

QCD and dynamical phase transitions, I

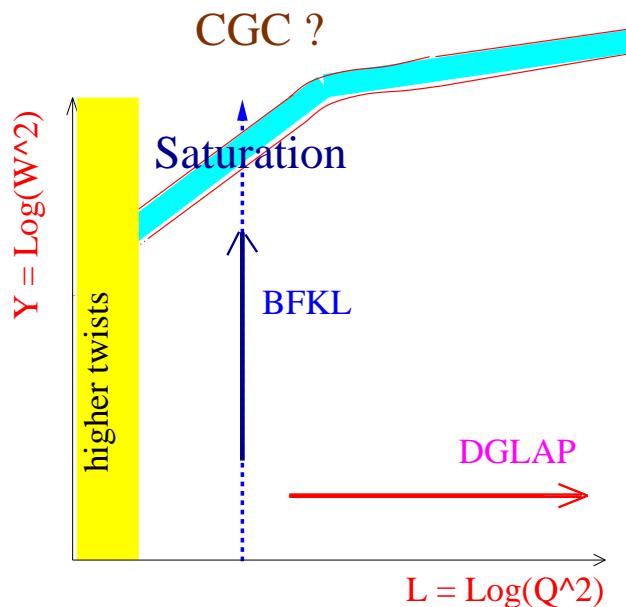
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Cracow School of Theoretical Physics XLVI, 2006

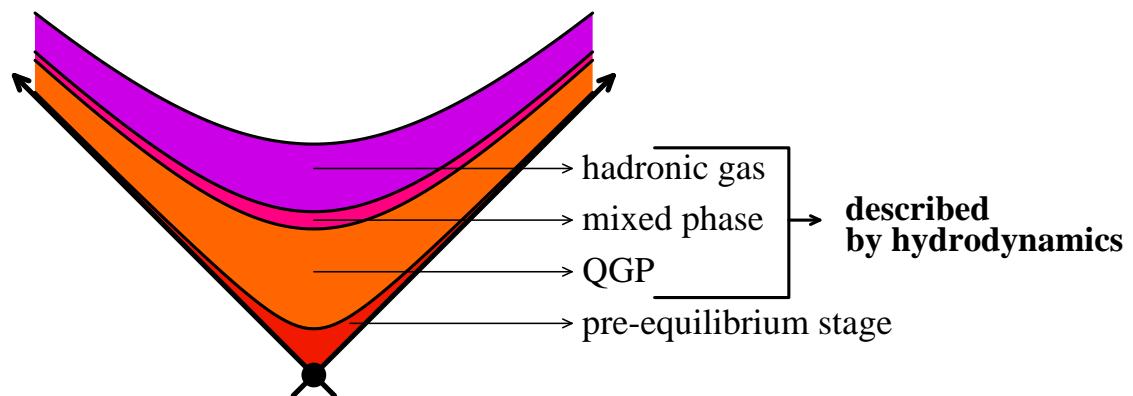
- **Introduction: QCD “*in situ*”**
Phase transitions in collision
- *Saturation and Non-Linear QCD Evolution:*
The Gluon Phase-space at saturation
- *Mapping to random polymers:*
A Spin-Glass Phase transition
- *Gluons near saturation:*
Clustering Phase and “Hot Spots”

