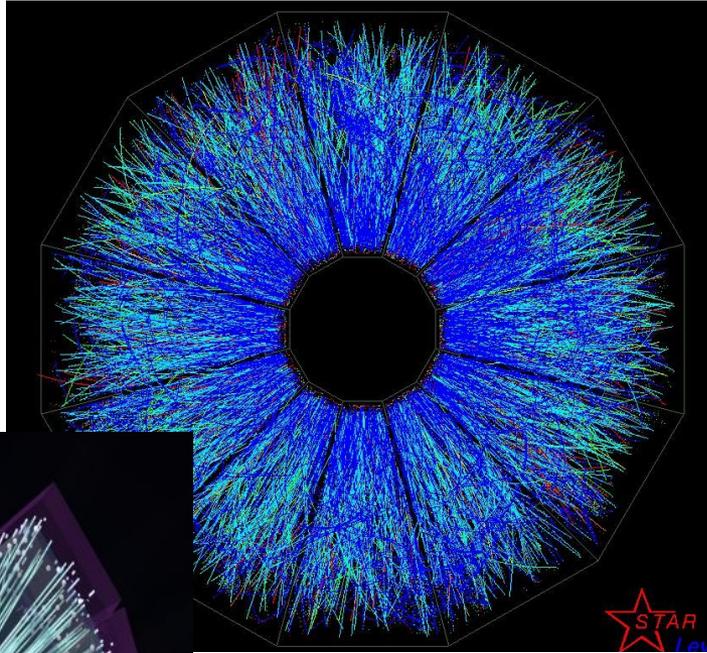


# Nature's First Liquid: The Quark Gluon Plasma

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*Barbara Jacak*

*Krakow School  
of  
Theoretical Physics*

*May 28, 2006*



OCTOBER 2003

# PHYSICS TODAY

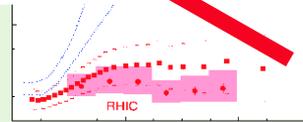
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LETTERS

# SCIENTIFIC AMERICAN

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~~Do the nucleons melt  
into quarks and gluons?~~



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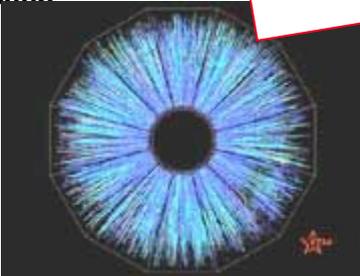
Stopping  
Alzheimer's  
Birth of  
the Amazon  
Future



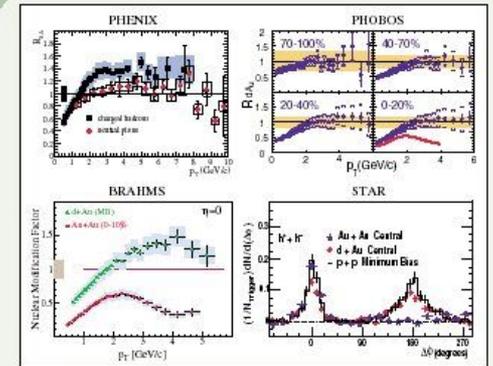
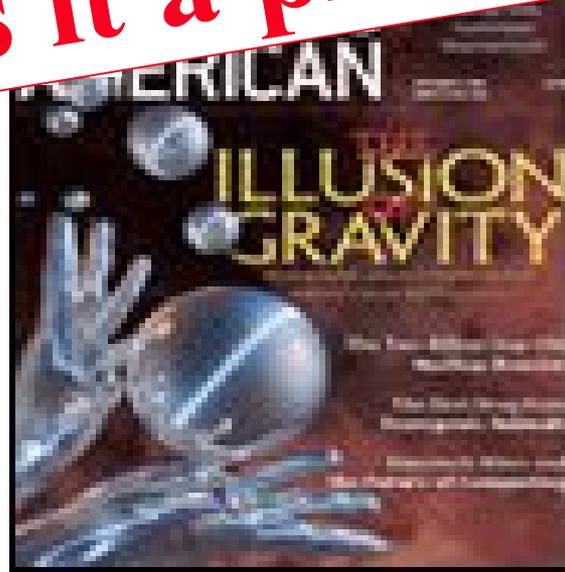
Articles published week ending  
15 AUGUST 2003  
Volume 91, Number 7

~~Is it a plasma?~~

● Little Bang



*Discover*  
*January 2002*



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APS Published by The American Physical Society

A 3-D digital camera tracks the paths of the thousands of new subatomic particles created when two gold ions are smashed together in a collider at Brookhaven National Laboratory. Photograph courtesy of Brookhaven National Laboratory

# outline

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- **introductory remarks**
- **transmission of color-charged probes by the plasma**
- **collective effects**
  - **applying hydrodynamics**
- **heavy quark probes**
  - **diffusion and viscosity**
- **strongly-coupled plasmas**
- **(color) screening**
  - **the ever-surprising  $J/\psi$**

# what is a plasma?

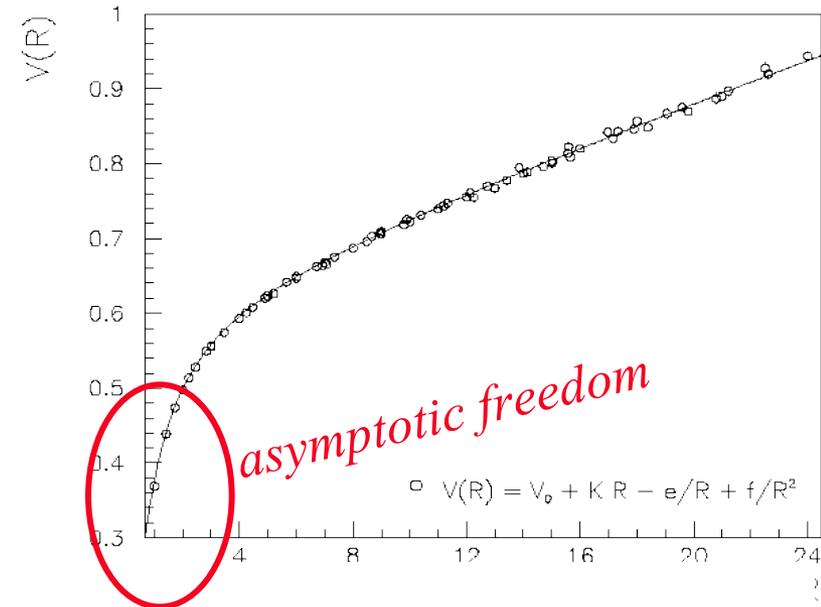
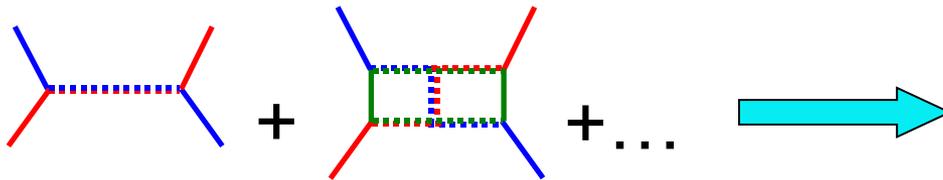
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- 4<sup>th</sup> state of matter (after solid, liquid and gas)
- a plasma is:
  - ionized gas which is macroscopically neutral
  - exhibits collective effects
- interactions among charges of multiple particles
  - spreads charge out into characteristic (Debye) length,  $\lambda_D$
  - multiple particles inside this length
  - they screen each other
  - plasma size  $> \lambda_D$
- “normal” plasmas are electromagnetic (e + ions)
  - quark-gluon plasma interacts via strong interaction
  - color forces rather than EM
  - exchanged particles: g instead of  $\gamma$

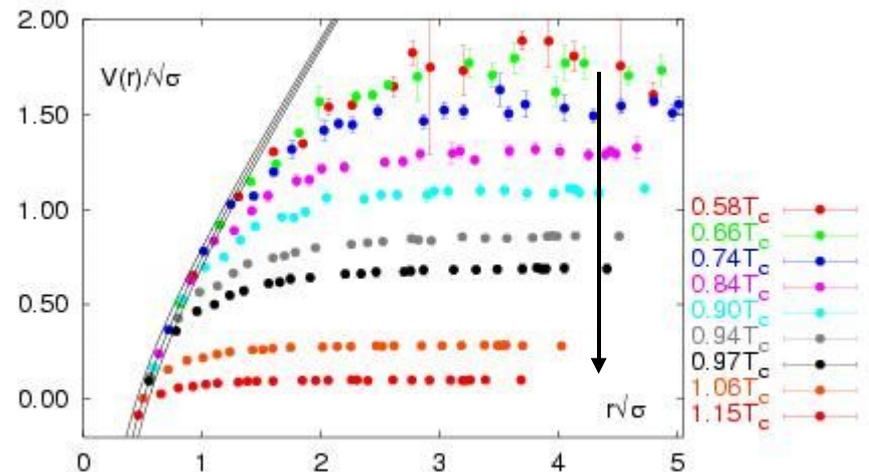
# QCD: expect a phase transition

gluons carry color charge  $\Rightarrow$  gluons interact among themselves  
theory is non-abelian

special properties at large distance:  
confinement of quarks in hadrons

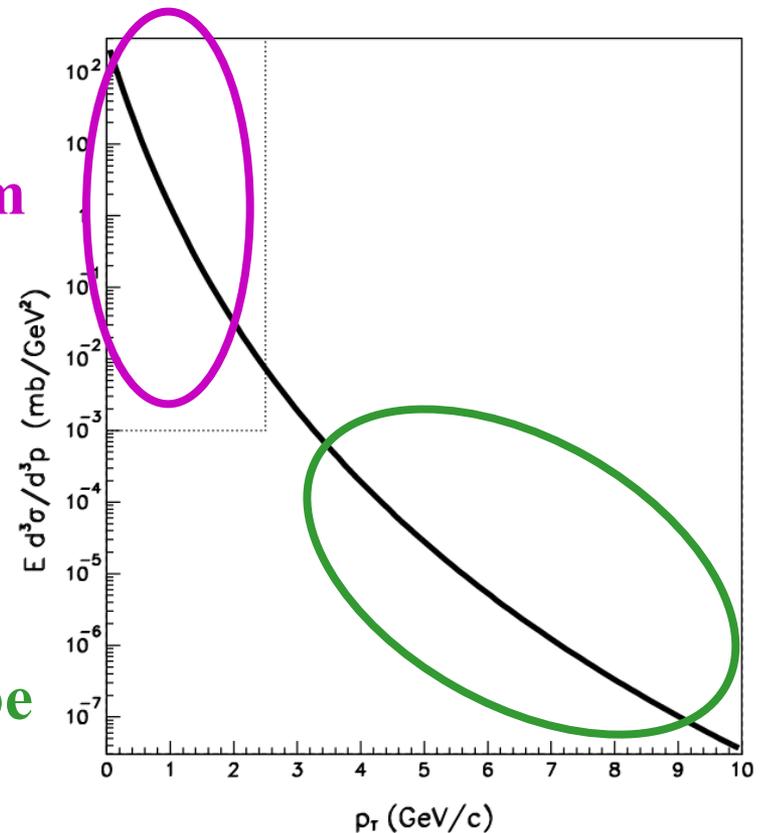


*At high temperature and density:*  
force is screened by produced  
color-charges  
expect transition to *gas of free  
quarks and gluons*

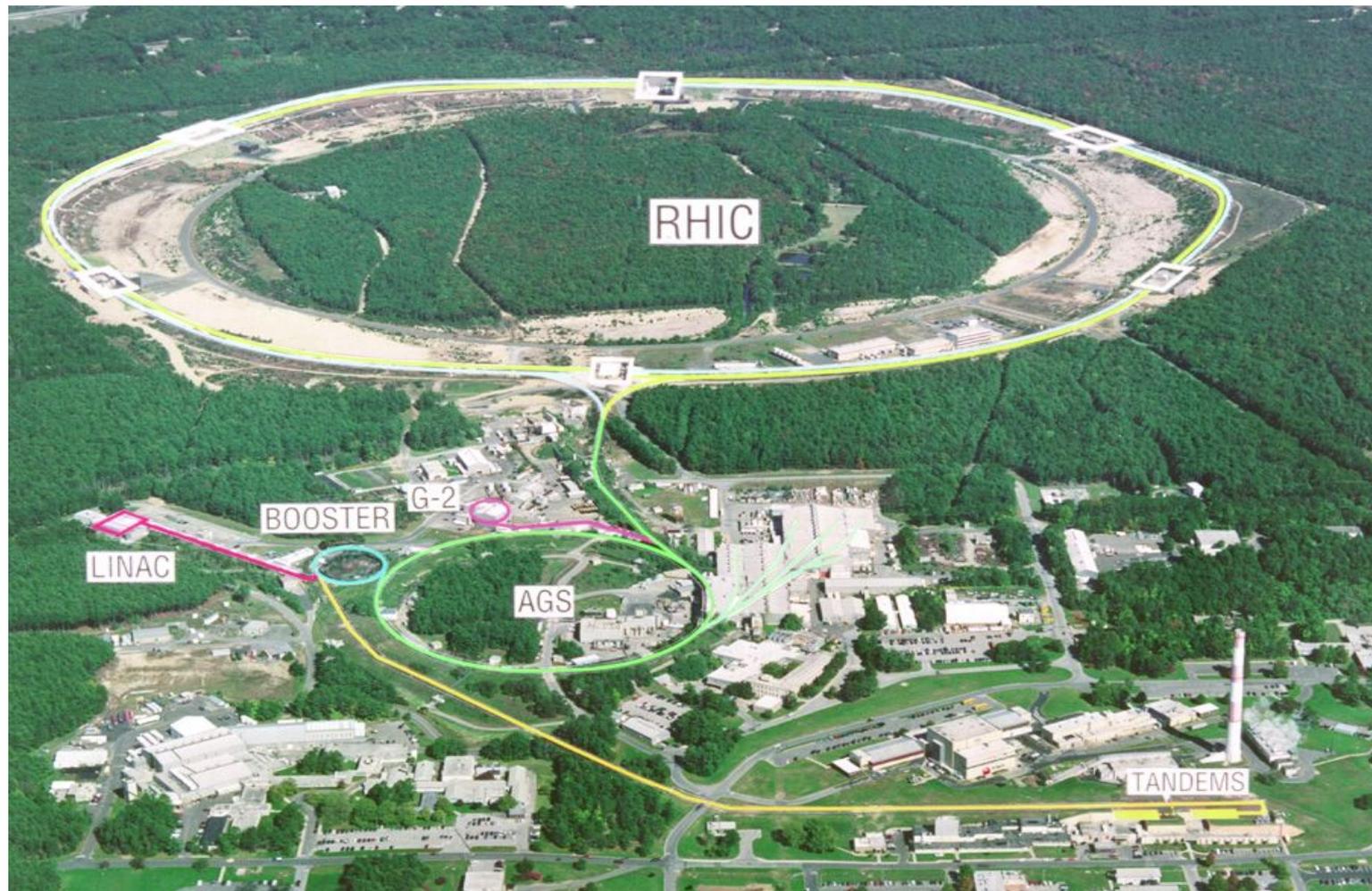


# look at radiated & “probe” particles

- as a function of transverse momentum  
 $p_T = p \sin \theta$  (with respect to beam direction)  
 $90^\circ$  is where the action is (max T,  $\rho$ )  
midway between the two beams!
- $p_T < 1.5 \text{ GeV}/c$   
“thermal” particles  
radiated from bulk of the medium  
internal plasma probes
- $p_T > 3 \text{ GeV}/c$   
jets (hard scattered q or g)  
heavy quarks, direct photons  
produced early → “external” probe



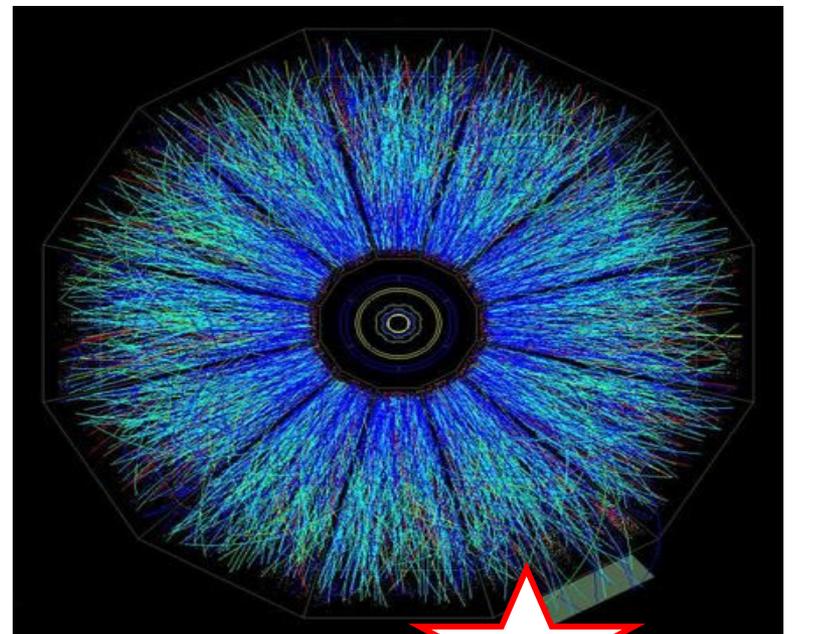
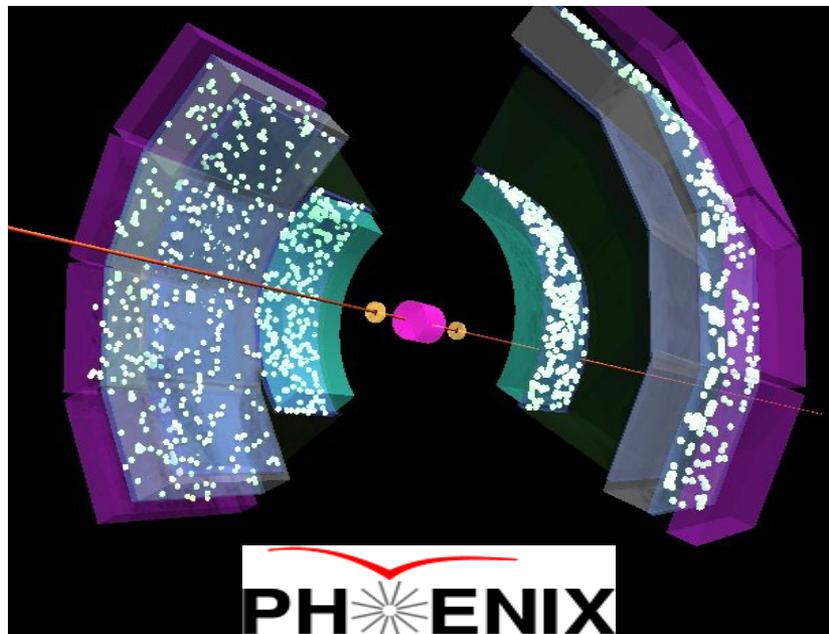
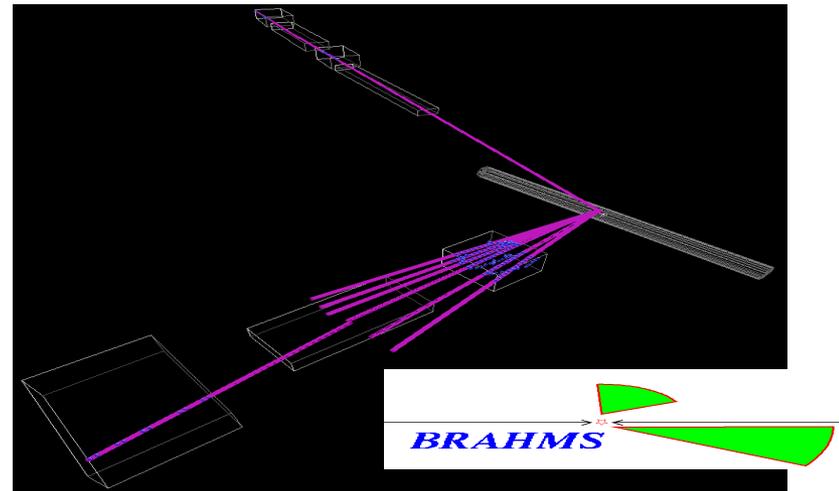
# RHIC at Brookhaven National Laboratory



**Collide Au + Au ions for maximum volume**

**$\sqrt{s} = 200 \text{ GeV/nucleon pair}$ , p+p and d+A to compare**

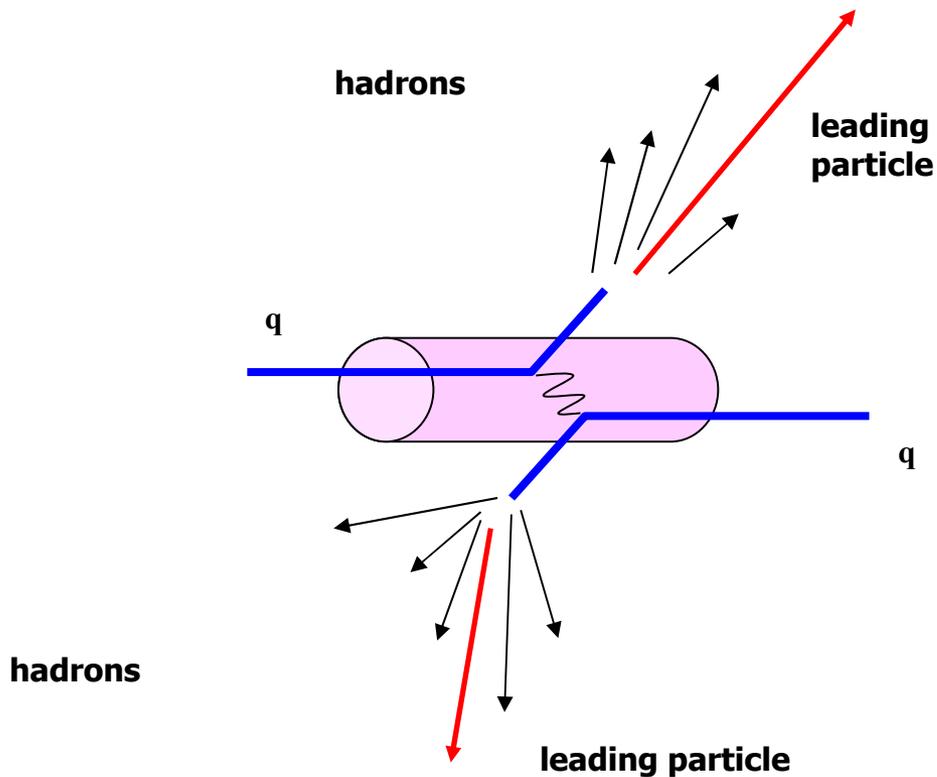
# 4 complementary experiments



# a common plasma technique

## transmission of probes which interact with plasma for QGP: fast $g$ and light quarks

schematic view of jet production

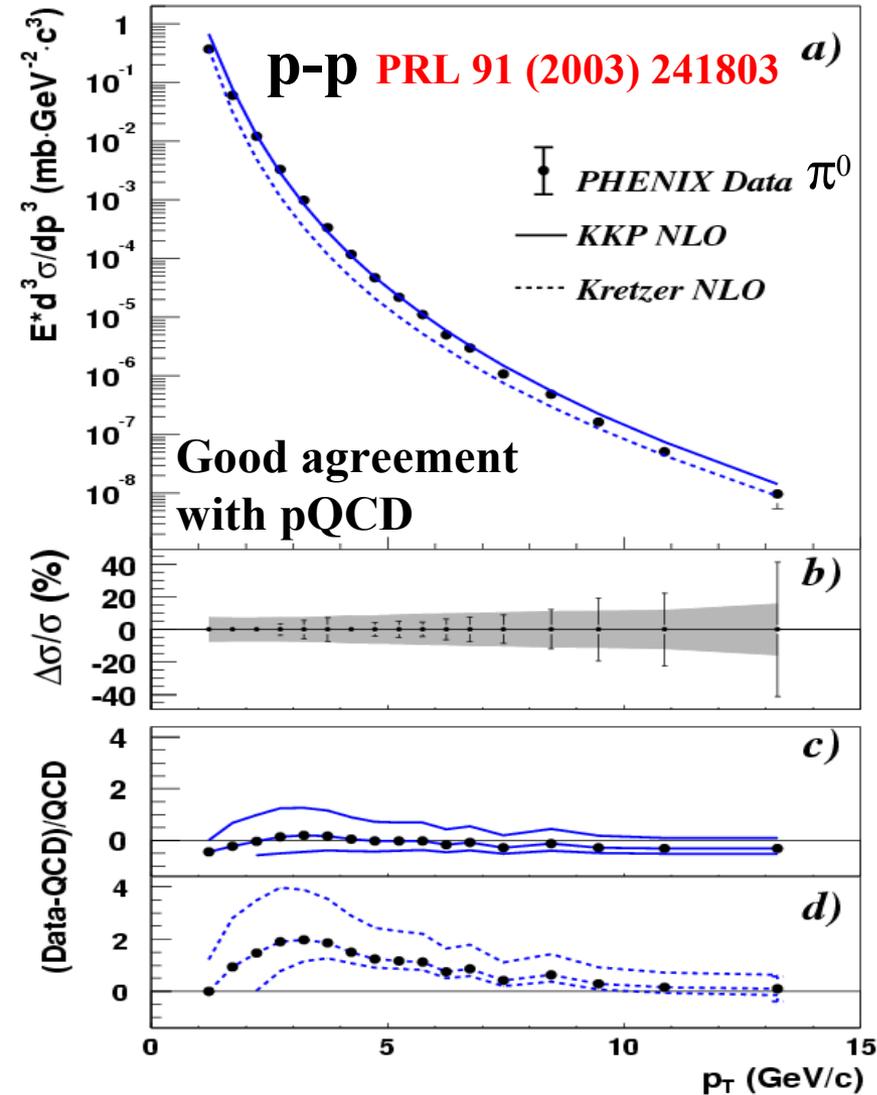


calculate probe  
rate & distribution  
with pQCD

look for modification  
by medium

d+Au collisions  
provide the control

# step 1: benchmark in $p+p$ collisions

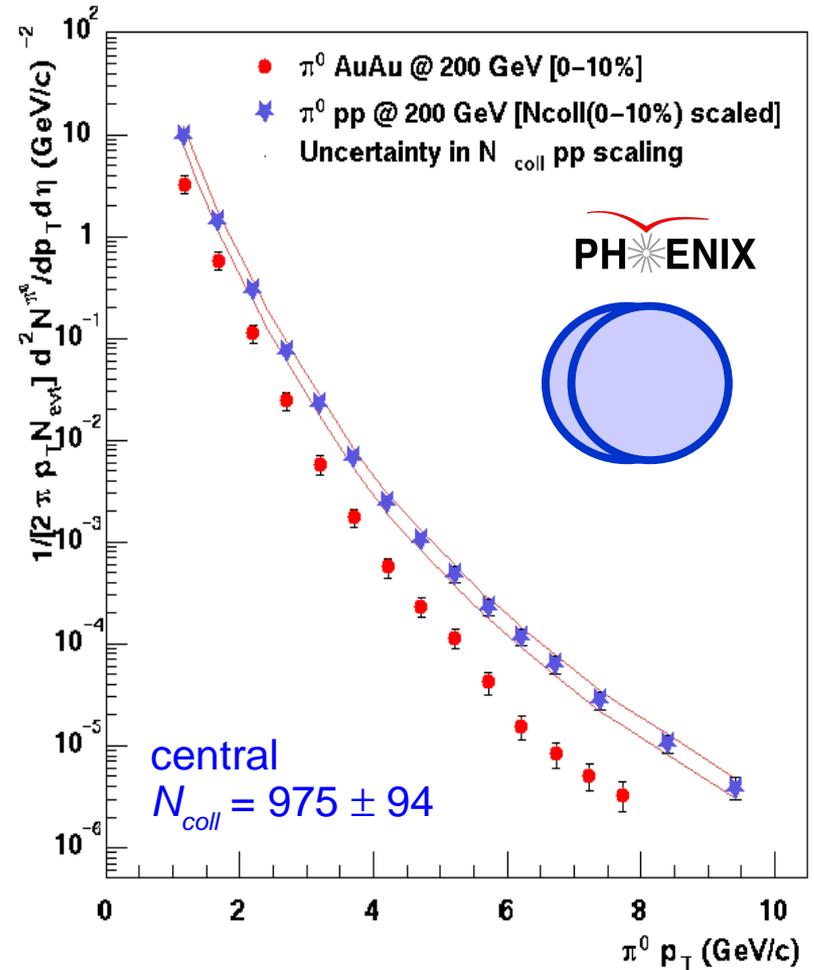
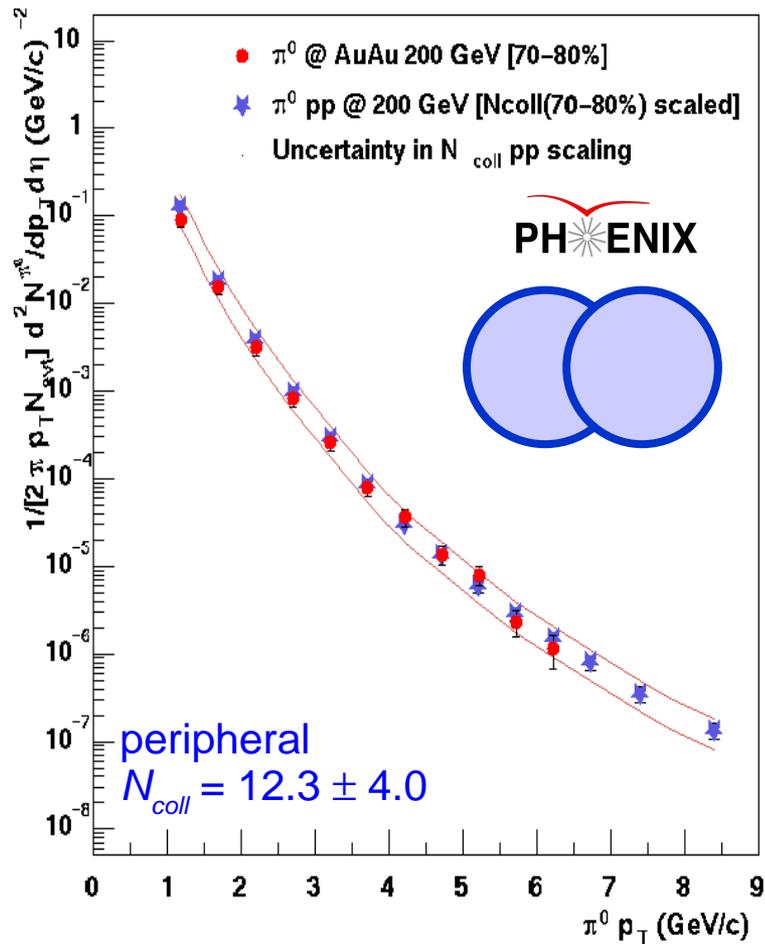


