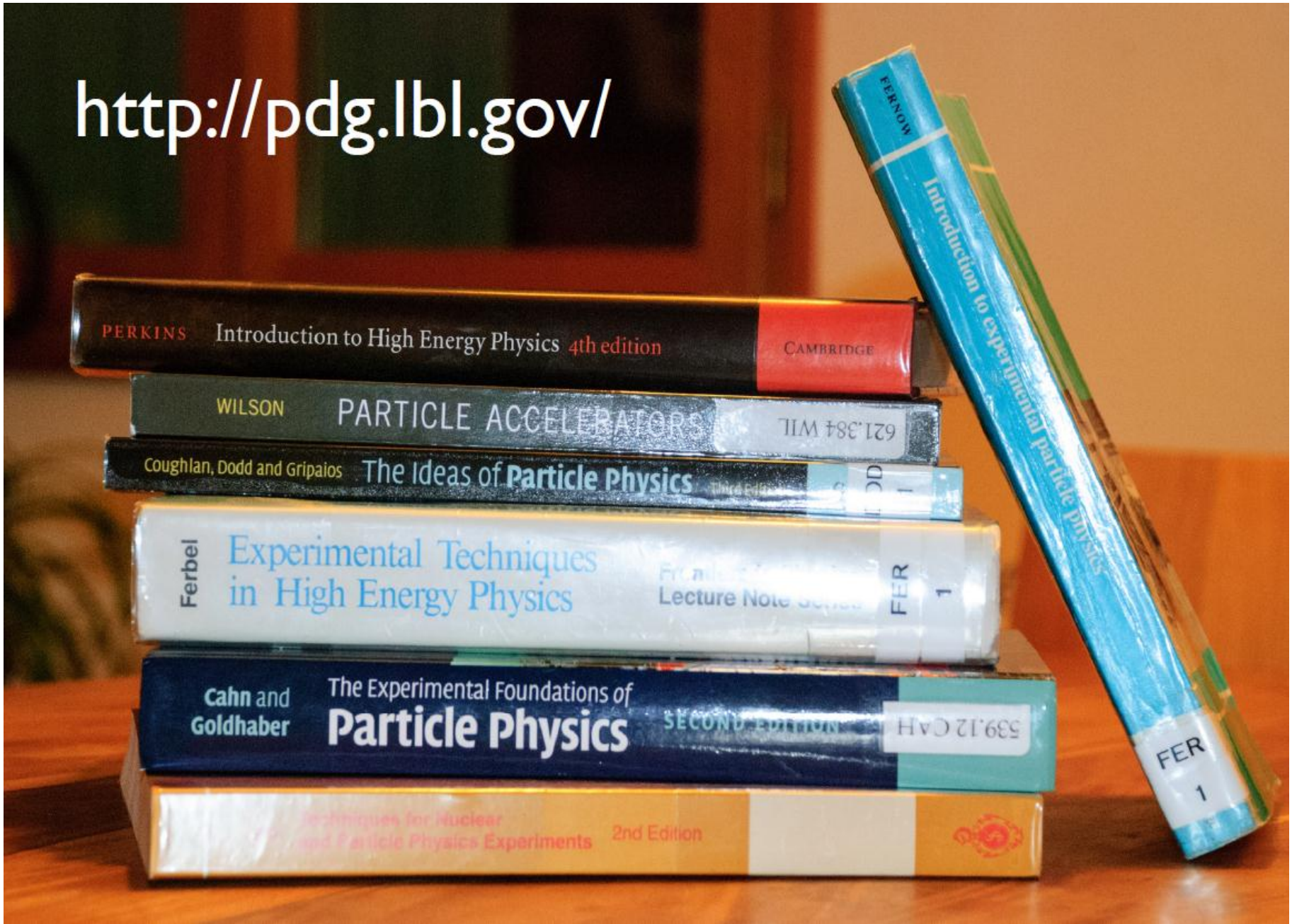


Introduction to particle physics: experimental part

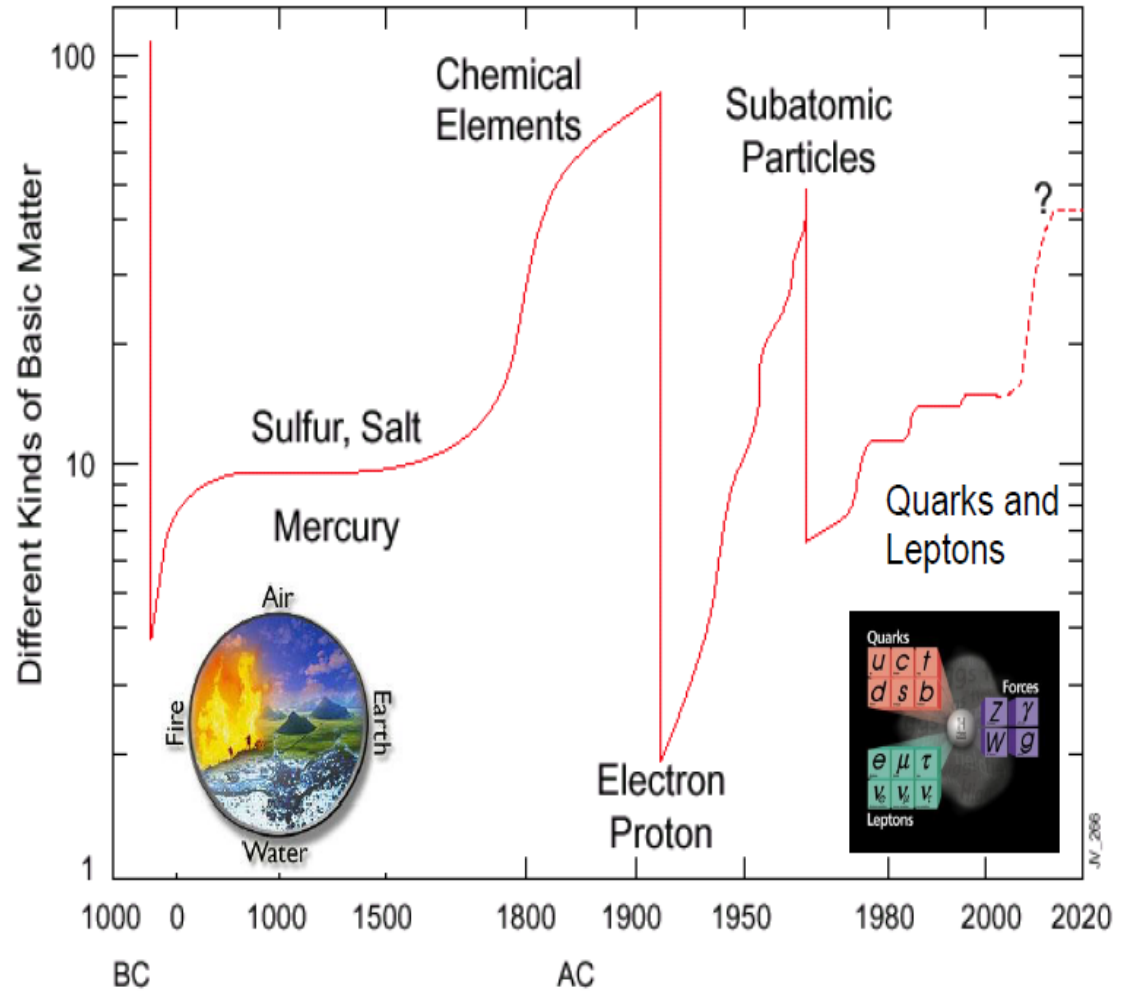
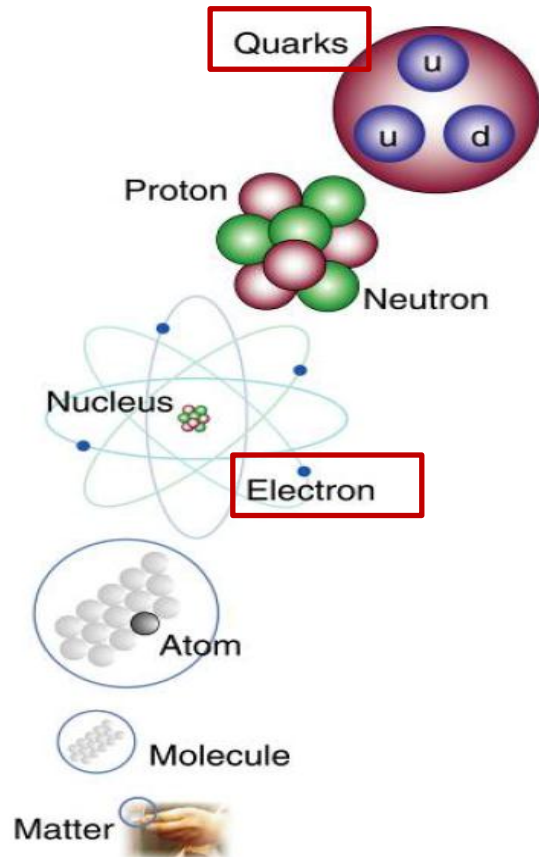
Introduction (ones more) to Standard Model
Accelerators
CERN and LHC

<http://pdg.lbl.gov/>



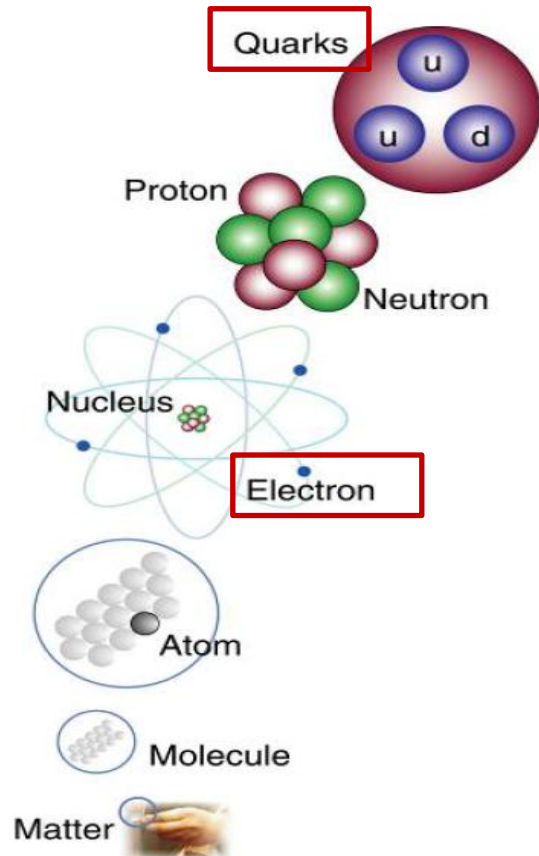
Constituents of matter along History

Quantum mechanics



Particles of the Standard Model

Quantum mechanics



Matter particles
($< 10^{-16}$ cm)

Interaction particles

<p>2.4M $\frac{2}{3}$ u up $\frac{1}{2}$</p>	<p>1.27G $\frac{2}{3}$ c charm $\frac{1}{2}$</p>	<p>171.2G $\frac{2}{3}$ t top $\frac{1}{2}$</p>	<p>g gluon 1</p> <p>strong nuclear force (color charge)</p>
<p>4.8M $-\frac{1}{3}$ d down $\frac{1}{2}$</p>	<p>104M $-\frac{1}{3}$ s strange $\frac{1}{2}$</p>	<p>4.2G $-\frac{1}{3}$ b bottom $\frac{1}{2}$</p>	
<p>0.511M -1 e electron $\frac{1}{2}$</p>	<p>105.7M -1 μ muon $\frac{1}{2}$</p>	<p>1.777G -1 τ tau $\frac{1}{2}$</p>	<p>γ photon 1</p> <p>electromagnetic (charge)</p>
<p>< 2.2 0 ν_e e-neutrino $\frac{1}{2}$</p>	<p>$< 0.17M$ 0 ν_μ μ-neutrino $\frac{1}{2}$</p>	<p>$< 15.5M$ 0 ν_t t-neutrino $\frac{1}{2}$</p>	<p>W^{+} $+1$</p> <p>Z 0</p> <p>weak nuclear force</p>



Higgs particle
Is not a matter particle and
not an interaction particle

Nobel Prizes in Elementary Particle Physics



Sin-Itiro Tomonaga



Julian Schwinger



Richard P. Feynman



Sheldon Lee Glashow



Abdus Salam



Steven Weinberg

GREEN - theoretical
BLUE - experimental

1964: „Higgs mechanism”
was born



Leon M. Lederman



Melvin Schwartz



Jack Steinberger

1957 – C. N. Yang, T. Lee

1965 – S. I. Tomonaga, J. Schwinger, R.P Feynman

1969 – M. Gell-Mann

1976 – B. Richter and S. Ting

1979 – S.L. Glashow, A. Salam, S. Weinberg

1980 – J. Cronin, V. Fitch

1984 – C. Rubbia, S. van der Meer

1988 – L. M. Lederman, M. Schwartz, J. Steinberger

1990 – J. Friedman, J. Kendall, R. Taylor

1992 - G. Charpak

1995 – M. Perl, F. Reines

1999 - G. tHooft, M. J. Veltman

2004 - D. J. Gross, H. D. Politzer, F. Wilczek

2008 – Y. Nambu, M. Kobayashi, T. Masakawa

2013 – F. Englert and P. Higgs

2012: „Higgs particle”
was discovered

2015 - T. Kajita and A. B. McDonald



Carlo Rubbia



Simon van der Meer



Georges Charpak



Gerardus 't Hooft



Martinus J.G. Veltman



M. Gell-Mann

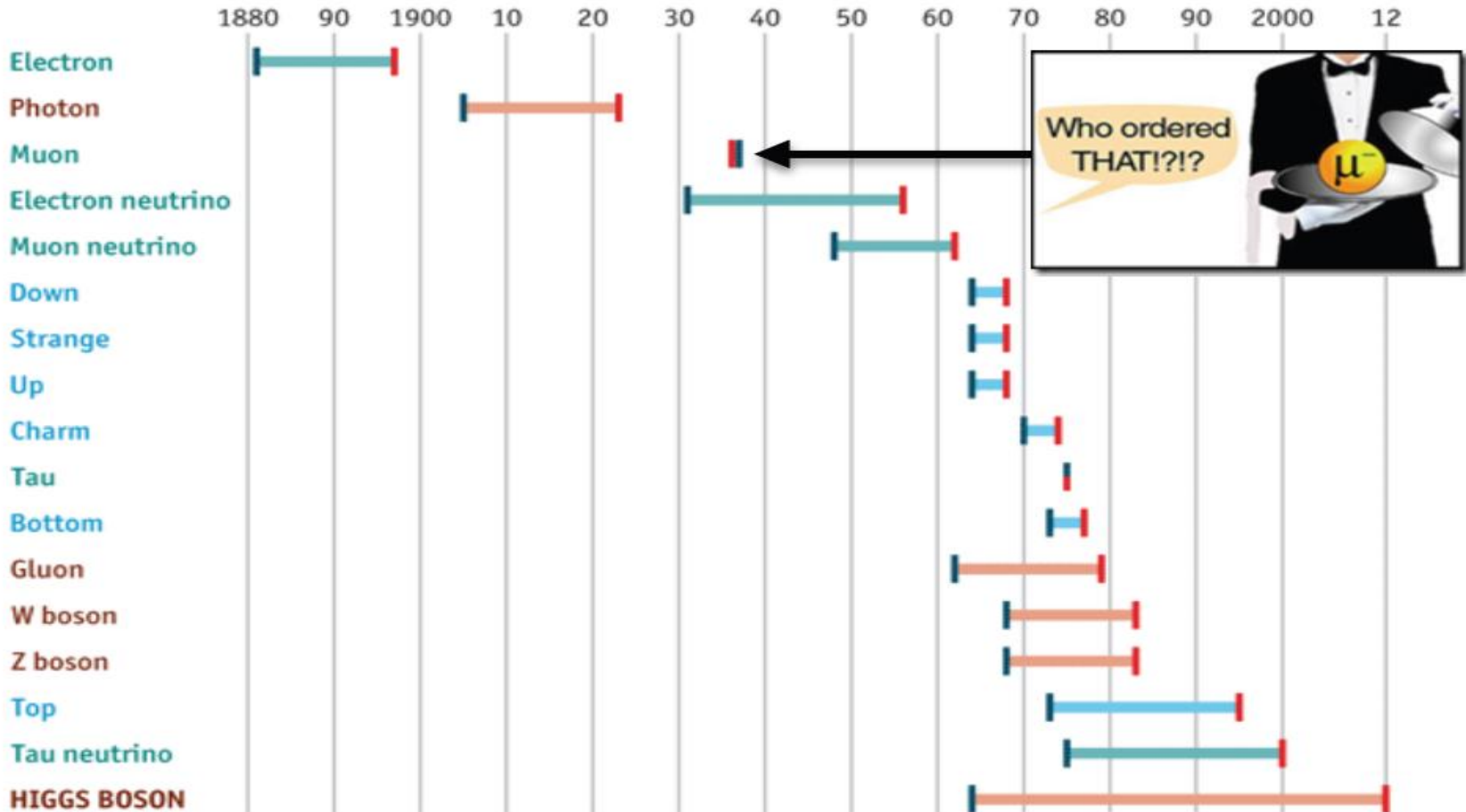
Uncharted discoveries?