QCD Dynamical Phase transitions

- The transition to Saturation at HERA

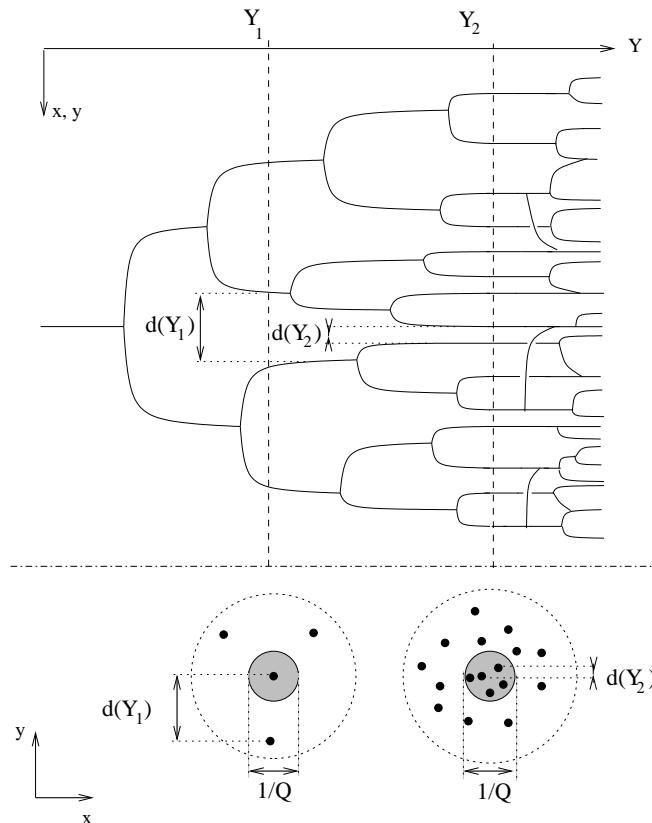


- The transition to QGP at RHIC



The QCD transition to Saturation

The “Tree” of Gluons / Dipoles



$d(Y) \rightarrow 0 \Rightarrow$ what happens?

$Y \sim Y_1$: Exponential growth: BFKL

$Y \sim Y_2$: Transition to Saturation

BFKL \rightarrow BK, JIMWLK, Fluctuations

$Y > Y_2$: Beyond: the CGC?

The QCD rapidity evolution

- The BFKL Evolution Operator

Diffusive Approx. (1d momenta $l \equiv \log k^2$):

$$\chi(-\partial_l) \equiv 2\psi(1) - \psi(-\partial_l) - \psi(1 + \partial_l) \sim A_0 + A_1 \partial_l + A_2 \partial_l^2 + \mathcal{O}(\partial_l^3)$$

A_0 : Exponential term

A_1 : “Drift” term

A_2 : Diffusion term

- “Mean-Field” Amplitude : $\mathcal{T} \sim cst. \mathbb{I}$

$$\partial_Y \mathcal{T} = \bar{\alpha} \chi(-\partial_l) \mathcal{T} - \bar{\alpha} \mathcal{T}^2$$

$$\mathcal{T}(l, Y) \sim \log \left[\frac{k^2}{Q_s^2(Y)} \right] \left(\frac{k^2}{Q_s^2(Y)} \right)^{-\sqrt{A_0/A_2}} \exp \left\{ -\frac{1}{A_2} \frac{\log^2 \left[\frac{k^2}{Q_s^2(Y)} \right]}{Y} \right\}$$

- Saturation Scale:

$$\log Q_s^2(Y) \propto [2\sqrt{A_0 A_2} - A_1] Y - \frac{3}{2\sqrt{A_0/A_2}} \log Y + \mathcal{O}(1)$$

In red: geometric scaling

The Balitskiĭ-Kovchegov Equation

- The Non-Linear BK Equation for \mathcal{T} :

$$\partial_Y \mathcal{T} = \bar{\alpha} \chi (-\partial_L) \mathcal{T} - \bar{\alpha} \mathcal{T}^2$$

- Equation BK \Rightarrow F-KPP

S.Munier, R.P., 2003,2004

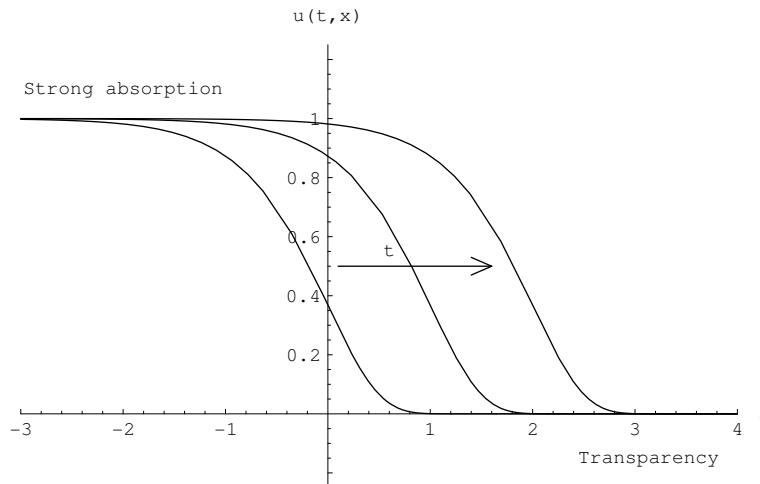
$$\partial_t u(t, x) = \partial_x^2 u(t, x) + u(t, x)(1 - u(t, x))$$

