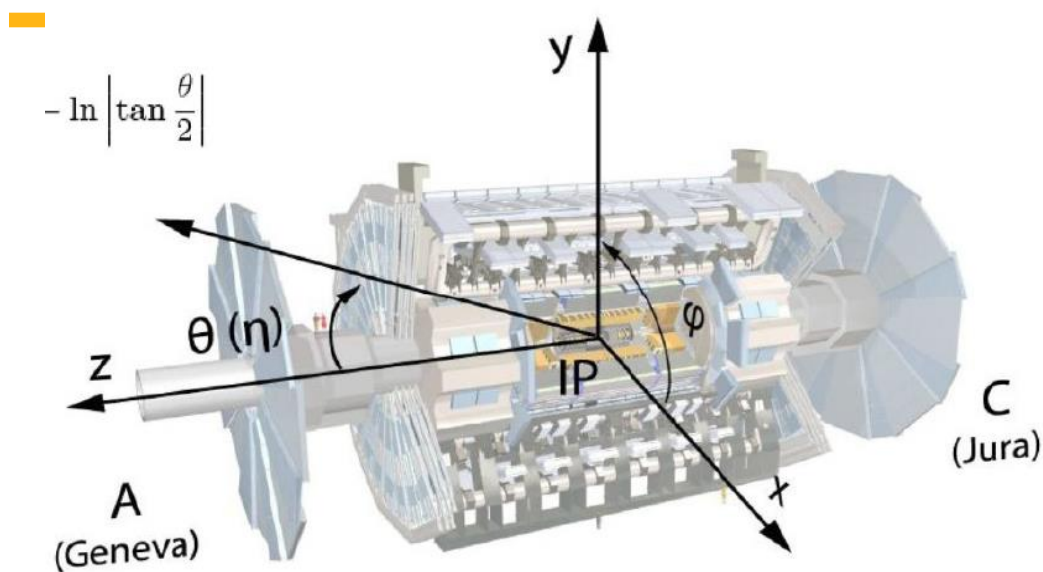


ATLAS Detector

THE ATLAS DETECTOR IS REALLY BIG!

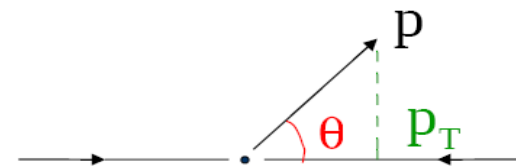


- Length : ~ 46 m
- Radius : ~ 12 m
- Weight : ~ 7000 tons
- $\sim 10^8$ electronic channels
- 3000 km of cables

Transverse momentum

(in the plane perpendicular to the beam)

$$p_T = p \sin\theta$$



Rapidity: $\eta = -\log(\operatorname{tg} \frac{\theta}{2})$

$$\theta = 90^\circ \rightarrow \eta = 0$$

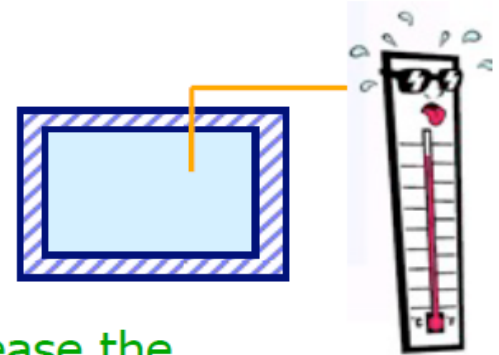
$$\theta = 10^\circ \rightarrow \eta \cong 2.4$$

$$\theta = 170^\circ \rightarrow \eta \cong -2.4$$

What is calorimeter

Concept comes from thermo-dynamics:

A leak-proof closed box containing a substance which temperature is to be measured.



Temperature scale:

1 calorie (4.185J) is the necessary energy to increase the temperature of 1 g of water at 15°C by one degree

At hadron colliders we measure GeV (0.1 - 1000)

1 GeV = 10^9 eV $\approx 10^9 * 10^{-19}$ J = 10^{-10} J = $2.4 * 10^{-9}$ cal

1 TeV = 1000 GeV : kinetic energy of a flying mosquito

Required sensitivity for our calorimeters is
~ a thousand million time larger than
to measure the increase of temperature by 1°C of 1 g of water

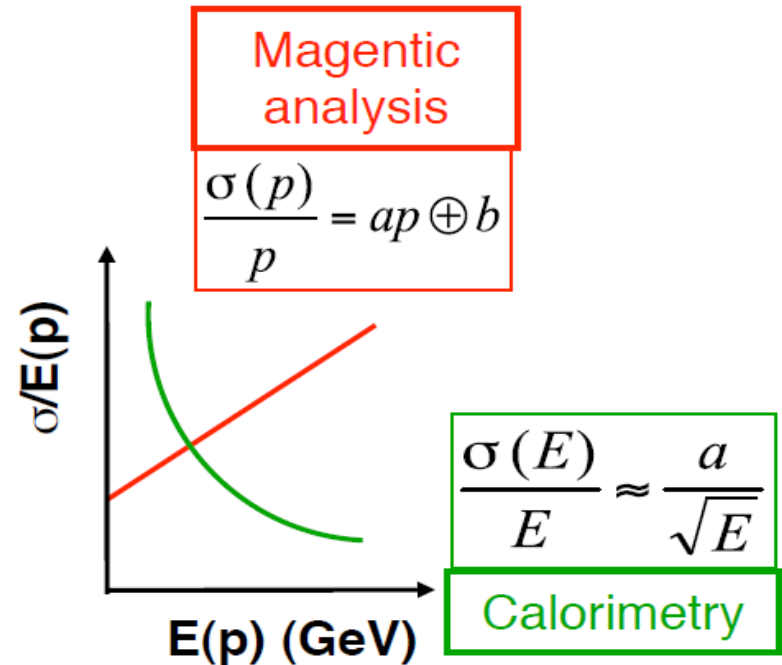
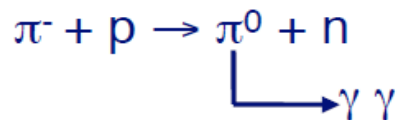
Why calorimeters

First calorimeters appeared in the 70's:

need to measure the energy of all particles, **charged** and **neutral**.

Until then, only the momentum of **charged particles** was measured using **magnetic analysis**.

The measurement with a calorimeter is destructive e.g.

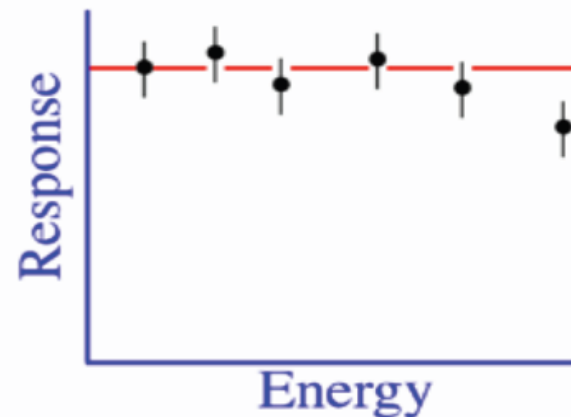
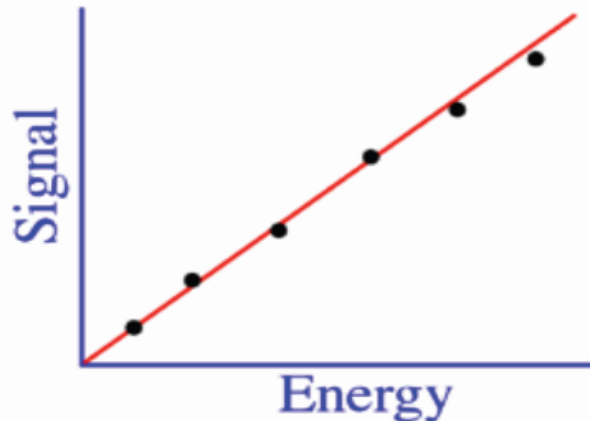


