

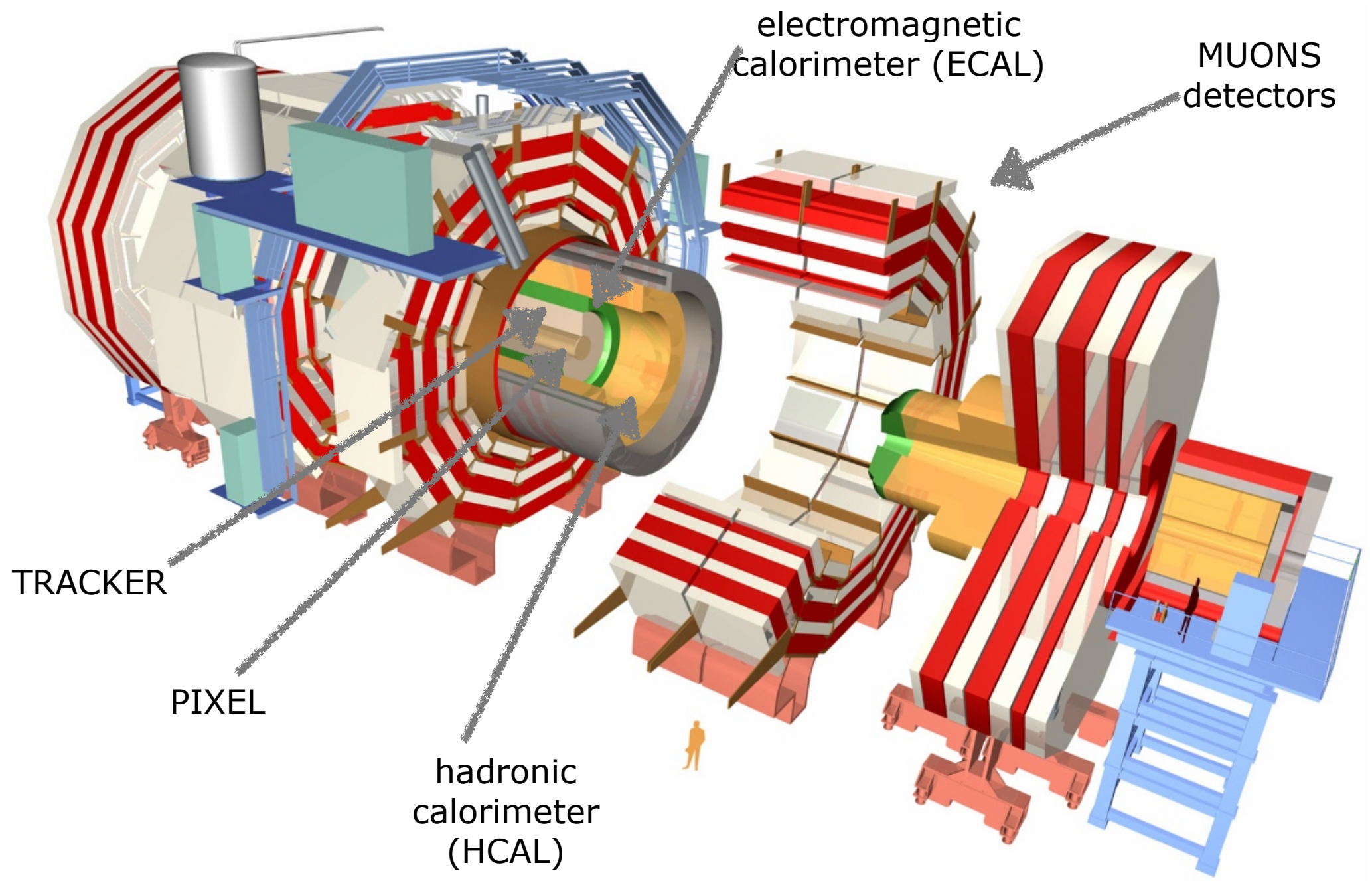
introduction

- my interpretation of “data analysis techniques” is here “doing a data analysis”
- follow the steps from the beginning (data taking) to the end (the result)
 - ▶ the luminosity
 - ▶ the trigger, from the point of view of the analysis
 - ▶ the reconstruction and detector response
 - ▶ the simulation
 - ▶ differential cross-section measurement: a di-jet correction
 - ▶ searches: the $H > WW > l\nu l\nu$
 - ▶ multivariate techniques

thanks to the following people, for interesting discussions, for liberally “borrowing” slides, or both: D. Benedetti, C. Bernet, T. Camporesi, G. Cowan, K. Cranmer, K. Ellis, S. Gennai, A. Ghezzi, A. Hoecker, R. Van Kooten, M. Nguyen, M. Paganoni, M. Pelliccioni, E. Rizvi, R. Rossin ...

physics objects reconstruction

the cms detector



TRACKER

PIXEL

hadronic
calorimeter
(HCAL)

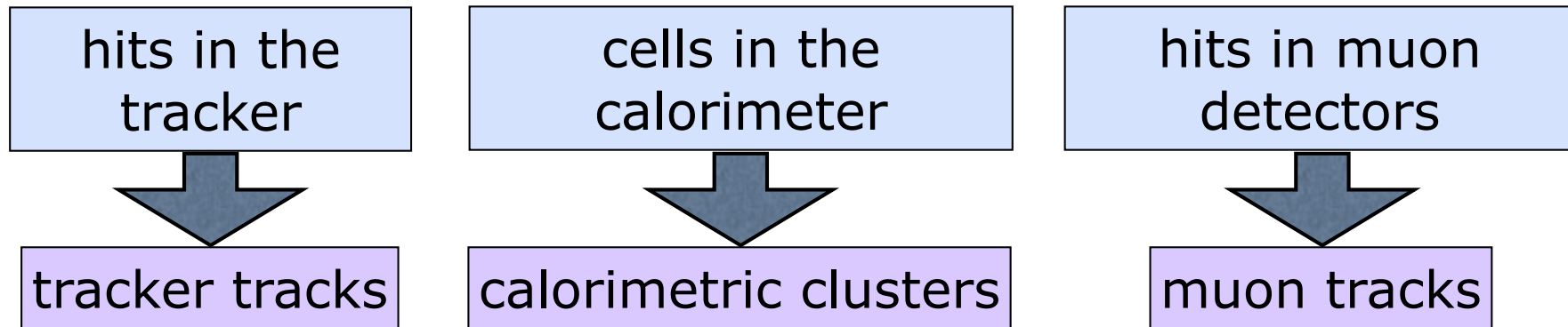
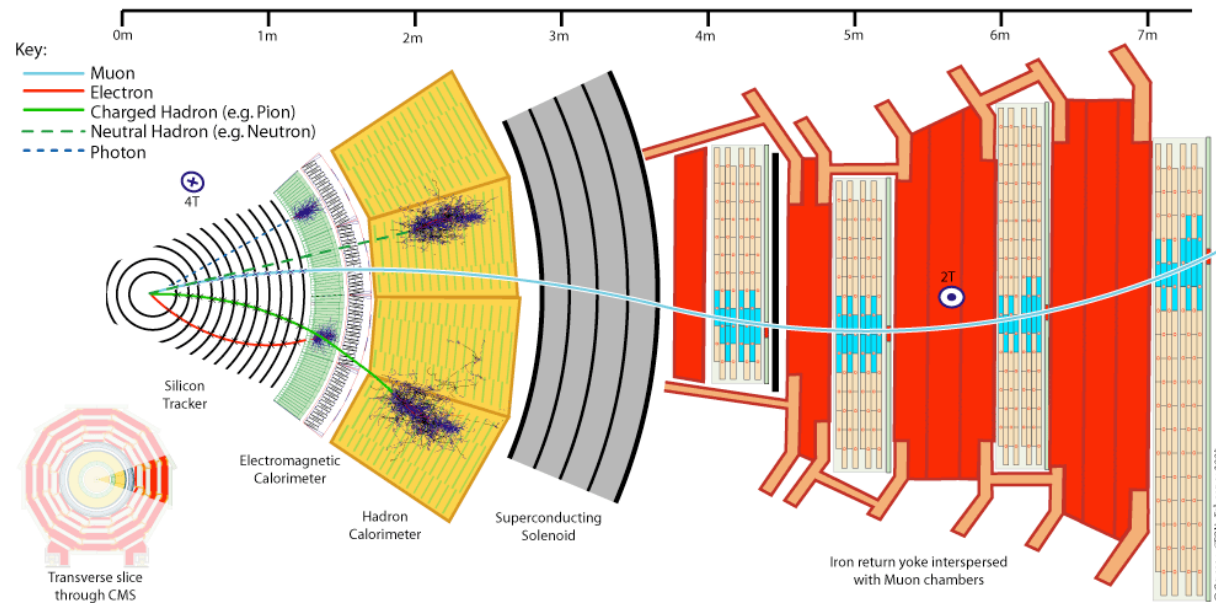
electromagnetic
calorimeter (ECAL)

MUONS
detectors

physics objects reconstruction

- obtain **physics objects from the detector response**
 - hits in the tracker and muon detectors
 - energy deposits in the calorimeters
- two ways are available in CMS
 - **single objects** reconstruction: build final objects (e.g. muons, electrons, jets) from the detector response
 - **particle-flow** reconstruction: build a coherent list of stable particles and produce the analysis objects on top of them

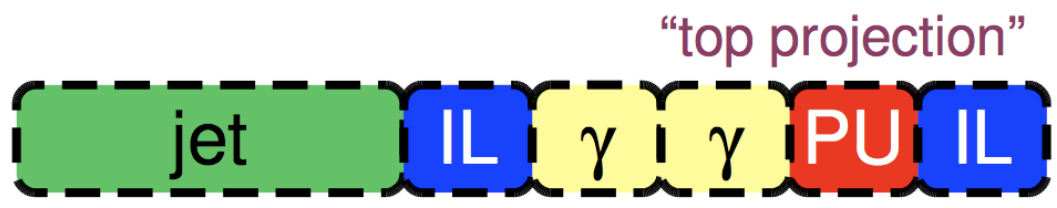
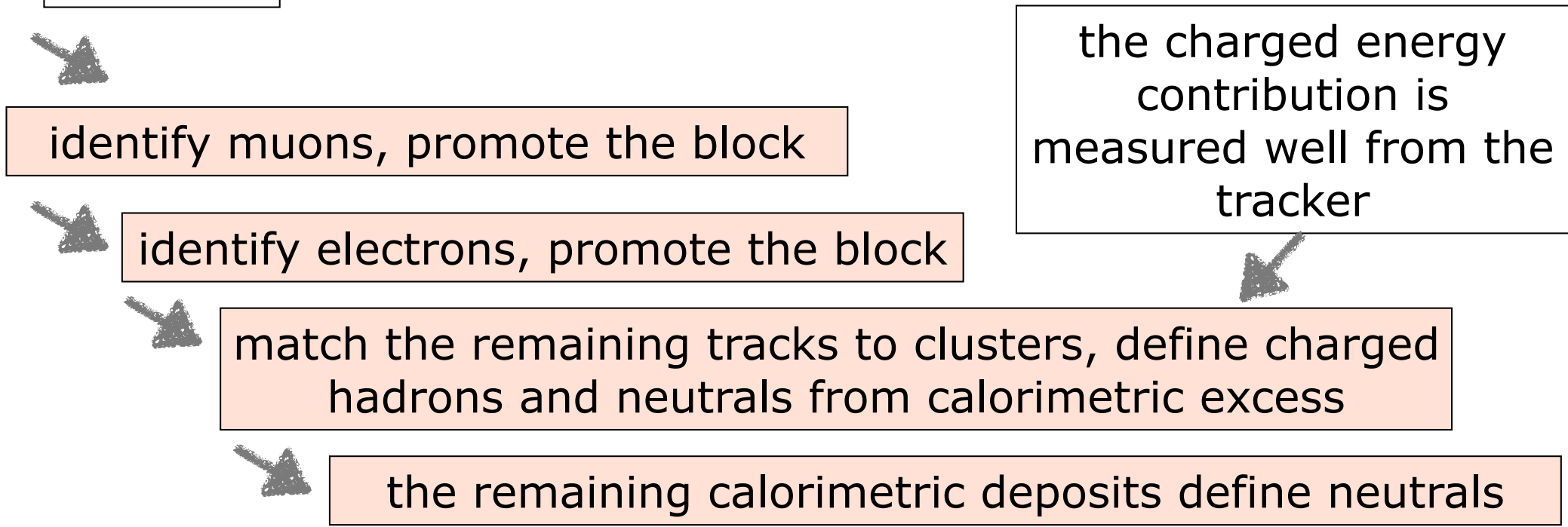
the particle flow