The Standard Model of particle physics

Years from concept to discovery

Leptons
Bosons
Quarks

Theorised/explained
Discovered



Source: *The Economist*

Matter particles

2.4M u up B: 2/3 L: 0 S: 1/2	1.27G c charm B: 2/3 L: 0 S: 1/2	171.2G t top B: 2/3 L: 0 S: 1/2
4.8M d down B: 1/3 L: 0 S: 1/2	104M s strange B: 1/3 L: 0 S: 1/2	4.2G b bottom B: 1/3 L: 0 S: 1/2
0.511M e electron B: 0 L: -1 S: 1/2	105.7M μ muon B: 0 L: -1 S: 1/2	1.777G τ tau B: 0 L: -1 S: 1/2
< 2.2 ν_e e-neutrino B: 0 L: 1 S: 1/2	< 0.17M ν_μ μ-neutrino B: 0 L: 1 S: 1/2	< 15.5M ν_τ τ-neutrino B: 0 L: 1 S: 1/2

- **u, d, e**
ordinary matter

- **s, μ**
cosmic rays

- **c, b, t, τ**
accelerators

Matter particles

1900

2000



1897

Electron

J.J. Thomson, *Philosophical magazine* **44**:293

1969

up, down, strange quarks

E.D. Bloom *et al.* *Physical Review Letters* **23** (16): 930

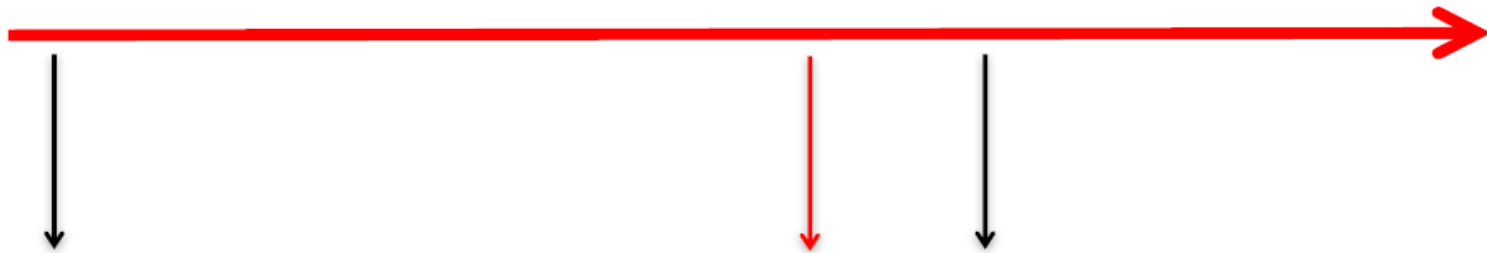
J. M. Breidenbach *et al.* *Physical Review Letters* **23** (16): 235

u,d proposed 1960s, discovered ~1968
e discovered 1897

Matter particles

1900

2000



1956

Electron neutrino

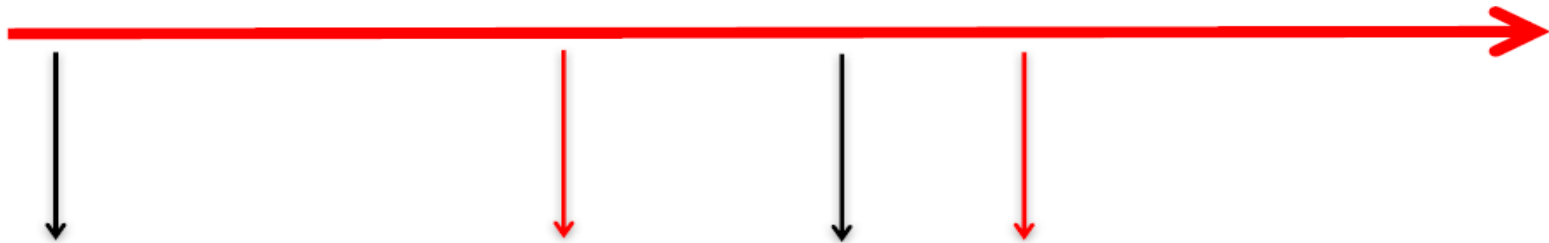
F. Reines, C.L. Cowan, *Nature* **178** (4531): 446

Radioactive decay (inferred 1930s, seen 1956)

Matter particles

1900

2000



1937
Muon
S.H. Neddermeyer, C.D. Anderson,
Physical Review **51** (10): 884

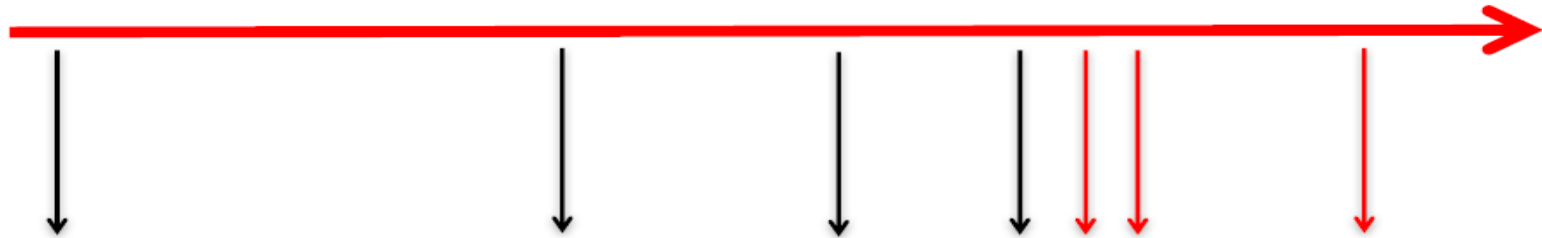
1969
up, down, strange quarks
E.D. Bloom *et al.* *Physical Review Letters* **23** (16): 930
J. M. Breidenbach *et al.* *Physical Review Letters* **23** (16): 235

Cosmic ray experiments (1930s, 1940s)

Matter particles

1900

2000



1974

Charm quarks

J.J. Aubert *et al.* *Physical Review Letters* **33** (23): 1404

J.-E. Augustin *et al.* *Physical Review Letters* **33** (23): 1406

1977

Bottom quarks

S.W. Herb *et al.* *Physical Review Letters* **39** (5): 252.

1995

Top quarks

F. Abe *et al.* ([CDF collaboration](#)) *Physical Review Letters* **74** (14): 2626–2631.

S. Arabuchi *et al.* ([D0 collaboration](#)) *Physical Review Letters* **74** (14): 2632–2637.

Collider experiments (1960s -)

And ... antimatter

Einstein's equation of motion*: $E^2 = p^2 c^2 + m^2 c^4$

Two energy solutions for the same mass;

- Matter
- Antimatter

Every fermion has an antimatter version.

Same mass, opposite charge


eg. antiquark \bar{q} , antimuon μ^+ , antineutrino $\bar{\nu}$

*(and others, more famously Dirac)

Forces in particle physics

Electromagnetic

$0 \text{ eV}/c^2$
0
1
photon




U(1)

Strong (QCD)

8 x

$0 \text{ eV}/c^2$
0
1
gluon




SU(3)


Weak

2 x

$80.4 \text{ GeV}/c^2$
 ± 1
1
W boson



, $91.2 \text{ GeV}/c^2$
0
1
Z boson



SU(2)

Note:
No gravity!!

Forces in particle physics

EM force

Electric charge (1)

Massless photon

Coupling g

Weak force

Weak charge (2)

Massive W^\pm, Z

Coupling g_W

Strong force

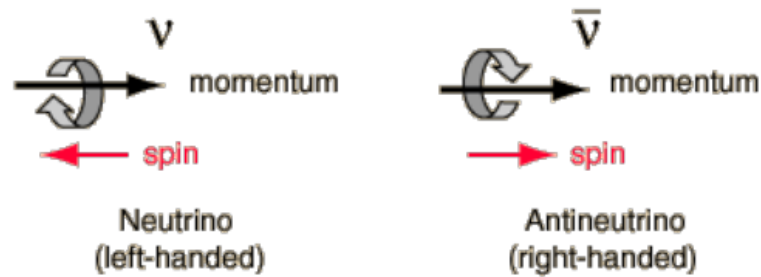
Colour charge (3)

8 massless gluons

Coupling g_s

Value unknown/
not predicted

Forces in particle physics



EM force

Abelian

Only charged particles couple

Weak force

Non-abelian

Only left handed particles couple

quark mixing (3 generations, CP)

Neutrino mixing (3 generations, CP)

Strong force

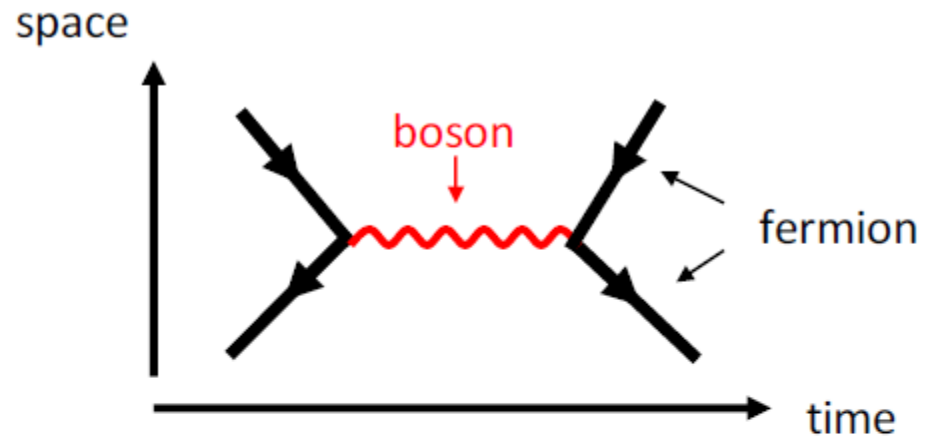
Non-abelian

Only quarks couple

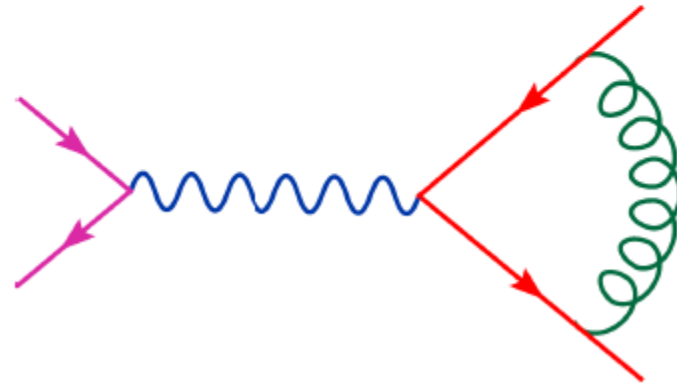
Value unknown/
not predicted

Aside: Feynman diagrams

“tree” level
Lowest order

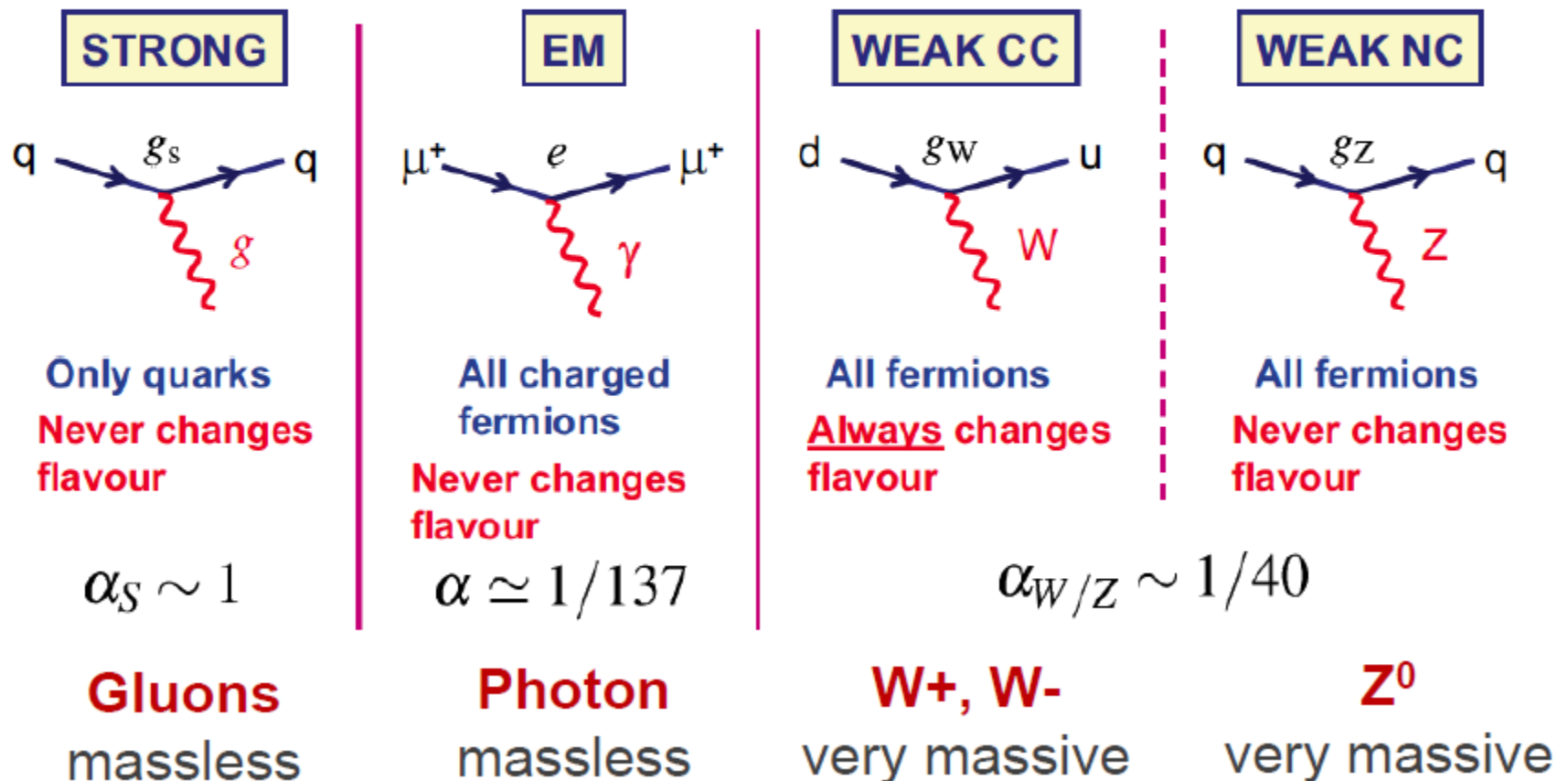


Higher orders possible
Loops



Interactions

The interaction of gauge bosons with fermions is described by the Standard Model

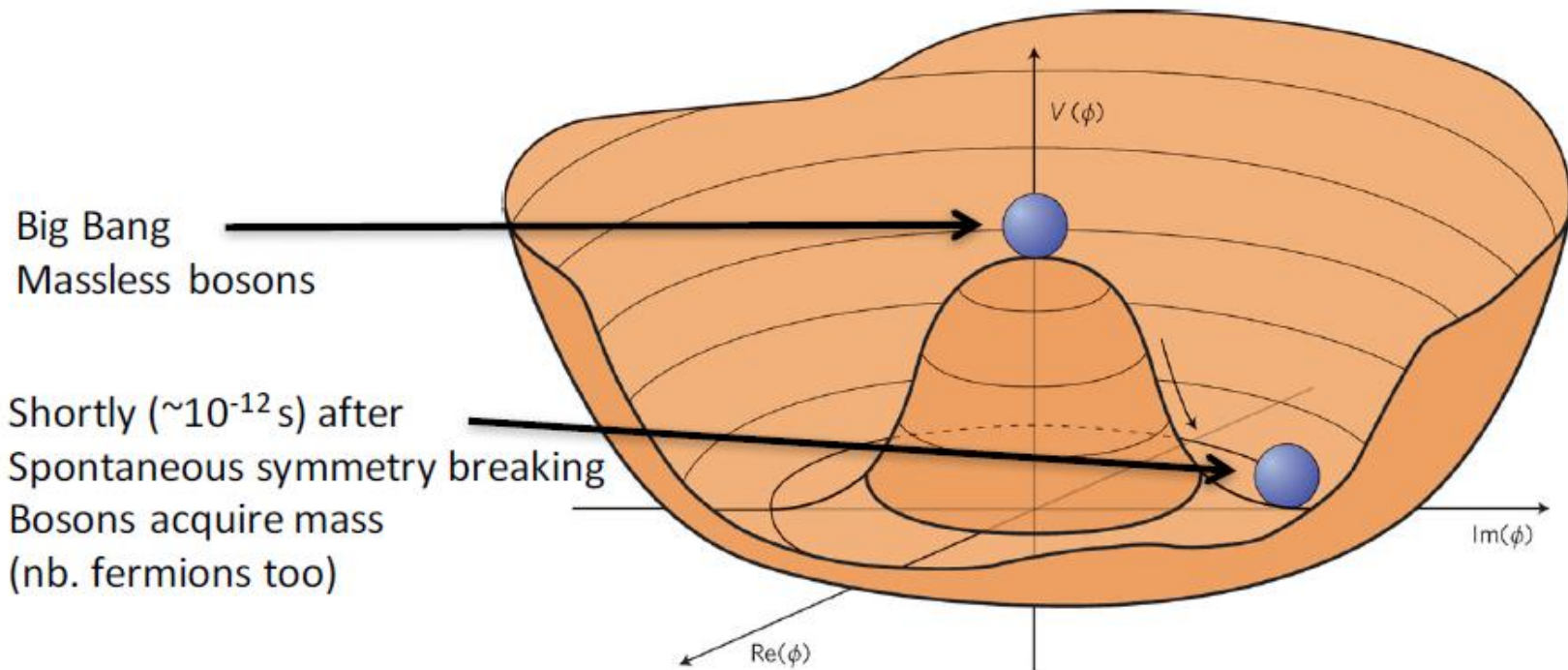


The Higgs boson

Introduce Higgs field:

Complex doublet (but 1d case shown here to get idea)

$$V(\phi) = -0.5\mu^2|\phi|^2 + \lambda|\phi|^4$$



The Higgs boson

Introduce Higgs field :

Couples to particles to give mass (amount \sim coupling strength)

Complex doublet has **4 free parameters**

3 absorbed into W^+ , W^- , Z boson mass

W^+ , W^- , Z , γ admixtures of original weak, em massless bosons.

1 manifested as a massive Higgs boson (m_H)

Connection between weak and electromagnetic forces

(note: Higgs field gives mass to fermions by a different mechanism)

Yukawa coupling; yet to be fully tested.

- No deep explanation; motivated by simplicity.

The Higgs boson

Introduce Higgs field :

After symmetry breaking, Higgs sector properties are:

- spinless Higgs boson (m_H)
- vacuum expectation value (mean field value) (v)

Consequences:

Weak and electromagnetic forces connected

Massive Z is mixture of massless em + weak bosons



Relates M_W , M_Z and weak, electromagnetic couplings:

$$\tan \theta_W = g_W / g$$

$$M_W = M_Z \cos \theta_W$$



QED and QCD

Quantum Electrodynamics: QED

Electric charge  Atoms  Molecules

Interaction of electric charges and photons

Quantum Chromodynamics: QCD

Colour charge  Baryons  Nucleus

Interaction of colour charges and gluons

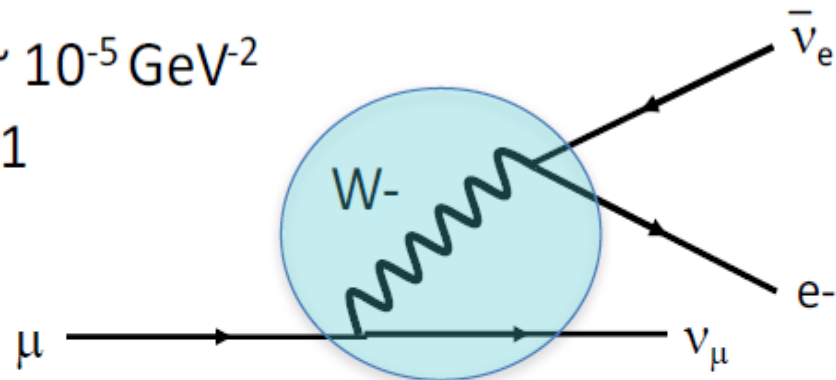
Different forces, but **similar** (mathematical) structure/behaviour

Weak force vs QED and QCD

Muon decay:

Strength of weak force $\sim G_F \sim 10^{-5} \text{ GeV}^{-2}$

cf. strength of em force ~ 0.01



W boson **massive**

Factor involved in boson exchange $\sim 1/(E^2 + M^2)$ (hence units)

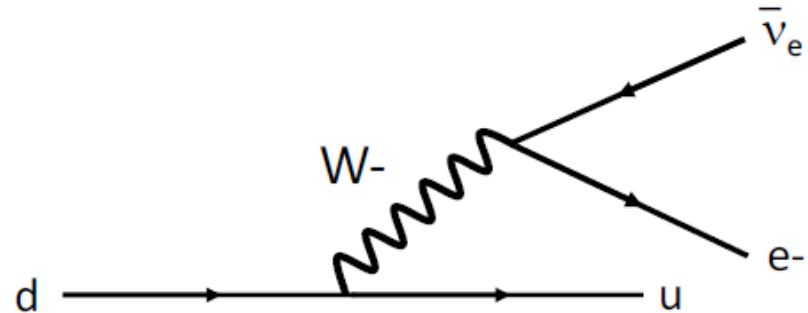
Strength of weak force = em force if $M \sim 30 \text{ GeV}$ ($M_W \sim 80 \text{ GeV}$)

Weak force vs QED and QCD

W couples to:

Upper and lower members
of a fermion generation.

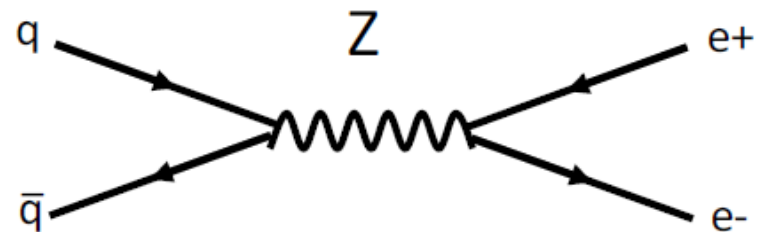
L- (R-) handed (anti)particles



Z couples to:

Matter and antimatter
versions of a fermion.

