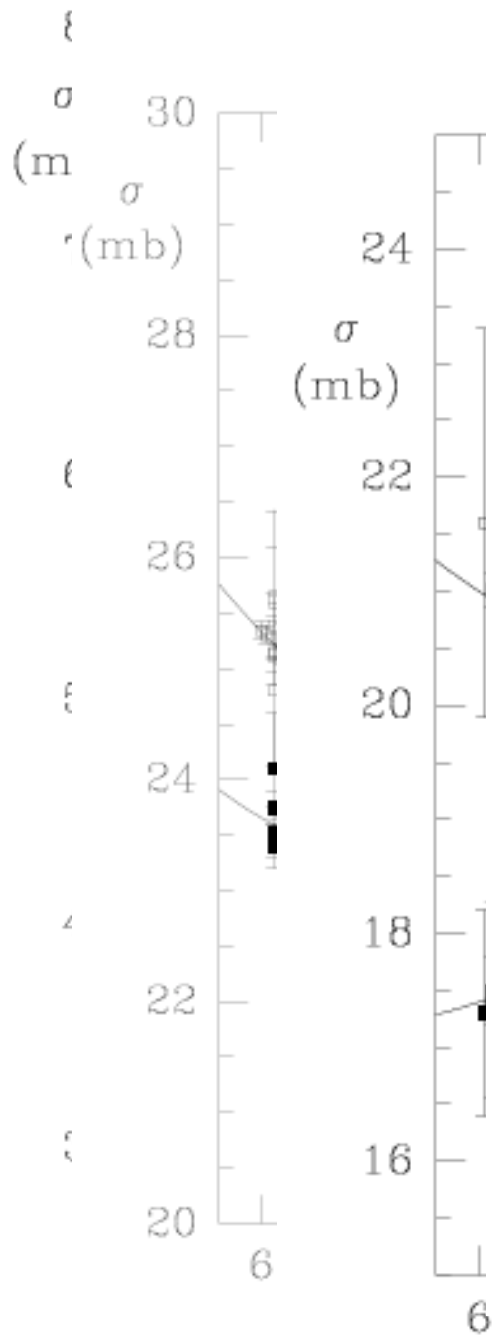
Note: odderon

$$\alpha_O(0) \leq 1$$

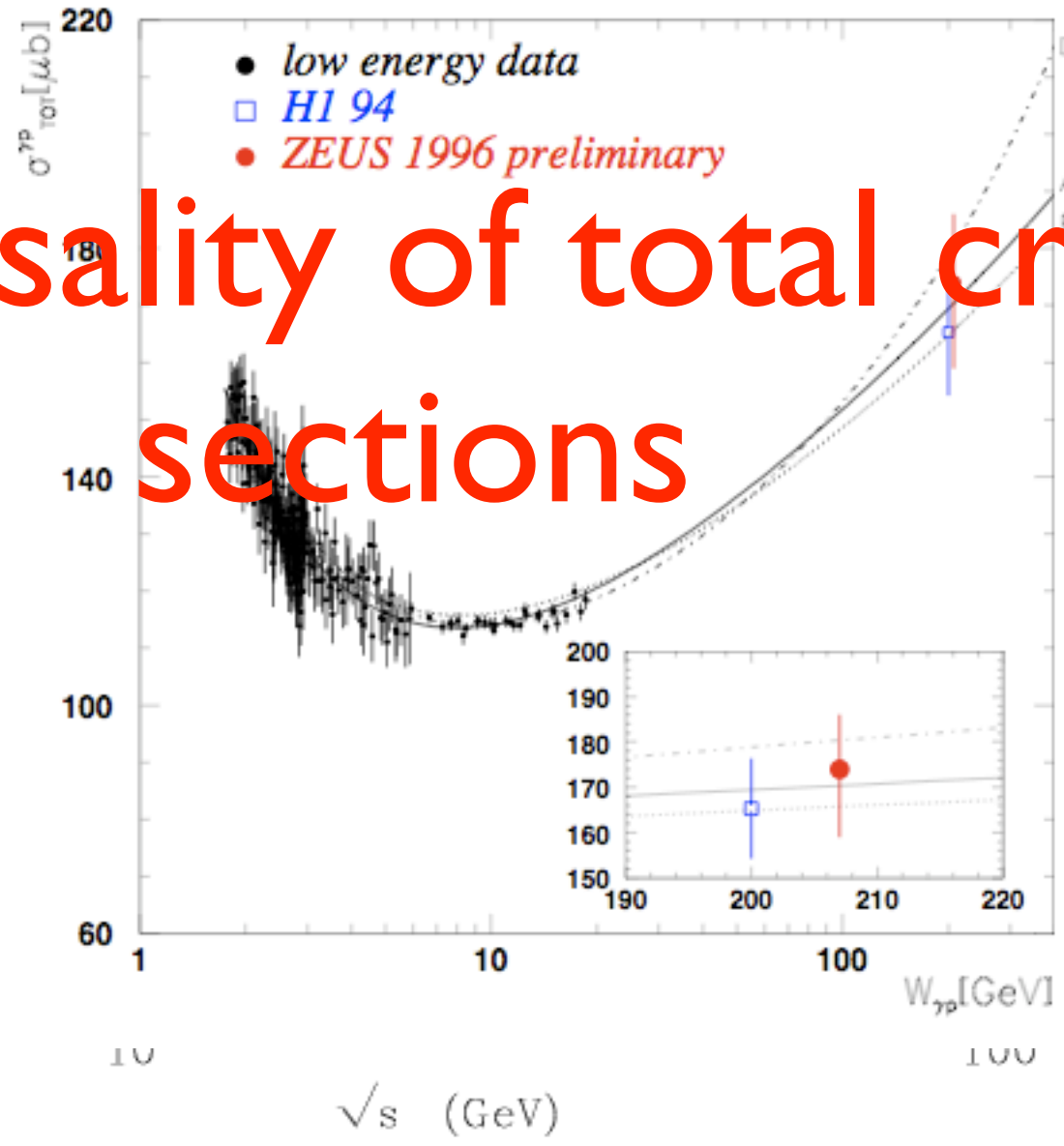
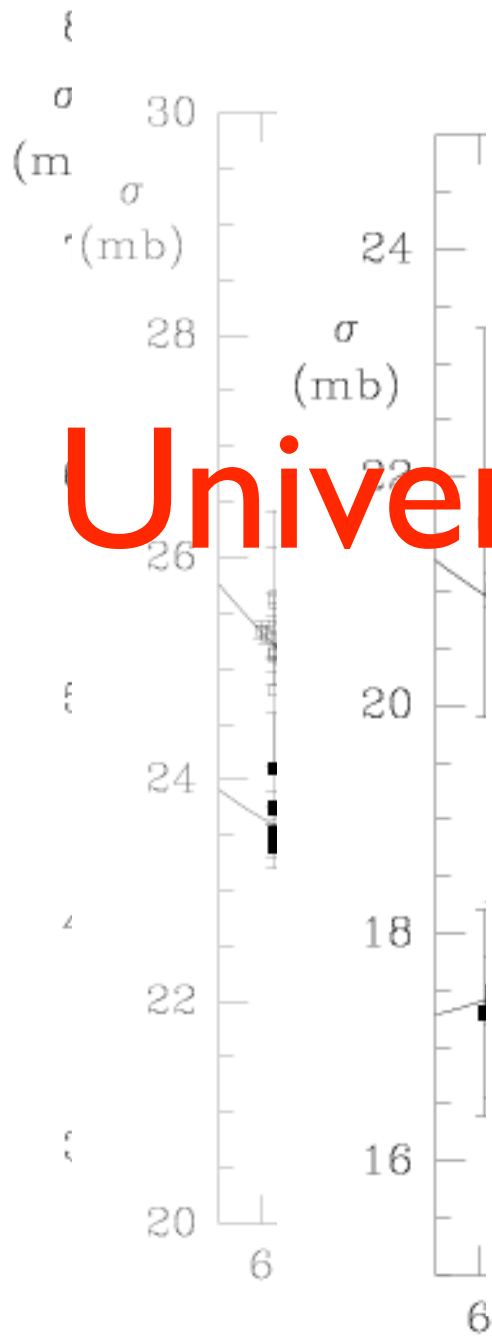


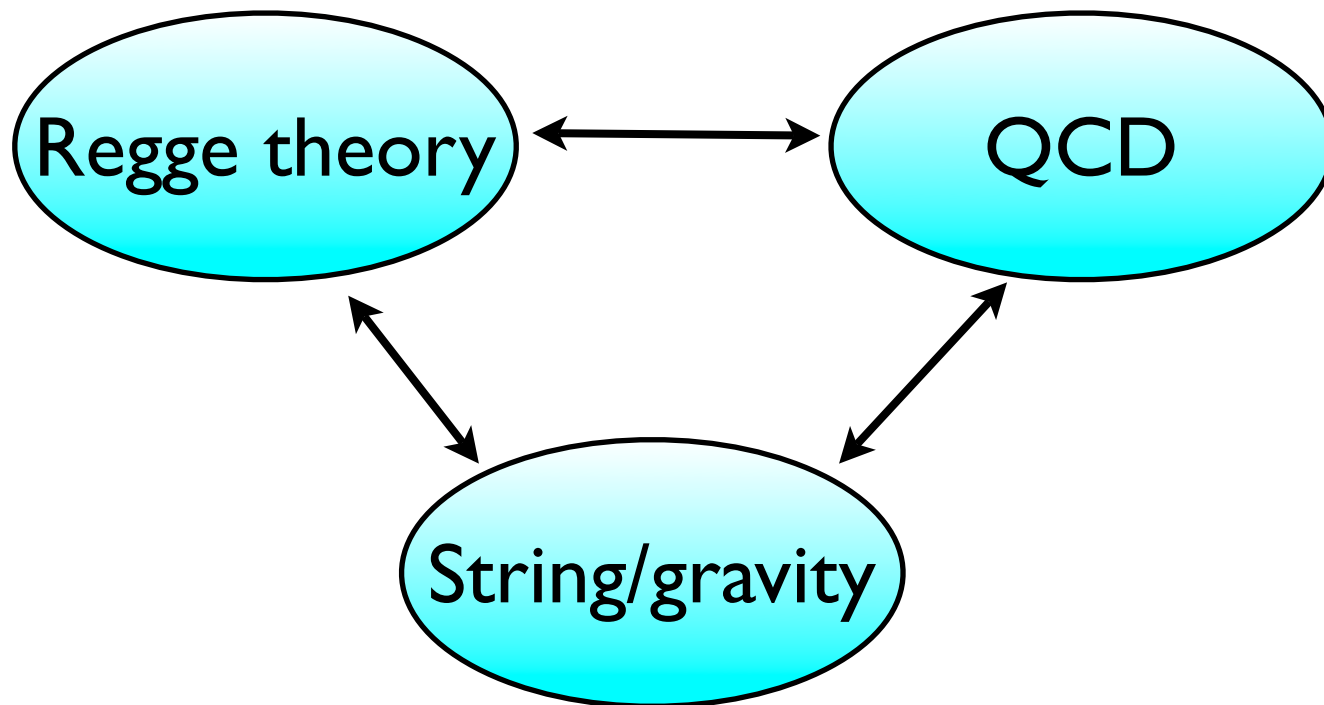


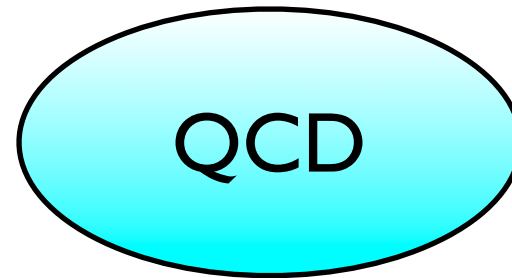




Universality of total cross sections





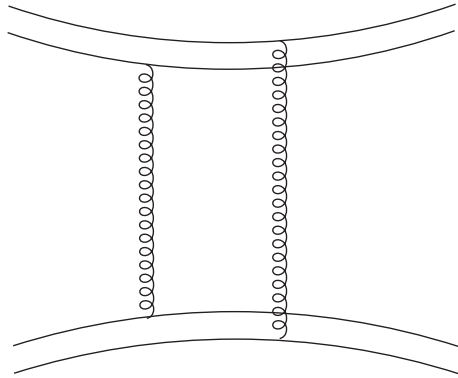


QCD

Pomeron in gauge theory

Low-Nussinov model

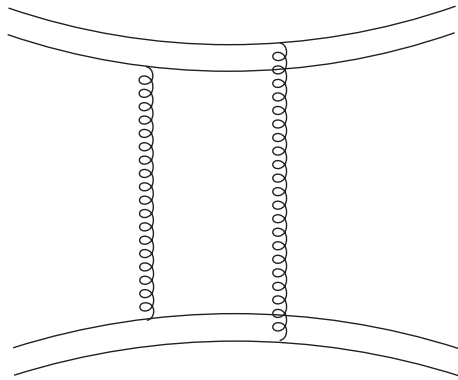
2-gluon exchange



Pomeron in gauge theory

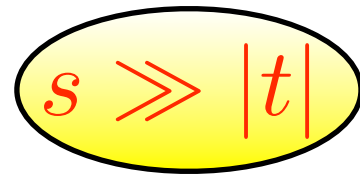
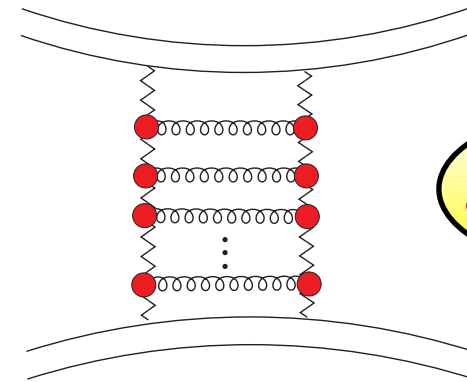
Low-Nussinov model

