

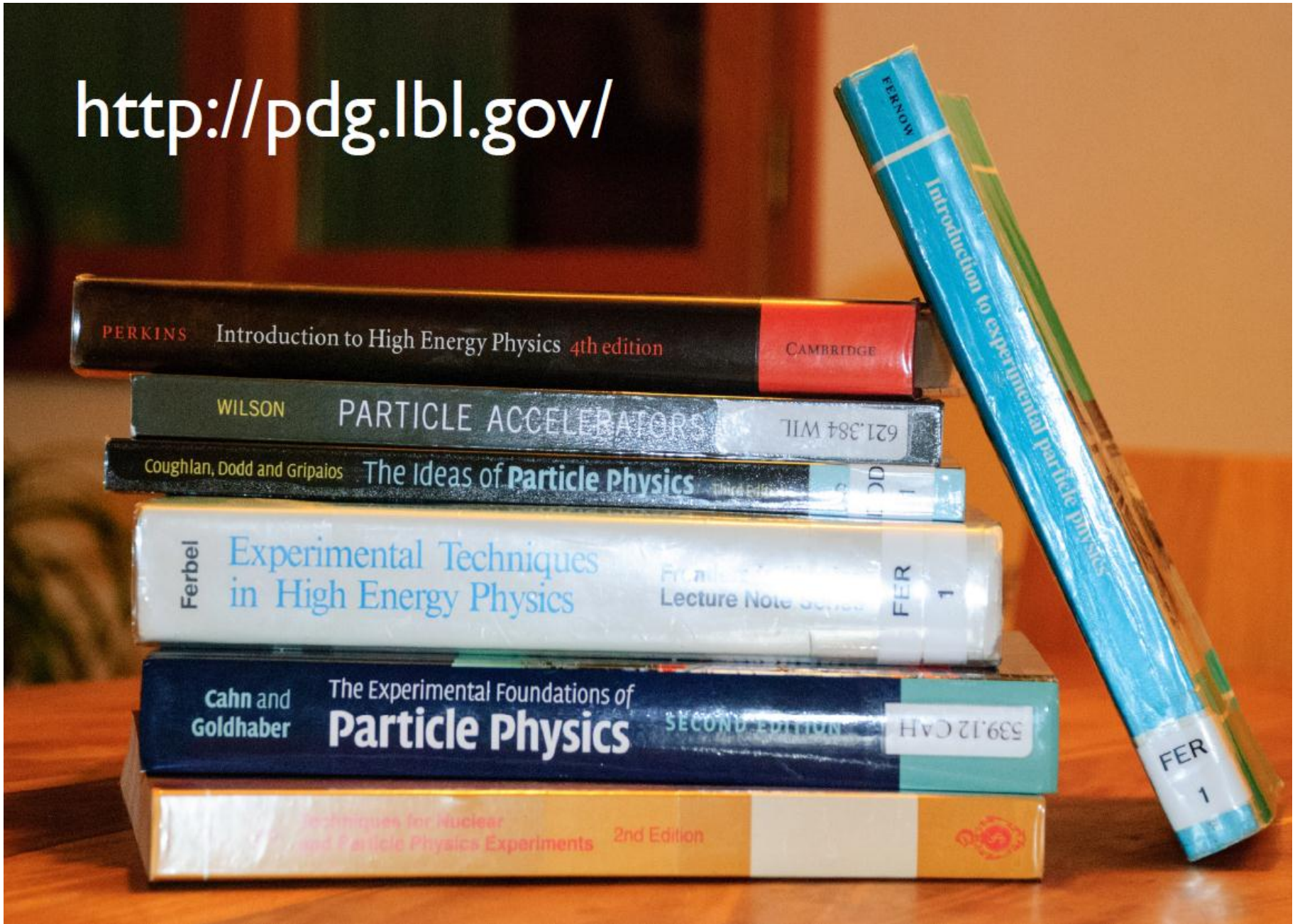
# Introduction to particle physics: experimental part

**Introduction**

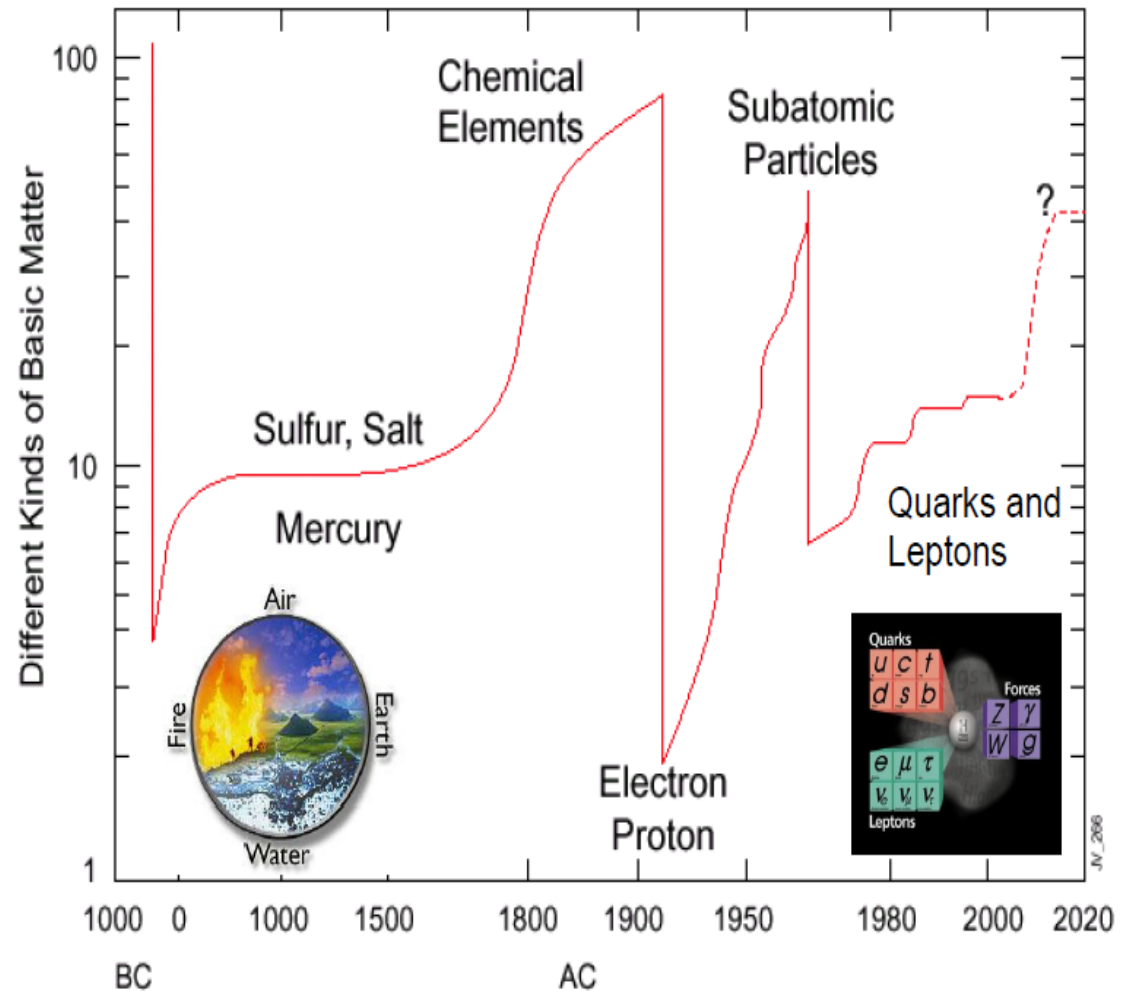
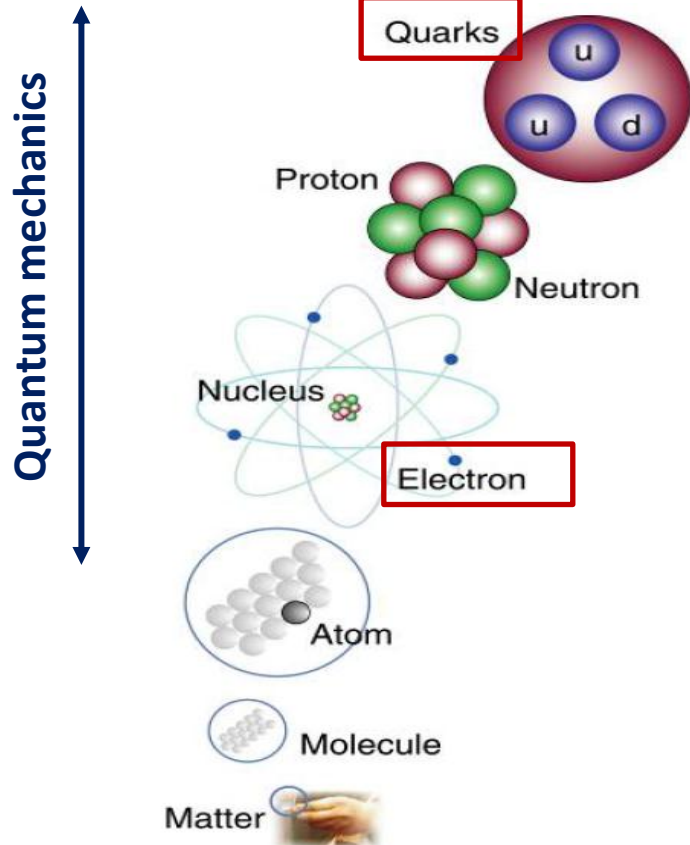
**Units, kinematics**

Large fraction of those slides from M. Delmastro lectures at ESIPAP school

<http://pdg.lbl.gov/>

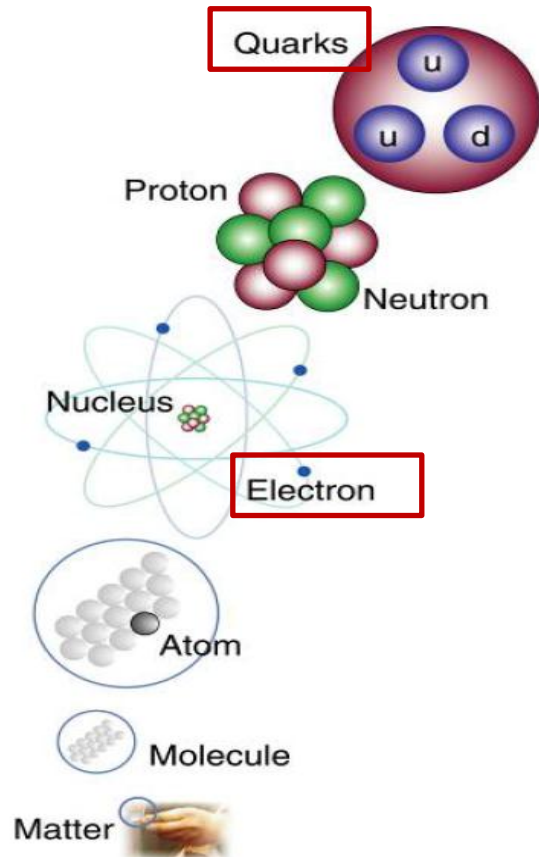


# Constituents of matter along History



# Particles of the Standard Model

Quantum mechanics



**Matter particles**  
( $< 10^{-16}$  cm)

**Interaction particles**

2.4M <b>u</b> up 2/3 1/2	1.27G <b>c</b> charm 2/3 1/2	171.2G <b>t</b> top 2/3 1/2	strong nuclear force (color charge)
4.8M <b>d</b> down -1/3 1/2	104M <b>s</b> strange -1/3 1/2	4.2G <b>b</b> bottom -1/3 1/2	
0.511M <b>e</b> electron -1 1/2	105.7M <b>μ</b> muon -1 1/2	1.777G <b>τ</b> tau -1 1/2	
< 2.2 <b>ν<sub>e</sub></b> e-neutrino 0 1/2	< 0.17M <b>ν<sub>μ</sub></b> μ-neutrino 0 1/2	< 15.5M <b>ν<sub>τ</sub></b> τ-neutrino 0 1/2	electromagnetic (charge)
			<b>γ</b> photon 1
			weak nuclear force
			<b>Z</b> 1