Particles do not come out alive of a calorimeter

Important characteristic: linearity

Response: mean signal per unit of deposited energy
e.g. # of photons electrons/GeV, pC/MeV, $\mu\text{A}/\text{GeV}$

→ A linear calorimeter has a constant response



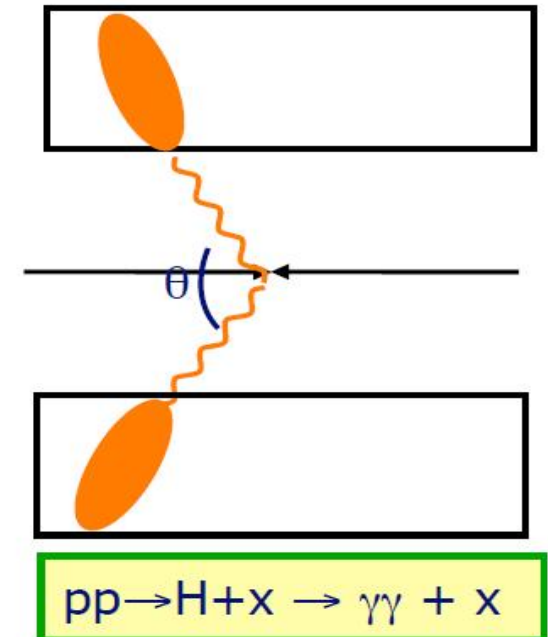
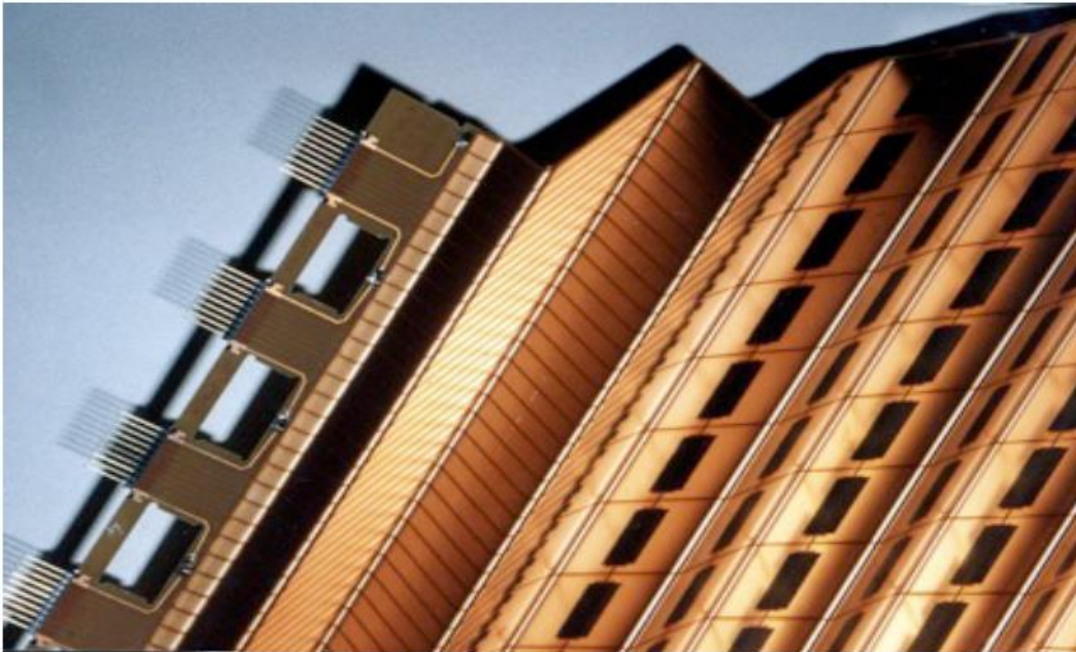
Electromagnetic calorimeters are in general linear.
All energies are deposited via ionisation/excitation of the absorber.

Important characteristic: position resolution

Higgs Boson search in ATLAS

if $M_H \sim 120$ GeV search in channel $H \rightarrow \gamma\gamma$

$$\sigma(M_H) / M_H = \frac{1}{2} [\sigma(E_{\gamma 1})/E_{\gamma 1} \oplus \sigma(E_{\gamma 2})/E_{\gamma 2} \oplus \cot(\theta/2) \sigma(\theta)]$$

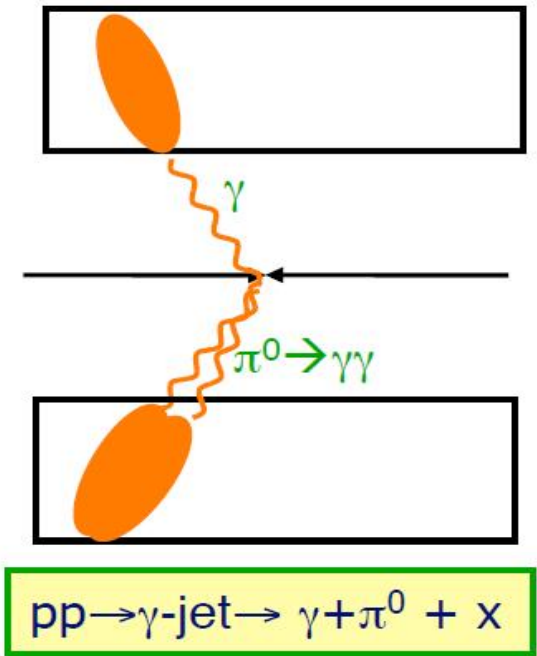


Important characteristic: particle identification

Higgs boson search in ATLAS

if $M_H \sim 120$ GeV search in channel $H \rightarrow \gamma\gamma$

Background: π^0 looking like a γ



General characteristic



Calorimeters have the following properties:

Sensitive to charged and neutral particles

Precision improves with Energy (opposite to magnetic measurements)

No need of magnetic field

Containment varies as $\ln(E)$: compact

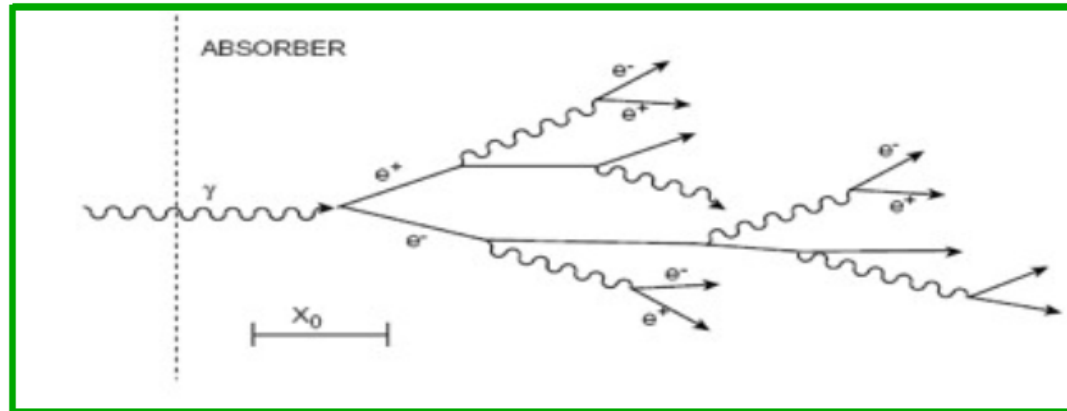
Segmentation: position measurement and identification

