

Collapse of Primordial Clouds in the Presence of UV Radiation Field

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Introduction

- First stars formation and the cooling mechanism of primordial clouds
- The UV background from the first stars and quasars
- H_2 formation and destruction processes in the primordial clouds
- Spherically symmetric collapse of primordial cloud under UV background radiation

Hubble Ultra Deep Field Details
Hubble Space Telescope • Advanced Camera for Surveys

Formation of the first objects in the Universe

The first objects are a direct consequence of the growth of primordial density fluctuations.

- primordial density fluctuations
- decoupling of the perturbation from the Hubble flow
- subsequent collapse
- formation of a virialized cloud (cloud in the hydrostatic equilibrium)

Now we need some cooling mechanism.

- the successive fragmentation and contraction processes

Primordial star formation can be viewed as the successive fragmentation and contraction processes of collapsed cosmological objects.

Cooling mechanism in the metal-free gas clouds

- Large primordial clouds ($T > 10^4$ K, $M > 10^8 M_{\text{SUN}}$)
Emission of the radiation from the excited atomic hydrogen H.
- Small primordial clouds ($T < 10^4$ K, $M < 10^8 M_{\text{SUN}}$)
Emission of the radiation from the excited rotational and vibrational states of H_2 molecules. This mechanism is more efficient in the temperature of the order 100 K.

The presence of initial mass fraction of the molecular hydrogen H_2 of only 10^{-6} is enough for triggering the final collapse of low mass clouds.

Formation of primordial objects has been investigated mainly in the following two contexts:

- first luminous objects (first stars, quasars, globular clusters and proto-galaxies)
 - little influences of the external radiation field except for that of the CMB (the cosmic microwave background)
- ‘second generation objects’
 - largely affected by the external UV radiation produced by „the first luminous objects”

Sources of the UV radiation

- First stars
 - From the detailed 3D simulations of the formation process of the first stars we can conclude that these stars were very massive ($M \sim 100 M_{\text{SUN}}$).
- quasars
- The decay of very massive exotic particles