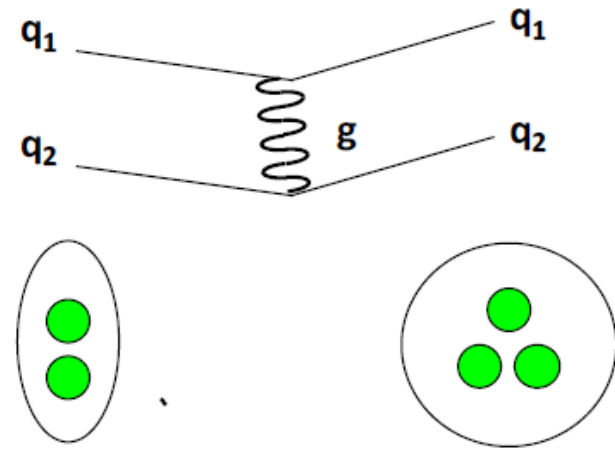
Complicated mix of L-, R-
particles.



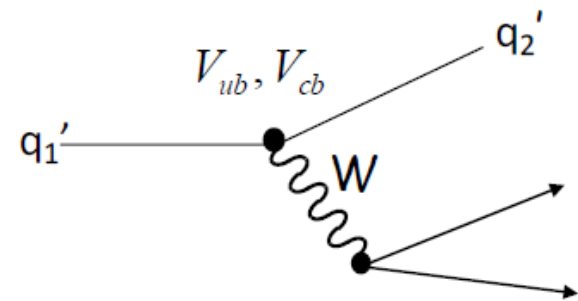
“vector, axial couplings”; Higgs mechanism.

Weak vs. mass quark eigenstates

Mass eigenstates of quarks form hadrons



W couples to weak quark eigenstates q'
 q' admixture of q and vice versa



Quark mixing

$$\begin{pmatrix} d_W \\ s_W \\ b_W \end{pmatrix} = \begin{pmatrix} & & \\ & & \\ & & \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Weak, mass eigenstates related by mixing matrix in SM (3x3 matrix)

Mixing matrix is unitary (inverse = complex conjugate)

CKM matrix

CKM matrix (1973 – before charm! Predicted 3rd generation)

Elements describe every weak quark transition

SM does not predict existence of or values for matrix elements (couplings of W to quarks).

Input by experimental data

$$V_{CKM} = \begin{matrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{matrix}$$

CKM matrix

- Need 3 generations of quarks to introduce CP violation into theory

$$\begin{pmatrix} d_W \\ s_W \\ b_W \end{pmatrix} = \begin{pmatrix} & & \\ & & \\ & & \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Mixing matrix is 3x3.

Unitarity constraints \Rightarrow 4 independent parameters

3 angles quantify mixing between (1,3) (2,3) (1,2) generations, **1 complex phase** (mechanism for introducing CP)

CP violation

C = charge operator

P = parity operator

CP operation changes particle q to antiparticle \bar{q} (and vice versa)

CP **violation** if $q \rightarrow q'$ rate different to $q' \rightarrow q$ ie. $V_{qq'} \neq V_{qq'}^*$

CP violation observed in weak decays.

Note:

- **SM does not predict** CP violation.
- **SM does not explain** CP violation.
- CP violation **must be added** to SM.

Aside: neutrino CP violation, mixing

- Similar framework adopted for neutrinos (PMNS matrix).
Weak (ν_e, ν_μ, ν_τ) related to mass eigenstates (ν_1 etc):

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} & & \\ & & \\ & & \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

3 angles quantify mixing between (1,3) (2,3) (1,2) generations, **1 complex phase** (mechanism for introducing CP)

Note: parameters investigated in dedicated neutrino experiments

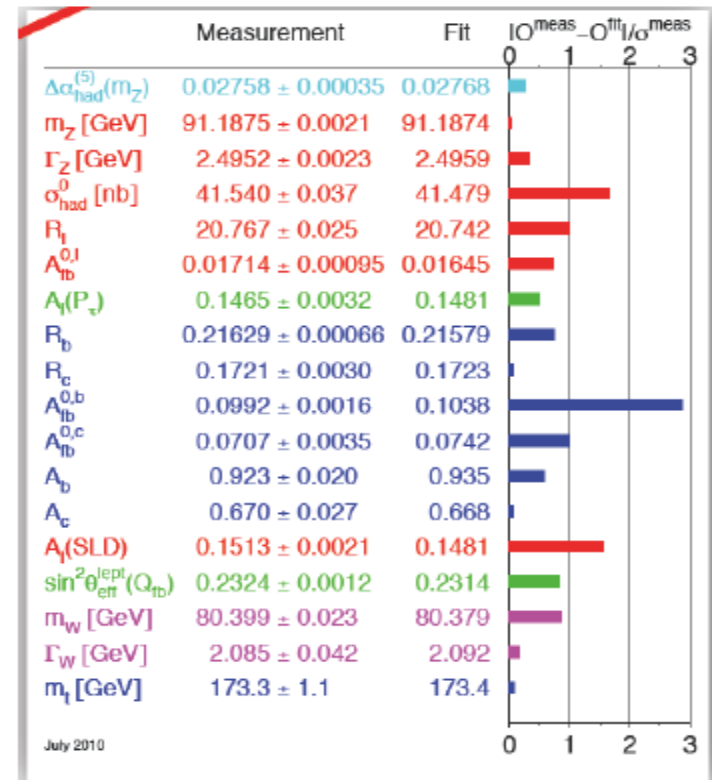
Standard Model confirmed by the data

	I	II	III	
mass →	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²	0
charge →	2/3	2/3	2/3	0
spin →	1/2	1/2	1/2	1
name →	u up	c charm	t top	γ photon
	4.8 MeV/c ²	104 MeV/c ²	4.2 GeV/c ²	0
	-1/3	-1/3	-1/3	0
	1/2	1/2	1/2	1
Quarks	d down	s strange	b bottom	g gluon
	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²	91.2 GeV/c ²
	0	0	0	0
	1/2	1/2	1/2	1
Leptons	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z⁰ Z boson
	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²	80.4 GeV/c ²
	-1	-1	-1	±1
	1/2	1/2	1/2	1
	e electron	μ muon	τ tau	W[±] W boson

Gauge bosons

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi} \not{D} \psi + h.c.$$

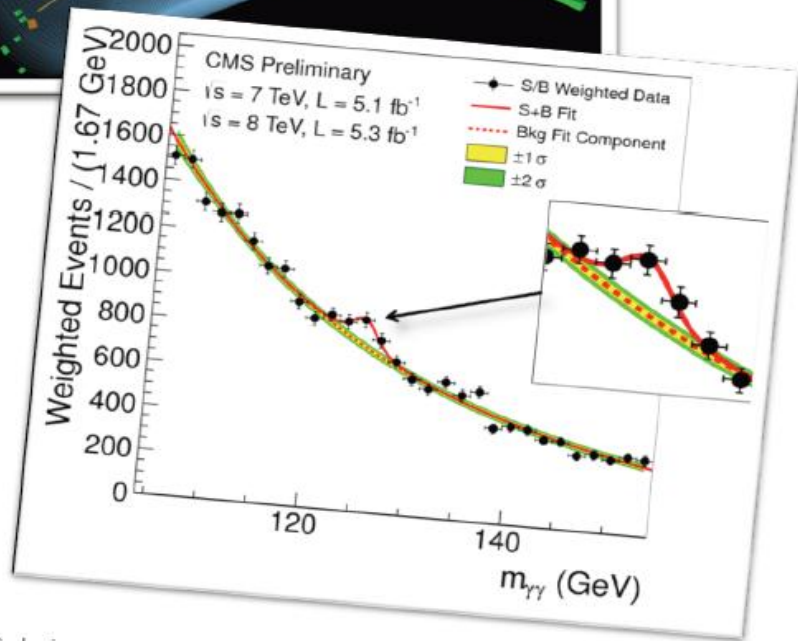
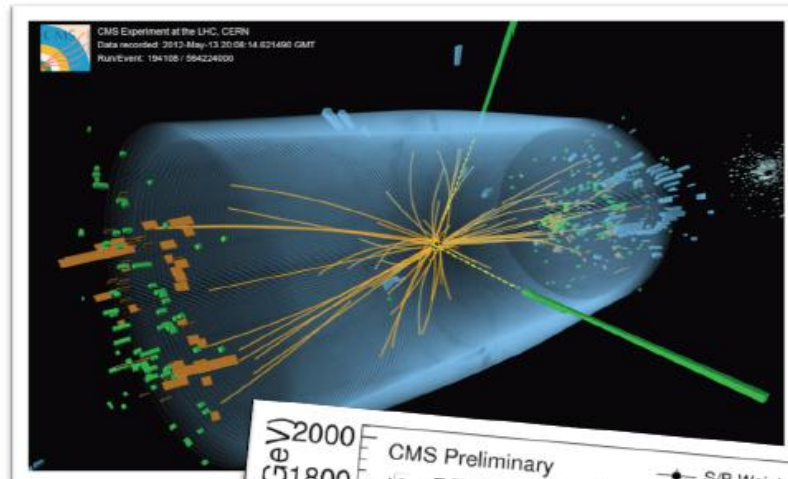
STANDARD MODEL OF ELEMENTARY PARTICLES



Confirmed at sub 1% level!

Experiment = probing theories with data

$$\begin{aligned}
 & -\frac{1}{2}\partial_\nu g_\mu^\nu \partial_\rho g_\mu^\rho - g_s f^{abc} \partial_\mu g_\nu^\mu \partial_\rho g_\nu^\rho - \frac{1}{2}g_s^2 f^{abc} f^{def} g_\mu^\mu g_\nu^\nu g_\rho^\rho g_\sigma^\sigma + \\
 & \frac{1}{2}g_s^2 (q_i^\mu \gamma^\mu q_j^\nu) g_\mu^\mu + G^a \partial^\mu G^a + g_s f^{abc} \partial_\mu G^a G^b G^c - \partial_\mu W_\nu^+ \partial_\rho W_\mu^- - \\
 & M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2}M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\nu H - \\
 & \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2}M^2 \phi^0 \phi^0 - \beta_h \frac{[2M^2]}{2} + \\
 & \frac{2M^2}{\Lambda} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) + \frac{2M^2}{\Lambda} \alpha_h - ig_{c_w} [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
 & W_\mu^- W_\nu^+) - Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
 & W_\nu^- \partial_\nu W_\mu^+) - ig_{s_w} [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\mu^- W_\nu^+) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
 & W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+) - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\
 & \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^- W_\nu^+ + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^-) + g^2 s_w^2 (W_\mu^+ W_\nu^-) + \\
 & g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
 & W_\nu^- W_\mu^+) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^2 + H\phi^0 \phi^0 + 2(H\phi^+ \phi^-) - \\
 & \frac{1}{2}g^2 \alpha_h (H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2) - \\
 & gMW_\mu^+ W_\mu^- H - \frac{1}{2}g\frac{M}{c_w} Z_\mu^0 Z_\nu^0 H - \frac{1}{2}igW_\mu^+ (\partial_\nu \phi^0 \phi^0 - \phi^- \partial_\nu \phi^0) - \\
 & W_\mu^- (\phi^0 \partial_\nu \phi^+ - \phi^+ \partial_\nu \phi^0) + \frac{1}{2}ig[W_\mu^+ (H\partial_\nu \phi^0 - \phi^- \partial_\nu H) - W_\mu^- (H\partial_\nu \phi^+ - \\
 & \phi^+ \partial_\nu H)] + \frac{1}{2}g\frac{1}{c_w} (Z_\mu^0 (H\partial_\nu \phi^0 - \phi^0 \partial_\nu H) - ig\frac{2s_w}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \\
 & ig_{s_w} MA_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig\frac{1-2s_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) - \\
 & ig_{s_w} A_\mu (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) - \frac{1}{2}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
 & \frac{1}{4}g^2 \frac{1}{c_w} Z_\mu^0 Z_\nu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)\phi^+ \phi^-] - \frac{1}{2}g^2 \frac{2s_w}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) - \frac{1}{2}ig^2 \frac{2s_w}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 \frac{2s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\
 & g^2 s_w^2 A_\mu A_\nu \phi^+ \phi^- - e^2 (\gamma \partial + m_\nu^2) e^\lambda - \rho^2 \gamma \partial \nu^\lambda - u_\nu^2 (\gamma \partial + m_\nu^2) u_\nu^\lambda + \\
 & d_\nu^2 (\gamma \partial + m_\nu^2) d_\nu^\lambda + ig_{s_w} A_\mu [-(e^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(u_\nu^2 \gamma^\mu u_\nu^\lambda) - \frac{1}{3}(d_\nu^2 \gamma^\mu d_\nu^\lambda)] + \\
 & \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (u_\nu^2 \gamma^\mu (\frac{2}{3}s_w^2 - \\
 & 1 - \gamma^5) u_\nu^\lambda) + (d_\nu^2 \gamma^\mu (1 - \frac{2}{3}s_w^2 - \gamma^5) d_\nu^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + \\
 & (\bar{d}_\nu^2 \gamma^\mu (1 + \gamma^5) C_{\lambda\nu} d_\nu^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(e^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + (\bar{d}_\nu^2 \gamma^\mu (1 + \gamma^5) \nu^\lambda) - \\
 & \frac{ig}{2\sqrt{2}} M d_\nu^\lambda] + \frac{ig}{2\sqrt{2}} M d_\nu^\lambda [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (e^\lambda (1 + \gamma^5) \nu^\lambda)] - \\
 & \frac{ig}{2\sqrt{2}} M [H (e^\lambda e^\lambda) + i\phi^0 (e^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2\sqrt{2}} M \phi^+ [-m_\nu^2 (\bar{u}_\nu^2 C_{\lambda\nu} (1 - \gamma^5) d_\nu^\lambda) + \\
 & m_\nu^2 (\bar{d}_\nu^2 C_{\lambda\nu} (1 + \gamma^5) u_\nu^\lambda) - \frac{ig}{2\sqrt{2}} M d_\nu^\lambda (\bar{d}_\nu^2 C_{\lambda\nu} (1 + \gamma^5) u_\nu^\lambda) - \\
 & \frac{ig}{2\sqrt{2}} M d_\nu^\lambda H (\bar{u}_\nu^2 u_\nu^\lambda) - \frac{ig}{2\sqrt{2}} M H (\bar{d}_\nu^2 d_\nu^\lambda) + \frac{ig}{2\sqrt{2}} M \phi^0 (\bar{u}_\nu^2 \gamma^5 u_\nu^\lambda) - \\
 & \frac{ig}{2\sqrt{2}} M \phi^0 (\bar{d}_\nu^2 \gamma^5 d_\nu^\lambda) + \bar{X}^+ (\partial^\mu - M^2) X^+ + \bar{X}^- (\partial^\mu - M^2) X^- + \bar{X}^0 (\partial^\mu - \\
 & \frac{M^2}{c_w}) X^0 + \bar{Y} \partial^\mu Y + ig_{c_w} W_\mu^+ (\partial_\nu \bar{X}^0 X^- - \partial_\nu \bar{X}^+ X^0) + ig_{s_w} W_\mu^+ (\partial_\nu \bar{Y} X^- - \\
 & \partial_\nu \bar{X}^+ Y) + ig_{c_w} W_\mu^- (\partial_\nu \bar{X}^- X^0 - \partial_\nu \bar{X}^0 X^+) + ig_{s_w} W_\mu^- (\partial_\nu \bar{X}^- Y - \\
 & \partial_\nu \bar{Y} X^+) + ig_{c_w} Z_\mu^0 (\partial_\nu \bar{X}^+ X^- - \partial_\nu \bar{X}^- X^+) + ig_{s_w} A_\mu (\partial_\nu \bar{X}^+ X^- + \\
 & \partial_\nu \bar{X}^- X^+) - \frac{1}{2}igM[\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\
 & \frac{1-2c_w^2}{2c_w} igM[\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM[\bar{X}^0 X^- \phi^+ - \bar{X}^+ X^+ \phi^-] + \\
 & igMs_w[\bar{X}^0 X^- \phi^+ - \bar{X}^+ X^+ \phi^-] + \frac{1}{2}igM[\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
 \end{aligned}$$

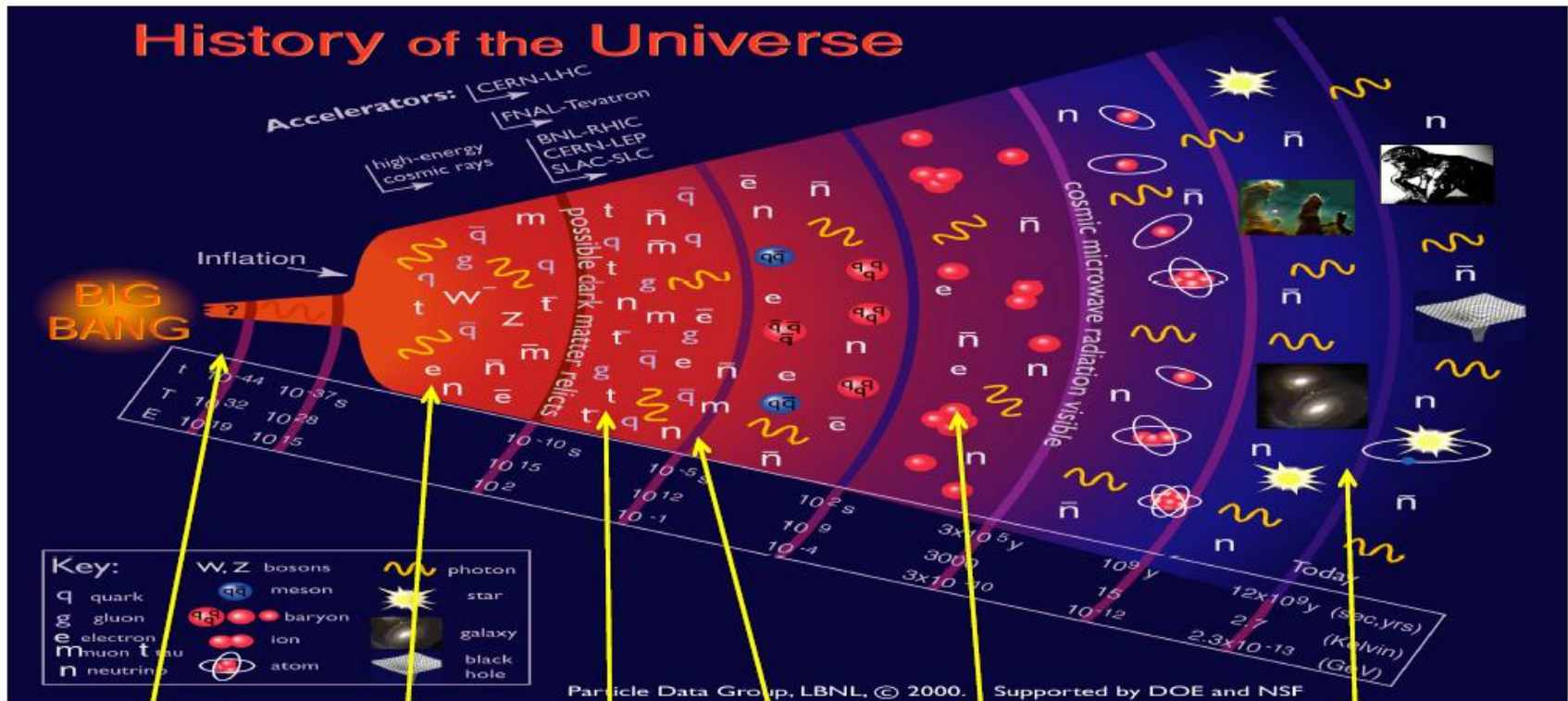


- Delmastro

(experimental) LHC physics

Accelerators for high energy physics experiments

History of the Universe



Cosmology

Cosmic rays

LHC

Quark/gluon plasma

Nuclear physics

Astrophysics

Why accelerators?