Fast response

Triggering capabilities

Electromagnetic shower

Electromagnetic showers result from electrons and photons undergoing **bremsstrahlung** and **pair creation**



For high energy (GeV scale) **electrons** **bremsstrahlung** is the dominant energy loss mechanism

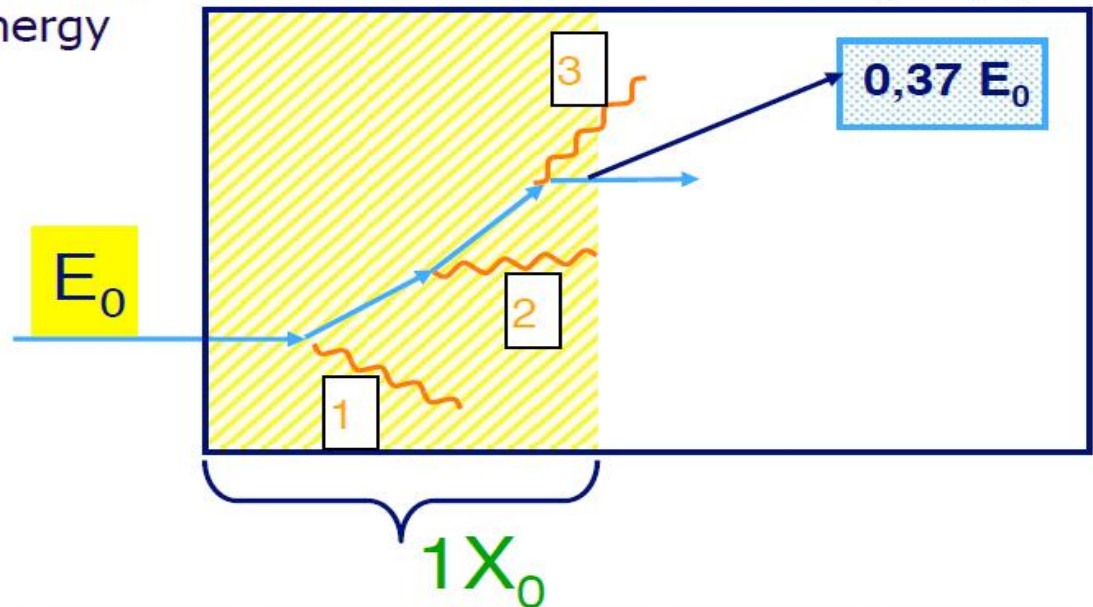
For high energy **photons** **pair creation** is the dominant absorption mechanism

Shower development is governed by these processes

Radiation length

The radiation length is a “universal” distance, very useful to describe electromagnetic showers (electrons & photons)

X_0 is the distance after which the incident electron has radiated $(1-1/e)$ 63% of its incident energy



	Air	Eau	Al	LAr	Fe	Pb	PbWO ₄
Z	-	-	13	18	26	82	-
X_0 (cm)	30420	36	8,9	14	1,76	0.56	0.89

Development of EM shower

The shower develops as a **cascade** by **energy transfer** from the incident particle to a **multitude of particles** (e^\pm and γ).

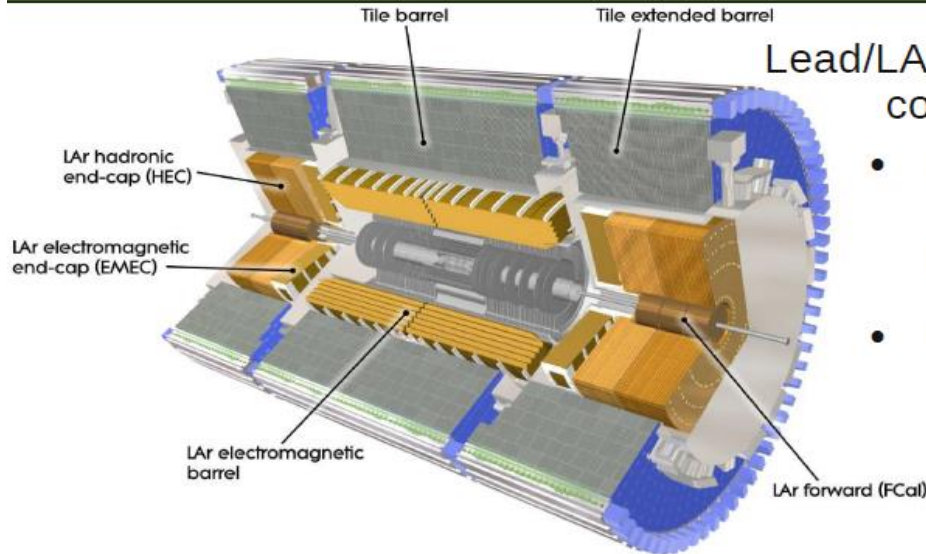
The **number of cascade particles** is **proportional** to the **energy deposited** by the incident particle

The role of the calorimeter is to **count** these cascade particles

The relative occurrence of the various processes briefly described is a function of the material (Z)

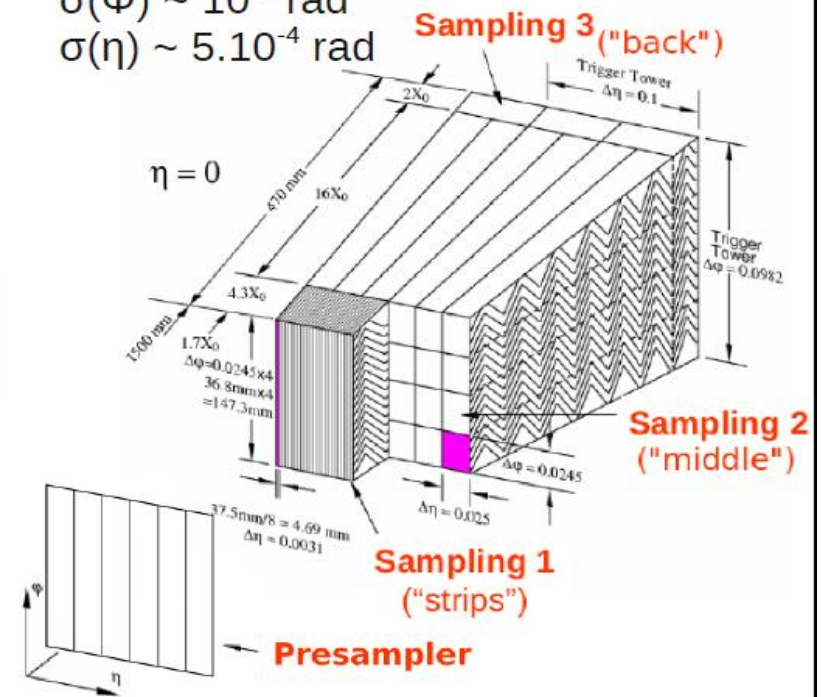
The radiation length (X_0) allows to universally describe the shower development

