Elementary Particle Physics: theory and experiments

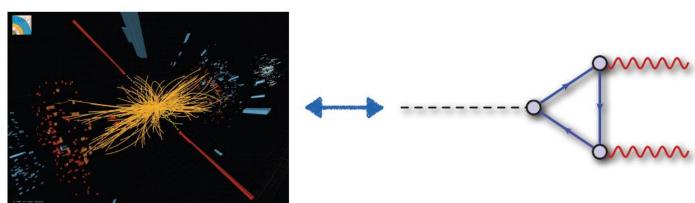
Introduction: highlights of last 100 years

Elementary particle physics

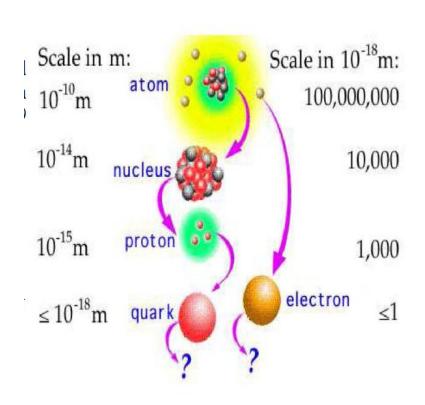
This course will introduce the Standard Model of particle physics and both theoretical foundations and experimental techniques. Main aim is to introduce phenomenology side of particle physics.

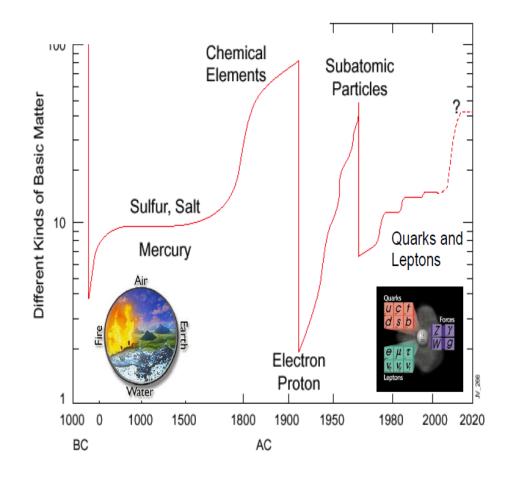
- Most of the presented experimental material will come from upto-date measurements at LHC experiments.
- The theoretical foundations will follow book by D. Griffiths, "Introduction to Elementary Particles".

When preparing those slides, I have borrowed some material from lectures or courses by: F. Krauss, C. Williams, P. Bechtle, R. Otario, Ch. Sanders, and others.



Constituents of matter along History



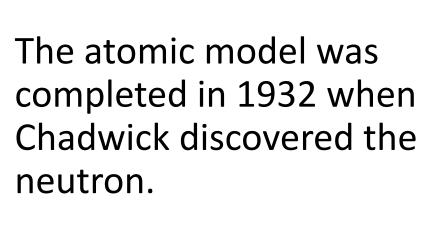


Early fundamental particles

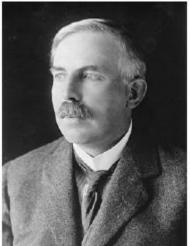
The earlies fundamental particle discovery was the electron, by J. J. Thomson in 1987



Followed by the discovery of proton by Rutherford in 1917



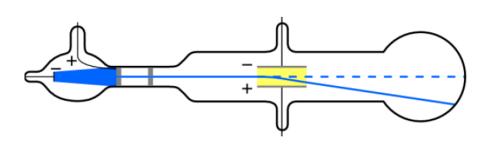




The Electron

Observation: 1897

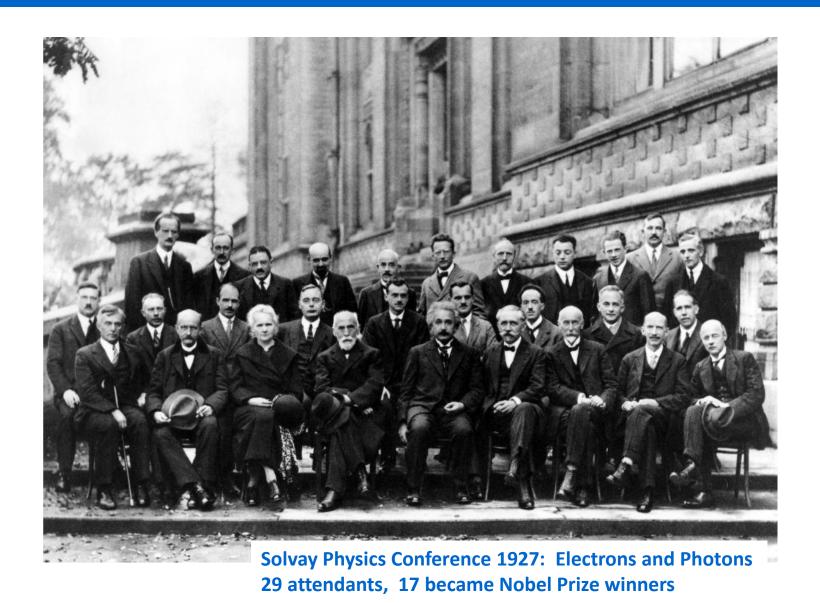
- Constituents of cathode rays deflected by electric field
- Constituents of cathode rays deflected by magnetic fields + heating of thermal junction -> first mass/charge ratio
- Higher precision of mass/charge from comparing deflection by electric and magnetic field





Joseph J. Thomson, *1856, †1940 (NP 1906)

Quantum Mechanics



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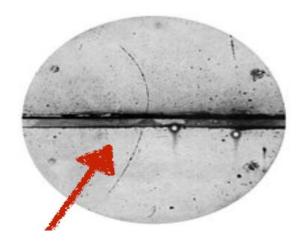
Anti-particles

This was achieved by Dirac in 1927,

$$i\gamma^{\mu}\partial_{\mu}\psi = m\psi$$

It permits (demands) solution with negative energy!

- Dirac proposed that the vacuum is made up of a "sea" with energy electrons
- What in 1930 was a huge issue for Dirac's equation became in 1931 its greatest's triumph, Andersen discovered the positron

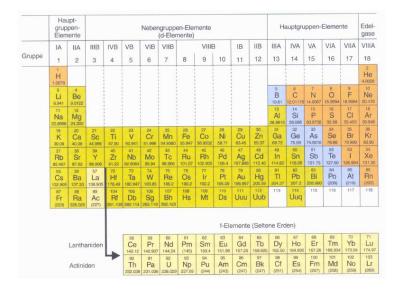






The strong force

What force binds protons and neutrons together to form the elements of the periodic table and just the world around us

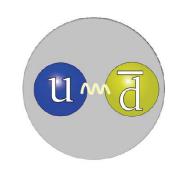


We know it has to be a force which is considerably stronger than the electric force since atoms made of 10's-100's protons stick together

We also know, from daily experience that a long distance ($^{\sim}$) a few 10⁻¹³ m) we do not experience the nuclear force

The mesons

Yukawa was the first to propose a model of strong force, in 1934. Proposed a new field, which mediated the interaction between the proton and the neutron



Since the force is short range, Yukawa made the mediator heavy, so it cannot propagate to long distances

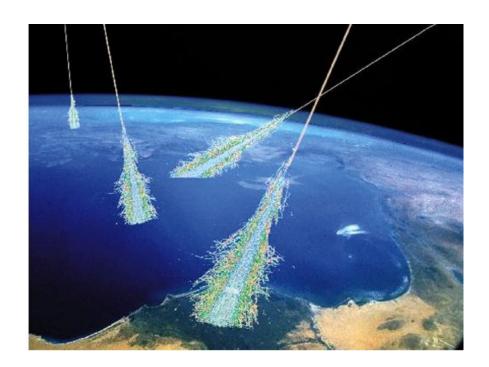
Yukawa reasoned this bosonic meson must be about 300 heavier than electron

The proton and neutron are examples of baryons which are fermionic

Pion discovery?

