

Quantum Kaluza-Klein Theory with $M_2(\mathbb{C})$ and particle physics

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Noncommutative geometry: metric and spectral aspects

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Elements of Quantum Riemannian Geometry(QRG)

Quantum Riemannian Geometry on $M_2(\mathcal{C})$

Kaluza-Klein model and scalar fields on $C^\infty(M) \otimes M_2(\mathcal{C})$

Elements of Quantum Riemannian Geometry(QRG)

We work with an algebra A , typically a $*$ -algebra over \mathbb{C} .

Differential structure $d: A \rightarrow \Omega^1$ $d(ab) = adb + (da)b$, where $\Omega^1 = \text{span}\{a db \mid a, b \in A\}$

$d: \Omega^n \rightarrow \Omega^{n+1}$ $d(\omega \wedge \eta) = (d\omega) \wedge \eta + (-1)^n \omega \wedge d\eta$, $d^2 = 0$

where $\omega, \eta \in \Omega = \bigoplus_{n \geq 0} \Omega^n$ with $\Omega^0 = A$

Connection $\nabla: \Omega^1 \rightarrow \Omega^1 \otimes_A \Omega^1$, $\nabla(a.\omega) = a.\nabla\omega + da \otimes_A \omega$, where $\omega \in \Omega^1$

σ map: $\Omega^1 \otimes_A \Omega^1 \rightarrow \Omega^1 \otimes_A \Omega^1$, $\nabla(\omega.a) = (\nabla\omega).a + \sigma(\omega \otimes_A da)$

Connection $\nabla: \Omega^1 \otimes_A \Omega^1 \rightarrow \Omega^1 \otimes_A \Omega^1 \otimes_A \Omega^1$ $\nabla(\omega \otimes_A \eta) = \nabla\omega \otimes_A \eta + (\sigma \otimes_A \text{id})(\omega \otimes_A \nabla\eta)$

Elements of Quantum Riemannian Geometry(QRG)

Metric $g \in \Omega^1 \otimes_A \Omega^1$, which is invertible in the sense of a bimodule map $(\ , \) : \Omega^1 \otimes_A \Omega^1 \rightarrow A$ obeying

$$((\omega, \) \otimes_A \text{id})g = \omega = (\text{id} \otimes_A (\ , \omega))g$$

quantum symmetry condition:

$$\wedge(g)=0$$

Quantum Levi-Civita connection(QLC):

$$\nabla g = 0 \text{ and } \wedge \nabla - d=0$$

Riemann curvature $R_\nabla : \Omega^1 \rightarrow \Omega^2 \otimes_A \Omega^1$

$$R_\nabla = (d \otimes_A \text{id} - \text{id} \wedge \nabla)\nabla$$

Ricci tensor :

$$\text{Ricci} = ((\ , \) \otimes_A \text{id})(\text{id} \otimes i \otimes \text{id})(\text{id} \otimes R_\nabla)g, \quad \text{where } i: \Omega^2 \rightarrow \Omega^1 \otimes_A \Omega^1$$

Ricci scalar S:

$$S = (\ , \)\text{Ricci} \in A$$

Reality condition:

$$g^\dagger = g \text{ and } \nabla \circ * = \sigma \circ \dagger \circ \nabla$$

Recall QRG on algebra $A = M_2(\mathbb{C})$

In terms of a self-adjoint basis s^i (where $s^{i*} = s^i$), the exterior algebra is

$$s^1 \wedge s^1 = s^2 \wedge s^2 = i \text{Vol}, \quad s^1 \wedge s^2 = s^2 \wedge s^1 = 0$$

Using Pauli matrices and $1 = \text{id}$ as a basis of $M_2(\mathbb{C})$, the differentials are then

$$d\sigma^1 = \sigma^3 s^2, \quad d\sigma^2 = -\sigma^3 s^1, \quad d\sigma^3 = \sigma^2 s^1 - \sigma^1 s^2, \quad ds^i = -\sigma^i \text{Vol}, \quad \text{Vol} = -i s^1 \wedge s^1$$