QCD works at RHIC

can use perturbation theory for high  $p$  transfer processes

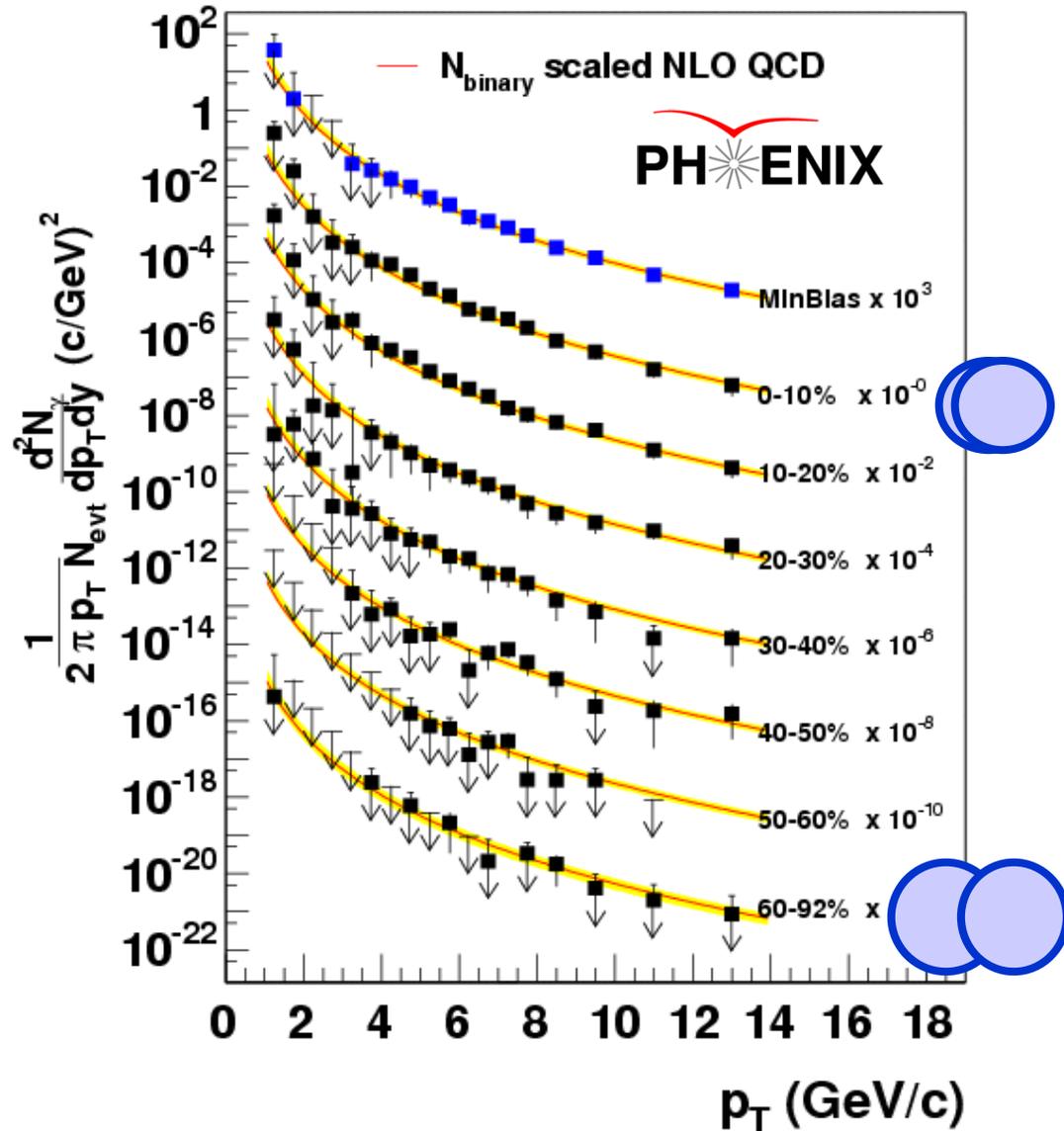
understand initial hard interactions in Au+Au: scattering of  $q, g$  inside N

# jet fragments in Au+Au vs. p+p



**jets are quenched by the plasma!**

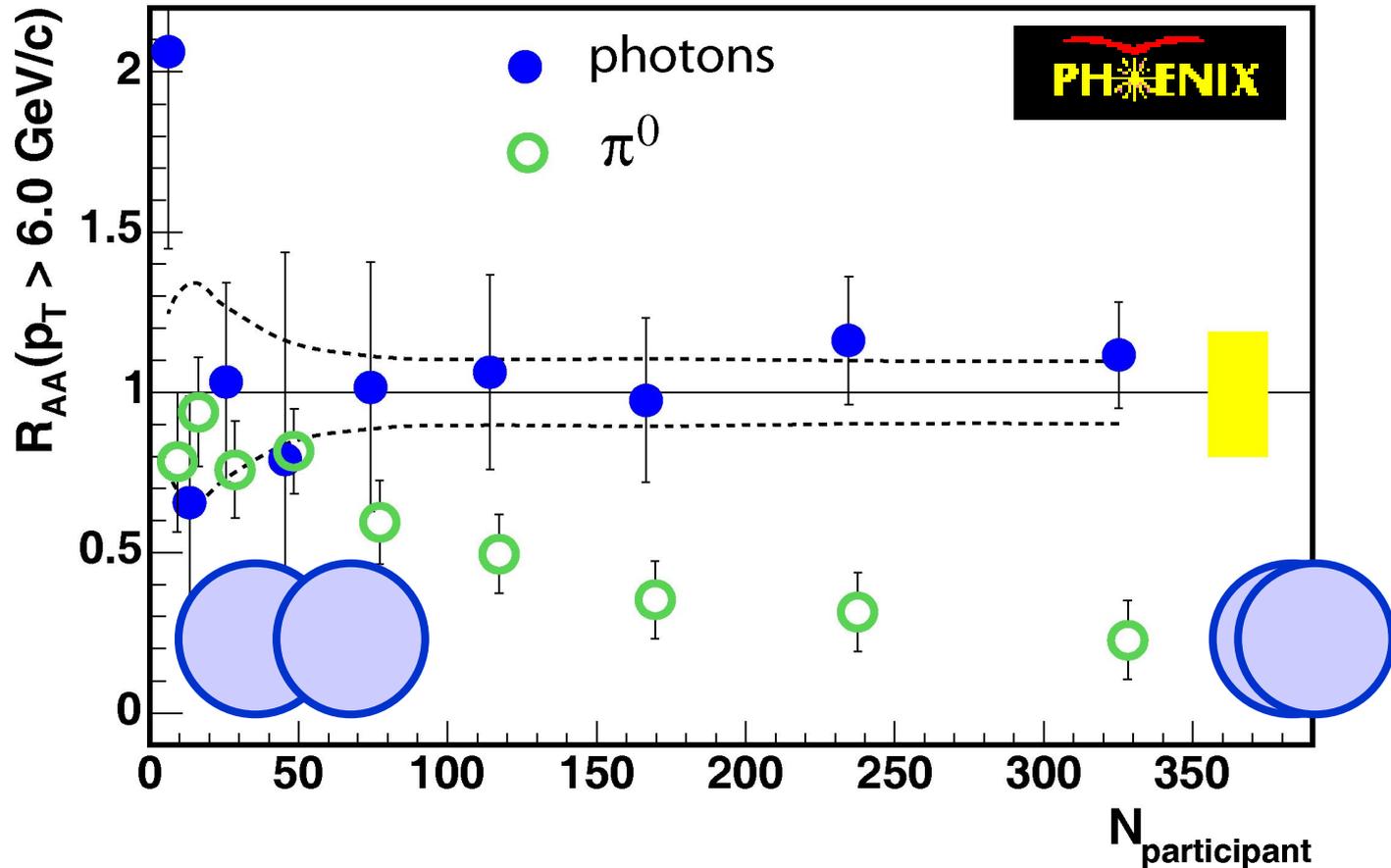
# Direct Photon Spectra in Au+Au



- $q + g \rightarrow q + \gamma$
- $\gamma$  should not interact with the color charges
- data agree with QCD  $\rightarrow$  calibrated probes
- pQCD works in the complex environment of two nuclei (Au+Au) colliding at high energies!

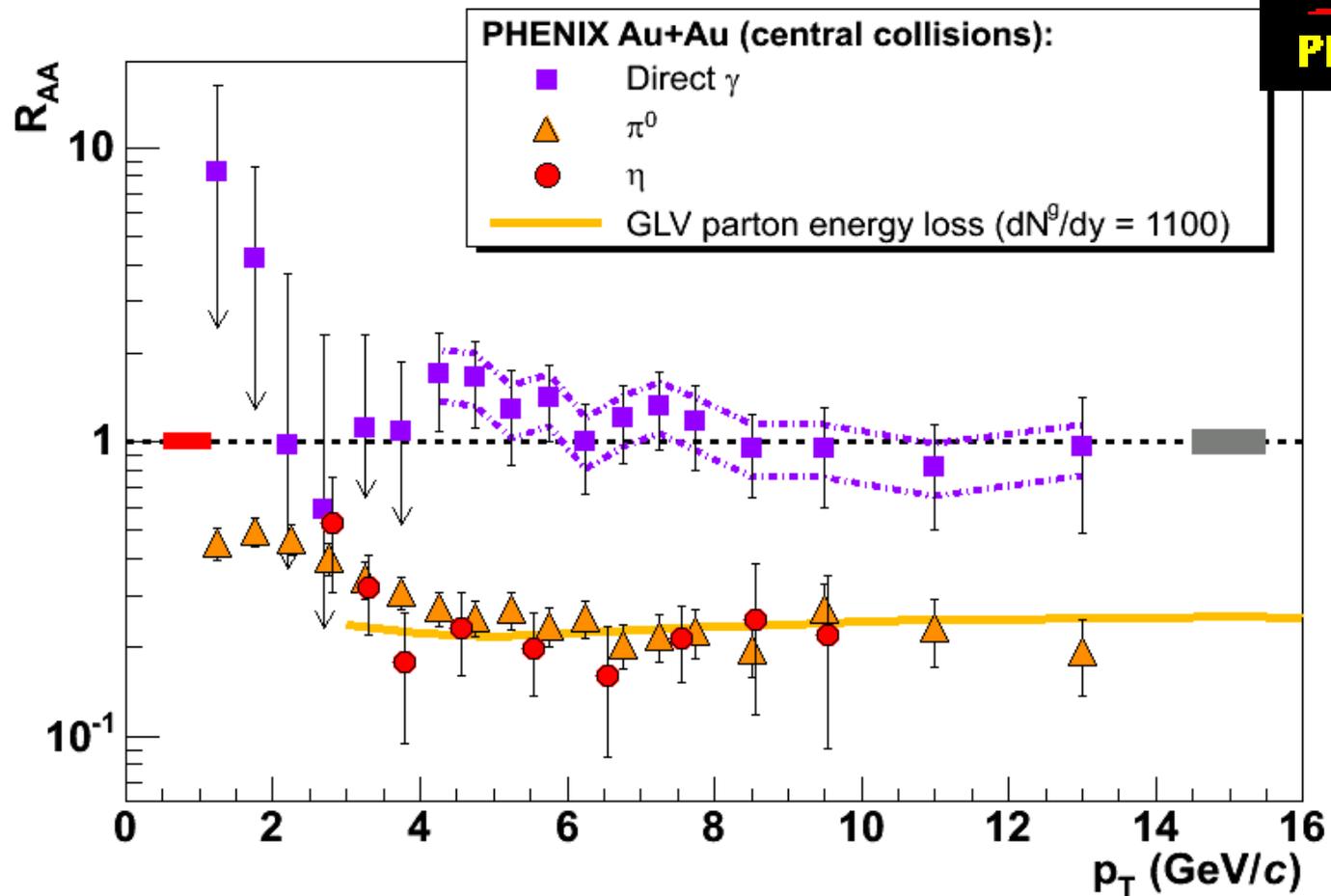
# nuclear modification factor

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$

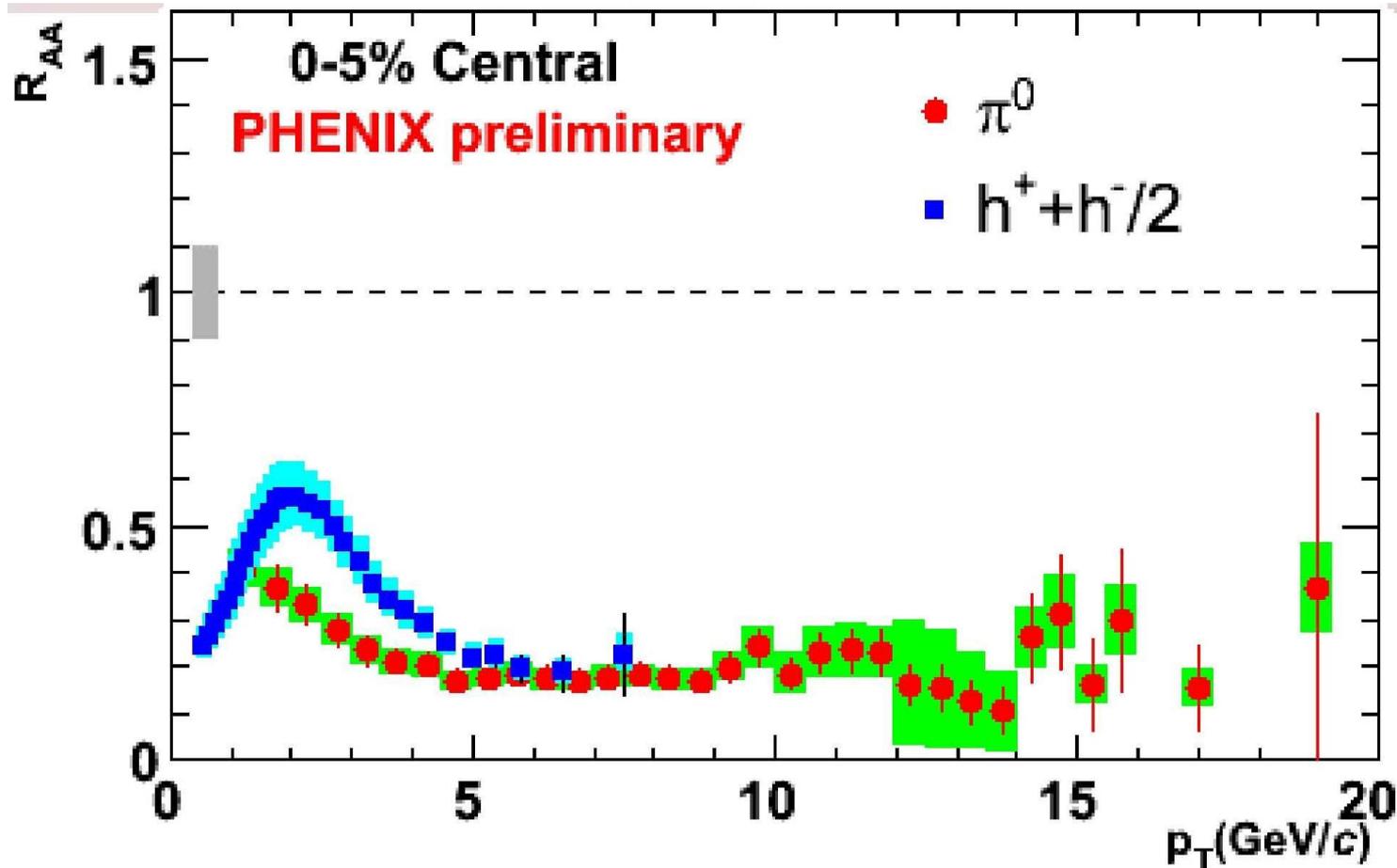


- photons escape plasma
- pions and other hadrons: strong interaction, absorbed

# mesons show common $p_T$ dependence



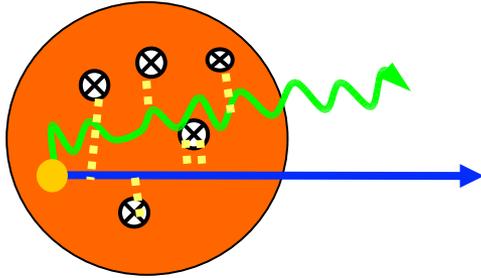
# suppression persists to 20 GeV/c!



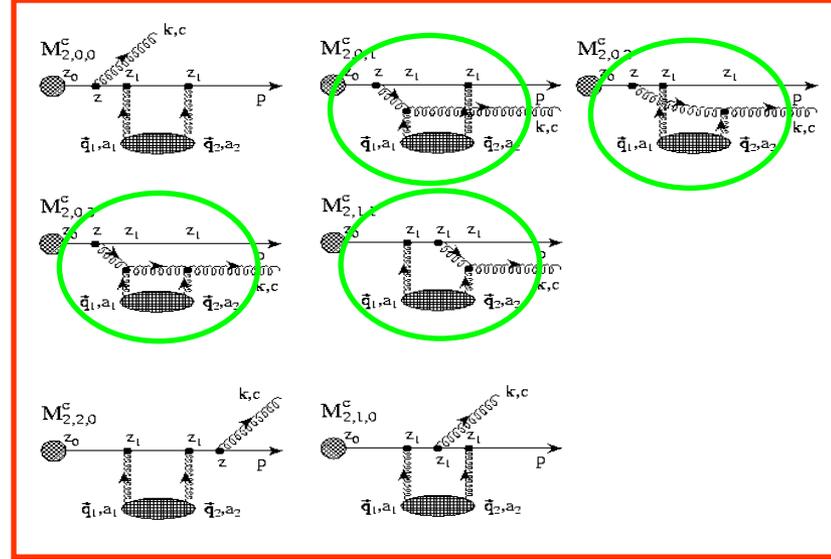
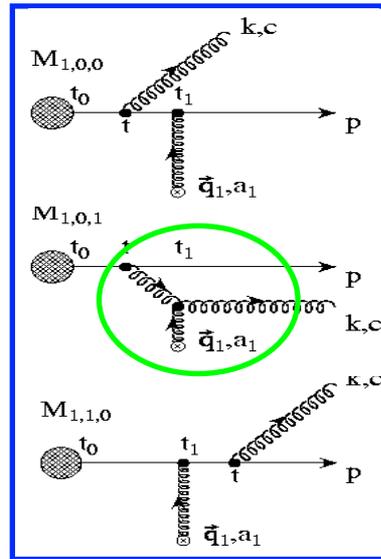
question to theory: accidental cancellations or a physics message?

flatness not expected at LHC – what will the data say?

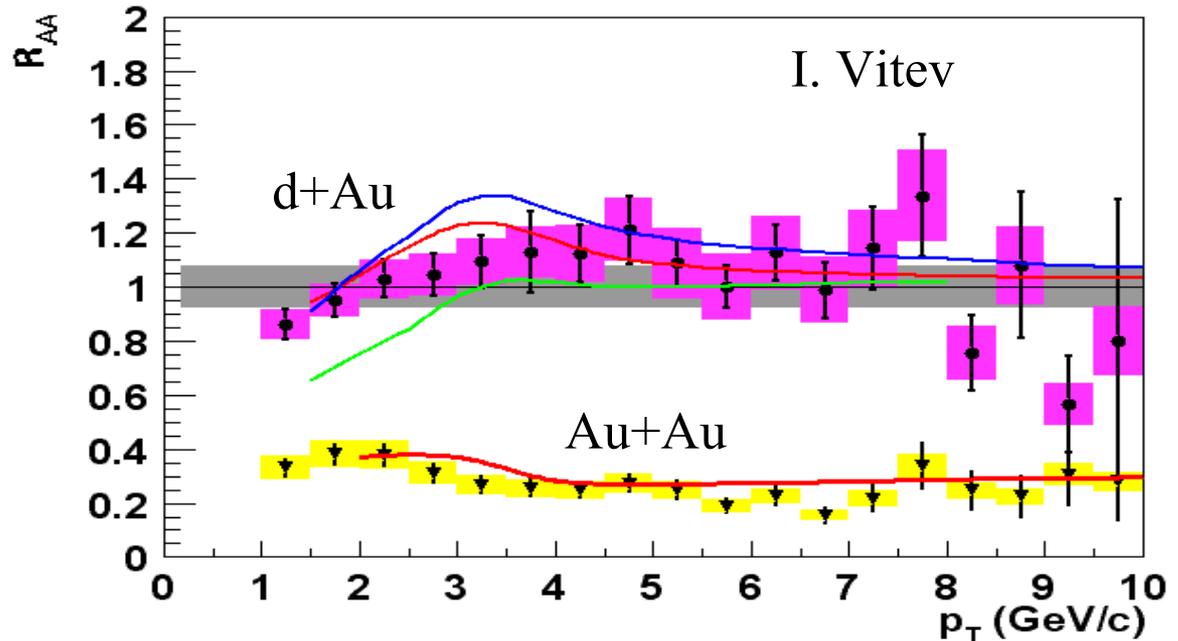
# learn about QGP density from this



interaction of radiated gluons with gluons in the plasma greatly enhances the amount of radiation

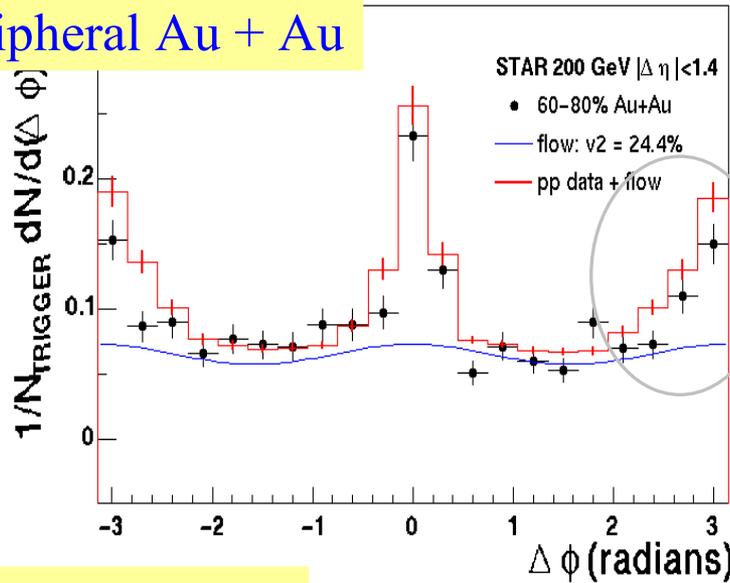


calculate using an opacity expansion  
 answer:  $L/mfp \sim 3.5$   
*model dependent!*  
 $\rho \geq 1000$  gluons/dy

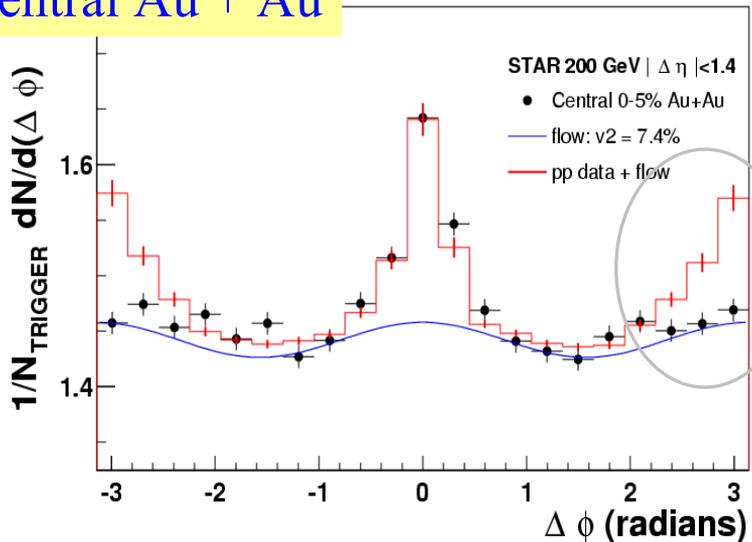


# look for the jet on the other side

## Peripheral Au + Au

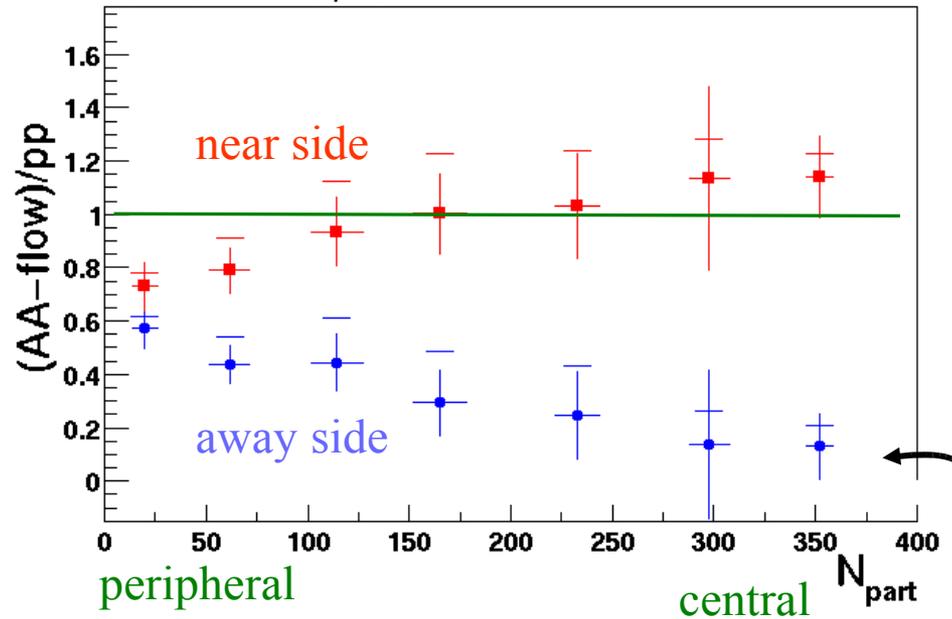


## Central Au + Au

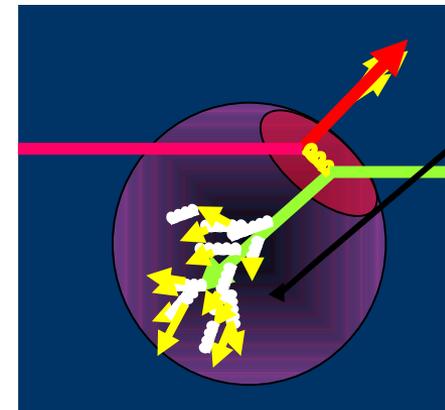


STAR PRL 90, 082302 (2003)

$$D_2(Au + Au) = D_2(p + p) + B(1 + v_2^2 \cos(2\Delta \phi))$$



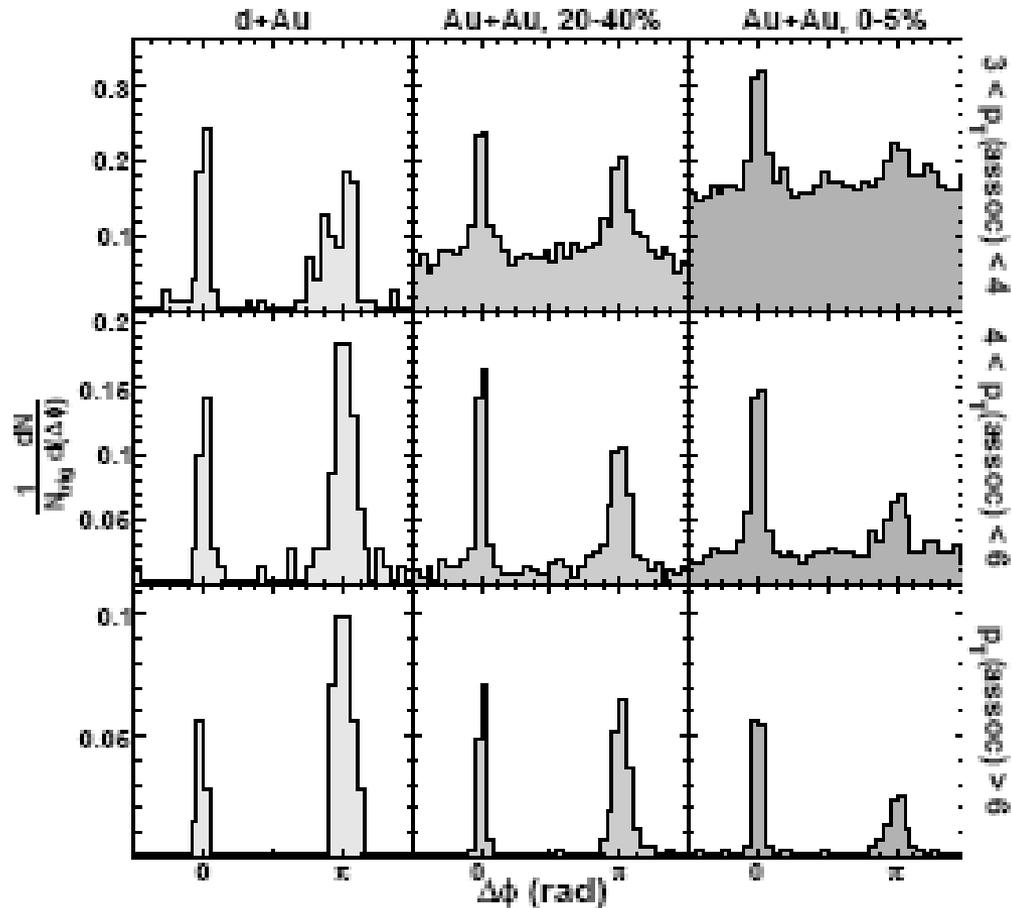
Medium  
is  
opaque!