Around this time, one of the primary particle physics laboratories was the study of energetic particles hitting the earth's upper atmosphere.

In 1937 Anderson (again) and Neddermeyer discovered a particle which appeared consistent with Yukawa's meson.



This particle has a mass of around 200 times the electron.

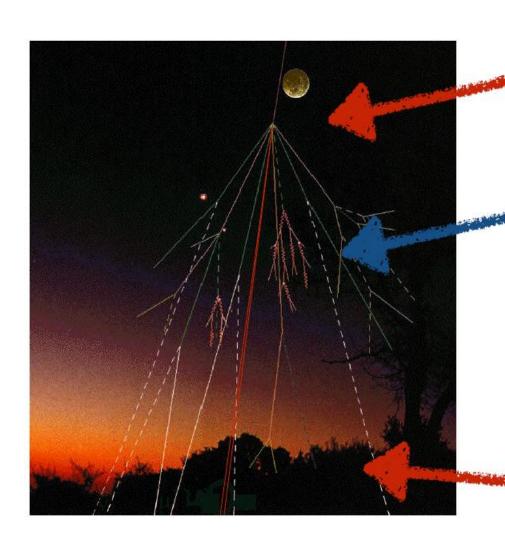
A puzzle

The problem with this new particle was that it did not interact strongly enough with nuclei to be the carrier of the strong force.

In fact, what was discovered was the muon, a heavier copy of the electron.

The actual pion (lightest mezon) was found in 1947 by Powel

The muon



Pions produced in initial cosmic ray interaction

Pions decay, e.g. through

$$\pi^- \to \mu^- + \overline{\nu}$$

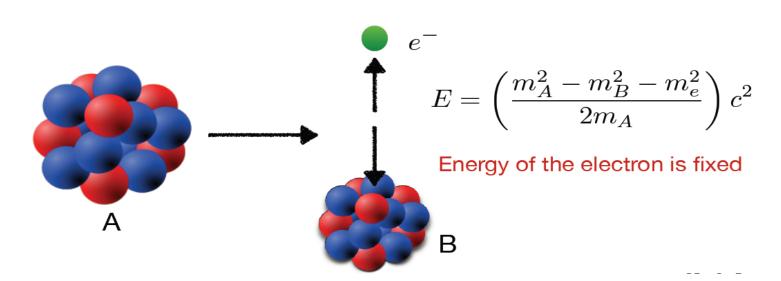
At lower altitude we detect mostly muons.

A puzzle

The discovery of the muon was completely unexpected.

The world of leptons was becoming increasingly more interesting ...

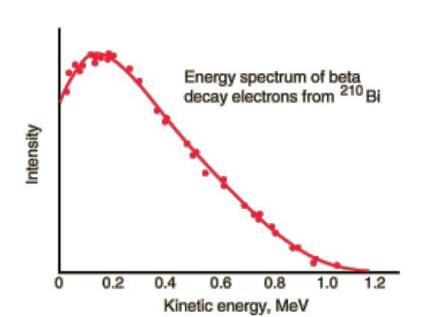
Radioactive decays provided a mystery, if a heavy nuclei A decays at rest into an electron and a lighter nuclei B then the conservation of energy fixes the energy of the electron



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Neutrinos

This is not what is actually observed!



G. J. Neary, Roy. Phys. Soc. (London), A175, 71 (1940).

Pauli proposed that a new, very weakly interacting neutral particle is present, and restores the conservation of energy.

Fermi constructed a beautiful description of beta decay which incorporated Pauli's neutrino in

1933.

Neutrinos

Neutrinos appear everywhere

Beta decay
$$n o p + e^- + \overline{\nu}$$

Pion decays
$$\pi^- o \mu^- + \overline{\nu}$$

Muon decay
$$\mu^-
ightarrow e^- + 2
u$$

This is all nice theoretically, but is there any experimental evidence?

Are neutrinos and anti-neutrinos different particles?

Neutrinos

In the 1950's Cowan and Reines used the inverse beta decay

$$\overline{\nu} + p^+ \rightarrow n + e^+$$

To conclusively prove the existence of anti-neutrinos

The non-observation of process

$$\overline{\nu} + n \rightarrow p + e^-$$

proved that neutrinos and antineutrinos are fundamentally different particles



Frederick Reines (left) and Clyde L. Cowan, Jr. with the control equipment used in their first tentative observations of the neutrino at Hanford, Washington, in 1953. Their definitive detection of the (anti) neutrino was performed at Savannah River, Georgia, three years later. (Courtesy General Electric Co.)

Lepton Number

The non-observation of the crossed process was theoretically expected based upon the idea of Lepton number conservation.

Matter particles (electrons, muons, neutrinos) are assigned lepton-number L=+1, whilst anti-mater particles (positrons, anti-muons and anti-neutrinos) are assigned L=-1, all other particles have L=0

Lepton number conservation then ensures that both LHS and RHS of reactions have the same total Lepton number.

$$\overline{\nu} + p^+ \to n + e^+$$
 $\downarrow : +1 +0 => 0 + 1$
 $\overline{\nu} + n \Rightarrow p + e^ \downarrow : +1 +0 => 0 - 1$

Lepton Number

Lepton number conservation allows us to rewrite the decay of muon

$$\mu^- \rightarrow e^- + 2\nu \longrightarrow \mu \rightarrow e^- + \nu + \overline{\nu}$$

What about?

$$\mu^- \rightarrow e^- + \gamma$$

$$Br(\mu^+ \to e^+ \gamma) < 5.7 \times 10^{-13}$$
 MEG experiment in PSI

Lepton Number

The lack of muon decays to electron and a photon suggests a conservation law of "mu-iness"

But what about regular muon decays?