ATLAS electromagnetic calorimeter



Lead/LAr EM calorimeter divided in 3 longitudinal compartments + Pre-sampler in front

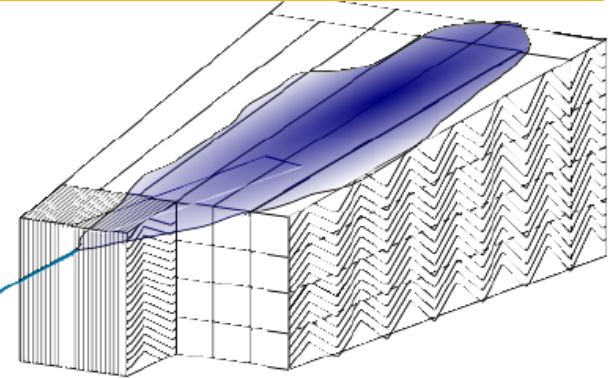
- Good energy resolution :
 $\sigma(E)/E = a/E \oplus b/\sqrt{E} \oplus c$ (with $a \sim 0.3$ GeV, $b \sim 10\%$, $c \sim 0.7\%$)
- Good angular resolution :
 $\sigma(\Phi) \sim 10^{-3}$ rad
 $\sigma(\eta) \sim 5 \cdot 10^{-4}$ rad



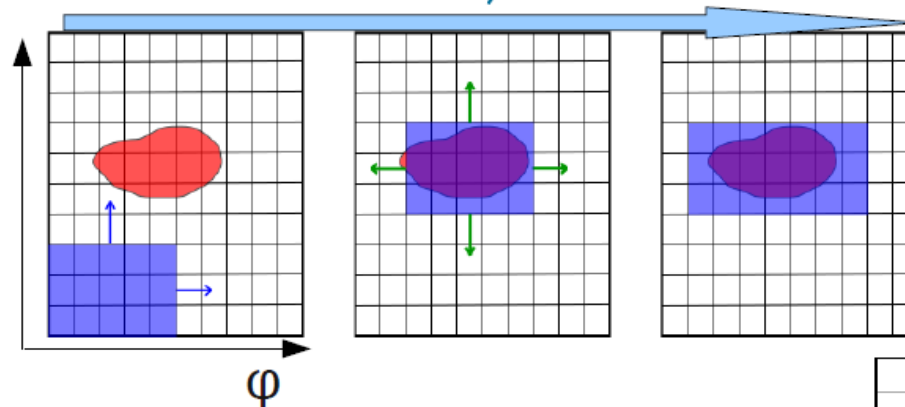
Layer	Granularity $\Delta\eta \times \Delta\phi$	Radiation length
Pre-sampler	0.025 x 0.1	
Strips	0.003 x 0.1	4.3 X_0
Middle	0.025 x 0.025	16 X_0
Back	0.05 x 0.025	2 X_0

ATLAS electromagnetic object

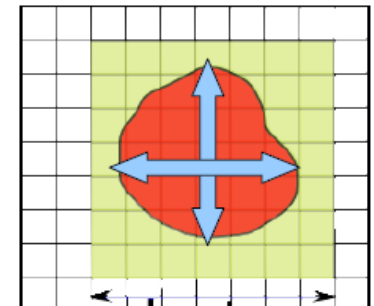
- In ATLAS an electron or a photon candidate is defined as a cluster of cells in the calorimeters representing the energy deposit to which we can associate tracks reconstructed in the inner detector



- Sliding window algorithm to reconstruct the energy deposits : η

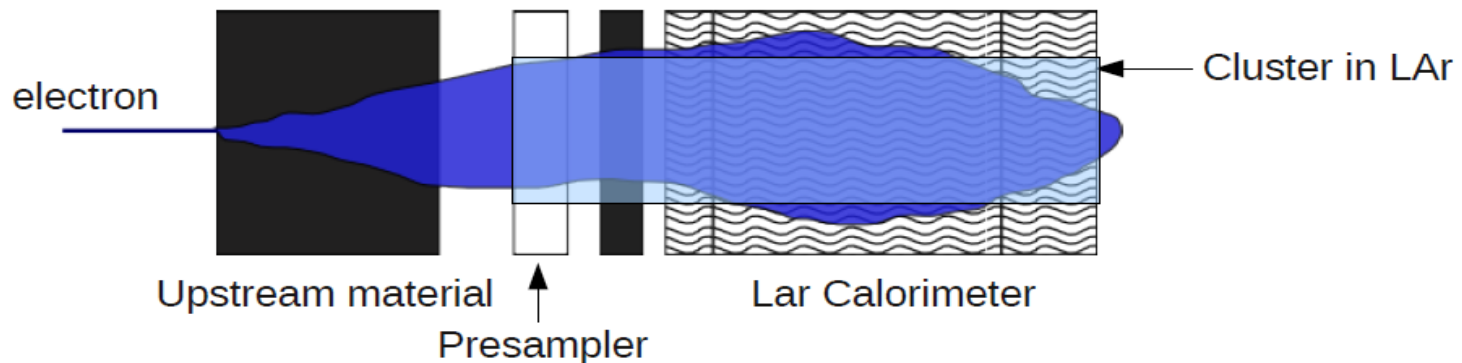


- The identification of such objects is then based on :
 - The shower shape in the calorimeter
 - Track quality (number of hits, direction wrt the cluster,...)
 - Transition radiation (TRT "high threshold hits")



Energy calibration

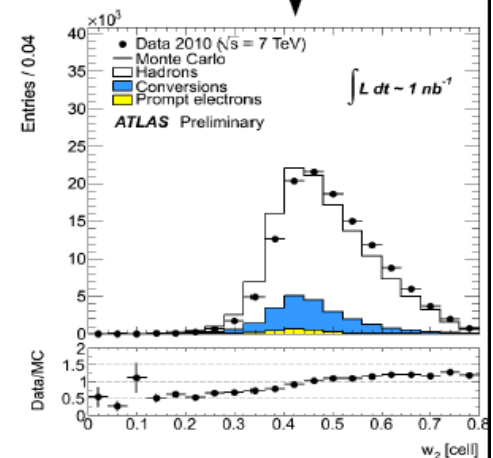
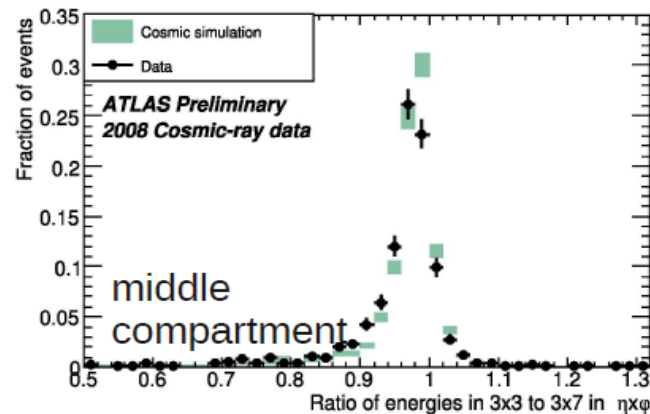
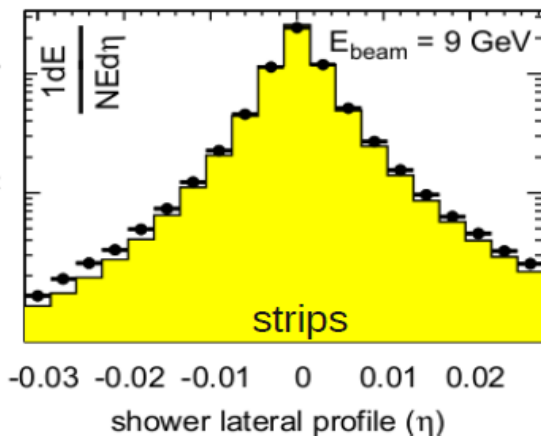
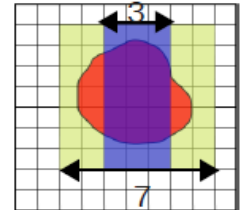
- One reason why the simulation is sensitive to the knowledge of material is the energy calibration scheme
- As the initial energy does not fully deposit within the electron/photon cluster, it is important to correct the cells energy sum to improve the energy scale and resolution
- Our calibration procedure is based on calibration hits
 - Store all GEANT4 energy deposits (in active, inactive material or escaping)
 - Parametrize the energy leaks (outside the cluster, in the dead material,...) in function of the position, the energy and the shower depth using this simulation