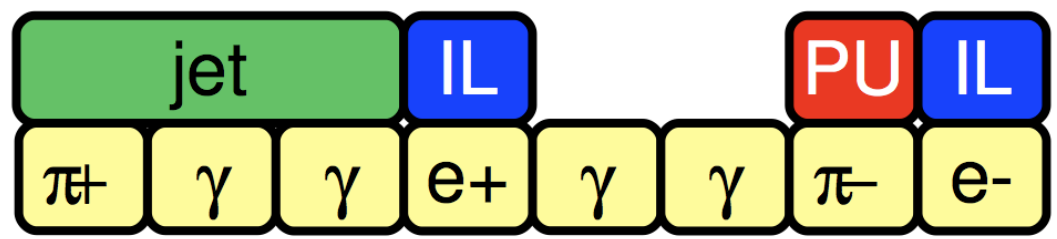
link the single objects with geometrical requirements on the extrapolated trajectories and create **blocks**

the particle flow

blocks



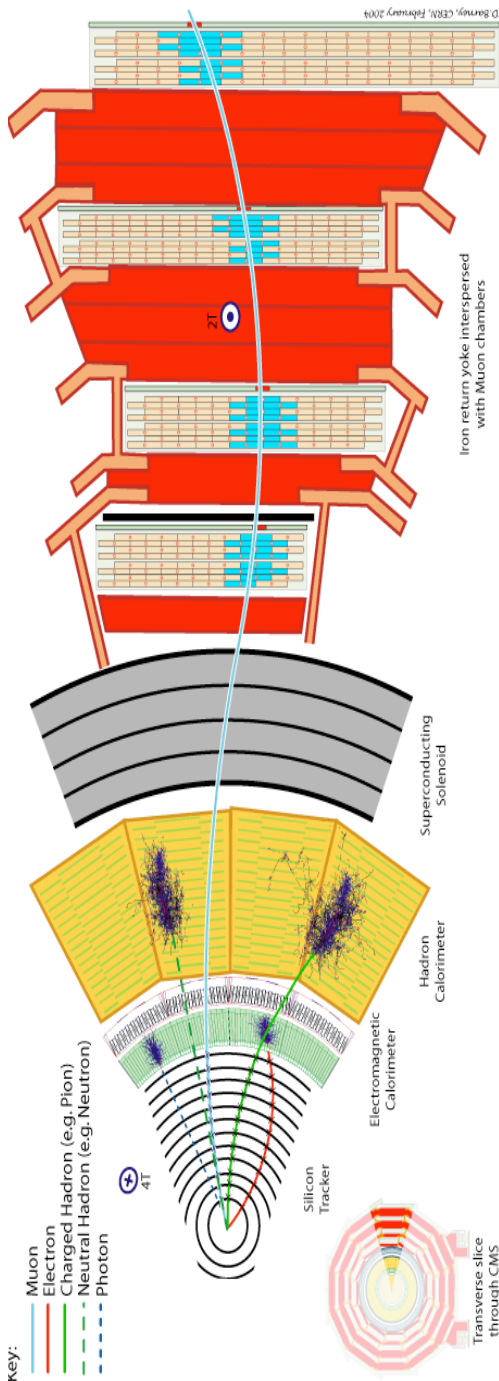
$$+ \vec{MET} = - \sum_{i=0}^{N_{particles}} \vec{E}_T^i$$



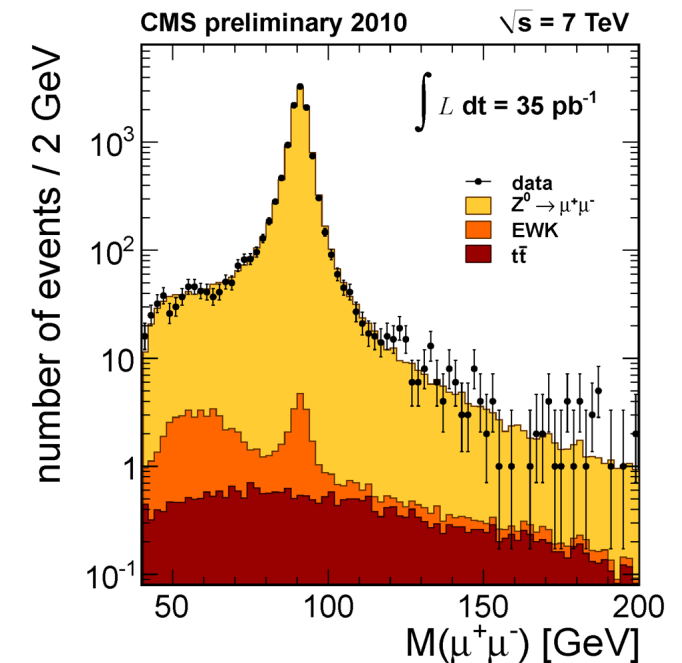
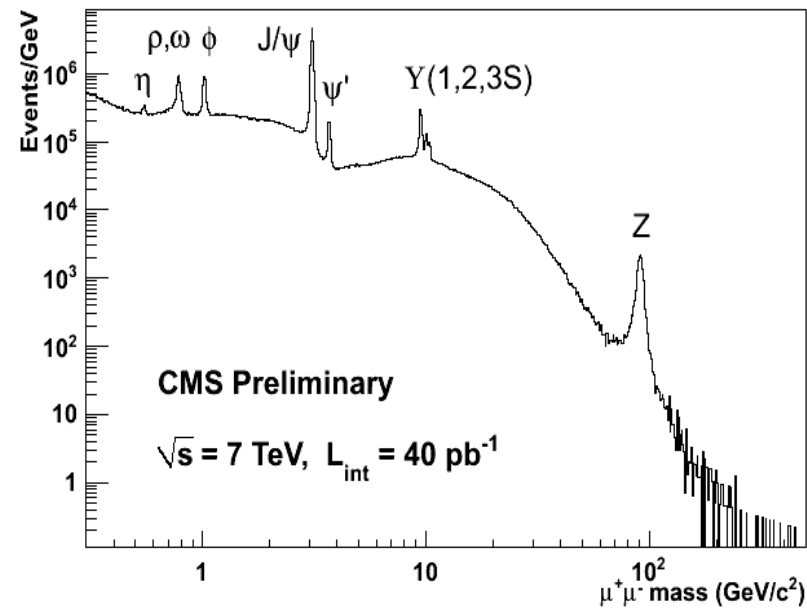
The list of PFCandidates

the list of particles obtained (candidates) is used for high level objects classification and reconstruction, to be used in the analysis

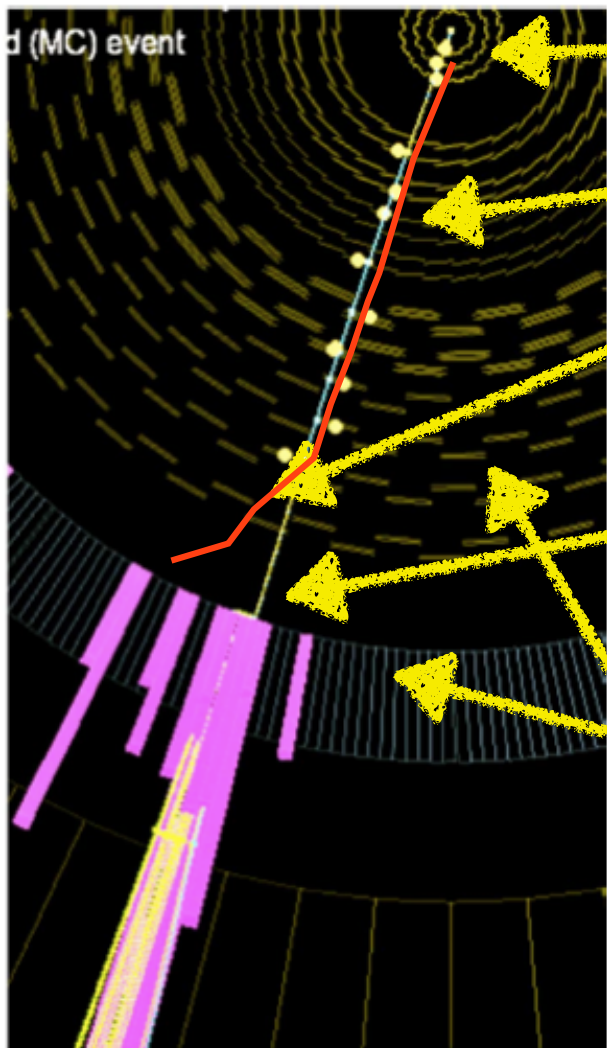
muon reconstruction



- high **purity** = fit with hits in both tracker and muon
- high **efficiency** = fit in the tracker + confirmation in the muon detector
- **momentum determination** from both tracker and muons information: best resolution from the tracker for $p_T < 200$ GeV, from the muons above (effect of multiple scattering)
- above 1 TeV, the **bremstrahlung** is significant



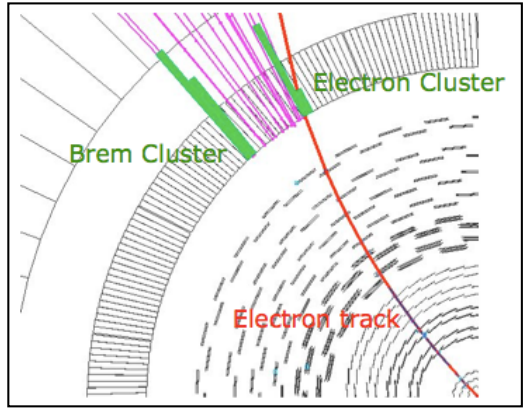
electron reconstruction



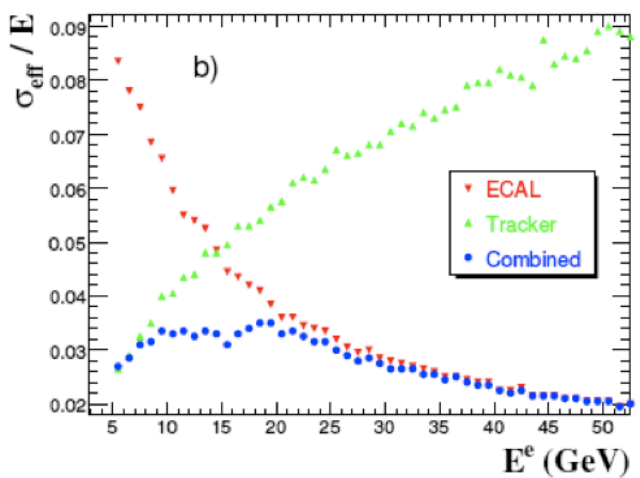
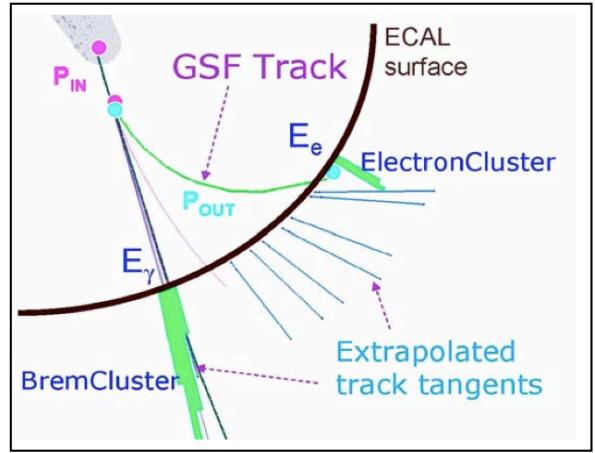
- Seeding
- electron tracking: GSF
- cleaning of GSF track duplicates
- Identification of all the electron energy deposits in the ECAL
- electron 4-mom determination.