Accelerators are instruments that increase the energy of particles to study smaller and smaller structures and create heavy short-lived objects in collision with

$$E = m \cdot c^2$$

Wavelength of probe radiation needs to be smaller than object to be resolved

$$\lambda = \frac{h}{p} = \frac{h \cdot c}{E}$$

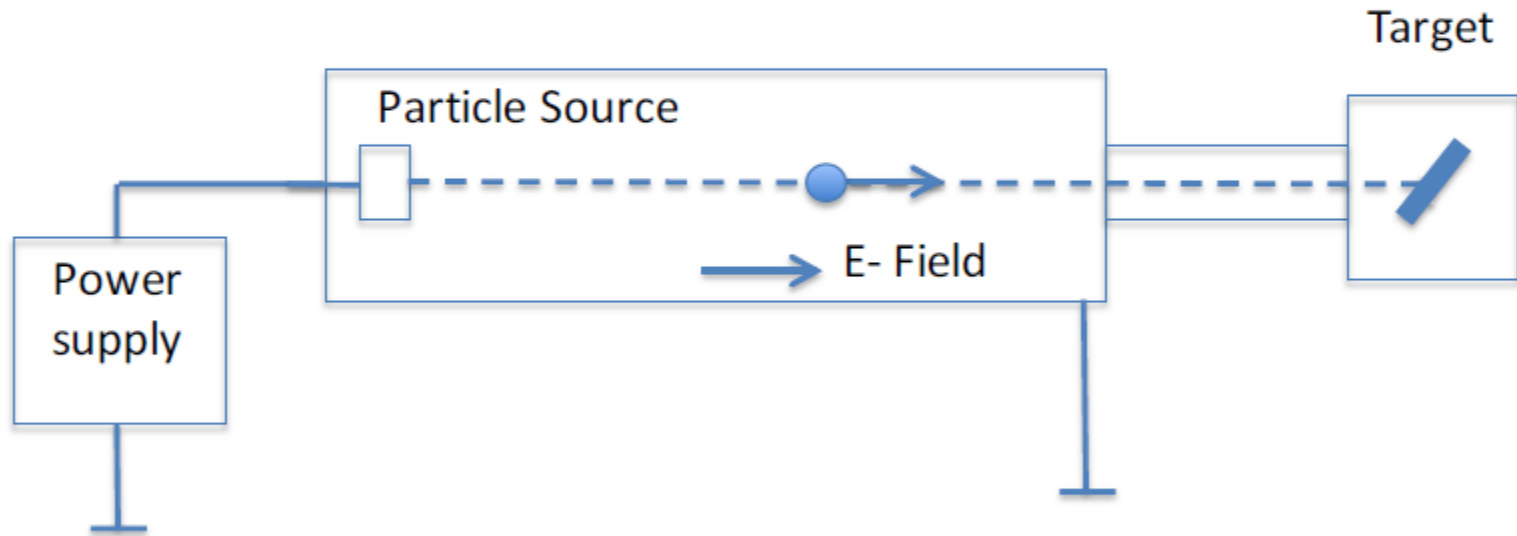


Object	size	Radiation energy
Atom	10^{-10} m	0.00001 GeV
Nucleus	10^{-14} m	0.01 GeV
Nucleon	10^{-15} m	0.1 GeV
Quarks	-	> 1 GeV

The basic accelerator

Electrostatic accelerator:

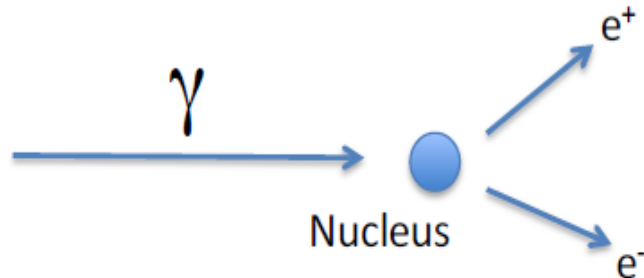
Charged particles go through the accelerating voltage gap **once** and then hit the target.



Limited by the maximum reachable voltage: ~ 10 MV

Why collisions

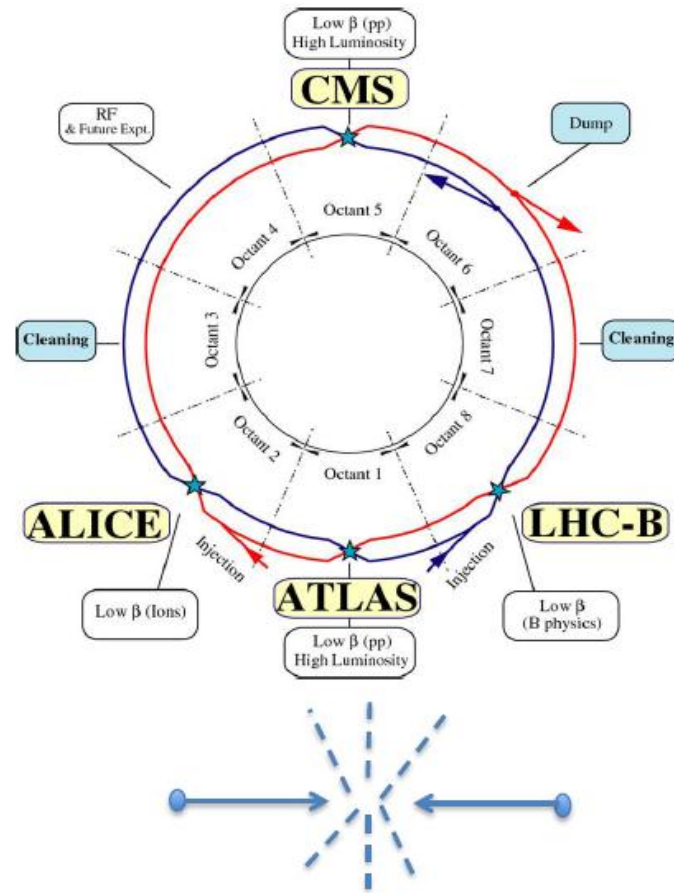
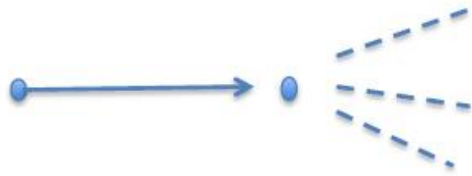
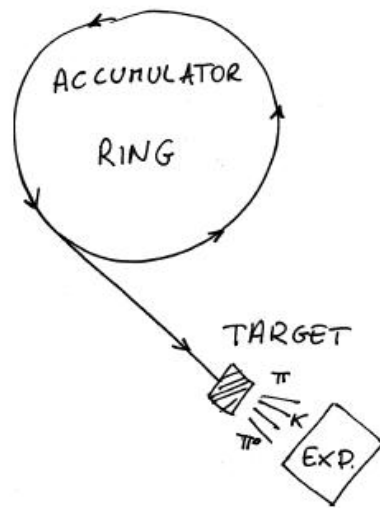
- Conservation laws: e.g. momentum and energy conservation



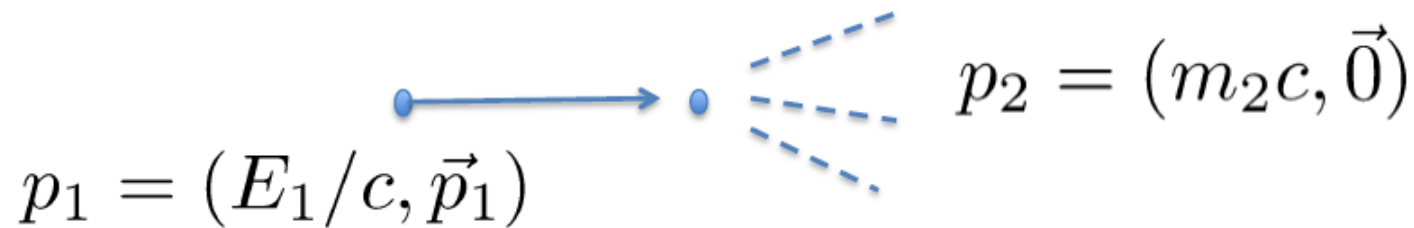
Photon into e^+, e^- only in proximity of nucleus. Nucleus takes part of momentum (and part of available energy...)

- Center-of-mass Frame and Center-of-mass Energy (E_{CM})
 - Center-of-mass frame defined where:
$$\sum \vec{p}_i = \vec{0}$$
 - The energy available for creation of particles corresponds to E_{CM}

Fixed target vs Colliders



E_{CM} in Fixed Target Experiment




$$p_{tot} = (E_1/c + m_2c, \vec{p}_1)$$

$$E_{CM}^2 = (m_1^2 + m_2^2)c^4 + 2E_1m_2c^2$$

$$E_{CM} \propto \sqrt{E_1}$$

E_{CM} in Collider Experiment

Laboratory Frame = CM Frame


$$p_1 = (E_1/c, \vec{p}_1) \quad p_2 = (E_2/c, -\vec{p}_1)$$

$$E_{CM} = E_1 + E_2$$

➔ Collider more energy efficient;
But also more complex: two beams to be accelerated and to be brought into collision

Acceleration

Lorentz force law

$$\mathbf{F} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

Electric field Velocity Magnetic field

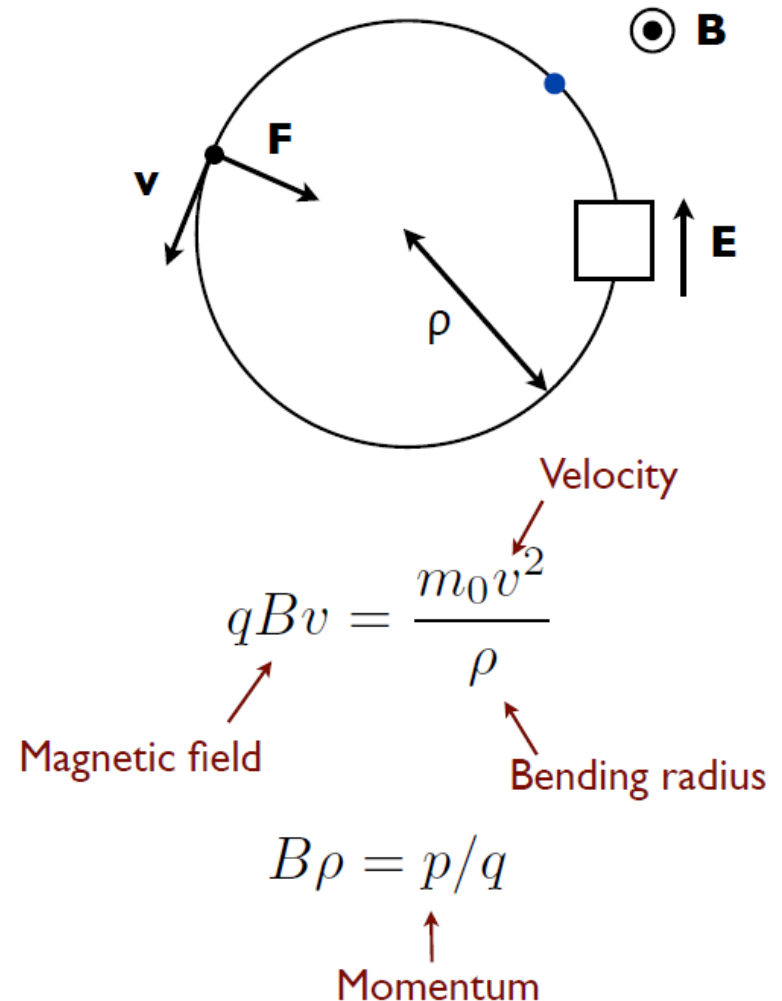
Energy change

$$\Delta E = \int_{\mathbf{r}_1}^{\mathbf{r}_2} \mathbf{F} \cdot d\mathbf{r}$$

- Electric field (either static or more commonly, time varying) to accelerate, or more appropriately, increase energy of beam
- Magnetic part of Lorentz force used to guide and focus
 - Dipole magnets: to bend
 - Quadrupole: to focus or defocus

Synchrotron

- Workhorse of modern particle physics
 - Huge legacy of discovery
 - Increase energy whilst synchronously increasing bending magnet strength
 - Stable storage of high beam current/power
- Magnetic field proportional to momentum



Storage ring Colliders

Make use of all the particles' energy. 2-beam synchrotrons.

The first one: AdA (Frascati), 1961-64, e^+, e^- , 250 MeV, 3m circumference

Many examples to come at DESY, SLAC, KEK, Fermilab with the Tevatron (980 GeV), BNL with RHIC

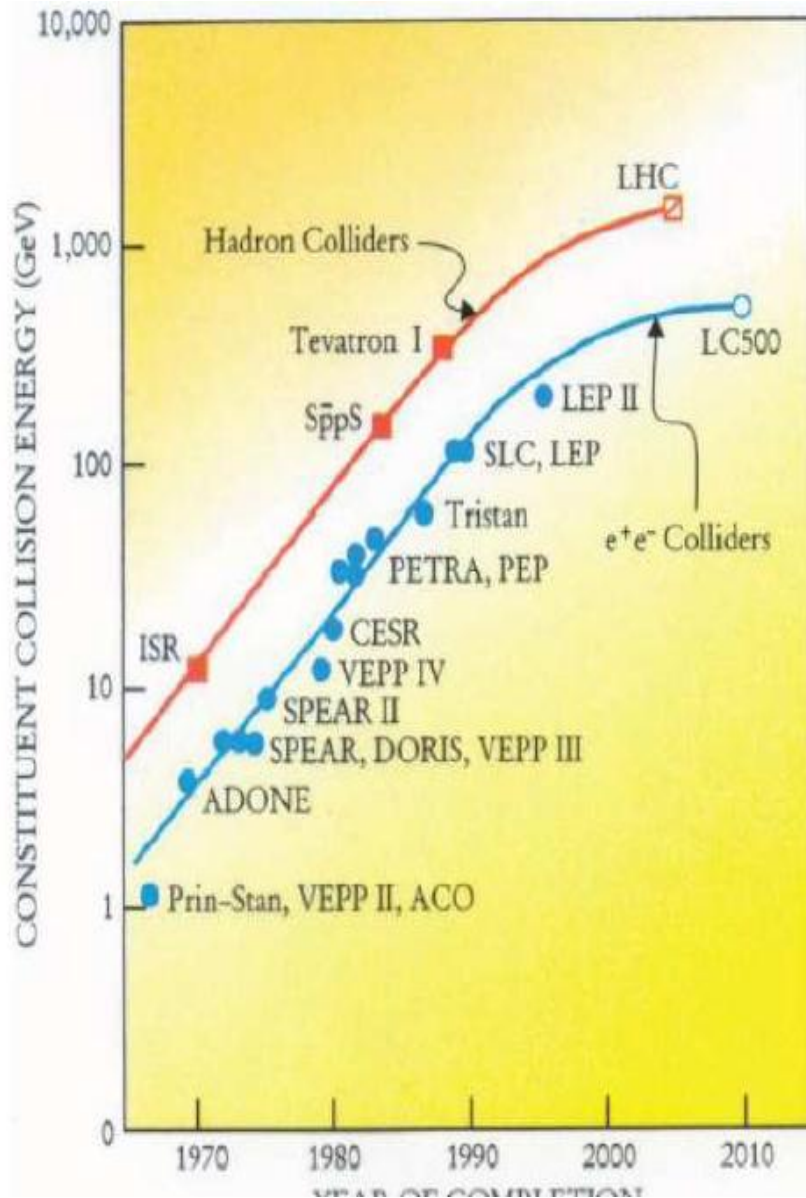
1971-1984: ISR (CERN), p^+, p^+ , 31.5 GeV, 948 m circumference

1981-1991: SPS running as $Sp\bar{p}S$, p^+ , p^- , 270 – 315 GeV, 6.9 km circumference; discovery of W and Z Bosons

1989-2000: LEP highest energy electron synchrotron, e^+, e^- , 104 GeV, 27 km circumference; three generations of quarks, gluons and leptons

2008 - : LHC highest energy proton synchrotron, p^+, p^+ , heavy ions, 6.5 TeV (2.76 TeV per nucleon for $^{208}\text{Pb}^{82+}$); Discovery of Higgs

Energy frontier



- The interplay between electron and hadron machines has a long and fruitful tradition
 - J/ψ at SPEAR (e^+e^-) and AGS (proton fixed target)
 - Υ discovery at E288 (p fixed target), precision B studies at the e^+e^- B factories
 - ...
 - top quark at LEP and Tevatron
- To be continued in the form of LHC and ILC

Complementarity between pp and ee machines



- Proton-(Anti-)Proton Colliders
 - Higher energy reach (limited by magnets)
 - Composite particles: unknown and different colliding constituents, energies in each collision
 - Confusing final states
- Discovery machines (W , Z , t)
- In some cases: precision measurements possible (W mass at the Tevatron)



- Electron-Positron-Colliders
 - Energy reach limited by RF
 - Point like particles, exactly defined initial system, quantum numbers, energy, spin polarisation possible
 - Hadronic final states with clear signatures
- Precision machines
- Discovery potential, but not at the energy frontier

Luminosity

- What luminosity is required for measurement?
- Need some knowledge of x-section
- Simple relationship between number of particles, frequency of collision and beam sizes

$$\mathcal{L} = f \frac{N_1 N_2}{4\pi\sigma_x\sigma_y}$$

Luminosity [$s^{-1} m^{-2}$] Bunch populations
 Frequency of collisions [Hz] Beam r.m.s. sizes [m]

$$\sigma = \sqrt{\epsilon\beta}$$

Emittance [m] Beta function [m]

$$\mathcal{L} = f \frac{N_1 N_2}{4\pi\sqrt{\epsilon_x\beta_x^*\epsilon_y\beta_y^*}}$$

Luminosity frontier

- Need corresponding rise in luminosity (beam intensity)

$$N = \sigma L = \sigma \int \mathcal{L} dt$$

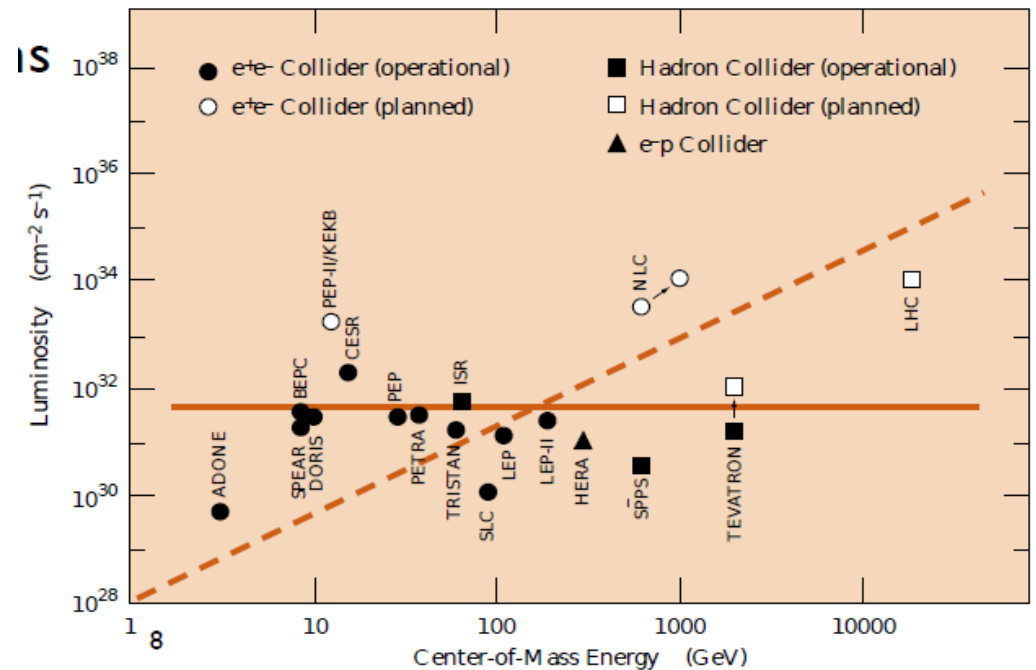
Number of events
Instantaneous luminosity

↓
↓

↑
↑

Cross section
Integrated luminosity

- High luminosity brings all the challenges for the detectors:
 - High event rates
 - Pile up
 - Beam –beam interactions
 - Beamstrahlung



Designing a machine

- Particle species
 - Electron/positrons
 - Protons/antiprotons
 - Muons/antimuons
- Beam energy
- Spin
- Luminosity
- How do you produce antiparticles?
- Ones produced how ones keep them (muon collider)?
- Ones collided what ones does with spent beams?
- Accelerator and detector protection

Accelerator is much more than just....