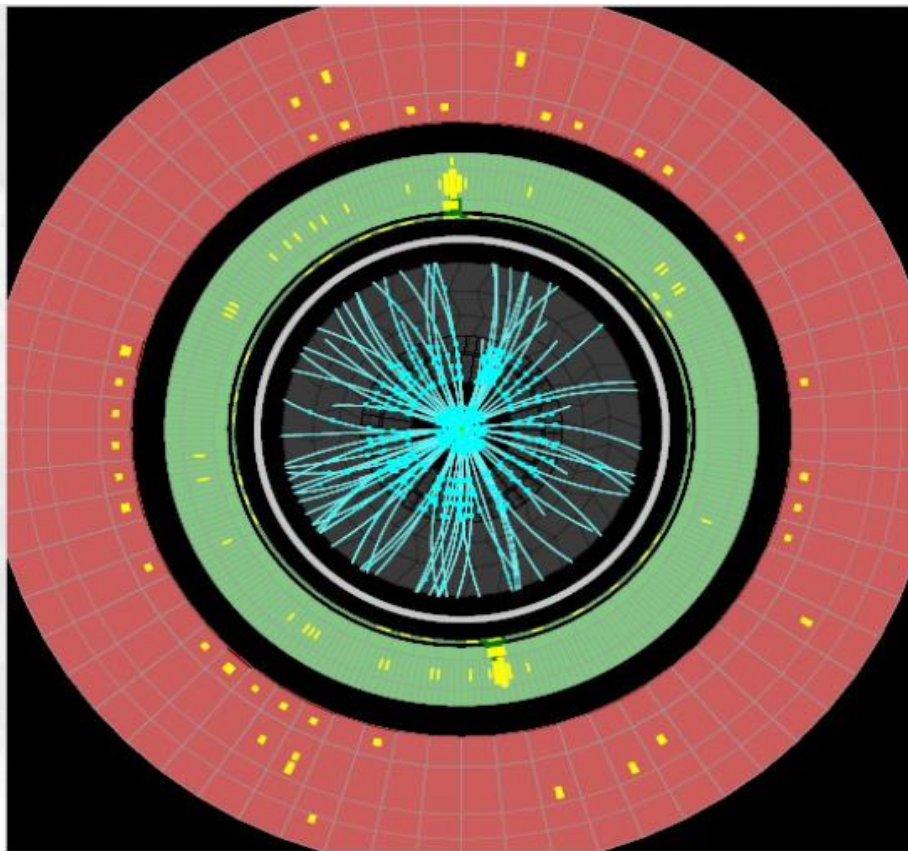


- Of course this calibration is strongly dependent on the knowledge of the upstream material, this is why we need to map it

Shower shape variables

- The shower shape in the calorimeter allows for the rejection of a large fraction of background ($O(1000)$)
- Benefiting from the thin granularity and the segmentation of the calorimeter, ATLAS defined a few variables illustrating the shower width in eta/phi and its longitudinal extension
- Even if the agreement is fairly good, the simulation does not perfectly predict the key distributions for the lateral development
- This has been observed during the test beam, the cosmics, and the collisions





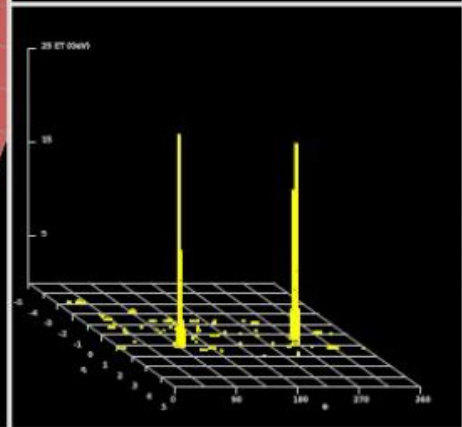

ATLAS
EXPERIMENT
 Run Number: 203779, Event Number: 56662314
 Date: 2012-05-23 22:19:29 CEST

$$\sqrt{s} = 8 \text{ TeV.}$$

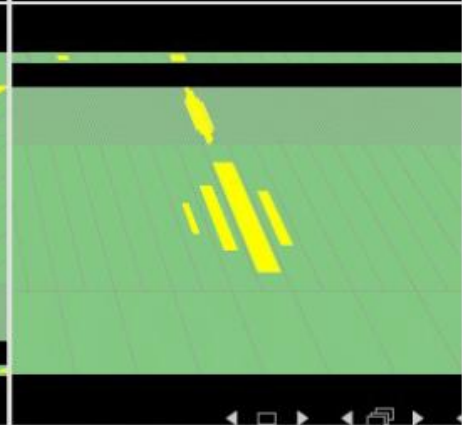
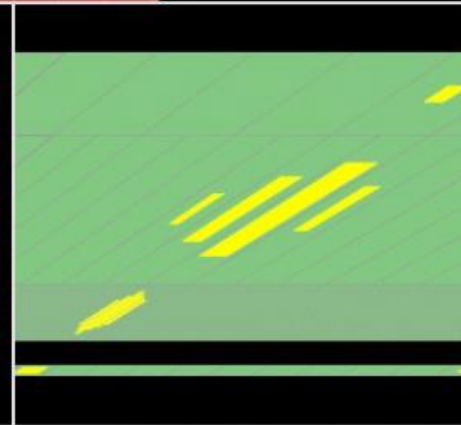
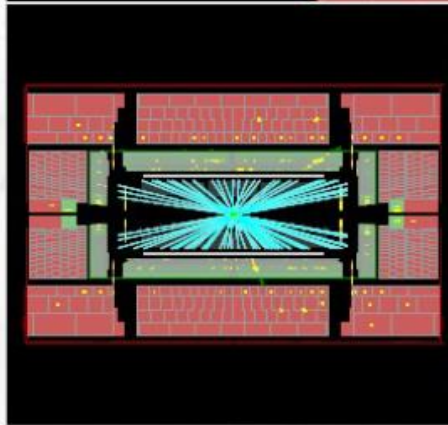
Leading γ :
 $E_T = 62.2 \text{ GeV,}$
 $\eta = 0.39$

Subleading γ :
 $E_T = 55.5 \text{ GeV}$
 $\eta = 1.18$

$$m_{\gamma\gamma} = 126.9 \text{ GeV.}$$

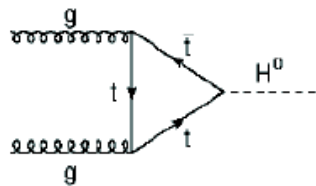


Only reconstructed tracks with $P_T > 1 \text{ GeV}$, hits in the pixel and SCT layers and TRT hits with a high threshold are shown.