# do see some away side particles at higher $p_T$

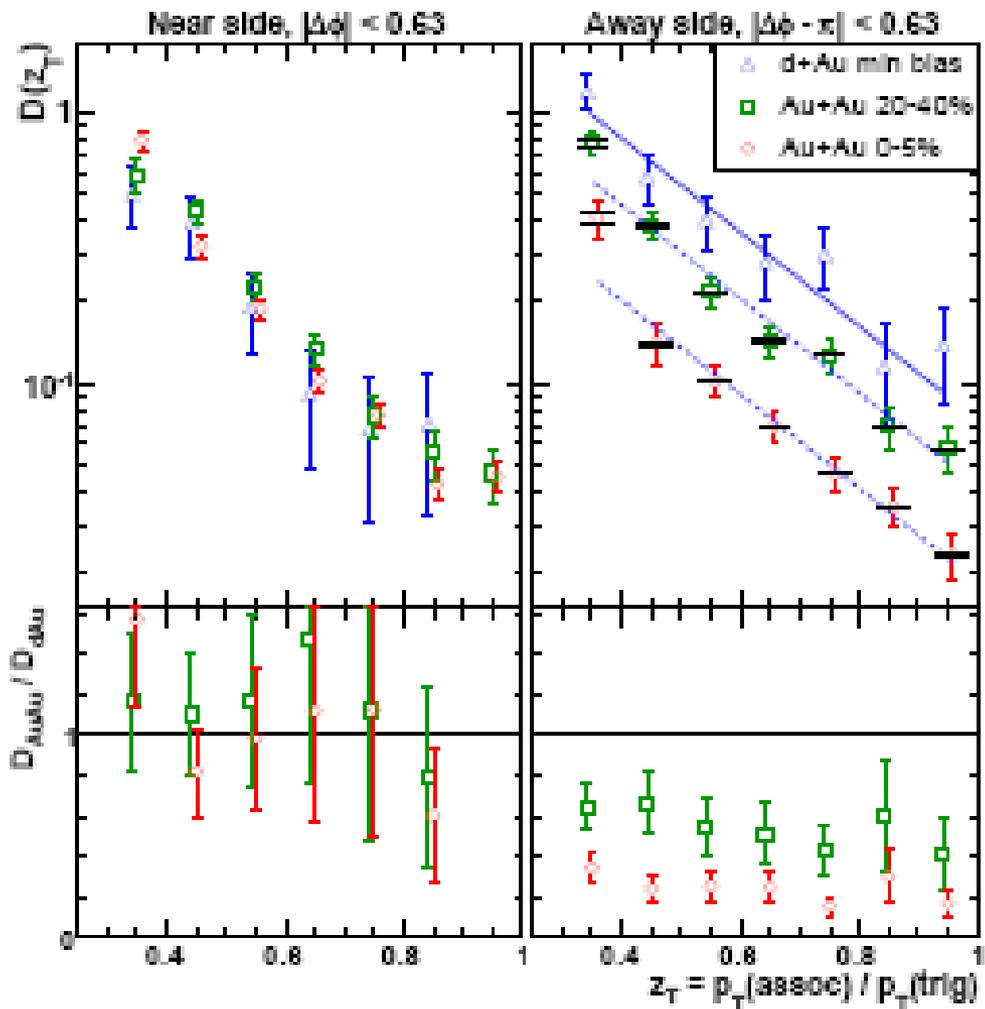
STAR nucl-ex/0604018

$p_T$  trigger  $> 8$  GeV/c

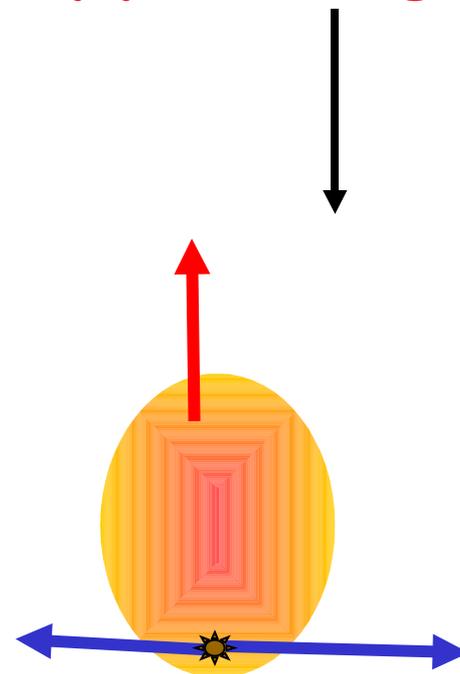


# away side yield: some jets escape, some eaten

STAR nucl-ex/0604018



**Note similarity of  
away side jet  
fragmentation.  
Only yield changes**



# summary of transmission measurements

---

**we have seen that:**

- **the medium is very opaque to hard probes  
opacity (L/mfp)  $\sim 3.5$**
- **medium is very dense and/or interaction cross section is very high  
perturbative calculations say  $dN_g/dy \sim 1000$**
- **but detailed calculations of expected medium modification of fragmentation function don't quite reproduce the data quantitatively...**

**Let's turn to another feature of plasmas:  
collective motions**

# collective effects

---

*a basic feature distinguishing plasmas from ordinary matter*

- **simultaneous interaction of each charged particle with a considerable number of others**

**due to long range of electromagnetic forces**

**both charge-charge and charge-neutral interactions**

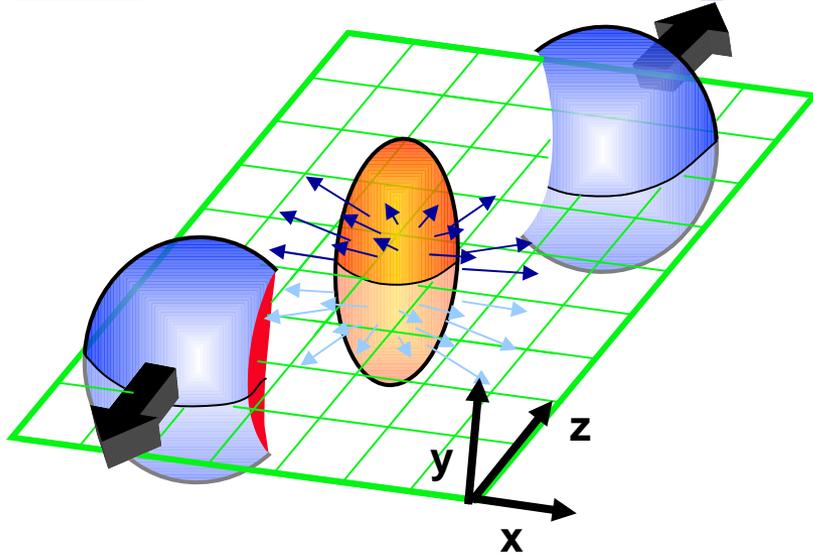
**charge-neutral dominates in weakly ionized plasmas**

**neutrals interact via distortion of e cloud by charges**

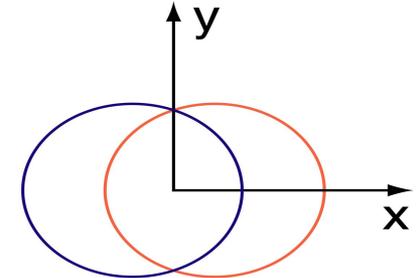
- **magnetic fields generated by moving charges give rise to magnetic interactions**

# search for collectivity in QGP

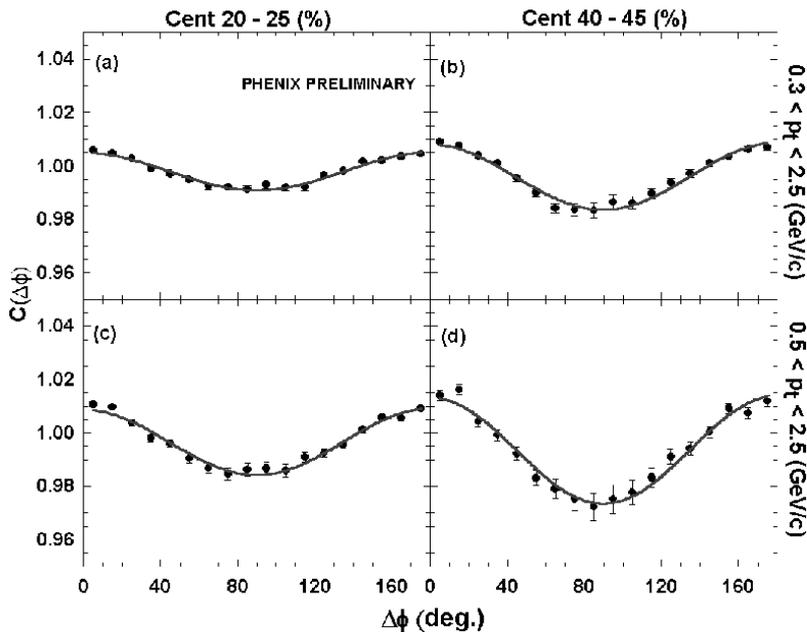
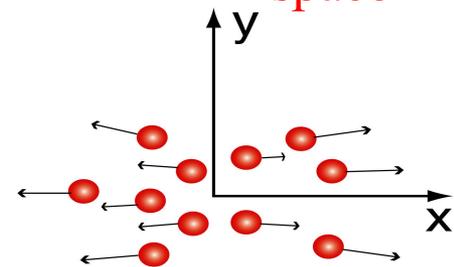
## use “internal” probes – emitted particles



Almond shape  
overlap region  
in **coordinate**  
**space**



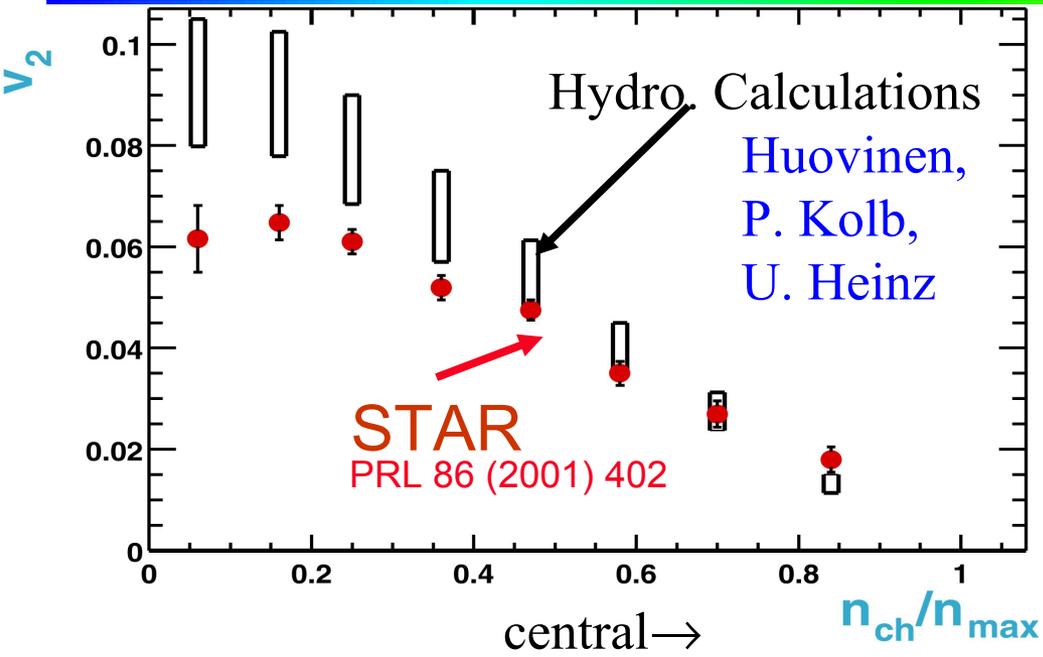
**momentum**  
**space**



$$dN/d\phi \sim 1 + 2 v_2(p_T) \cos(2\phi) + \dots$$

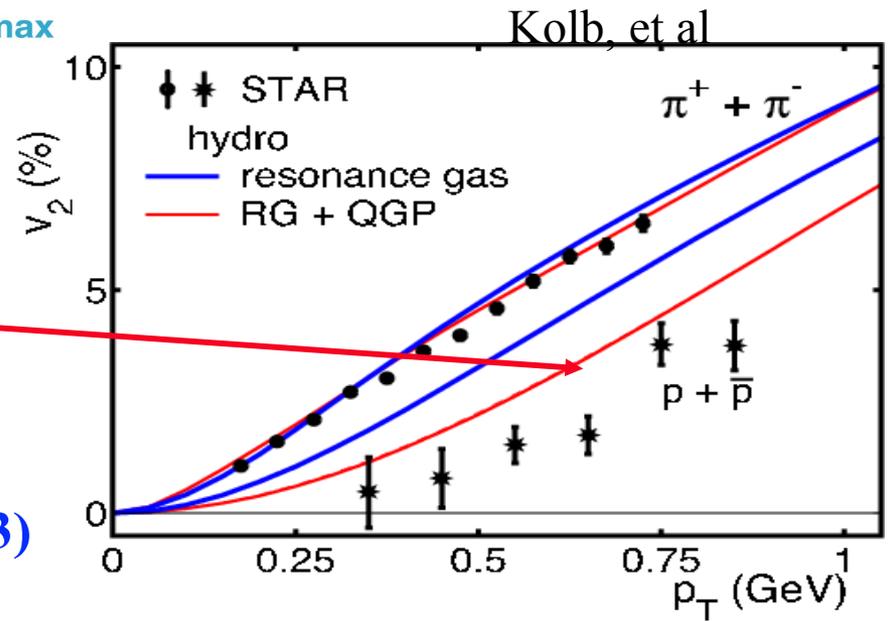
“elliptic flow”

# $v_2$ is large & reproduced by hydrodynamics



- large pressure buildup
- anisotropy  $\rightarrow$  happens fast
- fast equilibration!

Hydrodynamics reproduces elliptic flow  $q-\bar{q}$  and  $3q$  states  
 Mass dependence requires *softer than hadronic EOS!!*



NB: these calculations have viscosity  $\sim 0$   
 "perfect" liquid (D. Teaney, PRC68, 2003)

# Truth in advertising! (details matter)

proton

pion

nucl-ex/0410003

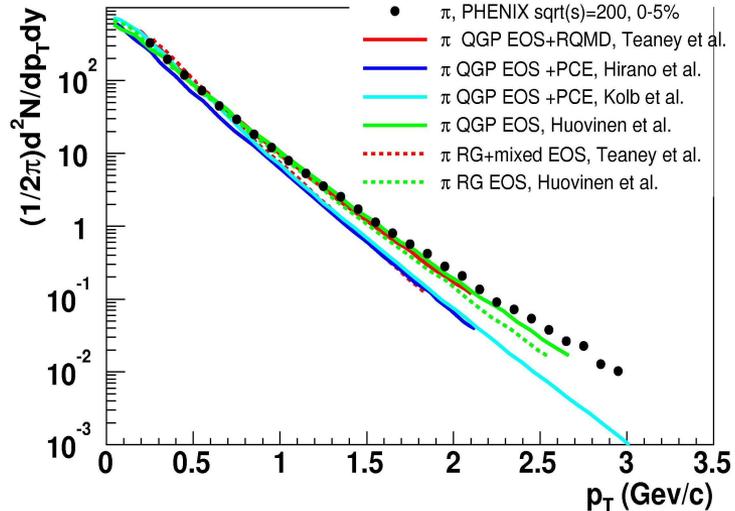
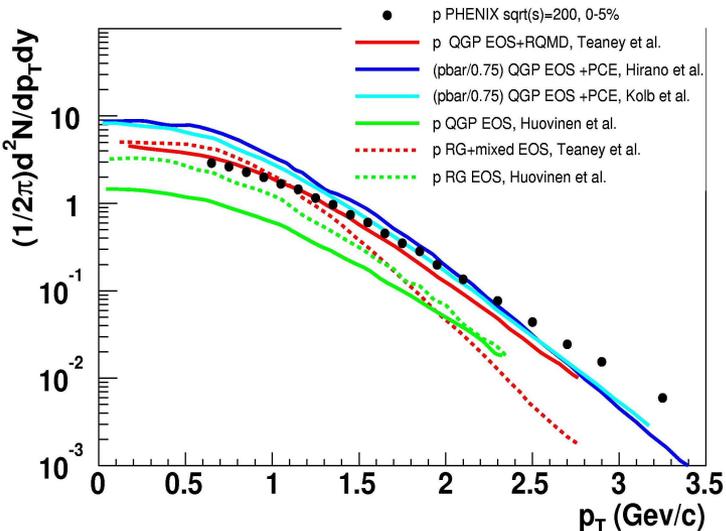
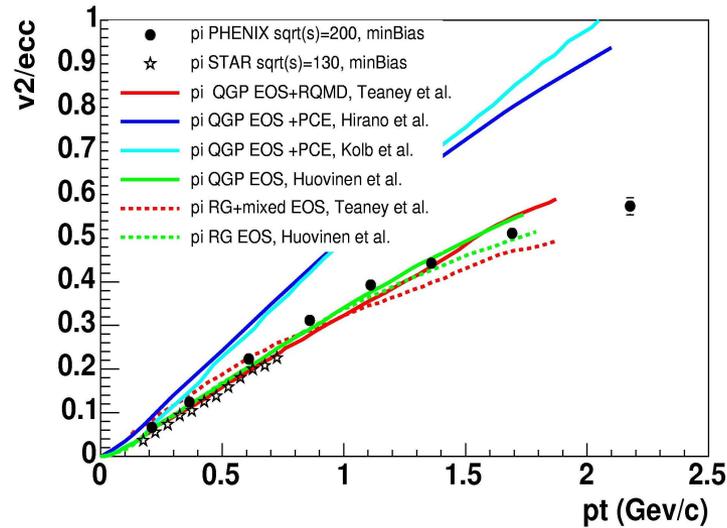
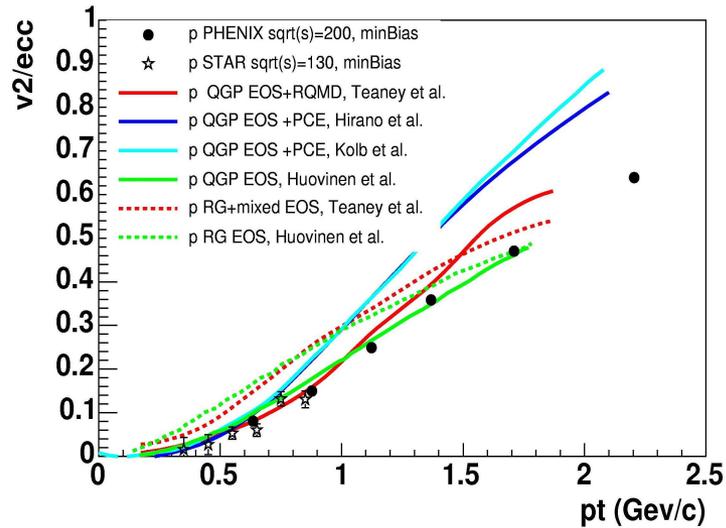
Hydro models:

**Teaney**  
(w/ & w/o  
*RQMD*)

**Hirano**  
(3d)

**Kolb**

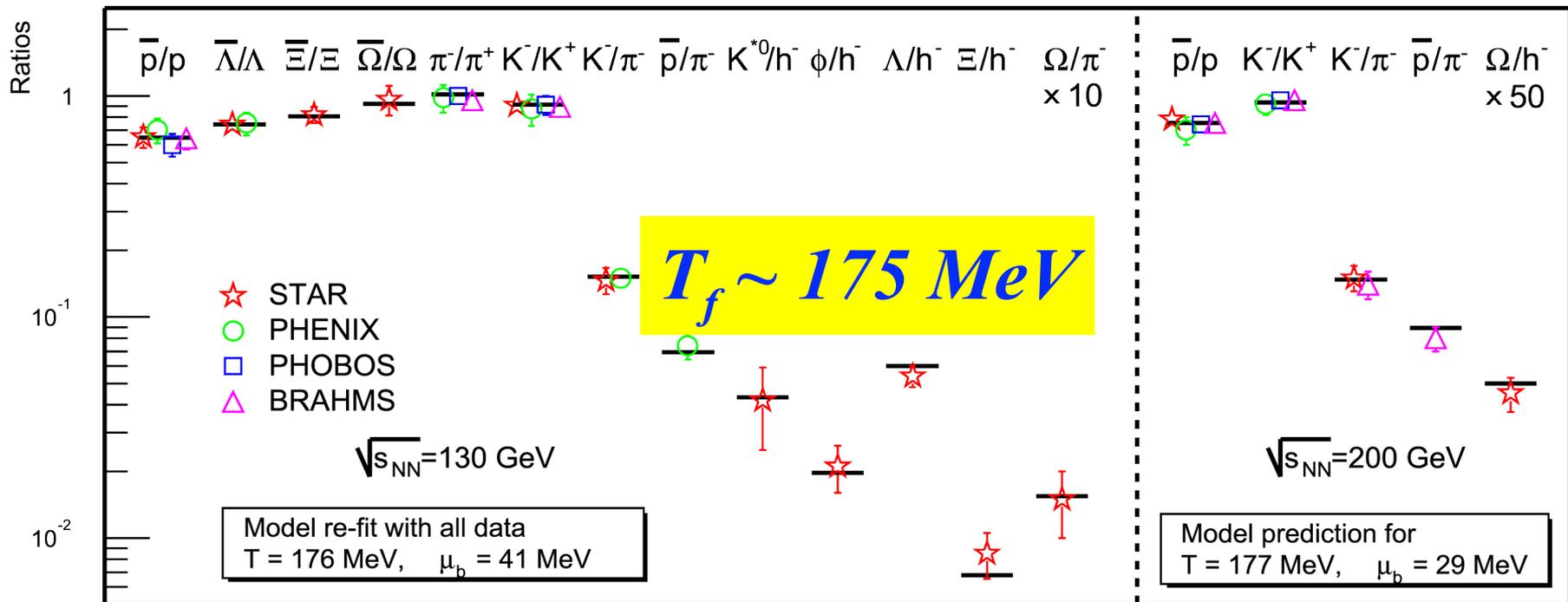
**Huovinen**  
(w/ & w/o  
*QGP*)



# hadron population arises from thermal system

Assume all distributions described by temperature  $T$   
and (baryon) chemical potential  $\mu$  :

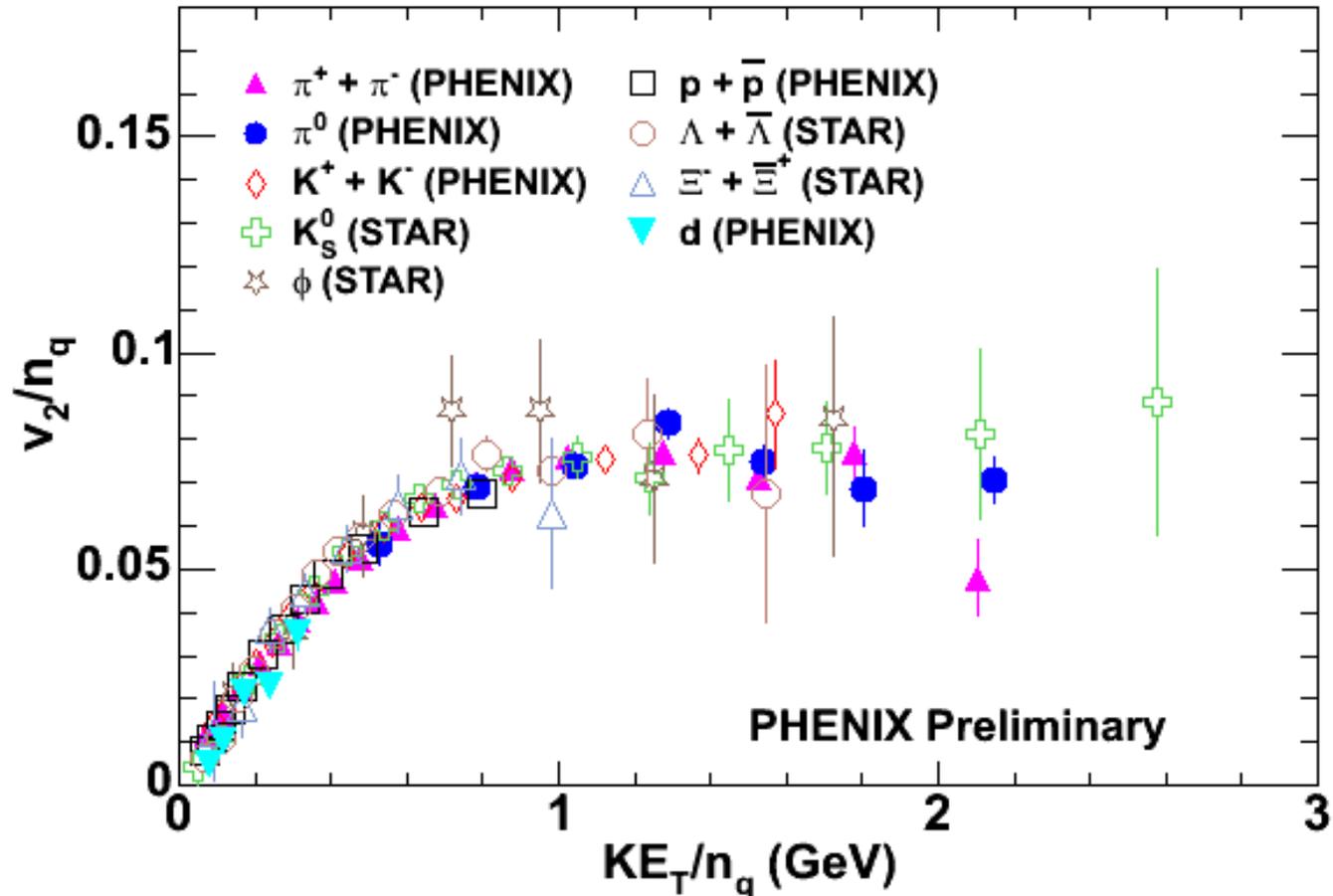
$$dn \sim e^{-(E-\mu)/T} d^3p$$



Braun-Munzinger et al., PLB 518 (2001) 41

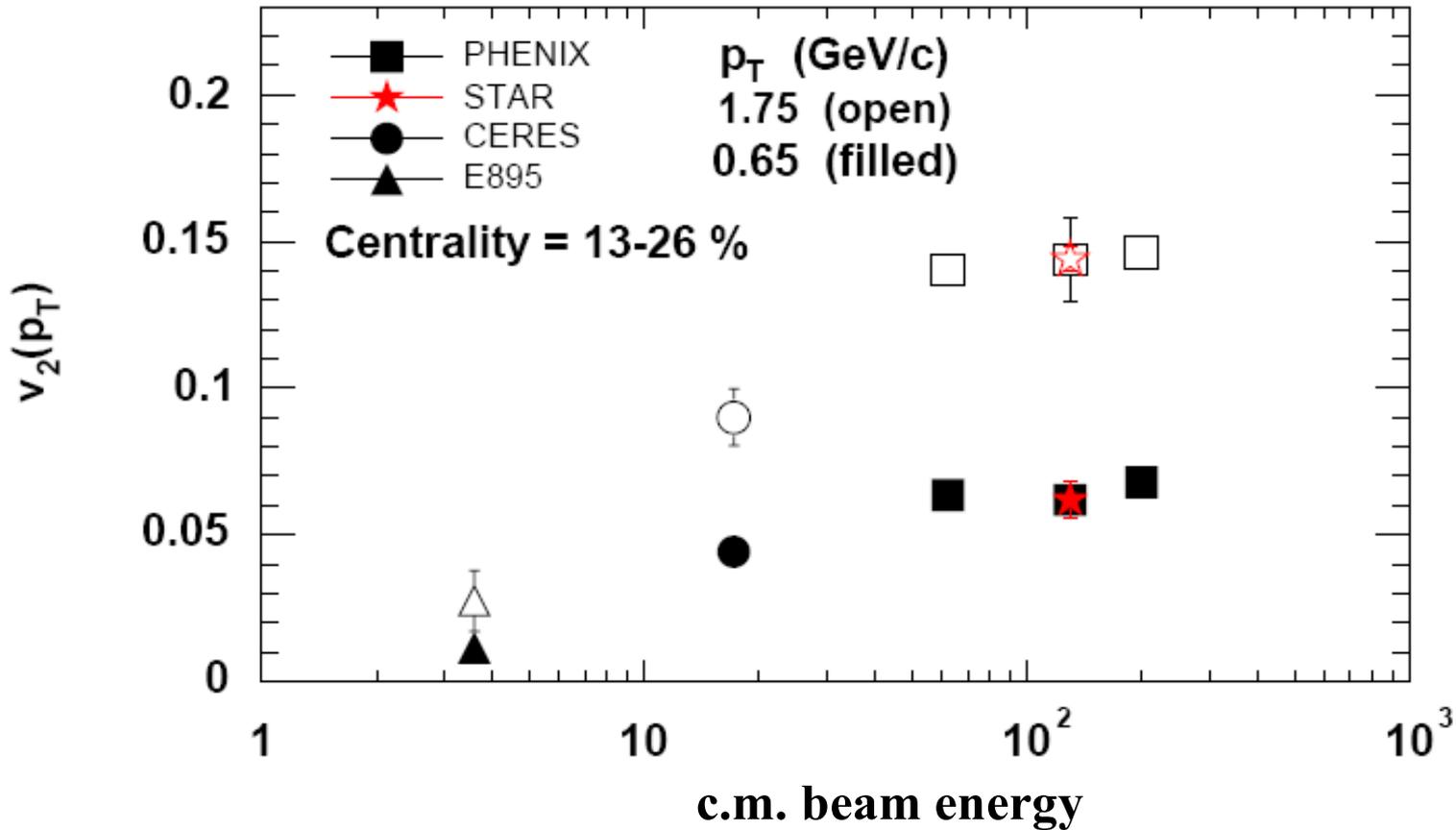
D. Magestro (updated July 22, 2002)