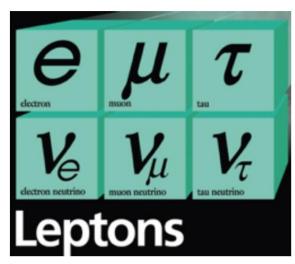
$$\mu \rightarrow e + \nu + \overline{\nu}$$

There must be (at least) two types of neutrinos

$$\mu^- \rightarrow e^- + \nu_\mu + \overline{\nu}_e$$

The Lepton Family

We know now that there are 3 families of leptons, the tau lepton was discovered in 1975 and tau neutrino in 2000 (DONUT experiment)

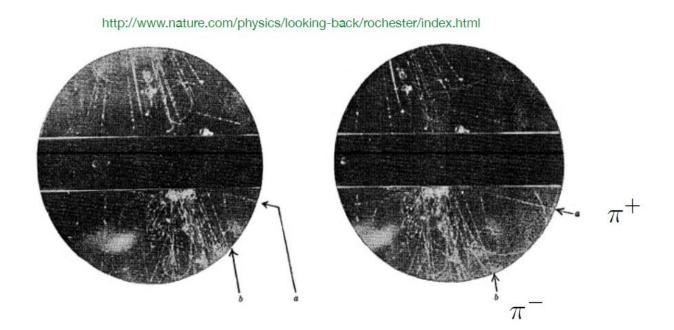


$$m_e = 0.511$$
 MeV $m_{\mu} = 105.7$ MeV $m_{\tau} = 1.78$ GeV

However in 1947 it was no idea "what for" are muons (I. Rabi)



The strange tide



In December 1947, Rocheter and Butler observed a "strange" new particle in cosmic ray experiments. $K^0 \to \pi^- + \pi^+$

The particle appeared to be a new type of meson!

The strange tide

The Kaon was not the last "strange" particle to be discovered, in 1950 Anderson (again!) discovered a new heavy baryon through its decay products.

$$\Lambda \to p + \pi^-$$

How do we know this is a Baryon? Similarly why are we here?

Or

$$p \to e^+ + \gamma$$

 $p \to e^+ + \gamma + \nu_e$

These are pretty catastrophic processes for the Universe.

Baryon Number

In a similar fashion to what we saw for leptons, we introduce the concept of baryon number conservation

Proton and neutrons have a Baryon number of +1 (antibaryons have a Baryon number of -1)

There is no corresponding meson number conservation, for example we already now how pions decay

$$\pi^- \to \mu^- + \overline{\nu}$$

Baryon number thus give a mechanism to quantify whether hadrons are baryons, or mesons, based on the production and decay mechanism, and baryon number conservation (you can use spin too, baryons are fermions, mesons are bosons).

The strange particles

The new strange particles are indeed "strange"

- They are copiously produced in energetic collisions (with timescales around 10⁻²³ seconds). But relatively long-lived, with lifetimes around 10⁻¹⁰ second.
- This suggests that maybe there is different production mechanism and decay mechanism (i.e. two separate forces)
- This is further backed up by noting that strange particles are produced in pairs e.g.

$$\pi^- + p \rightarrow K^+ + \Sigma^-$$

• But we can never singly produce a strange particle $\pi^- + p \not\to \pi^+ + \Sigma^-$

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The weak force and strangeness

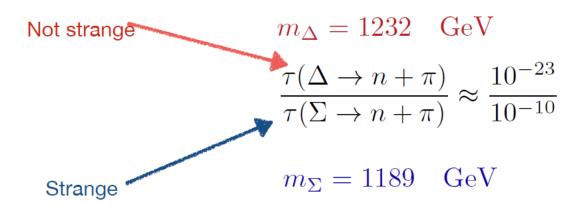
- All this suggests yet another quantum number, strangeness, which is conserved by the strong nuclear force, but not by the weak force.
- The K's have strangenes S=+1 and the ∑ has S=-1, ordinary matter has S=0.
- The strong and electromagnetic force conserve strangeness, but the weak force doesn't which allows strange particles to decay.

$$\Sigma^+ \to p^+ + \pi^0$$

 The introduction of a new force should not be surprising, we know that neutrinos, exist, but they are a) not charged, b) not strongly interacting, so one expect some new type of force, this is indeed the same weak force we see above.

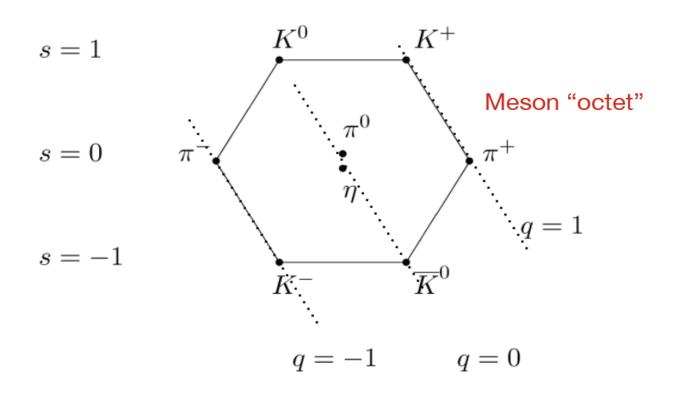
We are around 1950'ties

- We learned that there were strange hadrons, which are pair produced in strong interactions, but decay via a new weak force.
- On the other hand, there are some hadrons, which are not strange, and decay via the strong force, for example lifetime ratio below:



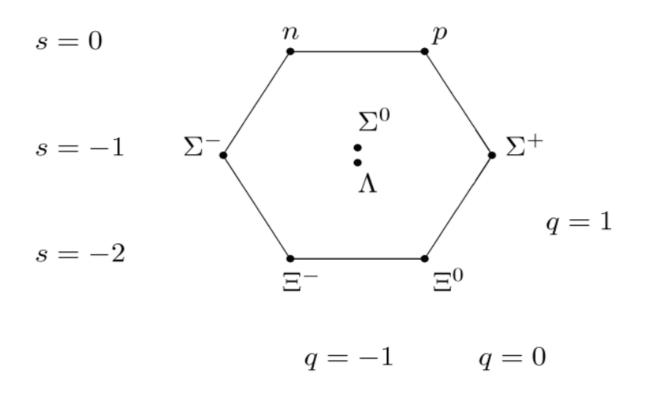
We want to understand this better: The Quark Model

The "Eightfold way"



The first person who made significant progress on forming a "periodic table" for hadrons was Murray Gell-Mann 1961. The mesons can be grouped into the octed, based on charge and strangeness

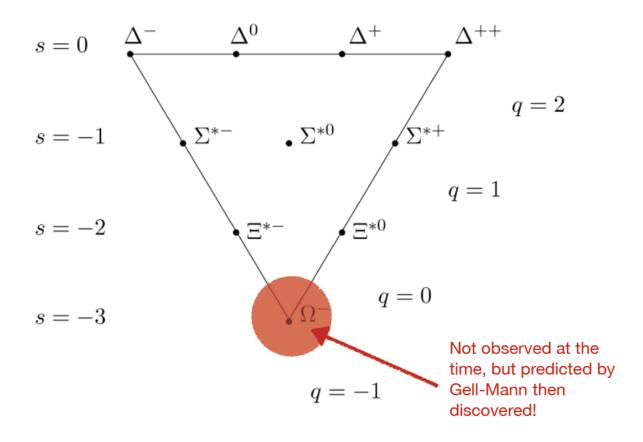
The "Eightfold way"



A similar octed appears for baryons, which includes also our friends: p and n

(note the shift in scale for the strangeness)

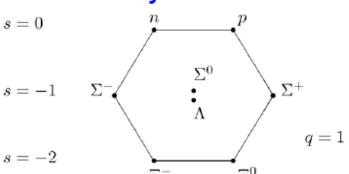
More baryons



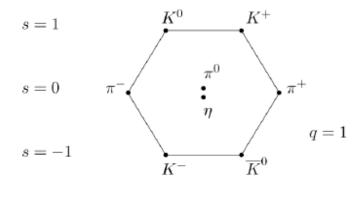
Nonet is not the only structure, spin 3/2 baryons form a decoupled.