- Particle production
- Damping, cooling or preparation
- Injection and extraction
- Acceleration
- Collimation (betatron, energy etc.)
- Diagnostics and controls
- Machine (and detector protection)
- Beam delivery and luminosity production
- Technology spin off
 - Lower energy machines, medical applications, applied physics, materials,

CERN laboratory (founded in 1954)

Mont Blanc

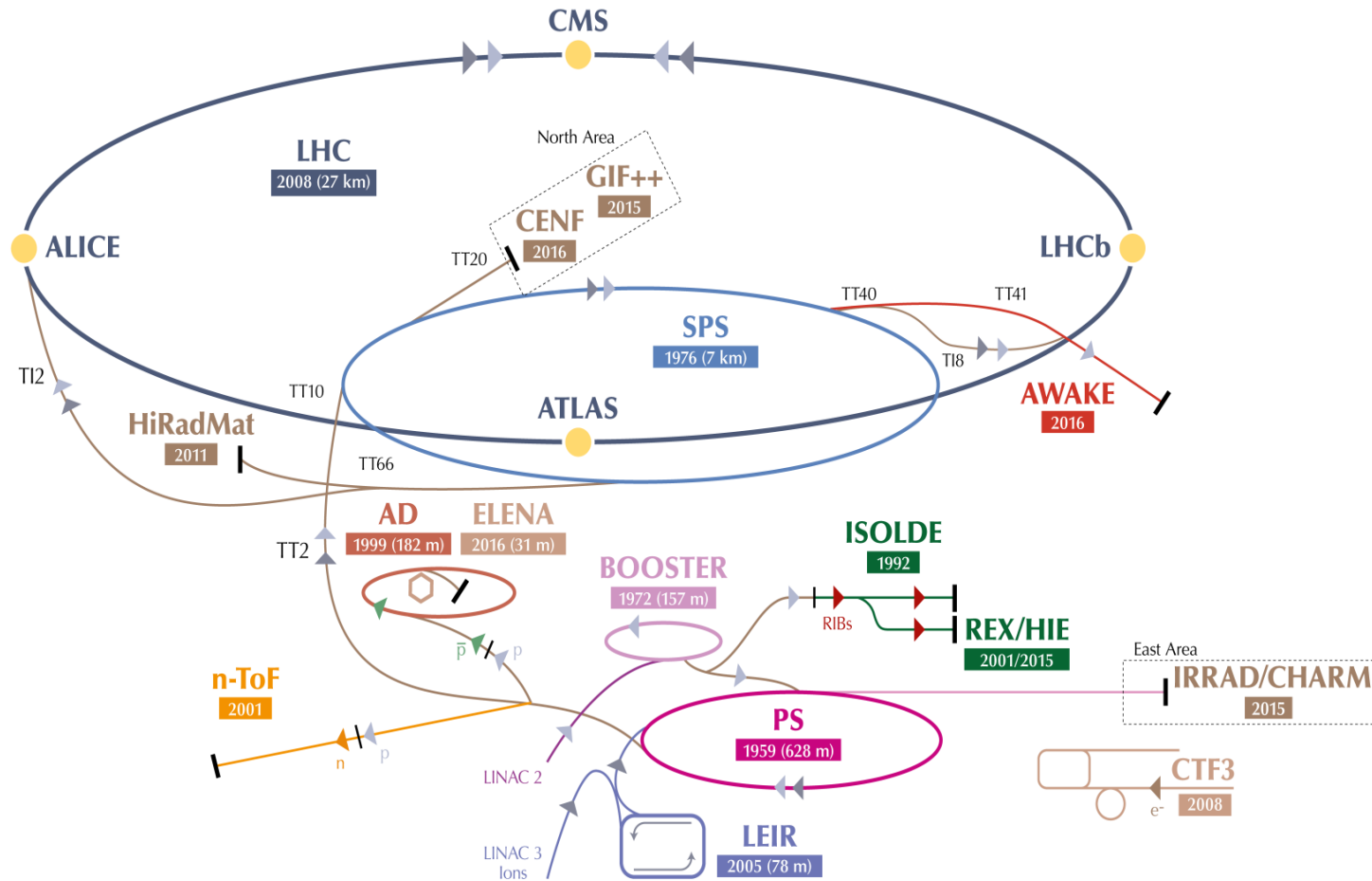
20 krajów członkowskich (Polska od 1992)
40 krajów stowarzyszonych
2 300 zatrudnionych osób
+ 10 000 naukowców

Jezioro Genewskie



CERN accelerator complex

CERN 2017



LHC pp and ions

7 TeV/c –up to
now 4 TeV/c

26.8 km
Circumference

The confusion with 7 TeV: energy of one
proton or two protons ? ...watch out

Switzerland
Lake Geneva

LHC Accelerator
(100 m down)

CMS, TOTEM

CERN-
Prevezin

ALICE

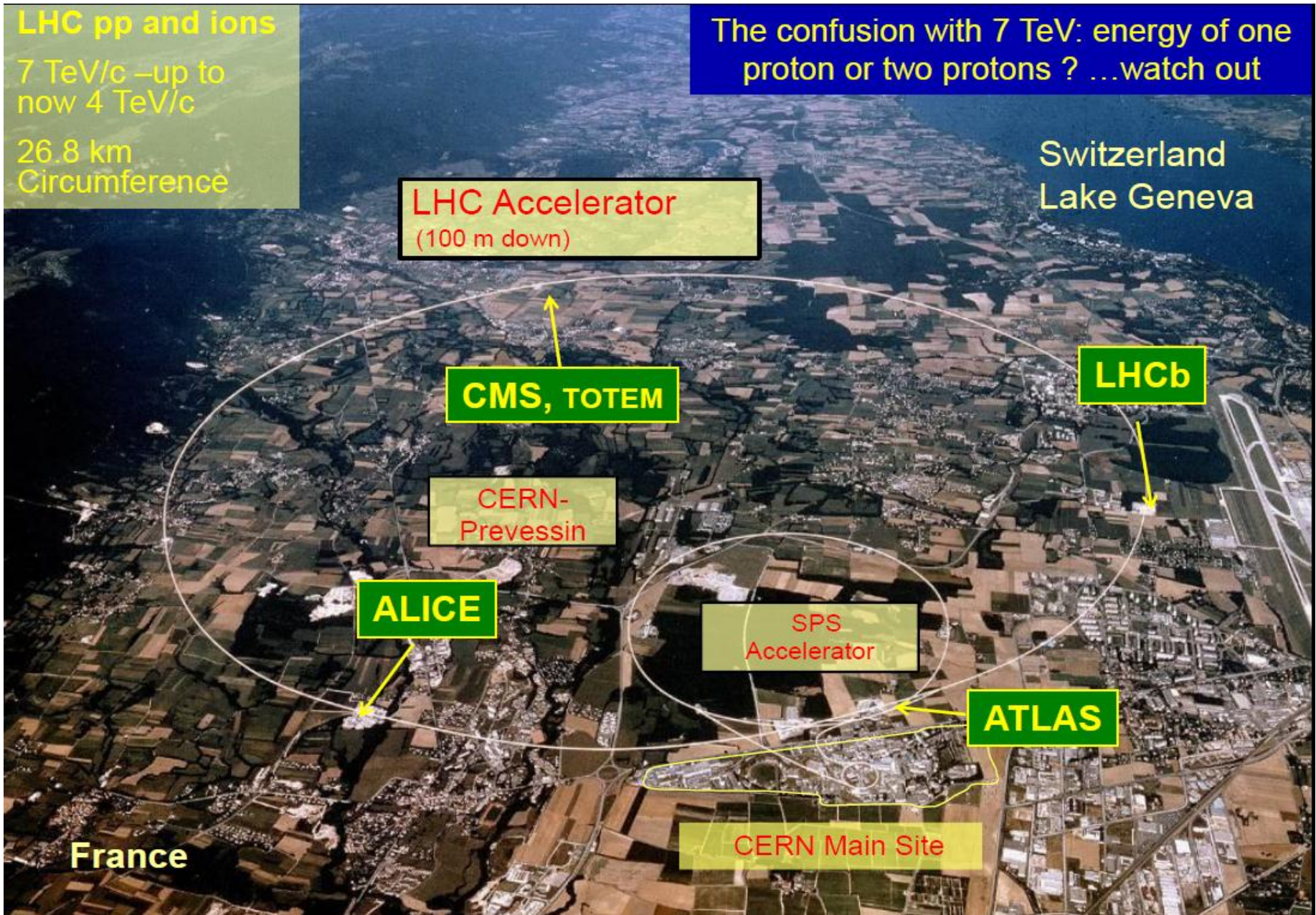
SPS
Accelerator

ATLAS

CERN Main Site

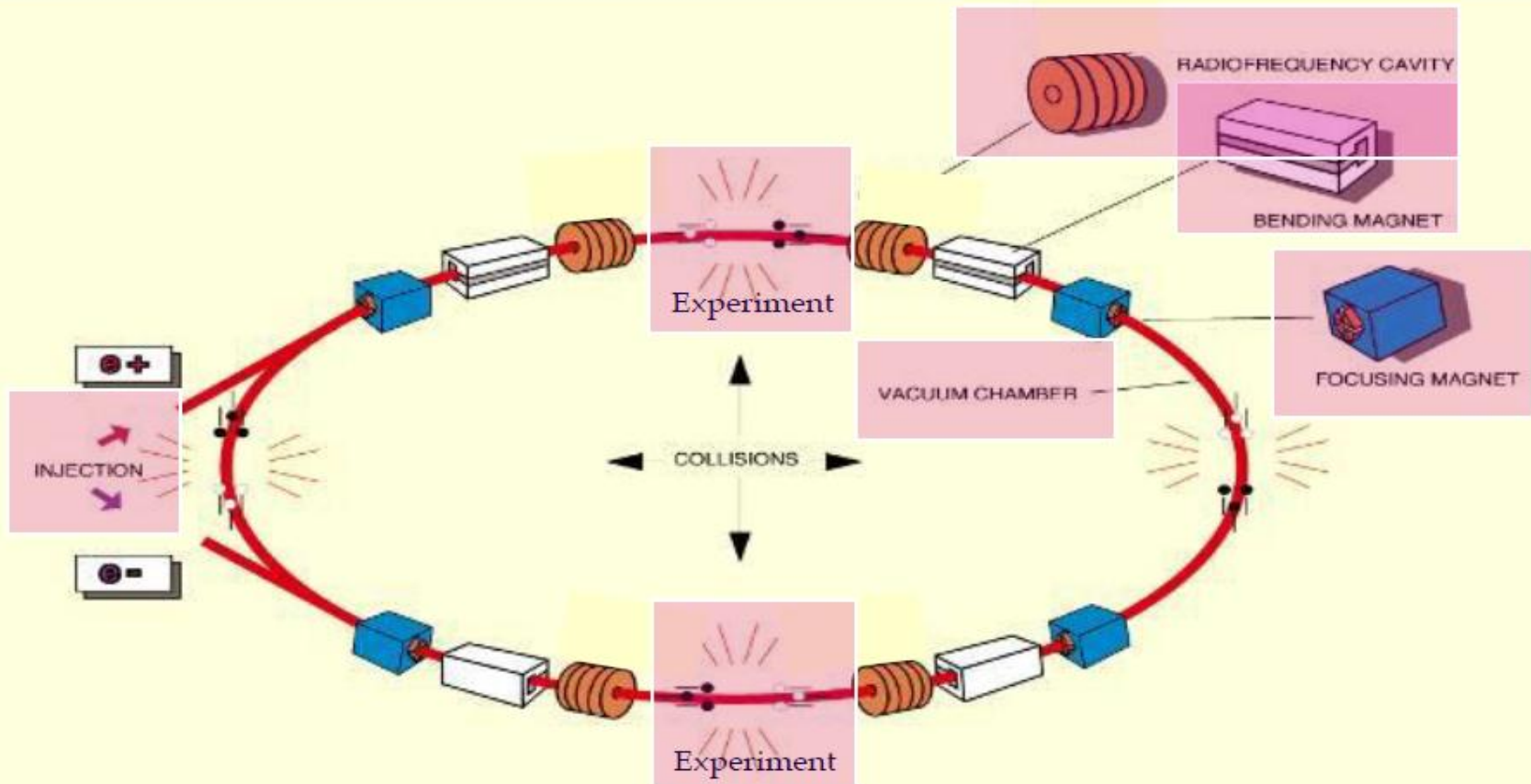
LHCb

France

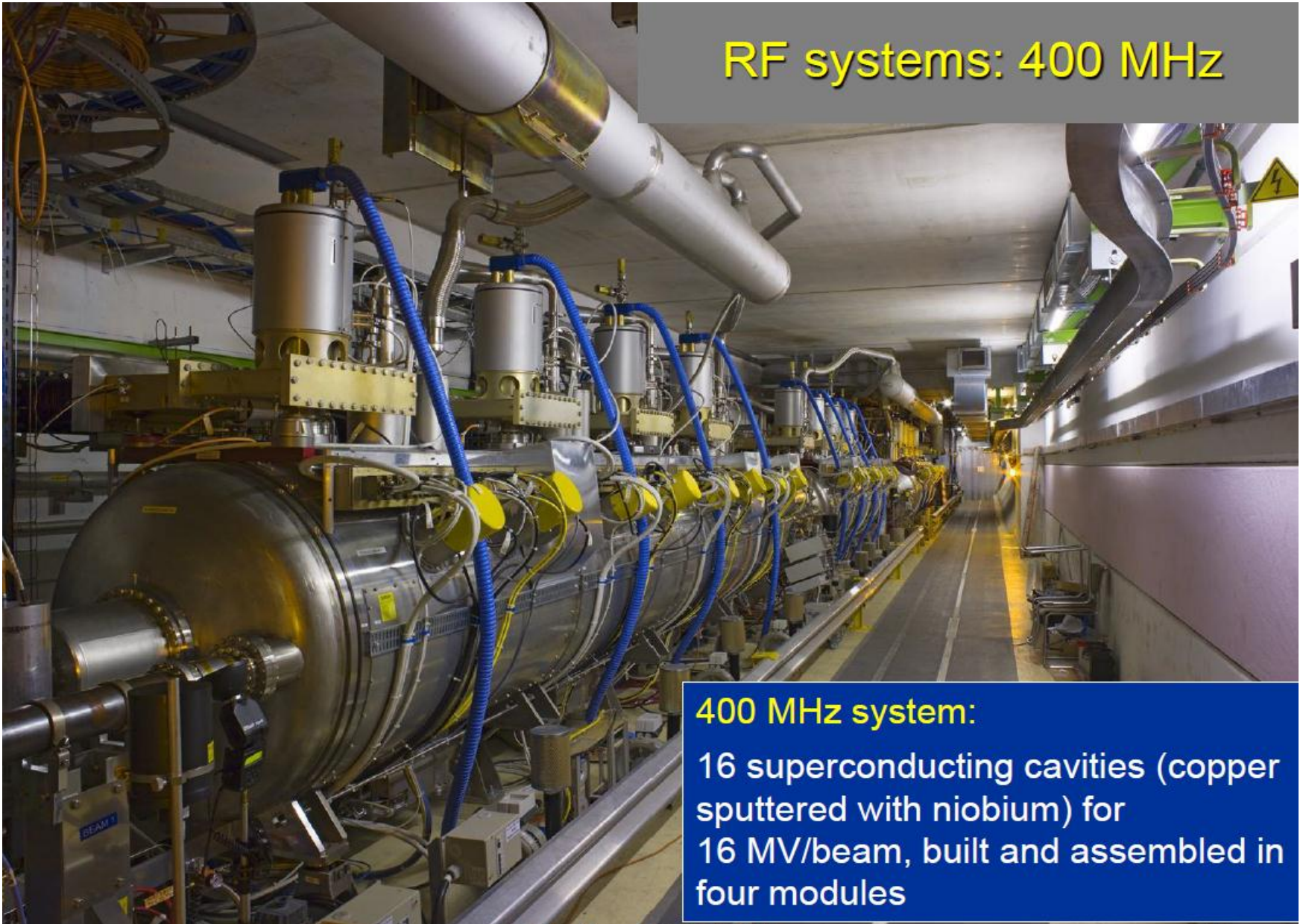


Synchrotron + many passages in RF cavities

LHC **circular machine** with energy gain per turn ~ 0.5 MeV
acceleration from 450 GeV to 7 TeV will take about 20 minutes



RF systems: 400 MHz



400 MHz system:

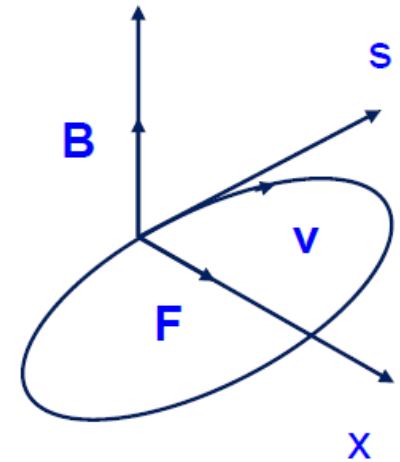
16 superconducting cavities (copper sputtered with niobium) for 16 MV/beam, built and assembled in four modules

Particle deflection: superconducting magnets

The force on a charged particle is proportional to the charge, the electric field, and the vector product of velocity and magnetic field given by Lorentz Force:

$$\vec{F} = q \cdot (\vec{E} + \vec{v} \times \vec{B})$$

$$B = \frac{p}{e_0 \cdot R}$$



Maximum momentum 7000 GeV/c

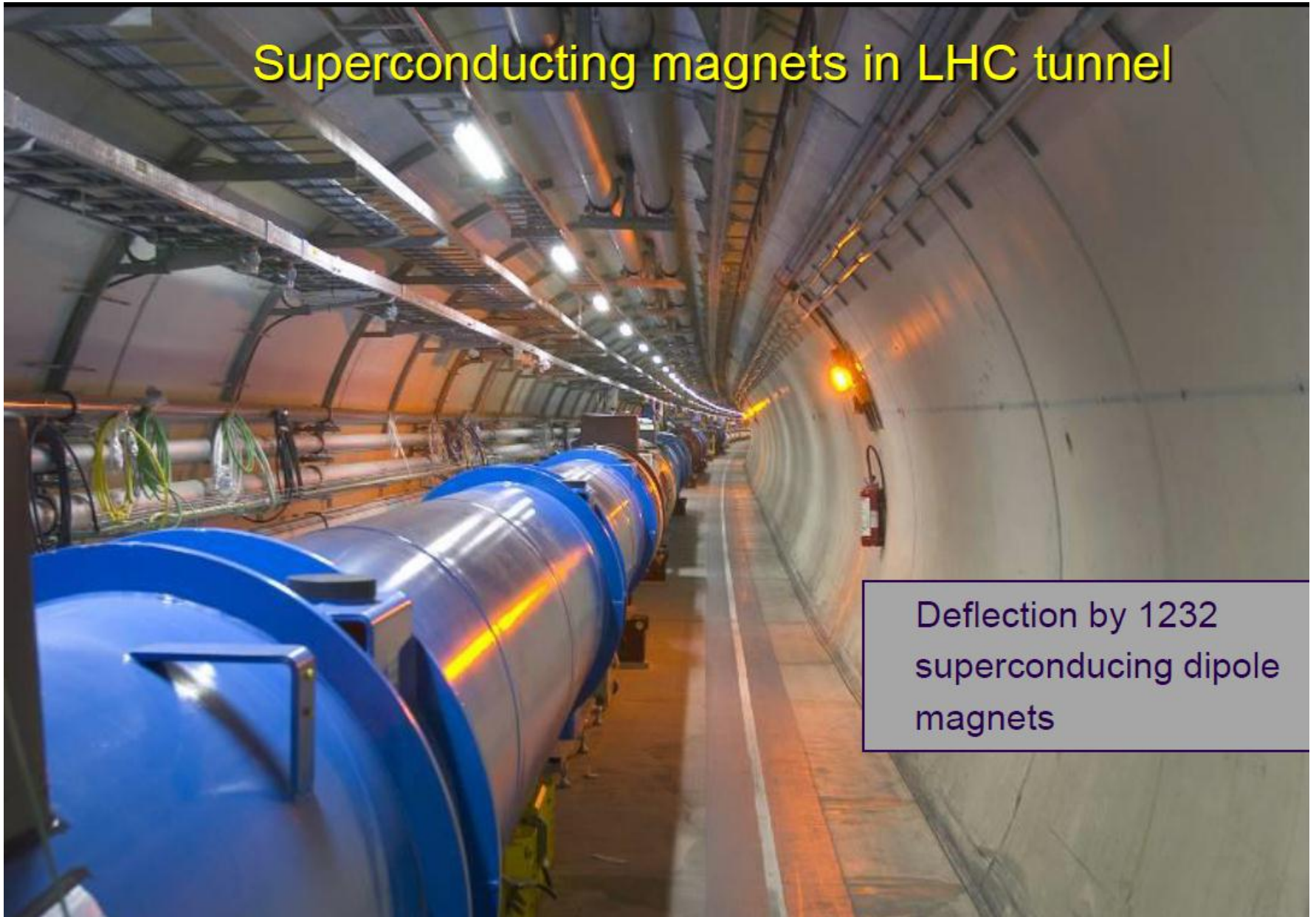
Radius 2805 m fixed by LEP tunnel

Magnetic field B = 8.33 Tesla

Iron magnets limited to 2 Tesla, therefore superconducting magnets are required

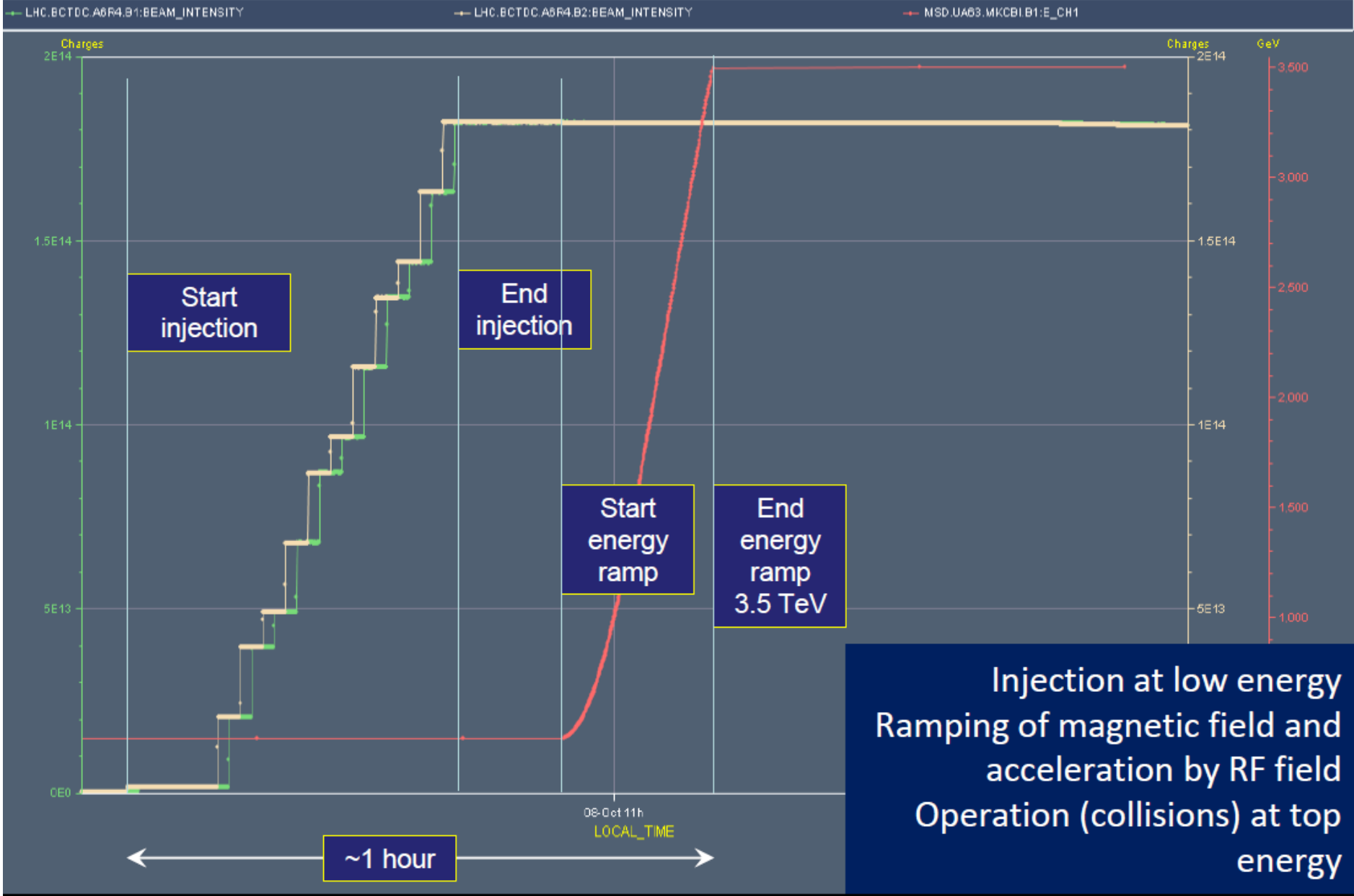
Deflecting magnetic fields for two beams in opposite directions

Superconducting magnets in LHC tunnel

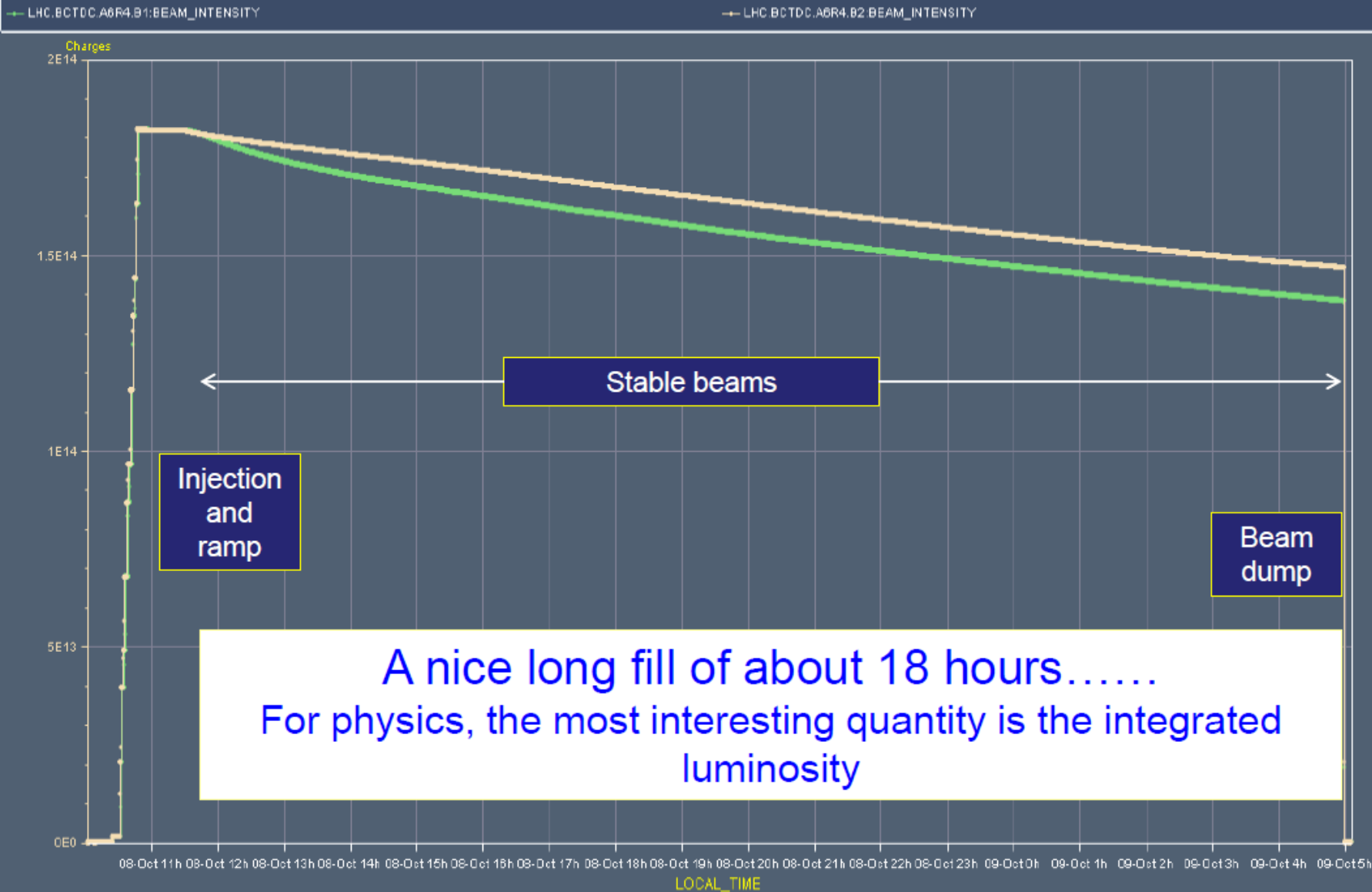


Deflection by 1232
superconducting dipole
magnets

Synchrotron principle: LHC fill (2011)



Excellent fill (2011)



LHC Layout

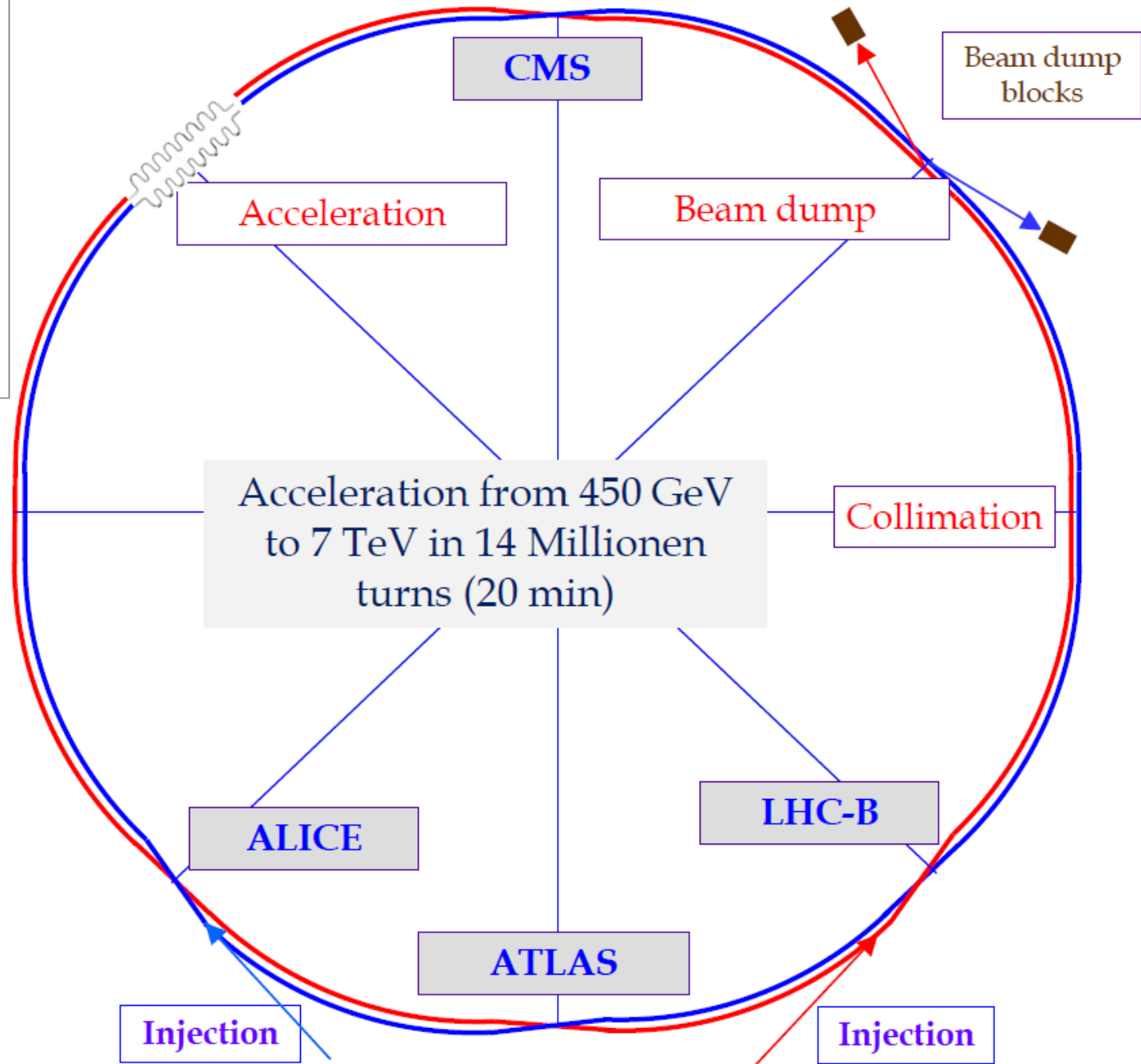
eight arcs (sectors)

eight long straight section (about 700 m long)

1232 deflecting dipole magnets

Collimation

- 27 km
- 2 beams
- 11246 turns/second
- 8 arcs
- 8 straight sections



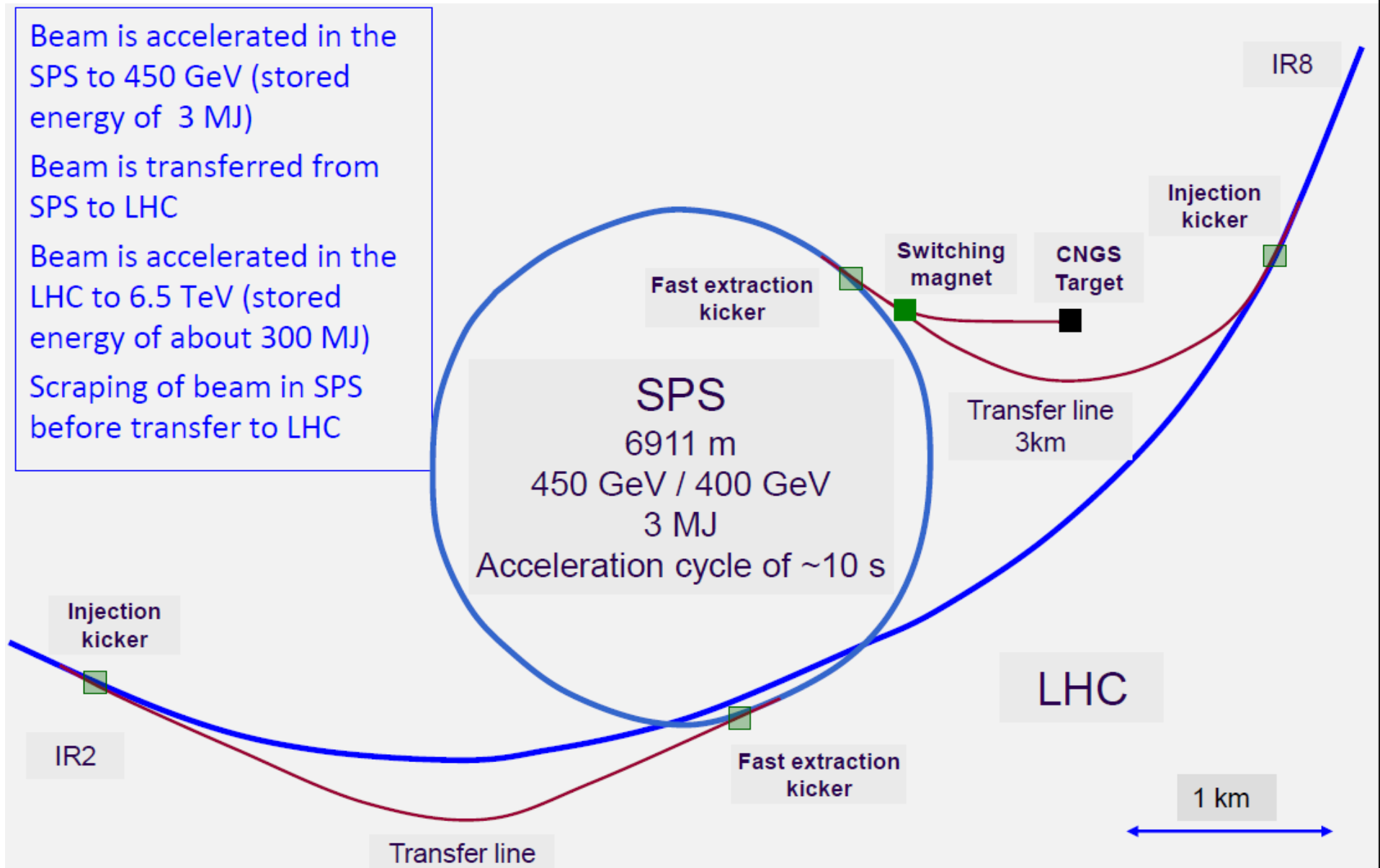
SPS, transfer line and the LHC

Beam is accelerated in the SPS to 450 GeV (stored energy of 3 MJ)

Beam is transferred from SPS to LHC

Beam is accelerated in the LHC to 6.5 TeV (stored energy of about 300 MJ)

Scraping of beam in SPS before transfer to LHC



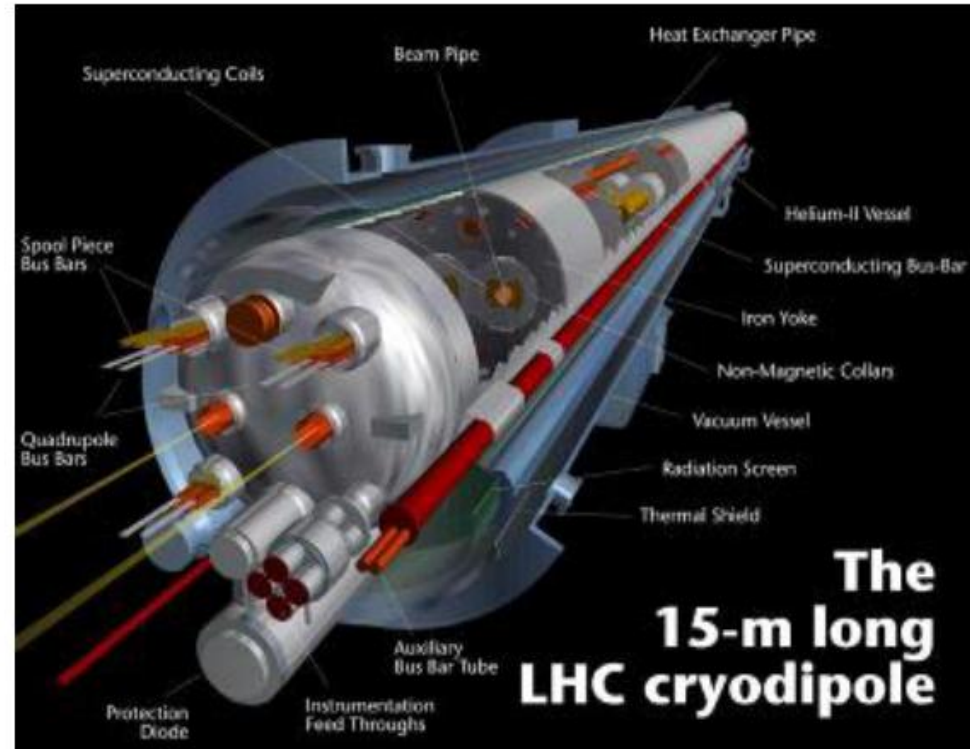
Dipole magnets for the LHC

1232 Dipole magnets
Length about 15 m

Magnetic Field 8.3 T for
7 TeV

Two beam tubes with an
opening of 56 mm

plus many other magnets, to ensure
beam stability (1700 main magnets and
about 8000 corrector magnets)



Colliding trains of bunches

Number of „New Particles“
per unit of time:

$$\frac{N}{\Delta T} = L[\text{cm}^{-2} \cdot \text{s}^{-1}] \cdot \sigma[\text{cm}^2]$$