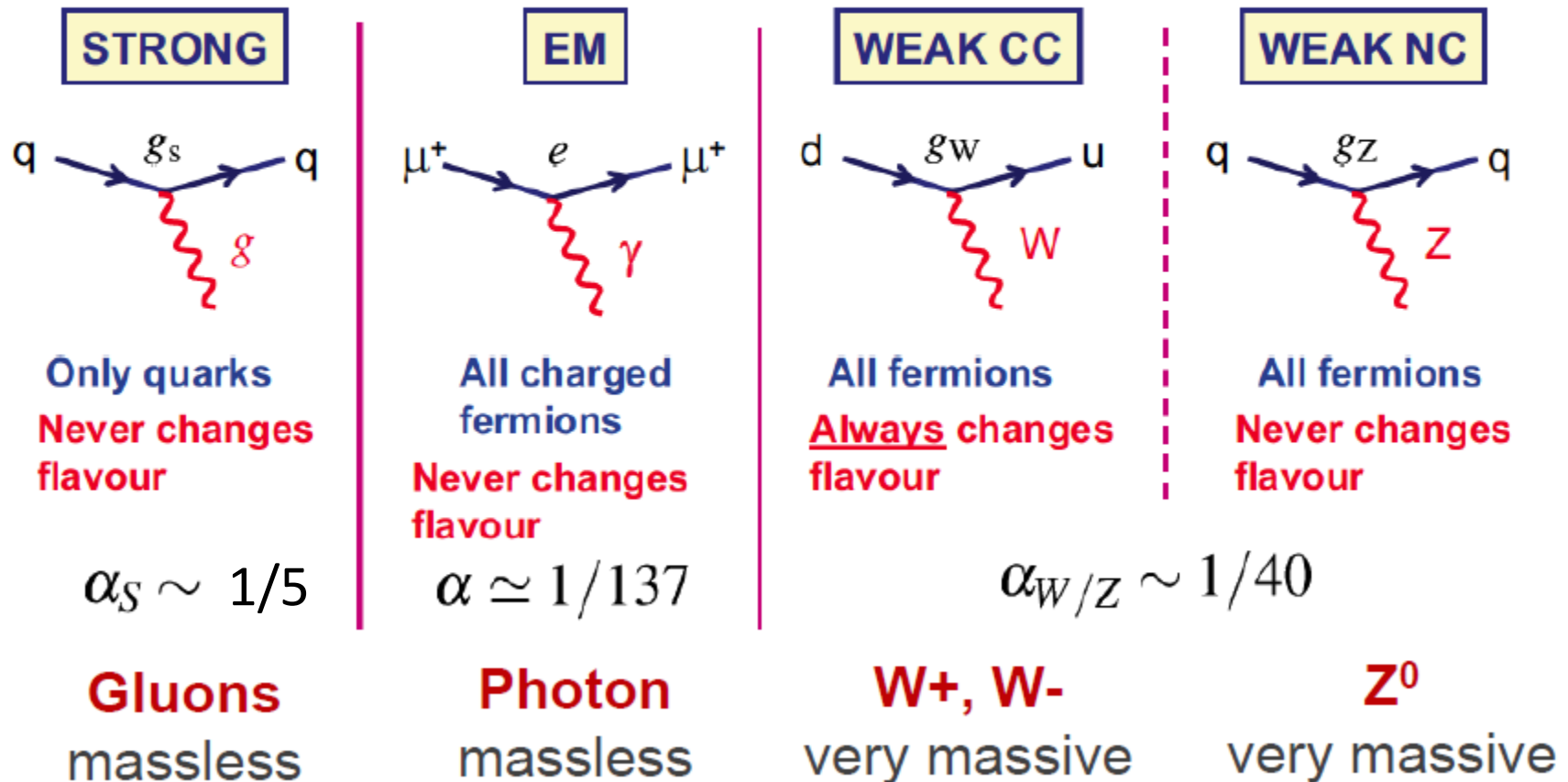


**Higgs particle**  
Is not a matter particle  
nor an interaction particle

$$L_H = \frac{1}{2}(\partial_\mu H)^2 - m_H^2 H^2 - h\lambda H^3 - \frac{h}{4}H^4 + \frac{g^2}{4}(W_\mu^+ W^\mu + \frac{1}{2\cos^2\theta_W} Z_\mu Z^\mu)(\lambda^2 + 2\lambda H + H^2) + \sum_{l,q,q'} (\frac{m_l \bar{l}l}{\lambda} + \frac{m_q \bar{q}q}{\lambda} + \frac{m_{q'} \bar{q}'q'}{\lambda})H$$

# Interactions

The interaction of gauge bosons with fermions is described by the Standard Model



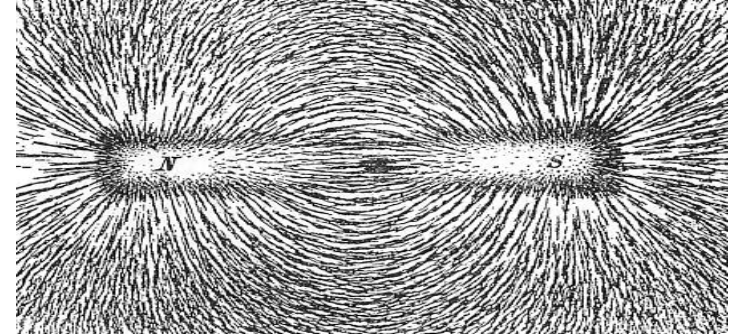
# Vacuum → Concept of the field → Higgs mechanism



Michael Faraday (1845)

Introduced concept of field into description of magnetic interactions. Magnetic field has a source.

Force lines of the magnetic field



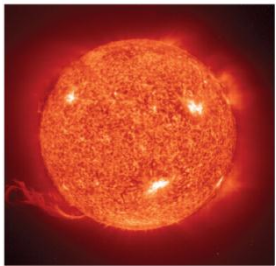
In classical understanding vacuum = „empty space”, or state of „energy = 0”. In quantum mechanics vacuum is „full of life”. Just bubbling with creation and annihilation of matter particles. Symmetries dictate laws of interactions. Stable state of vacuum does not necessarily has an energy = zero.

## Higgs mechanism (1964):

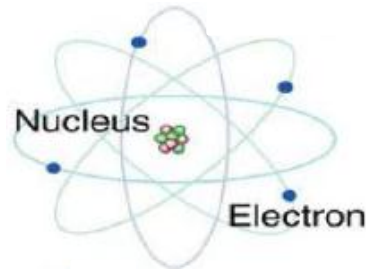
Quantum vacuum is not empty, is full of the „Higgs field”. This field has no structure nor sources. After „symmetry is spontaneously broken”, field is not „neutral” anymore and interaction of elementary particles with this field gives masses to the elementary particles . This mechanism predicted also the existence of a massive scalar particle, so called, „Higgs particle” which is the quantum fluctuation of the Higgs field.

# Mass Spectrum of Elementary Particles

If the mass of the W boson was smaller, time of combustion of the sun (nuclear fusion) would be shorter and at lower temperature.

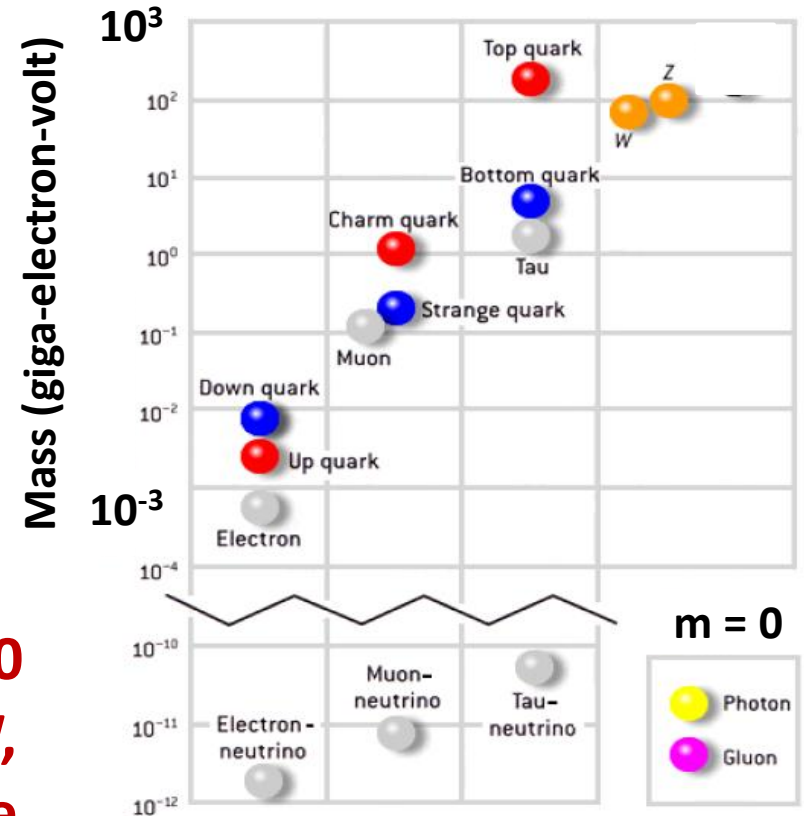


Sun shield



Mass of the electrons is 25000 times smaller than mass of W, but if it was exactly zero there would be no atomic binding.

If the masses of elementary particles were different the Universe, as we know it, would not exist.

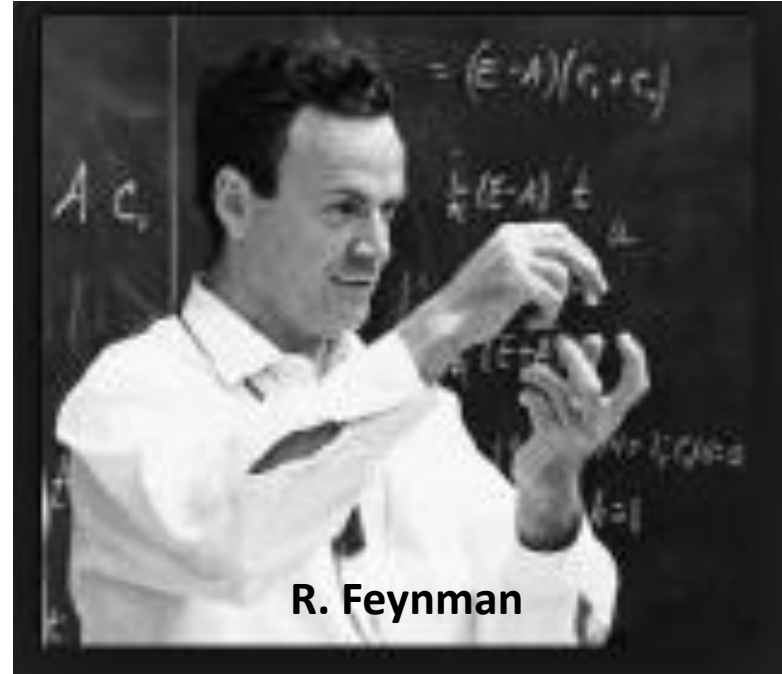


# Theory has to be confirmed by experiment

***„It doesn't matter how beautiful your theory is, it doesn't matter how smart you are. If it doesn't agree with experiments, it's wrong.”***