The "Eightfold Way"

Baryon octet





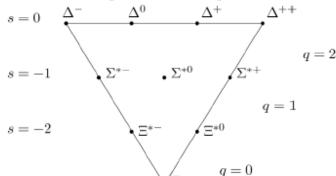


Meson octet

q = 0

q = -1

Baryon decuplet



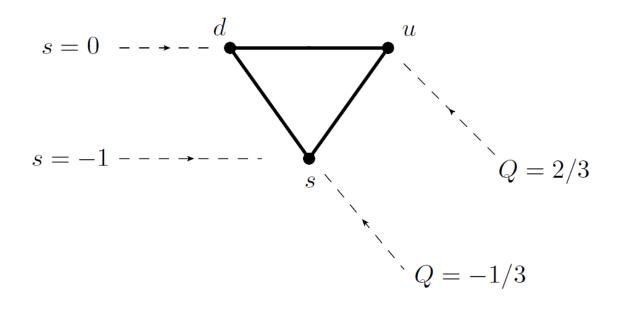


M. Gell-Mann (NP 1969)

$$q = -1$$

- Ordering Mesons and Baryons in octets (spin 1/2) and decuplets (spin 3/2), according to quantum numbers (i.e. quark composition)
- Prediction of baryons Ω^{-} with three *s* quarks (observed in 1964)
 - → Pauli's exclusion principle
 - → new quantum number (color)
- Number of colors from hadronic to leptonic branching ratio of e⁺e⁻ collisions

The Quark Model



In 1964 Gell-Mann proposed a model to explain observed in the hadron spectrum. He postulated that there are 3 quarks, with the quantum numbers shown above. Using these quarks and corresponding anti-quarks) he was able to reproduce the observed spectrum of hadrons.

The Quark Model

In the quark model the hadrons are constructed according to the following rules:

1) Baryons are made of three quarks (qqq), anti-baryons are made of anti-quarks

$$p = uud$$

$$n = udd$$

$$\Omega^{-} = sss$$

$$\Sigma^{*+} = uus$$

2) Mesons are made of quark anti-quark pair.

$$\pi^+ = u\overline{d}$$

$$K^+ = u\overline{s}$$

Problems for the Quark Model

The quark model has a serious flaw as it currently stands, namely it predicts:

$$\Delta_{3/2}^{++} = u_{\uparrow} u_{\uparrow} u_{\uparrow}$$

Worse still, such a doubly charged spin 3/2 baryon was observed by Fermi in 1951, and lives in our decoupled

As it stands our states are completely symmetric indistinguishable quantum state made up of spin ½ fermions, and is in complete violation of our understanding of spin statistics

A second less catastrophic problem, is the lack of existence of qq and other allowed states. Foe example n uu charge 4/3 meson has ever been observed, but there is no reason to not construct it in our simple quark model.

Color: the "charge" of a strong force

We solve both of these problems with the introduction of the color charge.

We re-write the wave function as follows,

$$\Delta_{3/2}^{++} = u_{\uparrow}^R u_{\uparrow}^G u_{\uparrow}^B$$

We introduce three color charges, and assign a different one to each up quark, now the particles ar distigushable and the spin-statistics theorem is satisfied.

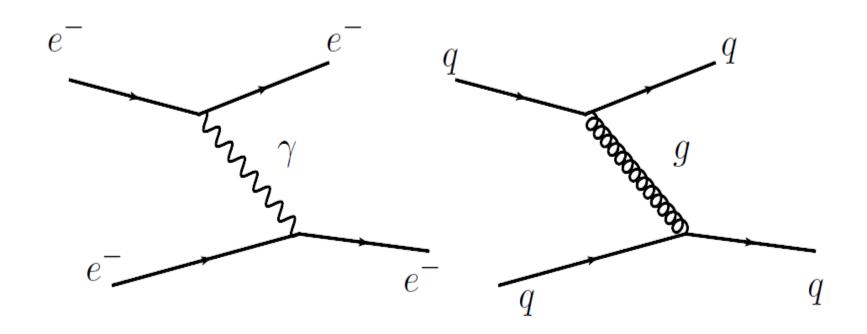
Color: the "charge" of a strong force

Since we don't "see" color like we do charge, we also enforce that all observed hadrons are "white", i.e. color neutral combinations

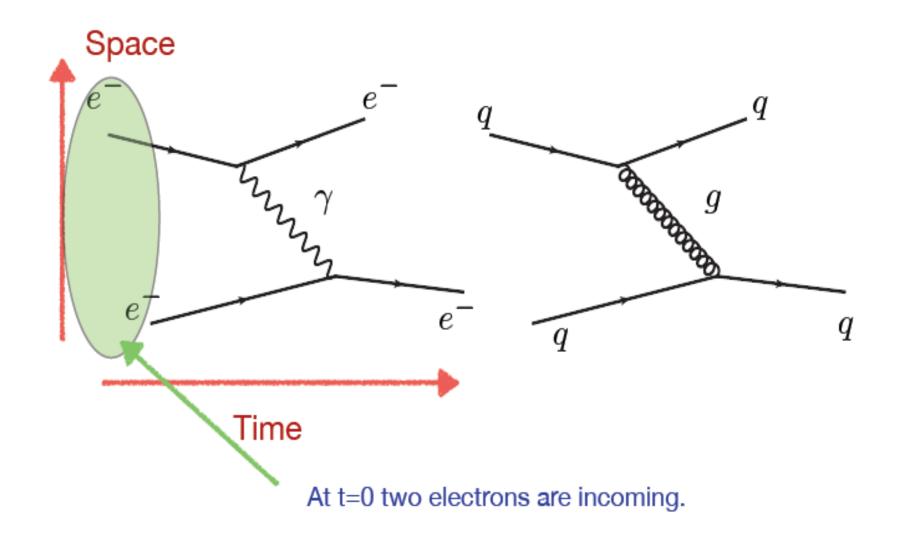
Color neutral (white) combinations are given by either combining RGB, a=ro a color wit its anticolor (in a meson) e.g. R \overline{R}