- “Dictionary”

$$\begin{aligned} Time &= t \propto Y \\ Space &= x \propto L + \frac{\bar{\alpha}D}{2} Y \\ Traveling\ Wave &= u(t, x) \sim u(t - vx) \propto \mathcal{T}(Y, k) \end{aligned}$$

Traveling wave Solutions

Bramson (1983)



- Traveling wave \rightarrow Geometric Scaling

$$u(t, x) \xrightarrow[t \rightarrow \infty]{} u(x - v(\gamma_c)t) \Rightarrow \boxed{\mathcal{T}(Y, k) = \mathcal{T}\left(\frac{k}{Q_s(Y)}\right)}$$

- Saturation Scale: \rightarrow “Universal terms”

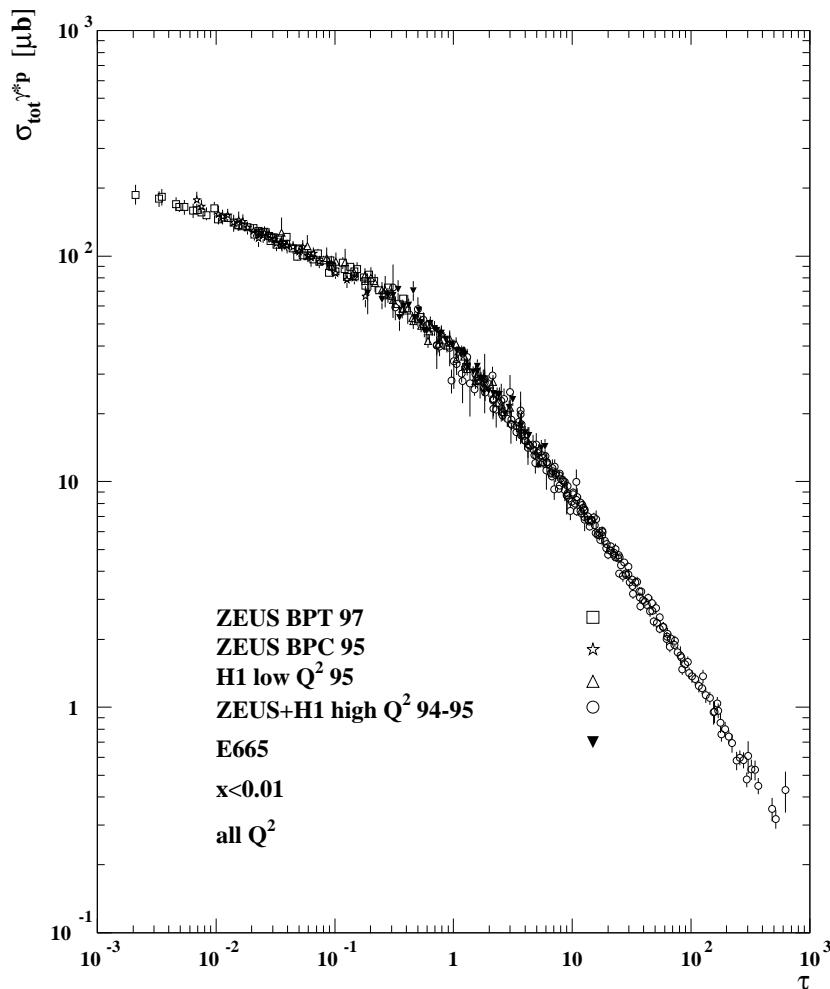
$$\log Q_s^2(Y) = \bar{\alpha} \frac{\chi(\gamma_c)}{\gamma_c} Y - \frac{3}{2\gamma_c} \log Y - \frac{3}{(\gamma_c)^2} \sqrt{\frac{2\pi}{\bar{\alpha} \chi''(\gamma_c)}} \frac{1}{\sqrt{Y}} + \mathcal{O}(1/Y)$$

Geometric Scaling

K.Golec-Biernat, J.Kwiecinski, A.Stasto (2000)