# Elliptic flow scales with number of quarks



*implication: quarks are the relevant degrees of freedom when the pressure is built up*

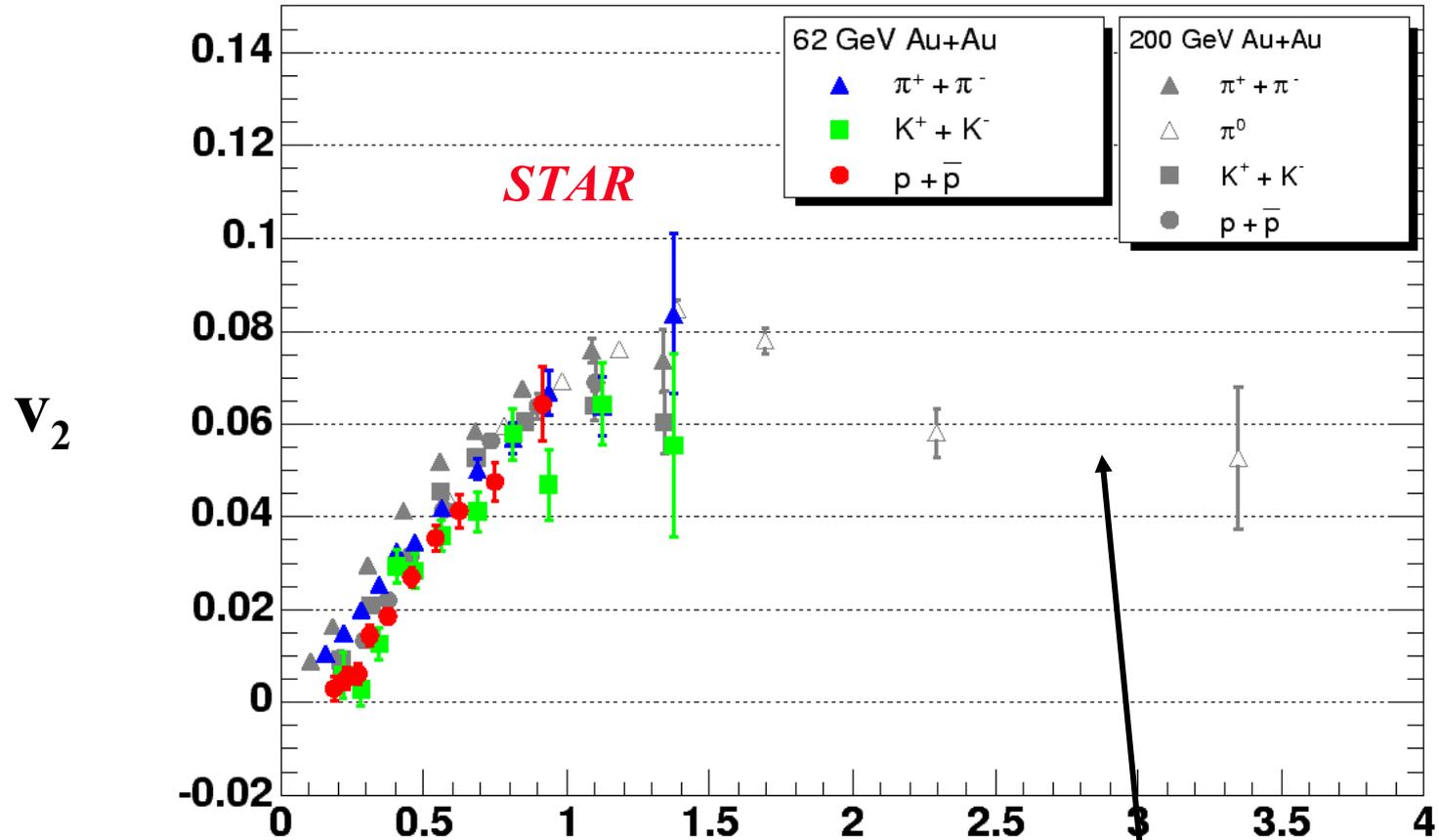
# beam energy dependence



**Anisotropy amplitude grows with beam energy, then flattens.**

*For LHC first guess – use same  $v_2$  at same  $p_T$*

# at high $p_T$ $v_2$ reflects opacity of medium



approximately expected level from jet quenching

## other probes? *consider leptons in matter*

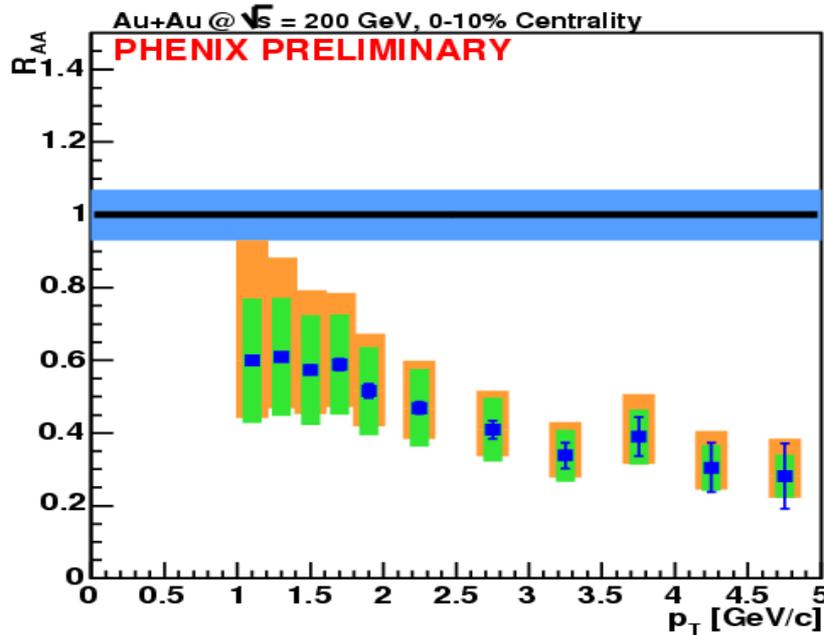
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- electrons vs. muons
- electrons radiate  $\gamma$  and stop very quickly  
the radiation is bremsstrahlung
- muons have large range because they DON'T radiate!  
radiation is suppressed by the large mass  
dominant energy loss mechanism is via collisions
- suggests that we should use heavy quarks to probe  
should we expect collisional energy loss for heavy quarks?  
is it reasonable to expect ONLY radiative energy loss for  
light quarks?

*EM plasmas suggest answer = no*

*radiation: blackbody, bremsstrahlung, collisional, recombination*

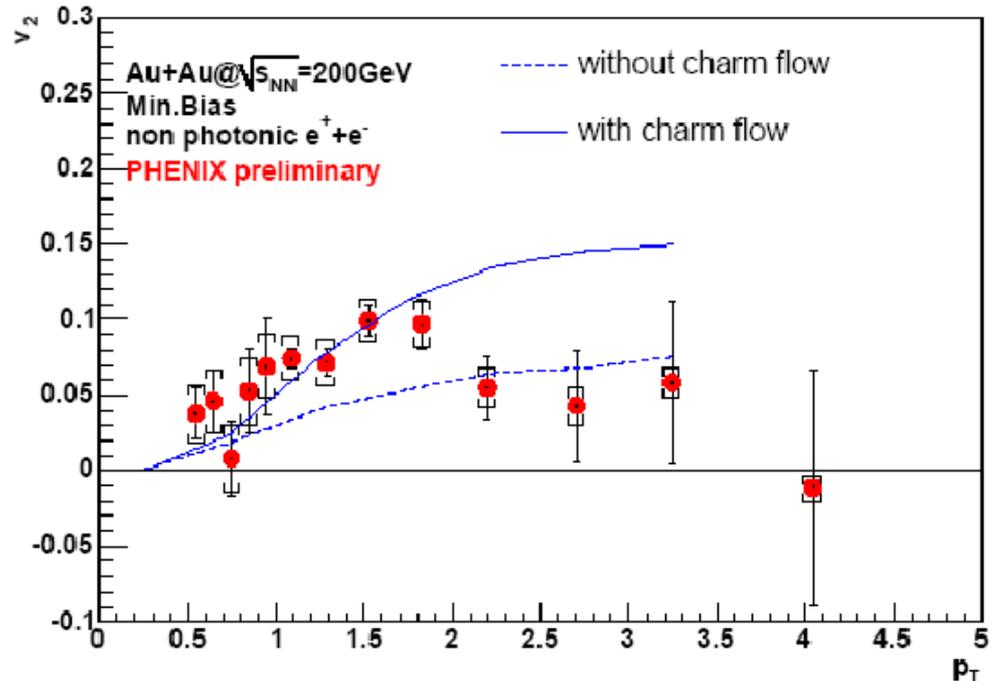
# further test interaction with charm quarks ( $m=1.3 \text{ GeV}/c^2$ )



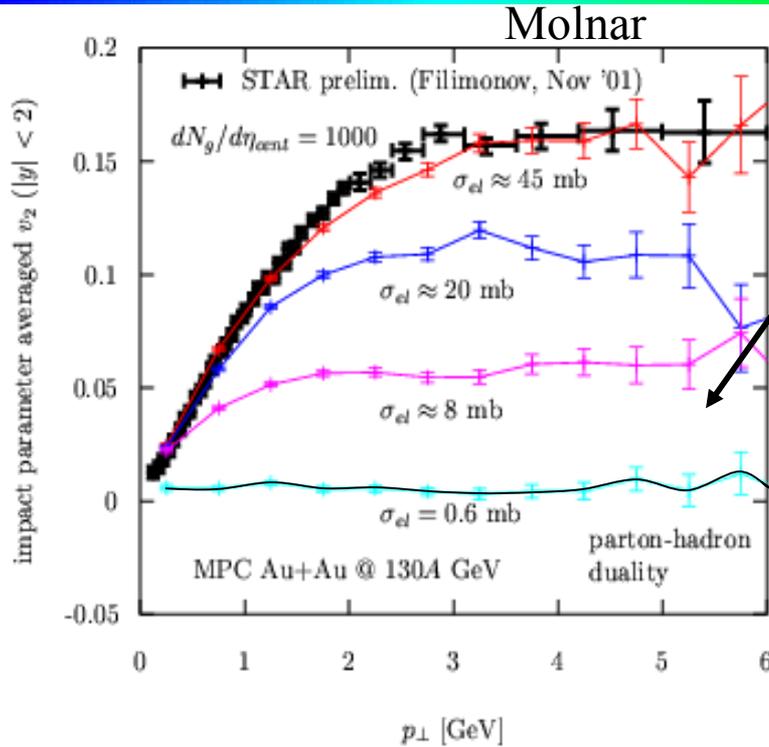
$\sim$  same E loss as u,d quarks  
 $\therefore$  energy loss not all radiative  
need collisions!

charm also flows

thermalization with  
the light quarks?  
*not so easy!*



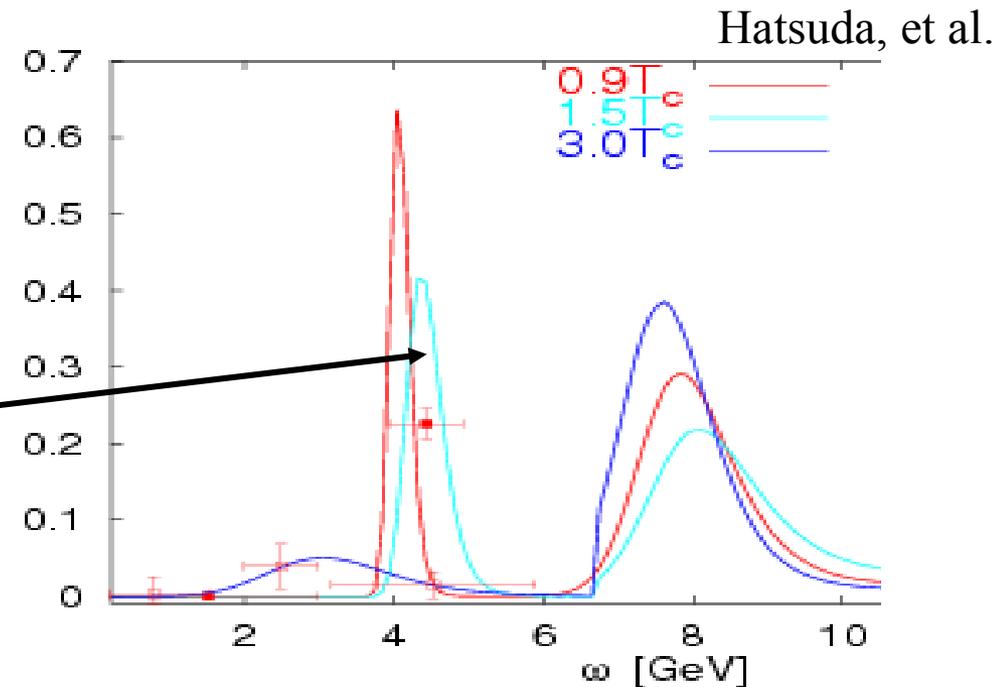
# Fast equilibration, high opacity (even for charm): how?



multiple collisions using free  $q, g$  scattering cross sections doesn't work!

need  $\sigma \times 50$  in the medium

Lattice QCD shows  $q\bar{q}$  resonant states at  $T > T_c$ , also implying high interaction cross sections



# What is going on?

---

- **The objects colliding inside the plasma are not baryons and mesons**
- **The objects colliding also do not seem to be quarks and gluons totally free of the influence of their neighbors**
- **So, what kind of a plasma should we expect the QGP to be?**

# Plasma Coulomb coupling parameter $\Gamma$

---

- ratio of mean potential energy to mean kinetic energy

$$\Gamma_e = (e^2 / 4\pi\epsilon_0 a) / k_B T_e$$

**a = interparticle distance**

**e = charge**

**T = temperature**

$$a = (4\pi n / 3)^{-1/3}$$

- typically a small number in a normal, fully shielded plasma  
 $\Gamma = 1/(\text{number particles in Debye sphere})$
- when  $\Gamma > 1$  have a strongly coupled, or non-Debye plasma  
many-body spatial correlations exist  
behave like liquids, or even crystals when  $\Gamma > 150$   
 $\lambda_D < a$

# plasma basics – Debye Length

---

- distance over which the influence of an individual charged particle is felt by the other particles in the plasma
- charged particles arrange themselves so as to effectively shield any electrostatic fields within a distance of order  $\lambda_D$
- $\lambda_D = \left[ \frac{\epsilon_0 kT}{n_e e^2} \right]^{1/2}$
- Debye sphere = sphere with radius  $\lambda_D$
- number electrons inside Debye sphere is typically large

$$N_D = N/V_D = \rho V_D \quad V_D = \frac{4}{3} \pi \lambda_D^3$$

# Debye screening in QCD: a tricky concept

---

- in leading order QCD (O. Philipsen, hep-ph/0010327)

$$V(r) \sim \frac{e^{-m_D^0 r}}{4\pi r}, \quad m_D^0 = \left( \frac{N}{3} + \frac{N_f}{6} \right)^{1/2} gT.$$

- However, at next-to-leading order the problem becomes non-perturbative. The general form of the series in  $g$  can be shown to be <sup>2</sup>

$$m_D = m_D^0 + \frac{N}{4\pi} g^2 T \ln \frac{m_D^0}{g^2 T} + c_N g^2 T + \mathcal{O}(g^3 T), \quad (6)$$

which is non-analytic in the coupling constant. While the coefficient of the logarithm is fixed perturbatively,  $c_N$  is entirely non-perturbative. The reason is that, starting from this order in  $g$ , the non-abelian  $A_0$  couples to the soft magnetic gluons  $A_i \sim g^2 T$ , and hence becomes sensitive to the non-abelian infrared divergencies in the magnetic sector for which there is no perturbative cure <sup>3</sup>:  $\Pi_{ii}(k_0 = 0, \mathbf{k} \rightarrow 0) \sim g^2 T \neq 0$ , with contributions from all loop orders two and larger.

This raises the conceptual problem whether a perturbative definition of Debye screening in QCD is at all sensible.

# give up on the concept?

---

*Of course not!!!*

- **Two options proposed by Philippsen:**
  - 1) **assume a pole in the propagator and attempt to measure its value from the exponential fall-off in some fixed gauge (done with lattice QCD)**
  - 2) **seek a manifestly gauge invariant definition**
- **Lattice says: “interactions weak @  $l = 1/T$ , but screening function is not exponential until  $1/2T$ ”**
- *idea: calculate  $\lambda_D$  for strongly coupled plasma & convert  $\varepsilon$  inside to particle density to get  $\Gamma$*

# screening and thermal masses

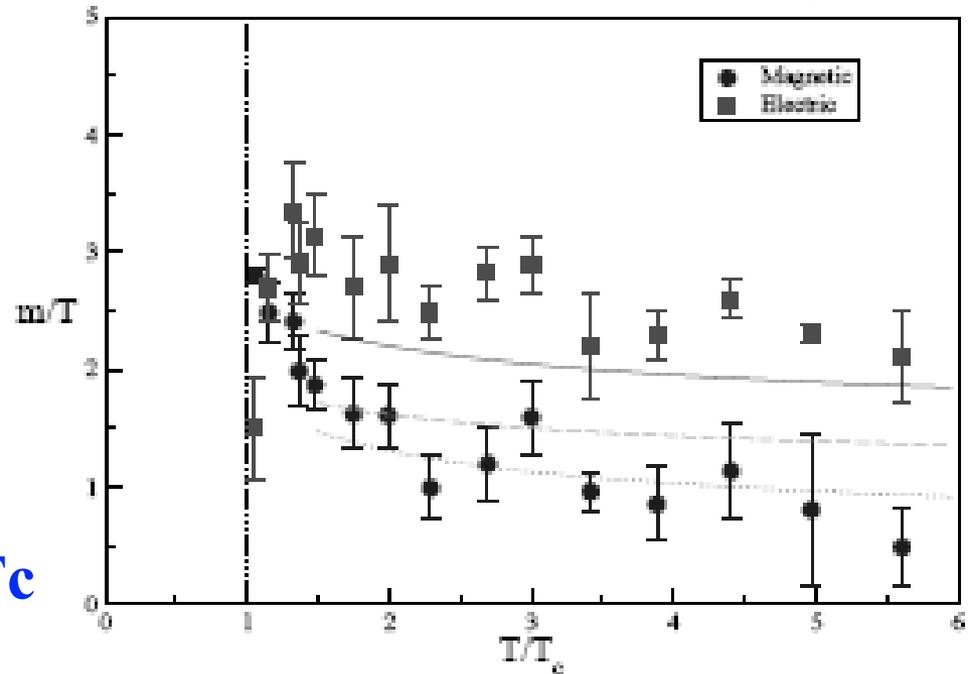
*Screening mass,  $m_D$ , defines inverse length scale*

*Inside this distance, an equilibrated plasma is sensitive to insertion of a static source*

*Outside it's not.*

Nakamura, Saito & Sakai, hep-lat/0311024

**T dependence of electric & magnetic screening masses**  
**Quenched lattice study of gluon propagator**



**figure shows:**

**$m_{D,m} = 3T_c$ ,  $m_{D,e} = 6T_c$  at  $2T_c$**

**$\lambda_D \sim 0.4$  &  $0.2$  fm**

*magnetic screening mass significant*  
*not very gauge-dependent, but DOES*  
*grow w/ lattice size (long range is important)*

# compare to oversimplified $\Gamma$ estimate

estimate  $\Gamma = \langle PE \rangle / \langle KE \rangle$

using QCD coupling strength  $g$

$\langle PE \rangle = g^2/d$   $d \sim 1/(4^{1/3}T)$

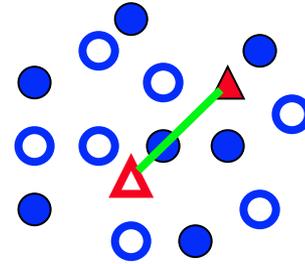
$\langle KE \rangle \sim 3T$

$\Gamma \sim g^2 (4^{1/3}T) / 3T$

$g^2 \sim 4-6$  (value runs with  $T$ )

for  $T=200$  MeV plasma parameter  $\Gamma \sim 3$

*see M. Thoma, J.Phys. G31(2005)L7*



*$\Gamma > 1$ : strongly coupled, few particles inside Debye radius*

*use  $\lambda=0.2$  fm and  $\epsilon=15$  GeV/fm<sup>3</sup> get 0.5 GeV inside Debye sphere*

*FEW particles!*

**$\therefore$  quark gluon plasma should be a strongly coupled plasma**

**As in warm, dense plasma at lower (but still high)  $T$**

**dusty plasmas, cold atom systems**

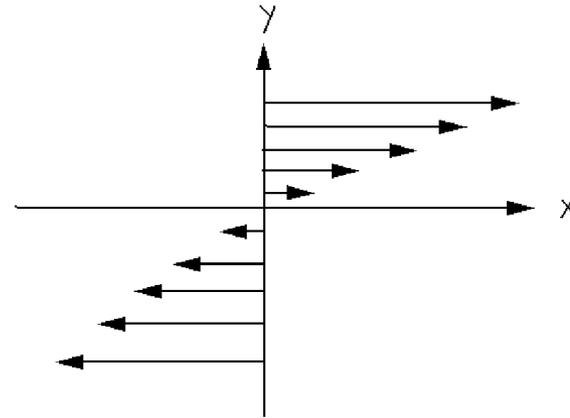
*such EM plasmas are known to behave as liquids!*

# so, what is a “perfect fluid?”

- one that exhibits ideal, non-dissipative hydrodynamics  
“not-quite-ideal” = can support a shear stress

Viscosity  $\eta$  is 
$$\frac{F_x}{A} = -\eta \frac{\partial v_x}{\partial y}$$

- $\eta \sim \langle v \rangle / \sigma \sim \sqrt{kT} / \sigma$   
so viscosity *increases* with  $T$   
*decreases* with  $\sigma$



Ideal hydrodynamics ( $\eta \sim 0$ ) reproduces data at RHIC  
does this make sense to QCD?