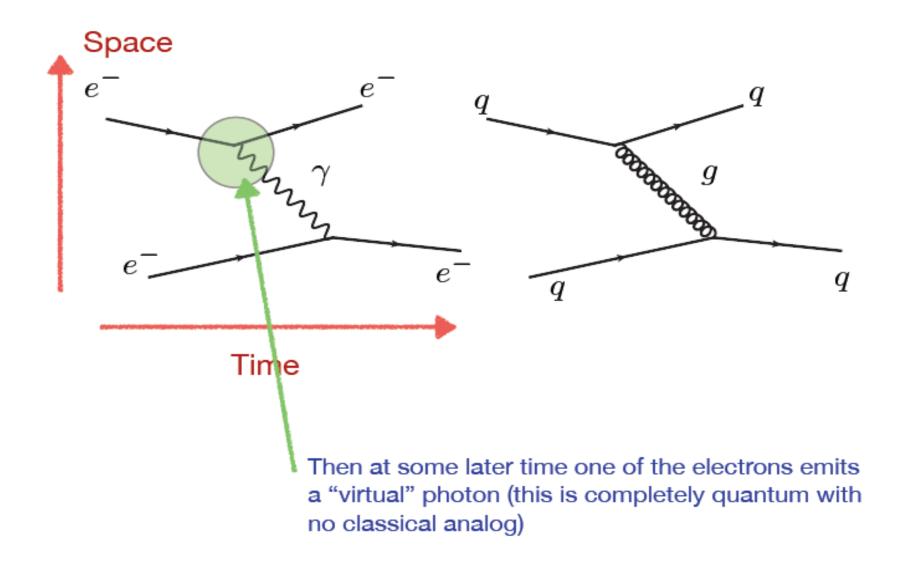
$$p = "RGB"$$
 $n = "RGB"$ $\pi = "R\overline{R} + G\overline{G} + B\overline{B}"$

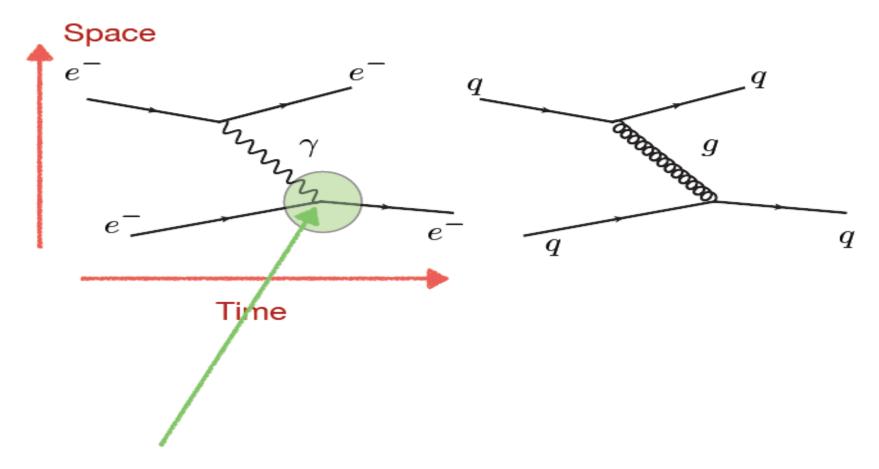
The gluon: the "photon" of the strong force



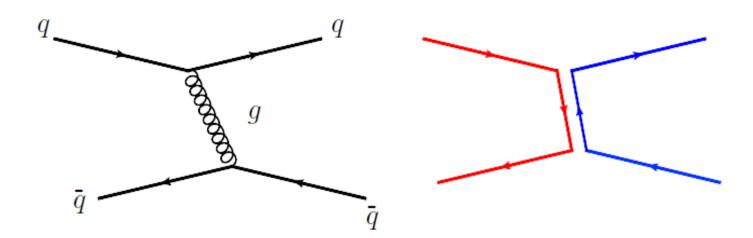
Here we discuss it only heuristically, I will come back to it latter







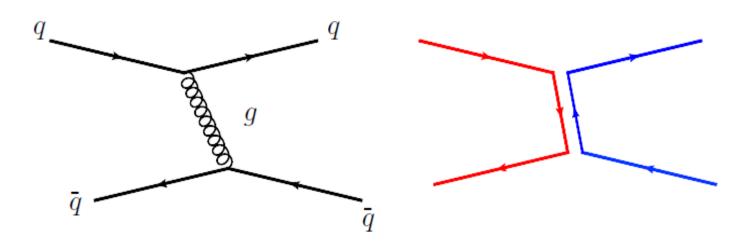
Finally the second electron, reabsorbs the photon, and feels the Coulomb repulsion from the first electron.



There must be an equivalent messenger of the strong force. Since it is responsible for "gluing" quarks together inside hadrons, we call it the gluon.

However, unlike the case of electron scattering, there are multiple color charges that interact with each other. In the above a red quark scatter off an anti-red anti-quark, producing new blue quark/anti-quark in the final state.

Therefore in order to mediate the strong force, it is clear that the gluon itself must be colored, and it carriers two indices.



Simple counting suggest there should be nine combinations of gluons pairings

RR, RG, RB, BR, BG, BB, GR, GG, GB

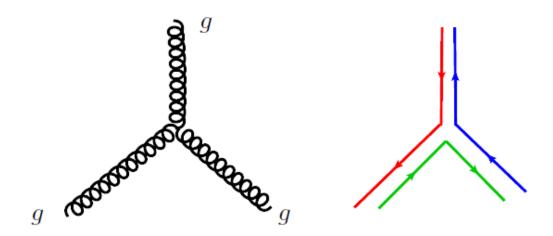
However, we already know that the combination

$$RR + BB + GG$$

Is a color singlet (this is a color charge of a pion), so we can actually eliminate one color combinations from the basis.

Therefore ther are 8 gluons!

Self interactions of gluons



A further, fundamental difference between photons and gluons, is that gluons, which are charged under the strong force (i.e. are colored objects) can interact with themselves.

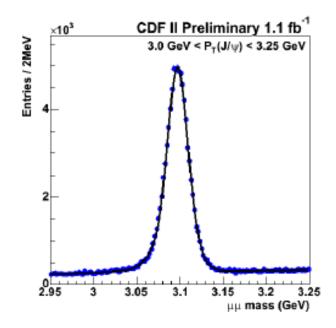
This results in dramatically different physics from electromagnetic interactions and ultimately leads to the bounding of nuclei together.

Theories, like electromagnetism, in which the boson mediator is not charged under the interaction are called Abelian, theories like the strong force, in which gluon boson itself carries the color charge, are called Non-Abelian

The J/Ψ discovery

Initially reaction to the quark model was mixed, its simplicity and predictivity was admired. But the inability to produce free quarks, and the ad hoc nature of the color interpretation were serious issues.

In 1974 the game changed with the discovery of a new very heavy meson (3 times heavier than the proton). This particle was (almost) simultaneously discovered at Brookhaven and SLAC



It was easy to incorporate this new meson into the quark model, by adding in a new heavy quark, the **charm quark**

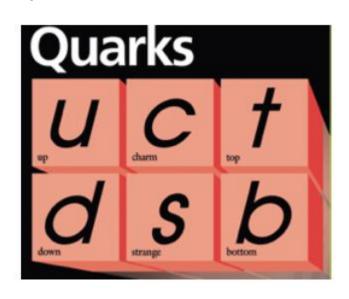
$$J/\Psi = c\overline{c}$$

The quark model then predicts many new baryons and mesons which contain charm quark

Completing the quark story

Since then two additional quarks have been discovered, the bottom quark (1977) and top quark (1995) were discovered at Fermilab.

The third generation of quarks are very heavy, the bottom quark has a mass around 5 x the mass of the proton, whilst the top quark has a mass around 175 times that of the proton!

