

# Mechanika Kwantowa dla doktorantów

## zestaw 17 na dzień 16.03.2017 godz. 8:15

1. Consider radial Schrödinger equation for a potential  $V(r) \rightarrow 0$  for  $r \rightarrow \infty$  and for energy  $E = \hbar^2 k^2 / (2m) > 0$  assuming that the wave function factorizes into a radial and angular parts

$$\psi_{klm}(\vec{r}) = R_{kl}(r) Y_l^m(\vartheta, \varphi)$$

where  $Y_l^m$  denotes spherical harmonics. For large  $r$  where  $V(r) = 0$  (or equivalently for a free particle), this is so called modified Bessel equation. First find solutions for  $l = 0$  defining a new function  $u(r) = r R_{k0}(r)$ . Note that there are two possible solutions of this equation, one physical that is finite at the origin, and the other one that explodes for  $r = 0$ . They correspond to the spherical Bessel functions  $j_l(kr)$  and  $y_l(kr)$ , respectively.

Next, in order to find solutions for  $l \neq 0$  do the following:

- define  $\chi_{kl}(r) = R_{kl}(r)/r^l$  and derive equation for  $\chi_{kl}(r)$ ,
- differentiate the above equation over  $r$ ,
- define a new function  $f_{kl}(r) = r \chi'_{kl}(r)$ ,
- compare equation for  $f_{kl}$  with the initial equation for  $\chi_{kl}$  and read off the recurrence formula relating  $\chi_{kl+1}$  to  $\chi_{kl}$ ,
- solve the recurrence formula for  $R_{kl}(r)$  in the case  $l = 1, 2$ ,
- find asymptotic behaviour of  $R_{kl}(r)$  for large  $r$  (that is find asymptotics for  $j_l$  but also for  $y_l$ ),
- using the initial equation for  $R_{kl}(r)$  find asymptotic behaviour for small  $r$  (also both  $j_l$  and for  $y_l$ ).

2. Show that for an infinite "hard ball" potential:

$$V(r) = \begin{cases} 0 & \text{dla } a < r \\ \infty & \text{dla } r < a \end{cases}$$

the radial wave function  $R_{kl}^{\text{ball}}$  differs from the one of a free equation  $R_{kl}$  by a phase shift  $\delta_l(k)$

$$R_{kl}^{\text{ball}}(r) = R_{kl}(r + \delta_l(k)). \quad (1)$$

Calculate the phase shifts from the condition  $R_{kl}(a) = 0$ . Find low energy behaviour of  $\delta_l(k)$ .

HINT: The solution for  $l = 0$  is straightforward. However for  $l \neq 0$  more work is needed. Use (1) in the asymptotic region  $r \rightarrow \infty$  where one can use explicit asymptotic form of  $j_l(kr + \delta_l(k))$ . Use trigonometric identity to separate dependence on  $\delta_l(k)$  and on  $rk$ . Replace asymptotic forms by full Bessel functions  $j_l(kr)$  and  $y_l(kr)$  and apply condition  $R_{kl}(a) = 0$ .

3. Prove by substituting to the spherical Bessel equation that

$$j_l(z) = \frac{1}{2} \frac{1}{i^l} \int_{-1}^1 dt e^{izt} P_l(z).$$

To this end use the identity

$$P_l(z) = \frac{1}{2^l l!} \frac{d^l}{dz^l} (z^2 - 1)^l.$$