→  $\eta \geq (h/4\pi) S$  : conjectured quantum mechanical limit

entropy density

“A Viscosity Bound Conjecture”,  
P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231

# collisions → transport in the plasma

---

- **transport of particles → diffusion**
- **transport of energy by particles → thermal conductivity**
- **transport of momentum by particles → viscosity**
- **transport of charge by particles → electrical conductivity**  
**is transport of color charge an analogous question for us?**

# what exactly is diffusion?

---

- **diffusion = brownian motion of particles**

**definition: flux density of particles  $J = -D \text{ grad } n$**

*particle concentration*

- **integrating over forward hemisphere:**

**$D = \text{diffusivity} = 1/3 \langle v \rangle l$**

*$l = \text{mean free path}$*

**so  $D = \langle v \rangle / 3n\sigma$**

**$D \propto \text{collision time}$**

**determines relaxation time for the system**

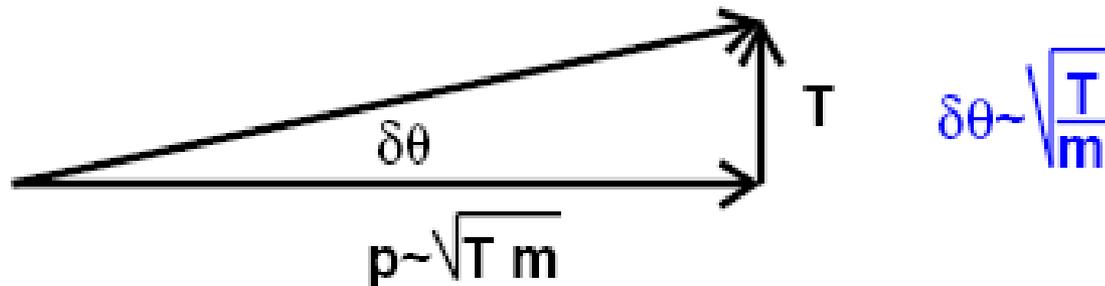
**note:  $\eta = 1/3 \rho \langle v \rangle l$  so  $D = \eta/\rho$**

**nice implication: measure  $D$  get  $\eta$ !**

**$\rho$  from  $T$ , or maybe transmission**

## Estimate of transport times with Heavy Quarks

- Put a heavy quark in this medium

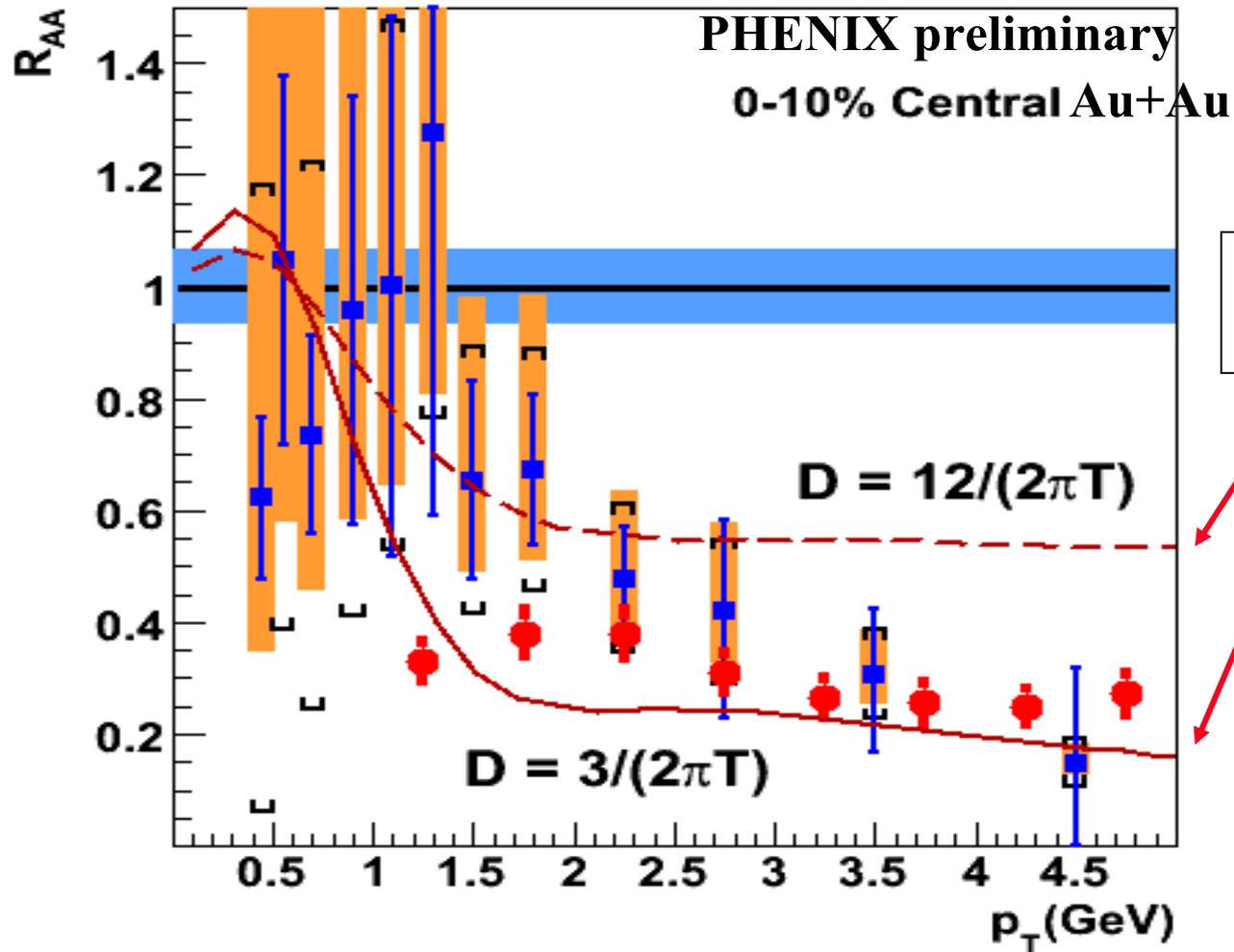


- The charm quark undergoes a random walk suffering many collisions
- The relaxation time of the heavy quark is:

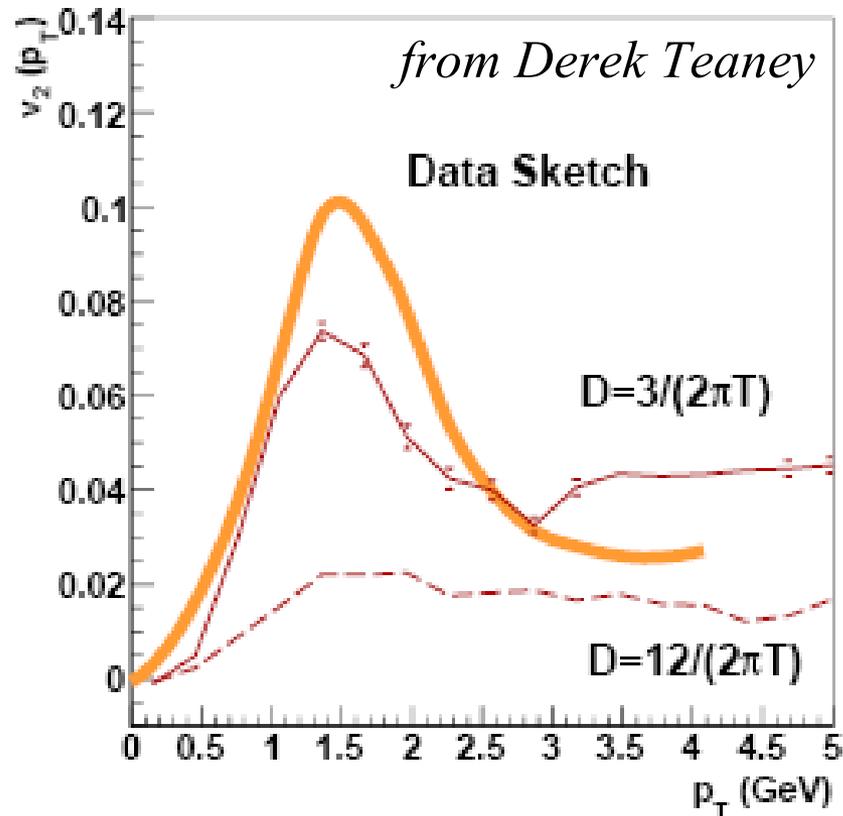
$$\tau_R^{\text{charm}} \sim \frac{M}{T} \tau_R^{\text{light}}$$

Will the heavy quark follow the flow of the underlying medium

# can we measure the diffusion coefficient?



# collisional energy loss also implies flow



$$D \sim 3/(2\pi T)$$

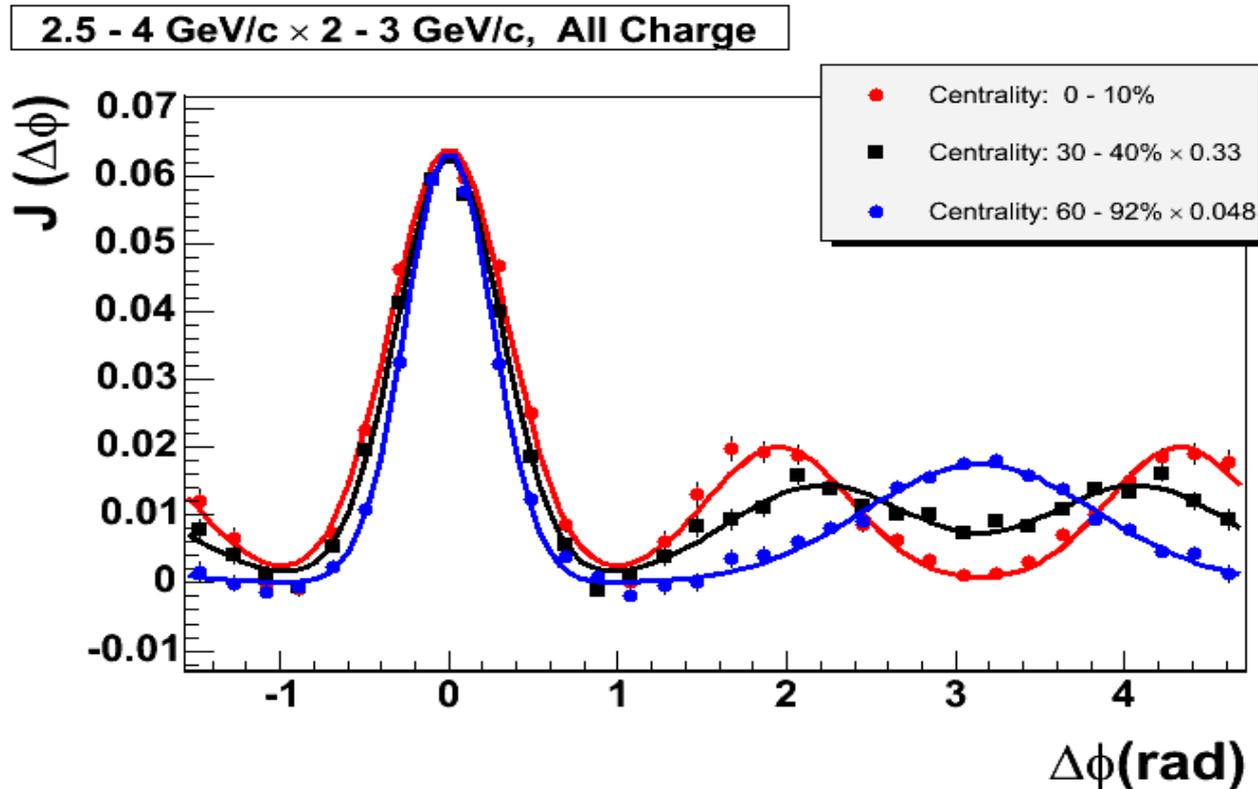
small!

→ strong interaction  
of charm quarks  
with the plasma

larger  $D$  would mean  
less charm  $e$  loss  
fewer collisions with  
plasma, smaller  $v_2$

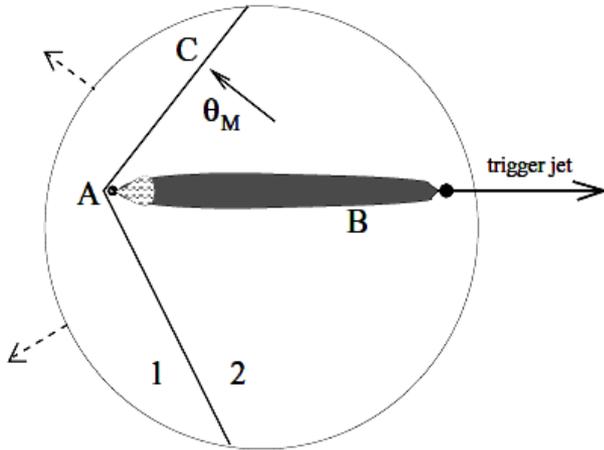
# can we measure the energy transport?

- Plasma physicists do by applying a shear stress, or shock
- Jets do that when they dump energy into QGP  
where does the energy go?



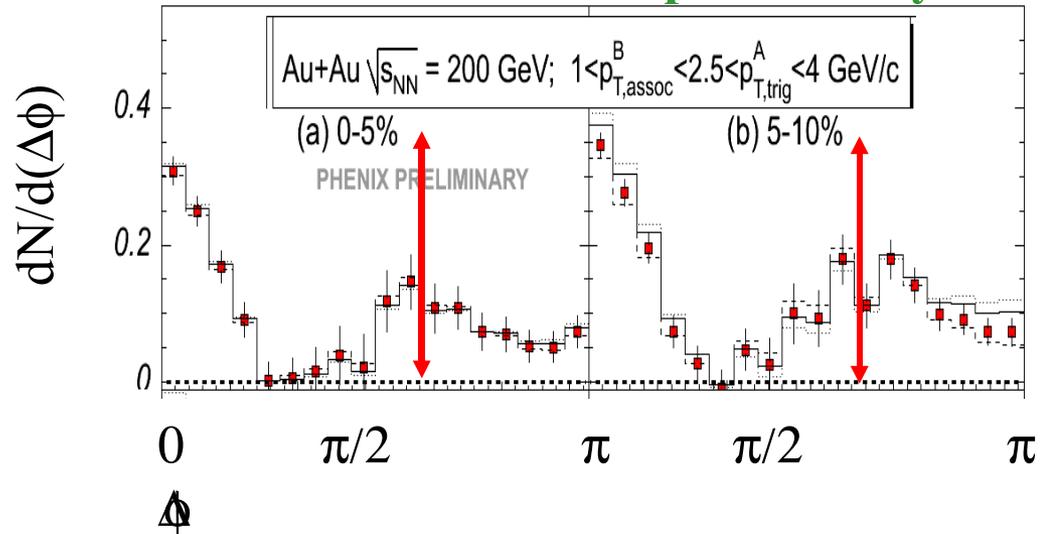
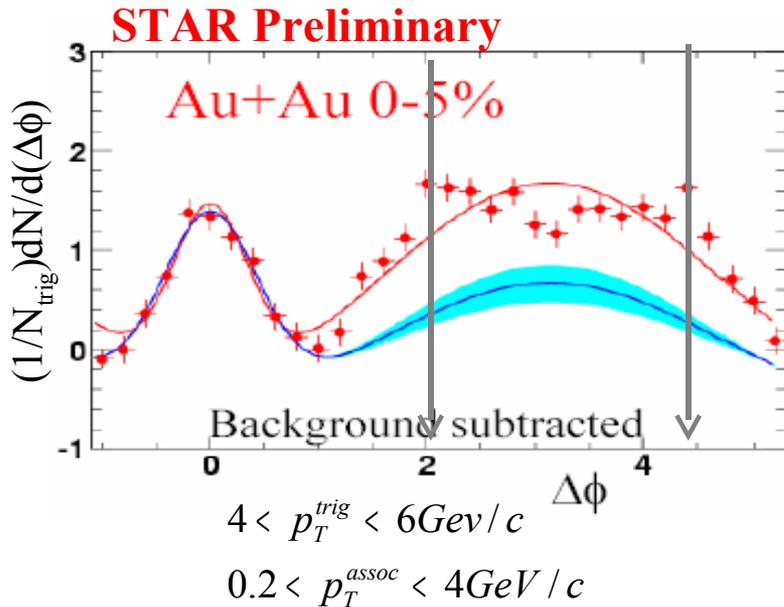
*look at lower  
 $p_T$  particles  
where away  
side jet should  
be*

# excite a density wave in the plasma?



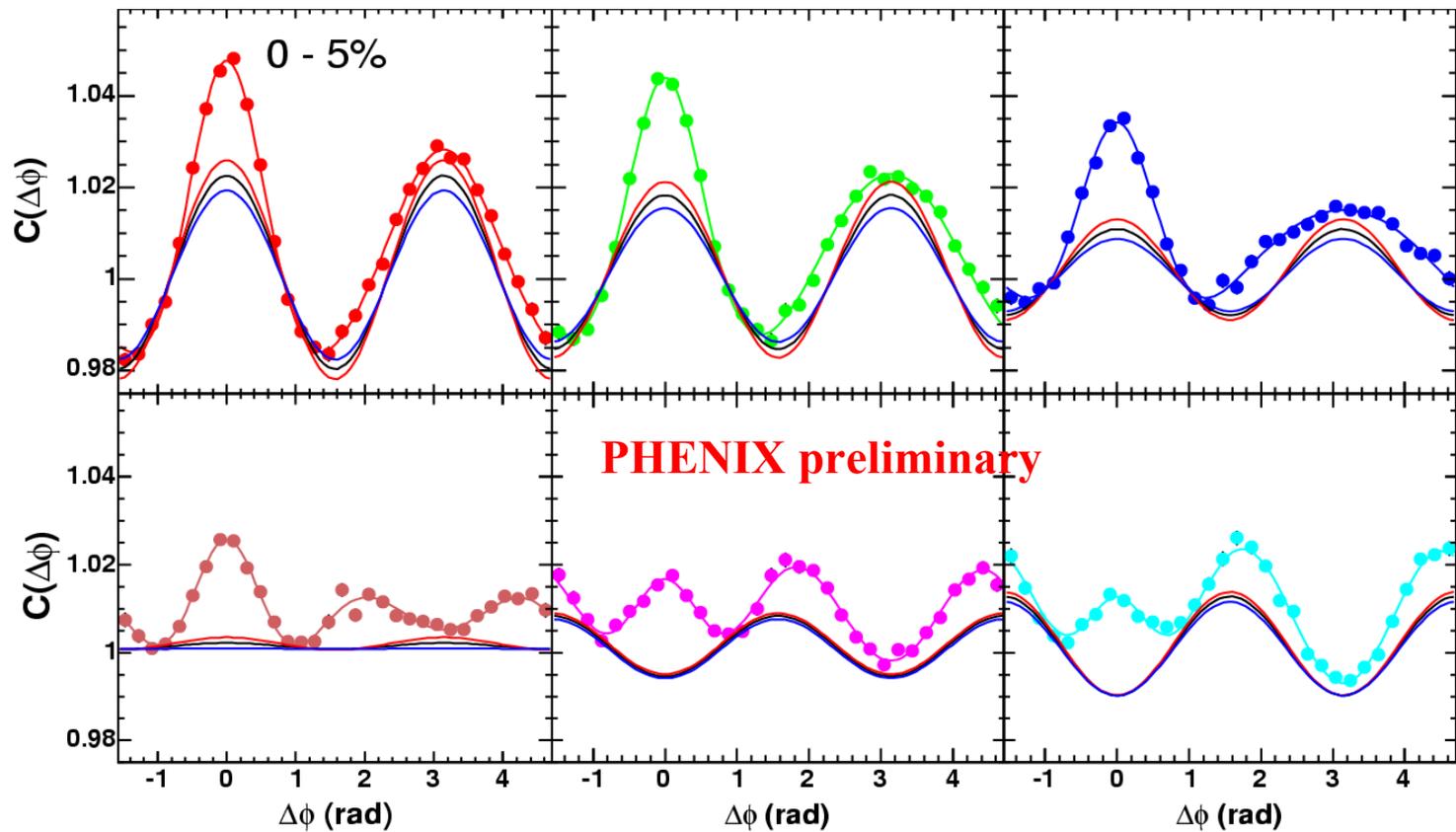
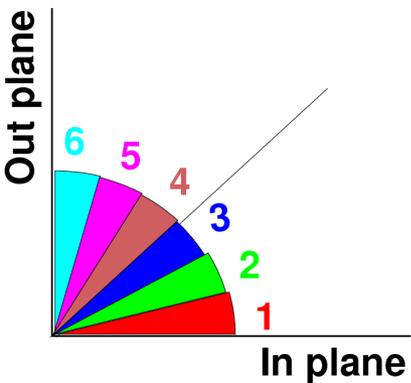
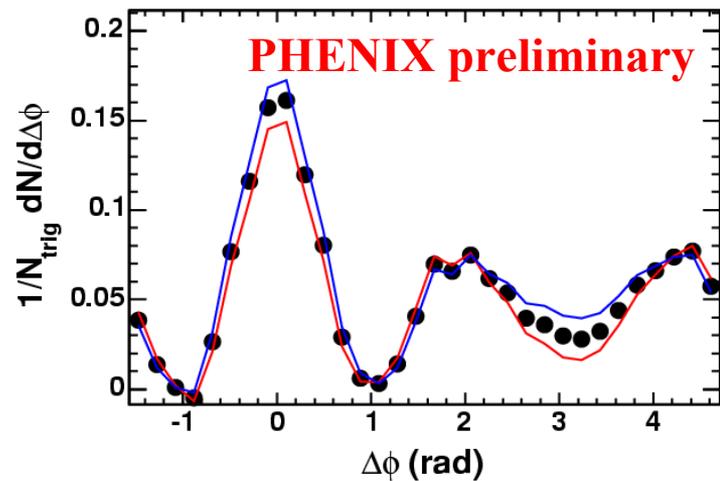
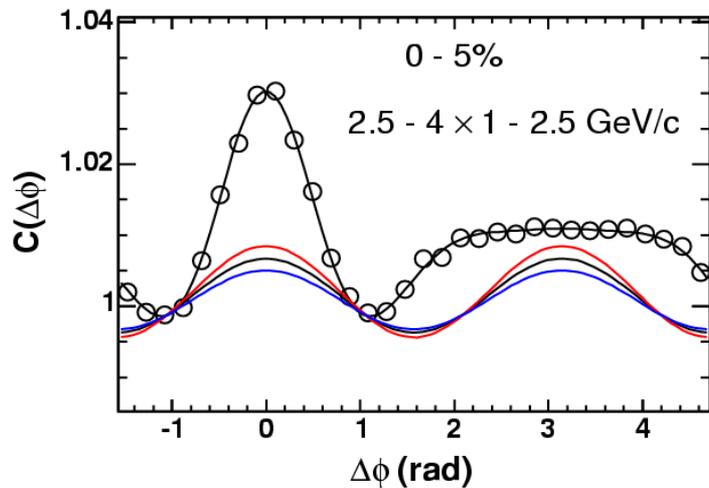
**g radiates energy**  
**kick particles in the plasma**  
**accelerate them along the jet**

**PHENIX preliminary**



$\phi = \pi \downarrow \propto \nearrow 1.23 = 1.91, 4.37 \rightarrow c_s \sim 0.33$   
 ( $\sqrt{0.33}$  in QGP, 0.2 in hadron gas)

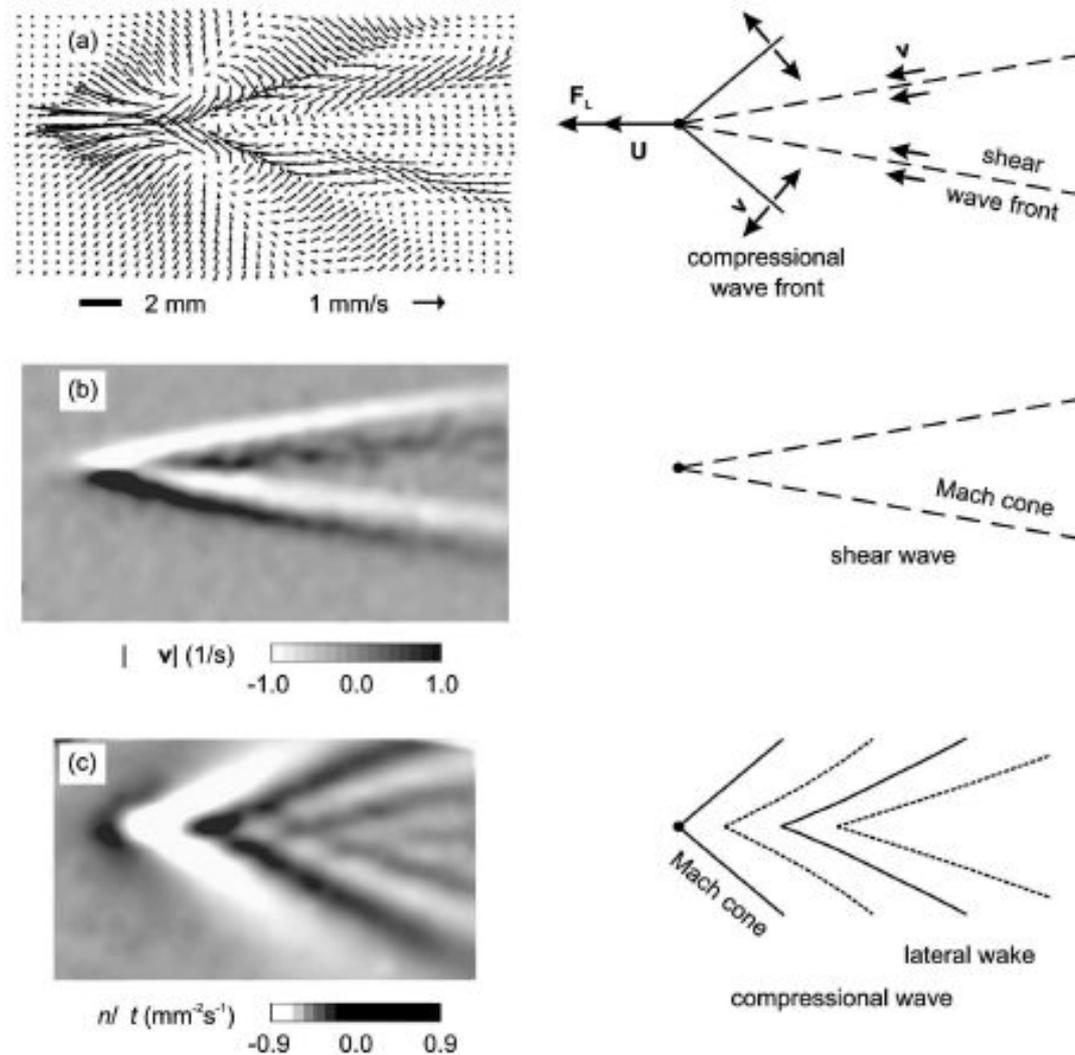
0-5%



# Compressional and shear wakes in a two-dimensional dusty plasma crystal

NOSENKO *et al.*

PHYSICAL REVIEW E 68, 056409 (2003)

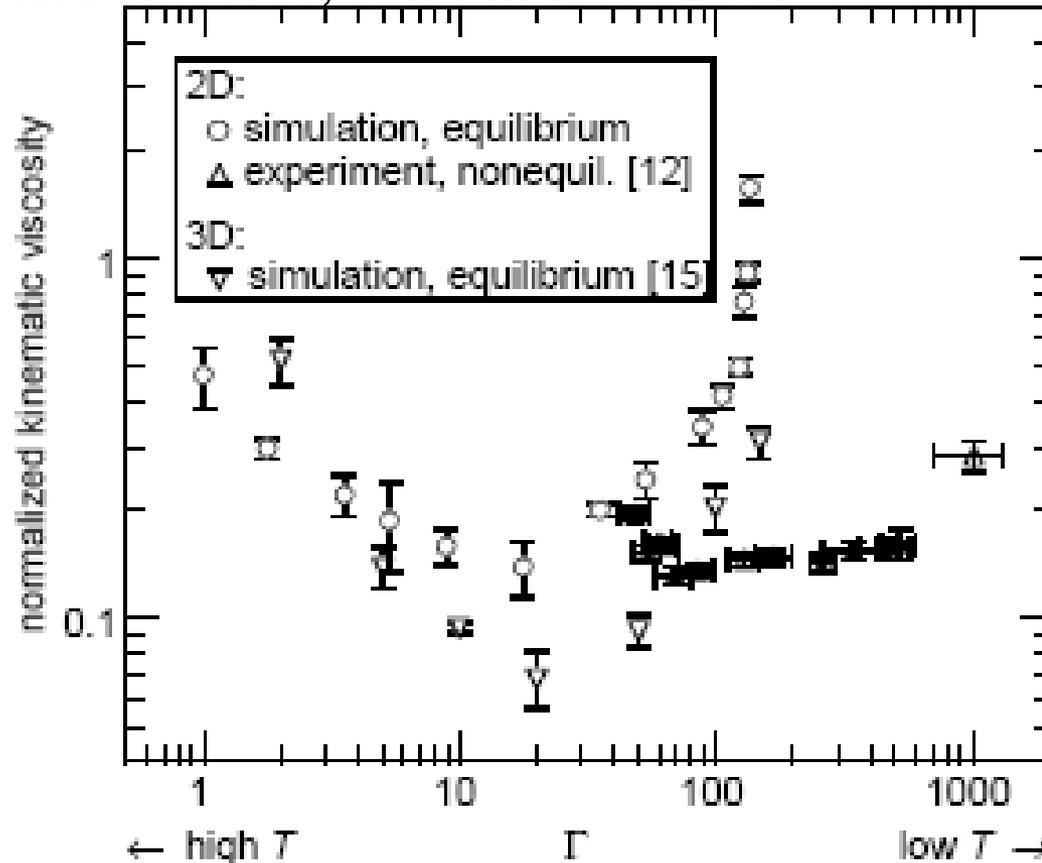


*generally  
a phenomenon  
in crystals but  
not liquids*

FIG. 4. The compressional- and shear-wave Mach cones, excited simultaneously. The scanning speed  $U$  is higher than the sound speed for both the compressional and shear waves. Maps are shown for (a) particle velocity  $\mathbf{v}$ , (b) vorticity  $|\nabla \times \mathbf{v}|$ , and (c)  $\partial n / \partial t$ , where  $n$  is the particle areal number density.