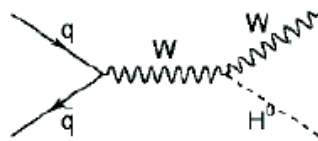


SM predictions for $H \rightarrow \gamma\gamma$

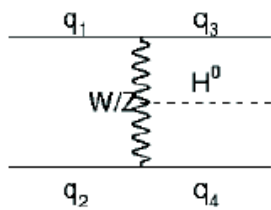
➤ SM Higgs production channels



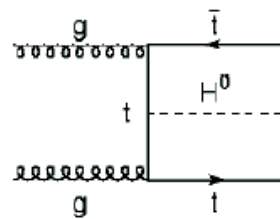
Gluon-gluon fusion (~87%)



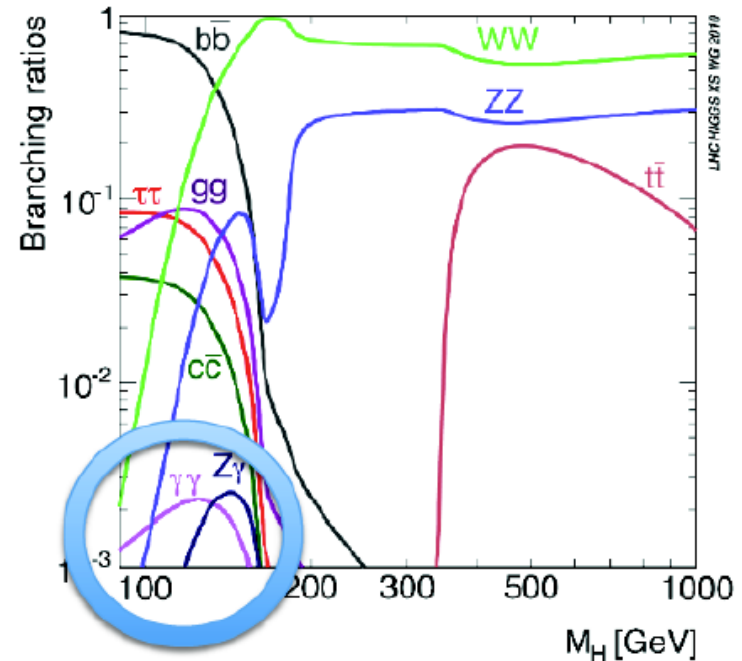
Associated Higgs (< 5%)



Vector-Boson Fusion (~7%)

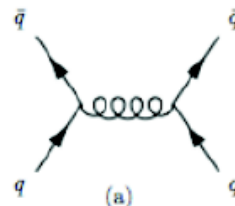
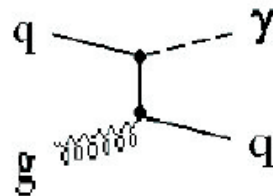
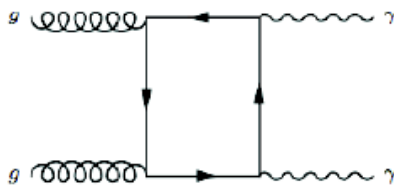


ttH (< 5%)



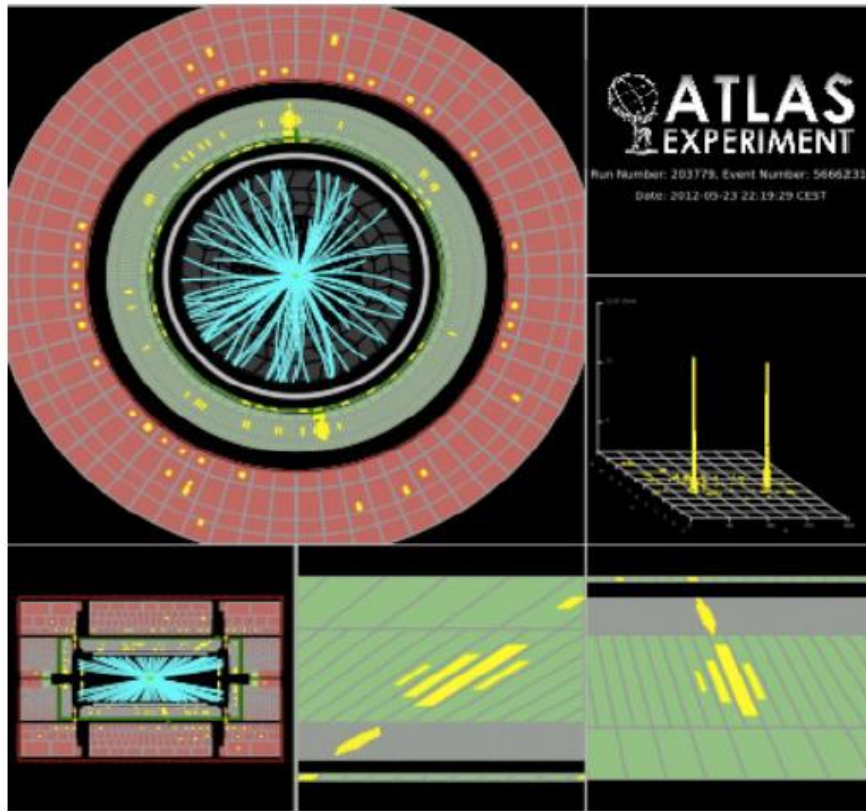
➤ Branching fraction small **but simple signature (two high p_T photons in final state)**

Main backgrounds to $H \rightarrow \gamma\gamma$ are SM diphoton, jet- γ and jet-jet events



➤ Signal expected as **narrow resonance over smooth decaying background**

H- $\rightarrow\gamma\gamma$ event signature



Simple event signature

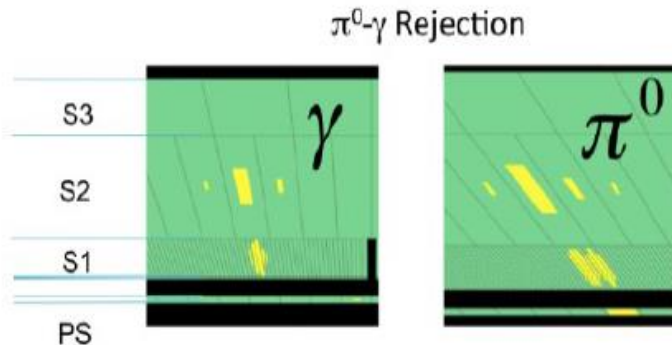
- Two high p_T photons
 $p_{T_1} > 40$ GeV and $p_{T_2} > 30$ GeV
- High trigger efficiency
 $\sim 99\%$
- High event selection efficiency
despite high jet-jet & γ -jet
production
 $\sim 40\%$
- High signal over background
 $\sim 3-10\%$ (depending on sub-category)

Invariant mass reconstruction $m_{\gamma\gamma}^2 = 2 * E_1 E_2 (1 - \cos \alpha)$

- Good energy calibration
- Robust primary vertex reconstruction
- \rightarrow Excellent invariant mass resolution ~ 1.6 GeV with 90% of events within $\pm 2\sigma$**

Shower shapes and vertex reconstr.