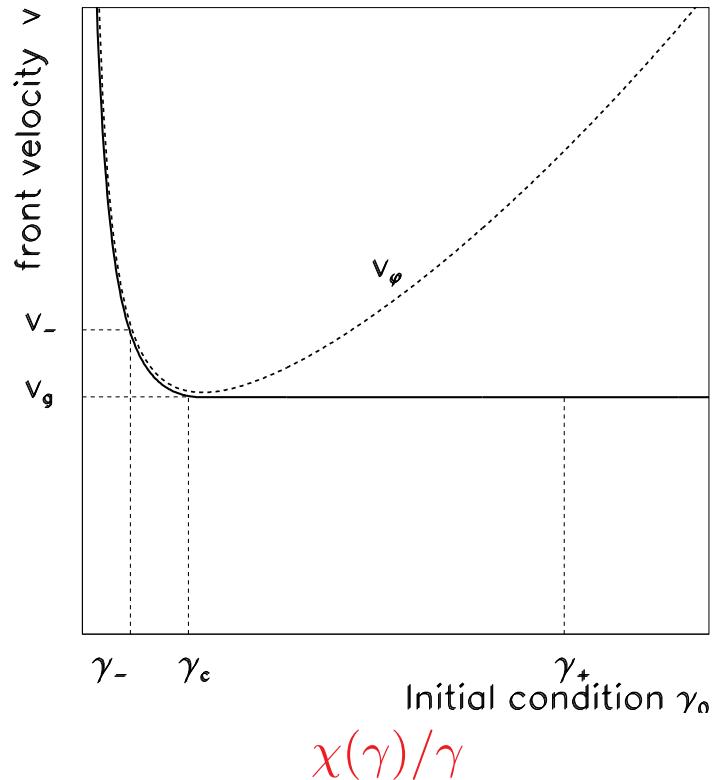
The Dipole Tree Observed in DIS:

$$\sigma^{\gamma^*}(Y, Q) = \int_0^\infty x_{01}^3 dx_{01} |\psi(x_{01}Q)|^2 \int kdk J_0(kx_{01}) \mathcal{T}(Y, k)$$



More on the (Super)Critical Regime

Nonlinear Dynamics



- Sub-critical regime: phase velocity

$$\gamma_0 = \gamma_- \Rightarrow v \equiv v_\varphi(\gamma) = \bar{\alpha} \chi(\gamma)/\gamma$$

- Critical regime: phase \equiv group velocity

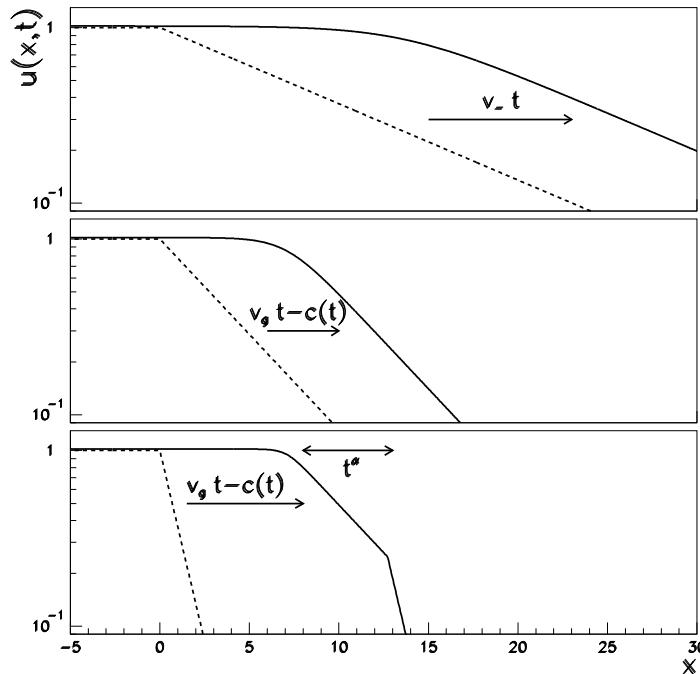
$$\gamma_0 = \gamma_c = .6275\dots \Rightarrow v_\varphi(\gamma_c) \equiv v_g(\gamma_c) = \bar{\alpha} \chi'(\gamma_c)$$

- Super-critical regime (cf. QCD at $\gamma_+ = 1$)

$$\gamma_0 = \gamma_+ \Rightarrow \bar{v} \equiv v_g(\gamma_c) = \min[v_\varphi(\gamma)]$$

The (Super)Critical Wave Front

Derrida, Van Saarlos: “Pulled vs. Pushed fronts”