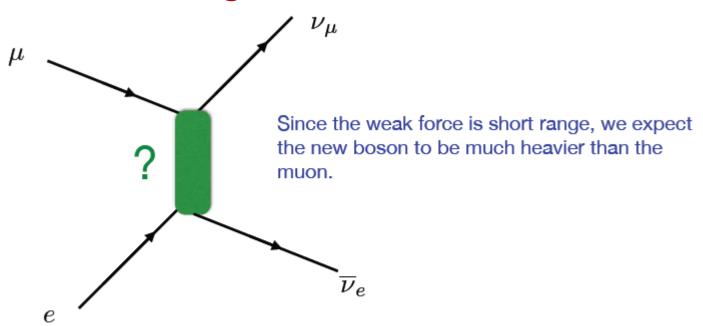


We still believe that quarks are fundamental and that Standard Model includes a fundamental particle heavier than a Gold nuclei!

Massive vector bosons

Fermi's theory for muon decay related the decay to a 4 fermion contact interaction.

However, we know that this is only true at low energies, we know that there has to be force carrying boson, like the photon and the gluon.

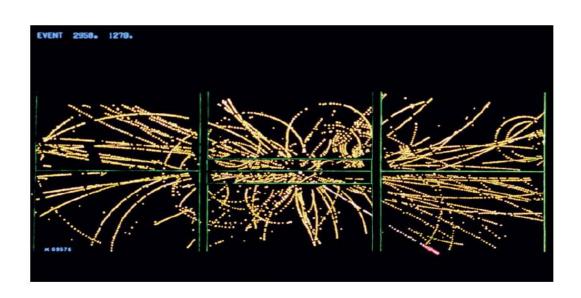


The search for massive vector bosons, to mediate the weak force, was on!

The W and Z boson

The hunt for the W and Z bosons was finished in 1983, when they were discovered at CERN. The bosons are heavy, with masses,

$$m_W = 80.4 \text{ GeV}$$
 $m_Z = 91.2 \text{ GeV}$

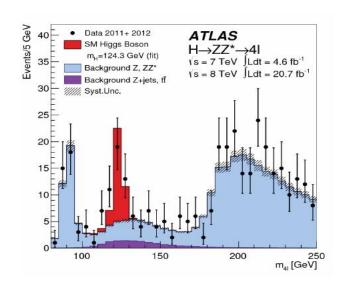


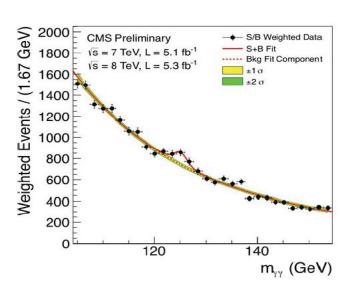
The Higgs boson

The final piece of the Standard Model was the Higgs boson.

Theoretically it was desperately needed, it provided a mechanism in which the W and Z could require masses without spoiling theory.

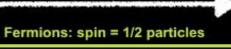
You've probably heard, but the Higgs was discovered in 2012 at CERN.

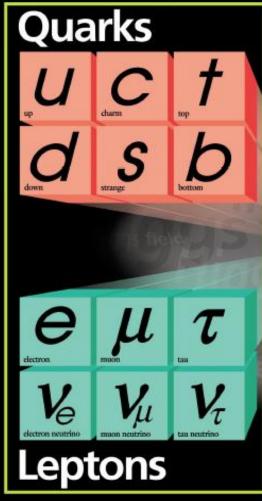


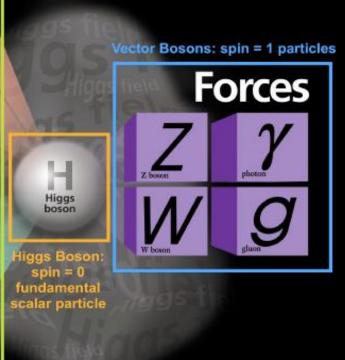


The Standard Model

Heavier copies







Fermions and Boson

Fundamental characteristics of particles: Spin ("self angular momentum")

- Integer values (0, ±1, ...) → Bosons
- Half integer values (±1/2, ±3/2, ...) → Fermions

Bosons (Cooper pairs ...) can be described by **common** wavefunction → Funny effects (super conductivity, super

fluidity, ...)



Fermions (electrons or protons ...) must be in different states

→ Pauli's exclusion principle (basis of all chemistry ... and much more)

Satyendranath Bose



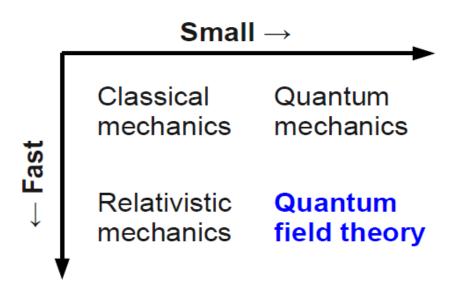
Enrico Fermi (NP 1938)



Wolfgang Pauli (NP 1945)



Quantum Field Theory



First major achievement: Dirac's equation for free electrons (and positrons)

$$E^{2} - \mathbf{p}^{2}c^{2} = m^{2}c^{4}$$

 $E = \pm \sqrt{\mathbf{p}^{2}c^{2} + m^{2}c^{4}}$

Interpretation of negative energies: sea of electrons → holes in sea act as positively charged electrons → confirmed by Anderson 1932



Paul Dirac (NP 1933)



Carl Anderson (NP 1936)

QFT – Gauge Interactions

Example:

Requirement: Lagrangian

$$\mathcal{L} = \frac{1}{2} (\partial_{\mu} \Phi)^{T} \partial^{\mu} \Phi - \frac{1}{2} m^{2} \Phi^{T} \Phi$$

invariant under special local transformations G(x)

$$\Phi \mapsto \Phi' = G\Phi$$

Invariance is in general not guaranteed, since

$$\partial_{\mu}(G\Phi)
eq G(\partial_{\mu}\Phi)$$

Introduce covariant derivatives (with gauge fields $A_{_{11}}$) $D_{\mu}=\partial_{\mu}+gA_{\mu}$

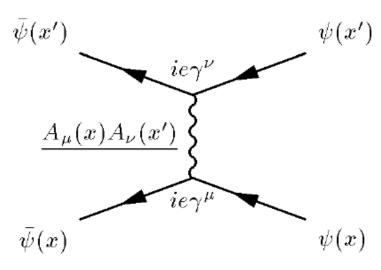
$$D_{\mu}=\partial_{\mu}+\mathsf{g}\mathsf{A}_{\mu}$$

- ightharpoonup Locally gauge invariant Lagrangian $\mathcal{L}_{loc} = \frac{1}{2} (D_{\mu} \Phi)^T D^{\mu} \Phi \frac{1}{2} m^2 \Phi^T \Phi$
- \rightarrow Interaction terms: $\mathcal{L}_{int} = \frac{g}{2} \Phi^T A_{\mu}^T \partial^{\mu} \Phi + \frac{g}{2} (\partial_{\mu} \Phi)^T A^{\mu} \Phi + \frac{g^2}{2} (A_{\mu} \Phi)^T A^{\mu} \Phi$
- → Gauge fields (gauge bosons) are interaction "particles"

General recipe for a QFT:

- Write down Lagrangian of mass, kinetic, and interaction terms
- Quantize fields
- Decribe scattering theory with quantized fields