Identification against charged-particles interacting in the ECAL

from ECAL clusters or tracks:

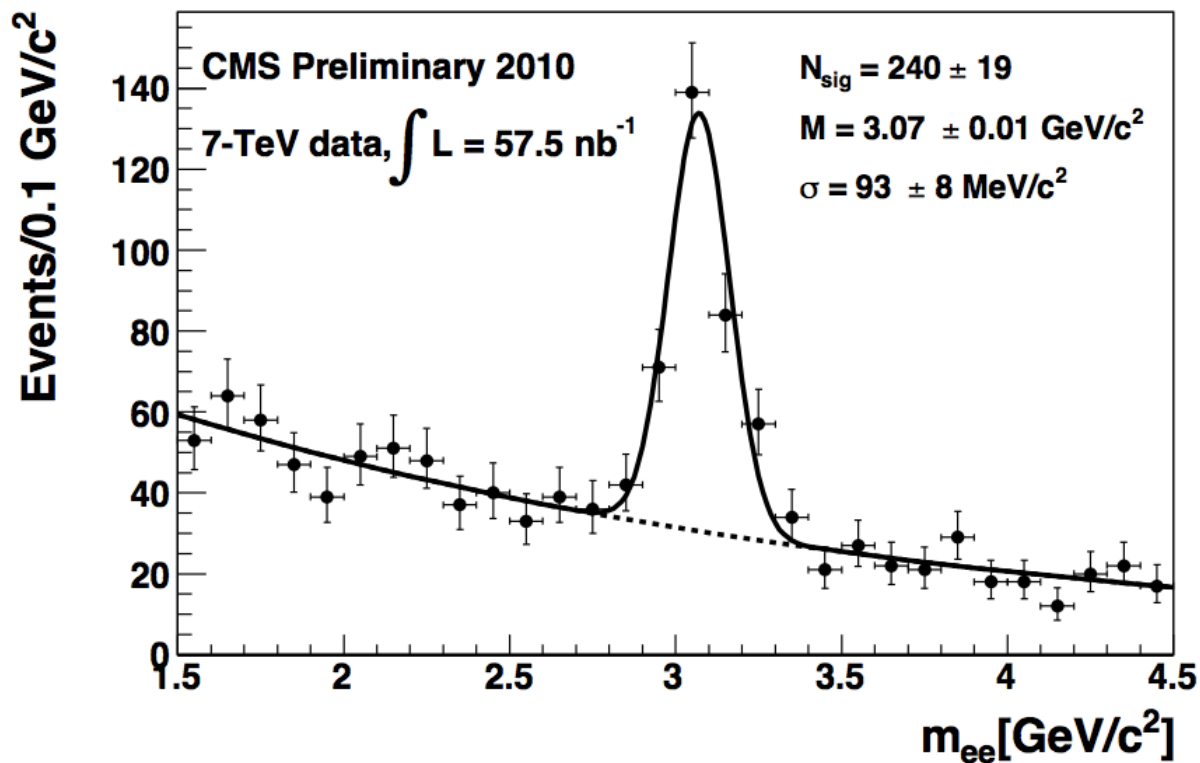


from ECAL footprint or tracks extrapolation:



use ECAL at high p_T,
tracker at low p_T

electron reconstruction



search for the decay:

$$J/\Psi \rightarrow e^+ e^-$$

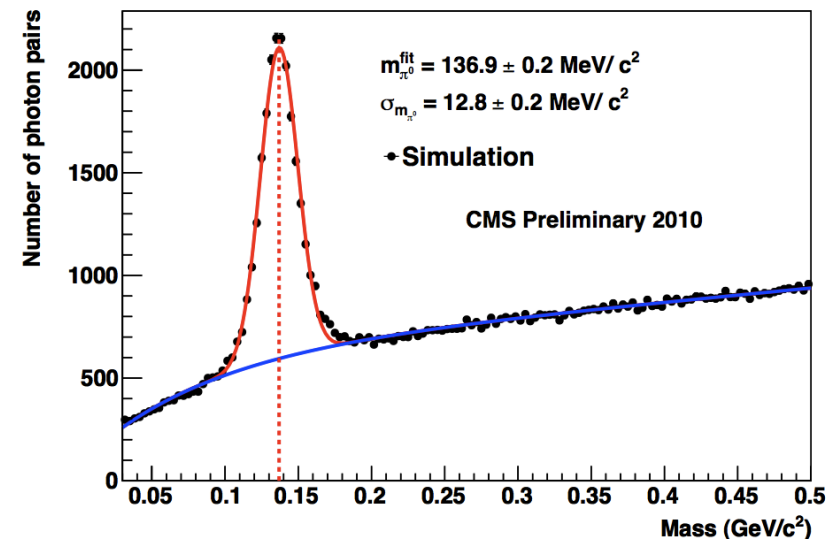
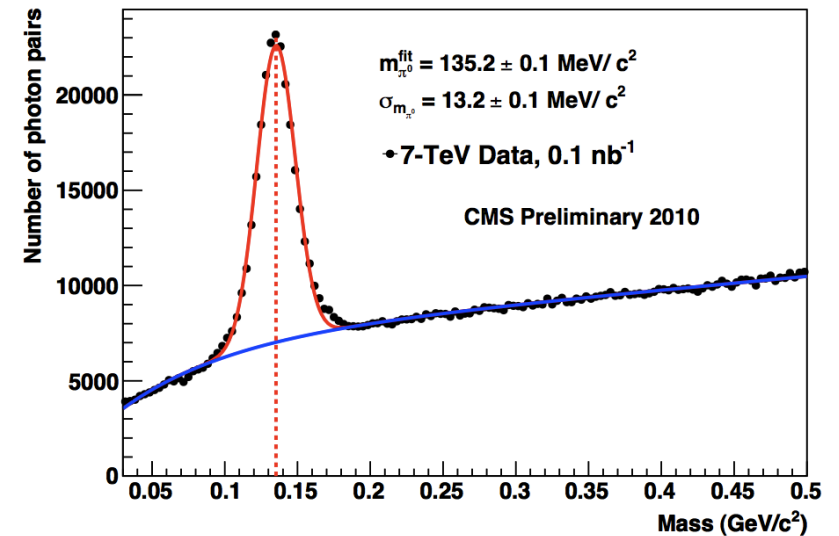
contamination sources:

- real electrons, either from photon conversions or from semi-leptonic b-hadron decays,
- mis-identified charged hadrons.

- at most one hit missing in the pixel detector (reduce conversions)
- electrons originate from the same vertex (reduce the b-decay background)
- quality cuts to reject charged hadrons contamination
- opposite charge

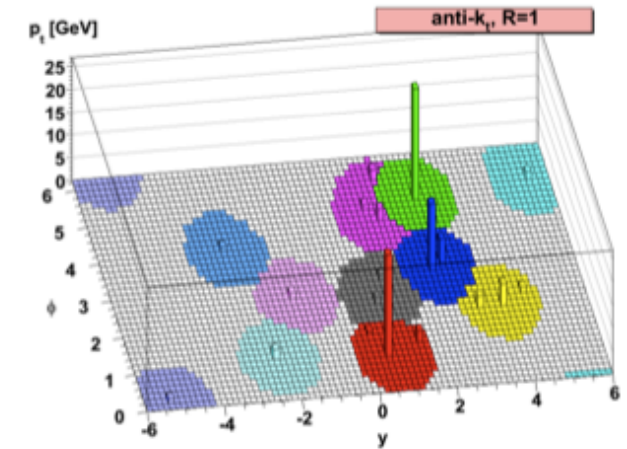
photon reconstruction

- **ECAL clusters** not associated to a track, nor a deposit in the hadronic calorimeter
- ECAL detector response is **calibrated**, to account for the effect of the noise cut on the single crystals readout
- **check the photons energy scale** calibration with 2010 data, by looking at the π^0 peak position
- pair all photons with at least 400 MeV energy
- determine the peak position with a combined fit of signal + background

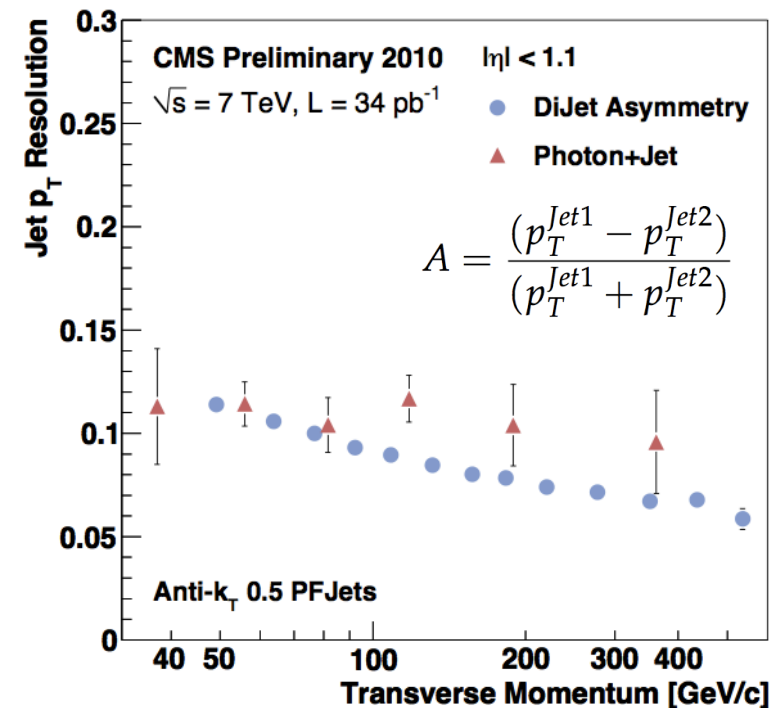
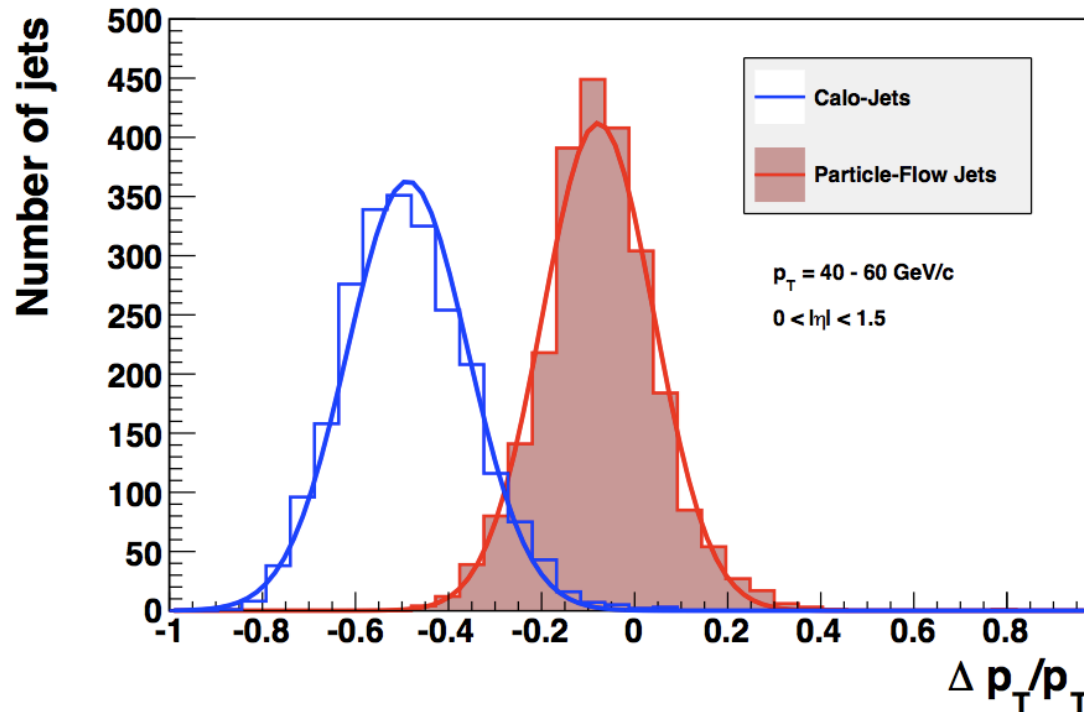


jet reconstruction

- jets are reconstructed with the AKT5 algorithm
- for the single object reconstruction: with calorimetric deposits
- for the particle-flow: with particle flow candidates