The quantum metric of interest is

$$g = -s^1 \otimes s^1 + s^2 \otimes s^2$$

A natural 1-parameter QLC has braiding (where $\bar{1} = 2$ and $\bar{2} = 1$),

$$\sigma(s^i \otimes s^j) = \begin{cases} s^i \otimes s^j & \text{if } i \neq j \\ -s^{\bar{i}} \otimes s^{\bar{i}} - 2i\rho s^{\bar{i}} \otimes s^i & \text{if } i = j \end{cases}$$

Recall QRG on algebra $A = M_2(\mathbb{C})$

The connection is then

$$\nabla s^i = \frac{\iota}{2} \sigma^i (s^1 \otimes s^1 + s^2 \otimes s^2) + \frac{\iota}{2} \epsilon_{ij} \sigma^j (s^2 \otimes s^1 - s^1 \otimes s^2) - \rho \sigma^i s^{\bar{i}} \otimes s^i.$$

Choosing the canonical symmetric lift

$$i(\text{Vol}) = \frac{1}{2\iota} (s^1 \otimes s^1 + s^2 \otimes s^2)$$

Riemann curvature and Ricci tensor are

$$R_{\nabla} s^i = -\iota(1 + \rho^2) \text{Vol} \otimes s^i, \quad \text{Ricci} = -\frac{1}{2}(1 + \rho^2)(s^1 \otimes s^1 + s^2 \otimes s^2), \quad S = 0.$$

QRG on algebra $A = C^\infty(M) \otimes M_2(\mathbb{C})$

Follow a quantum Kaluza-Klein formulation, we solve for the quantum Riemannian geometry on $A = C^\infty(M) \otimes M_2(\mathbb{C})$,

We take the graded tensor product exterior algebra $\Omega(A) = \Omega(M) \underline{\otimes} \Omega(M_2(\mathbb{C}))$

Proposition: For the reason that quantum metric g has to be central, the most general metric has the form

$$g = g_{\mu\nu}(x, t) dx^\mu \otimes dx^\nu + A_{i\mu}(x, t)(s^i \otimes dx^\mu + dx^\mu \otimes s^i) + h_{ij}(x, t)s^i \otimes s^j$$

where $g_{\mu\nu}$ are symmetric, $h_{11} + h_{22} = 0$ and $A_{i\mu} = A_{\mu i}$, and all coefficients are functions on spacetime M .

QRG on algebra $A = C^\infty(M) \otimes M_2(\mathbb{C})$

From now on, we assume for simplicity that if σ exists, it obeys

- (i) $\sigma(dx^\mu \otimes s^i)$ is the flip;
- (ii) The result of $\sigma(s^i \otimes s^j)$ doesn't have any terms with dx^μ .

Proposition: If ∇ is a torsion free bimodule connection then it has the form

$$\nabla dx^\mu = -\Gamma_{\alpha\beta}^\mu dx^\alpha \otimes dx^\beta + B_{i\alpha}^\mu (dx^\alpha \otimes s^i + s^i \otimes dx^\alpha) + C_{ij}^\mu s^i \otimes s^j$$

$$\nabla s^k = D_{\alpha\beta}^k dx^\alpha \otimes dx^\beta + E_{i\alpha}^k (dx^\alpha \otimes s^i + s^i \otimes dx^\alpha) + \gamma_{ij}^k s^i \otimes s^j$$

where $\Gamma_{\alpha\beta}^\mu, D_{\alpha\beta}^k$ are symmetric for the subscripts, $\Gamma_{\alpha\beta}^\mu, B_{i\alpha}^\mu, C_{ij}^\mu, D_{\alpha\beta}^k, E_{i\alpha}^k \in C^\infty(M)$, $\gamma_{ij}^k \in C^\infty(M) \otimes M_2(\mathbb{C})$ but does not have an σ^3 component, and

$$C_{11}^\mu + C_{22}^\mu = 0, \quad \gamma_{11}^k + \gamma_{22}^k = \iota\sigma^k.$$