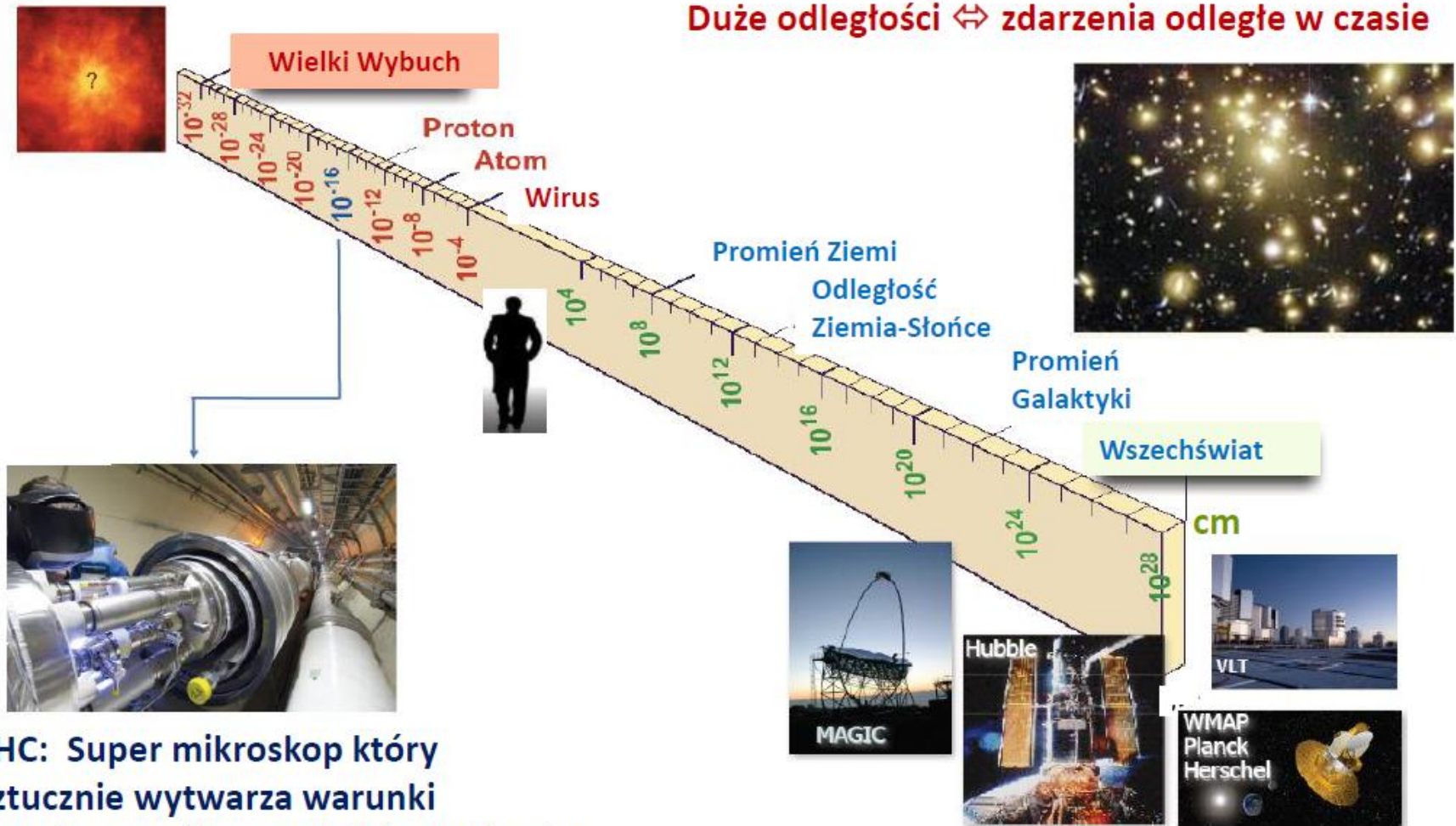
**R. Feynman**

***(To nie ma znaczenia jak piękna jest twoja teoria, nie ma znaczenia jaki jesteś inteligentny. Jeżeli nie zgadza się z eksperymentem to ta teoria jest nieprawdziwa.)***





# How we are probing structure of matter and how structure of the Universe



**Duże odległości ⇔ zdarzenia odległe w czasie**

# Nobel Prizes in Elementary Particle Physics



Sin-Itiro Tomonaga



Julian Schwinger



Richard P. Feynman



Sheldon Lee Glashow



Abdus Salam



Steven Weinberg

**GREEN** - theoretical  
**BLUE** - experimental

1964: „Higgs mechanism”  
was born



Leon M. Lederman



Melvin Schwartz



Jack Steinberger

1957 – C. N. Yang, T. Lee

1965 – S. I. Tomonaga, J. Schwinger, R.P Feynman

1969 – M. Gell-Mann

1976 – B. Richter and S. Ting

1979 – S.L. Glashow, A. Salam, S. Weinberg

1980 – J. Cronin, V. Fitch

1984 – C. Rubbia, S. van der Meer

1988 – L. M. Lederman, M. Schwartz, J. Steinberger

1990 – J. Friedman, J. Kendall, R. Taylor

1992 - G. Charpak

1995 – M. Perl, F. Reines

1999 - G. tHooft, M. J. Veltman

2004 - D. J. Gross, H. D. Politzer, F. Wilczek

2008 – Y. Nambu, M. Kobayashi, T. Masakawa

2013 – F. Englert and P. Higgs

2015 – T. Kajita and A. B. McDonald

2012: „Higgs particle”  
was discovered



Carlo Rubbia



Simon van der Meer



Georges Charpak



Gerardus 't Hooft



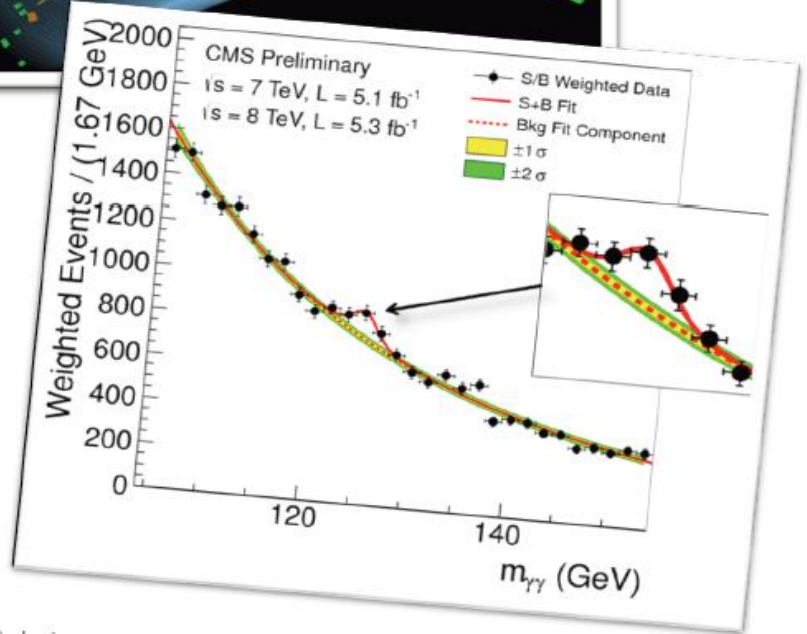
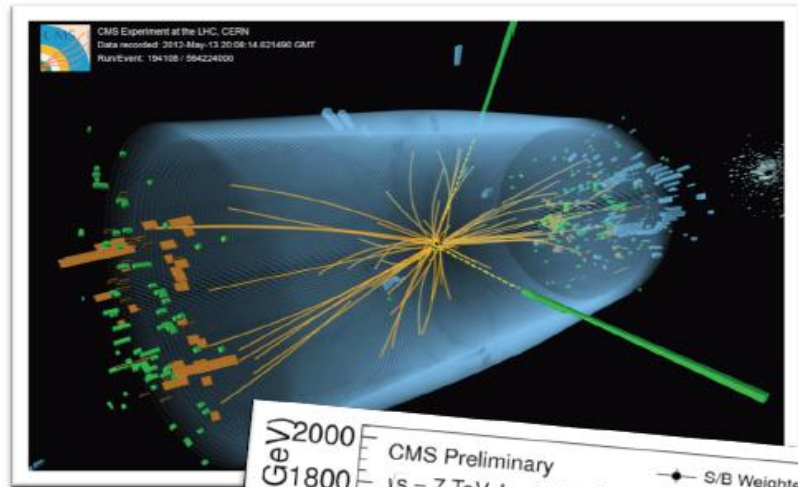
Martinus J.G. Veltman