Photon ID 2 – Photon shower shapes and background rejection



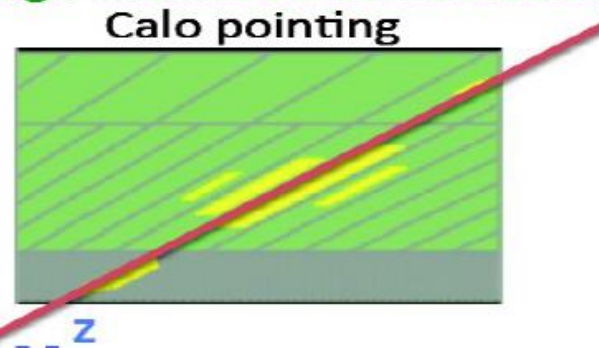
- Photons shower shape distributions in LAr sampling layers - different for signal and background (π^0)

Vertex Reconstruction

$$m_{\gamma\gamma}^2 = 2 * E_1 E_2 (1 - \cos \alpha)$$

☐ Vertex reconstructed through likelihood combination

- Calorimeter 'pointing'
- Σ tracks p_T^2
- Conversion vertex
- Mean vertex position



Event categorization

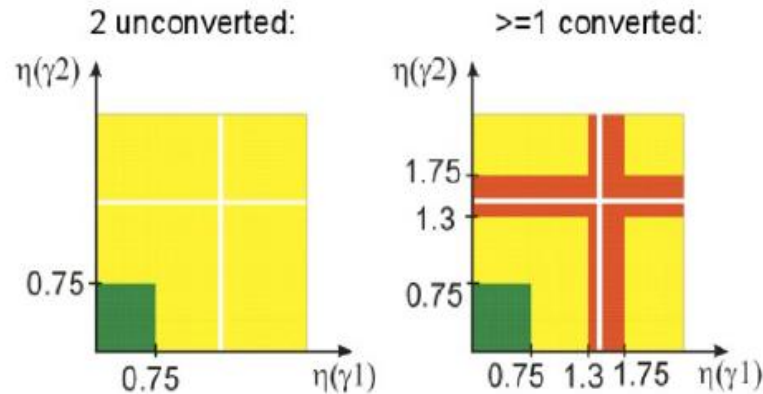
Event categories based on eta, pTt, and conversion

Both unconverted:

- Central
- Rest

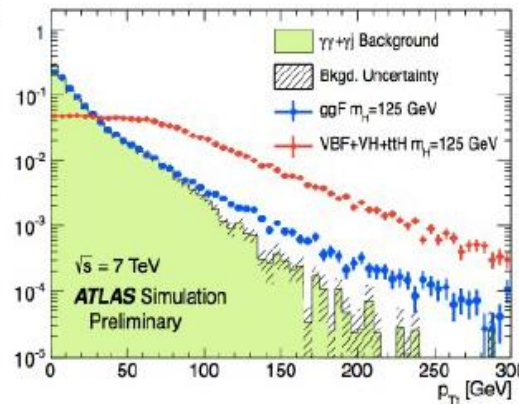
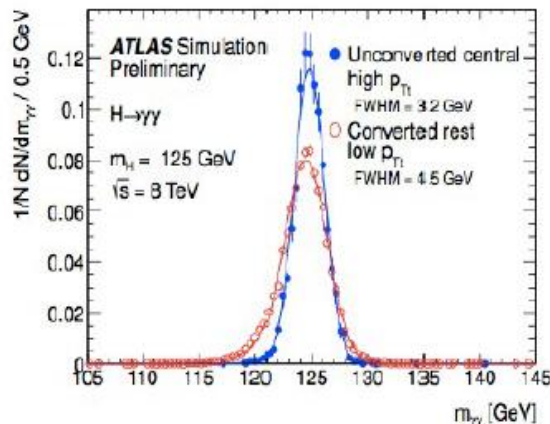
At least one converted:

- Central
- Transition
- Rest

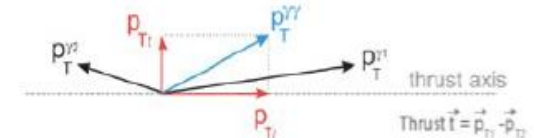


Resolution:

- Good
- Medium
- Poor



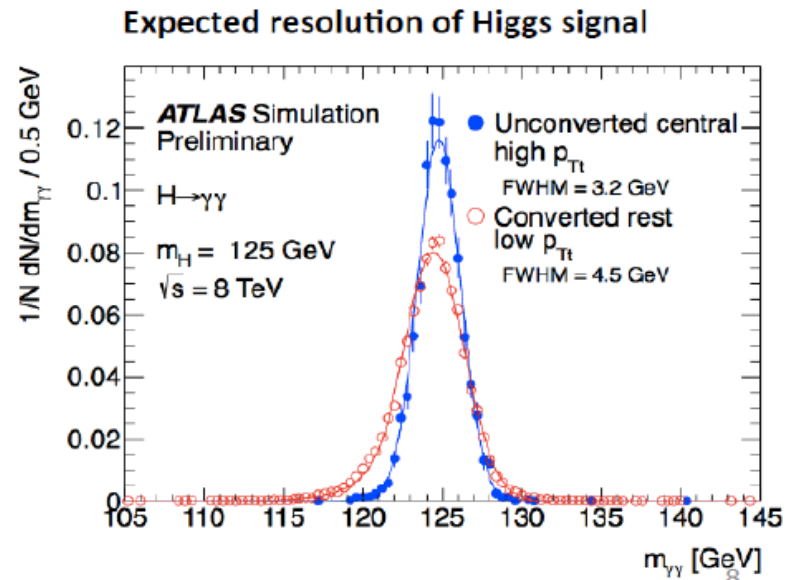
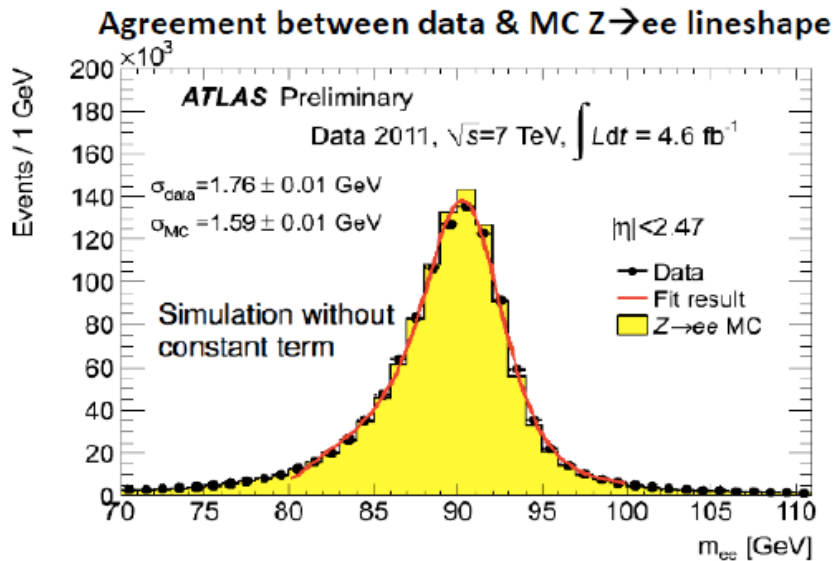
Central and Rest divided into $p_{Tt} < 60$ GeV and $p_{Tt} > 60$ GeV