Feynman Diagrams/Renormalisation



Feynman diagrams:

 $\psi(x')$ Powerful tool to write down scattering amplitudes (also for higher order perturbations)

Feynman rules:

Set of rules to get from a Feynman diagram to the mathematical expression



R. Feynman (NP 1965)

In QFTs perturbation theory is only valid for finite energy range → divergences in calculation of scattering amplitudes

Renormalization is a technique to remove these divergences





M. Veltman and G. 't Hooft (NP 1999)

Quantum Electro Dynamics

Lagrangian invariant under local ($\theta = \theta(x)$) phase transition: symmetry group U(1)

 $\psi \mapsto e^{i\theta}\psi$

Appropriate covariant derivative, with gauge field A

$$D_{\mu} = \partial_{\mu} - i \frac{e}{\hbar} A_{\mu}$$

→ QED Lagrangian

$$\mathcal{L}_{\mathsf{QED}} = ar{\psi} (ihc\,\gamma^{\mu}D_{\mu} - mc^2)\psi - rac{1}{4\mu_0} F_{\mu
u} F^{\mu
u}$$

 A_{μ} is the photon field!

Side remark: Generators of U(1) group commute (abelian group) → Gauge bosons have no charge (photons are electrically neutral)→ no self interaction

As in classical mechanics (Noether's Theorem):

Symmetry ← conserved quantity

Electroweak Theory

Local symmetry:

 $SU(2)_{L} \otimes U(1)$

Charges:

- Weak isospin I₃
- Weak hypercharge Y = 2Q 2I₃



A. Salam, S. Glashow, and S. Weinberg (NP 1979)

Gauge fields: W, B mix to W^{\pm} , Z^{0} , and γ

Generators of SU(2) do not commute (non-abelian group)

→ Self interaction

By construction all gauge bosons are massless (to keep local invariance)

However, short range of interaction implies heavy particles

Electroweak Gauge Bosons

- Observation of neutral current in 1973
- Observation of W in May 1983 (and a few month later also the Z) at UA1 and UA2

 $M_{W} \sim 80.4 \, \text{GeV}$

 $M_{z} \sim 91.2 \, \text{GeV}$

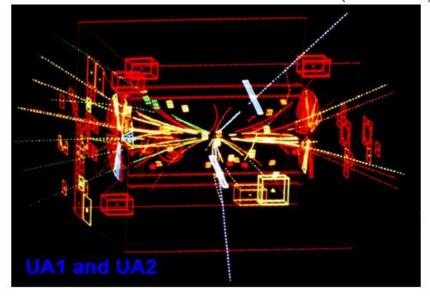




C. Rubbia and S. van der Meer (NP 1984)

→ Heavy particles!





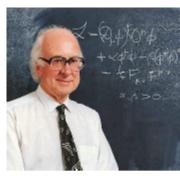
The Higgs mechanism

Formalism to obtain massive gauge bosons while keeping invariance of L intact by spontaneous symmetry breaking

Independently derived by Englert, Brout, Kibble, Hagen, Guralnik

Lagrangian of complex scalar Higgs field φ

$$\mathcal{L}_{\mathsf{Higgs}}(\phi, A) = (\hat{D}_{\mu}\phi)^{\dagger}(\hat{D}^{\mu}\phi) + \mu^{2}\phi^{\dagger}\phi - \lambda(\phi^{\dagger}\phi)^{2}$$



Peter Higgs

If
$$\mu^2 > 0 \rightarrow \text{trivial minimum at } \phi = 0$$

If
$$\mu^2 < 0 \rightarrow$$
 non trivial, degenerated minimum

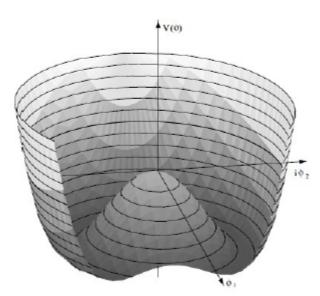
Vaccum expectation value
$$|\langle 0|\phi|0\rangle|=v=\sqrt{\frac{-\mu^2}{2\lambda}}$$

A. Choose arbitrary ground state in minimum

- A. Choose arbitrary ground state in minimum
- B. Expand Higgs field
- C. Covariant derivatives → massive gauge bosons

$$M_{W^{\pm}}^2 = \frac{g^2 v^2}{4}$$
 $M_{Z^0}^2 = \frac{(g^2 + g'^2)v^2}{4}$

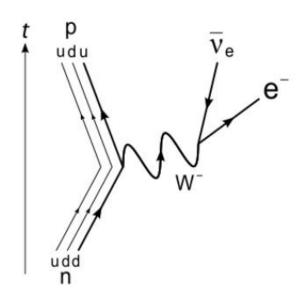
With g and g' as coupling constants of SU(2), and U(1)

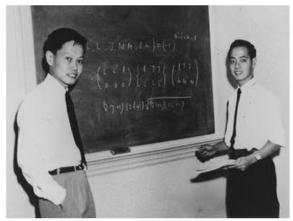


Mexican hat potential

EW Decay - C and P violation

- Classical gravitation, electromagnetism, and the strong interaction are invariant under charge C and parity P ("Mirror symmetry") transformations
- Yang and Lee: P could be violated in EW interactions
- Observation by Wu: cryogenic Co⁶⁰ in strong magnetic field → strong asymetry of direction of emitted electrons → maximal P violation
- C and P both violated by EW interactions, but CP (mostly) concerved

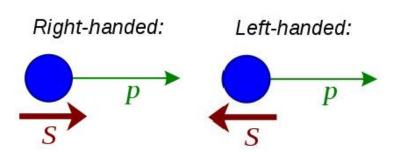




C. N. Yang and T.-D. Lee (NP 1956)



C. S. Wu

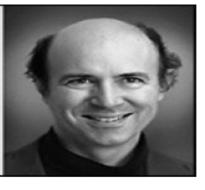


Quantum Chromodynamisc

- All ordinary matter consists of fermions and the largest fraction is carried by nucleons made of quarks and gluons
- Strong interaction: SU(3) → 8 massless gluons as gauge bosons
 - Quarks (and gluons) carry "color" charge
 - Gluon self interaction
 - → Coupling constant decreasing with energy: "assymptotic freedom" at high energies
 - → Quarks and gluons don't exist as free particles but as color neutral bound states: Mesons (quark-antiquark) or Baryons (three quarks)







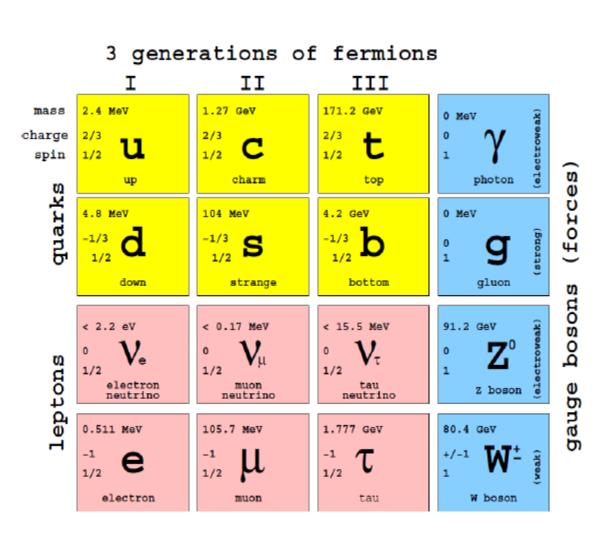
D. Politzer, D. Gross, F. Wilczek (NP 2004)