2-gluon exchange



BFKL resummation

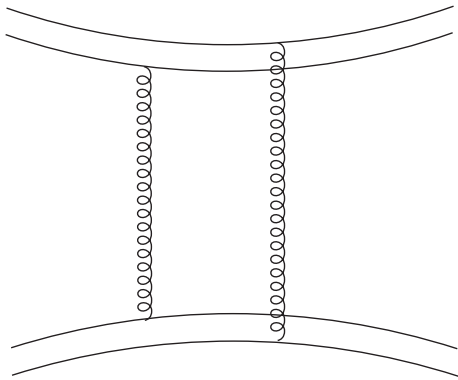
color singlet



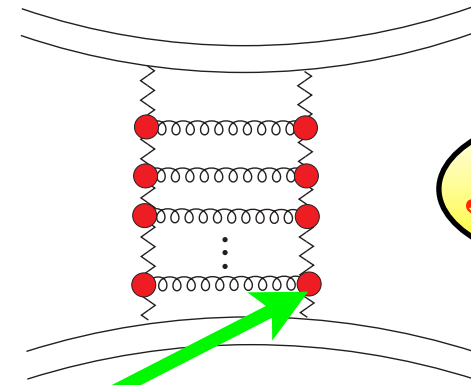
Balitskii
Fadin
Kuraev
Lipatov

Pomeron in gauge theory

Low-Nussinov model
2-gluon exchange



BFKL resummation
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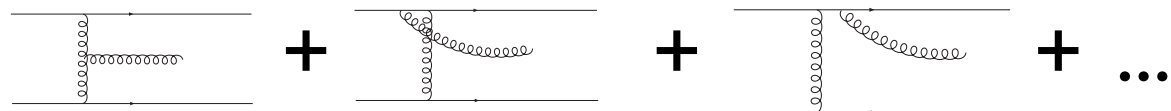


Balitskii
Fadin
Kuraev
Lipatov

$$s \gg |t|$$

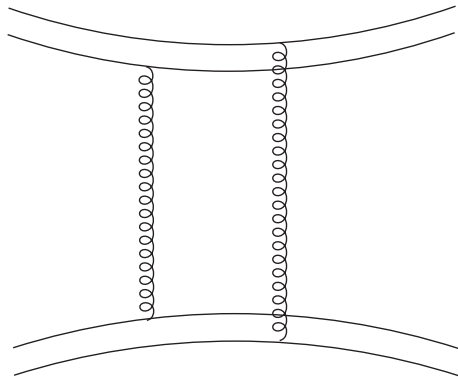
Effective vertex

$$\Gamma_{\mu\nu}^{\sigma}(k_i, k_{i+1})$$

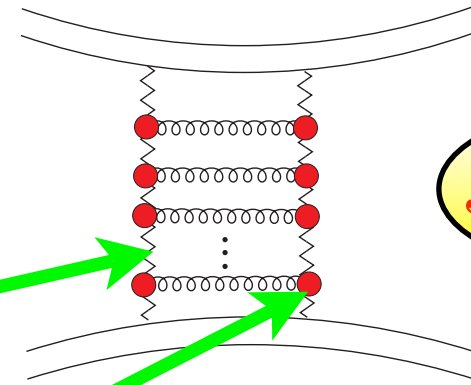


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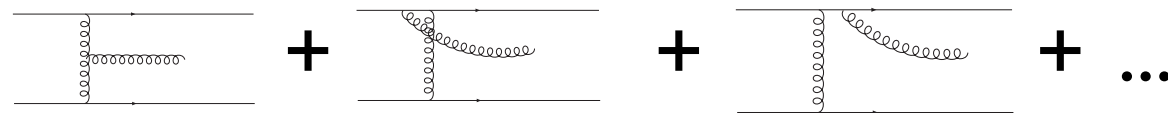
$$s \gg |t|$$

Reggeized gluon:

$$D_{\mu\nu}(\hat{s}, k_T^2) = \frac{ig_{\mu\nu}}{k_T^2} \left(\frac{\hat{s}}{k_T^2} \right)^{\epsilon_G(k_T^2)}$$

Effective vertex

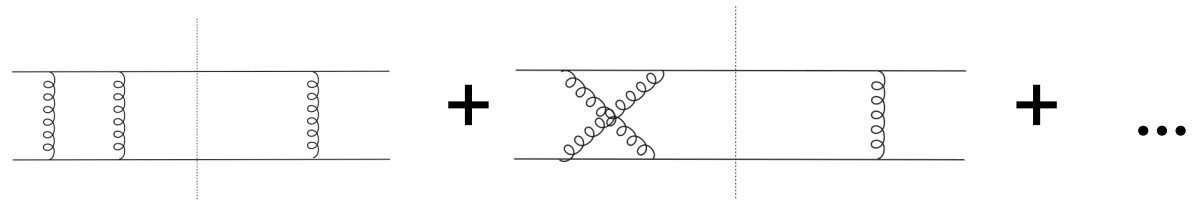
$$\Gamma_{\mu\nu}^\sigma(k_i, k_{i+1})$$



$$\hat{s}_i = (k_{i-1} - k_{i+1})^2$$

Regge trajectory: virtual diagrams

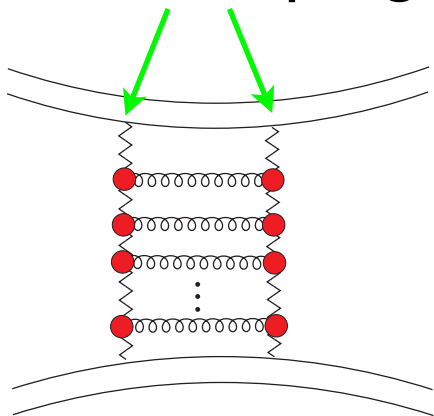
$$\epsilon_G(q_T^2) = \frac{N_c \alpha_s}{4\pi^2} \int_{\Lambda} d^2 k_T \frac{-q_T^2}{k_T^2 (k_T - q_T)^2}$$



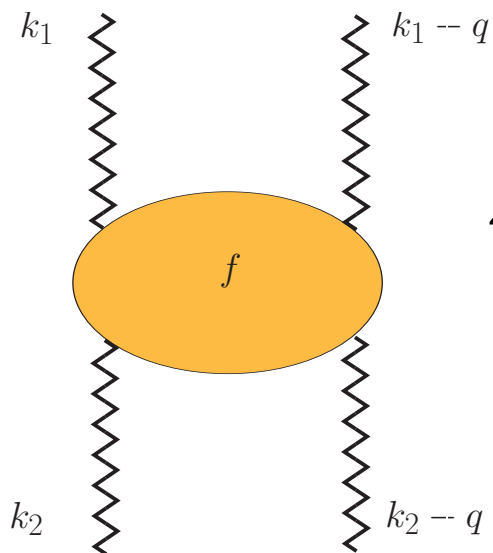
Infrared divergent!

Integral equation

Eikonal couplings



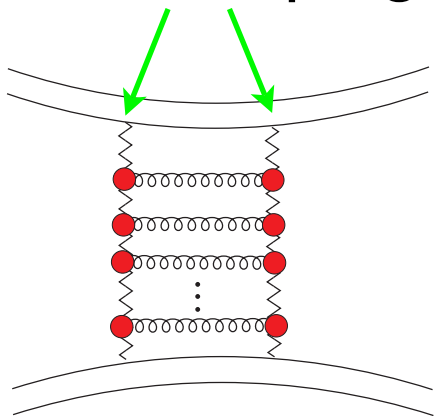
Universality



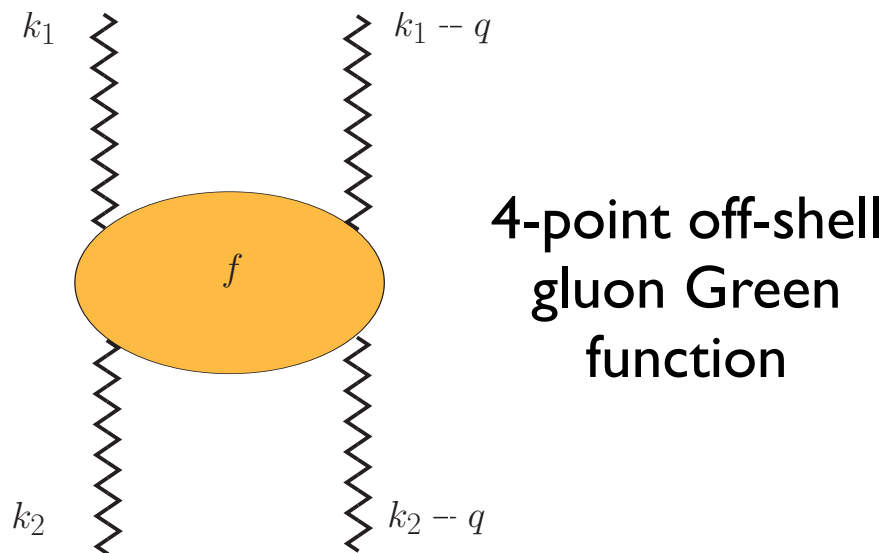
4-point off-shell
gluon Green
function

Integral equation

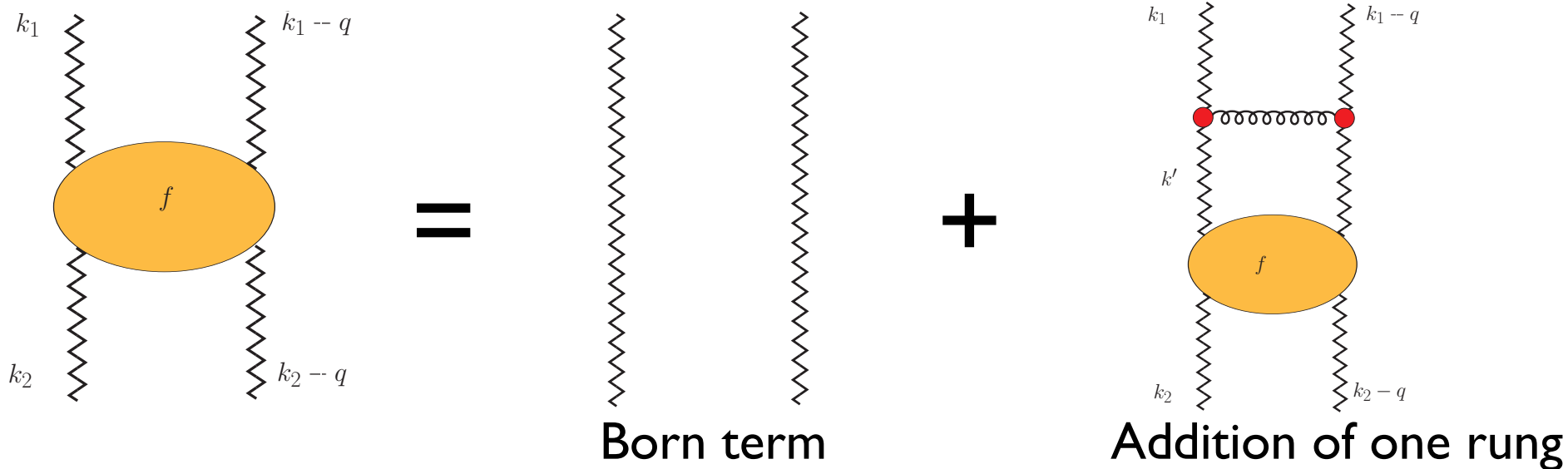
Eikonal couplings



Universality



Integral equation for the Pomeron



Integral equation

$$f(Y; k_{1T}, k_{2T}, q_T) = f^{(0)}(k_{1T}, k_{2T}, q_T) + \int_0^Y dy K(k_{1T}, k_{2T}, q_T) \otimes f(y; k_{1T}, k_{2T}, q_T)$$

Rapidity: $Y = \ln 1/x = \ln s/s_0$ Convolution in transverse momenta

! Scale choice (irrelevant at lowest order) !

Integral equation

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! Factorization of longitudinal and transverse components of momenta !

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Rapidity: $Y = \ln 1/x = \ln s/s_0$ Convolution in transverse momenta

! Scale choice (irrelevant at lowest order) !

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Mellin transform: $\int dY e^{(-\omega-1)Y} f(Y) dY = f(\omega)$

$$\omega f(\omega; k_{1T}, k_{2T}, q_T) = \delta^{(2)}(k_{1T} - k_{2T}) + K(k_{1T}, k_{2T}, q_T) \otimes f(\omega; k_{1T}, k_{2T}, q_T)$$

Integral kernel has Mobius invariance.