M. Gell-Mann

# Experiment = probing theories with data

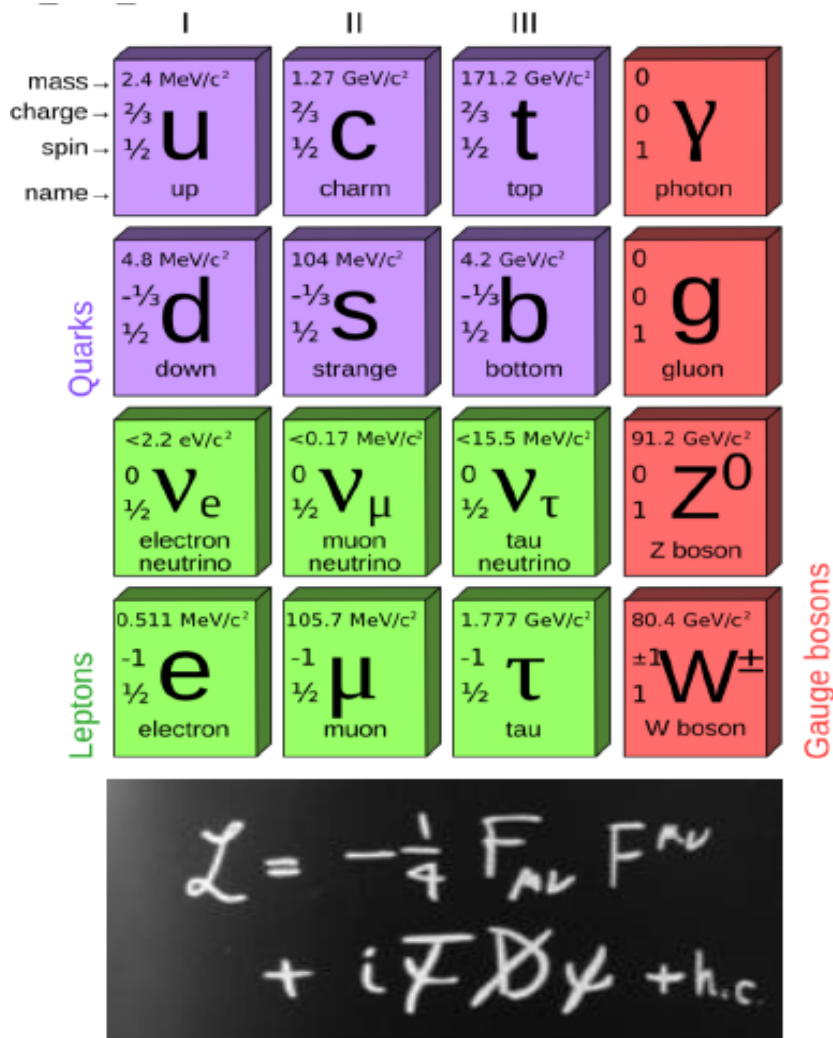
$$\begin{aligned}
 & -\frac{1}{2}g_{\phi}^2 g_{\phi}^2 g_{\phi}^2 - g_s f^{abc} \partial_{\mu} g_{\nu}^a g_{\rho}^b g_{\sigma}^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_{\mu}^a g_{\nu}^b g_{\rho}^c g_{\sigma}^d g_{\tau}^e + \\
 & \frac{1}{2}g_s^2 (g_1^{\mu\nu} g_2^{\rho\sigma}) g_{\mu\nu}^a + G^a \partial^{\mu} G^a + g_s f^{abc} \partial_{\mu} G^a G^b G^c - \partial_{\mu} W_{\nu}^+ \partial_{\rho} W_{\sigma}^{-} H - \\
 & M^2 W_{\nu}^{+} W_{\nu}^{-} - \frac{1}{2} \partial_{\nu} Z_{\mu}^0 \partial_{\nu} Z_{\mu}^0 - \frac{1}{2} M^2 Z_{\mu}^0 Z_{\mu}^0 - \frac{1}{2} \partial_{\mu} A_{\nu} \partial_{\rho} A_{\nu} - \frac{1}{2} \partial_{\mu} H \partial_{\nu} H - \\
 & \frac{1}{2} m_h^2 H^2 - \partial_{\mu} \phi^+ \partial_{\mu} \phi^{-} - M^2 \phi^+ \phi^{-} - \frac{1}{2} \partial_{\mu} \phi^0 \partial_{\mu} \phi^0 - \frac{1}{2} M^2 \phi^0 \phi^0 - \beta_h \frac{[2M^2]}{2} + \\
 & \frac{2M^2}{\Lambda} H + \frac{1}{2} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^{-}) + \frac{2M^2}{\Lambda} \alpha_h - ig_{\phi W} [\partial_{\nu} Z_{\mu}^0 (W_{\mu}^{+} W_{\nu}^{-} - \\
 & W_{\nu}^{+} W_{\mu}^{-}) - Z_{\mu}^0 (W_{\nu}^{+} \partial_{\nu} W_{\mu}^{-} - W_{\nu}^{-} \partial_{\nu} W_{\mu}^{+}) + Z_{\mu}^0 (W_{\nu}^{+} \partial_{\nu} W_{\mu}^{-} - \\
 & W_{\nu}^{-} \partial_{\nu} W_{\mu}^{+})] - ig_{\phi W} [\partial_{\nu} A_{\mu} (W_{\mu}^{+} W_{\nu}^{-} - W_{\nu}^{+} W_{\mu}^{-}) - A_{\nu} (W_{\mu}^{+} \partial_{\nu} W_{\mu}^{-} - \\
 & W_{\nu}^{-} \partial_{\nu} W_{\mu}^{+}) + A_{\mu} (W_{\nu}^{+} \partial_{\nu} W_{\mu}^{-} - W_{\nu}^{-} \partial_{\nu} W_{\mu}^{+})] - \frac{1}{2} g^2 W_{\mu}^{+} W_{\nu}^{-} W_{\nu}^{+} W_{\mu}^{-} + \\
 & \frac{1}{2} g^2 W_{\mu}^{-} W_{\nu}^{+} W_{\nu}^{-} W_{\mu}^{+} + g^2 c_w^2 (Z_{\mu}^0 W_{\nu}^{+} Z_{\nu}^0 W_{\mu}^{-} - Z_{\mu}^0 Z_{\nu}^0 W_{\nu}^{+} W_{\mu}^{-} - \\
 & g^2 s_w^2 (A_{\mu} W_{\nu}^{+} A_{\nu} W_{\mu}^{-} - A_{\mu} A_{\nu} W_{\nu}^{+} W_{\mu}^{-}) + g^2 s_w c_w [A_{\mu} Z_{\nu}^0 (W_{\mu}^{+} W_{\nu}^{-} - \\
 & W_{\nu}^{+} W_{\mu}^{-}) - 2A_{\mu} Z_{\nu}^0 (W_{\nu}^{+} W_{\mu}^{-}) - g\alpha [H^2 + H\phi^0 \phi^0 + 2(H\phi^+ \phi^{-}) - \\
 & \frac{1}{2} g^2 \alpha_h (H^4 + (\phi^0)^4 + 4(\phi^+ \phi^{-})^2 + 4(\phi^0)^2 \phi^+ \phi^{-} + 4H^2 \phi^+ \phi^{-} + 2(\phi^0)^2 H^2) - \\
 & g M W_{\mu}^{+} W_{\nu}^{-} H - \frac{1}{2} g \frac{M}{c_w} Z_{\mu}^0 Z_{\nu}^0 H - \frac{1}{2} ig W_{\mu}^{+} (\partial_{\nu} \phi^0 \phi^0 - \phi^0 \partial_{\nu} \phi^0) - \\
 & W_{\nu}^{-} (\phi^0 \partial_{\nu} \phi^+ - \phi^+ \partial_{\nu} \phi^0)] + \frac{1}{2} ig [W_{\mu}^{+} (H \partial_{\nu} \phi^0 - \phi^0 \partial_{\nu} H) - W_{\nu}^{-} (H \partial_{\nu} \phi^+ - \\
 & \phi^+ \partial_{\nu} H)] + \frac{1}{2} g \frac{1}{c_w} (Z_{\mu}^0 (H \partial_{\nu} \phi^0 - \phi^0 \partial_{\nu} H) - ig_{\phi W}^2 M Z_{\mu}^0 (W_{\mu}^{+} \phi^{-} - W_{\nu}^{-} \phi^{+}) + \\
 & ig_{\phi W} M A_{\mu} (W_{\mu}^{+} \phi^{-} - W_{\nu}^{-} \phi^{+}) - ig \frac{1-2c_w^2}{2c_w} Z_{\mu}^0 (\phi^+ \partial_{\nu} \phi^0 - \phi^0 \partial_{\nu} \phi^+) - \\
 & ig_{\phi W} A_{\mu} (\phi^+ \partial_{\nu} \phi^0 - \phi^0 \partial_{\nu} \phi^+) - \frac{1}{2} g^2 W_{\mu}^{+} W_{\nu}^{-} [H^2 + (\phi^0)^2 + 2\phi^+ \phi^{-}] - \\
 & \frac{1}{4} g^2 \frac{1}{c_w} Z_{\mu}^0 Z_{\nu}^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^{-}] - \frac{1}{2} g^2 \frac{2c_w}{c_w} Z_{\mu}^0 \phi^0 (W_{\mu}^{+} \phi^{-} + \\
 & W_{\nu}^{-} \phi^{+}) - \frac{1}{2} ig^2 \frac{2c_w}{c_w} Z_{\mu}^0 H (W_{\mu}^{+} \phi^{-} - W_{\nu}^{-} \phi^{+}) + \frac{1}{2} g^2 s_w A_{\mu} \phi^0 (W_{\mu}^{+} \phi^{-} + \\
 & W_{\nu}^{-} \phi^{+}) + \frac{1}{2} ig^2 s_w A_{\mu} H (W_{\mu}^{+} \phi^{-} - W_{\nu}^{-} \phi^{+}) - g^2 \frac{2c_w}{c_w} (2c_w^2 - 1) Z_{\mu}^0 A_{\nu} \phi^+ \phi^{-} - \\
 & g^1 s_w^2 A_{\mu} A_{\nu} \phi^+ \phi^{-} - e^{\lambda} (\gamma \partial + m_h^2) e^{\lambda} - \rho^{\lambda} \gamma \partial \nu^{\lambda} - u_3^{\lambda} (\gamma \partial + m_h^2) u_3^{\lambda} + \\
 & d_3^{\lambda} (\gamma \partial + m_h^2) d_3^{\lambda} + ig_{\phi W} A_{\mu} [-(e^{\lambda} \gamma^{\mu} e^{\lambda}) + \frac{2}{3} (u_3^{\lambda} \gamma^{\mu} u_3^{\lambda}) - \frac{1}{3} (d_3^{\lambda} \gamma^{\mu} d_3^{\lambda})] + \\
 & \frac{ig}{4c_w} Z_{\mu}^0 [(\nu^{\lambda} \gamma^{\mu} (1 + \gamma^5) \nu^{\lambda}) + (e^{\lambda} \gamma^{\mu} (4s_w^2 - 1 - \gamma^5) e^{\lambda}) + (u_3^{\lambda} \gamma^{\mu} (\frac{2}{3} s_w^2 - \\
 & 1 - \gamma^5) u_3^{\lambda}) + (d_3^{\lambda} \gamma^{\mu} (1 - \frac{2}{3} s_w^2 - \gamma^5) d_3^{\lambda})] + \frac{ig}{2\sqrt{2}} W_{\mu}^{+} [(e^{\lambda} \gamma^{\mu} (1 + \gamma^5) \nu^{\lambda}) + \\
 & (u_3^{\lambda} \gamma^{\mu} (1 + \gamma^5) C_{3\lambda} d_3^{\lambda})] + \frac{ig}{2\sqrt{2}} W_{\mu}^{-} [(e^{\lambda} \gamma^{\mu} (1 + \gamma^5) \nu^{\lambda}) + (d_3^{\lambda} \gamma^{\mu} C_{3\lambda}^{\lambda} (1 + \\
 & \gamma^5) u_3^{\lambda})] + \frac{ig}{2\sqrt{2}} \frac{m_h^2}{M} [-\phi^+ (\nu^{\lambda} (1 - \gamma^5) e^{\lambda}) + \phi^{-} (e^{\lambda} (1 + \gamma^5) \nu^{\lambda}) - \\
 & \frac{2}{3} \frac{m_h^2}{M} [H (e^{\lambda} e^{\lambda}) + i\phi^0 (e^{\lambda} \gamma^5 e^{\lambda})] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_h^2 (u_3^{\lambda} C_{3\lambda} (1 - \gamma^5) d_3^{\lambda}) + \\
 & m_h^2 (d_3^{\lambda} C_{3\lambda}^{\lambda} (1 + \gamma^5) u_3^{\lambda}) - m_h^2 (d_3^{\lambda} C_{3\lambda}^{\lambda} (1 - \\
 & \gamma^5) u_3^{\lambda}) - \frac{2}{3} \frac{m_h^2}{M} H (u_3^{\lambda} u_3^{\lambda}) - \frac{2}{3} \frac{m_h^2}{M} H (d_3^{\lambda} d_3^{\lambda}) + \frac{ig}{2} \phi^0 (\bar{u}_3^{\lambda} \gamma^0 u_3^{\lambda}) - \\
 & \frac{ig}{2} \frac{m_h^2}{M} \phi^0 (\bar{d}_3^{\lambda} \gamma^0 d_3^{\lambda}) + \bar{X}^+ (\partial^{\mu} - M^2) X^+ + \bar{X}^- (\partial^{\mu} - M^2) X^- + \bar{X}^0 (\partial^{\mu} - \\
 & \frac{M^2}{c_w} X^0 + \bar{Y} \partial^{\mu} Y + ig_{\phi W} W_{\mu}^{+} (\partial_{\nu} \bar{X}^0 X^- - \partial_{\nu} \bar{X}^+ X^0) + ig_{\phi W} W_{\mu}^{-} (\partial_{\nu} \bar{X}^+ X^- - \\
 & \partial_{\nu} \bar{X}^0 X^0) + ig_{\phi W} W_{\mu}^{-} (\partial_{\nu} \bar{X}^- X^0 - \partial_{\nu} \bar{X}^0 X^+) + ig_{\phi W} W_{\mu}^{-} (\partial_{\nu} \bar{X}^- X^- - \\
 & \partial_{\nu} \bar{X}^+ X^+) + ig_{\phi W} Z_{\mu}^0 (\partial_{\nu} \bar{X}^+ X^+ - \partial_{\nu} \bar{X}^- X^-) + ig_{\phi W} A_{\mu} (\partial_{\nu} \bar{X}^+ X^+ - \\
 & \partial_{\nu} \bar{X}^- X^-) - \frac{1}{2} ig M [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\
 & \frac{1-2c_w^2}{2c_w} ig M [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} ig M [\bar{X}^0 X^- \phi^+ - \bar{X}^+ X^+ \phi^-] + \\
 & ig M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^+ X^+ \phi^-] + \frac{1}{2} ig M [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
 \end{aligned}$$



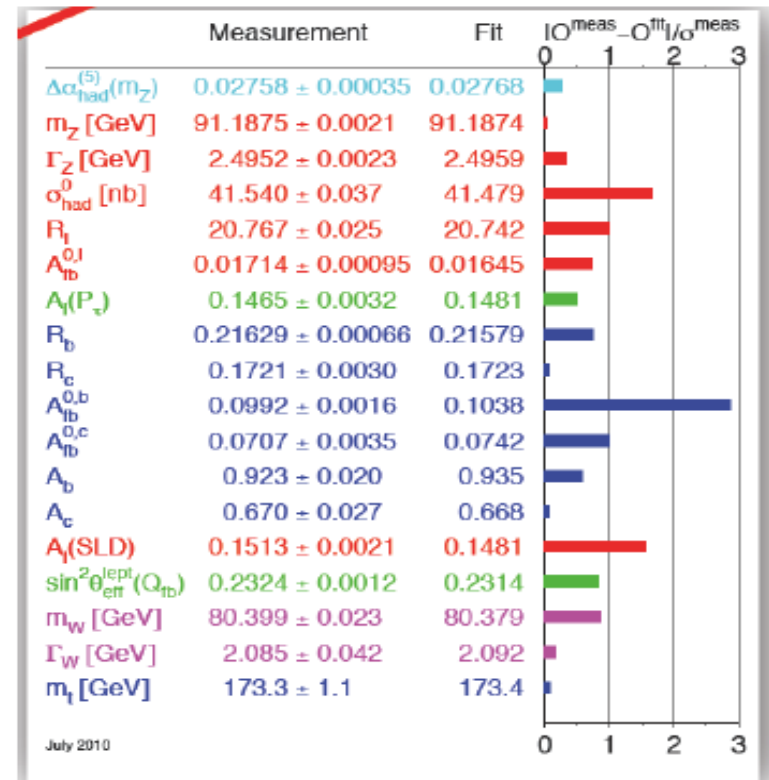
prof. E. Richter-Wąs

Wstęp do Fizyki Cząstek Wysokich Energii

# Standard Model confirmed by the data



## STANDARD MODEL OF ELEMENTARY PARTICLES



Confirmed at sub 1% level!