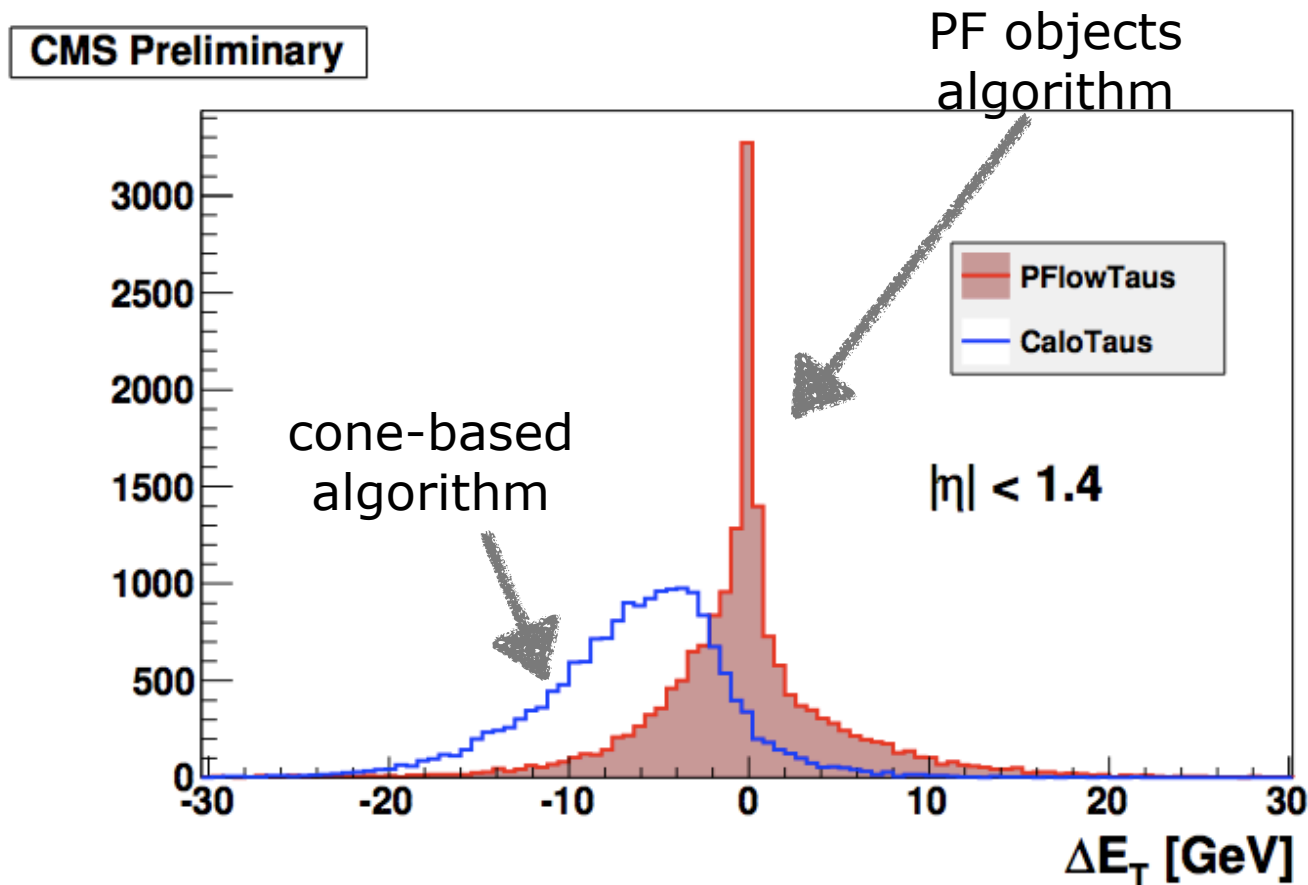


CMS Preliminary



tau reconstruction

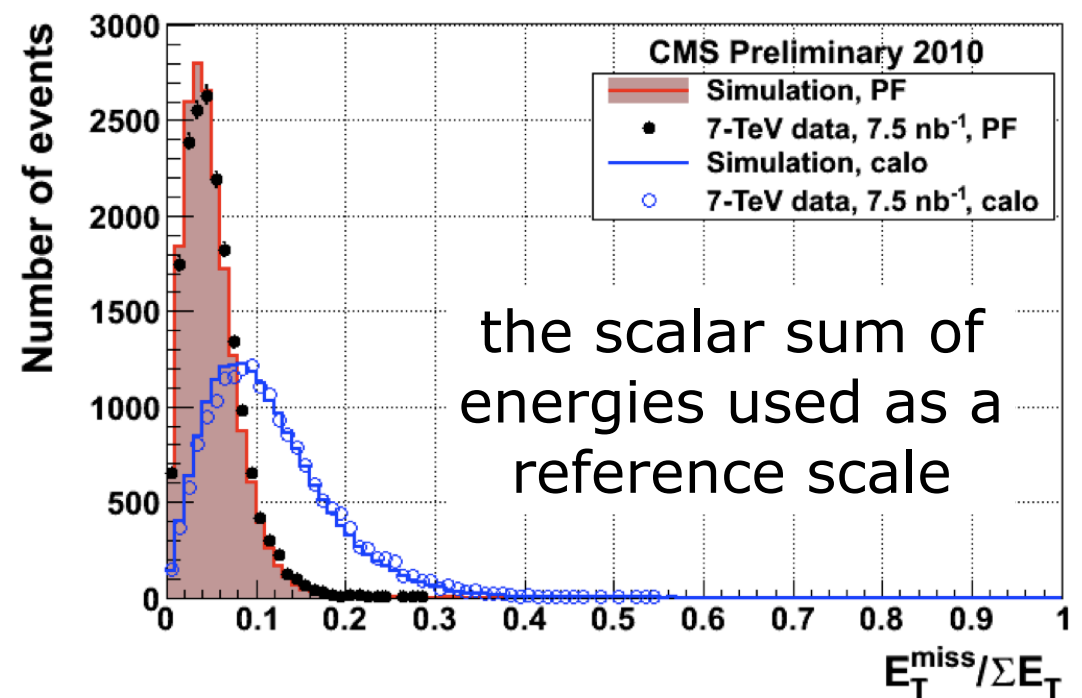
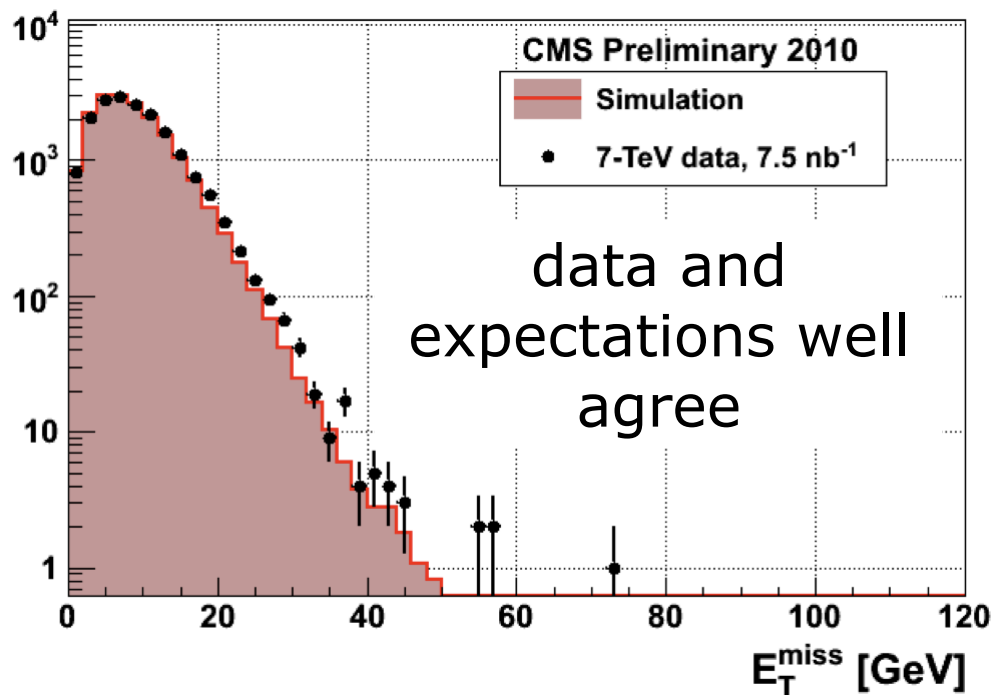
- reconstructed as narrow jets in the standard case, as the sum of the particles compatible with the tau decay in a narrow cone in the particle flow case



reconstructed taus E_T compared to the expected one,
test performed on a simulated $Z > \pi$ sample

missing energy reconstruction

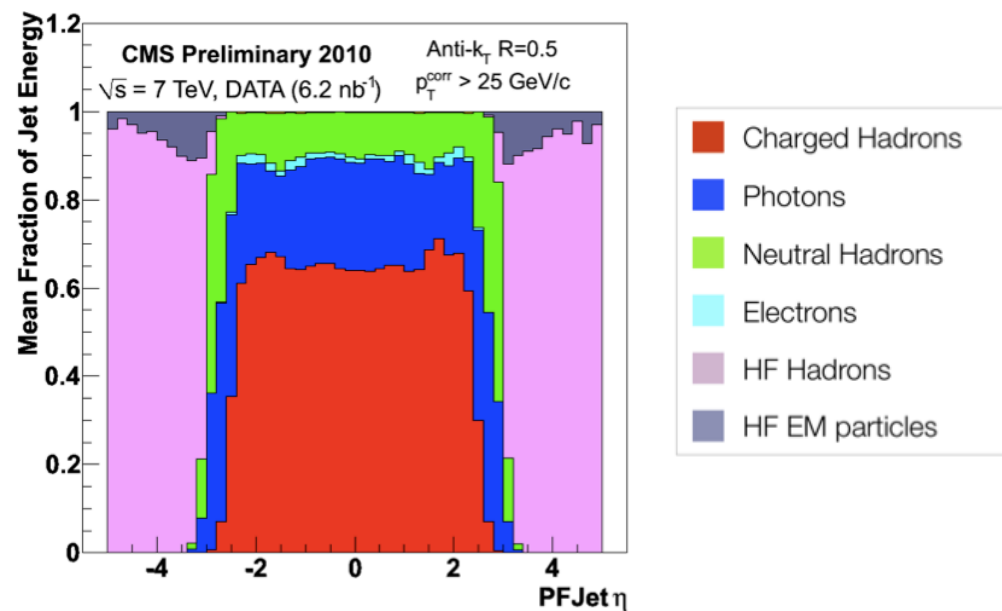
- derived from (minus) the sum of “all the rest”
- sensitive to uncertainties in all the other physics objects
- noise effects, mis-calibrations, etc. generate fake missing energy in events without missing energy
- perform a test on a di-jet sample