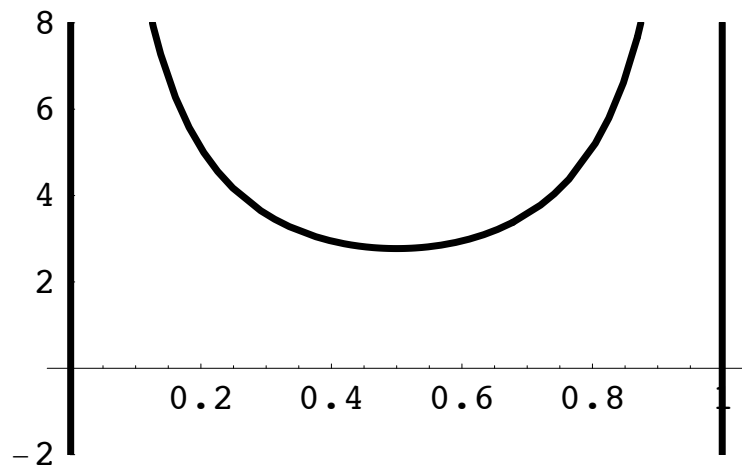
Solution of the BFKL equation

At zero momentum transfer: $q_T = 0$

Eigenfunctions: $\phi_\nu^n(k_T) = \frac{1}{\pi\sqrt{2}} (k_T^2)^{1/2+i\nu} e^{in\theta}$

Diagonalize equation: $K \otimes \phi_\nu^n = \frac{\alpha_s N_c}{\pi} \chi(\nu, n) \phi_\nu^n$

Eigenvalue (take $n=0$): $\chi(\nu, 0) = 2\psi(1) - \psi(1/2 + i\nu) - \psi(1/2 - i\nu)$



Simple poles:

$$\gamma = \dots, -2, -1, 0, 1, 2, \dots$$

$$\gamma = 1/2 + i\nu$$

Hard Pomeron

Hard Pomeron

Saddle point solution: around $\gamma = 1/2$

$$\chi(\nu) \simeq 4 \ln 2 - 14\zeta(3)\nu^2$$

Hard Pomeron

Saddle point solution: around $\gamma = 1/2$

$$\chi(\nu) \simeq 4 \ln 2 - 14\zeta(3)\nu^2$$

Approximate solution:

$$f(y = \ln s/s_0, k_{1T}, k_{2T}) \simeq \frac{1}{4\sqrt{k_{1T}^2 k_{2T}^2}} \frac{1}{\sqrt{14\zeta(3)\alpha_s N_c \pi^2 \ln s/s_0}} \left(\frac{s}{s_0}\right)^{4 \ln 2 \alpha_s N_c / \pi} \exp\left(-\frac{\pi \ln^2 \frac{k_{1T}^2}{k_{2T}^2}}{28\zeta(3)\alpha_s N_c \ln s/s_0}\right)$$

Hard Pomeron

Saddle point solution: around $\gamma = 1/2$

$$\chi(\nu) \simeq 4 \ln 2 - 14\zeta(3)\nu^2$$

Approximate solution:

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Diffusion pattern in transverse momenta

Hard Pomeron

Saddle point solution: around $\gamma = 1/2$

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Power-like growth with energy

Diffusion pattern in transverse momenta

Regge behavior from Feynman diagrams: $\alpha_P(0) = 1 + \frac{N_c \alpha_s}{\pi} 4 \ln 2$

Note: it is possible to compute Pomeron in electroweak theory also.

Pomeron phenomenology

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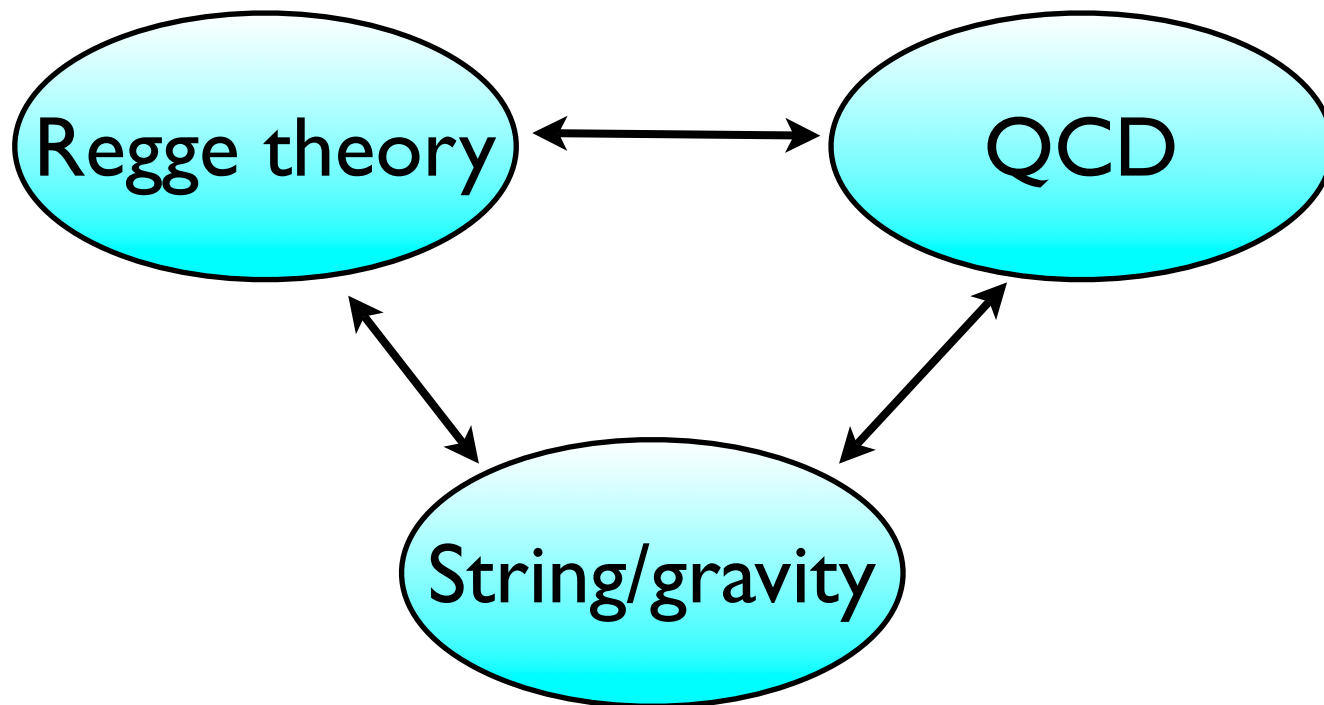
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 - Universal growth of the total cross sections.



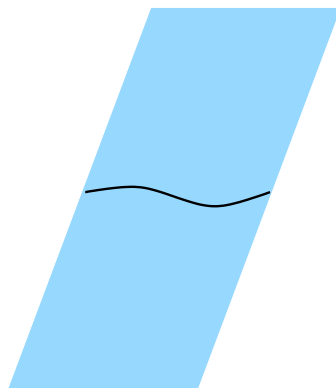
String/gravity

Graviton

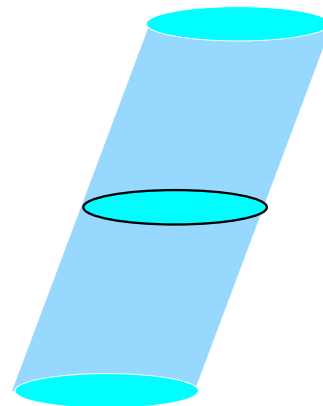
Graviton

Spin 2 massless (2 polarizations) particle:
symmetric rank 2 tensor.

In string theory: closed string state.



open



closed

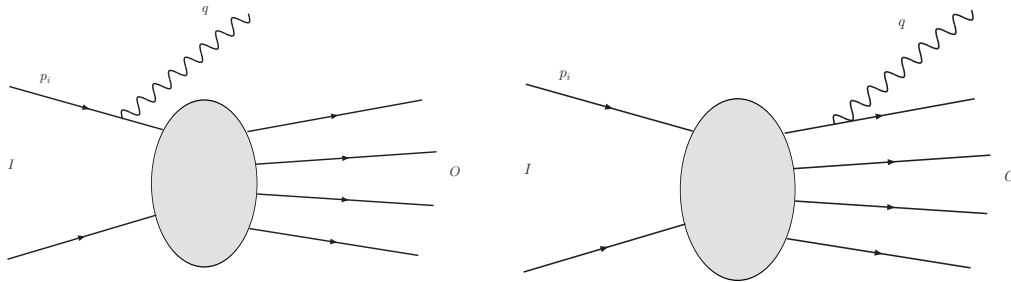
String theory includes gravity

Universality of couplings

Universality of couplings

Example: (Weinberg)

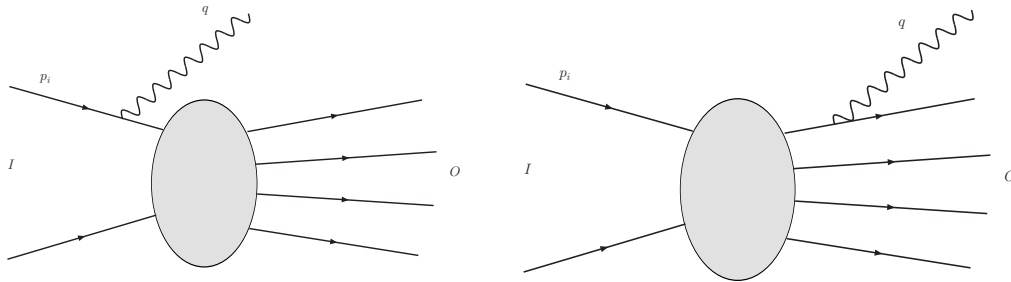
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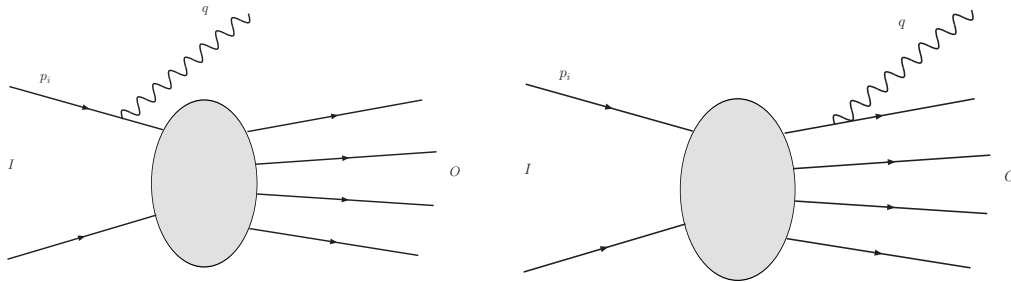
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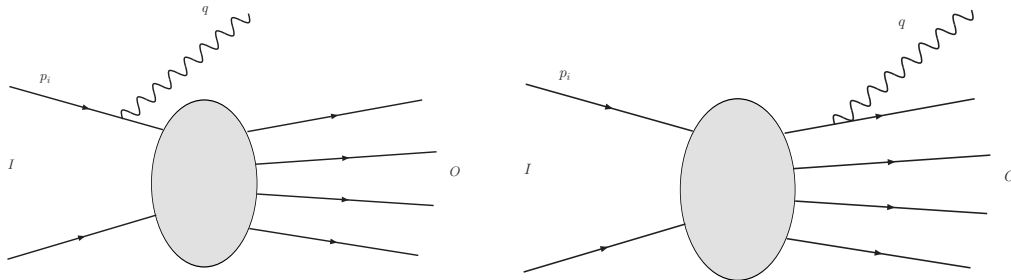
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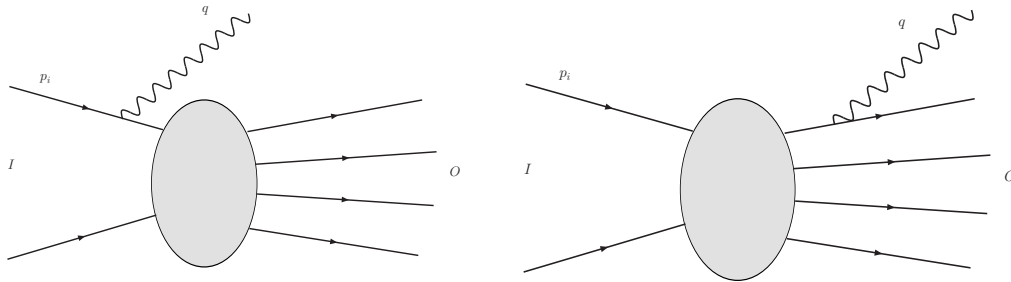
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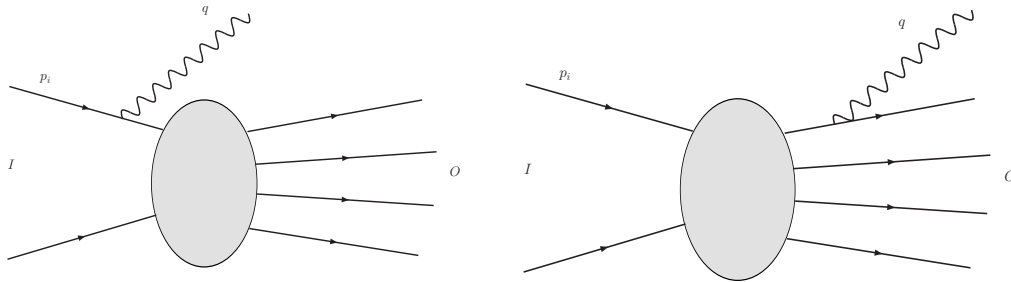
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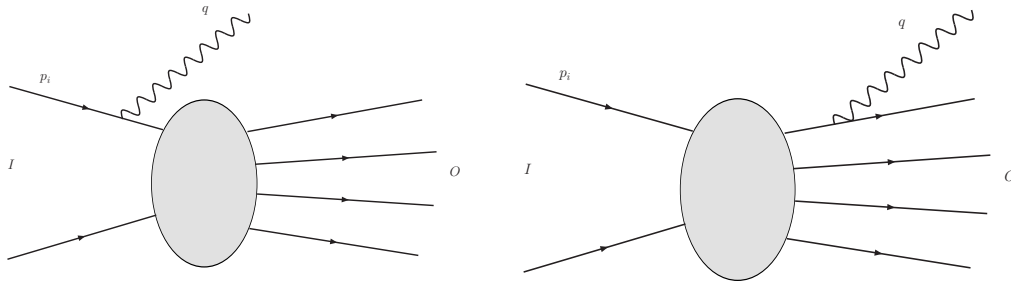
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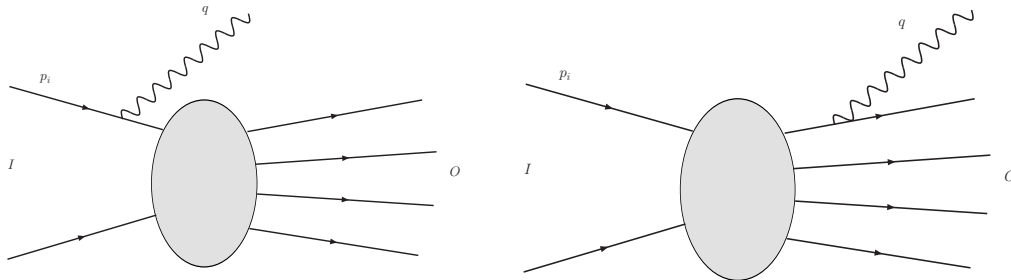
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Lorentz invariance for spin 2 particles gives principle of equivalence