# HEP, SI and „natural” units

Quantity	HEP units	SI units
length	1 fm	$10^{-15}$ m
charge	e	$1.602 \cdot 10^{-19}$ C
energy	1 GeV	$1.602 \times 10^{-10}$ J
mass	1 GeV/c <sup>2</sup>	$1.78 \times 10^{-27}$ kg
$\hbar = h/2\pi$	$6.588 \times 10^{-25}$ GeV s	$1.055 \times 10^{-34}$ Js
c	$2.988 \times 10^{23}$ fm/s	$2.988 \times 10^8$ m/s
$\hbar c$	197 MeV fm	...

## “natural” units ( $\hbar = c = 1$ )

mass	1 GeV
length	1 GeV <sup>-1</sup> = 0.1973 fm
time	1 GeV <sup>-1</sup> = $6.59 \times 10^{-25}$ s

# Measuring particles

- Particles are characterized by
  - ✓ **Mass** [Unit: eV/c<sup>2</sup> or eV]
  - ✓ **Charge** [Unit: e]
  - ✓ **Energy** [Unit: eV]
  - ✓ **Momentum** [Unit: eV/c or eV]
  - ✓ (+ spin, lifetime, ...)

Particle identification via measurement of:

e.g. (E, p, Q) or (p, β, Q)  
(p, m, Q) ...

- ... and move at **relativistic speed**

$$\beta = \frac{v}{c} \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

$$l = \frac{l_0}{\gamma} \quad \text{length contraction}$$

$$t = t_0 \gamma \quad \text{time dilatation}$$

$$E^2 = \vec{p}^2 c^2 + m^2 c^4$$

$$E = m \gamma c^2 = m c^2 + E_{\text{kin}}$$

$$\vec{\beta} = \frac{\vec{p}c}{E} \quad \vec{p} = m \gamma \vec{\beta} c$$

# Relativistic kinematics

$$\begin{aligned} E^2 &= \vec{p}^2 + m^2 \\ \ell &= \frac{\ell_0}{\gamma} & E &= m\gamma \\ t &= t_0\gamma & \vec{p} &= m\gamma\vec{\beta} \\ & & \vec{\beta} &= \frac{\vec{p}}{E} \end{aligned}$$

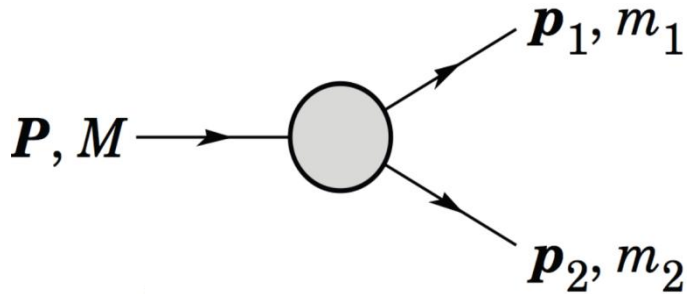
## Center of mass energy

- In the **center of mass frame** the total momentum is 0
- In **laboratory frame** center of mass energy can be computed as:

$$E_{\text{cm}} = \sqrt{s} = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$

# Kinematics

## 2-bodies decays

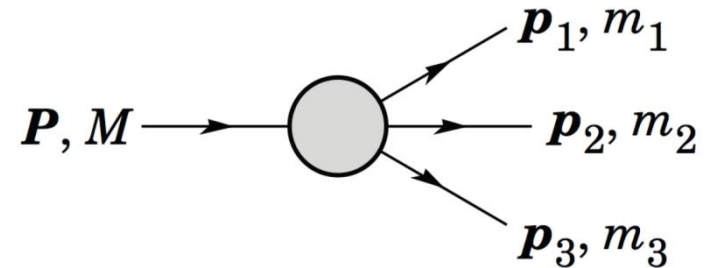


$$E_1 = \frac{M^2 - m_2^2 + m_1^2}{2M}$$

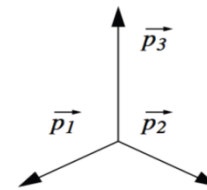
$$|\mathbf{p}_1| = |\mathbf{p}_2|$$

$$= \frac{[(M^2 - (m_1 + m_2)^2)(M^2 - (m_1 - m_2)^2)]^{1/2}}{2M}$$

## 3-bodies decays



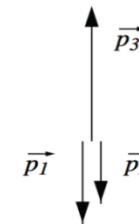
$$|\mathbf{p}_3| = \frac{[(M^2 - (m_{12} + m_3)^2)(M^2 - (m_{12} - m_3)^2)]^{1/2}}{2M}$$



(a)

$$\max(|\vec{p}_3|)$$

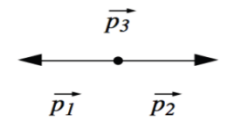
$$\min(|\vec{p}_3|)$$



(b)

$$(m_{12})_{min} = m_1 + m_2$$

$$(m_{12})_{max} = M - m_3$$



(c)

## Invariant mass

$$M = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$



# A real example: pion decays

pion decays at rest

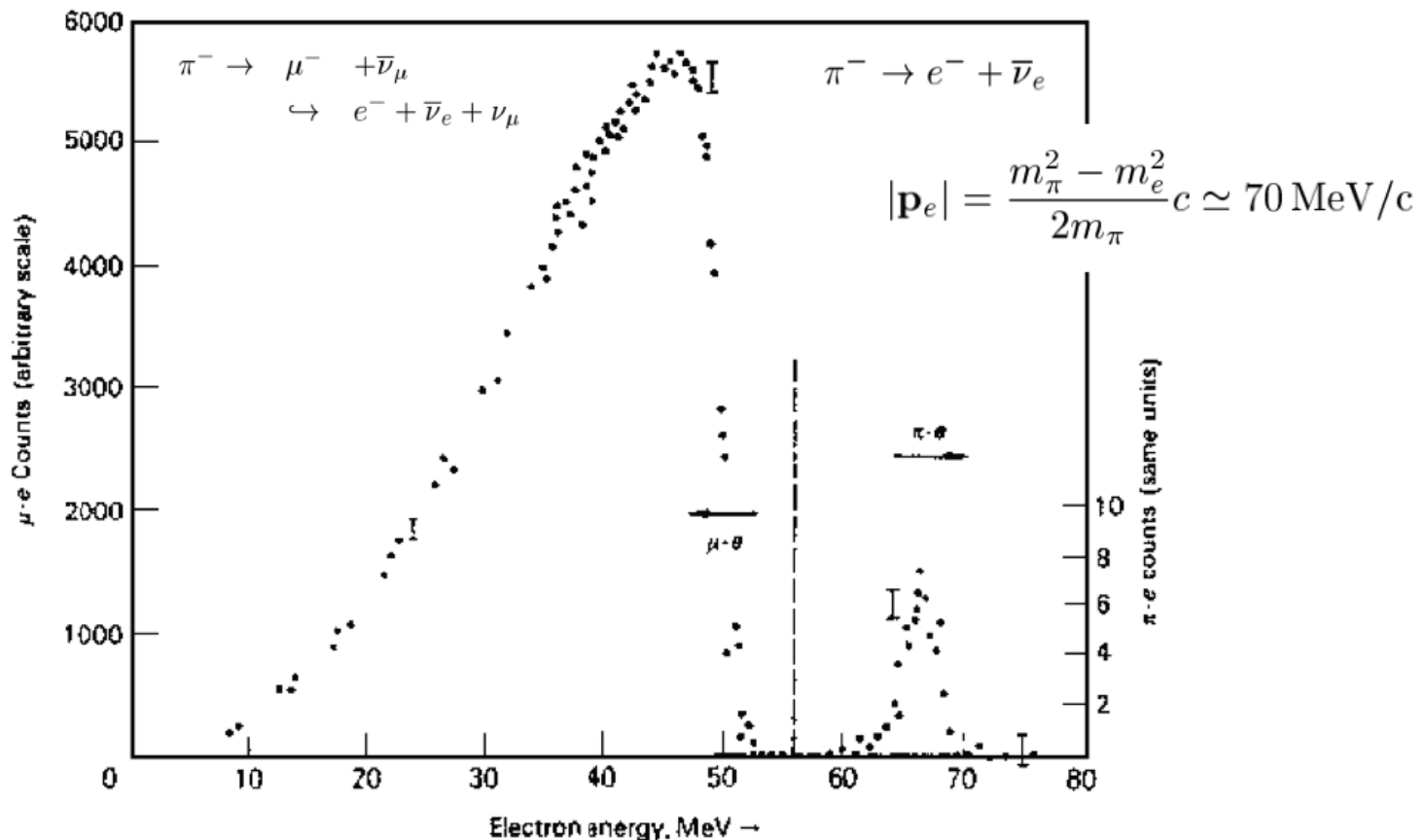
$$|\mathbf{p}_\mu| = \frac{m_\pi^2 - m_\mu^2}{2m_\pi} c \simeq 30 \text{ MeV}/c$$

$m_\nu = 0$ .

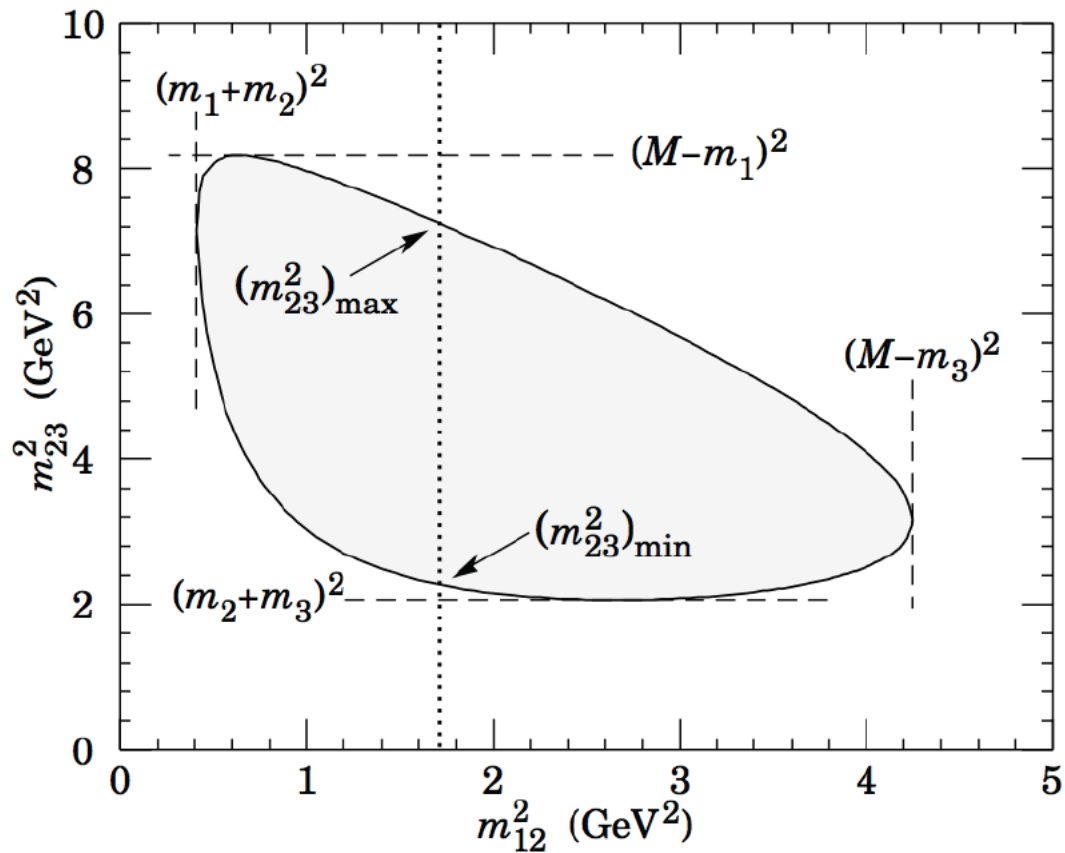
in most cases  
muon decays  
at rest

$$|\mathbf{p}_e|_{max} = \frac{m_\mu^2 - m_e^2}{2m_\mu} c \simeq 52 \text{ MeV}/c$$