reconstruction: in summary

- the reconstruction obtains from the detector measurements the physics objects in the final state
- in a coherent way, to close the kinematics (as much as possible)
- making use of the most precise sub-detector
- reconstruction and identification are not (always) disentangled, for example electrons need to be separated from jets
- data-driven techniques necessary to assess the performances

jet composition:
only for neutral
hadrons one cannot
profit of tracker
measurements



detector response

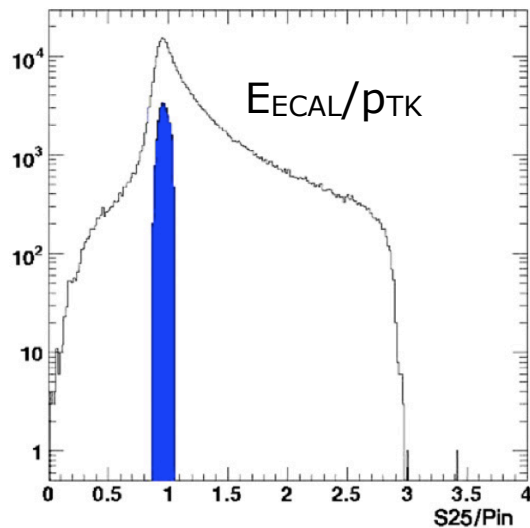
- the detector response is not perfect
- the output of the reconstruction needs to be **calibrated for the detector response**
- use **known physics processes** to get the calibrations and the relative uncertainty
- for example
 - **resonances** for leptons (energy scale, tag&probe)
 - **cosmic** rays (alignments)
 - transverse momentum **balances**
 -

ECAL calibration

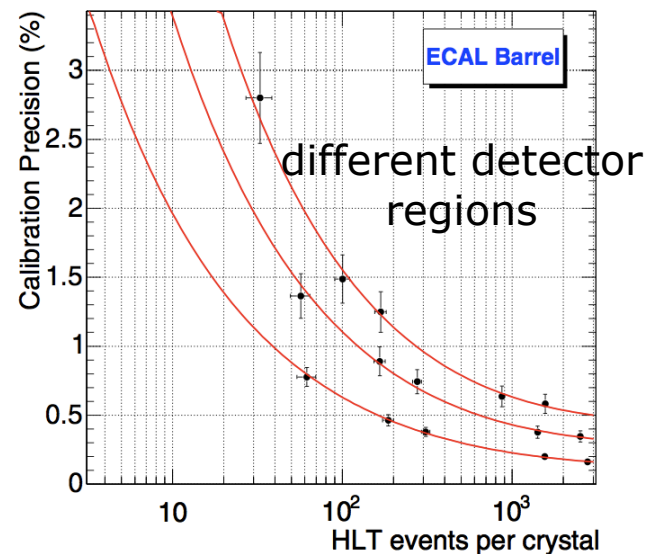
- each ECAL channel needs a calibration factor to equalize the response of all detector elements
- for electrons, the energy is measured in the tracker and in the ECAL
- find the calibration coefficients by minimizing a χ^2 of:

$$\text{energy in single elements} \rightarrow |\mathcal{E} \times \bar{c} - \mathcal{P}| \leftarrow \text{electrons momenta}$$

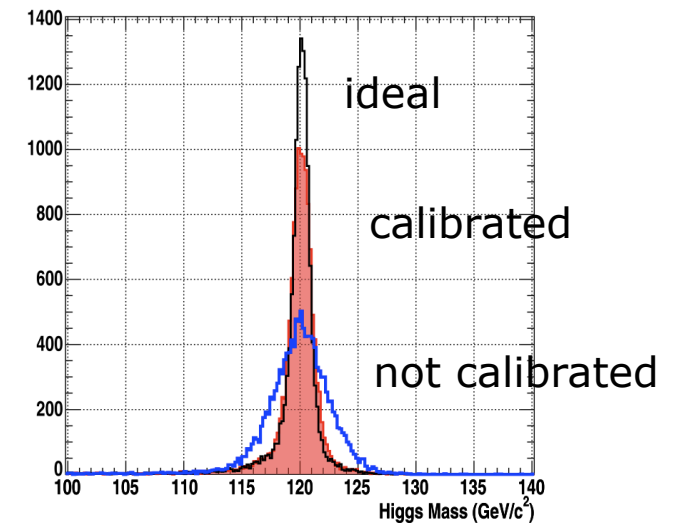
unknown coefficients



select good isolated electrons



statistical trend with luminosity



effect on $H \rightarrow \gamma\gamma$ invariant mass

jet energy corrections

- the jet energy scale needs to be calibrated, as a function of various variables

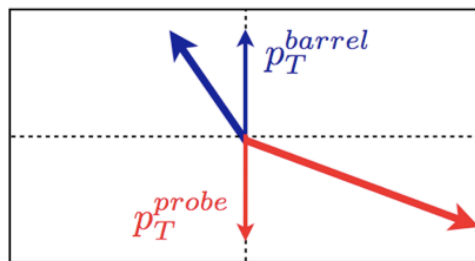


detector noise effects, pile-up

tag&probe like: di-jets events assumed to be balanced, get a relative correction

γ+jet balance in the transverse plane

Barrel Jet

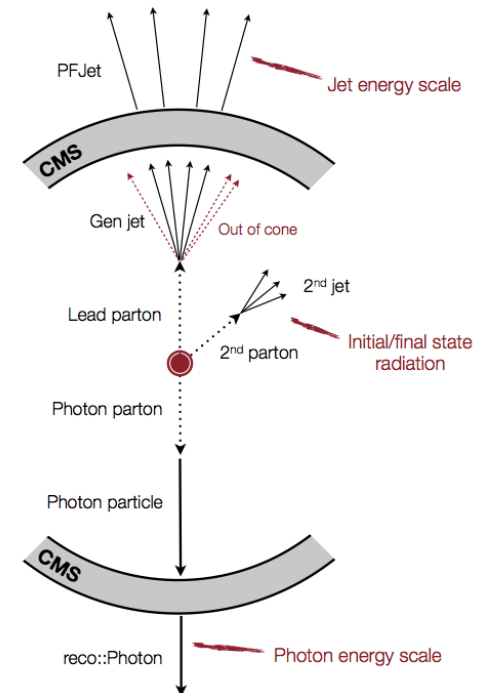


Probe Jet

$$p_T^{dijet} = \frac{p_T^{probe} + p_T^{barrel}}{2}$$

$$B = \frac{p_T^{probe} - p_T^{barrel}}{p_T^{dijet}}$$

$$r = \frac{2 + \langle B \rangle}{2 - \langle B \rangle}$$



the simulation

the simulation

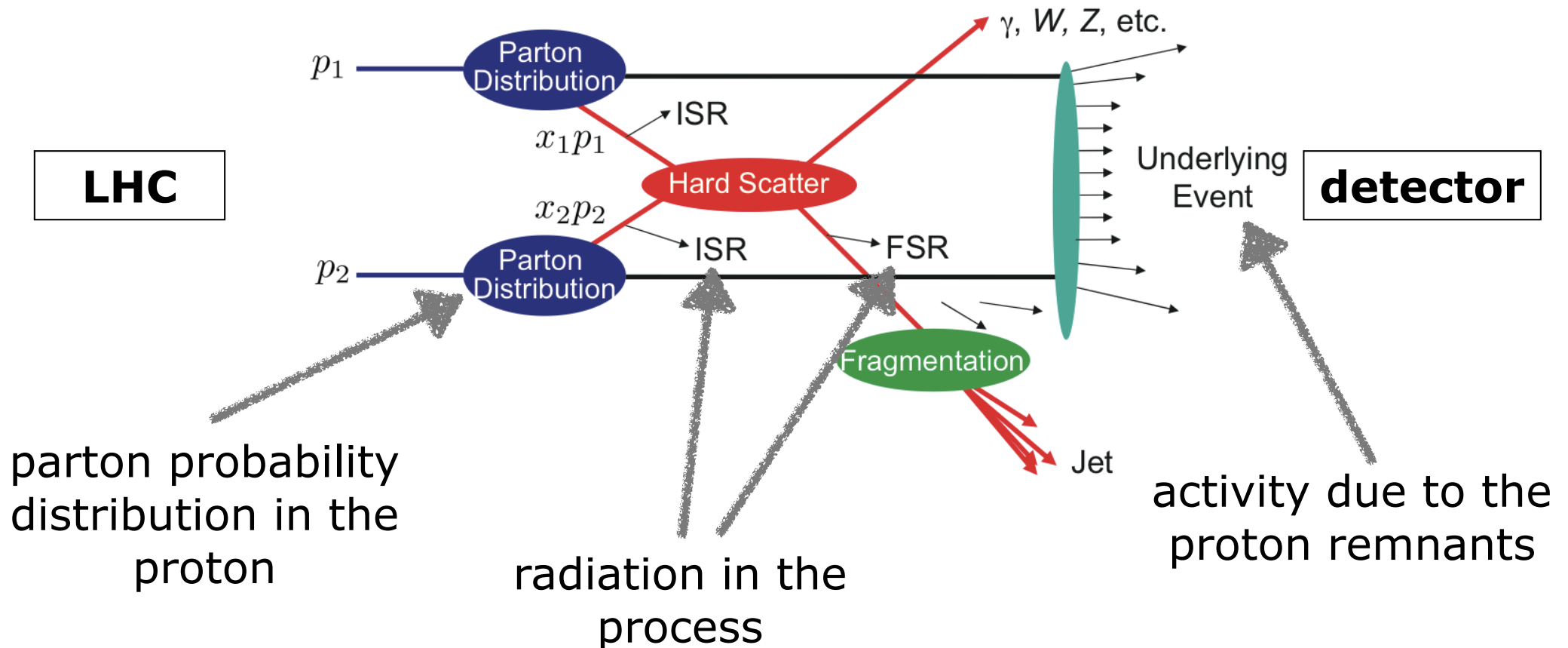
$$\sigma = \frac{N_{obs} - N_{bkg}}{\varepsilon \cdot \int \mathcal{L} dt}$$


$$\varepsilon = \varepsilon_{tr} \cdot \varepsilon_{reco} \cdot \varepsilon_{ID} \cdot \varepsilon_{sel}$$