QRG on algebra $A = C^\infty(M) \otimes M_2(C)$

Proposition: The only metrics that admit a QLC are of the form

$$g = g_{\mu\nu}(x, t) dx^\mu \otimes dx^\nu + h_{ij}(x, t) s^i \otimes s^j$$

and the QLC has the form in last proposition with the further restrictions $D = B = 0$, Γ is the usual LC for $g_{\mu\nu}$, γ is the usual QLC for h_{ij} in the sense

$$h_{ji} \gamma_{mn}^j + 2i h_{pj} \gamma_{mnq}^p \gamma_{qi}^j + h_{nj} \gamma_{mi}^j = 0$$

and C, E obey

$$C_{ij}^\mu = -g^{\mu\alpha} h_{mj} E_{i\alpha}^m,$$

$$\partial_\alpha h_{ij} + h_{mj} E_{i\alpha}^m + h_{im} E_{j\alpha}^m = 0,$$

$$2i h_{pm} E_{n\alpha}^m \gamma_{ijn}^p + (h_{jm} - h_{mj}) E_{i\alpha}^m = 0.$$

QRG on algebra $A = C^\infty(M) \otimes M_2(\mathbb{C})$

The inverse of the metric $(,)$ has the form

$$(dx^\mu, dx^\nu) = g^{\mu\nu}, \quad (s^i, s^j) = h^{ij}, \quad (dx^\mu, s^i) = 0, \quad (s^i, dx^\mu) = 0$$

where $g^{\mu\nu}, h^{ij}$ are the inverse of $g_{\mu\nu}, h_{ij}$

The Laplacian on $f = f_a(x, t)\sigma^a$ is

$$\Delta f = (,)\nabla df = \Delta_{LB}f + \Delta_{M_2}f - g^{\mu\alpha}h_{mj}E_{i\alpha}^m(\partial_\mu f_a)\sigma^a h^{ij}$$

where Δ_{LB} is the usual GR Laplacian on each component f_a and

$$\Delta_{M_2}f = h^{ij}(f_i\sigma^j - \epsilon_{kcb}f_c\sigma^b\gamma_{ij}^k)$$

is the Laplacian on $M_2(\mathbb{C})$ at each x, t .

QRG on algebra $A = C^\infty(M) \otimes M_2(C)$

If $h_{ij}s^i \otimes s^j = h(x,t)(-s^1 \otimes s^1 + s^2 \otimes s^2)$, then

$$\Delta f(x,t) = (\Delta_{LB} f_a(x,t))\sigma^a + \frac{1}{h(x,t)} \left(-f_1(x,t)\sigma^1 + f_2(x,t)\sigma^2 - g^{\mu\nu}(\partial_\nu h(x,t)\partial_\mu f_a(x,t))\sigma^a \right).$$

Especially, if $h(x,t)$ is constant in spacetime and we let $h^{-1}(x,t) = \delta_m$ then the KG equation operator reduces to

$$(\Delta + m^2)f = (\Delta_{LB} f_a)\sigma^a + m^2 f_0 + (m^2 - \delta_m)f_1(x,t)\sigma^1 + (m^2 + \delta_m)f_2(x,t)\sigma^2 + m^2 f_3\sigma^3,$$

This implies the field f usually regarded as one component field in classical spacetime now can be regarded as quadruplet of fields having different masses.

QRG on algebra $A = C^\infty(M) \otimes M_2(C)$

Proposition:

If the $h_{ij}s^i \otimes s^j = h(x, t)(-s^1 \otimes s^1 + s^2 \otimes s^2)$ and we let $h(x, t) = e^{\frac{2}{\sqrt{3-3\rho^2}}\phi(x, t)}$ then

$$S = -\frac{1}{2}S_M + (\nabla^\mu \phi)(\nabla_\mu \phi) + \frac{2}{\sqrt{3-3\rho^2}}\nabla^\mu \nabla_\mu \phi$$

where $\phi(x, t)$ is real scalar field and ρ is an imaginary parameter.

From here we can see that this result has the scalar field which also appears in KK theory.

Thank you



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