Gauge/Gravity duality

Maldacena

strings in AdS(D) \longleftrightarrow CFT(d=D-1)

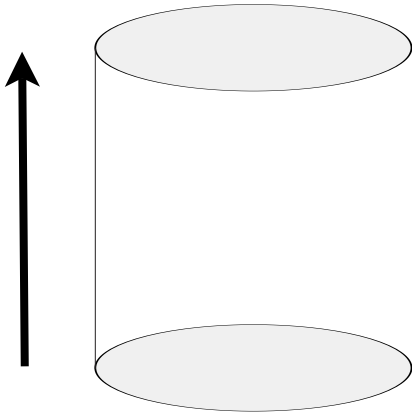
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time



(T,X): Minkowski coordinates

R: radius of curvature

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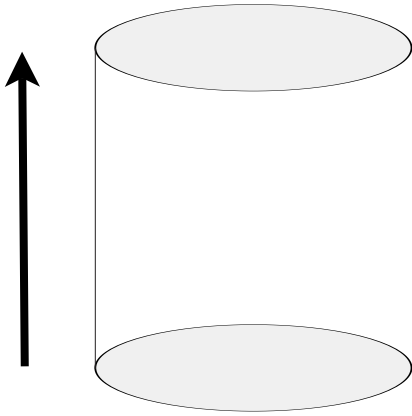
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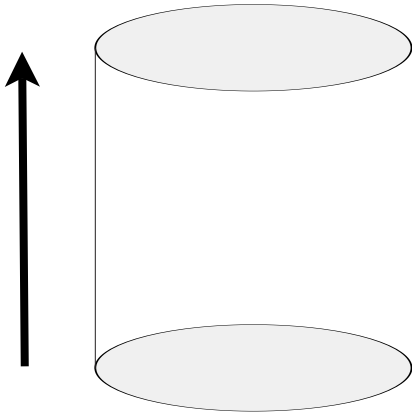
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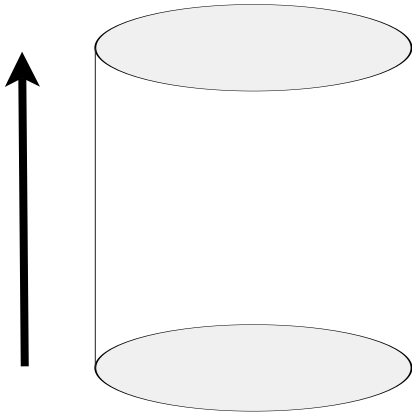
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Duality: different degrees of freedom in two different limits of the coupling $g^2 N_c$

$$g^2 N_c \gg 1$$

- Strongly coupled SYM
- Weakly coupled gravity

$$g^2 N_c \ll 1$$

- Weakly coupled SYM
- Strongly coupled gravity

String/N=4 SYM duality

Note that correspondence is expected to be valid for N=4 SYM:

- One gauge field A_μ
- Six scalars $\phi_i, i = 1, \dots, 6$
- Four fermions $\chi_k, k = 1, \dots, 4$
- Fields transform in the adjoint representation
- Conformal invariant $\beta \sim \mathcal{N} - 4 = 0$

SYM N=4 very different from QCD. Nevertheless a very good “laboratory” .

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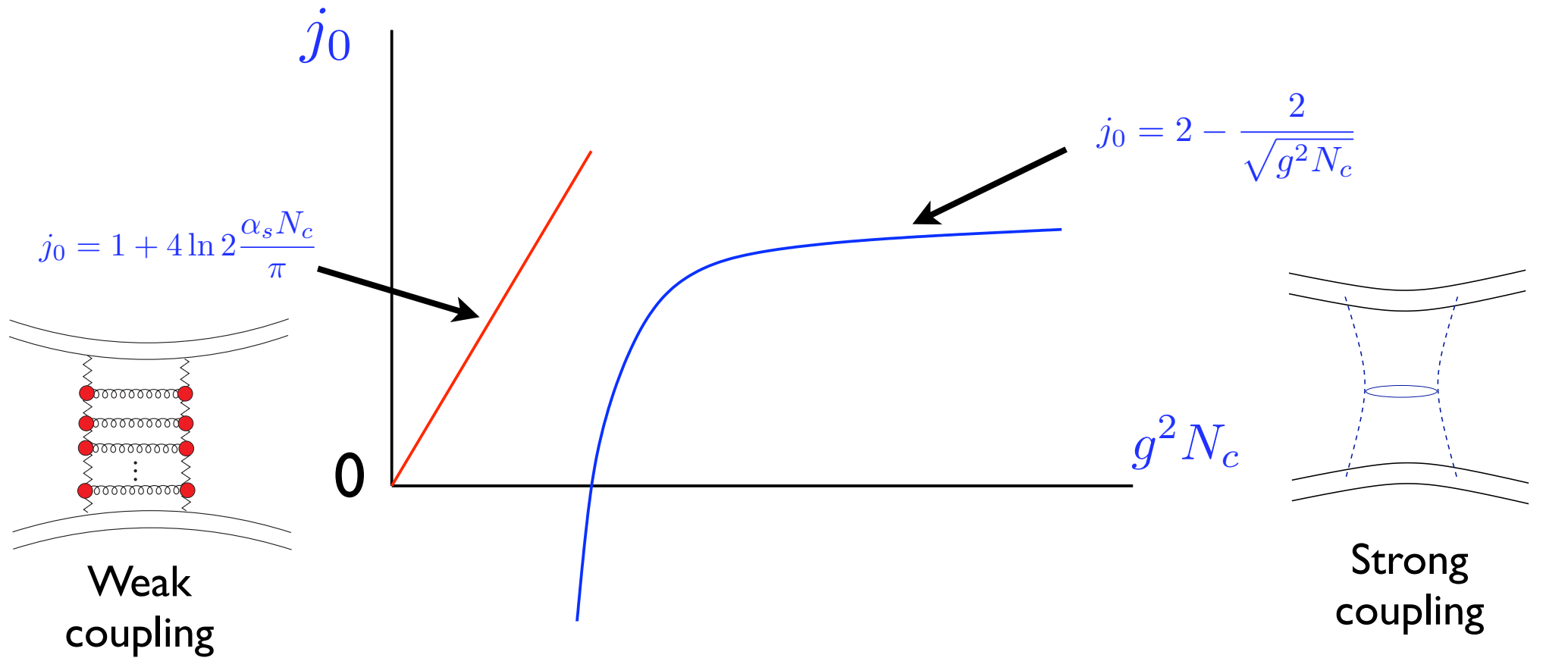
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Diffusion in transverse
(virtual) momenta



Diffusion in the fifth (radial)
dimension of AdS space

Pomeron/Graviton



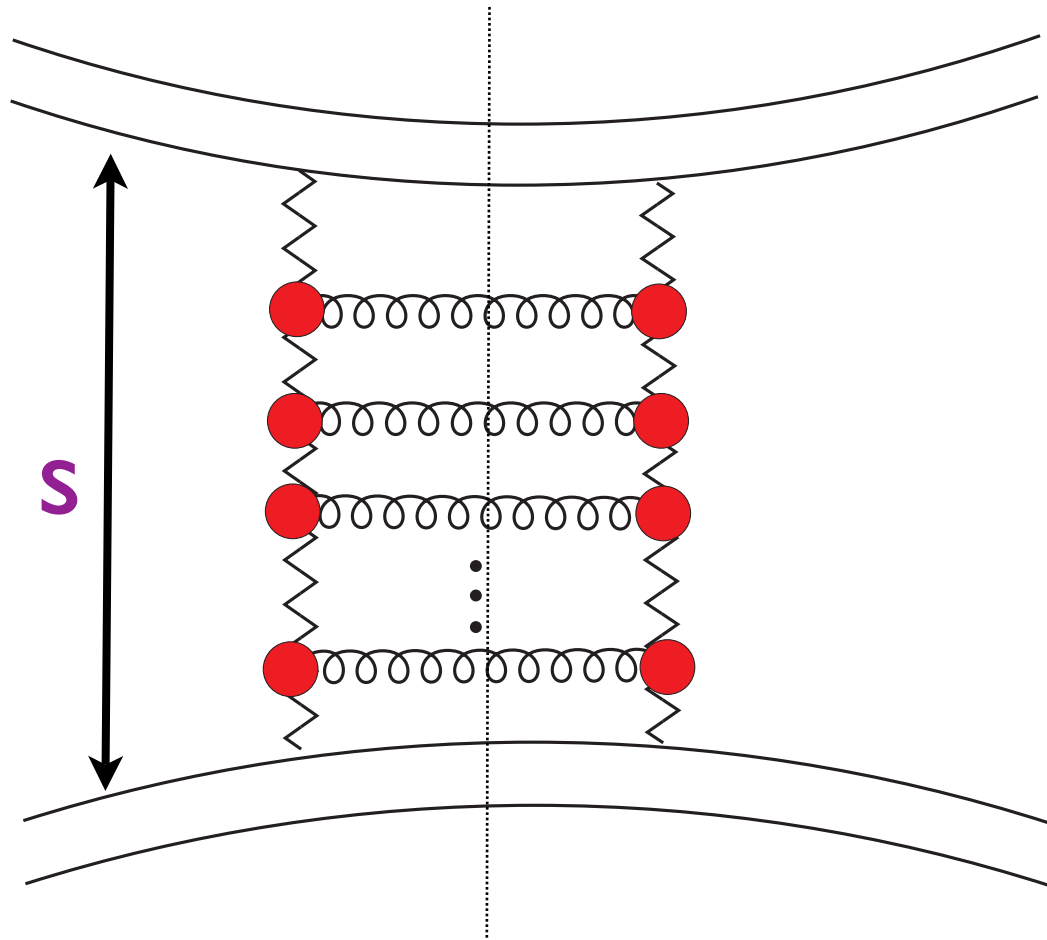
Pomeron: made out of many (reggeized) gluons. Growth of the cross section caused by dynamical effect: emission of many gluons.

Graviton: single object (closed string state). Growth of the cross section corresponds to the exchange of spin 2.

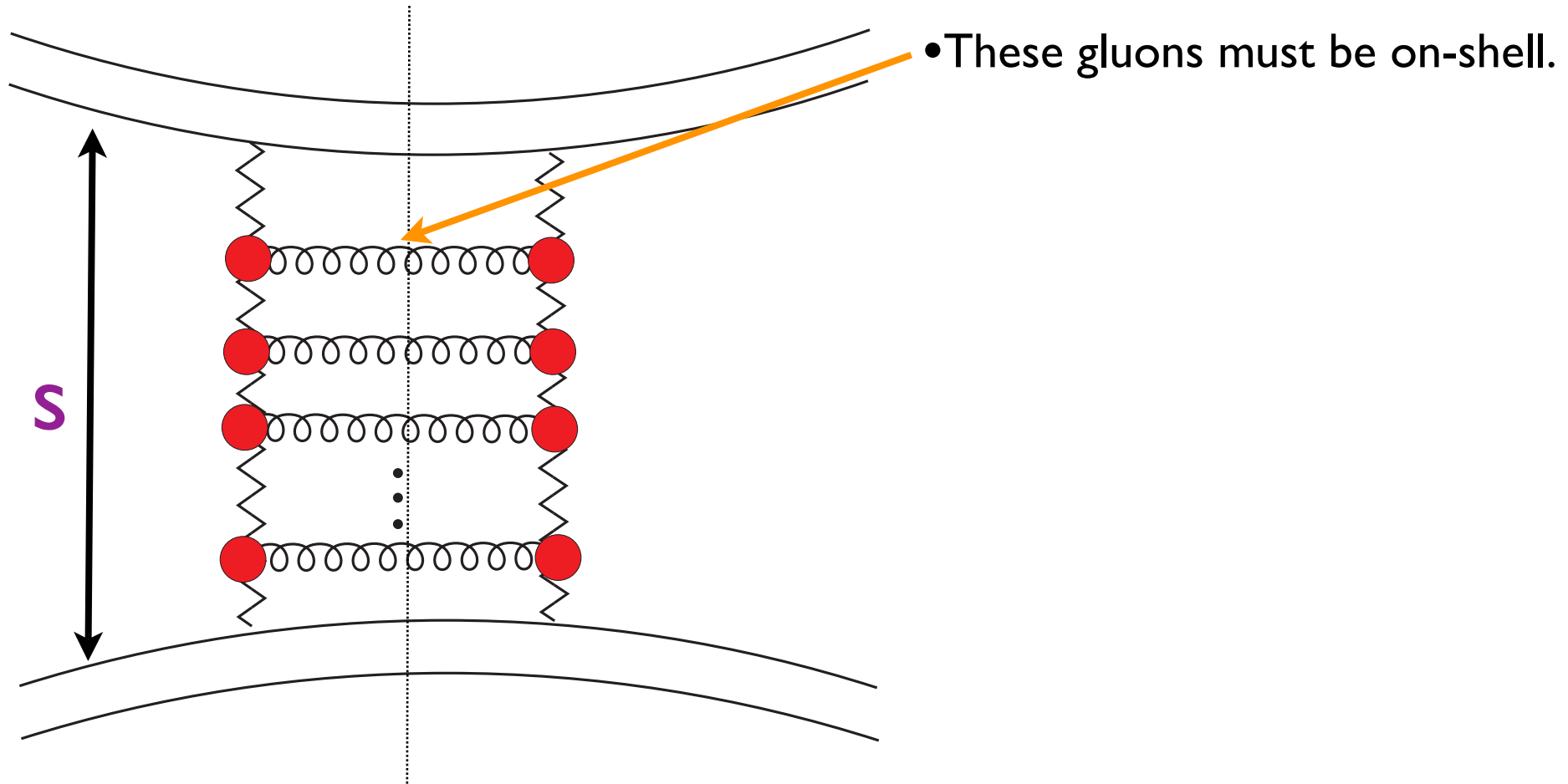
Resummation at high energies (small x)

- Next-to-leading order very large: $j_0 = 1 + 4 \ln 2 \frac{\alpha_s N_c}{\pi} \left(1 - 6.45 \frac{\alpha_s N_c}{\pi}\right)$
- Sources of large corrections:
 - Kinematical effects, energy momentum conservation. ← Common to QCD and SYM
 - Running of the coupling. QCD only
- Other corrections: quarks in the evolution.
- Need to take more than next-to-leading order:
all orders.