# use to get $\eta$ & compare to molecular dynamics

B. Liu and J. Goree, cond-mat/0502009



*MD: solve the equations of motion for massive particles subject to (screened) interaction potential*

*follow evolution of particle distribution function (& correlations)*

*solve coupled diff.eq's over nearby space*

*density-density correlations →  $\eta$*

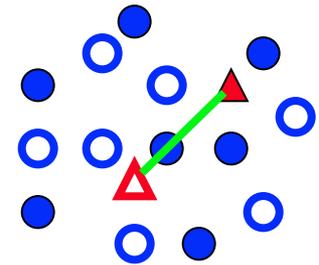
**minimum arises because kinetic part of  $\eta$  decreases with  $\Gamma$  & potential part increases**

# aim to measure the screening length

$J/\Psi$  (bound state of  $c$  and  $c$ bar quarks)

*Tests screening & confinement:*

*do bound  $c + c$  survive the medium?  
or does QGP screening kill them?*



*Look at  $R_{AA}$  for  $J/\psi$*

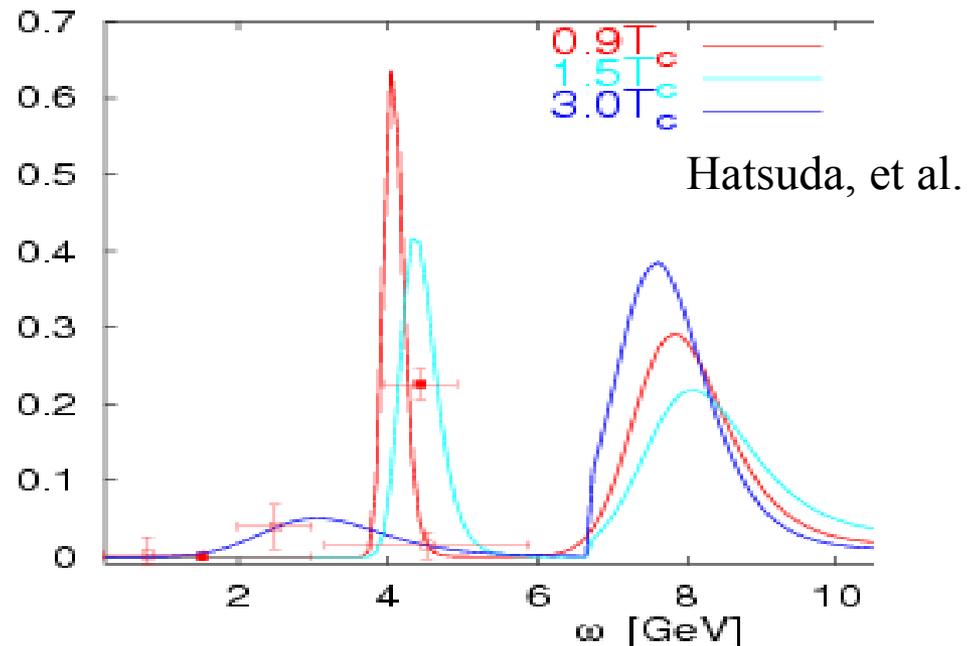
*different bound states*

*probe different lengths*

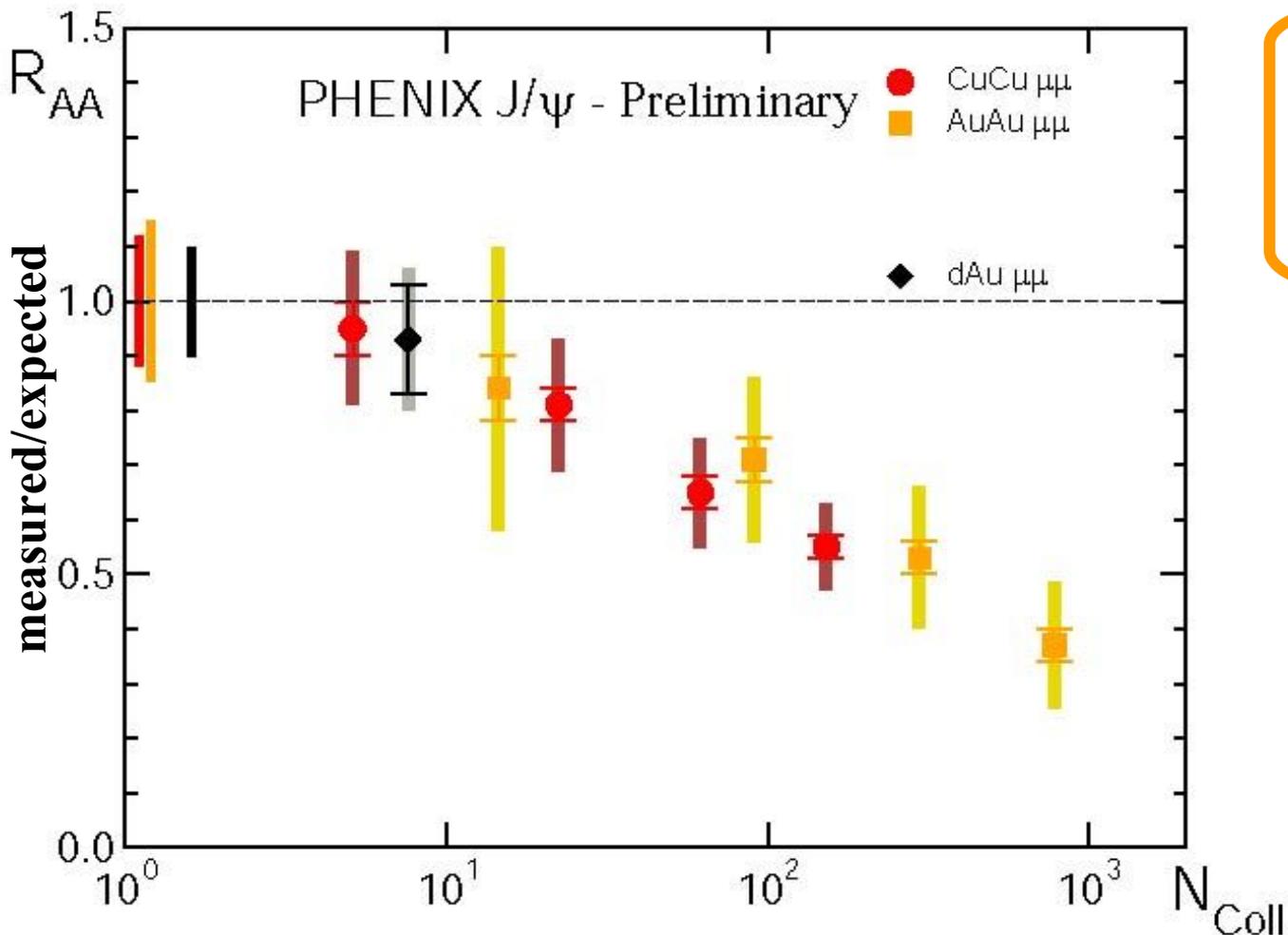
*note:*

*it's not so clear*

*what to expect...*



# At RHIC:



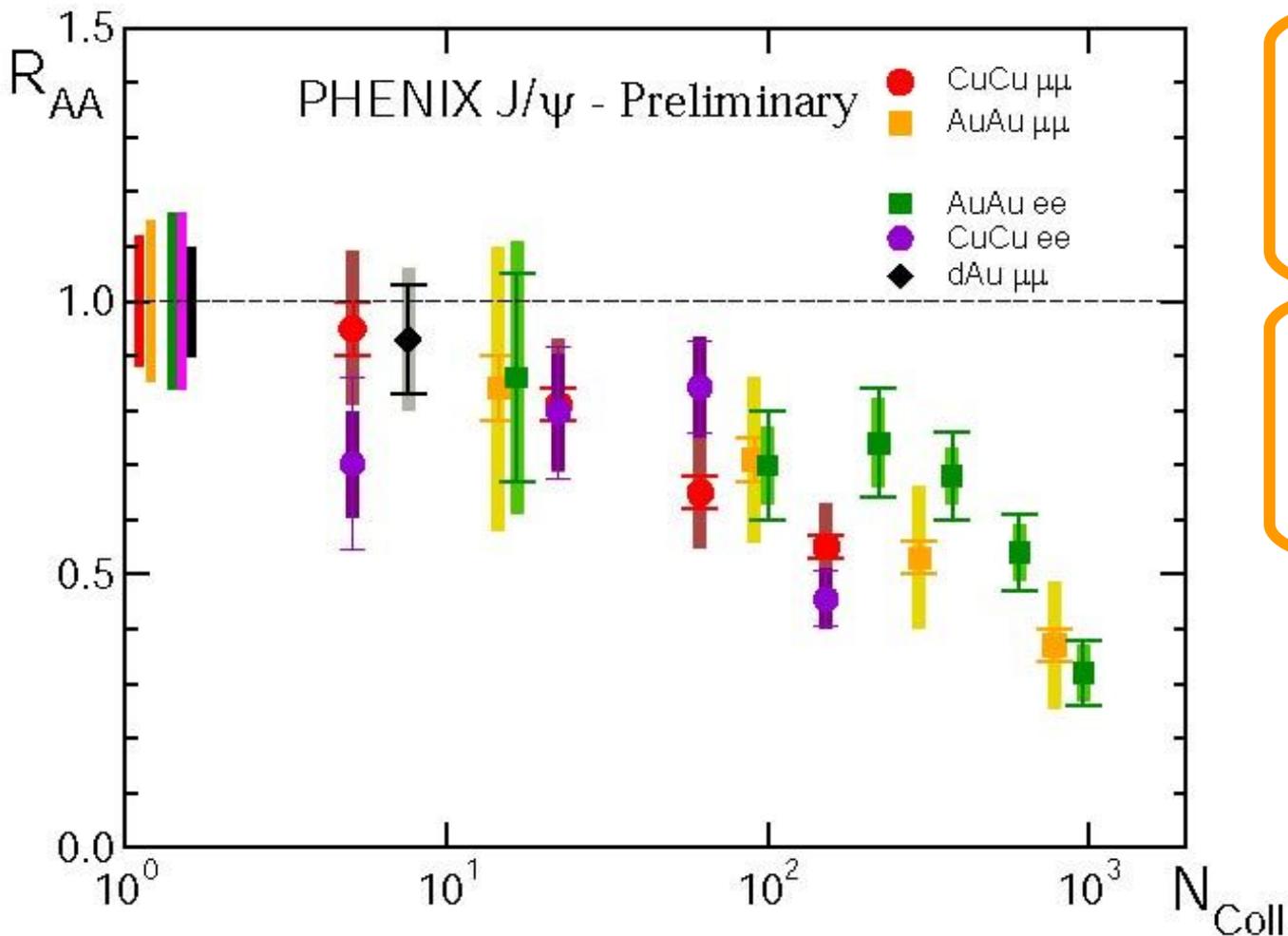
$J/\psi \rightarrow \mu\mu$   
muon arm  
 $1.2 < |y| < 2.2$

dAu  
 $\mu\mu$   
200 GeV/c

AuAu  
 $\mu\mu$   
200 GeV/c

CuCu  
 $\mu\mu$   
200 GeV/c

# At RHIC:



$J/\psi \rightarrow \mu\mu$   
muon arm  
 $1.2 < |y| < 2.2$

$J/\psi \rightarrow ee$   
Central arm  
 $-0.35 < y < 0.35$

dAu  
 $\mu\mu$   
200 GeV/c

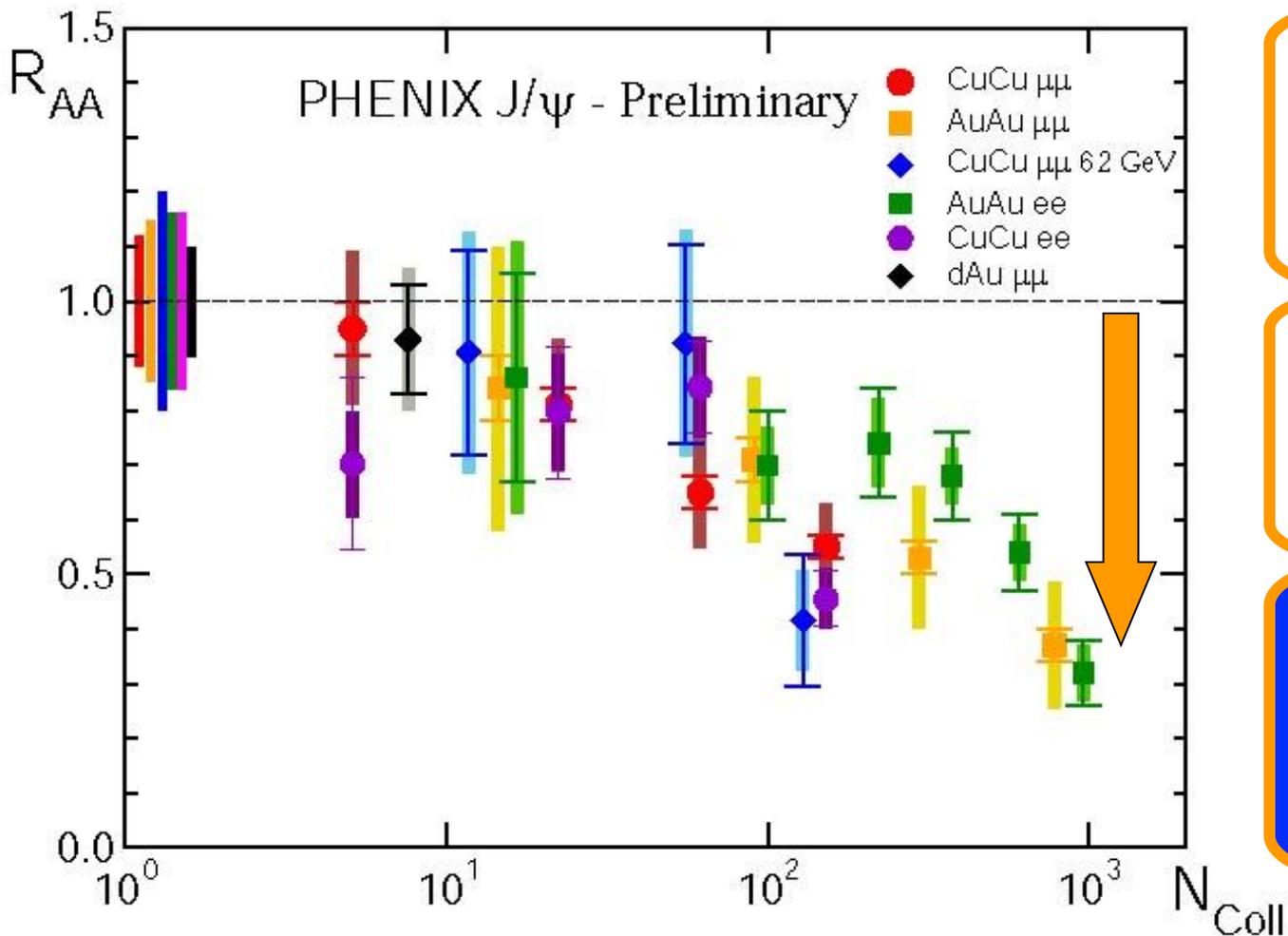
AuAu  
 $\mu\mu$   
200 GeV/c

CuCu  
 $\mu\mu$   
200 GeV/c

AuAu  
 $ee$   
200 GeV/c

CuCu  
 $ee$   
200 GeV/c

# At RHIC:



$J/\psi \rightarrow \mu\mu$   
muon arm  
 $1.2 < |y| < 2.2$

$J/\psi \rightarrow ee$   
Central arm  
 $-0.35 < y < 0.35$

Factor  $\sim 3$   
suppression  
in central events

dAu  
 $\mu\mu$   
200 GeV/c

AuAu  
 $\mu\mu$   
200 GeV/c

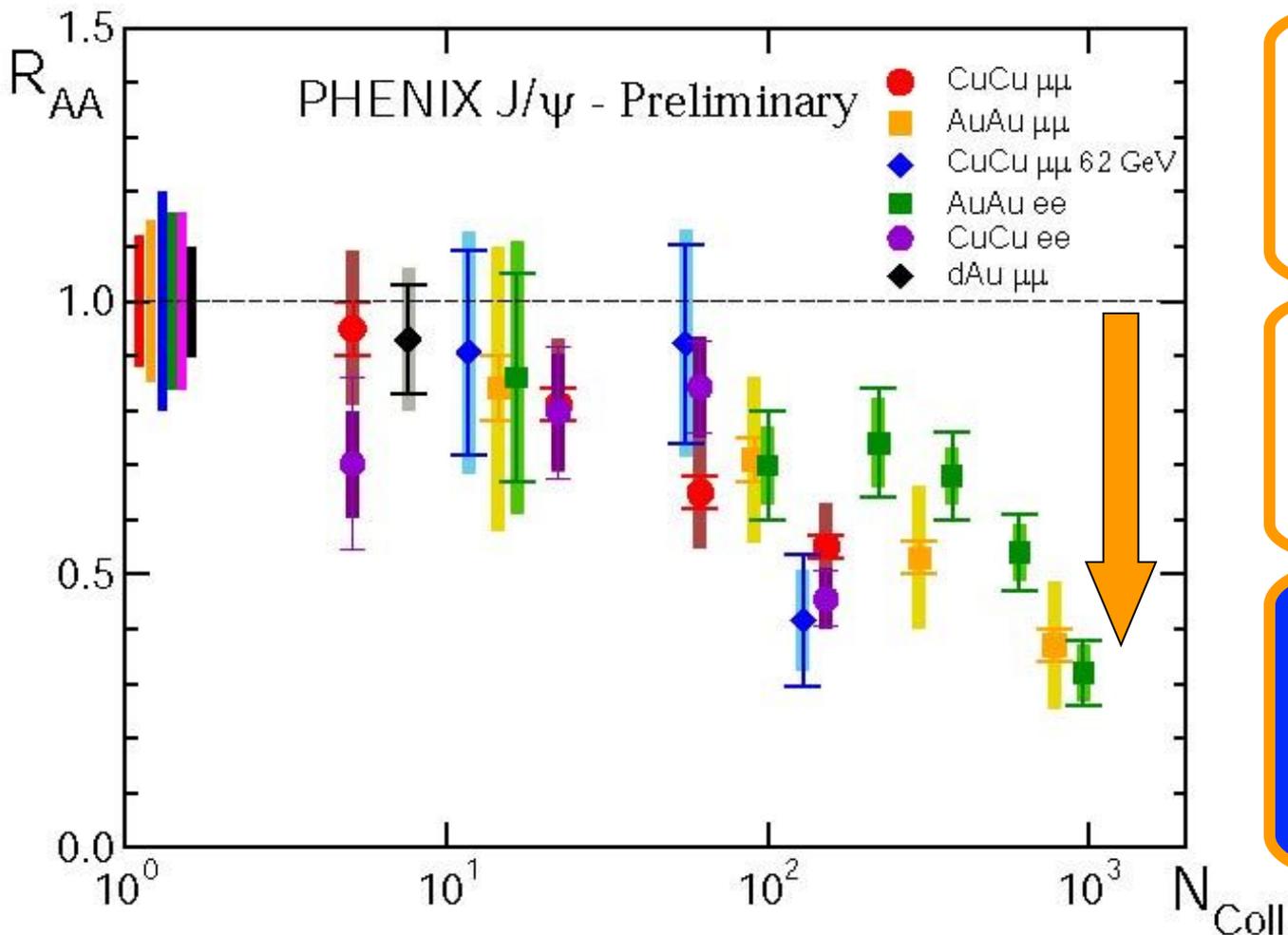
CuCu  
 $\mu\mu$   
200 GeV/c

AuAu  
 $ee$   
200 GeV/c

CuCu  
 $ee$   
200 GeV/c

CuCu  
 $\mu\mu$   
62 GeV/c

# At RHIC:



$J/\psi \rightarrow \mu\mu$   
muon arm  
 $1.2 < |y| < 2.2$

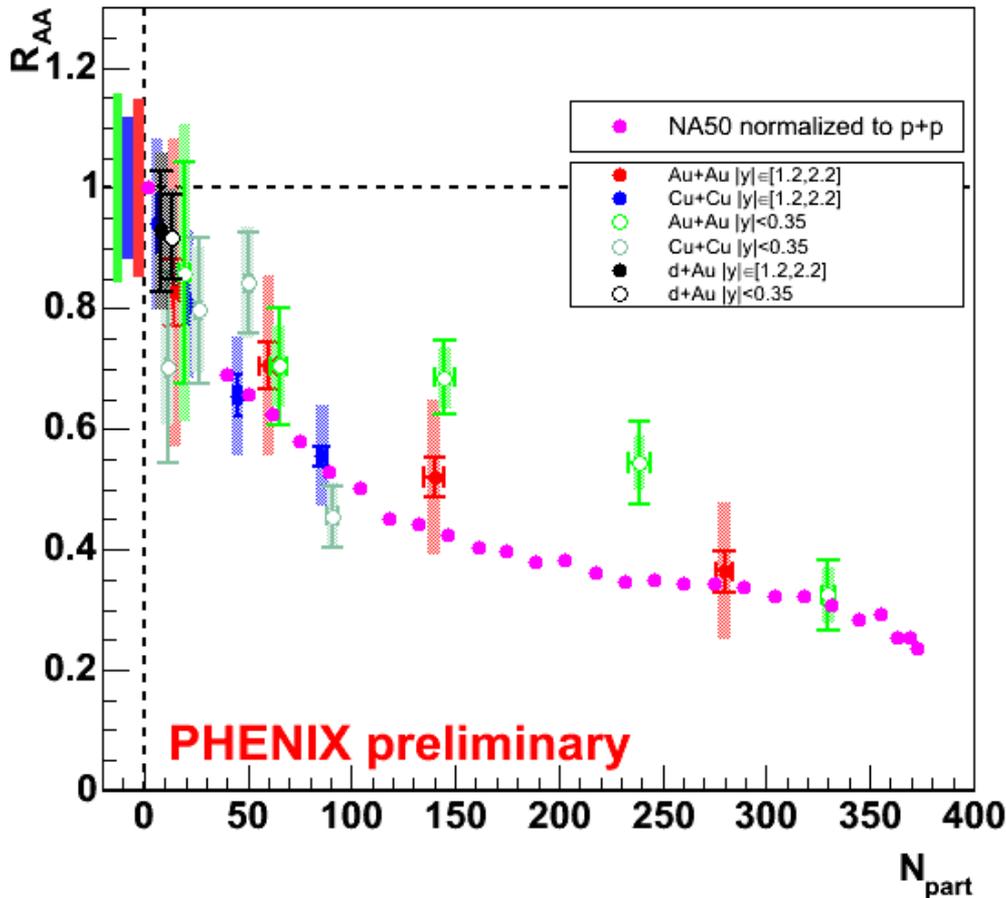
$J/\psi \rightarrow ee$   
Central arm  
 $-0.35 < y < 0.35$

Factor  $\sim 3$   
suppression  
in central events

Data show the same trend within errors  
for all beams and even at  $\sqrt{s}=62$  GeV

# $R_{AA}$ vs $N_{part}$ : PHENIX and NA50

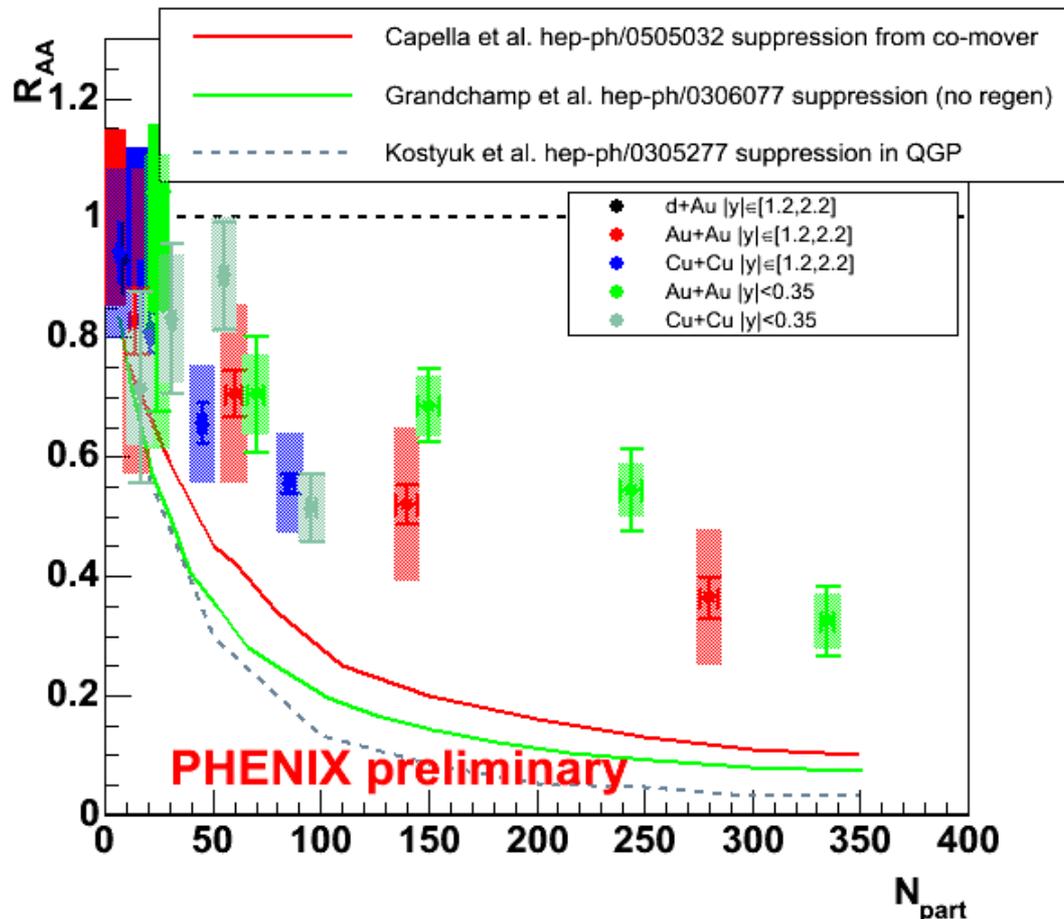
## $J/\psi$ nuclear modification factor $R_{AA}$



- NA50 data normalized at NA50 p+p point.
- Suppression similar in the two experiments, although the collision energy is 10 times higher (200GeV in PHENIX & 17GeV in NA50)

# What suppression should we expect?

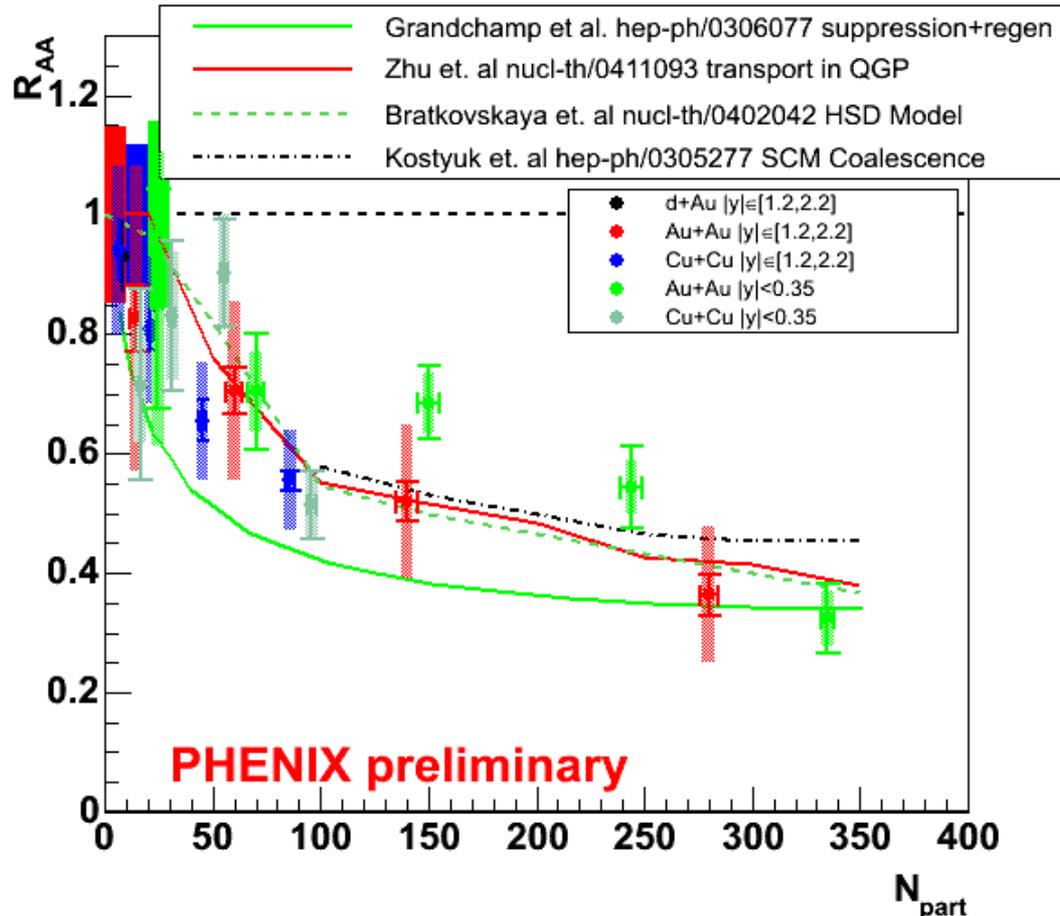
## J/ψ nuclear modification factor $R_{AA}$



Models that were successful in describing SPS data  
fail to describe data at RHIC  
- but lattice QCD says bound states until  $\sim 2T_c$  -

# can get better agreement with data

## $J/\psi$ nuclear modification factor $R_{AA}$



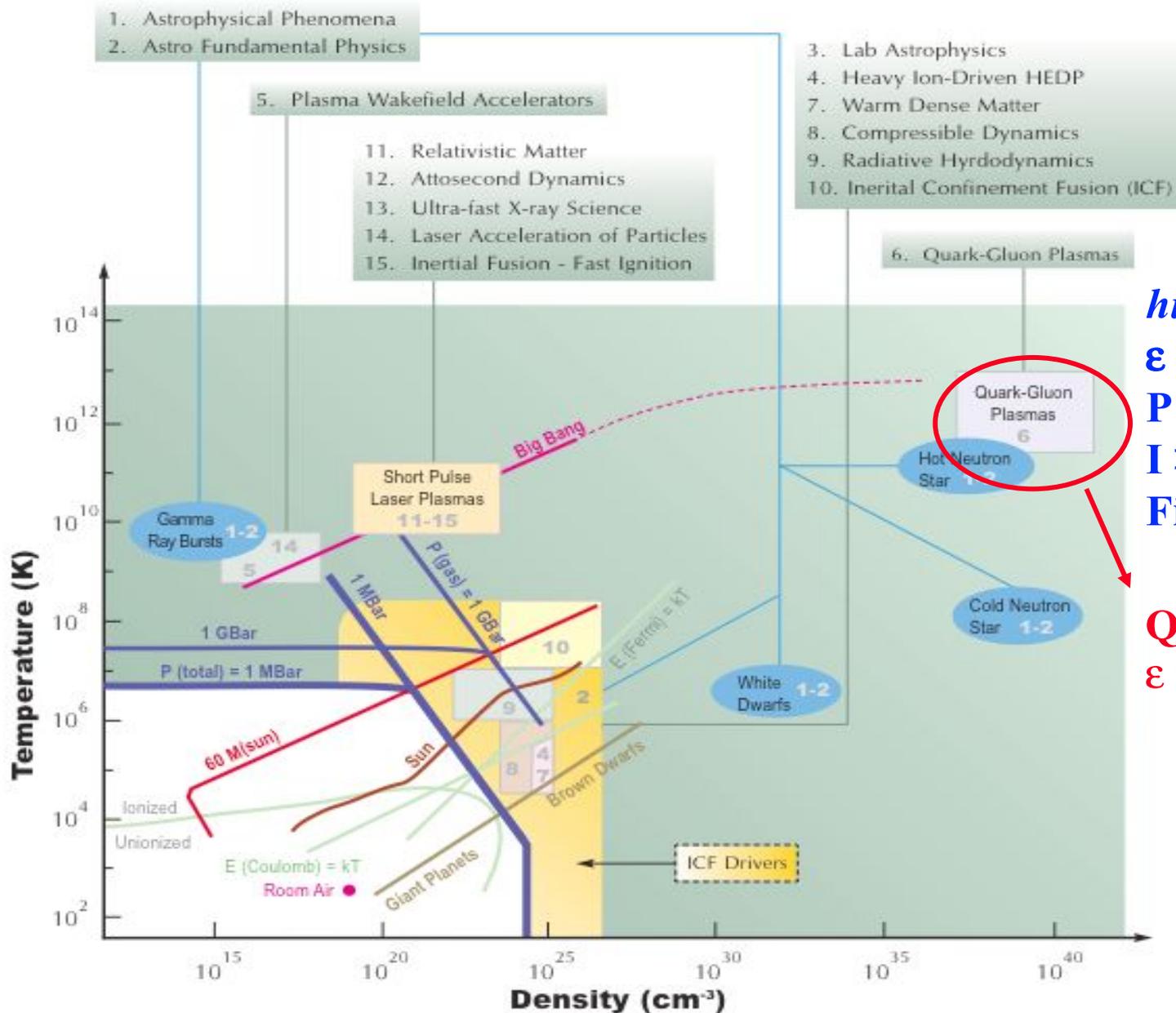
if add formation of “extra”  $J/\psi$  by coalescence of  
c and anti-c from the plasma  
caveat: not necessarily unique or correct explanation!