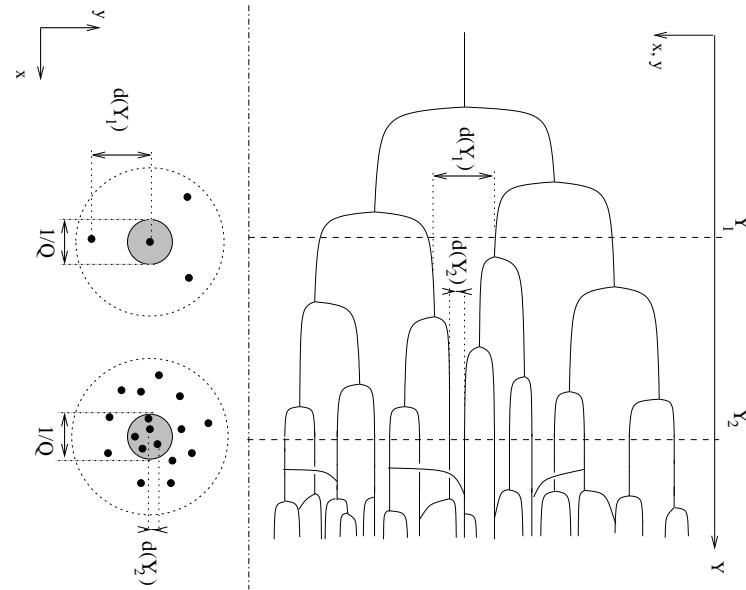


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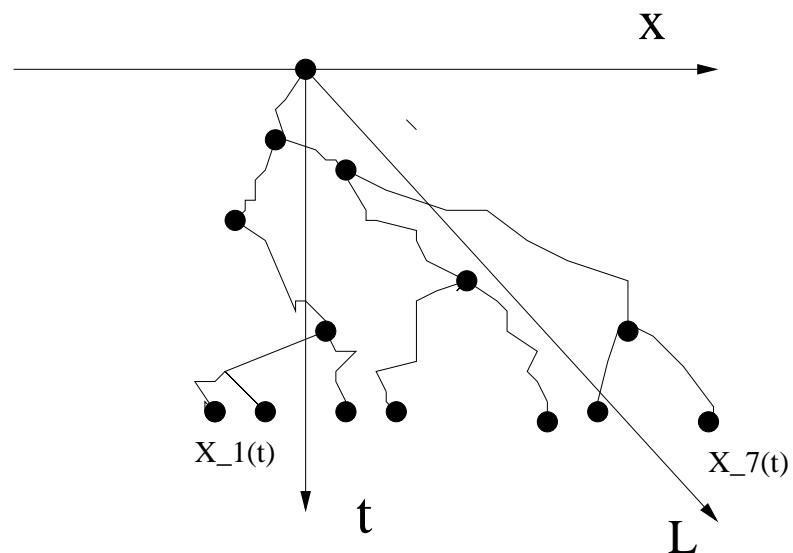
(Super)critical “Pulled” Fronts

Question: Is it the sign of a QCD dynamical phase transition?

Mapping to random polymers



$$y = \frac{t}{A_0} ; \quad L_i \equiv \log k_i^2(y) \Leftrightarrow -\beta (x_i(t) - x(0)) + (A_0 - A_1)y$$



$x_i(t) : \text{Branching} + \text{Diffusion} + \text{Shift}$

Thermodynamics of random polymers

- The Partition Function:

$$Z(t)_{\text{for one event}} \equiv \sum_{i=1}^n e^{-\beta x_i(t)} \Leftrightarrow \frac{1}{n} \sum_{i=1}^n k_i^2(y) \equiv \bar{k}^2(y)$$

- Generating functionale

$$G(x, t) \equiv \sum_p \frac{1}{p!} \langle \left[-e^{\beta x} Z(t) \right]^p \rangle_{\text{all events}} \Leftrightarrow 1 - u(x, t)$$

- “Dictionary”