$$|\mathbf{p}_e|_{min} = 0$$

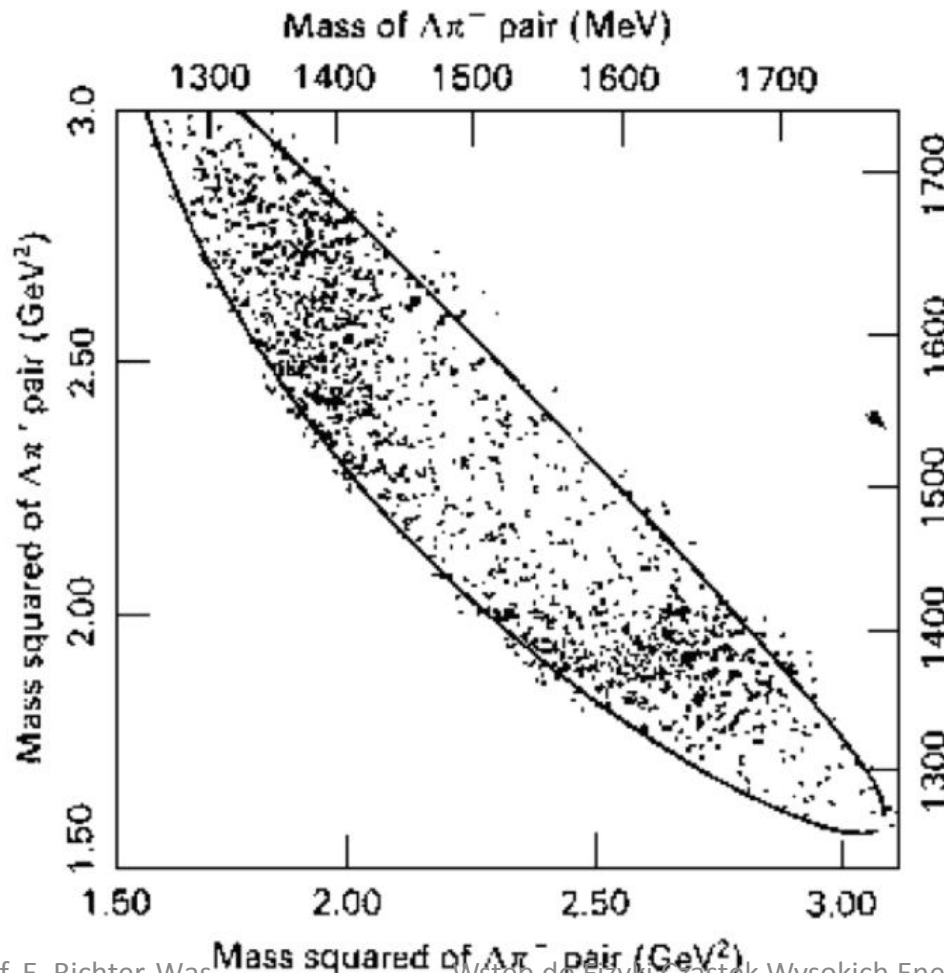
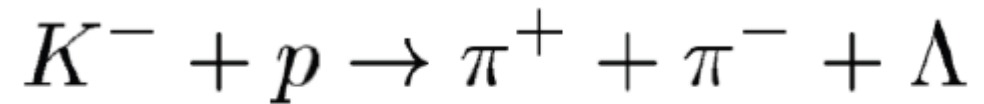


# 3-bodies decay: Dalitz plot

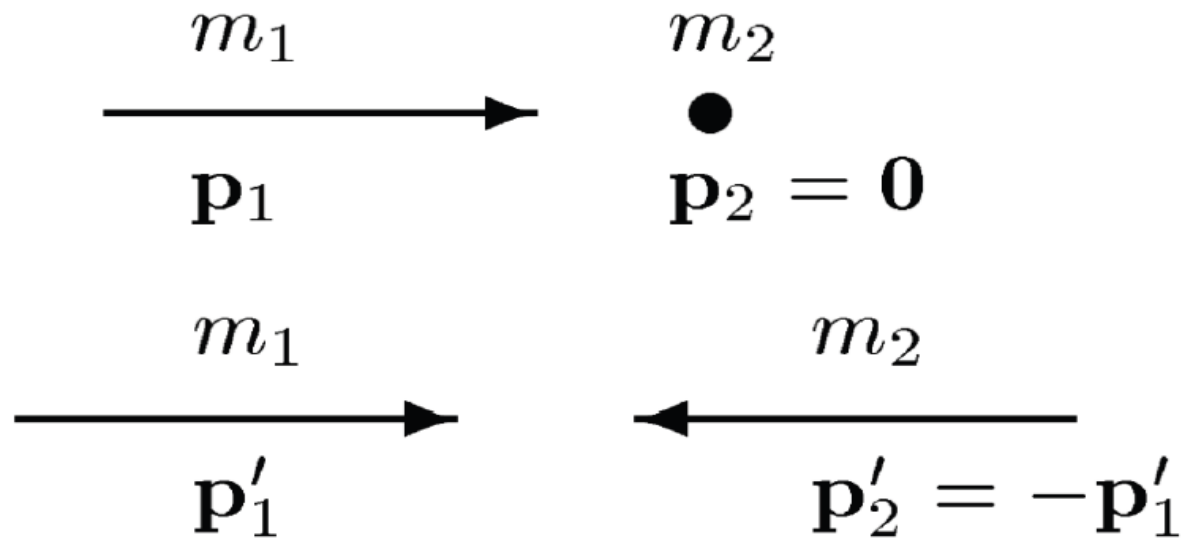


**Figure 45.3:** Dalitz plot for a three-body final state. In this example, the state is  $\pi^+ \bar{K}^0 p$  at 3 GeV. Four-momentum conservation restricts events to the shaded region.

# Multi-bodies decay



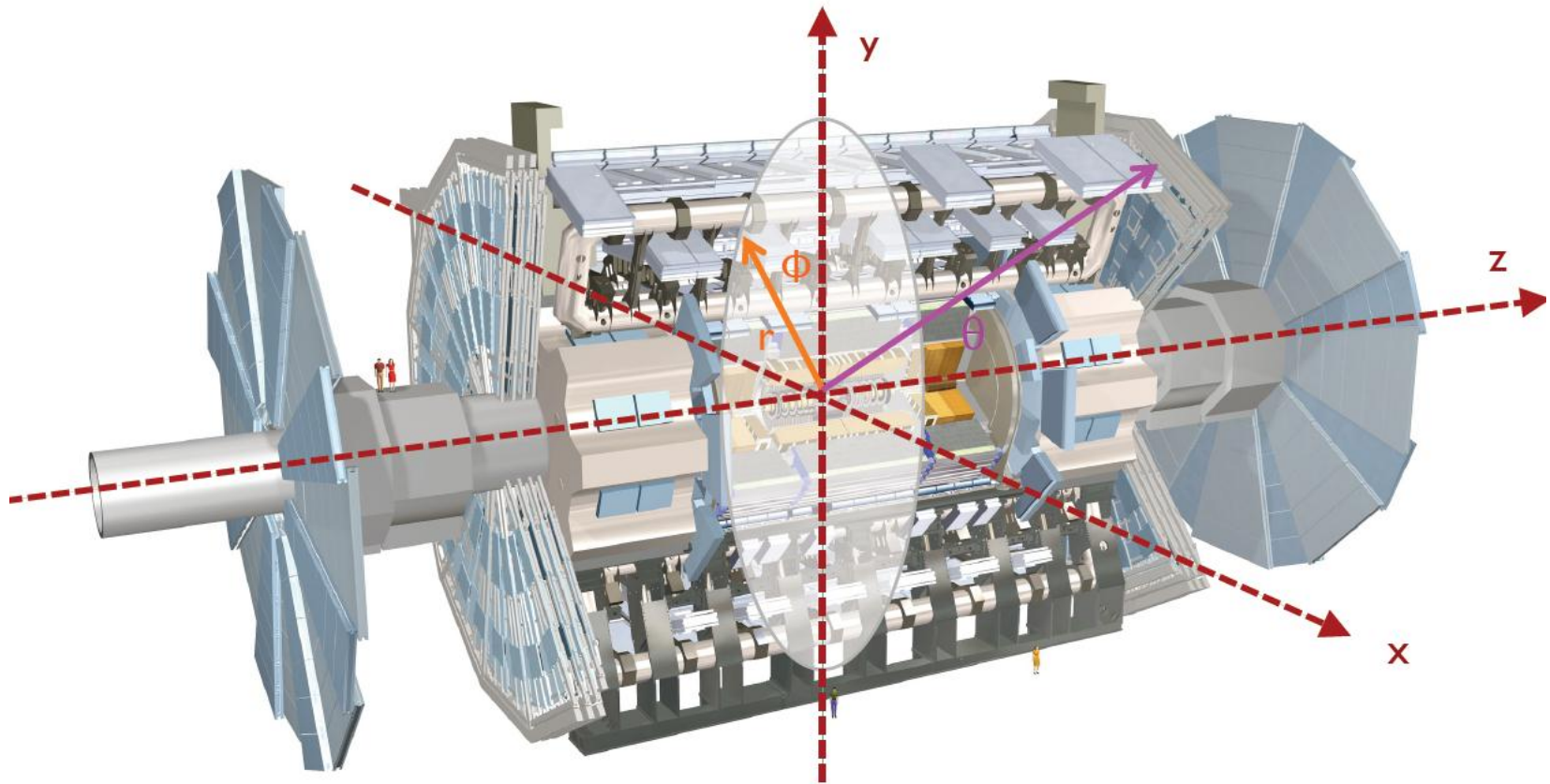
# Fixed target vs collider



How much energy should a fixed target experiment have to equal the center of mass energy of two colliding beam?

$$E_{\text{fix}} = 2 \frac{E_{\text{col}}^2}{m} - m$$

# Collider experiment coordinates



# Rapidity

Lorentz factor  $\gamma = \frac{1}{\sqrt{1 - \beta^2}} = \cosh \varphi$  Hyperbolic cosine of “rapidity”

$$\begin{aligned} E &= m \cosh \varphi \\ |\vec{p}| &= m \sinh \varphi \end{aligned} \quad \varphi = \tanh^{-1} \frac{E}{|\vec{p}|} = \frac{1}{2} \ln \frac{E + |\vec{p}|}{E - |\vec{p}|}$$