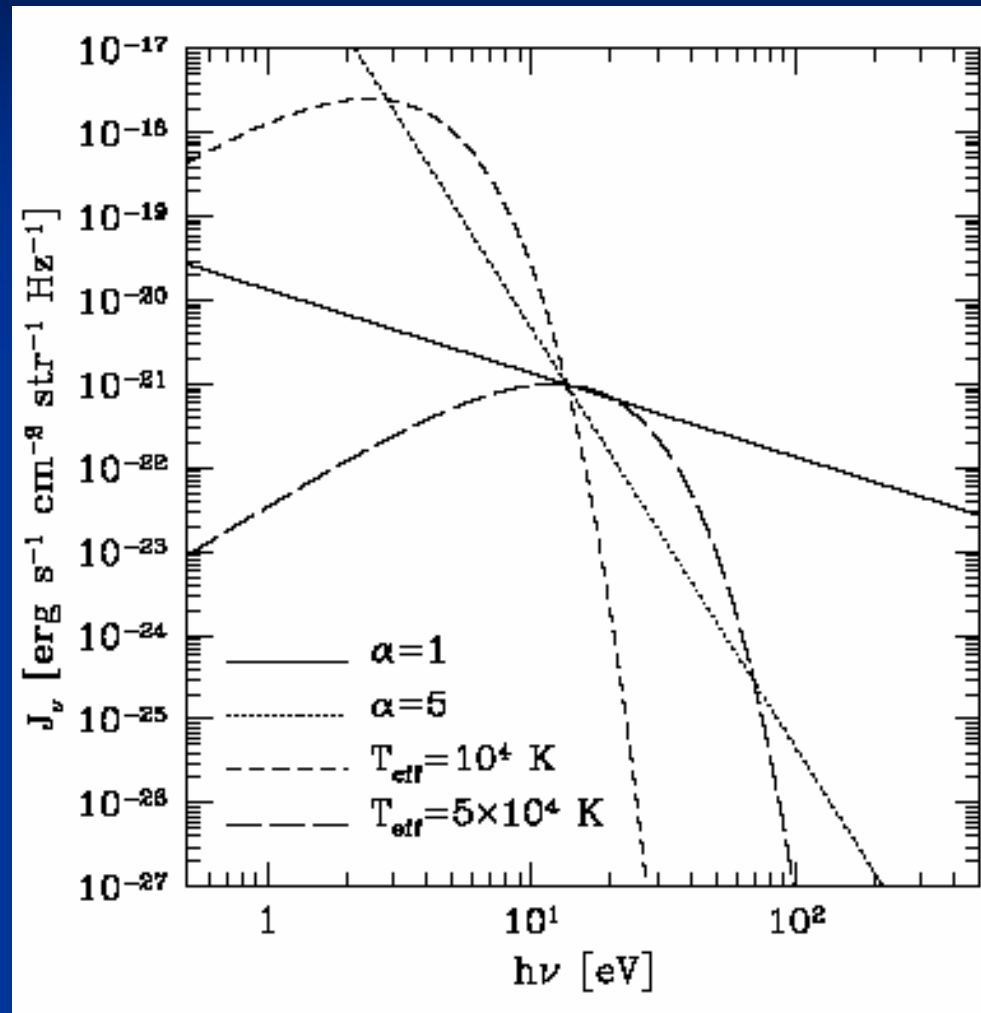
Spectrum of the ultraviolet background

- quasars

$$I_\nu \sim \nu^{-\alpha}$$

- First generation of stars
(black body spectrum)

$$I_\nu \sim \frac{\nu^3}{\exp(h\nu/k_B T_{\text{eff}}) - 1}$$



Feedback of the UV background on the star formation

- A key element is molecular hydrogen H_2 . In order for stars to form the gas needs to radiate energy efficiently and cool down to temperatures well below $T = 10^4$ K.
- In gas of primordial composition, without metals, cooling processes at this temperature are almost solely dominated by rotational-vibrational line excitation of H_2 .
- Formation and destruction of H_2 , however, is very sensitive to the presence of a radiation field.

H_2 formation and destruction processes in the primordial clouds

- Production of H_2



- Photo-dissociation of H_2 ($11,2 < h\nu < 13,6$ eV) (Lyman-Werner bands)



- Destruction of H^- and H_2^+ ($h\nu < 11,2$ eV)



Feedback of the UV background on the star formation

- Negative feedback on the star formation (first stars)
 - Dissociation of H_2 molecule by photons with energies from the Lyman-Werner bands
 - Destruction of H^- and H_2^+

- Positive feedback on the star formation (quasars, the decay of exotic particles)
 - enhancement of the ionized fraction.

Feedback of the UV background on the star formation

- In addition, self-shielding which takes place during dynamic collapse can also aid H_2 cooling. It means that **collapse and star formation in the high-density (more massive) clouds can take place efficiently even under high intensities of the UV background.**
- **We can see that it is by no means trivial in what circumstances the UVB has positive or negative feedback on star formation in the 'second generation' objects.**
- **We need the detailed simulation of the collapse of the primordial cloud**

'Difficulties' in taking into account the radiation feedback

- Non Local Thermal Equilibrium
the Boltzmann distribution function is not valid

$$\frac{N_u}{N_l} \neq \frac{g_u}{g_l} \exp\left(-\frac{E_{ul}}{kT}\right)$$

- Presence of thousands of H₂ lines
- Solving the radiative transfer equation in order to get correct cooling function

Evolution of spherically symmetric density contrast in the Lambda CDM Universe (Stachniewicz & Kutschera 2001; 2003)

- Equations of dynamic evolution of the dark matter and baryonic gas.

Cloud is divided into the concentric mass-shells. Simulation is tracing collapse of the individual mass-shells.

$$\frac{dm_B}{dr_B} = 4\pi r_B^2 \rho_B,$$

$$\frac{d^2 r_B}{dt^2} = -4\pi r_B^2 \frac{dP}{dm_B} - \frac{GM(< r_B)}{r_B^2},$$

$$\frac{d^2 r_{dm}}{dt^2} = -\frac{GM(< r_{dm})}{r_{dm}^2}.$$

Evolution of spherically symmetric density contrast in the Lambda CDM Universe (Stachniewicz & Kutschera 2001; 2003)

■ Energy equation

$$\frac{du}{dt} = \frac{P}{\rho_B^2} \frac{d\rho_B}{dt} - \frac{\Lambda_{cool}}{\rho_B}$$