# conclusions

---

- the matter formed at RHIC is a “perfect” fluid
  - shows collective flows with small viscosity
  - huge interaction cross sections, very opaque
  - multiple collisions affect even heavy charm quarks
- this is like other strongly coupled plasmas
  - as it *should* be → it *is* a plasma!
  - neutrality scale > interparticle distance
- How does this super high energy density plasma work?
  - map properties of the new stuff at RHIC
    - how does the plasma transport the “lost” energy?
    - radiation rate?
    - initial temperature achieved? (theory says ~380 MeV)
  - collide Pb+Pb at the LHC for higher  $T_{\text{initial}}$ 
    - reach ~ 800 MeV: is coupling strong or weak?

# Energy density of matter



**high energy density:**  
 $\epsilon > 10^{11} \text{ J/m}^3$   
 $P > 1 \text{ Mbar}$   
 $I > 3 \times 10^{15} \text{ W/cm}^2$   
 Fields > 500 Tesla

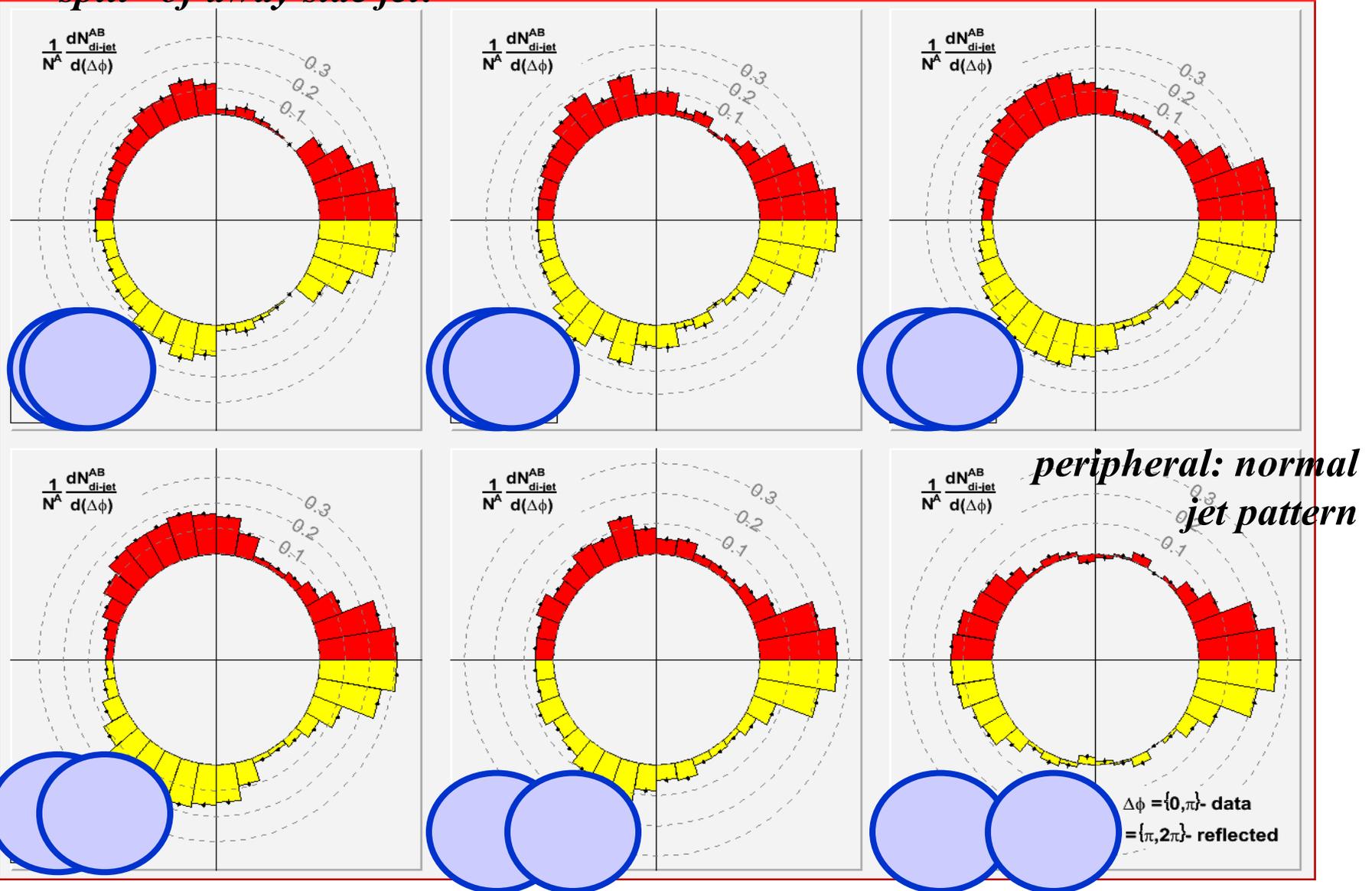
**QGP energy density**  
 $\epsilon > 1 \text{ GeV/fm}^3$   
 i.e.  $> 10^{30} \text{ J/cm}^3$

- 
- **backup slides**

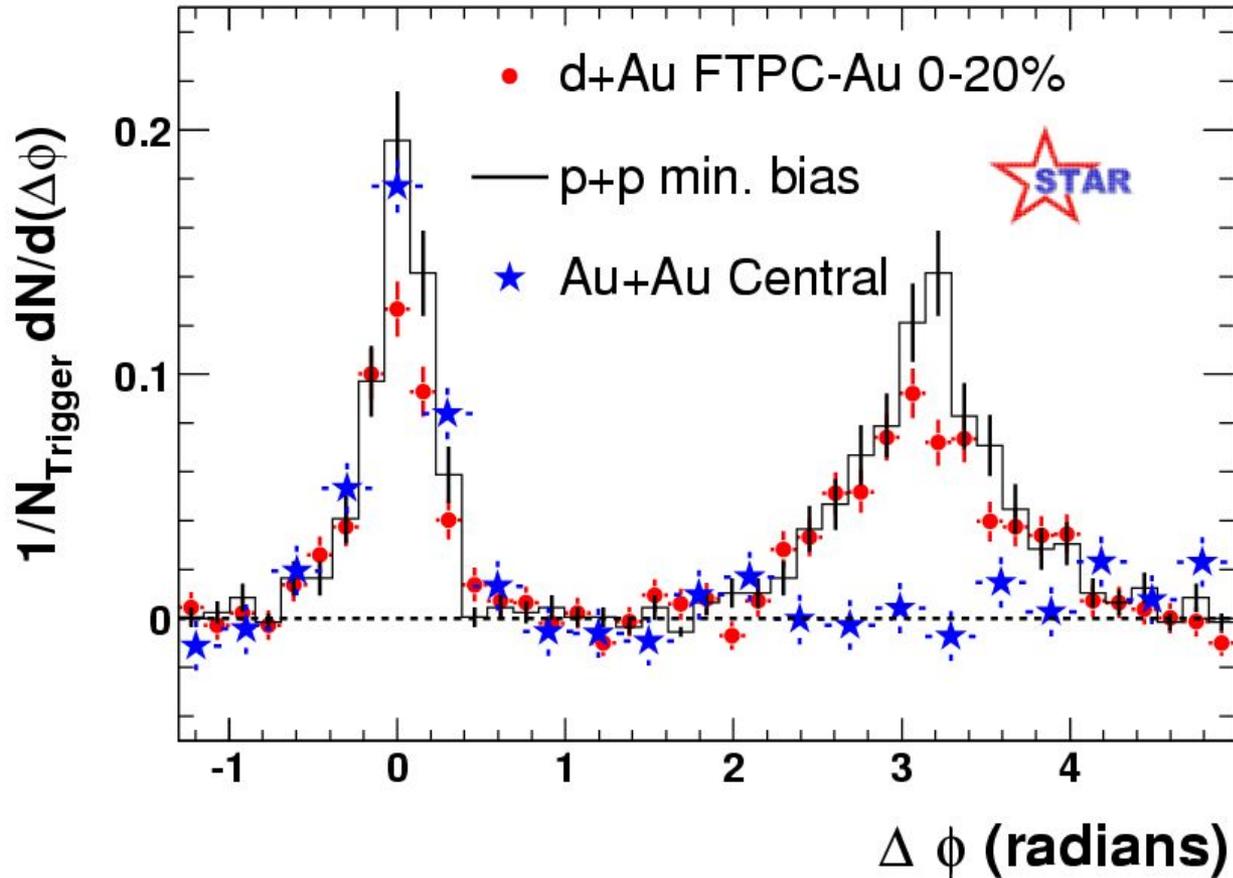
# Angular Correlation of $\sim 1\text{-}2$ GeV/c jet fragments

PHENIX (nucl-ex/0507004)

*“split” of away side jet!*



# Are back-to-back jets there in d+Au?



**Yes!**

**no medium**



**no jet  
quenching**

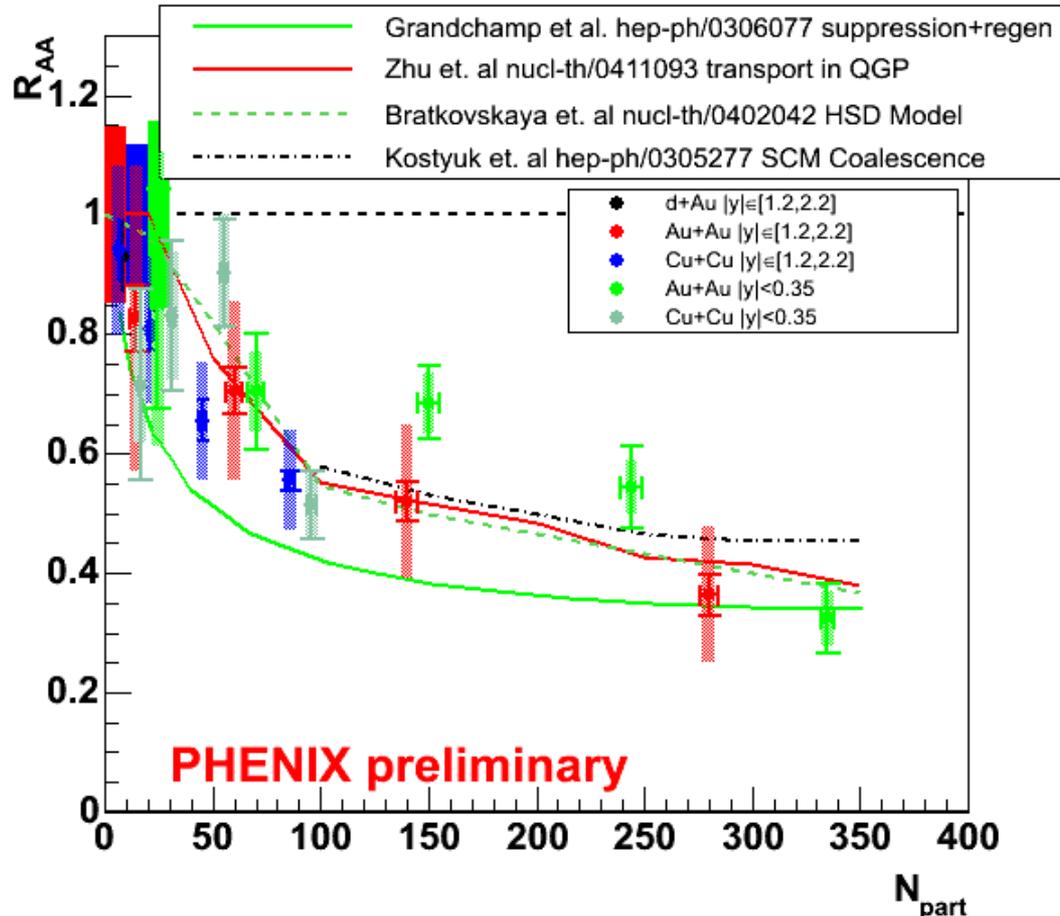
# plasma properties known, so far

## Extract from models, constrained by data

Energy loss $\langle dE/dz \rangle$ (GeV/fm)	7-10	0.5 in cold matter
Energy density (GeV/fm <sup>3</sup> )	14-20	>5.5 from $E_T$ data <i>above hadronic E density!</i>
$dN(\text{gluon})/dy$	$\sim 1000$	From energy loss, hydro <i>huge!</i>
T (MeV)	380-400	<i>Experimentally unknown as yet</i>
Equilibration time $\tau_0$ (fm/c)	0.6	From hydro initial condition; cascade agrees <i>very fast!</i> <i>NB: plasma folks have same problem &amp; use same technique</i>
Opacity (L/mean free path)	3.5	Based on energy loss theory

# can get better agreement with data

## $J/\psi$ nuclear modification factor $R_{AA}$



if add formation of “extra”  $J/\psi$  by coalescence of  
c and anti-c from the plasma  
caveat: not necessarily unique or correct explanation!

# A little more on coupling

potential  $V \propto \alpha_s/r$   $\langle KE \rangle \propto T$   $r$ =interparticle distance

QCD matter:  $\rho \propto 1/r^3$   $\rho \propto T^3$  and so we see that  $r \propto 1/T$

- $\Gamma = \langle PE \rangle / \langle KE \rangle \propto (\alpha_s/r)/T \propto \alpha_s T/T \propto \alpha_s$

$T$  cancels, but does affect  $\alpha_s$

- $\lambda_D = \{T/(4\pi\epsilon_0 e^2 \rho)\}^{1/2}$  so  $\lambda_D \propto \{T/(\alpha_s T^3)\}^{1/2} \propto 1/(T\alpha_s^{1/2})$   
 $\alpha_s$

- We know  $1/\Gamma \propto$  #particles inside Debye volume  $N_D$

- $N_D = N/V_D = \rho V_D$   $V_D = 4/3 \pi \lambda_D^3 \propto 1/(\alpha_s^{3/2} T^3)$   
so  $N_D \propto 1/\alpha_s^{3/2}$   $T$  cancels again

- for  $\alpha_s$  large,  $N_D$  is small ( $\lambda_D$  fairly small, but included in  $N_D$ )  
for  $\alpha_s$  small,  $N_D$  is large ( $\lambda_D$  largish)

# putting in some numbers

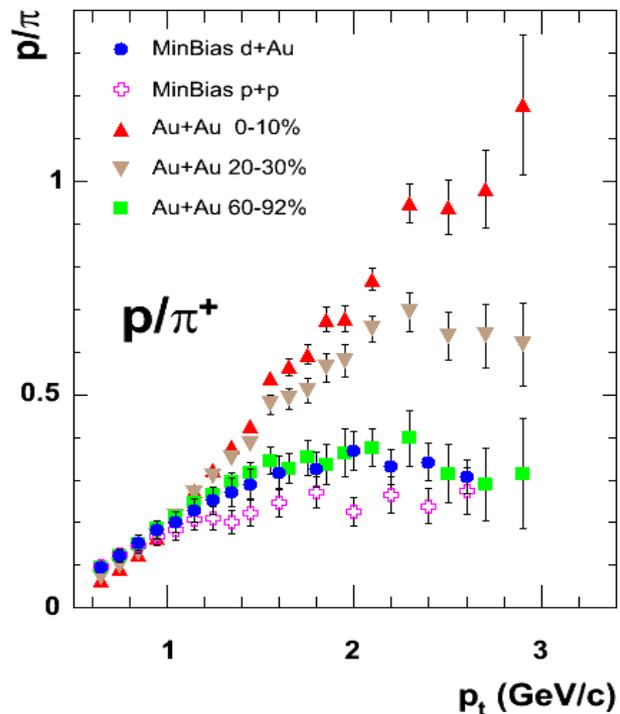
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- both  $\Gamma$  and  $N_D$  depend on  $\alpha_s$
- at RHIC  $dN_g/dy \sim 800$   
so  $\rho = 800 / (1 \text{ fm} * \pi R^2 \text{ fm}^2) = 800 / 100 = 8 \text{ /fm}^3$   
 $r = 0.5$
- from lattice at  $T \sim 200 \text{ MeV}$   $\alpha_s = 0.5-1$  for quarks  
for gluons multiply by  $3/(4/3) = 9/4$ . It's big!
- from pQCD  $\alpha_s = 0.3$  for quarks and  $\sim 0.7$  for gluons

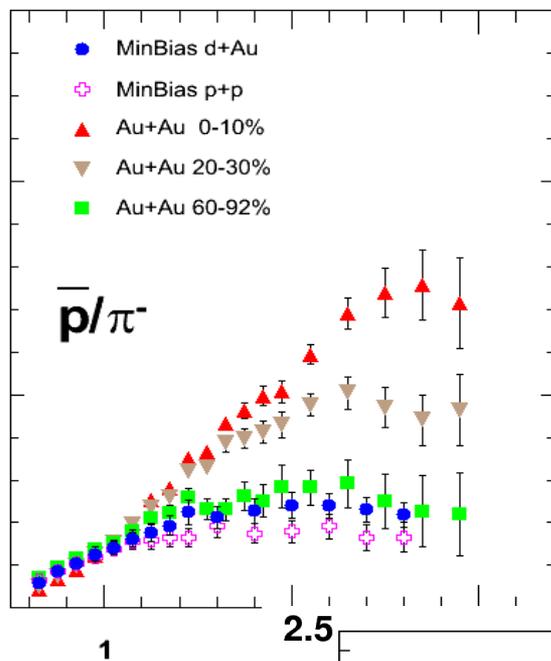
# baryon puzzle...

PHENIX

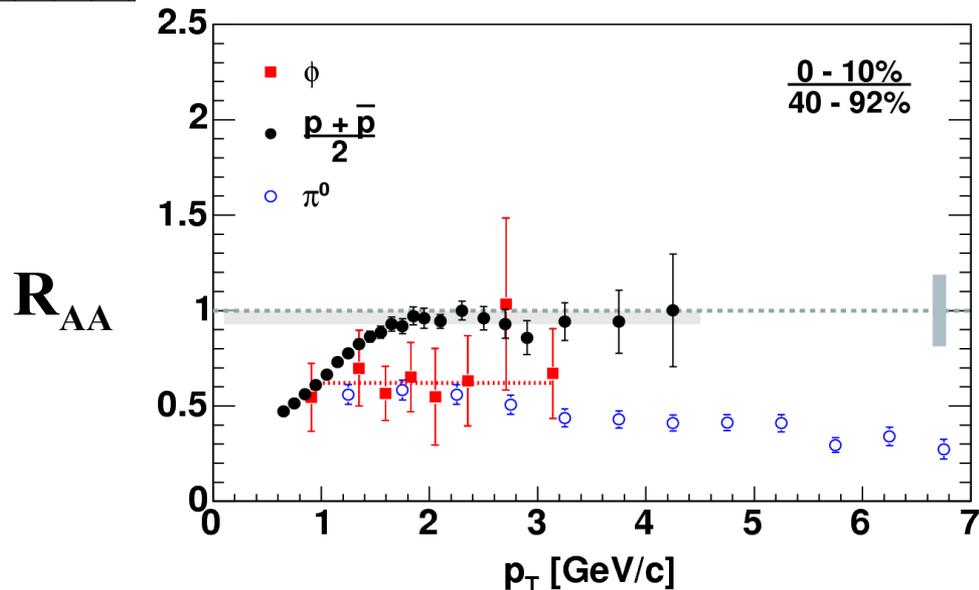
$\sqrt{s} = 200$  GeV



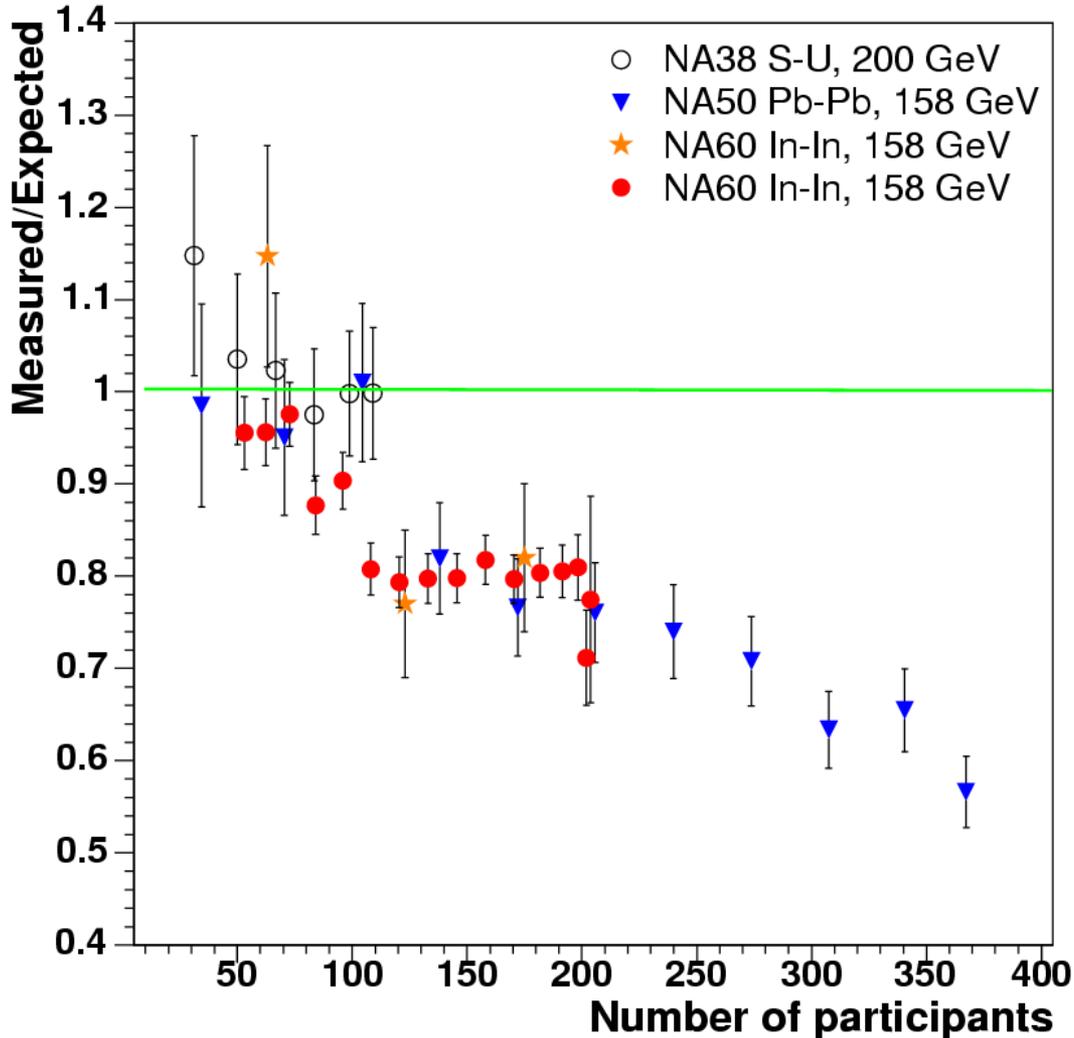
Collision Systems: p+p, d+Au, Au+Au



baryons enhanced for  $p_T < 5$  GeV/c



# At CERN ( $\sqrt{s} = 17$ GeV):



□ NA50 and NA60 show suppression in Pb+Pb & In+In

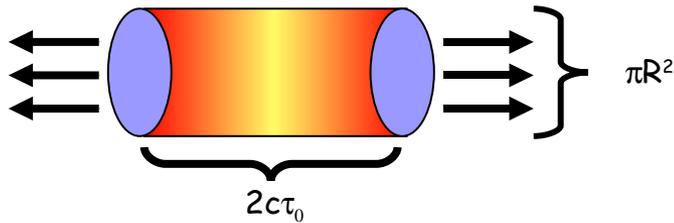
□ suppression follows system size

□ Normal nuclear absorption from p+A data:  $\sigma = 4.18 \pm 0.35$  mb

# Is the energy density high enough?

PRL87, 052301 (2001)

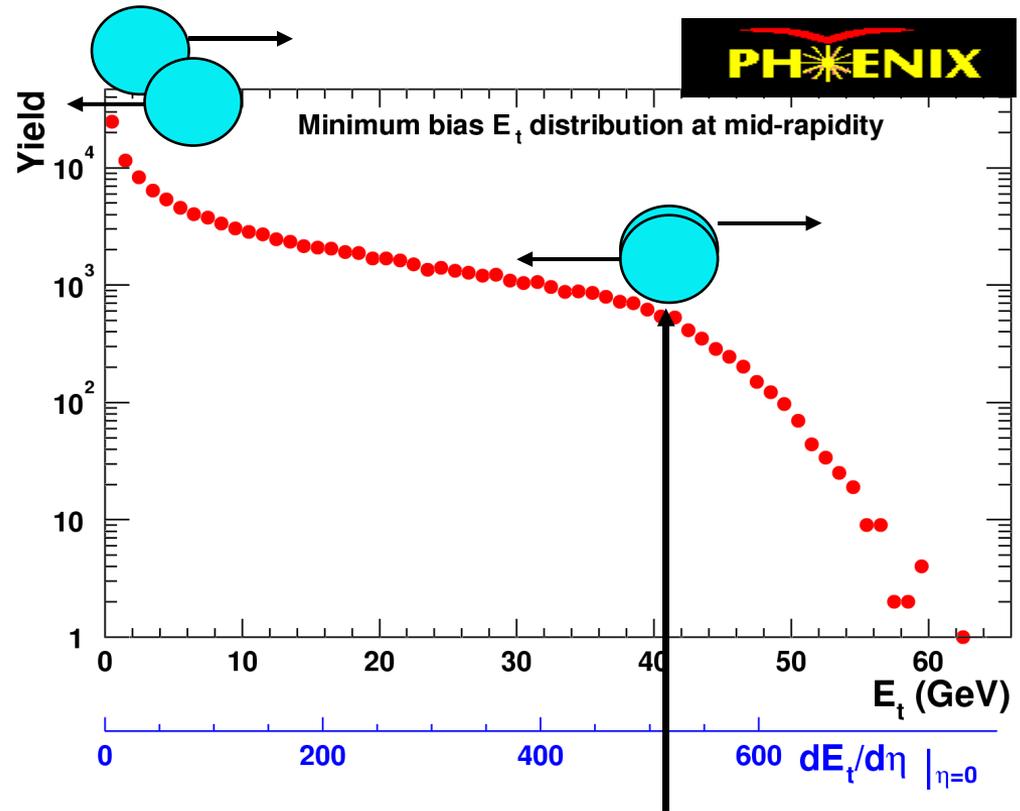
Colliding system expands:



Energy  $\perp$  to  
beam direction  $\downarrow$

$$\varepsilon_{Bj} = \frac{1}{\pi R^2} \frac{1}{2c\tau_0} \left( 2 \frac{dE_T}{dy} \right)$$

per unit  
velocity  $\parallel$  to beam  $\nearrow$



$\rightarrow \varepsilon \geq 5.5 \text{ GeV}/\text{fm}^3$  (200 GeV Au+Au)  
*well above predicted transition!*

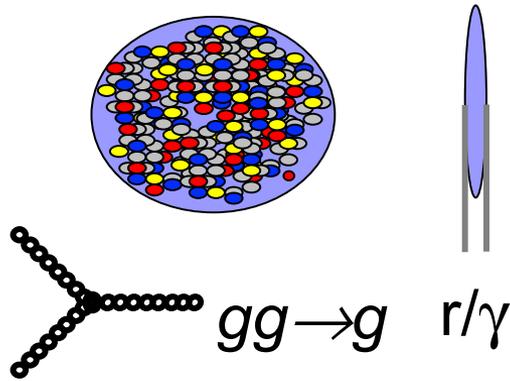
value is lower limit:

longitudinal expansion rate, formation time overestimated

# Saturation of gluons in initial state

## (colored glass condensate)

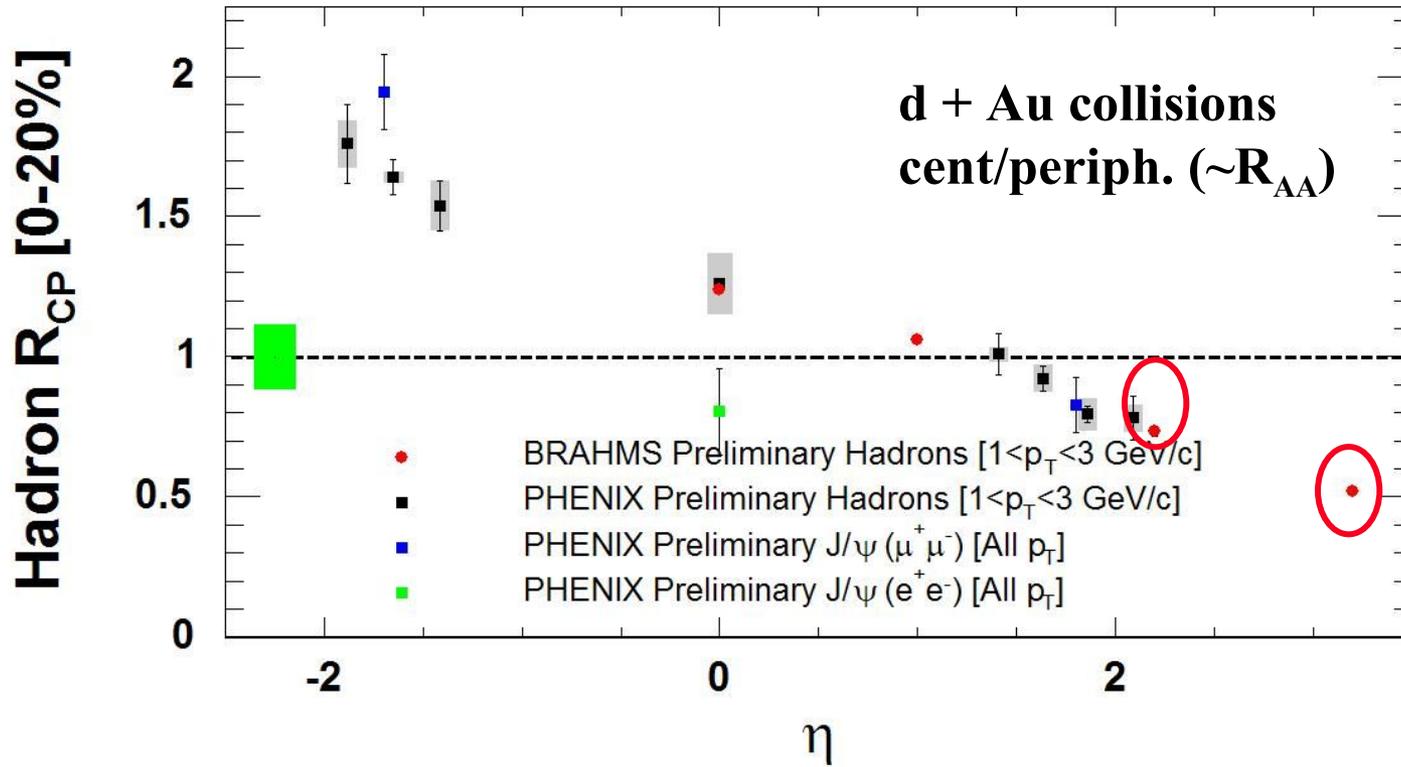
Mueller, McLerran, Kharzeev, ..