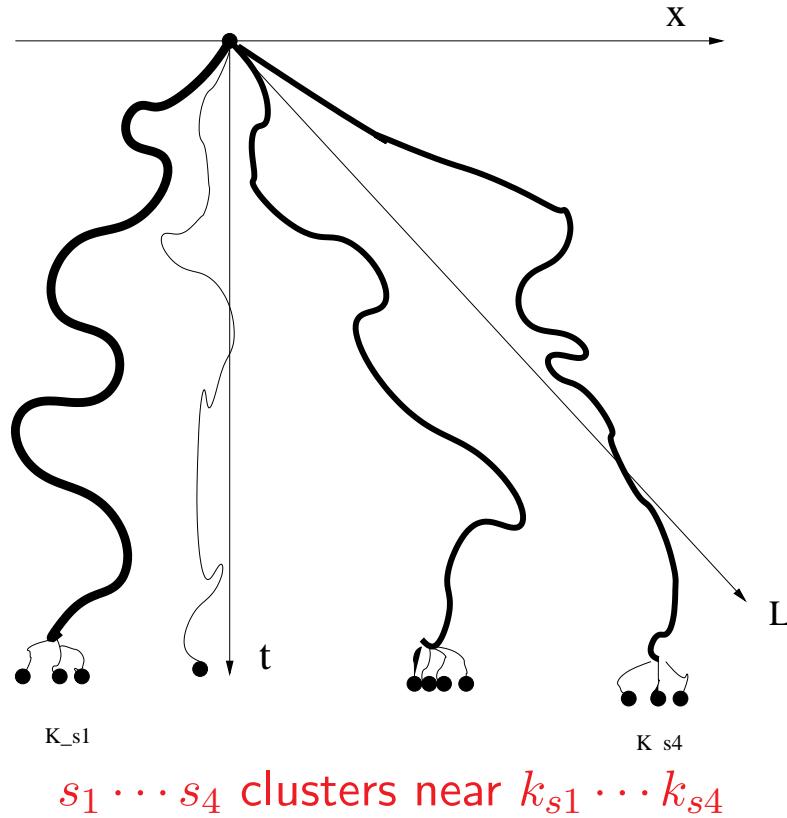
$$\text{Free energy : } F = -\frac{1}{\beta} \langle \log Z(t) \rangle \Leftrightarrow v \times t \sim \log(Q_s^2)$$

$$\text{Spectrum : } f = -\frac{1}{\beta} \log Z(t) + \frac{1}{\beta} \langle \log Z(t) \rangle \Leftrightarrow \mathcal{P}(f) \sim \mathcal{T}(Q^2/Q_s^2)$$

$$\text{Spin-Glass Phase : } \Leftrightarrow \beta > \beta_c : \gamma_0 > \gamma_c$$

The QCD spin-glass phase

Clustering structure



- Equivalent temperature $T < T_c$

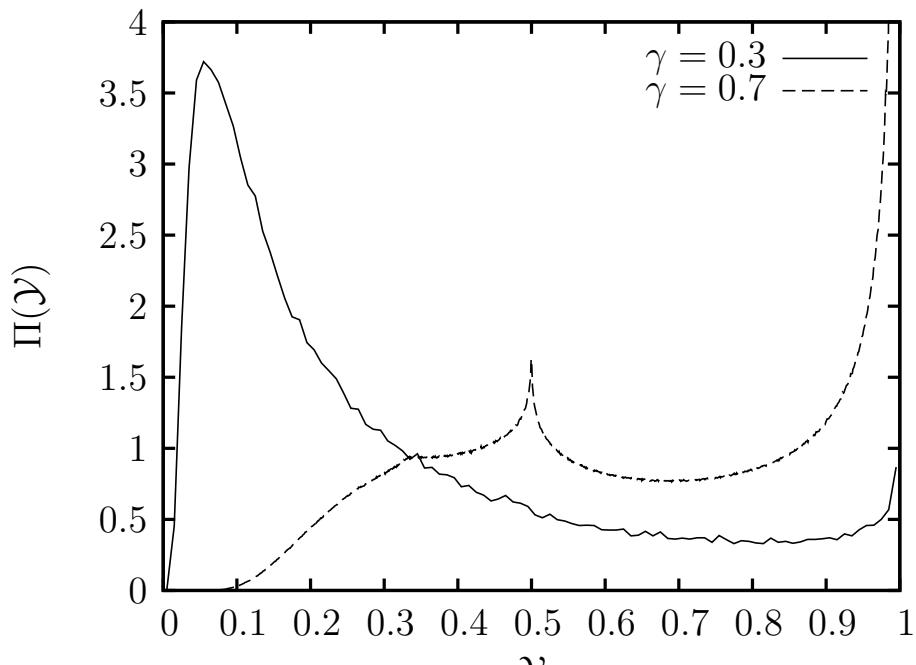
$$\frac{T_c - T}{T_c} \equiv 1 - \frac{\beta_c}{\beta} = 1 - \gamma_c$$