The Standard Model

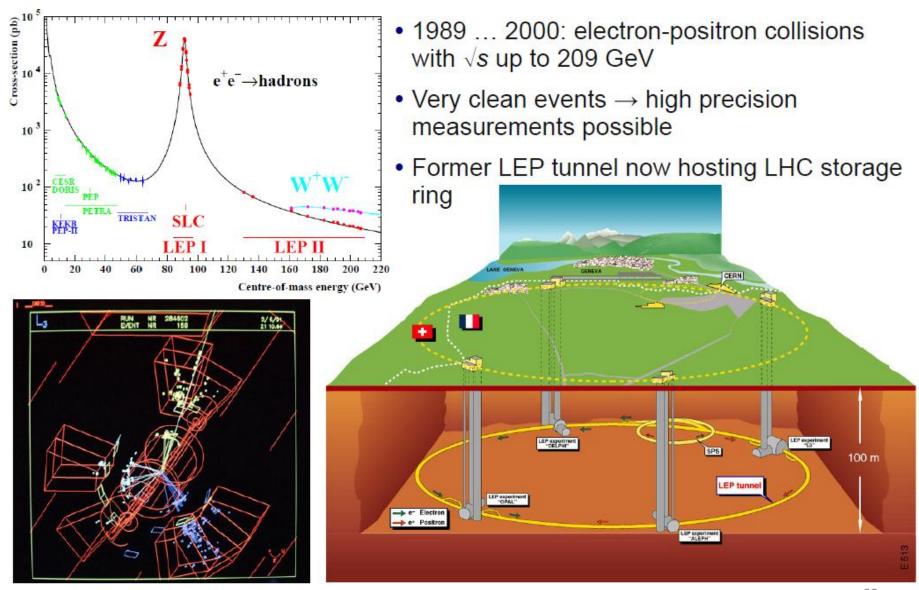
Interactions described by gauge groups:

U(1)⊗SU(2)_L⊗SU(3)_c

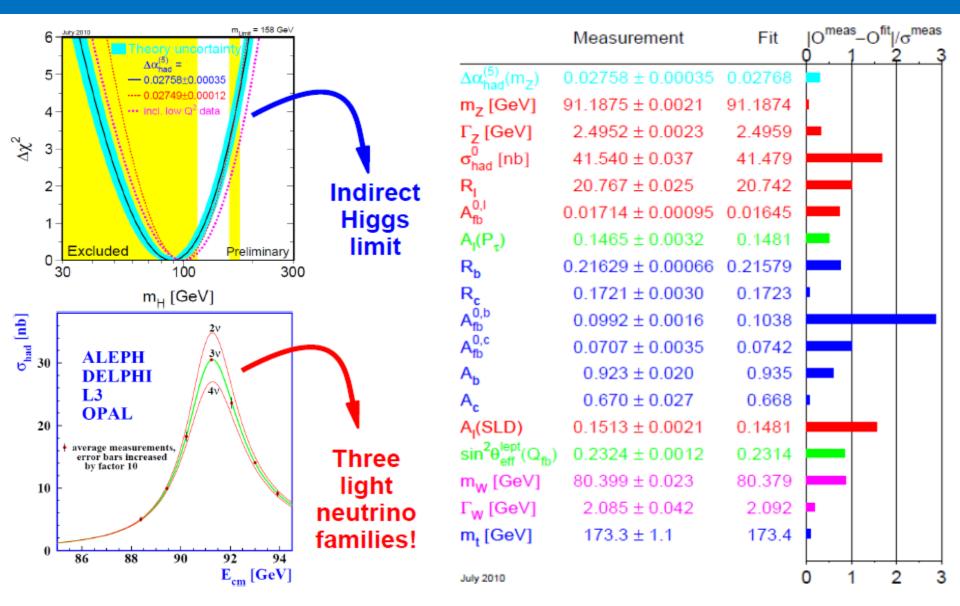
- Higgs boson last missing particle
- All measurements from colliders in excellent agreement with the SM



Large Electron Positron Collider



Electroweak precision data



4-th of July 2012, ATLAS and CMS experiments at Large Hadron Collider at CERN announced observation of a new particle, consistent with predicted by so called "Higgs mechanism".

2013 NOBEL PRIZE IN PHYSICS

François Englert Peter W. Higgs



THE BEH-MECHANISM, INTERACTIONS WITH SHORT RANGE FORCES AND SCALAR PARTICLES



8 October 2013

The Royal Swedish Academy of Sciences has decided to award the Nobel Prize in Physics for 2013 to

François Englert and Peter Higgs

"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

The neutrino

- Continous spectrum of electrons from β-decay
 - → new undetectable particle predicted by Pauli in 1930
- First observation by Poltergeist experiment in inverse β-decays in 1957

$$\bar{\nu}_e + p \rightarrow e^+ + n$$



F. Reimes (NP 1995)

Observation of the myon neutrino in 1962

And much more spectacular physics:

- Solar neutrino problem ...
- Neutrino oscillations ...
- Neutrino masses ...
- Cosmological neutrinos ...
- Neutrinos from super novae ...



L. Ledermann, M. Schwartz, J. Steinberger (NP 1988)

The neutrino oscilation

Nobel Prize in physics in 2015

Takaaki Kajita and Artur B. McDonald

for discovery of the neutrino oscillations, which shows that neutrino has mass.

Nobel Prizes in Elementary Particle Physics



Sin-Itiro Tomonaga



Julian Schwinger



Richard P. Feynman







GREEN - theoretical - experimental **BLUE**

1964: "Higgs mechanism" was born



Leon M. Lederman



Melvin Schwartz



Jack Steinberger





1969 - M. Gell-Mann

1976 – B. Richter and S. Ting

1979 – S.L. Glashow, A. Salam, S. Weinberg

1980 – J. Cronin, V. Fitch

1984 – C. Rubbia, S. van der Meer

1988 – L. M. Lederman, M. Schwartz, J. Steinberger

1990 – J. Friedman, J. Kendall, R. Taylor

1992 - G. Charpak

1995 – M. Perl, F. Reines

1999 - G. tHooft, M. J. Veltman

2004 - D. J. Gross, H. D. Politzer, F. Wilczek

2008 – Y. Nambu, M. Kobayashi, T. Masakawa

2013 – F. Englert and P. Higgs

2015 - T. Kajita and A. B. McDd

2012: "Higgs particle" was discovered



Gerardus 't Hooft



Simon van der Meer

Martinus I.G. Veltman





M. Gell-Mann

And then what?

Is the end of the story? For the SM yes, but for fundamental physics, No!

We know that SM is not the complete theory of nature, some of tis problem are:

- No dark matter
- Not enough CP violation to account for Matter/Anti-Matter asymmetry in the Universe.
- No explanation of the particle specrum, why is the Higgs so light?
- Neutrinos?

— . . .