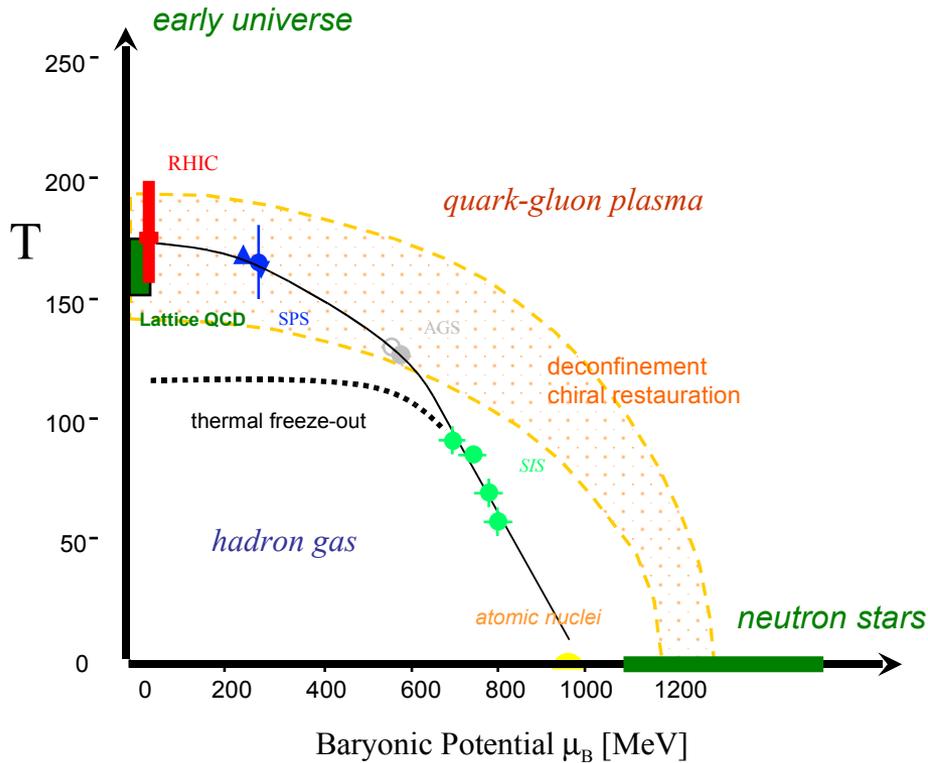
Wavefunction of low x (very soft) gluons overlap and the self-coupling gluons fuse.

Saturation at higher x at RHIC vs. HERA due to nuclear size

→ suppressed jet cross section; no back-back pairs



# Locate RHIC on phase diagram



From fit of yields vs. mass  
(grand canonical  
ensemble):

$$T_{\text{ch}} = 176 \text{ MeV}$$

$$\mu_B = 41 \text{ MeV}$$

These are the conditions  
when hadrons stop  
interacting

Observed particles “freeze out” at/near the deconfinement boundary!

# Why no energy loss for charm quarks?

---

- “dead cone” predicted by Kharzeev and Dokshitzer, Phys. Lett. B519, 199 (1991)
- Gluon bremsstrahlung:  
 $k_T^2 = \mu^2 t_{\text{form}}/\lambda$     transverse momentum of radiated gluon  
 $\mu = p_T$  in single scatt.  $\lambda =$  mean free path  
 $\theta \sim k_T / \omega$      $\omega =$  gluon energy
- But radiation is suppressed below angles  $\theta_0 = M_q/E_q$

soft gluon distribution is

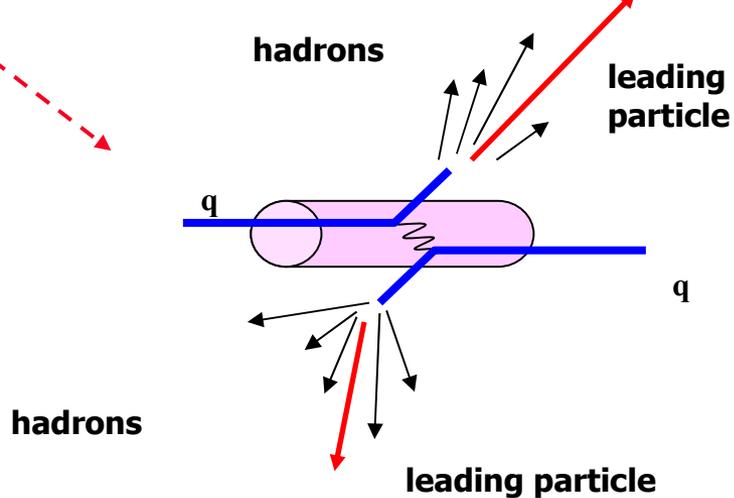
$$dP = \alpha_s C_F / \pi d\omega / \omega k_T^2 dk_T^2 / (k_T^2 + \omega^2 \theta_0^2)^2$$

not small for  
heavy quarks!  
causes a dead cone

# “external” probes of the medium

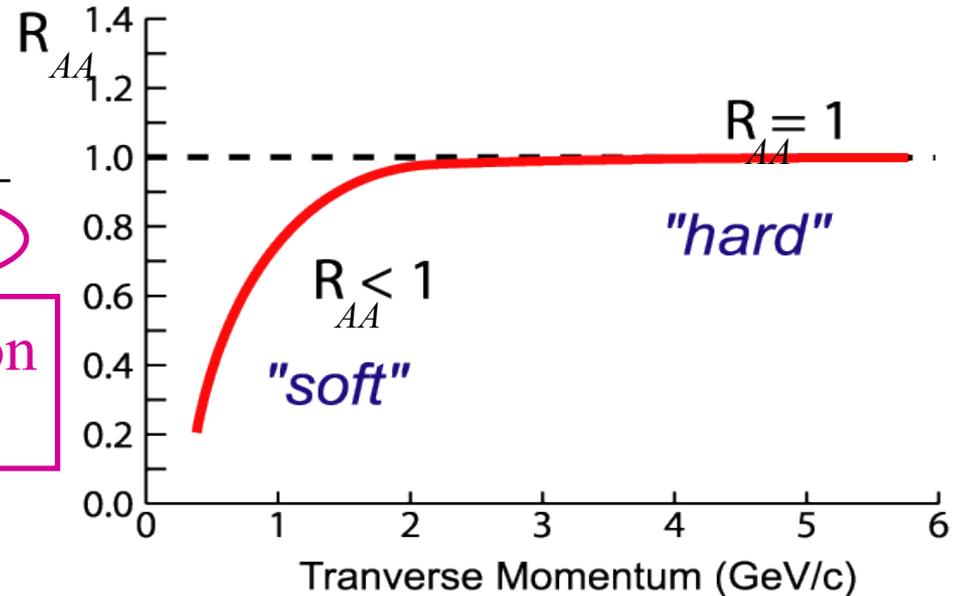
**Hard scattering of q,g early.**  
**Observe fast leading particles,**  
**back-back correlations**  
**Before creating hadron jets,**  
**scattered quarks induced to**  
**radiate energy ( $\sim \text{GeV}/\text{fm}$ )**  
**by the colored medium**  
**-> jet quenching**

schematic view of jet production



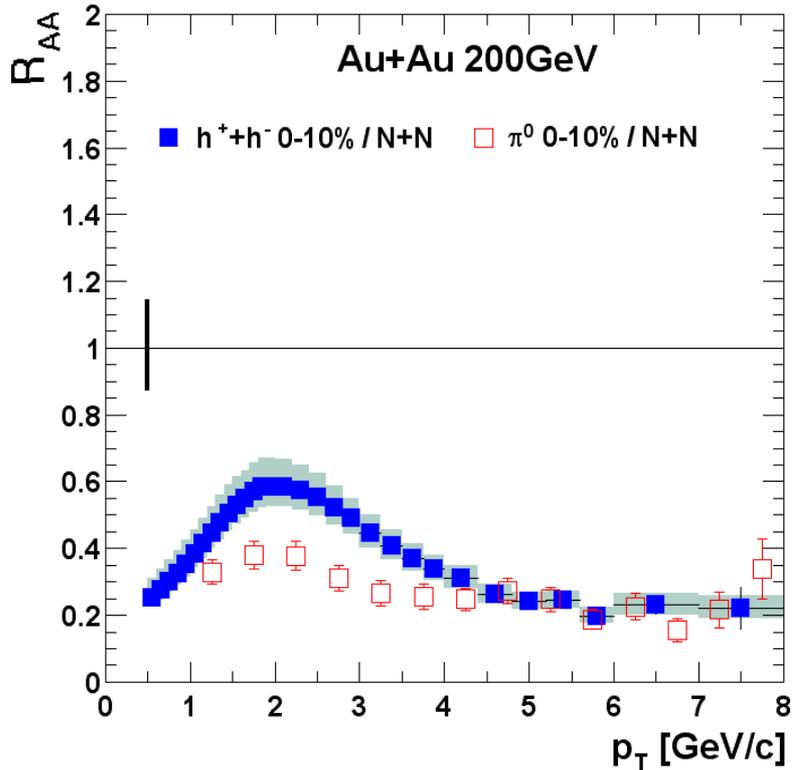
$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{\langle N_{\text{binary}} \rangle / \sigma_{\text{inel}}^{p+p} T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$

$\langle N_{\text{binary}} \rangle / \sigma_{\text{inel}}^{p+p}$  (blue text)  
 $T_{AA}$  (circled in blue)  
 $d^2 \sigma^{NN} / dp_T d\eta$  (circled in pink)  
 nucleon-nucleon cross section (pink box)

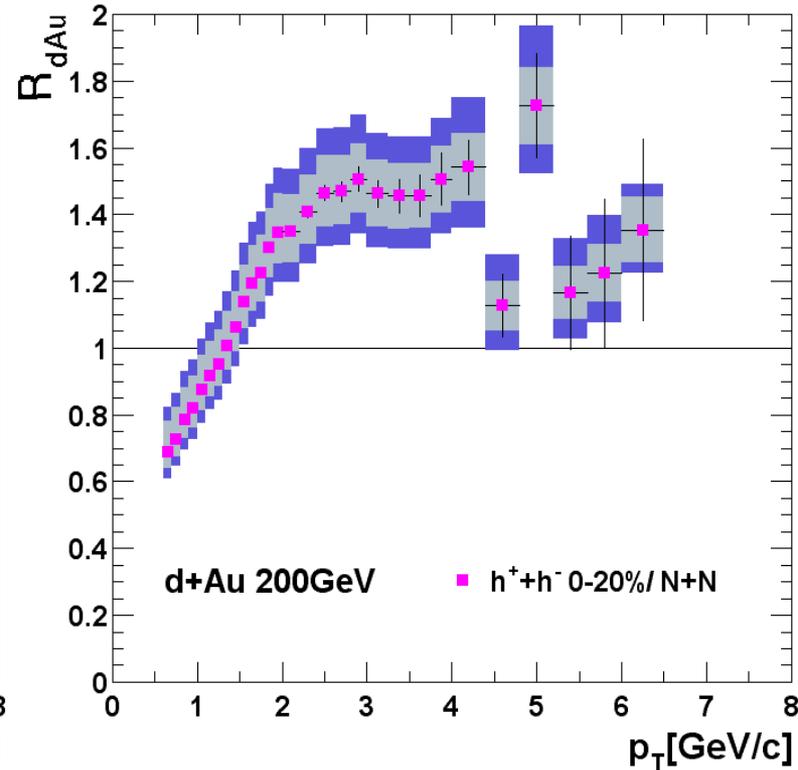


# Centrality Dependence

## Au + Au Experiment



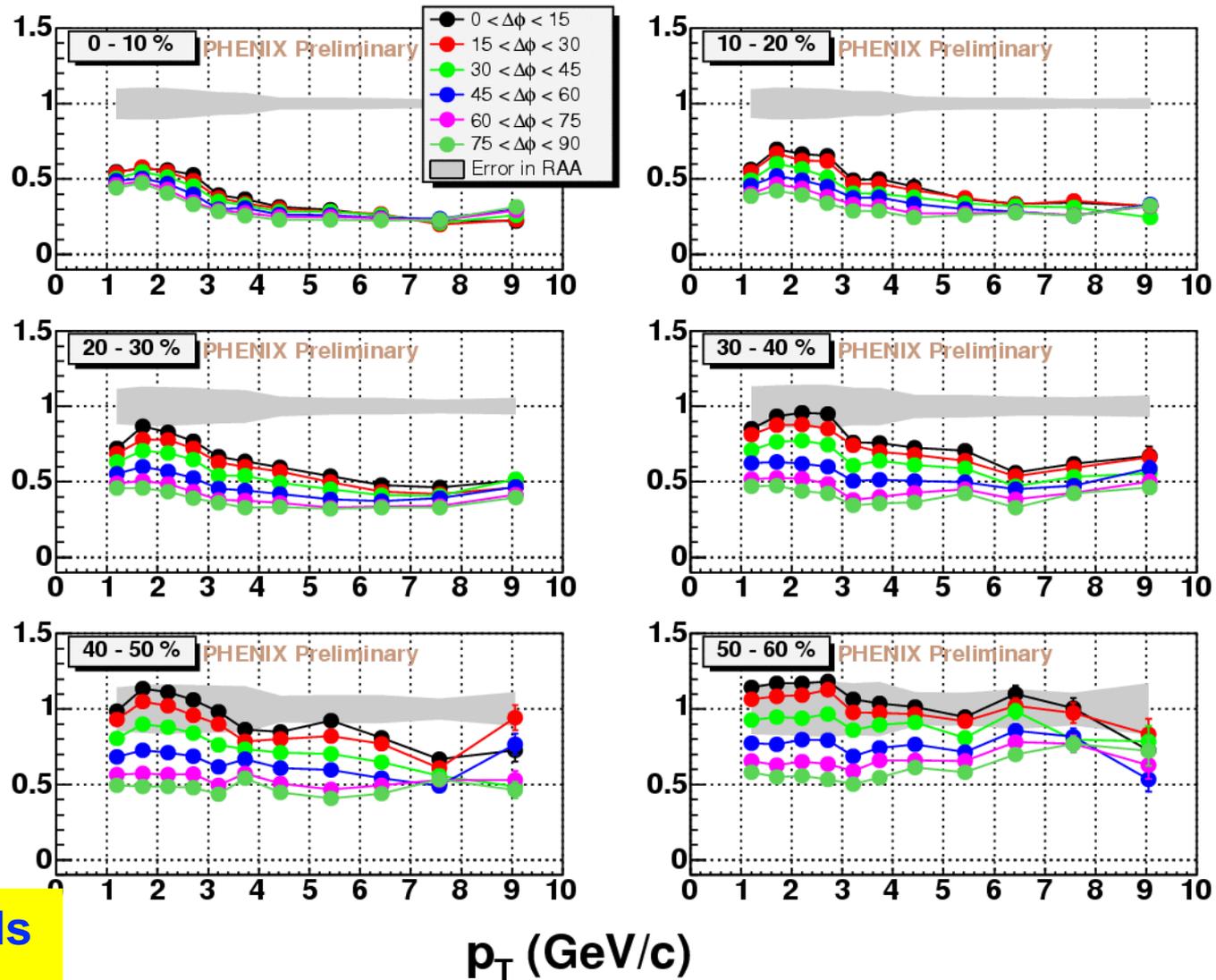
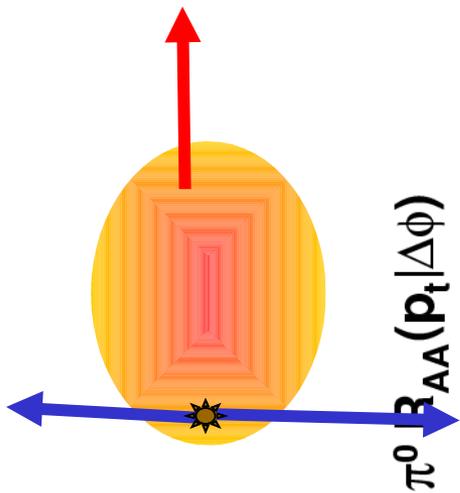
## d + Au Control



PHENIX  
preliminary

- Dramatically different and opposite centrality evolution of AuAu experiment from dAu control.
- Jet Suppression is clearly a final state effect.

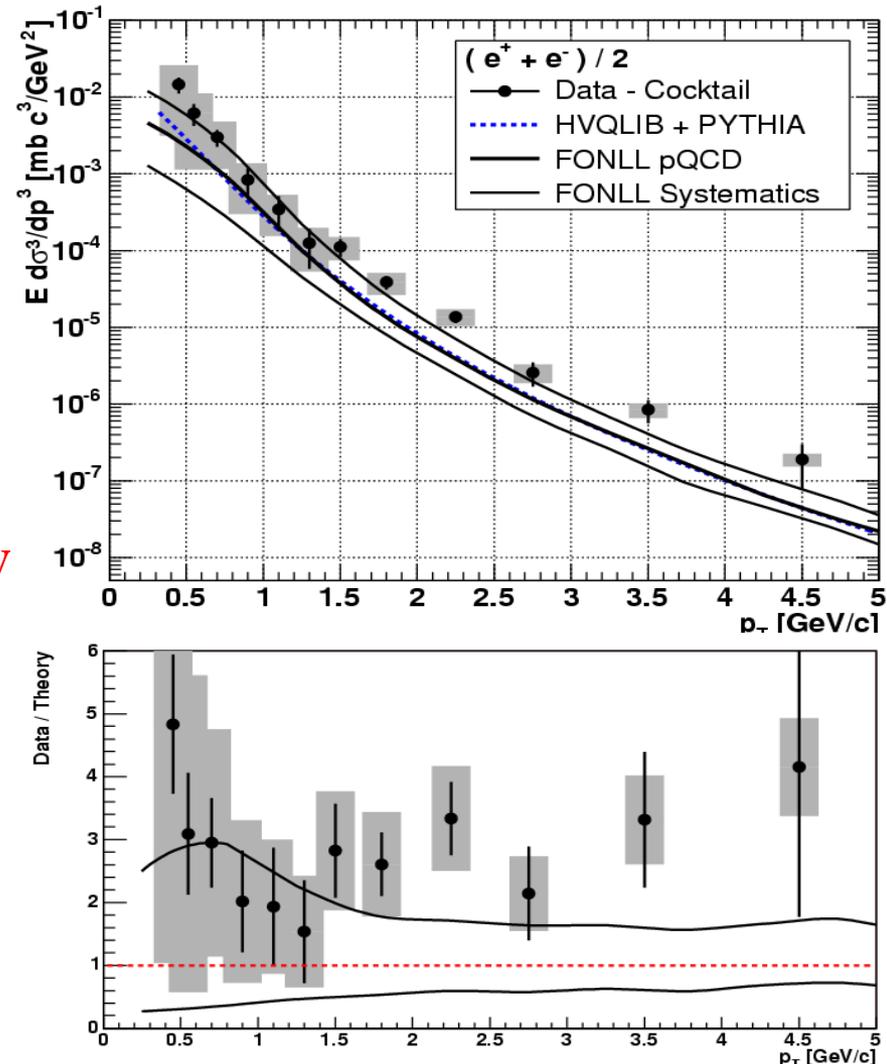
# $R_{AA}$ wrt reaction plane – more evidence



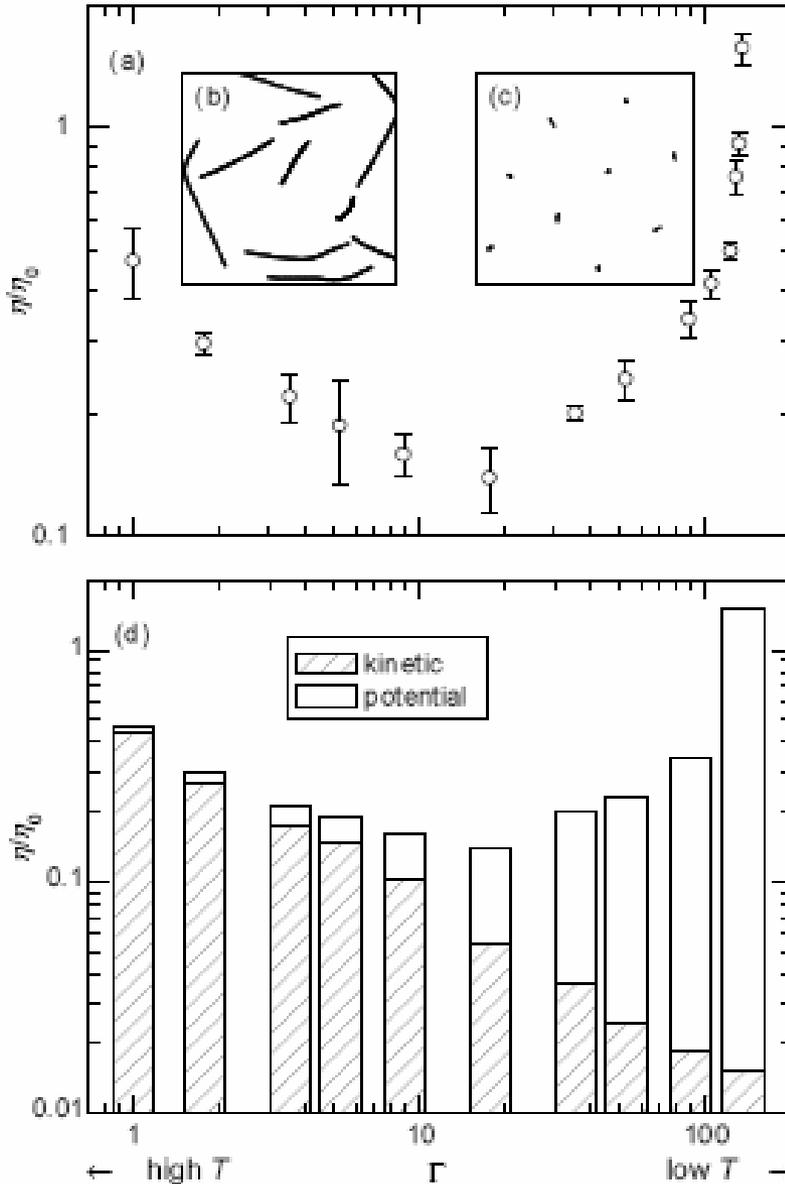
Energy loss depends on the path-length

# FONLL Predictions

- **Mateo Cacciari provided a prediction using the Fixed Order Next Leading Logarithm pQCD approach**
- **His calculation agrees perfectly with our “poor man’s” HVQLIB+PYTHIA predictions**
- **Data exceed the central theory curve by a factor of 2-3**
- **Possible explanations:  
NNLO contribution  
Fragmentation mechanisms  
need to be studied in more details**



# use this technique to measure viscosity



melt crystal with laser light  
induce a shear flow (laminar)  
image the dust to get velocity  
study:

spatial profiles  $v_x(y)$

moments, fluctuations  $\rightarrow T(x,y)$

curvature of velocity profile

$\rightarrow$  drag forces

viscous transport of drag in

$\perp$  direction from laser

compare to viscous hydro.

extract  $\eta/\rho$

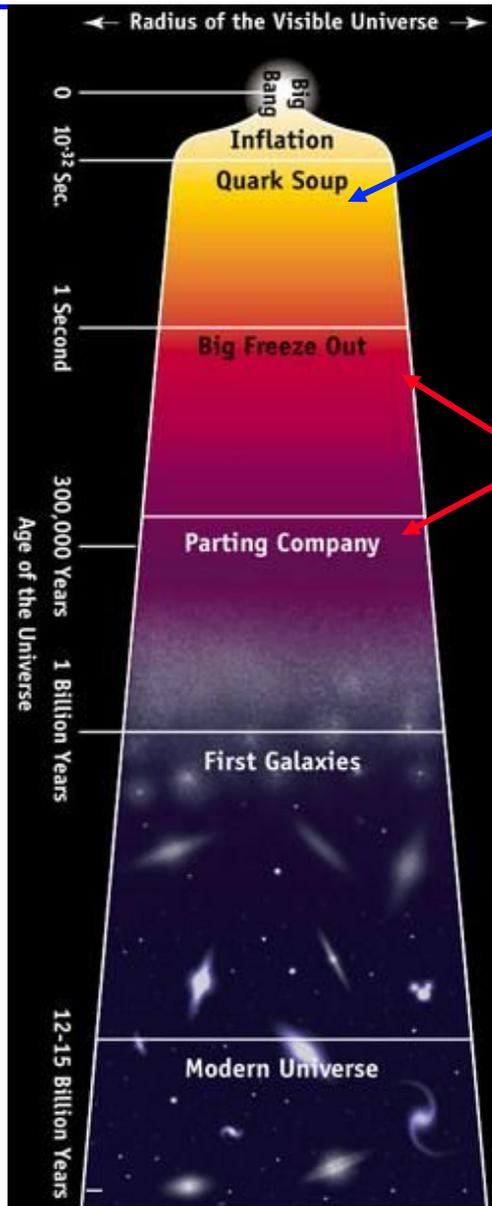
shear viscosity/mass density

**PE vs. KE competition governs**

**coupling & phase of matter**

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# conclusions



The stuff making up the universe ~1us after the Big Bang is **NOT** an ideal gas!  
It's a strongly coupled, low viscosity plasma

Does (should) cosmology (ists) care??

**NO!**

equilibrium → we can't see back that far

**YES!** A first order phase transition is not the right physics

**MAYBE...** don't know enough about QGP behavior to know how it matters

**This generation of experimenters & theorists will figure out the answer!**