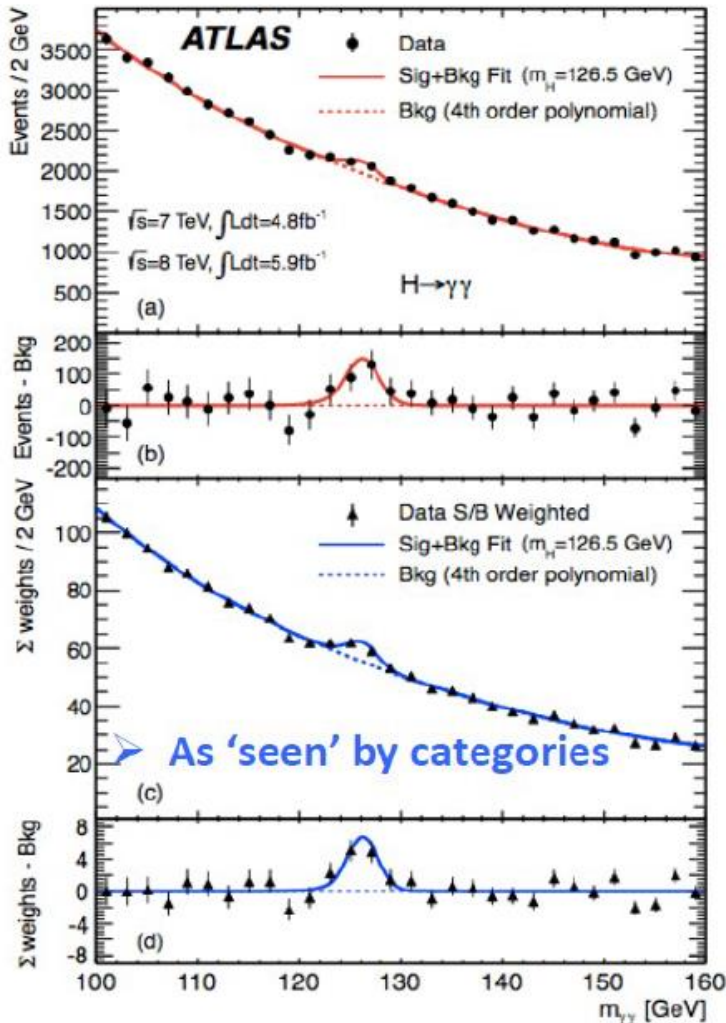
Energy calibration and resolution

$$m_{\gamma\gamma}^2 = 2 * E_1 E_2 (1 - \cos \alpha)$$

- MC based calibration improved with energy scale and resolution corrections based on in-situ analysis of $Z \rightarrow ee$, $W \rightarrow ev$ and $J/\psi \rightarrow ee$
- Energy scale at m_Z known to 0.3%, uniformity (constant term) 1% in barrel, 1.2 – 2.1% in endcap



Invariant mass distribution



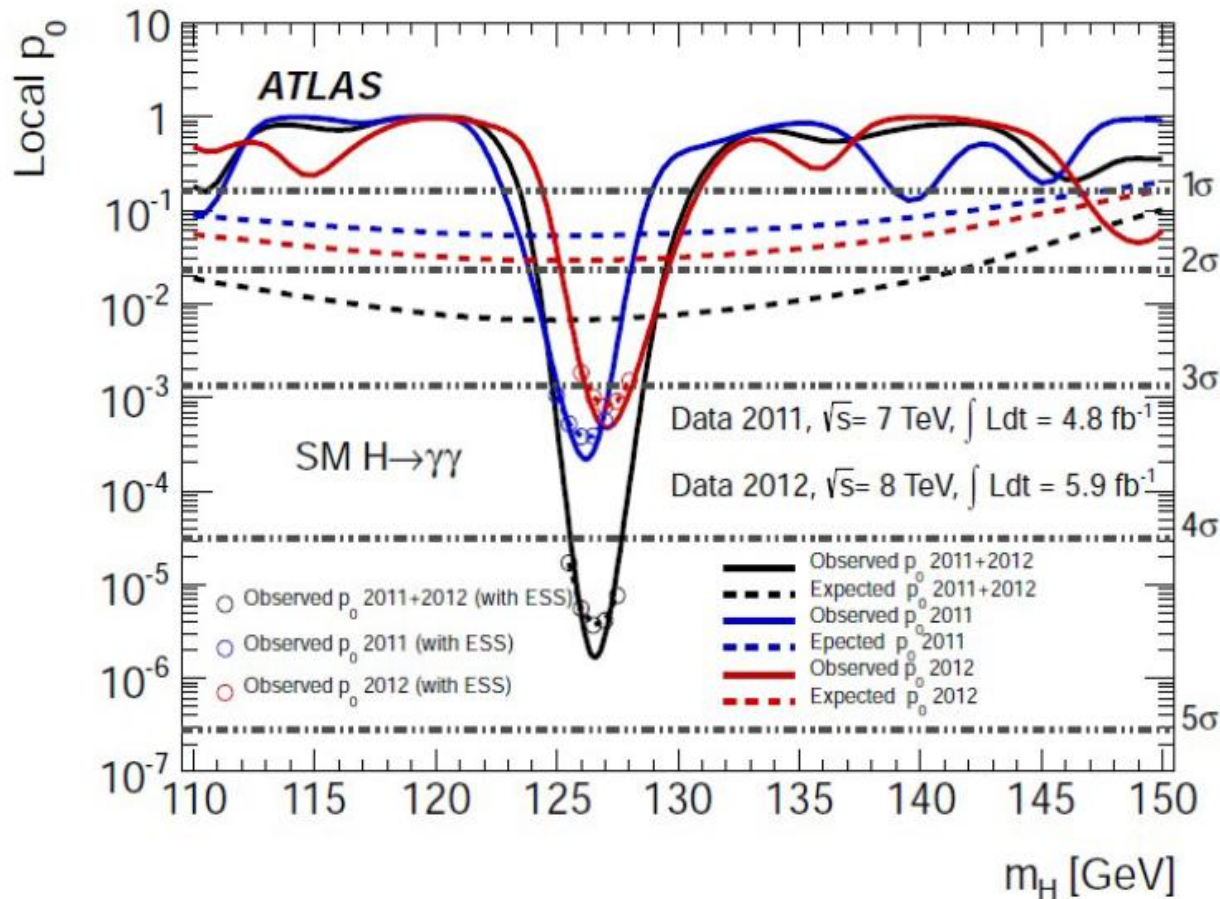
- Photon ID efficiency $\sim 10\%$
- Energy resolution $\sim 14\%$ and mass scale $\sim 0.6\%$
- Isolation $< 1\%$
- Pileup 4%
- Lumi 1-3.6 % (2011-2012)
- Theory cross section
 - \sim up to 25% (for VBF contribution)
 - \sim up to 12% (in other ggF)
 - (underlying event $\sim 5\%$ and PTt dist up to 12% at high PTt)
- Bkg Param (evts) 0.2-4.6 (0.3-6.8) for 2011(2012)

In VBF category

- Jet E-scale 9-10%
- Underl. Evt. 6-30%
- Higgs p_T up to 12.5%

23788 events (7 TeV) and 35251 events (8 TeV)
 Background+signal fit, signal fixed at 126.5 GeV

Quantifying the excess p0



- Excess (m_H):
126.5 GeV
- Expected (local-significance):
2.5 σ
- Observed (local-significance):
4.5 σ
- Fitted signal strength:
 $\hat{\mu} = 1.8 \pm 0.5$