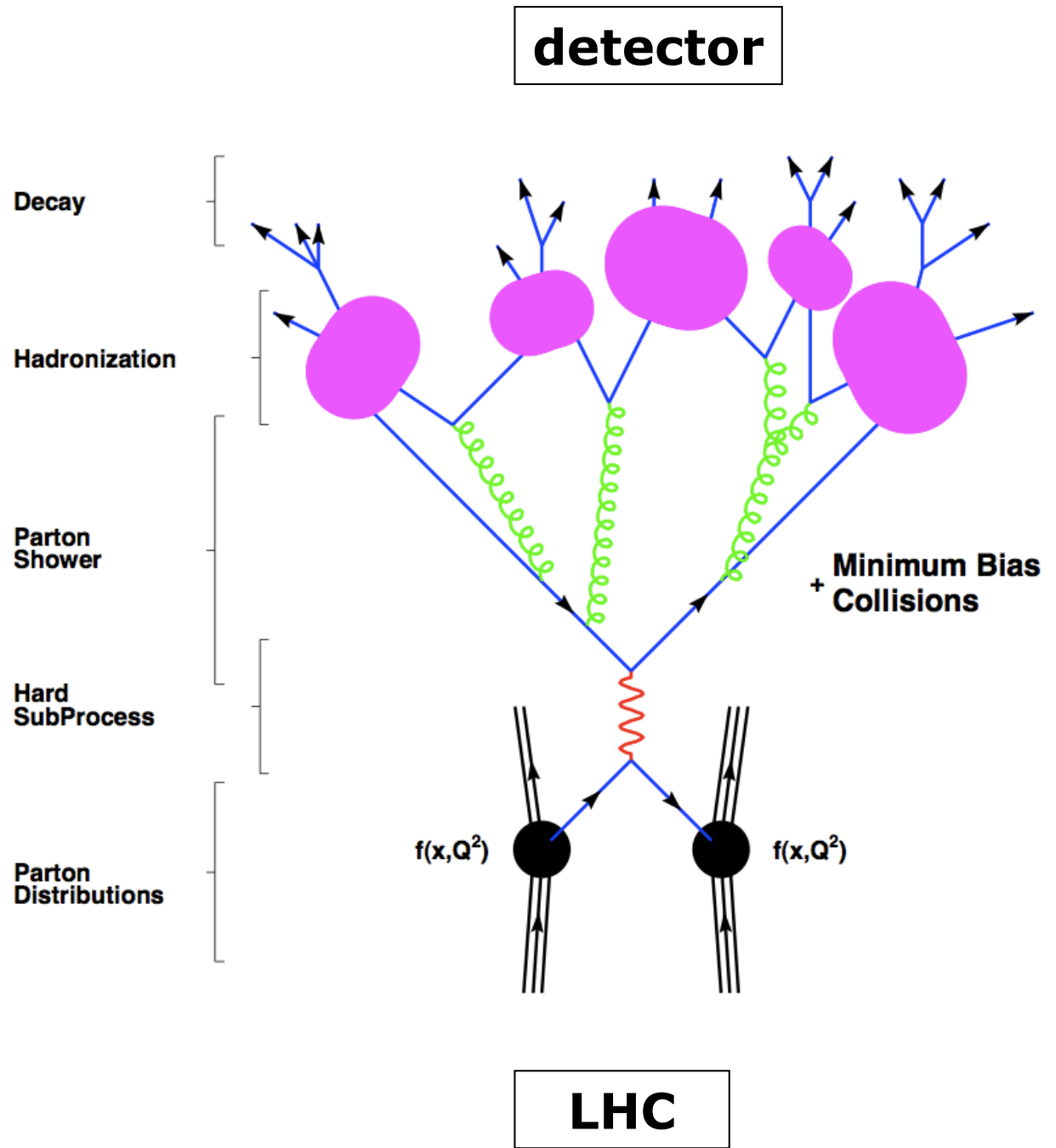
- calculate what fraction of events from a given decay **falls within the detector acceptance and the selections** of the analysis
- need a **forecast of how the event develops** in space, after the interaction
- the **simulations** are necessary both for known physics objects (Z, W production) and, of course, to build searches for new physics
- the **uncertainty** in the input parameters is source of systematics

the simulation

- calculate **inclusive cross-sections**
- calculate **differential cross sections** as a function of variables of interest in the analysis
- provide **simulated events**, that mimic Physics, and have on average the behaviour foreseen by the theoretical model



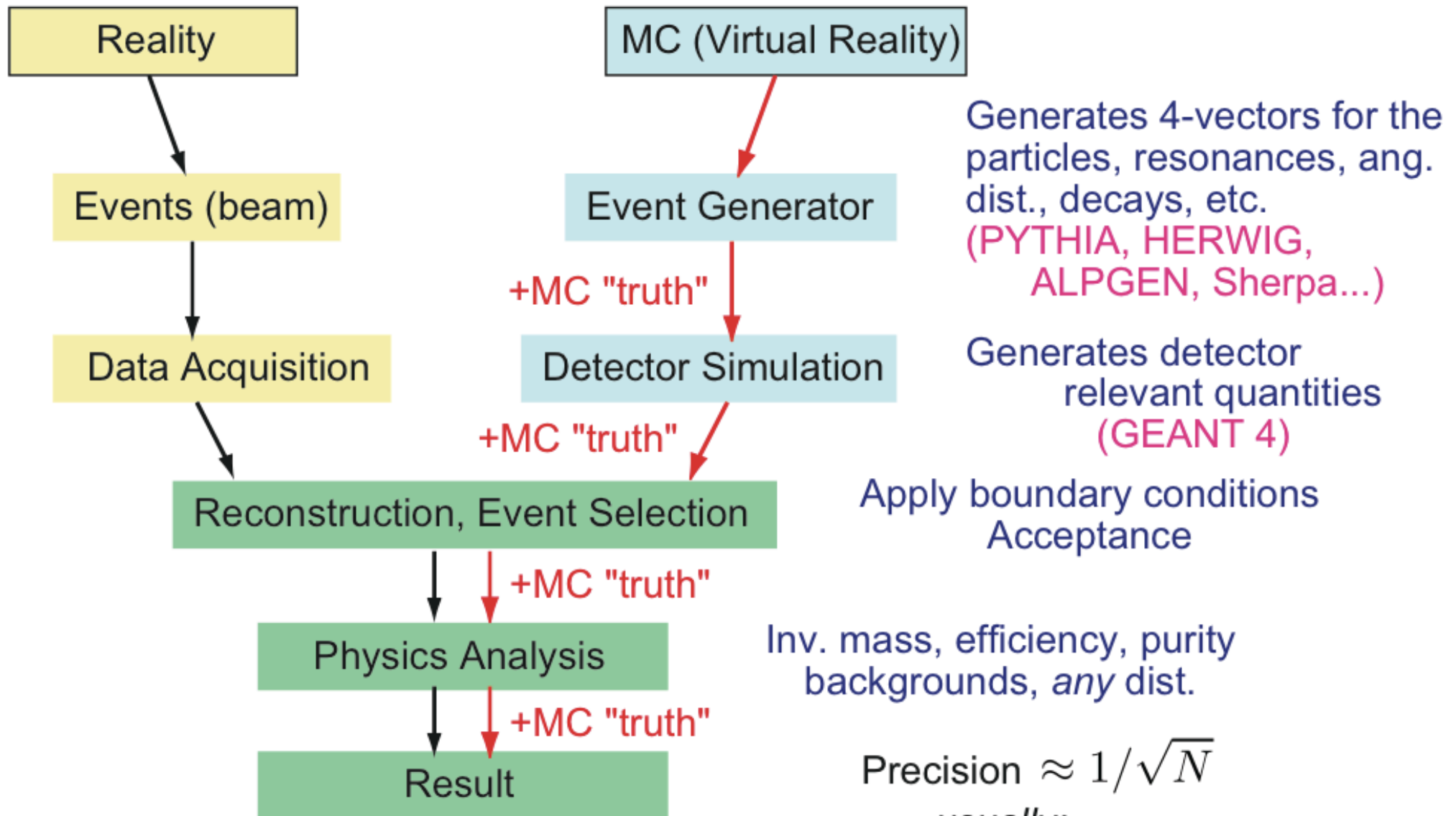
the physics event generation



the simulation of the detector

- each experiment creates a **simulation of the detector**
- the GEANT program uses generator output (4-vectors) and simulates the **interaction of particles within the detector volume** (need a good description of the geometry):
 - particle ionization in trackers
 - energy deposition in calorimeters
 - intermediate particle decays/radiation
- the GEANT code is merged with (experiment specific) **detector simulation**
- final output: the response of the electronics readout
- MC events are in the **same format as real raw data**

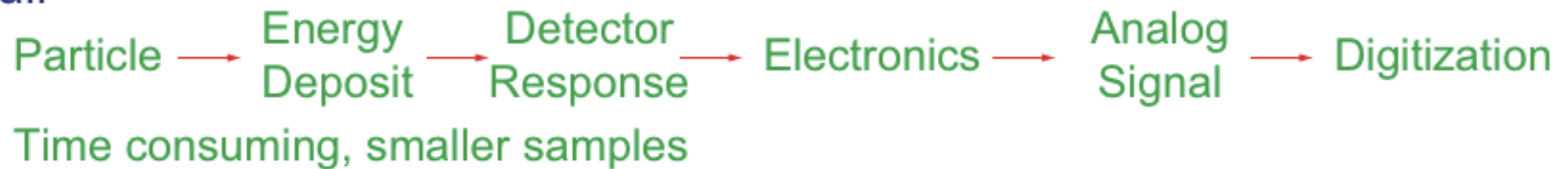
the samples processing



levels of simulation

Three typical levels of MC simulation:

- Full



- "Fast" or parameterized

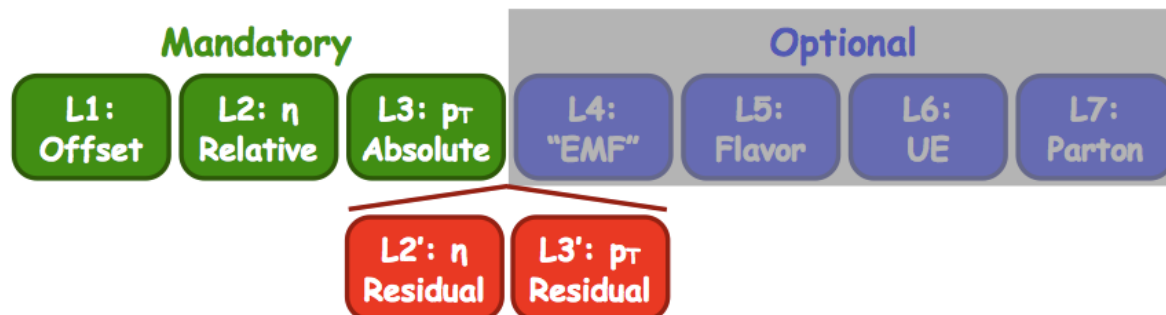
Intelligently smeared 4-vectors, efficiencies, noise (from data and full MC)
 And/or calorimeter shower libraries
 Larger samples

- Toy

Only throw from the handful of prob. dist. functions that you care about
 (with correlations)
 "Roll your own", usually write (easy in root!) and run yourself
 Crazy-large samples, quickly
 To determine probability of fluctuations, checks for systematic effects, etc..

beware of your simulation

- the simulation is a multi-dimensional **parametrization** of the knowledge of the detector and standard model predictions
 - is the theoretical simulation correct for the analysis?
 - additional jets production is crucial for analyses that apply a jet veto
 - spin correlations in the Higgs decay need to be treated correctly
- is the behaviour of the simulation in **agreement** with data, in the phase space of interest for the analysis?