■ Equation of state for ideal gas

$$P = (\gamma - 1) \rho_B u = \frac{k_B \rho_B T}{\mu m_p}, \quad \gamma = \frac{5}{3}$$

$(v', J') \quad 2p\sigma B^1\Sigma_u^+ \longleftrightarrow (v'', J'') \quad 1s\sigma X^1\Sigma_g^+ \quad \text{Lyman band}$

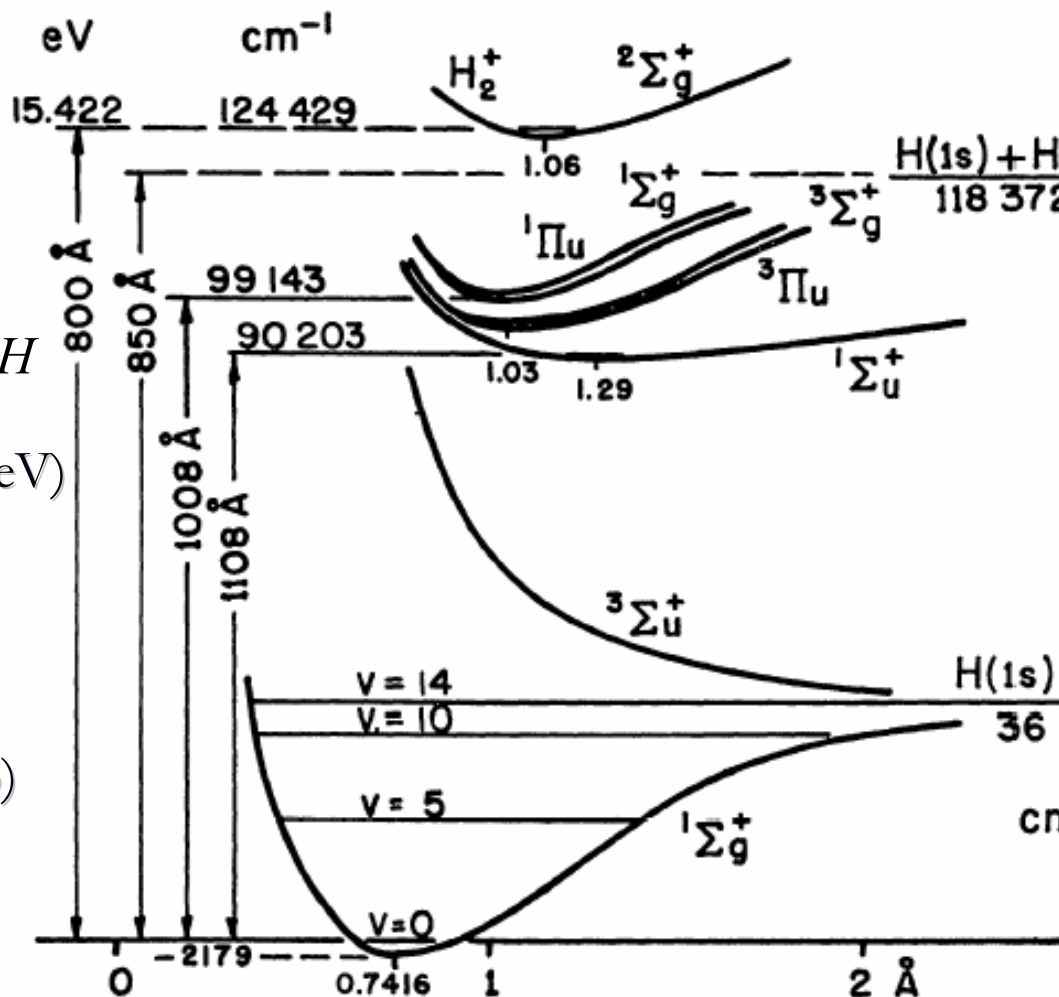
$(v', J') \quad 2p\pi C^1\Pi_u \longleftrightarrow (v'', J'') \quad 1s\sigma X^1\Sigma_g^+ \quad \text{Werner band}$