- Particle physicists prefer to use modified rapidity along beam axis

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

# Pseudorapidity

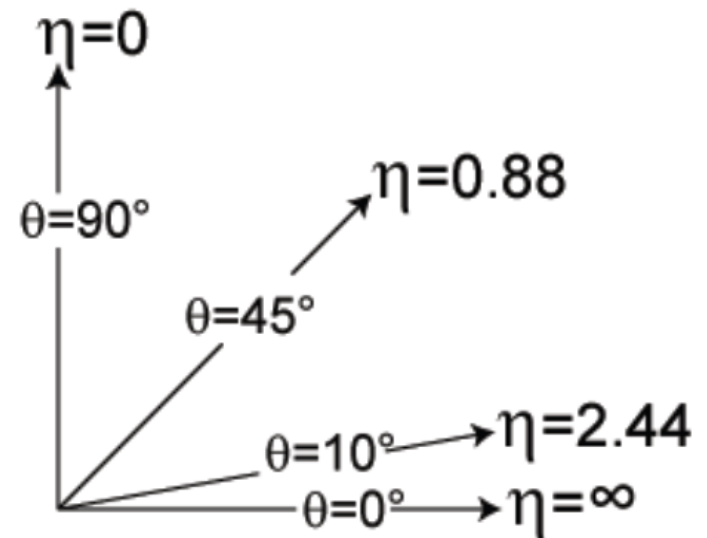
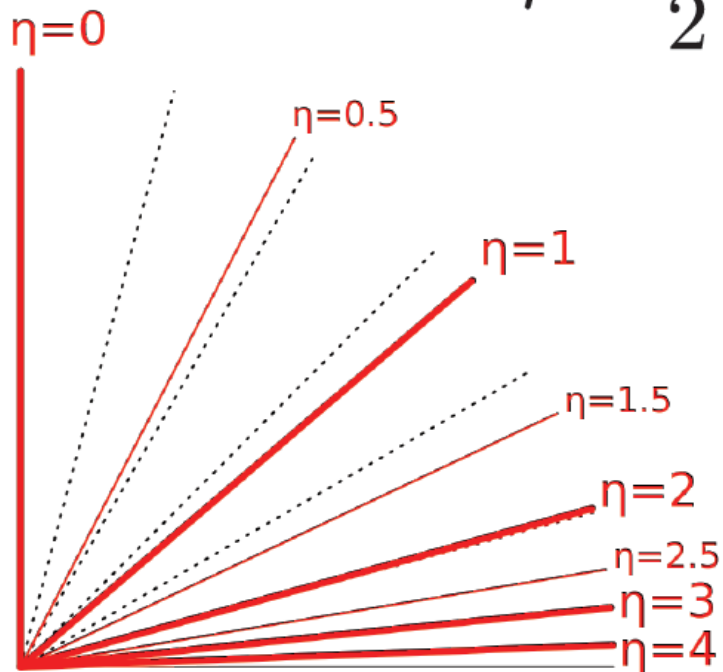
$$\eta = \frac{1}{2} \ln \frac{|\vec{p}| + p_z}{|\vec{p}| - p_z}$$

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

$$\eta \simeq y$$

if  $E \gg m$

$$\eta = \frac{1}{2} \ln \left( \tan \frac{\theta}{2} \right)$$



# Transverse variables

- At hadron colliders, a significant and unknown fraction of the beam energy in each event escapes down the beam pipe.
- Net momentum can only be constrained in the plane transverse to the beam z-axis!

$$p_T = \sqrt{p_x^2 + p_y^2} \quad \begin{aligned} p_x &= p_T \cos \phi \\ p_y &= p_T \sin \phi \\ p_z &= p_T \sinh \eta \end{aligned} \quad \begin{aligned} |p| &= p_T \cosh \eta \\ E_T &= \frac{E}{\cosh \eta} \end{aligned}$$

$$\sum p_x(i) = 0 \quad \sum p_y(i) = 0$$



# Missing transverse energy and transverse mass

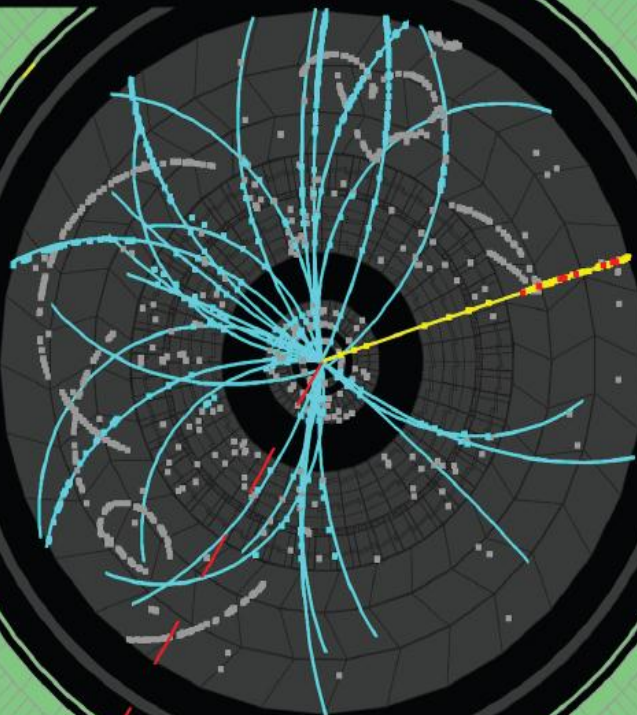
- If invisible particles are created, only their transverse momentum can be constrained: **missing transverse energy**

$$E_T^{\text{miss}} = \sum p_T(i)$$

- If a heavy particle is produced and decays into two particles one of which is invisible, the mass of the parent particle can be constrained with the **transverse mass quantity**

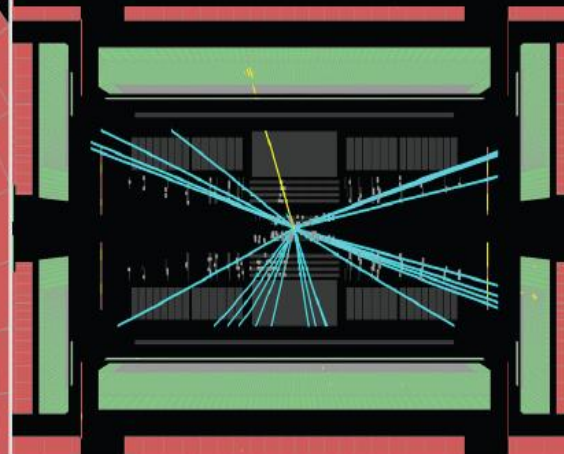
$$\begin{aligned} M_T^2 &\equiv [E_T(1) + E_T(2)]^2 - [\mathbf{p}_T(1) + \mathbf{p}_T(2)]^2 \\ &= m_1^2 + m_2^2 + 2[E_T(1)E_T(2) - \mathbf{p}_T(1) \cdot \mathbf{p}_T(2)] \end{aligned}$$

$$\text{if } m_1 = m_2 = 0 \quad M_T^2 = 2|\mathbf{p}_T(1)||\mathbf{p}_T(2)|(1 - \cos \phi_{12})$$



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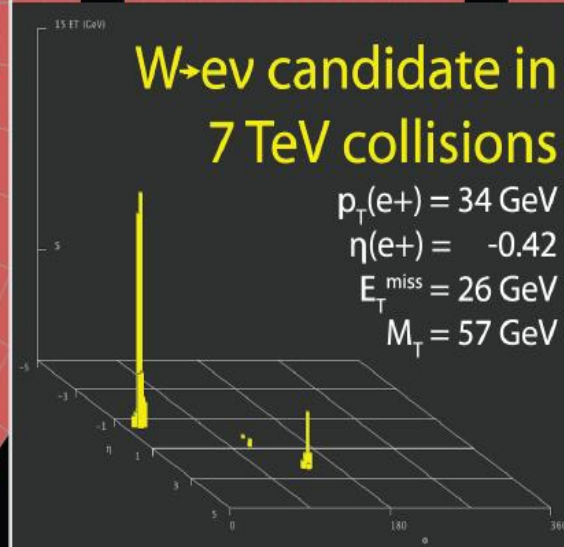
## W $\rightarrow$ ev candidate in 7 TeV collisions

$$p_T(e^+) = 34 \text{ GeV}$$

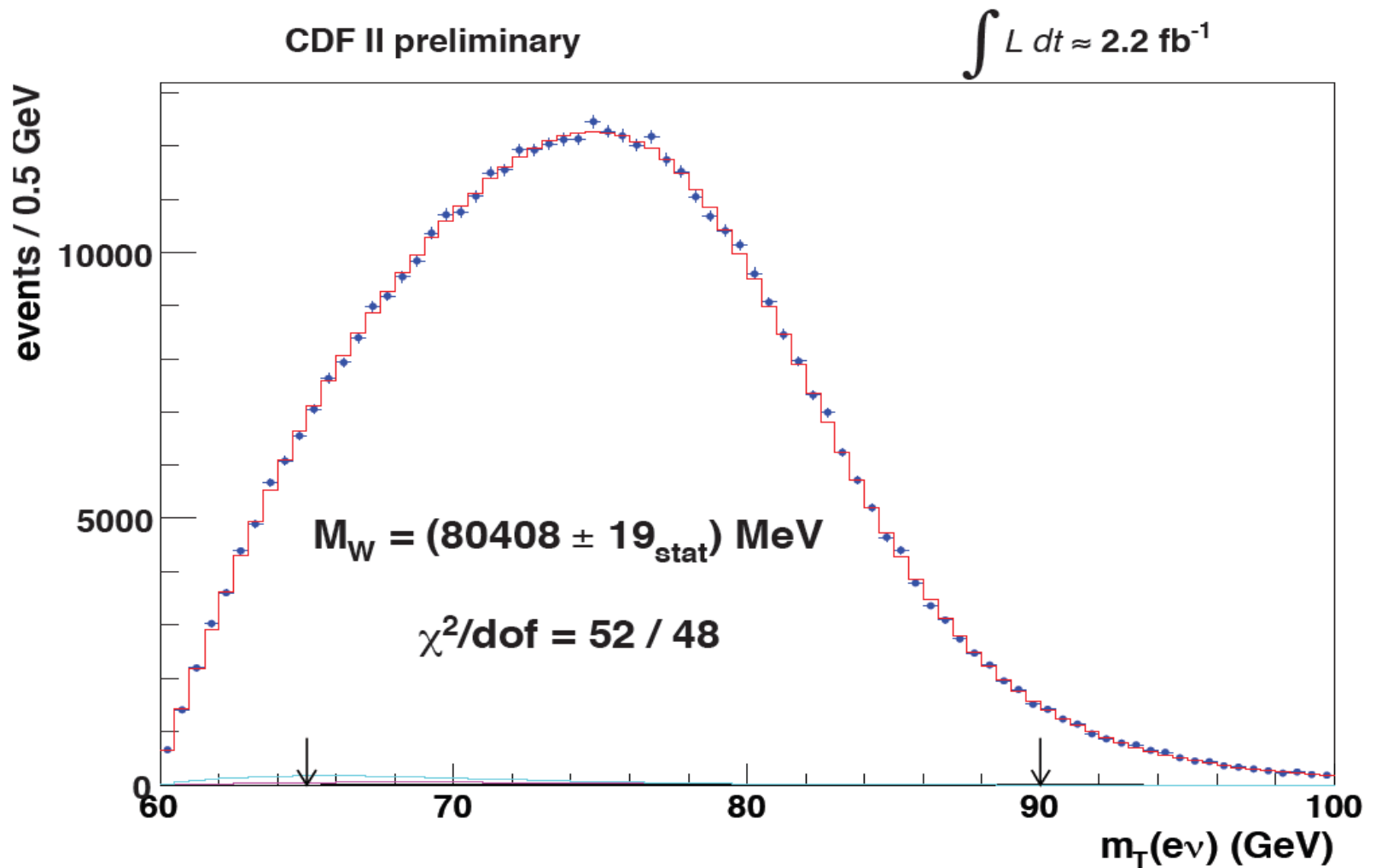
$$\eta(e^+) = -0.42$$

$$E_T^{\text{miss}} = 26 \text{ GeV}$$

$$M_T = 57 \text{ GeV}$$

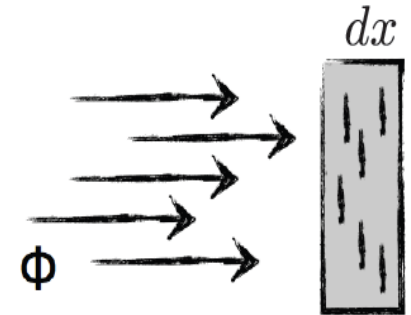


# Mass of the W boson



# Interaction cross-section

Flux  $\Phi = \frac{1}{S} \frac{dN_i}{dt}$   $[L^{-2} t^{-1}]$



Reactions per unit of time  $\frac{dN_{\text{reac}}}{dt} = \Phi \overbrace{\sigma N_{\text{target}} dx}^{\text{area obscured by target particle}}$   $[t^{-1}]$