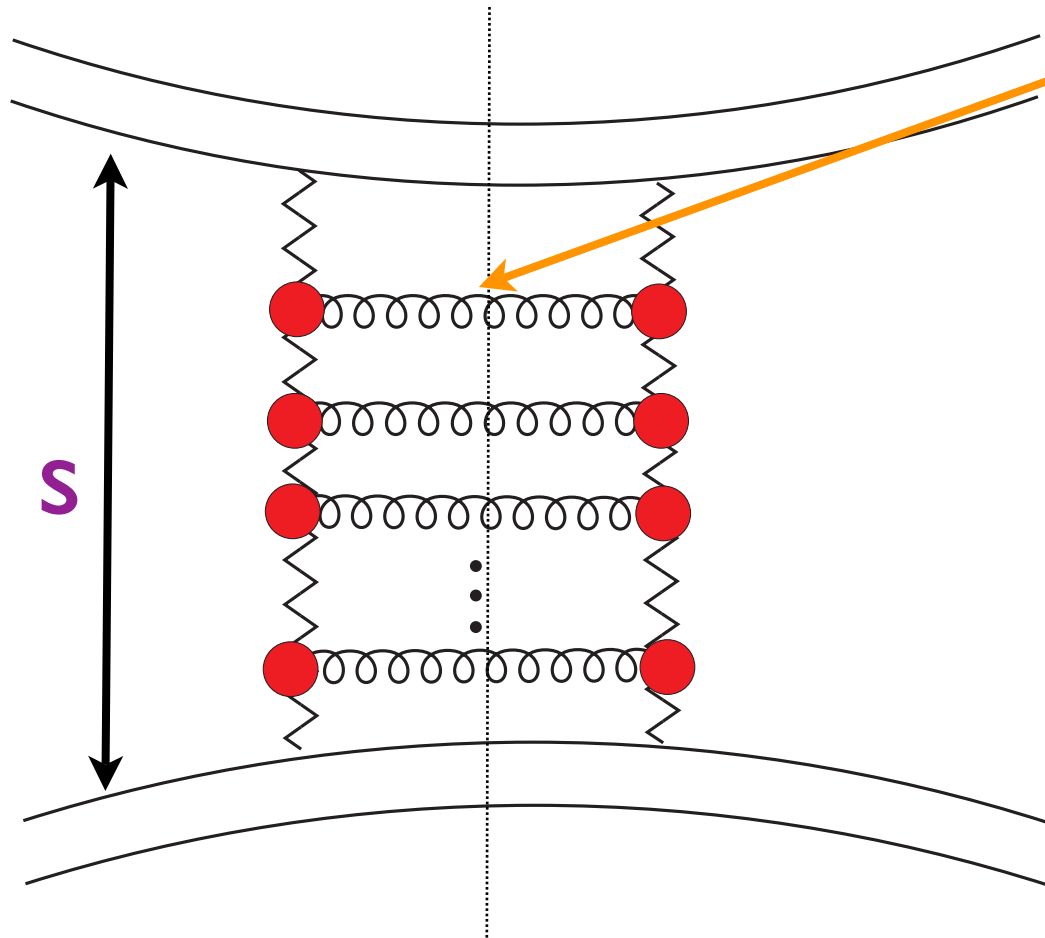
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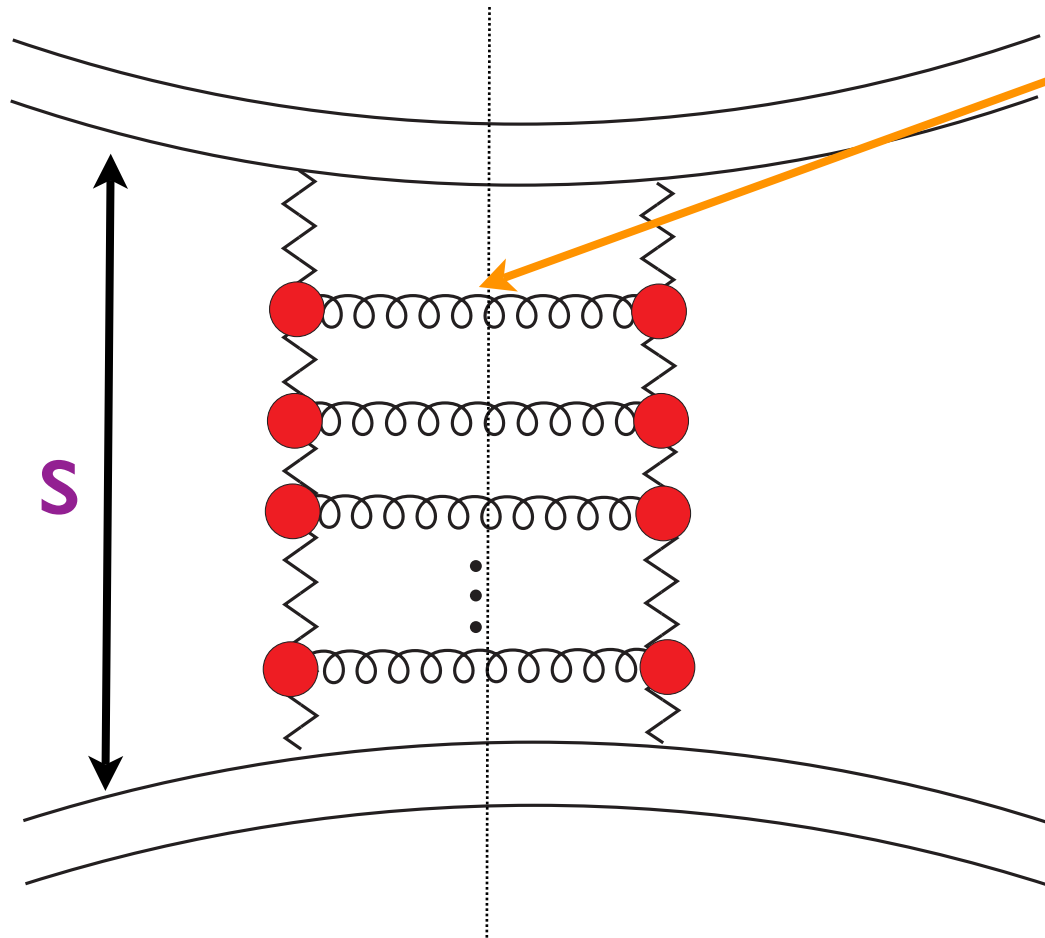