- Cluster weights

$$W_{si} = \frac{\sum_{i \in si} k_i^2}{\sum_i k_i^2}$$

The QCD spin-glass phase

Overlap Function



Overlap distribution $\Pi(\gamma)$

- Clustering Probability $\tilde{\mathcal{Y}}(q)$

$$\tilde{\mathcal{Y}}(q) dq = \frac{1}{\left\{ \sum (k_i^2) \right\}^2} \sum_{i,j=1}^n k_i^2 k_j^2 \Theta\{k_i(y=qY) = k_j(y=qY) \text{ in } [q, q+dq]\}$$

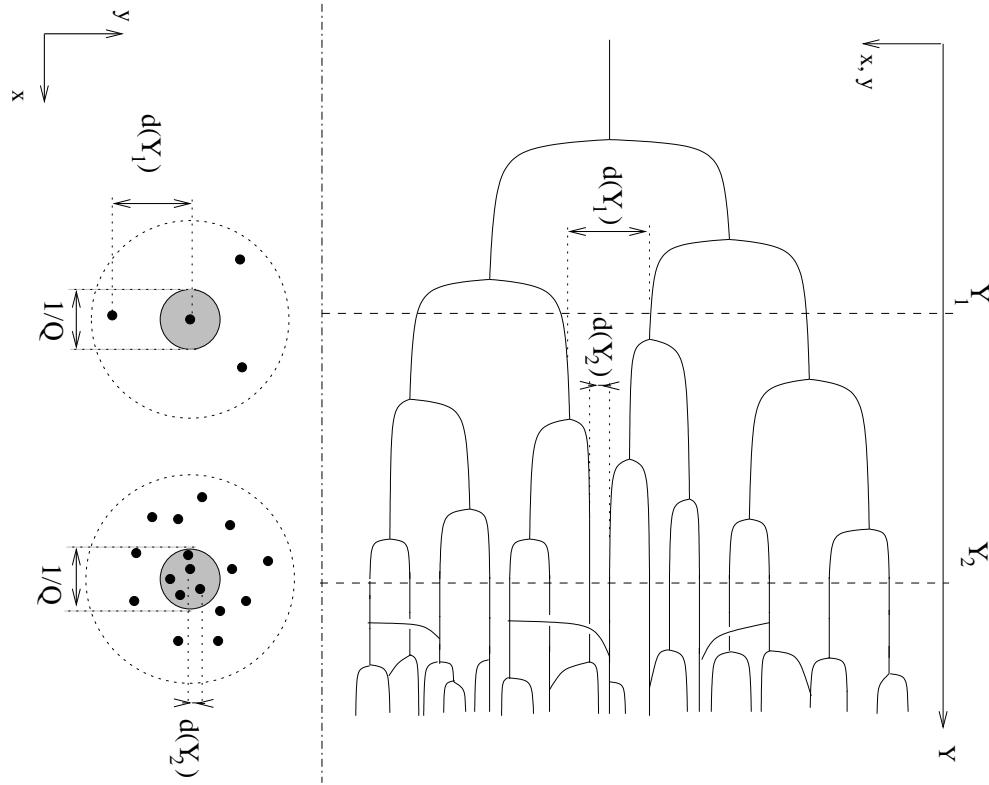
- Overlaps $\mathcal{Y} = \sum_{si} W_{si}^2$

$$\tilde{\mathcal{Y}}(q) dq = \delta(q-1) \mathcal{Y} + \delta(q) (1-\mathcal{Y})$$

- Cluster distribution $\Pi(\mathcal{Y})$

$$\langle \mathcal{Y} \rangle_\Pi = 1 - T/T_c ; \langle \mathcal{Y}^p \rangle_\Pi \Rightarrow \Pi(\mathcal{Y}) \text{ predicted}$$

Questions and Prospects



- Phenomenology
Connect with Experiments: “Hot Spots”
- BFKL kernel, Fluctuations, Pomeron Loops
~ Random polymers with constraints?
- The QCD phase beyond Saturation
Relation with Advanced Statistical Physics

Conclusions (part I)

- The QCD transition to Saturation
(Super)critical Dynamics
- Mapping to Random Polymers:
QCD Spin-Glass phase
- Results:
Clustering Structure: Overlaps
- Prospects
New Relation with Mathematics (non-linear Eqs.) and
Physics (Disordered systems, Polymer diffusion and Spin
glass phase transitions)