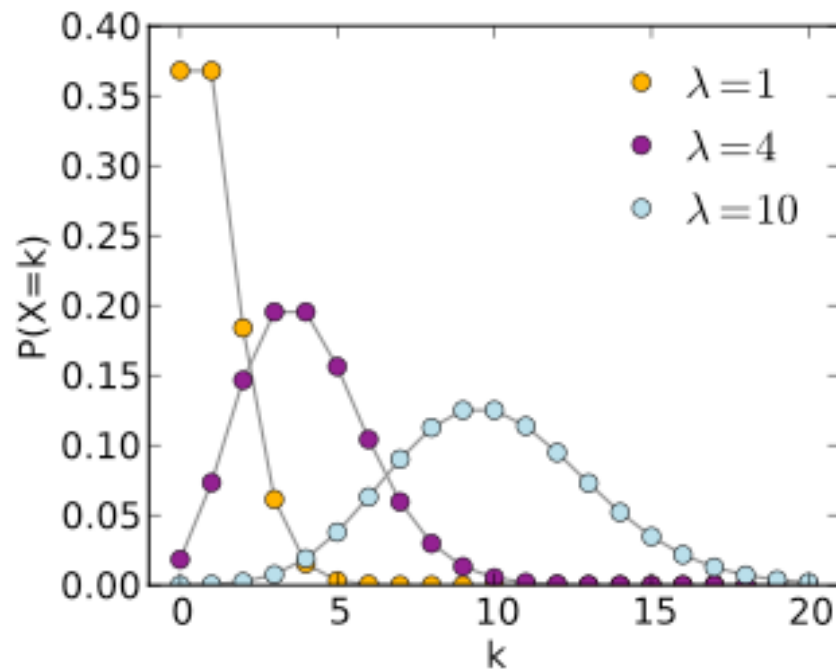


when there's agreement, **use it**:
the jet energy scale at CMS is calculated as a correction factor to the one obtained from simulation

the pile-up

the pile-up

- At LHC, the interaction rate is higher than the bunch crossing rate
- Within a bunch crossing in LHC, more interactions happen
- An event of interesting physics will be **recorded together with other events overlapped**, that are proton-proton interactions with low physics interest
- they are equivalent to a non-interesting event (**minimum bias**)



- given an average number of interactions, the number of PU events per bunch-crossing is expected to have roughly a poissonian distribution

measure the pile-up

- multiply the luminosity (per bunch) by the minimum bias cross-section (71.3 mb) gets the expected rate per bunch:

$$\text{Rate}_{\text{pileup}_{\text{xing,ls}}} = \mathcal{L}_{\text{xing,ls}} \cdot \sigma_{\text{minimum bias}}$$

- divide by the revolution frequency of a bunch to get the number of PU events:

$$\mathcal{N}_{\text{pileup}_{\text{xing,ls}}} = \frac{\mathcal{L}_{\text{xing,ls}} \cdot \sigma_{\text{minimum bias}}}{\text{circulation rate}}$$

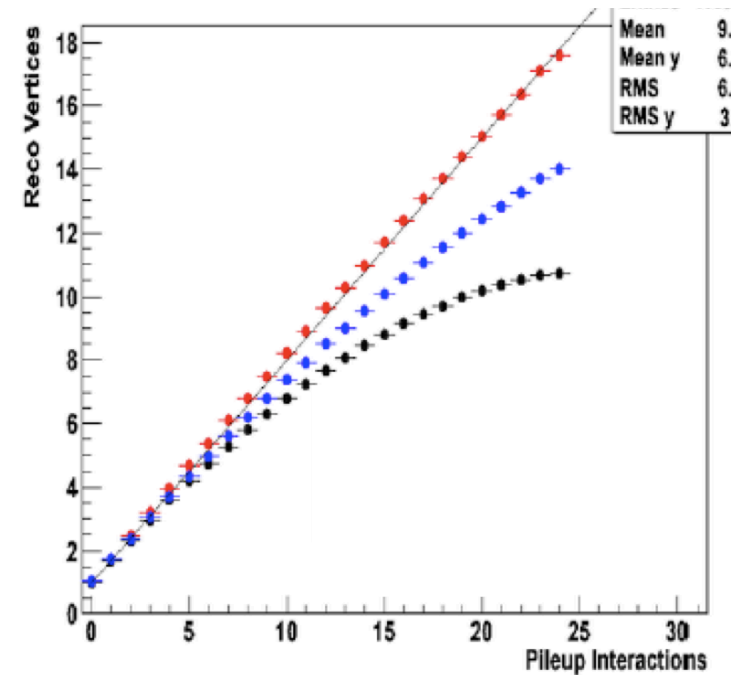
- calculate average distributions over longer periods, weighting by the luminosities

effects of pile-up

- fill in the detector with deposits:
 - **jet reconstruction** algorithms incorporate pile-up deposits
 - **lepton isolation** cones are filled in with pile-up deposits
 - **new jets** might appear in the event
 - more hits in the **tracker** appear
 - the **trigger** is affected
 - **MET** resolution worsens
 -

how to deal with it

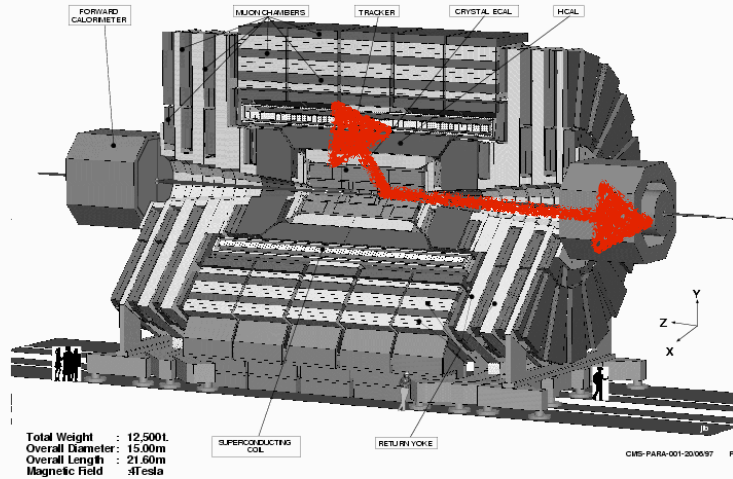
- apply strict **requirements on the vertexing of tracks** - need a precise vertex reconstruction algorithm
- measure the **pile-up density** event by event, and use it to subtract from the jets energy a pile-up term (FastJet)
- do the same with isolation cones
- subtract in the isolation cone the contribution of tracks that do not aim at the same vertex of the lepton
- reconstruct the MET only with particles that aim at a given vertex



dijet cross-section

dijet cross-section

- measure the production of one central + one forward jet in CMS



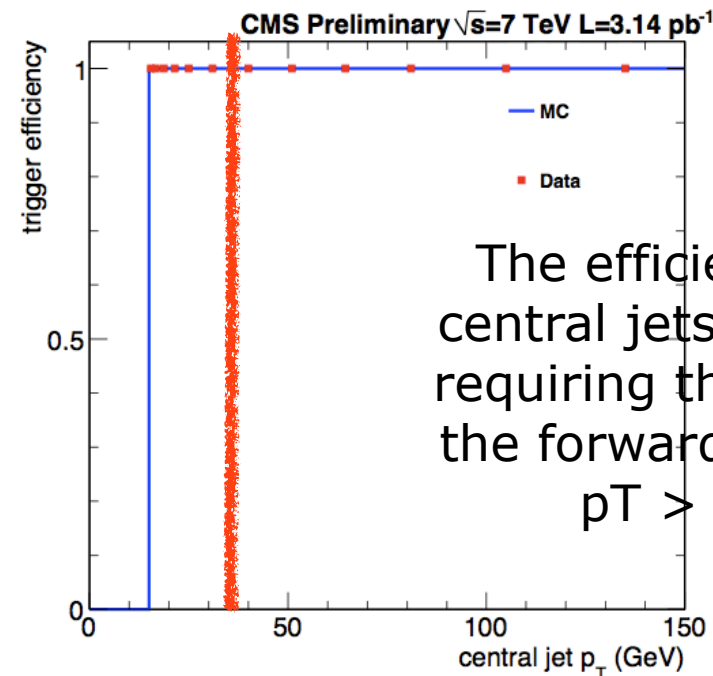
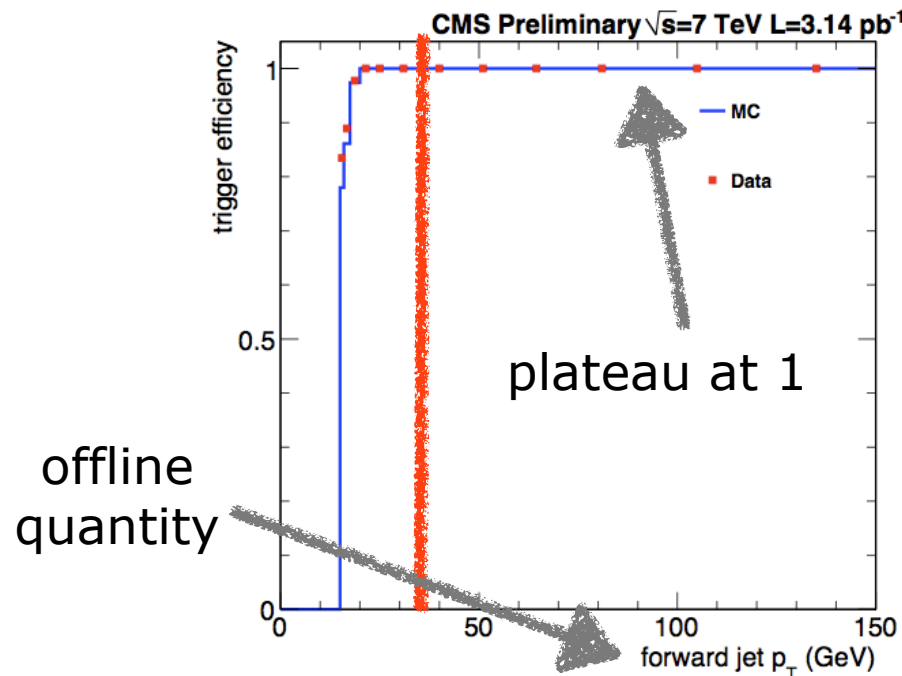
$$\frac{d^4\sigma}{dp_T^c dp_T^f d\eta^c d\eta^f}$$

- to cope with the statistics available, the measurement is done versus p_T only, integrated over the central and forward regions, averaged over η

$$\left\{ \begin{array}{l} \frac{d^2\sigma}{dp_T^f d\eta^f} = \frac{1}{\Delta\eta^f} \cdot \frac{d^4\sigma}{dp_T^c dp_T^f d\eta^c d\eta^f} \Big|_{p_T^c > 35\text{GeV} \wedge |\eta^c| < 2.8} \\ \frac{d^2\sigma}{dp_T^c d\eta^c} = \frac{1}{\Delta\eta^c} \cdot \frac{d^4\sigma}{dp_T^c dp_T^f d\eta^c d\eta^f} \Big|_{p_T^f > 35\text{GeV} \wedge 3.2 < |\eta^f| < 4.7} \end{array} \right.$$

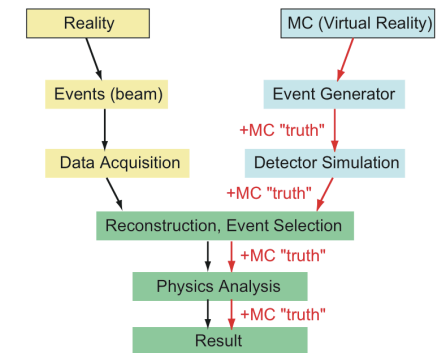
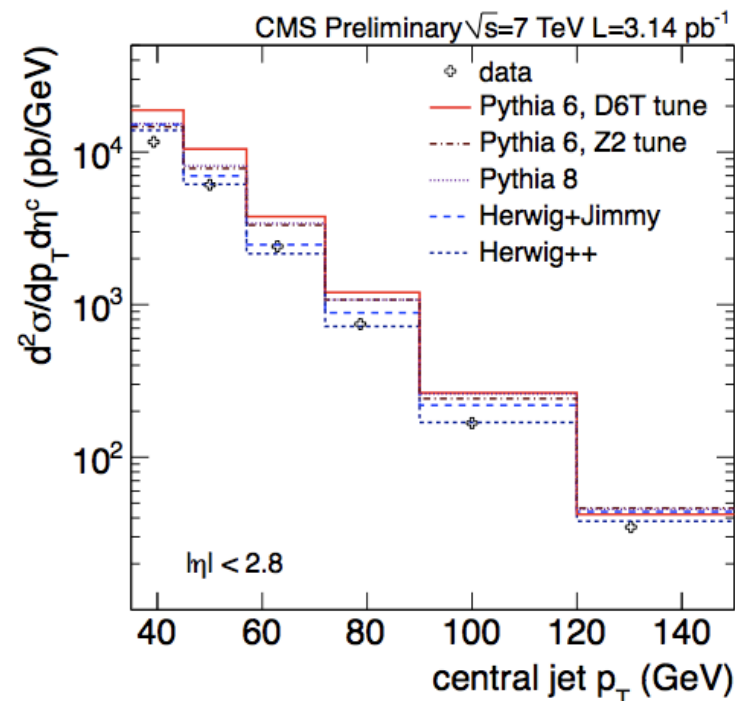
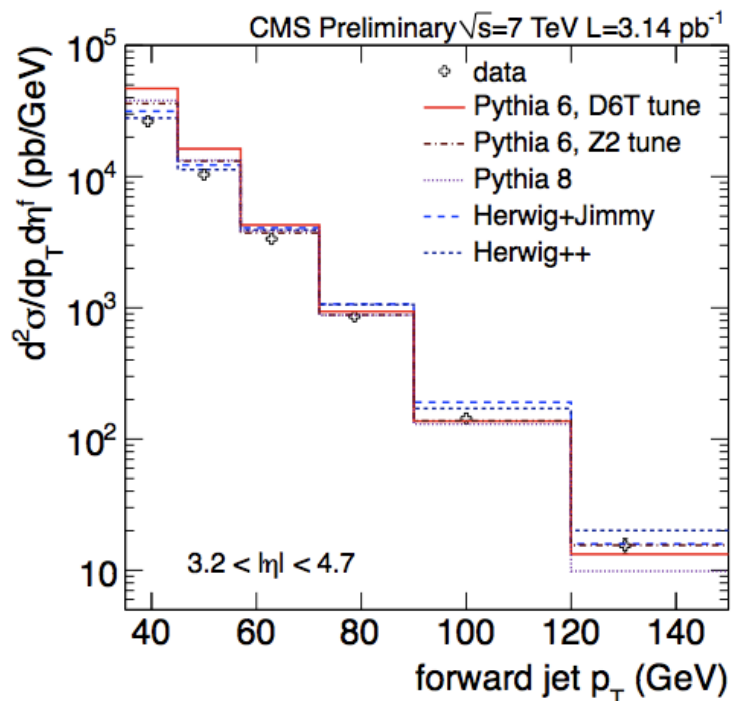
the analysis definition