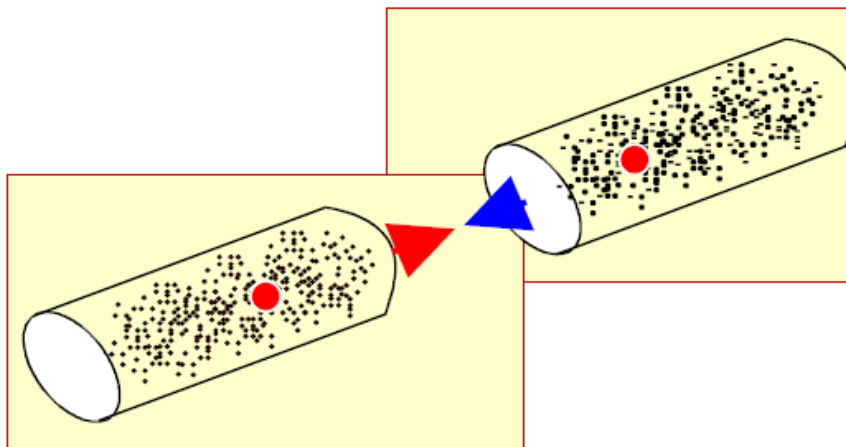
The objective for the LHC as proton – proton collider is a luminosity of about $10^{34} [\text{cm}^{-2}\text{s}^{-1}]$

LEP (e+e-)	:	3-4 $10^{31} [\text{cm}^{-2}\text{s}^{-1}]$
Tevatron (p-pbar)	:	some $10^{32} [\text{cm}^{-2}\text{s}^{-1}]$
B-Factories	:	$> 10^{34} [\text{cm}^{-2}\text{s}^{-1}]$

Luminosity parameters

$$L = \frac{N^2 \times f \times n_b}{4 \times \pi \times \sigma_x \times \sigma_y}$$

$N \dots$	number of protons per bunch
$f \dots$	revolution frequency
$n_b \dots$	number of bunches per beam
$\sigma_x \times \sigma_y \dots$	beam dimensions at interaction point



Luminosity parameters

Number of protons per bunch limited to about $1\text{-}3 \times 10^{11}$ due to the beam-beam interaction and beam instabilities

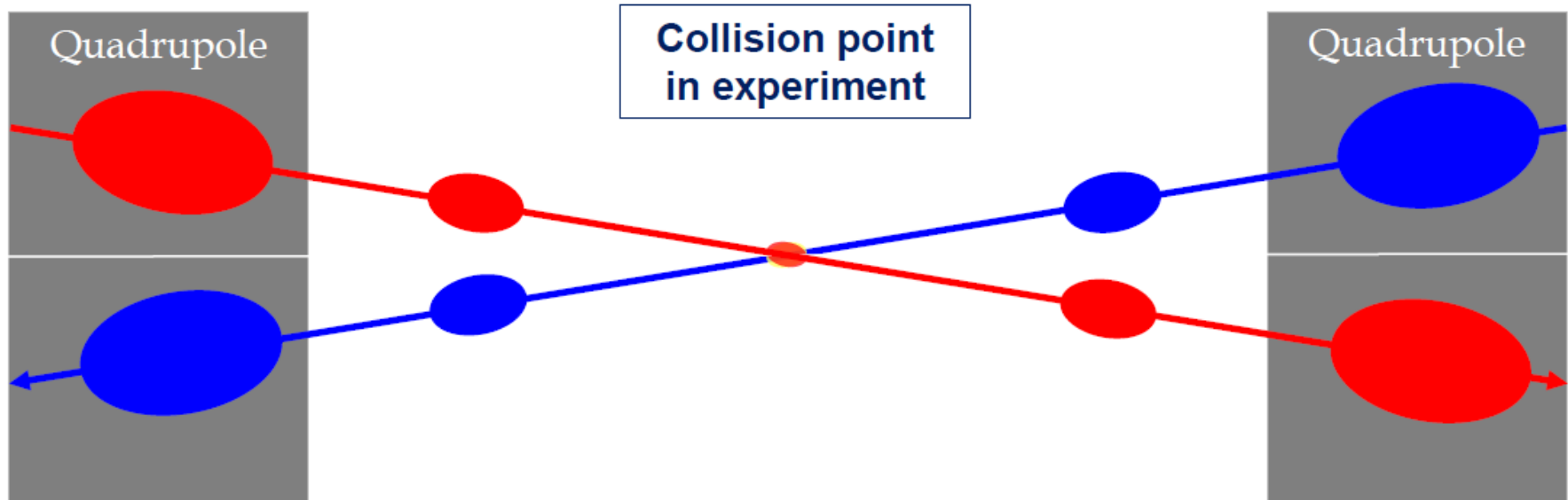
Beam size given by injectors and by space in vacuum chamber

$f = 11246 \text{ Hz}$

Beam size $16 \mu\text{m}$,
for $\beta = 0.5 \text{ m}$ (β is a function of the lattice)

$$L = \frac{N^2 \cdot f \cdot n_b}{4 \cdot \pi \cdot \sigma_x \cdot \sigma_y} = 10^{34} [\text{cm}^{-2}\text{s}^{-1}] \text{ for } 2808 \text{ bunches}$$

Beam size



- Large beam size in adjacent quadrupole magnets
- Separation between beams needed, about 10σ
- Limitation is the aperture in quadrupoles
- Limitation of β function at IP to 0.4 m (2017)

Event pileup in LHC experiments

Assuming nominal parameters, for one bunch crossing, the number of colliding proton pairs (events) is given by:

Event pile up for one bunch crossing:

$$L = \frac{N^2 \times f \times n_b}{4 \times \pi \times \sigma_x \times \sigma_y}$$

Total cross section: $\sigma_{\text{tot}} := 100 \text{mBarn}$

$$\sigma_{\text{tot}} = 1 \times 10^{-25} \text{cm}^2$$

Luminosity: $L = 1 \times 10^{34} \text{s}^{-1} \text{cm}^{-2}$

Number of events per second $L \cdot \sigma_{\text{tot}} = 1 \times 10^9 \frac{1}{\text{s}}$

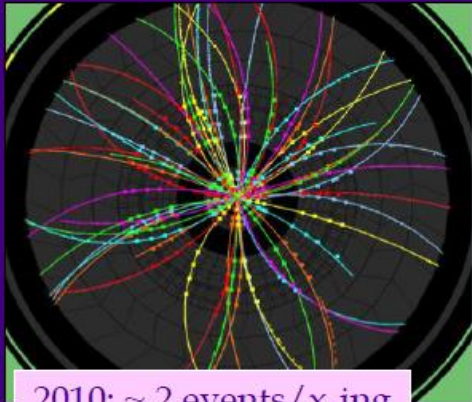
$\text{frev}_{\text{LHC}} = 1.1246 \times 10^4 \frac{1}{\text{s}}$ and $N_{\text{bunches_1beam}} = 2808$

Number of events per bunch crossing: $L \cdot \frac{\sigma_{\text{tot}}}{\text{frev}_{\text{LHC}} \cdot N_{\text{bunches_1beam}}} = 31.7$

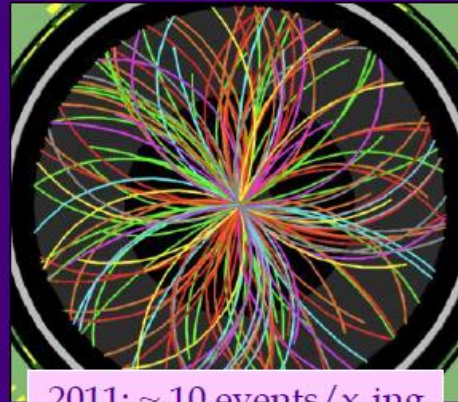
CMS

E
CMS Experiment at LHC, CERN
Data recorded: Mon May 28 01:16:20 2012 CERN
Run/Event: 195099 / 35488125
Lumi Section: 65
Orbit/Crossing: 16992111 / 2295

- ⇒ With the parameters of 2012 for each bunch crossing there are up to ~35 interactions (lower luminosity, less number of bunches)
- ⇒ 'Hats off' to ATLAS & CMS for handling this pile-up !!



2010: ~ 2 events/x-ing

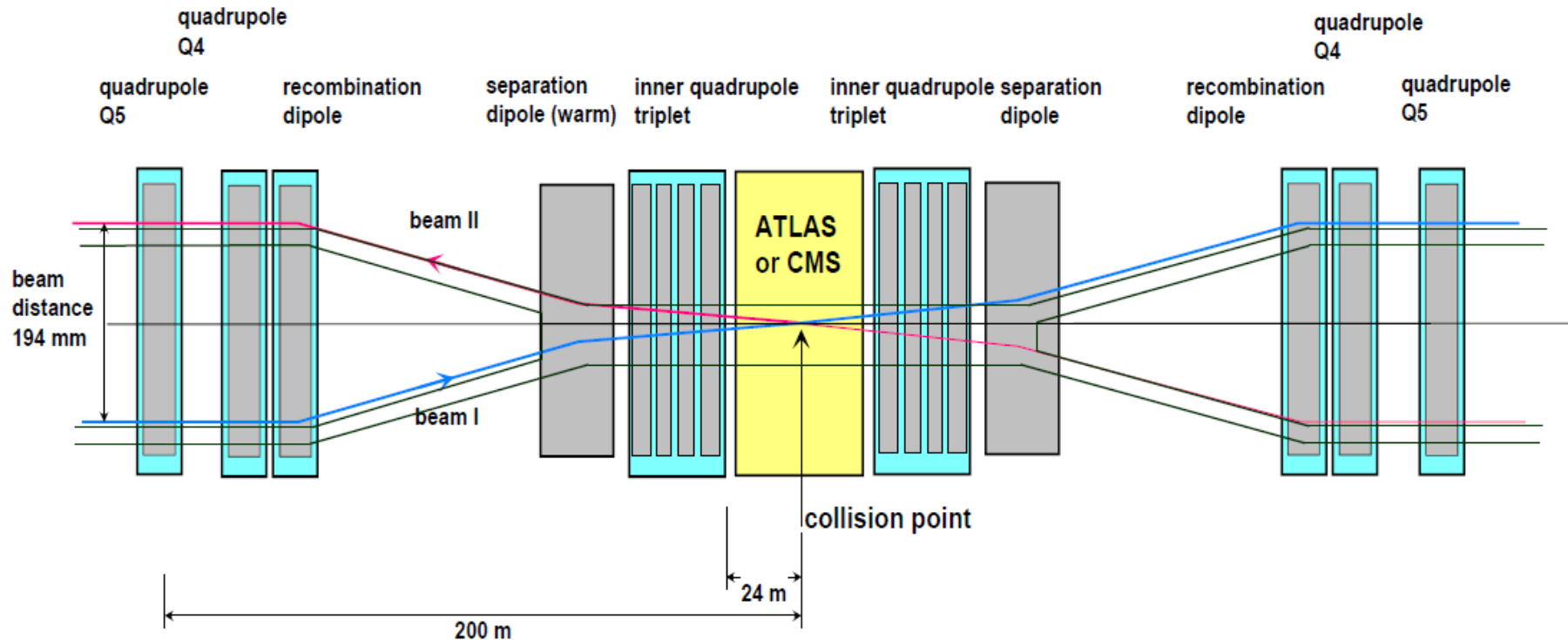


2011: ~ 10 events/x-ing



2012: ~ 20 events/x-ing

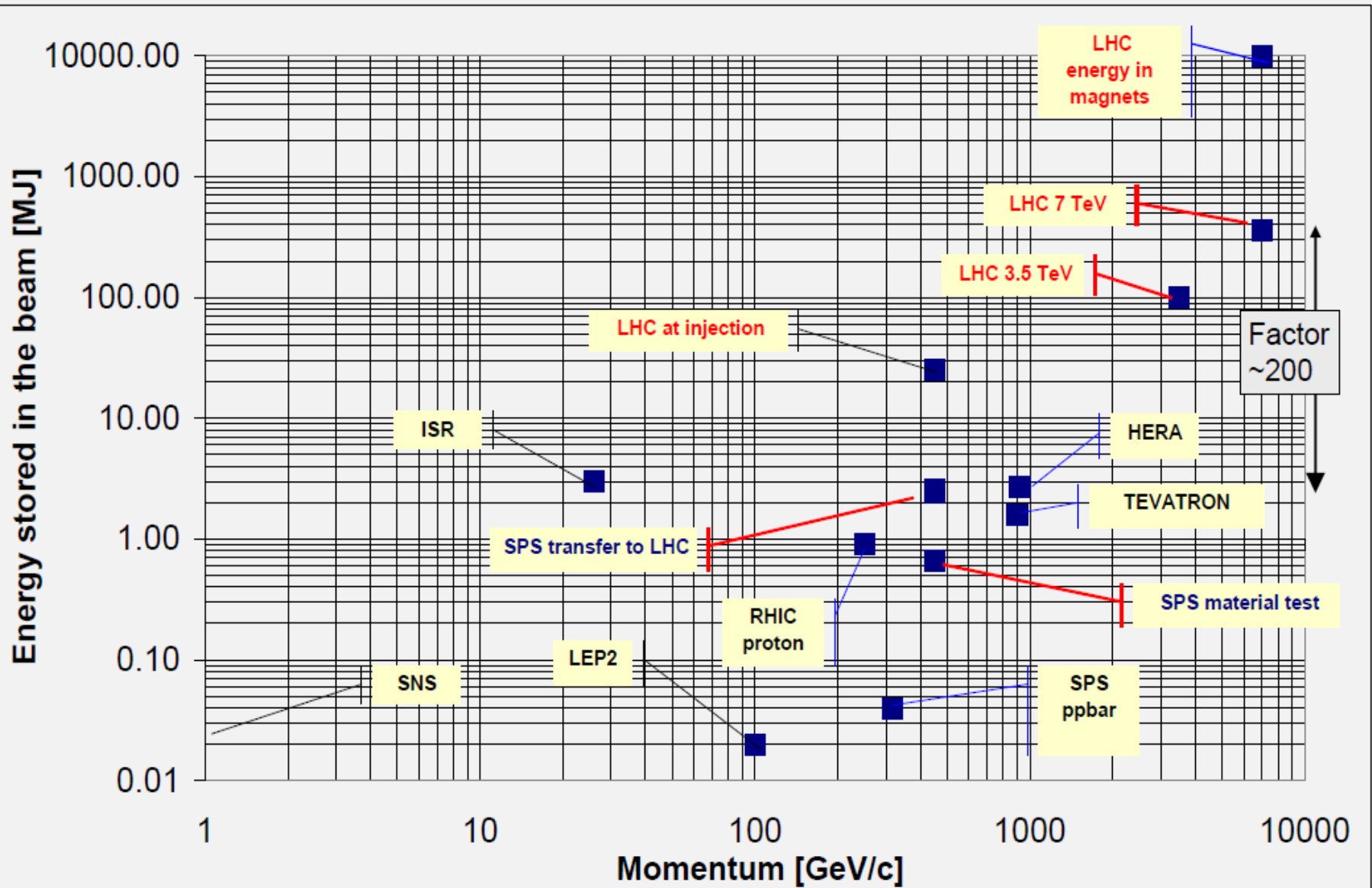
Experimental long straight section



Example for an LHC insertion with ATLAS or CMS

- The 2 LHC beams are brought together to collide in a 'common' region
- Over ~260 m the beams circulate in one vacuum chamber with 'parasitic' encounters (when the spacing between bunches is small enough)
- Total crossing angle of about 250 μrad

Energy stored in the beam



What does it mean?

The energy of an 200 m long fast train at 155 km/hour corresponds to the energy of 360 MJoule stored in one LHC beam



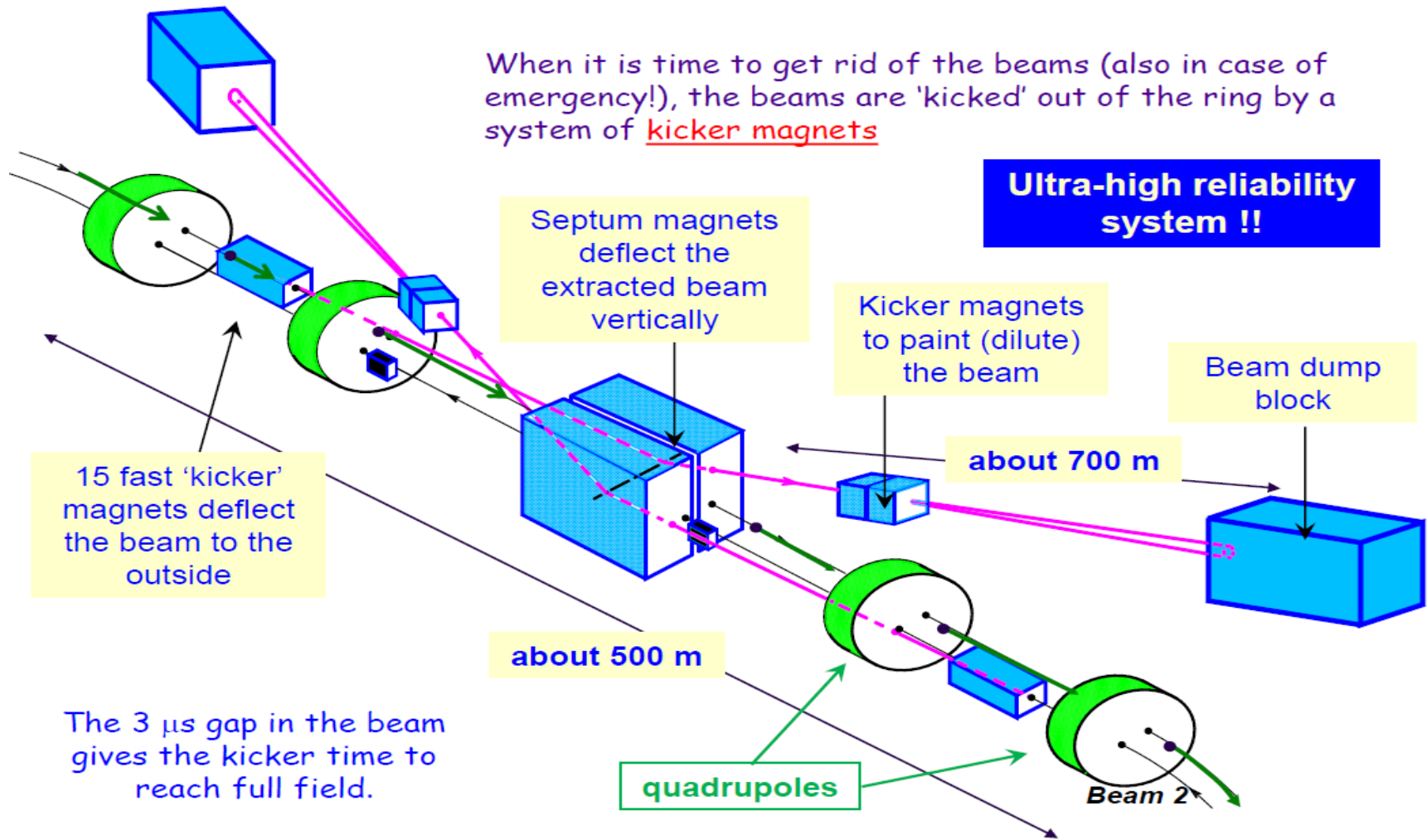
360 MJoule: the energy stored in one LHC beam corresponds approximately to...

- 90 kg of TNT
- 8 litres of gasoline
- 15 kg of chocolate



It's how ease the energy is released that matters most !!