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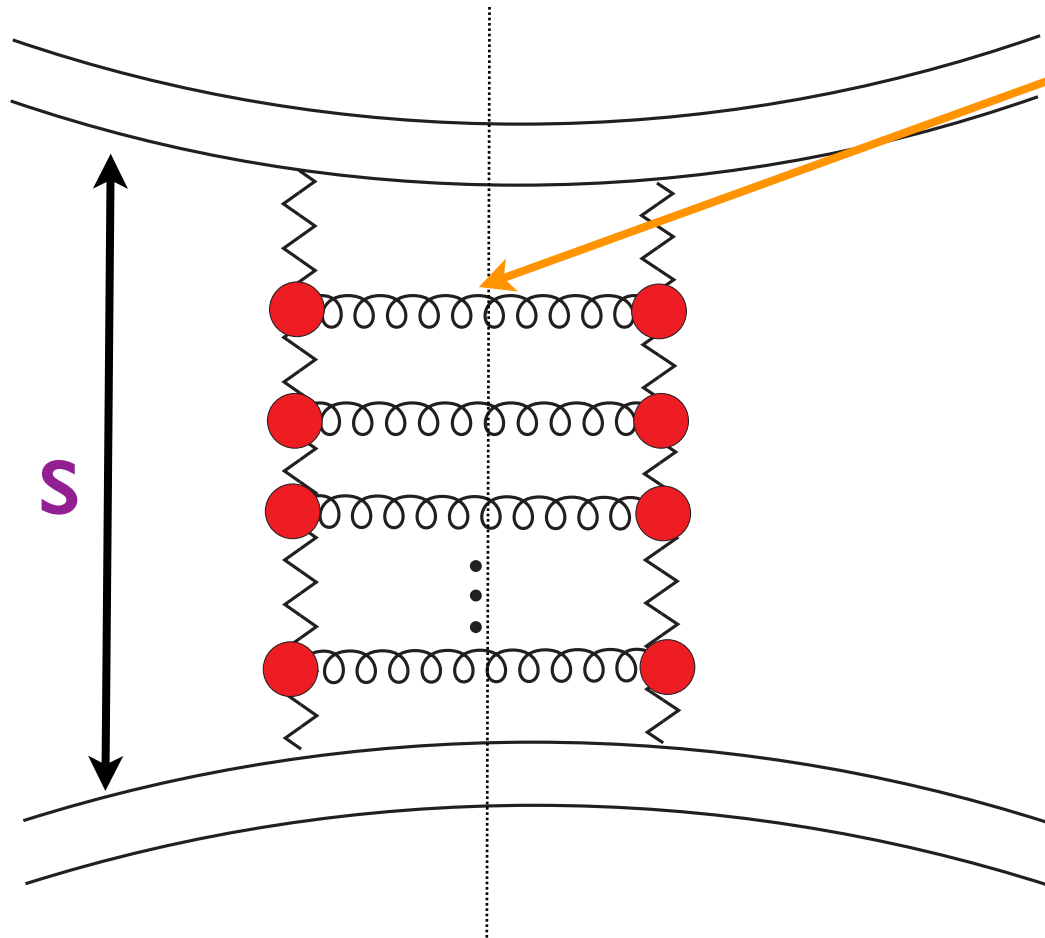
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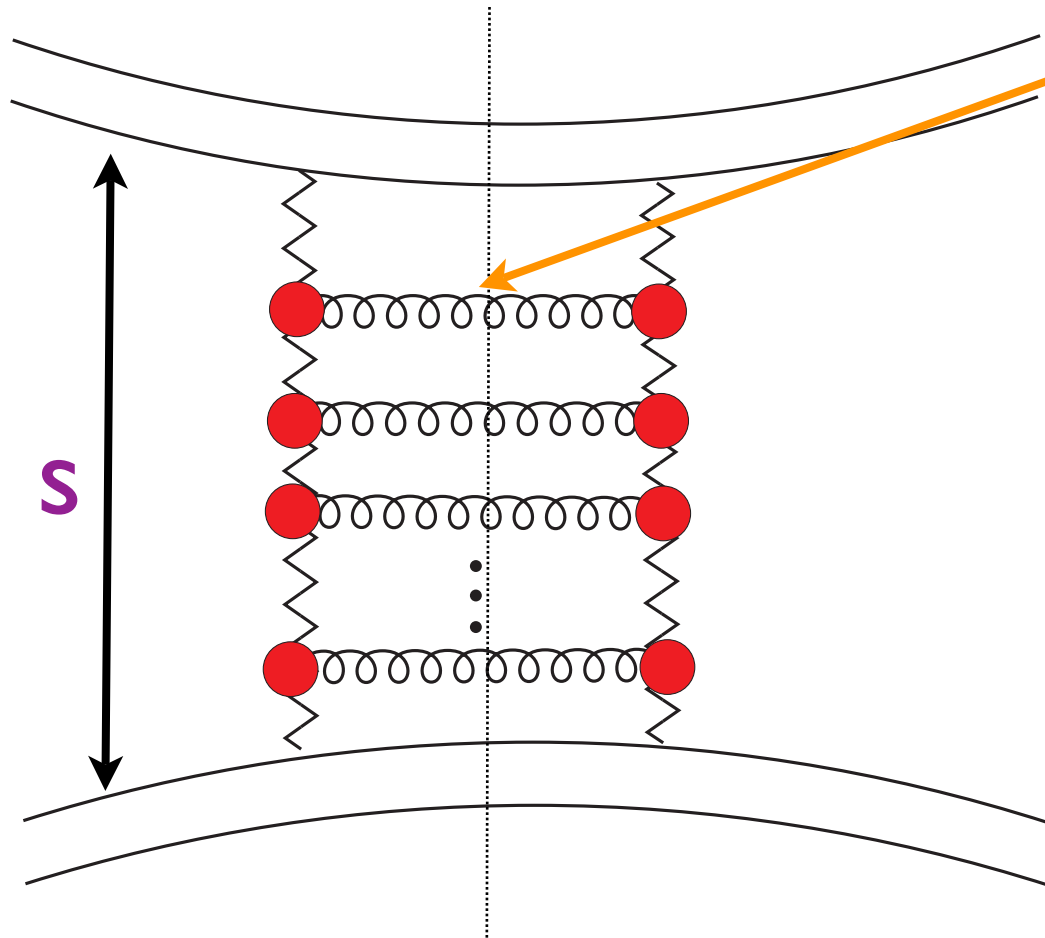
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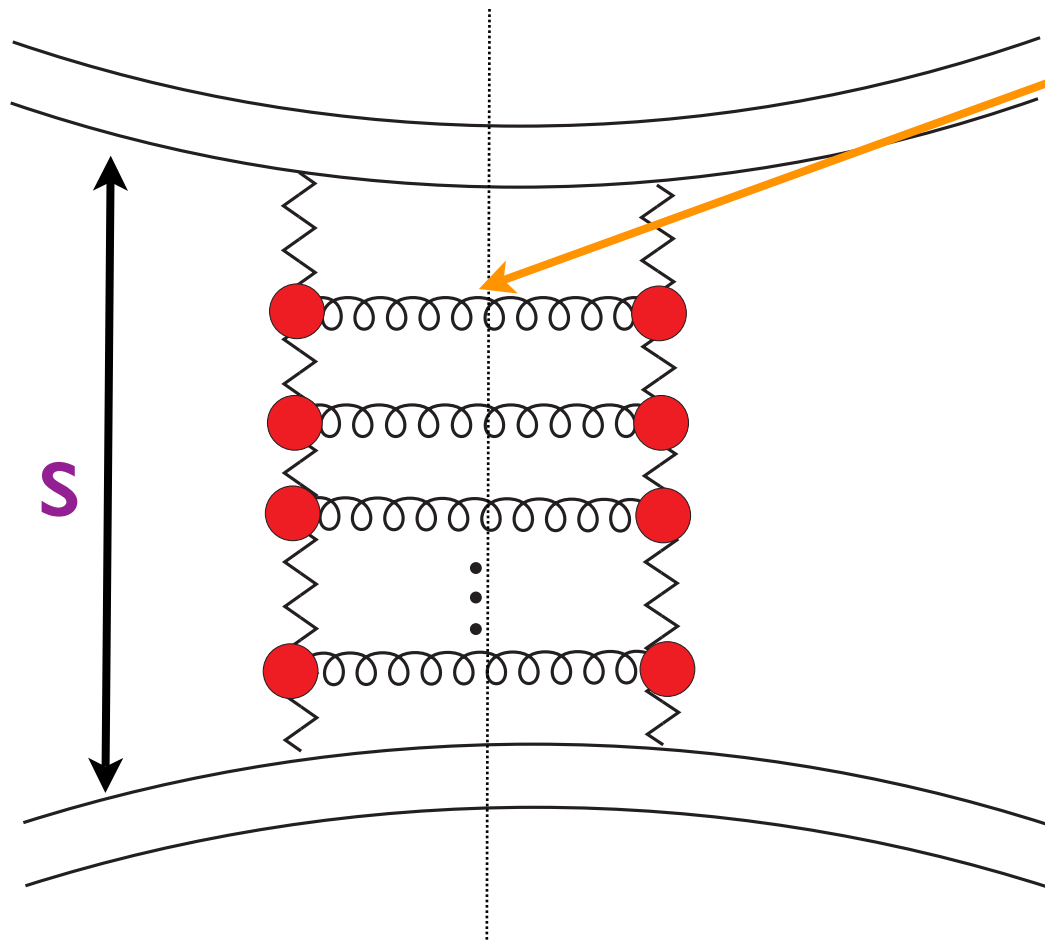
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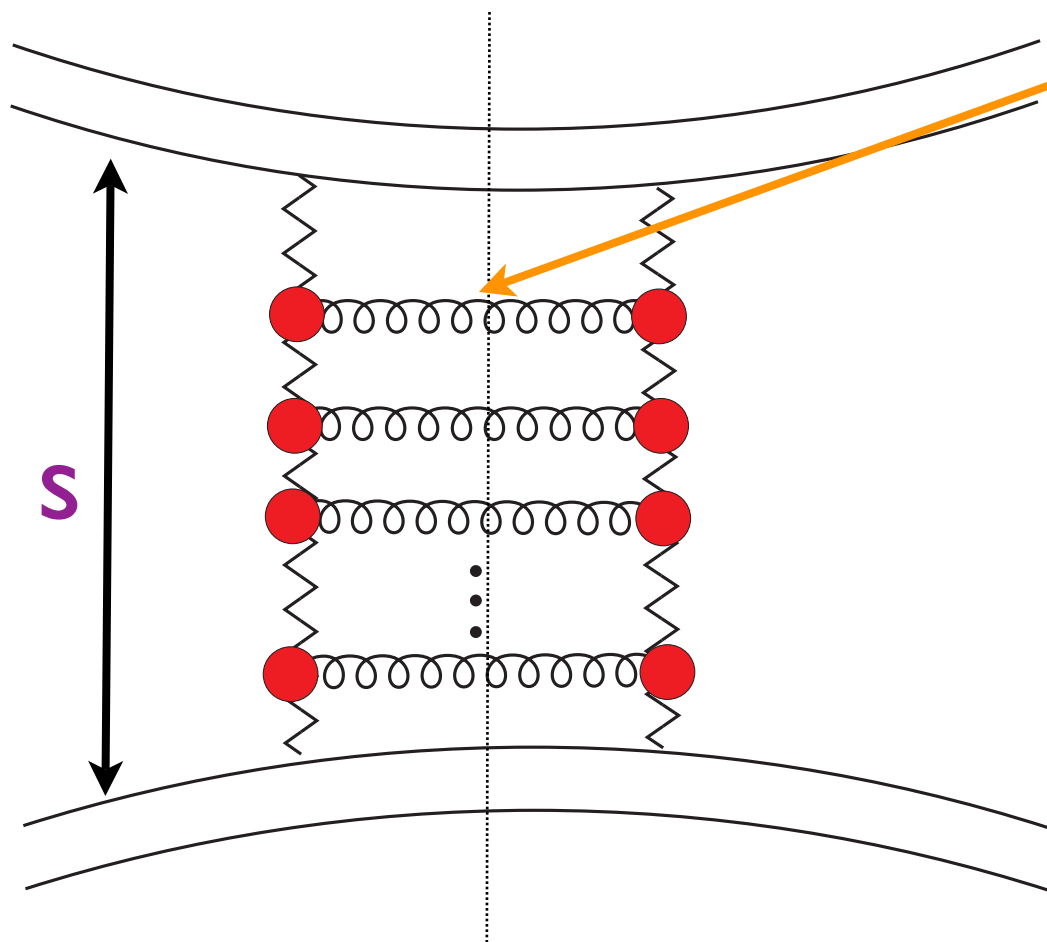


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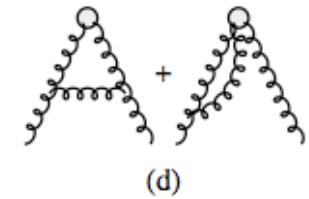
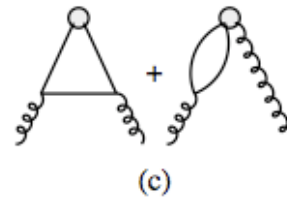
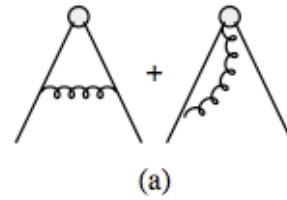
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Impose constraints to satisfy energy-momentum sum rule

(

A note on anomalous dimensions in QCD

In standard operator product expansion
approach to DIS
evaluate anomalous dimensions



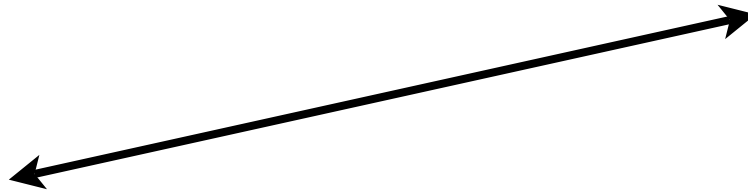
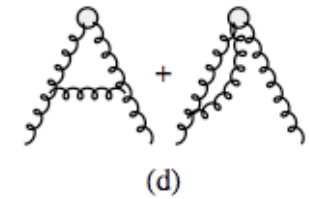
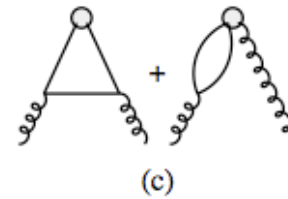
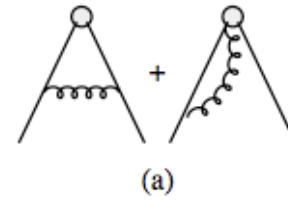
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RGE:

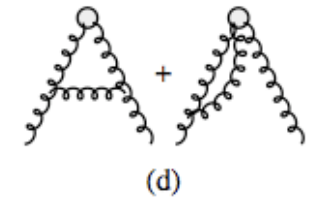
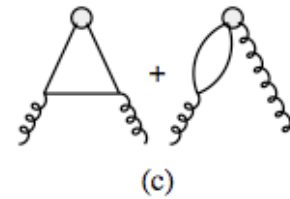
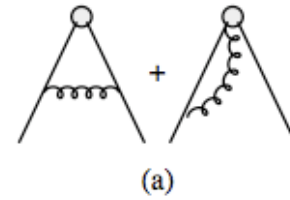
$$[\mathcal{D}\delta_{ab} - \gamma_{ab}^{(j)}] C_b^j(g, \mu, -q^2) = 0$$

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A note on anomalous dimensions in QCD

In standard operator product expansion
 approach to DIS
 evaluate anomalous dimensions



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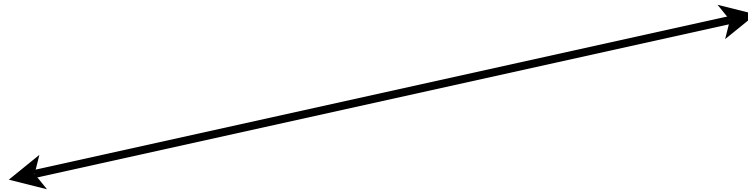
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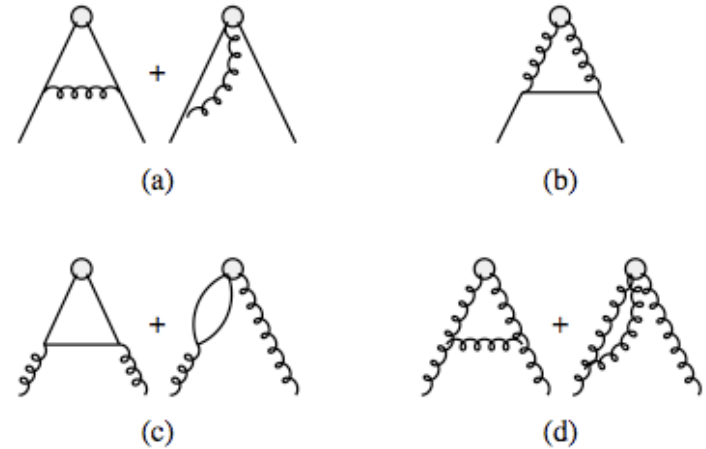
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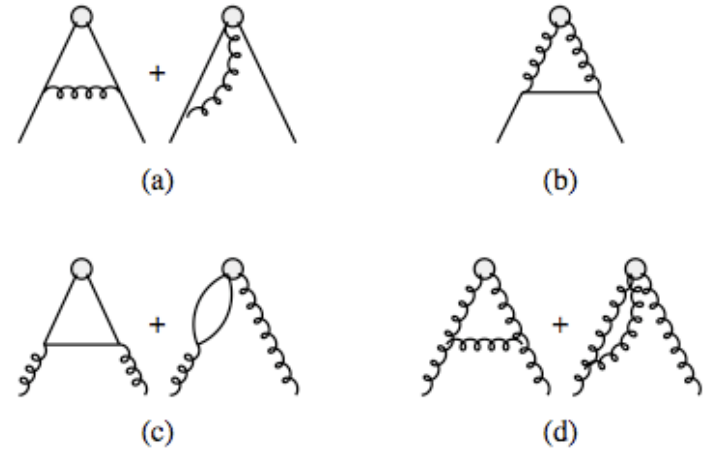
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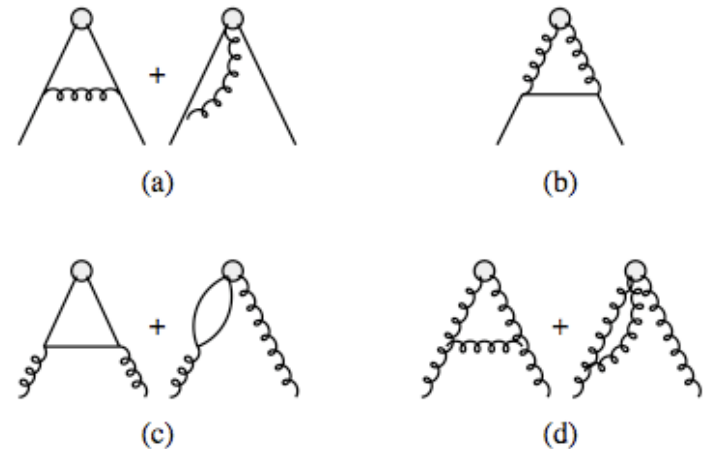
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Satisfied at each order of the perturbation theory