$[L^{-2} t^{-1}]$   $[?]$   $[L^{-1}]$   $[L]$

Reaction rate per target particle  $W_{if} = \Phi \sigma$   $[t^{-1}]$

Cross section per target particle  $\sigma = \frac{W_{if}}{\Phi}$   $[L^2]$  = reaction rate per unit of flux

$1b = 10^{-28} \text{ m}^2$  (roughly the area of a nucleus with  $A = 100$ )

# Fermi Golden rule

From non-relativistic perturbation theory...

transition probability      matrix element      energy density of final states

$$W_{if} = \frac{2\pi}{\hbar} |M_{if}|^2 \frac{dN}{dE_f}$$

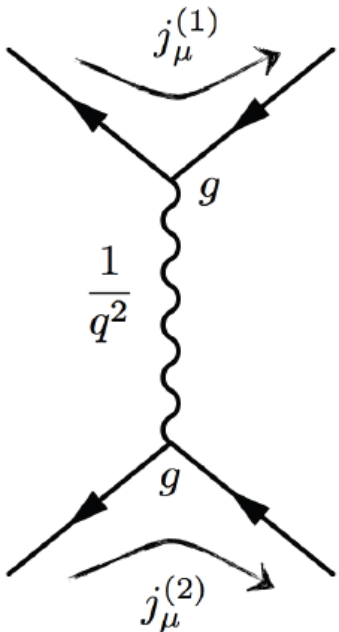
[t<sup>-1</sup>]

[E]

[E<sup>-1</sup>]

$$M_{if} = -i \int j_\mu^{(1)} \left( \frac{1}{q^2} \right) j_\mu^{(2)} d^4x$$

$$\sigma \sim |M_{if}|^2 \sim g^4 \left( \frac{1}{q^4} \right)$$



# Cross-section: magnitude and units

Standard

cross section unit:

$$[\sigma] = \text{mb}$$

with  $1 \text{ mb} = 10^{-27} \text{ cm}^2$

or in

natural units:

$$[\sigma] = \text{GeV}^{-2}$$

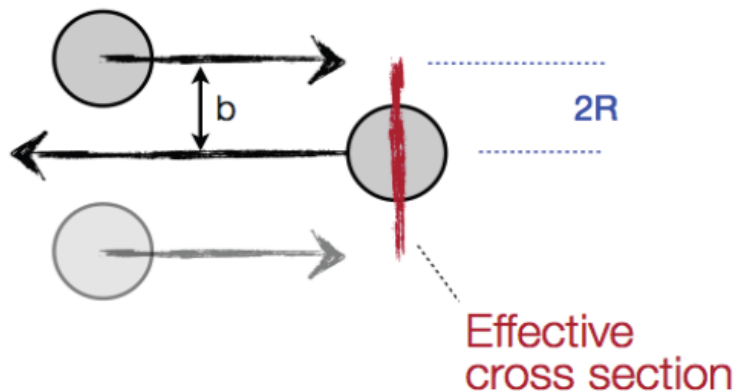
with  $1 \text{ GeV}^{-2} = 0.389 \text{ mb}$

$$1 \text{ mb} = 2.57 \text{ GeV}^{-2}$$

Estimating the  
proton-proton cross section:

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using:  $\hbar c = 0.1973 \text{ GeV fm}$   
 $(\hbar c)^2 = 0.389 \text{ GeV}^2 \text{ mb}$

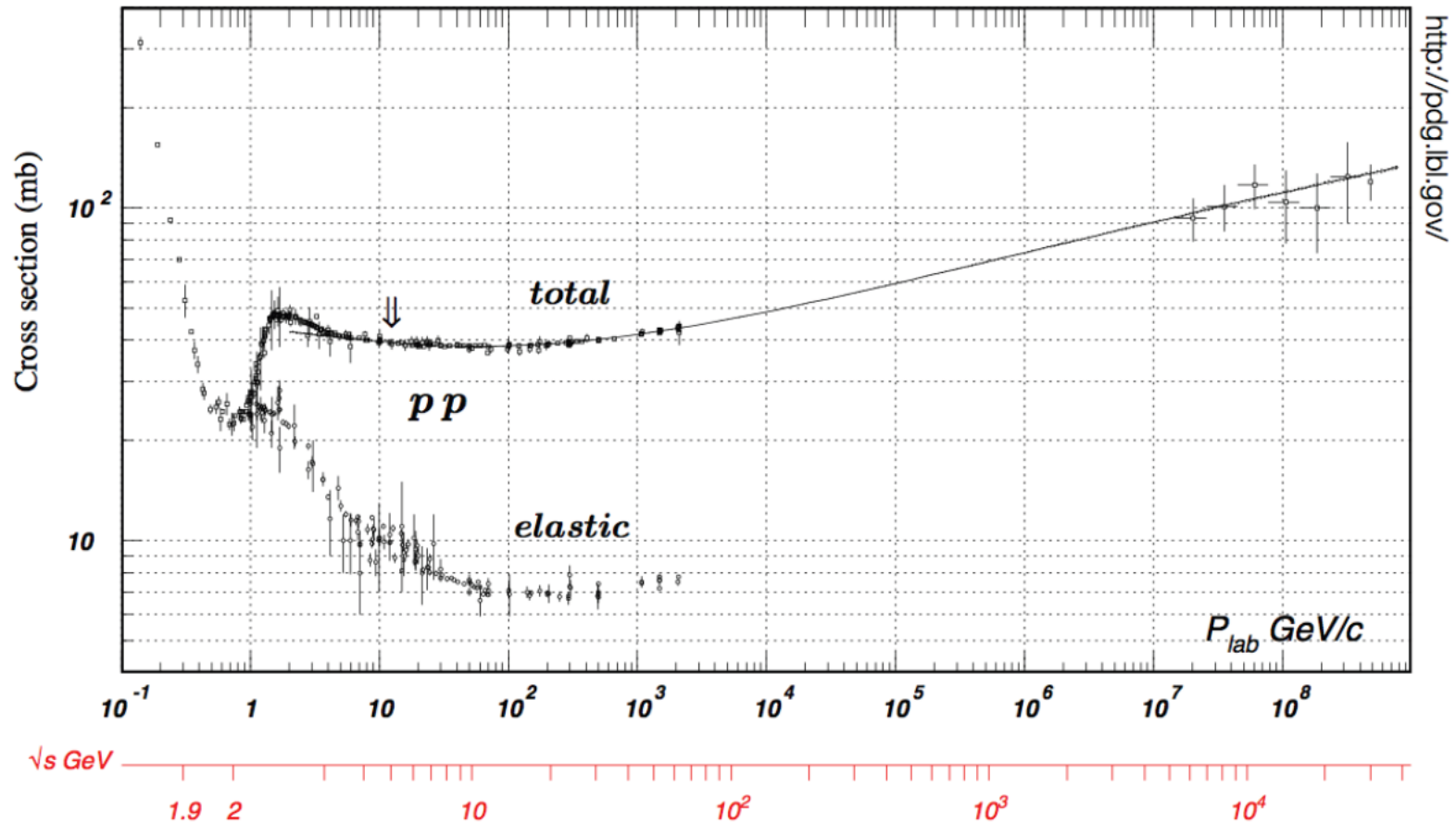


Proton radius:  $R = 0.8 \text{ fm}$

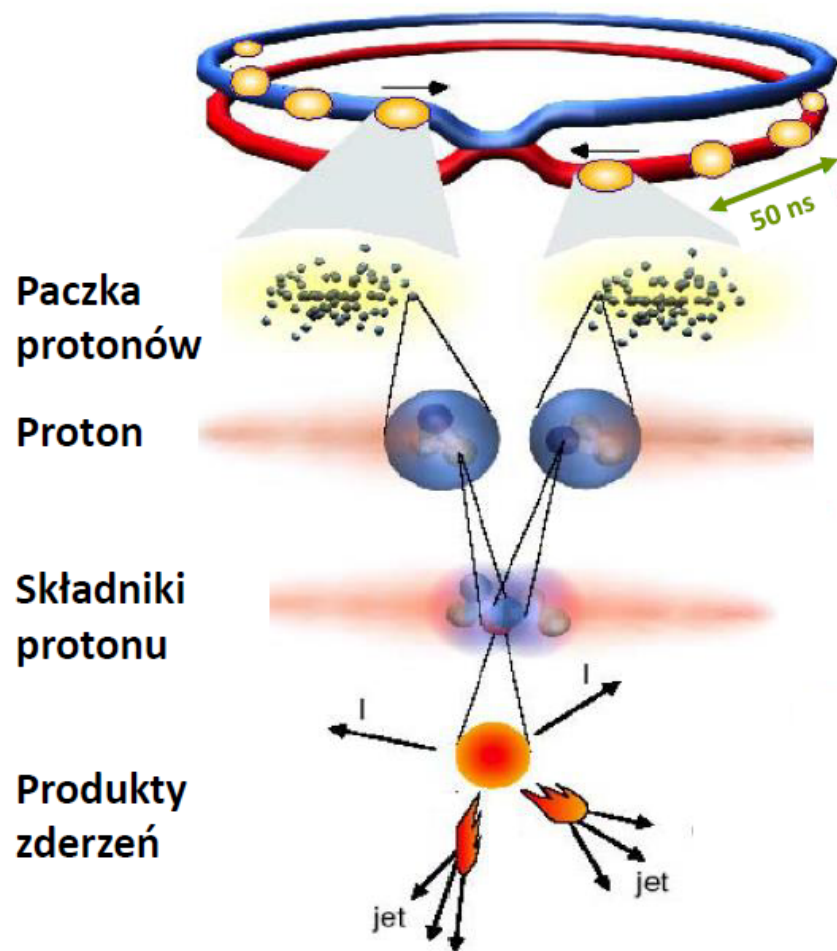
Strong interactions happens up to  $b = 2R$

$$\begin{aligned} \sigma &= \pi (2R)^2 = \pi \cdot 1.6^2 \text{ fm}^2 \\ &= \pi \cdot 1.6^2 \cdot 10^{-26} \text{ cm}^2 \\ &= \pi \cdot 1.6^2 \cdot 10 \text{ mb} \\ &= 80 \text{ mb} \end{aligned}$$

# Proton-proton scattering cross-section



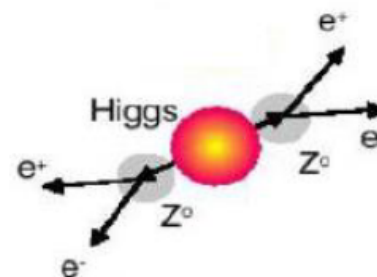
# Proton-proton collisions at LHC



<b>Proton-Proton</b>	1380 paczek/wiązkę
Protonów/paczka	$1.7 \cdot 10^{11}$
Energia wiązki	3.5, 4.0, 6.5 TeV

Każdy proton porusza się z prędkością bliską prędkości światła i niesie kinetyczną energię muchy w locie, okrąża pierścień akceleratora 1100 razy na sekundę.

Rozmiar poprzeczny wiązki:  $16 \mu\text{m}$  (4 razy mniejszy niż grubość ludzkiego włosa).  
Każda z wiązek niesie energię pociągu TGV o dł. 200 m i jadącego z prędkością 155km/godz (360M Jula).



**Takie zdarzenie pojawia się raz na 10 bilionów zderzeń**



# Cross-sections at LHC