Layout of beam system dump





Dump line



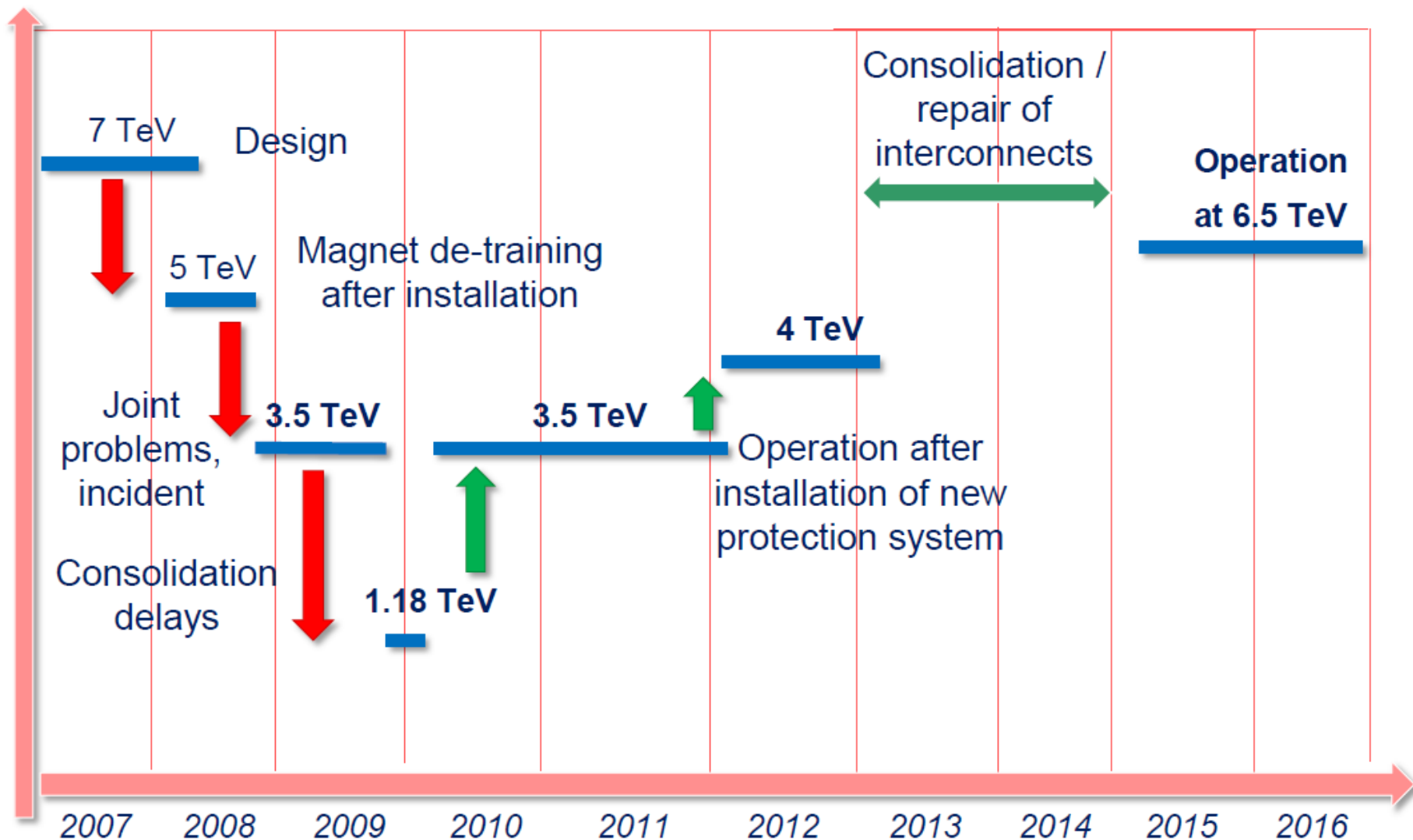
Beam Loss Monitors

- Ionization chambers to detect beam losses:
 - Reaction time $\sim \frac{1}{2}$ turn ($40 \mu\text{s}$)
 - Very large dynamic range ($> 10^6$)
- There are **~ 3600 chambers** distributed over the ring to detect abnormal beam losses and if necessary trigger a beam abort !
- Very important beam instrumentation!



LHC energy evolution

Energy (TeV)



The LHC: just another collider?

	Start	Type	Max proton energy [GeV]	Length [m]	B Field [Tesla]	Lumi [$\text{cm}^{-2}\text{s}^{-1}$]	Stored beam energy [MJoule]
TEVATRON Fermilab Illinois USA	1983	p-pbar	980	6300	4.5	$4.3 \cdot 10^{32}$	1.6 for protons
HERA DESY Hamburg	1992	p – e+ p – e-	920	6300	5.5	$5.1 \cdot 10^{31}$	2.7 for protons
RHIC Brookhaven Long Island	2000	Ion-Ion p-p	250	3834	4.3	$1.5 \cdot 10^{32}$	0.9 per proton beam
LHC CERN	2008	Ion-Ion p-p	7000 Now 4000	26800	8.3	10^{34} Now 7.7×10^{33}	362 per beam
Factor			7	4	2	50	100

LHC pp and ions

7 TeV/c –up to
now 4 TeV/c

26.8 km
Circumference

The confusion with 7 TeV: energy of one
proton or two protons ? ...watch out

Switzerland
Lake Geneva

LHC Accelerator
(100 m down)

CMS, TOTEM

CERN-
Prevezin

ALICE

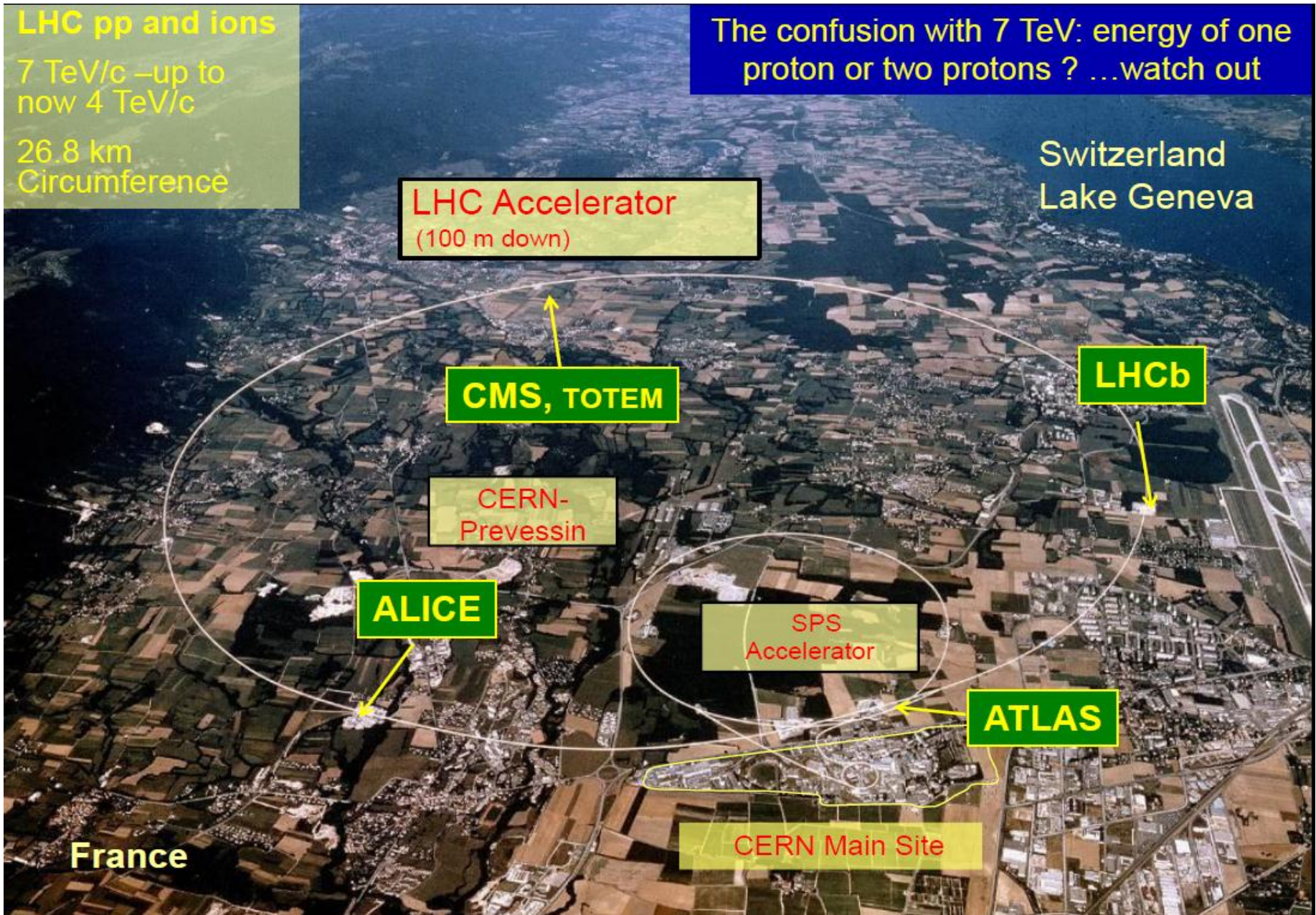
SPS
Accelerator

ATLAS

LHCb

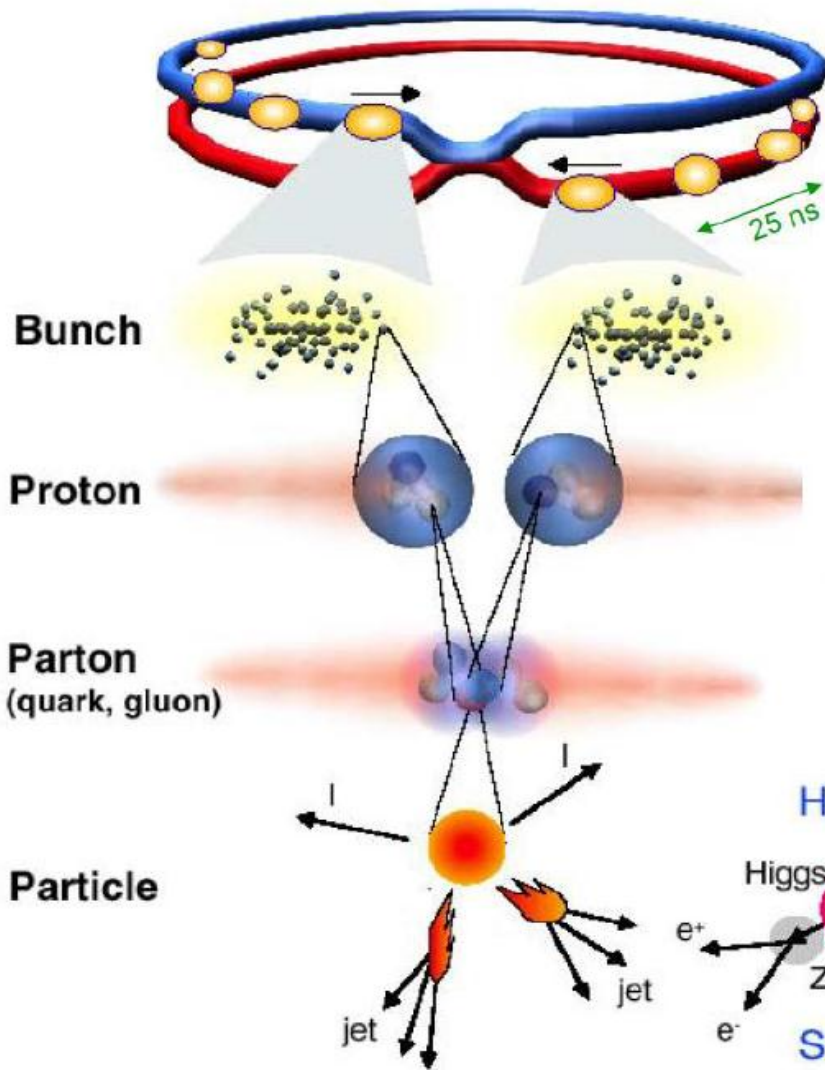
CERN Main Site

France



Collisions at LHC

Proton-Proton	2835 bunch/beam
Protons/bunch	10^{11}
Beam energy	7 TeV (7×10^{12} eV)
Luminosity	10^{34} cm ⁻² s ⁻¹

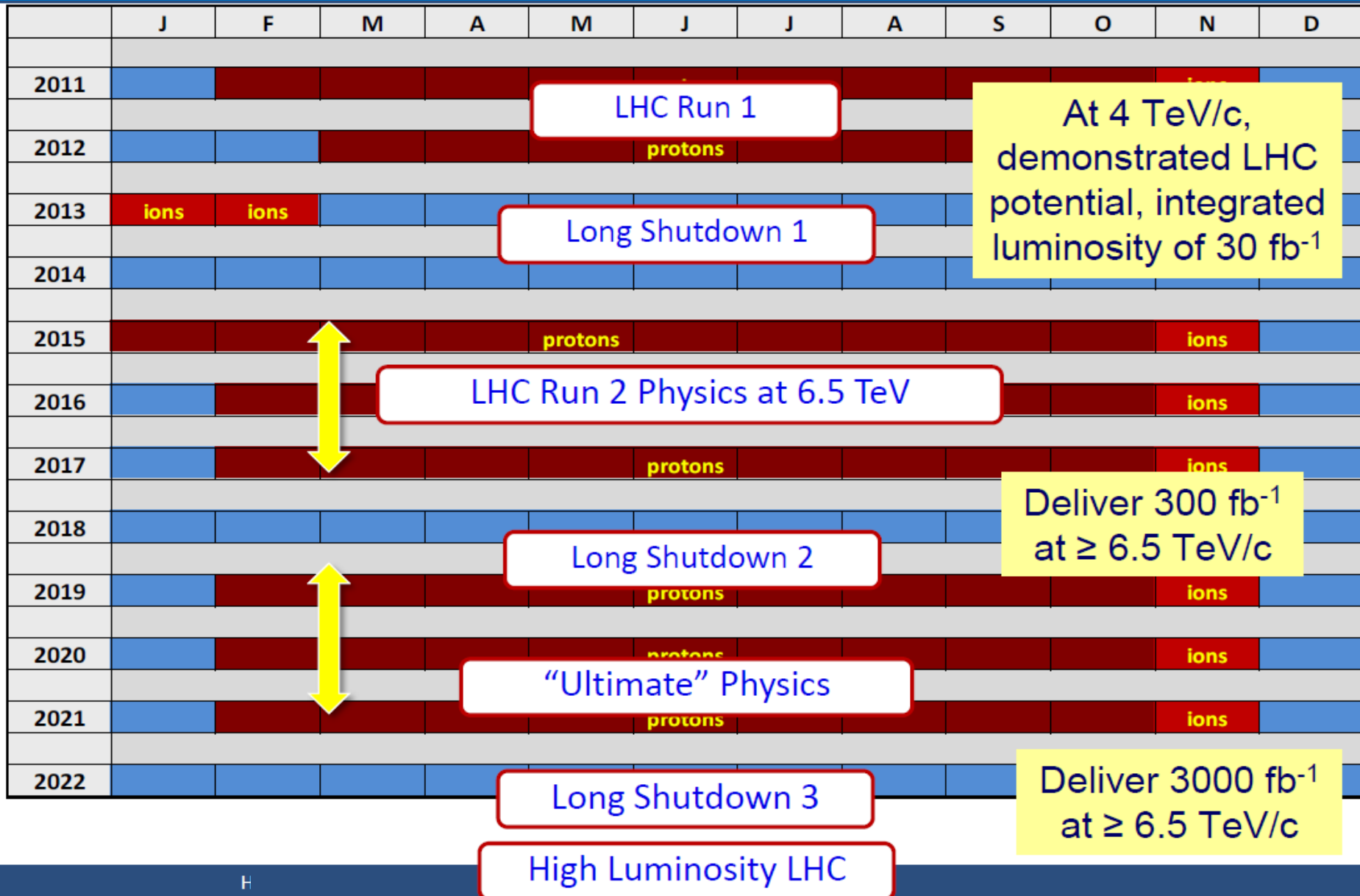


In the experiments:
 10^9 pp interactions per second
 ~ 1500 particles (p, n, π) produced in the detectors at each bunch-crossing

**Selection of 1 in
 10,000,000,000,000**

$$E = m c^2$$

The next years



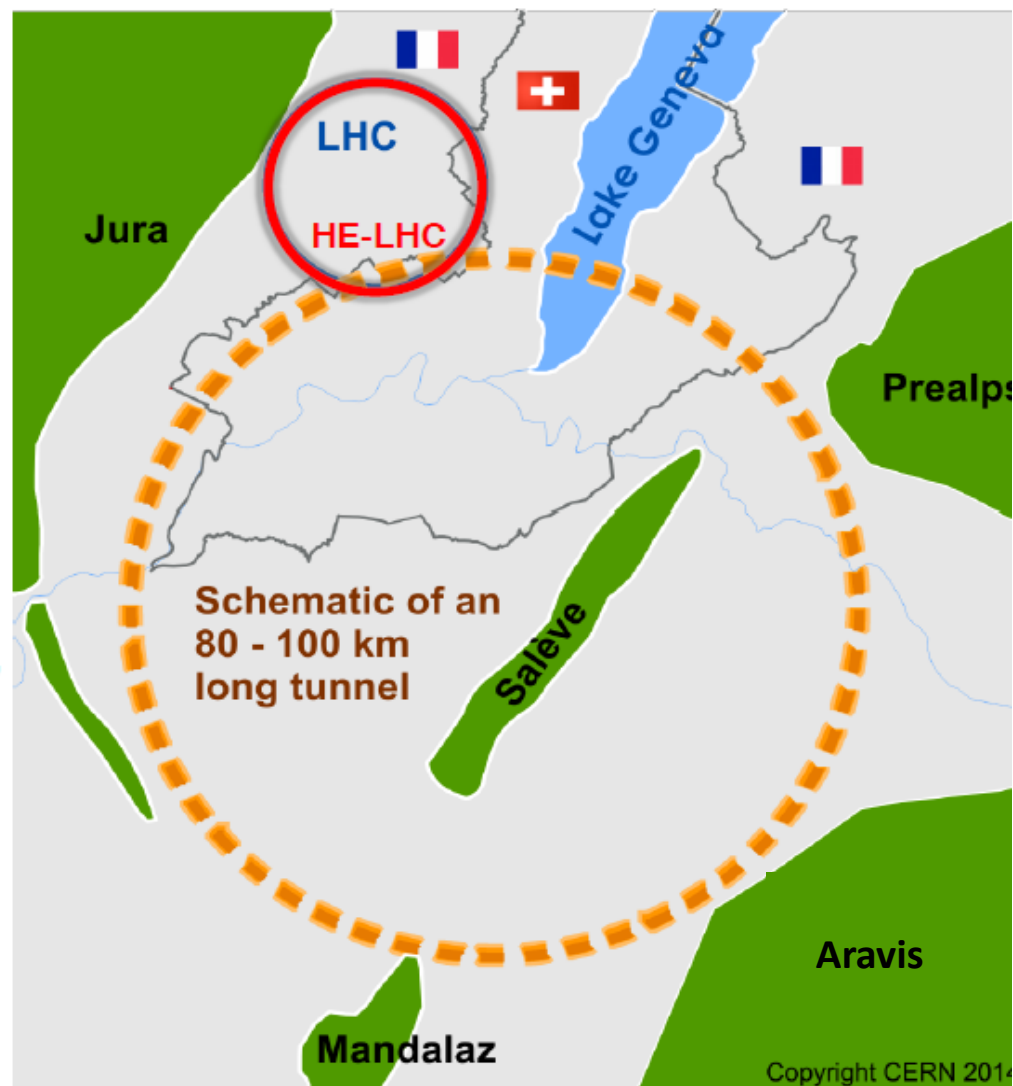
Future plans

international FCC collaboration (CERN as host lab) to design:

- pp -collider (*FCC-hh*)
→ main emphasis, defining infrastructure requirements

~16 T ⇒ 100 TeV pp in 100 km

- 80-100 km tunnel infrastructure in Geneva area, site specific
- e^+e^- collider (*FCC-ee*), as a possible first step
- $p-e$ (*FCC-he*) option, one IP, FCC-hh & ERL
- **HE-LHC** w *FCC-hh* technology



SPARES

Aside: units

Our scale	Particle Physics	Convert
Length m	Length fm	1 eV = 1.6×10^{-19} J
Mass kg	Mass eV/c ²	1 GeV = 10^9 eV
Time s	Time s	1 TeV = 10^3 GeV
Energy kg m ² s ⁻²	Energy eV	1 fm = 10^{-15} m

Note: often set $\hbar = c = 1$