$\frac{H(1s)+p+e}{145\,792} \quad 18.071$

$\frac{H(1s)+H(2s,2p)}{118\,372} \quad 14.671$

$\frac{H(1s)+H(1s)}{36\,113} \quad 4.476$

$\text{cm}^{-1} \quad \text{eV}$



$H_2 + h\nu \rightarrow H_2^* \rightarrow 2H$
($11,18 < h\nu < 13,6 \text{ eV}$)

(Field, et. al. 1966)

Transfer equation for spherically symmetric case

$$\frac{1}{c} \frac{\partial I_\nu}{\partial t} + \mu \frac{\partial I_\nu}{\partial r} + \frac{1 - \mu^2}{r} \frac{\partial I_\nu}{\partial \mu} = \rho \{j_\nu - \kappa_\nu I_\nu\},$$

where

- $I_\nu(r, \mu, t)$ $\left[\frac{\text{erg}}{\text{cm}^2 \text{ s Hz Sr}} \right]$ - radiation intensity,
- ρ - density,
- κ_ν, j_ν - opacity and emissivity for frequency ν ,
- $\mu = \cos \theta \in [-1, 1]$,
- θ - angle between radiation direction and radial direction,
- $r \in [r_{min}, r_{max}]$ - radius.

If we define „source function”

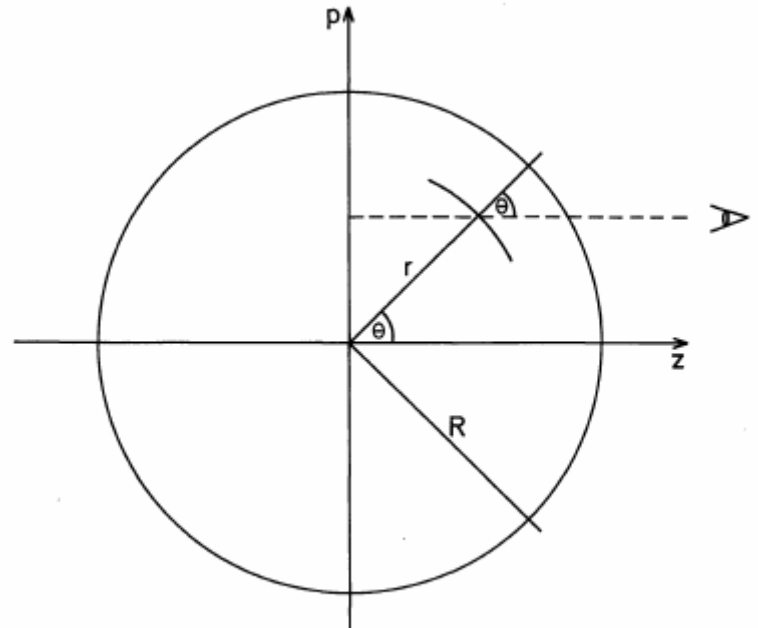
$$S_\nu = \frac{j_\nu}{\kappa_\nu},$$

than transfer equation will be

$$\frac{1}{c} \frac{\partial I_\nu}{\partial t} + \mu \frac{\partial I_\nu}{\partial r} + \frac{1 - \mu^2}{r} \frac{\partial I_\nu}{\partial \mu} = -\rho \kappa_\nu \{I_\nu - S_\nu\}.$$

$$I_\nu(r, \mu, t) \rightarrow \begin{cases} I_\nu^+(r, p, t) & \mu \geq 0 \\ I_\nu^-(r, p, t) & \mu < 0 \end{cases}$$

$$\begin{cases} \mu \frac{\partial I_\nu^+}{\partial r} = -\rho \kappa_\nu \{I_\nu^+ - S_\nu\} \\ \mu \frac{\partial I_\nu^-}{\partial r} = \rho \kappa_\nu \{I_\nu^- - S_\nu\} \end{cases}$$



$$(r, \mu) \rightarrow \left(r, p = r \sin \theta = r \sqrt{1 - \mu^2} \right)$$

We can define zero, first and second moment of radiation intensity as follows:

$$J_\nu(r) = \frac{1}{2} \int_0^1 (I_\nu^+ + I_\nu^-) d\mu$$

$$H_\nu(r) = \frac{1}{2} \int_0^1 (I_\nu^+ - I_\nu^-) \mu d\mu$$

$$K_\nu(r) = \frac{1}{2} \int_0^1 (I_\nu^+ + I_\nu^-) \mu^2 d\mu$$

$$L(r) = 16\pi^2 r^2 H(r)$$

Eddington factor:

$$f_\nu(r) = \frac{K_\nu(r)}{J_\nu(r)}$$

There are also the zero and first moment equation

$$\frac{\partial (f_\nu(r) J_\nu(r))}{\partial r} + \frac{3f_\nu(r) - 1}{r} J_\nu(r) + \rho \kappa_\nu H_\nu(r) = 0$$

$$\frac{\partial H_\nu(r)}{\partial r} + \frac{2H_\nu(r)}{r} + \rho \kappa_\nu^a J_\nu(r) - \rho j_\nu^t = 0$$

If we define

$$q_\nu(r) = \exp \left[\int_{r_c}^r \left(3 - \frac{1}{f_\nu(r')} \right) \frac{dr'}{r'} \right],$$

than we get

$$\frac{\partial (f_\nu(r) q_\nu(r) J_\nu(r))}{\partial r} = -\rho \kappa_\nu q_\nu(r) H_\nu(r)$$

$$\frac{\partial (H_\nu(r) r^2)}{\partial r} = r^2 \rho (j_\nu^t - \kappa_\nu^a J_\nu(r)).$$

Two systems of equations

- System I (solved for $f_\nu(r)$ and $q_\nu(r)$)

$$\begin{cases} \mu \frac{\partial I_\nu^+}{\partial r} = -\rho \kappa_\nu \{I_\nu^+ - S_\nu\} \\ \mu \frac{\partial I_\nu^-}{\partial r} = \rho \kappa_\nu \{I_\nu^- - S_\nu\} \end{cases}$$

- System II (solved for $J_\nu(r)$ and $H_\nu(r)$)

$$\begin{aligned} \frac{\partial (f_\nu(r) q_\nu(r) J_\nu(r))}{\partial r} &= -\rho \kappa_\nu q_\nu(r) H_\nu(r) \\ \frac{\partial (H_\nu(r) r^2)}{\partial r} &= r^2 \rho (j_\nu^t - \kappa_\nu^b J(r)) \end{aligned}$$

The numerical computational procedure to trace the dynamic evolution of the primordial cloud under the UV background

- We solve the hydrodynamic equations of motion along with equations for energy conservation, ionization, and dissociation of molecular and atomic species.
- We solve the system I equations with the initial value of the source function for f_ν and for q_ν .
- We solve the system II equations for J_ν and H_ν .
- Update the source function and solve the system I once more.
- Iterative procedure between system I and II is continued until convergence.
- After the convergence of the I and II system we calculate luminosity L_ν from the first moment H_ν and than cooling function from luminosity.
- We update abundance of different species and number densities of each atomic and molecular state.

Conclusions

- There is positive and negative feedback of UV radiation to the collapse of primordial clouds and star formation.
- It is by no means trivial in what circumstances the UV radiation has positive or negative feedback on star formation in the 'second generation' objects.
- We need the detailed simulation of the collapse of the primordial cloud
- Collapse and star formation in the high-density clouds can take place efficiently even under high intensities of the UV background.

Bibliography

- Bouwens, R., J., et. al., 2004, ApJ 616L, 79B (HUDF)
- Bromm, V., et. Al., 1999, ApJL 527, 5. (3D simulation of primordial cloud collapse)
- Bunker, A., J., et. al., 2004, MNRAS 355, 374 (HUDF)
- Field, G., B., et. al., 1966, ARA&A 4, 207F (properties of H₂ molecule)
- Pelló, R., et. al., 2004, A&A, 416, L35 (the farthest known galaxy)
- *Stachniewicz, S. & Kutschera, M. 2003, MNRAS 339,616*
- *Stachniewicz, S. & Kutschera, M. 2001, Acta Phys. Pol. B 32, 227*
- Stancil, P. C., 1994, ApJ, 430, 360 (H₂ reaction)
- Tegmark, M., et. al., 1997, ApJ, 474, 1 (H₂ reaction)
- Yan, H., Windhorst, R., A., 2004, ApJ 612, L93 (HUDF)