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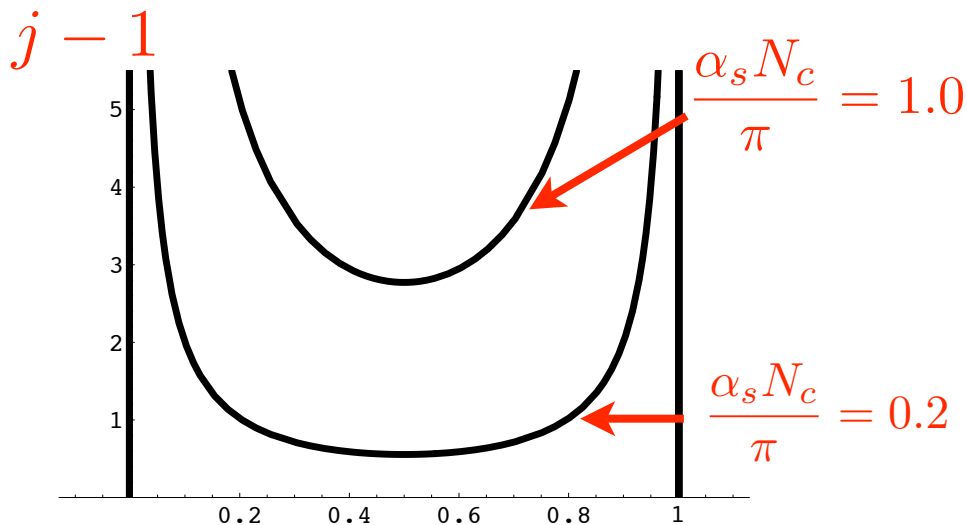
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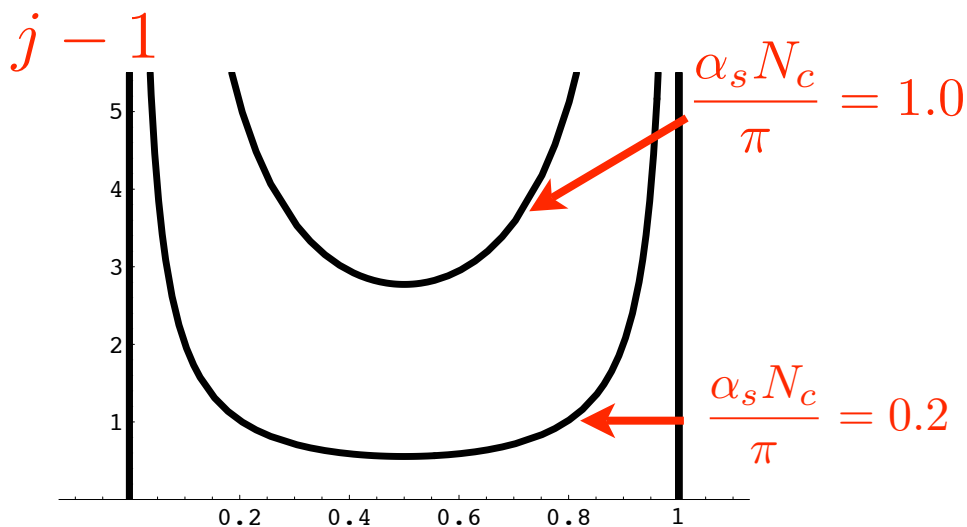
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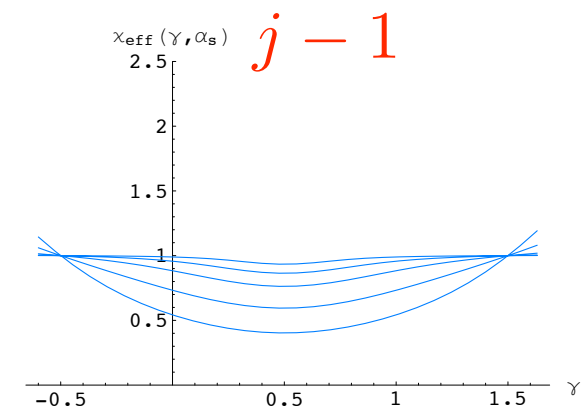
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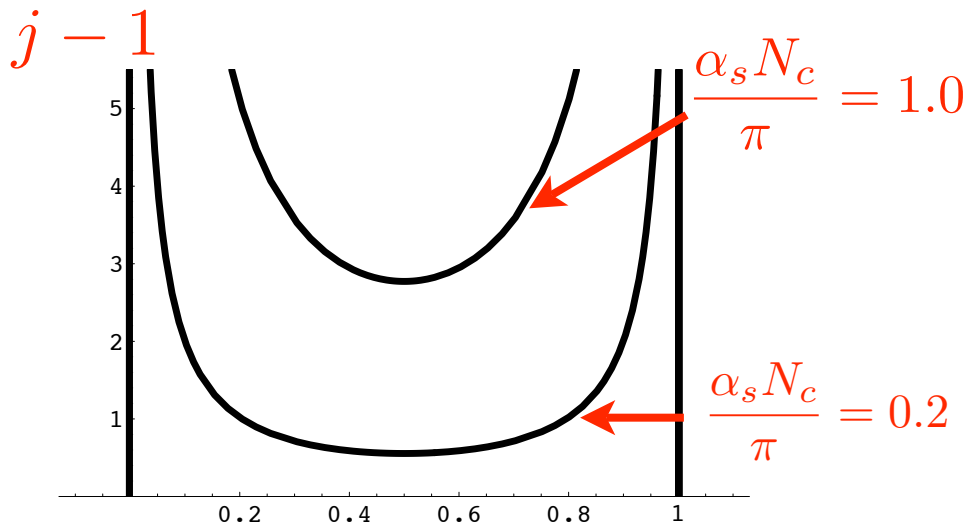
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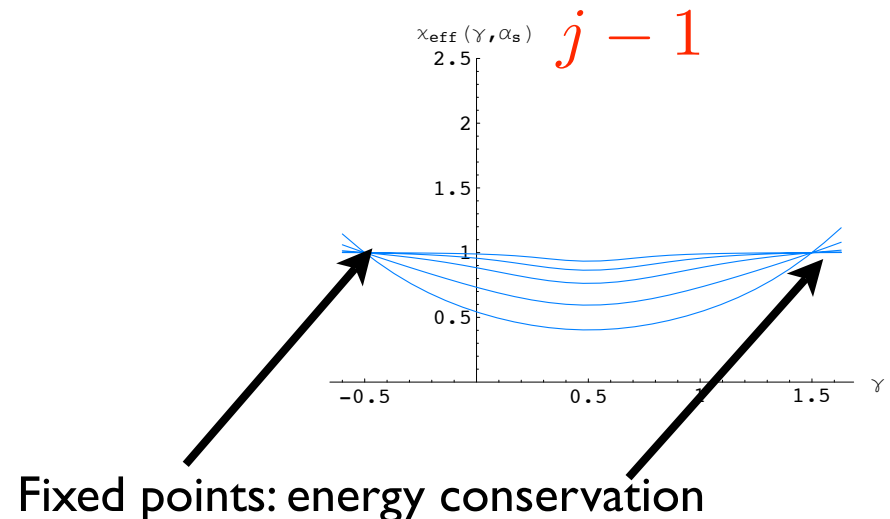
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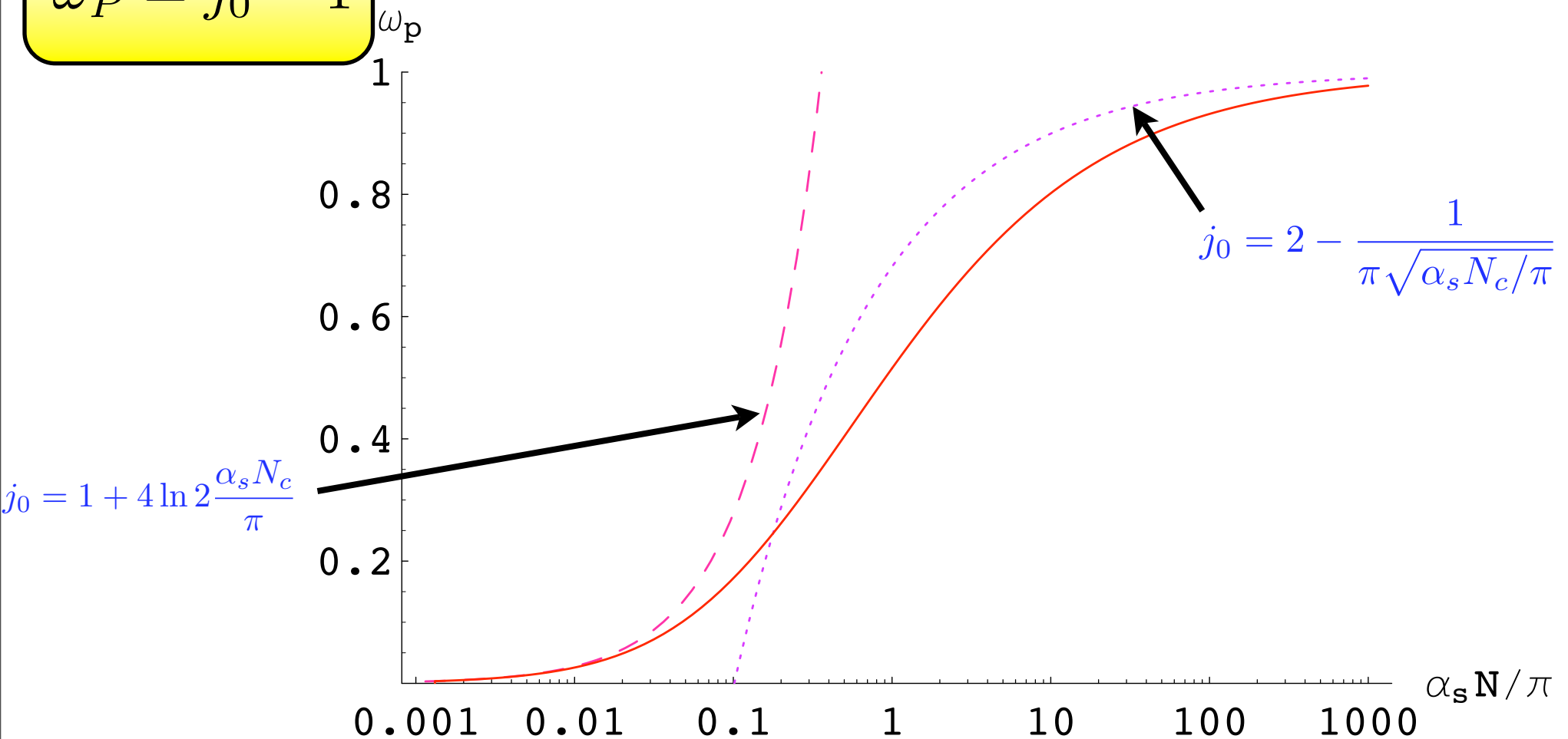


After:



Intercept in the resummed model

$$\omega_P = j_0 - 1$$



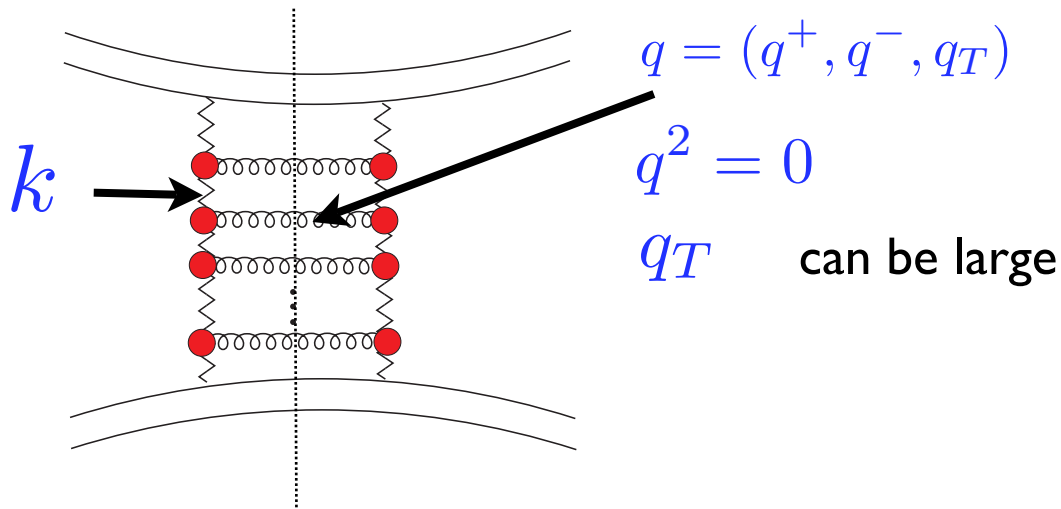
Note the logarithmic horizontal axis

Cross section: $\sigma \sim s^{j_0 - 1}$

Vanishing diffusion and soft gluons

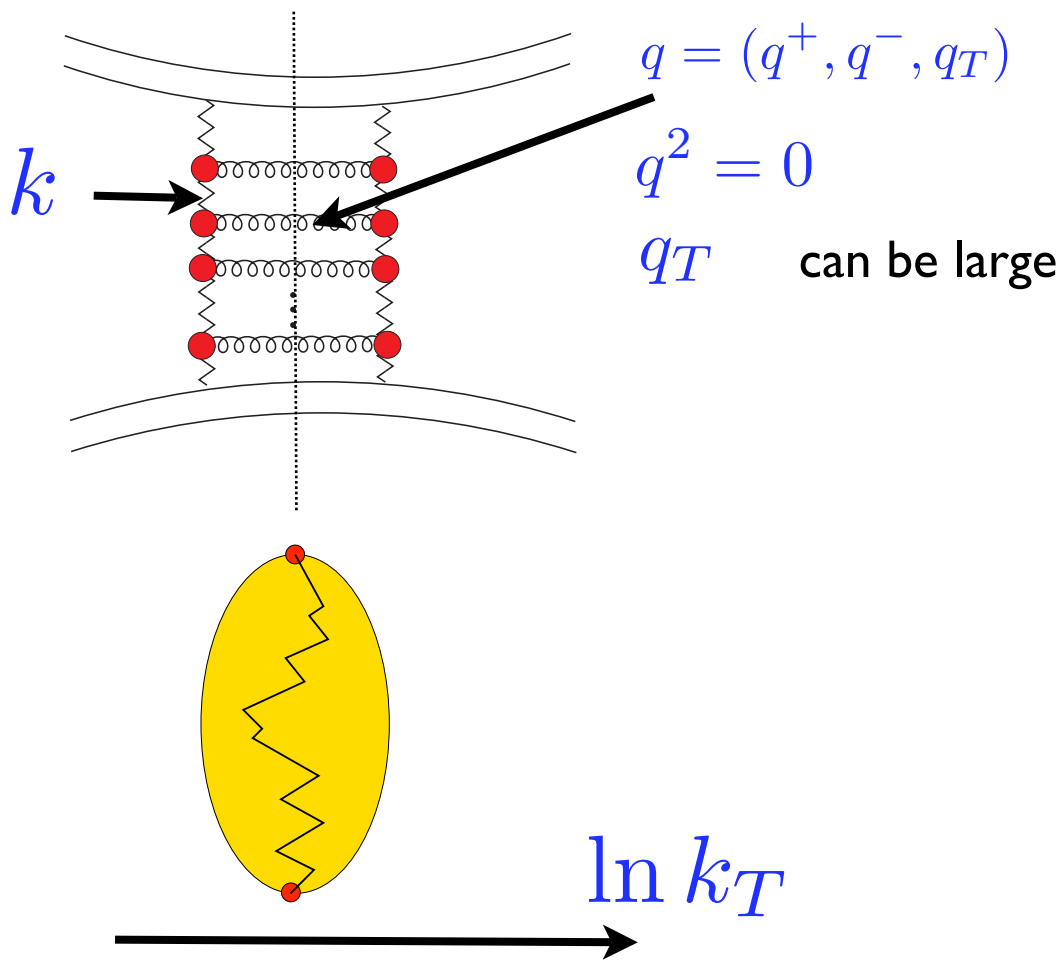
Vanishing diffusion and soft gluons

Small coupling



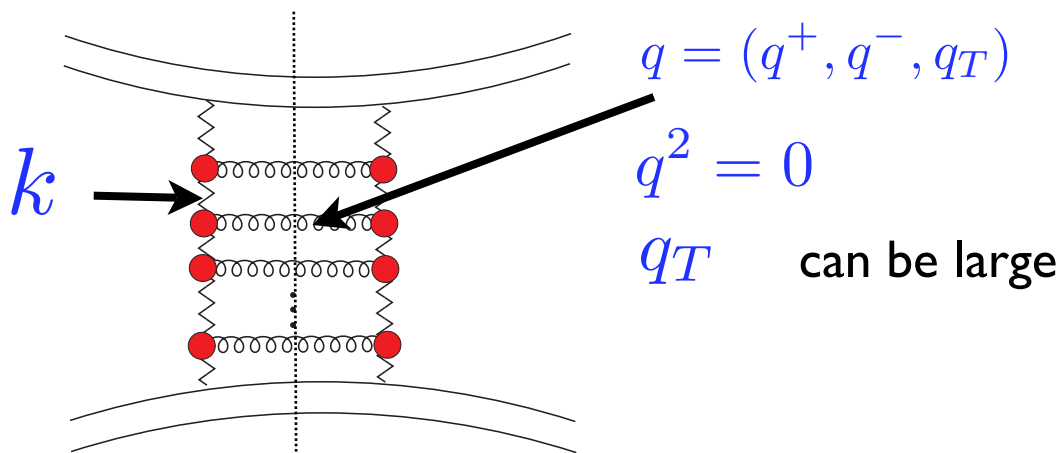
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$$q = (q^+, q^-, q_T)$$

$$q^2 = 0$$

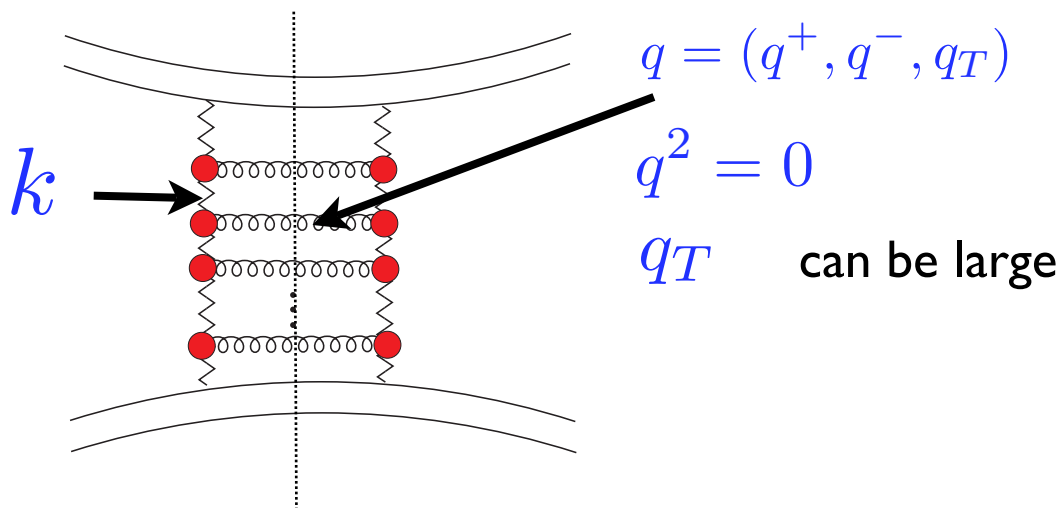
q_T can be large

$$\ln k_T$$

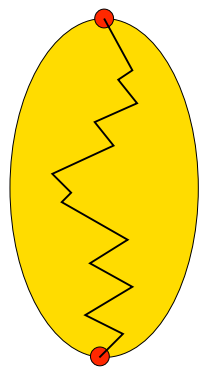
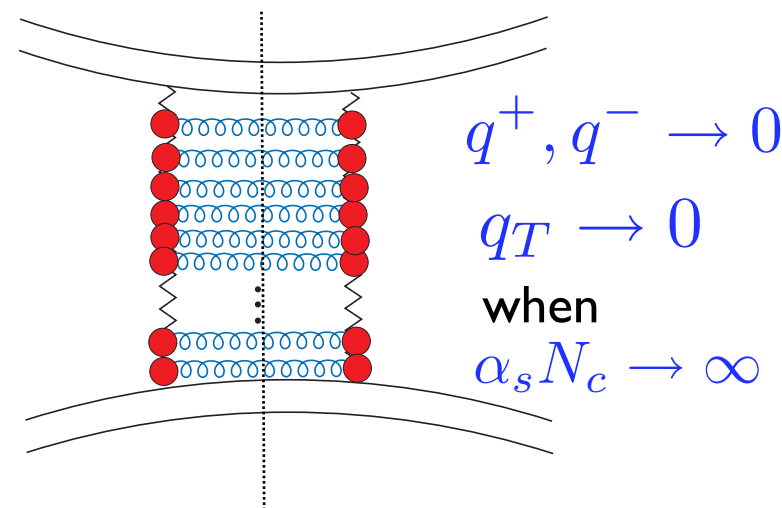
$$\chi''\left(\frac{\alpha_s N_C}{\pi} \ll 1\right) \text{ large}$$

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Large coupling

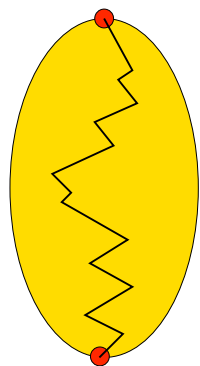
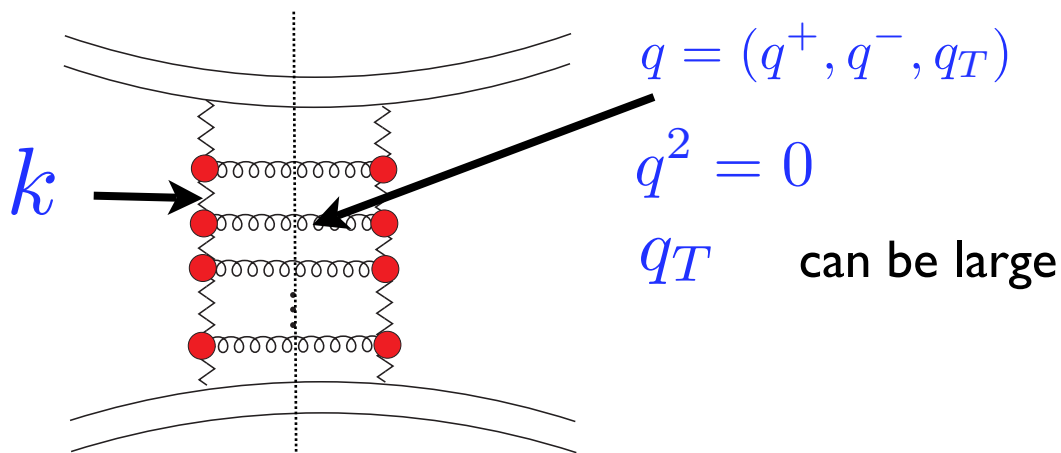


$\ln k_T$

$\chi''\left(\frac{\alpha_s N_c}{\pi} \ll 1\right)$ large

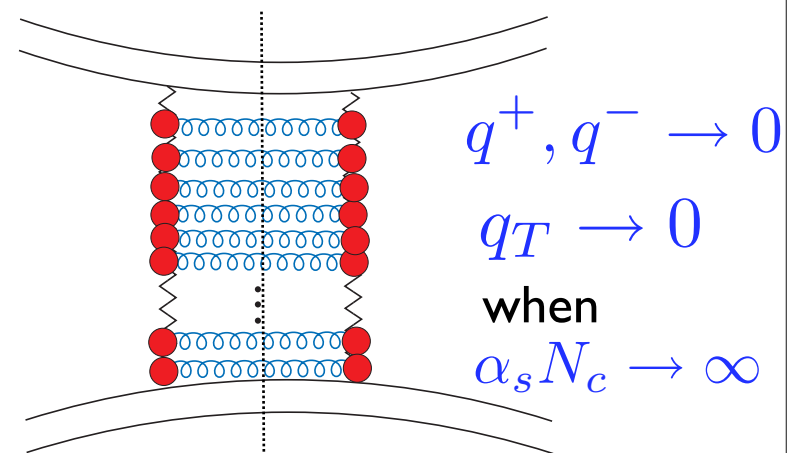
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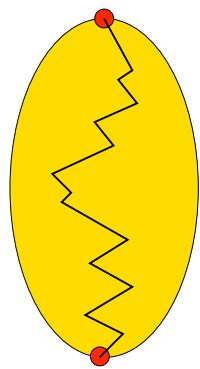
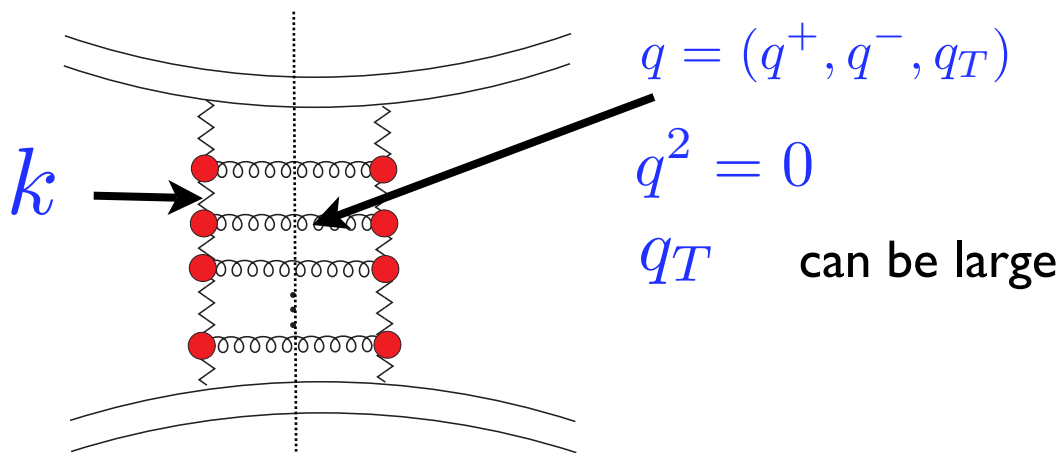
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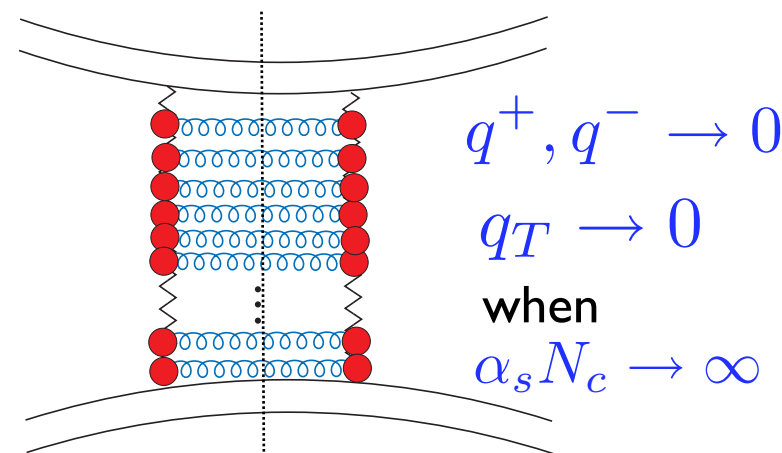
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Large coupling



$\chi''\left(\frac{\alpha_s N_C}{\pi} \gg 1\right)$ small

Summary

- Universal growth of hadronic cross sections.
- In QCD Pomeron: compound state of gluons, dominates the high energy behavior of cross sections.
- In string(gravity) theory: graviton dominates at high energies.
- Simple kinematic constraints lead to resummation: weak to strong coupling interpolation.
- In real QCD situation more complicated: running coupling, multi-Pomeron/graviton exchanges(interactions).