- simple topology: at least one central and one forward jet with $p_T > 35$ GeV in the event
- the one with the highest p_T is used in each region
- first question: do we trigger these events? di-jet trigger with a raw calorimeter energy threshold of $(E_{T,1} + E_{T,2})/2 > 15$ GeV
- measure the trigger efficiency with the bootstrapping method (wrt the minimum-bias one)

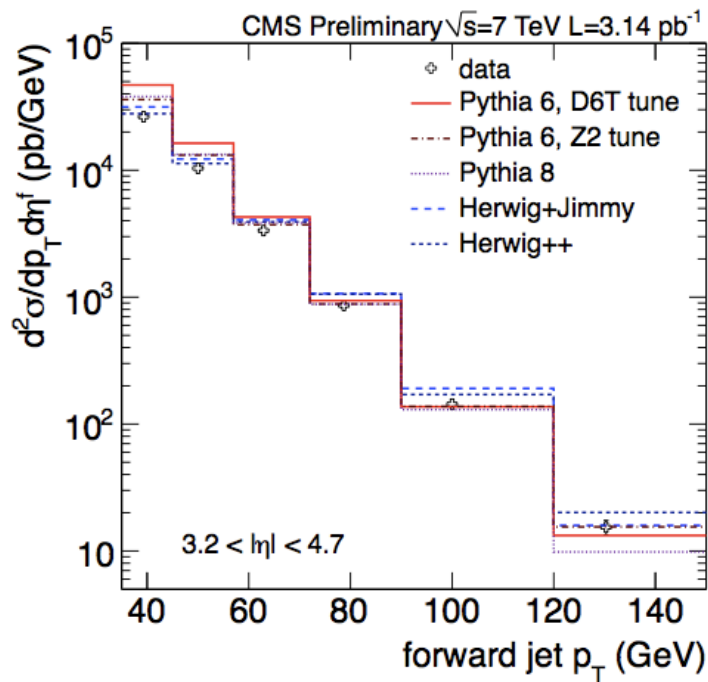


the observed cross sections

- count the number of events in bins of p_T , for the forward and central regions separately (at reconstruction level)
- comparison to some montecarlo predictions is possible, since the simulated events propagated though the whole chain
- what can be done for simulations that do not reach the end of the chain?

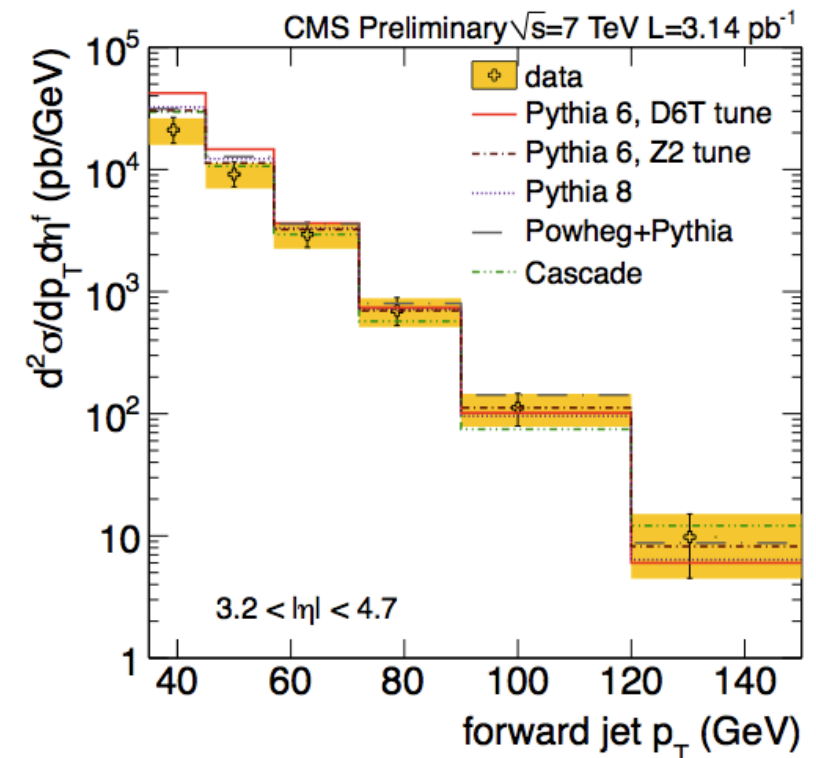


the unfolding



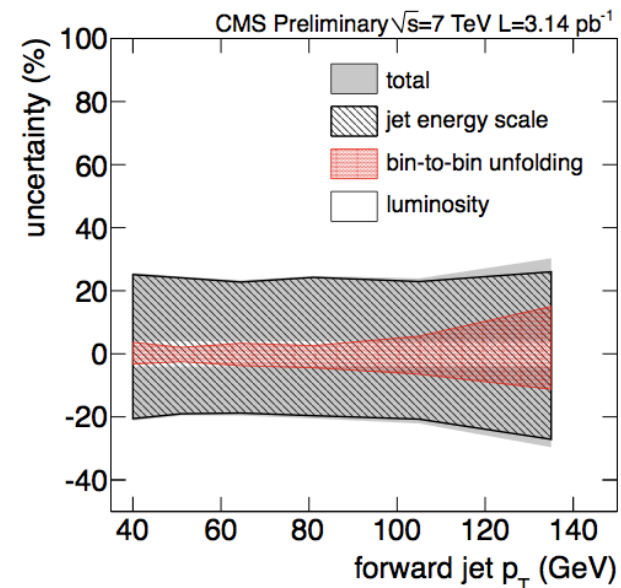
- the rather large jet energy resolution (10%) can give rise to migration effects among the bins
- the interaction with the detector can change the shape of the cross-section
- need to unfold the distribution to the hadron level

- pick one simulation
- reweight the simulated events on quantities at hadron level, to let it match the data
- use the ratio between hadron and detector level to get correction factors to be applied to the data



systematics

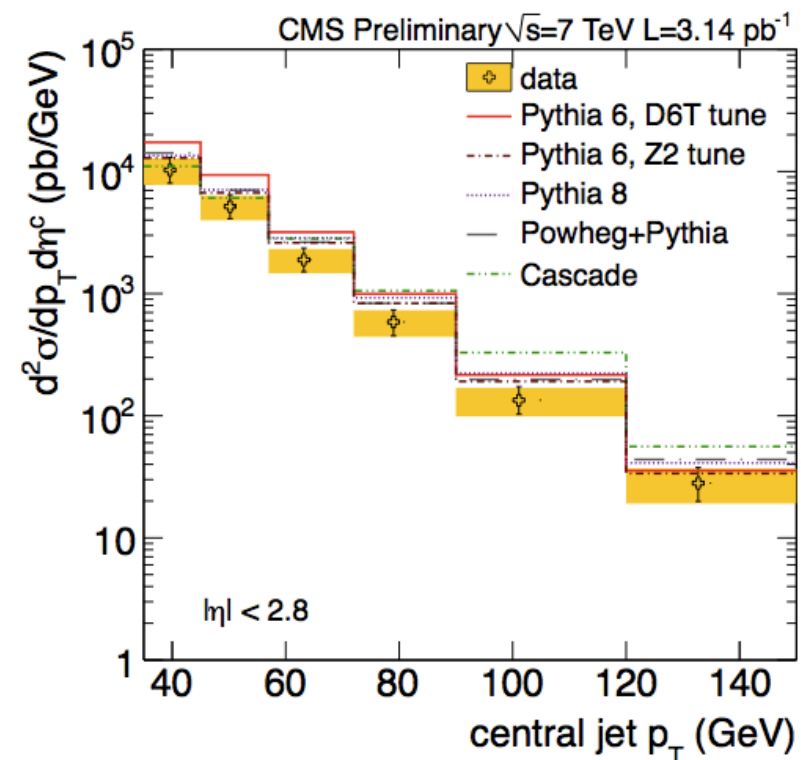
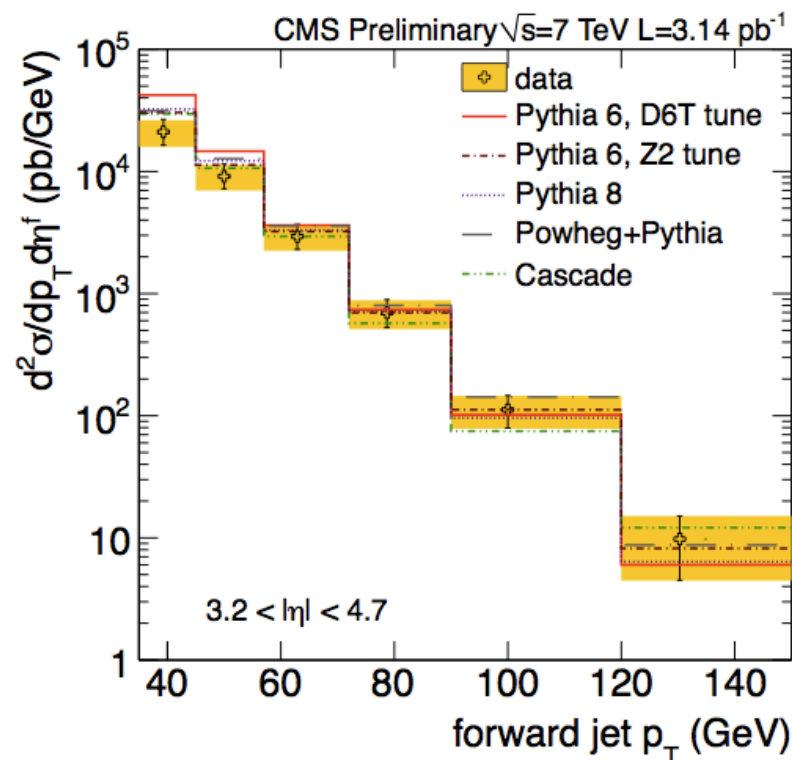
source	uncertainty [%]
JES	25
jet energy resolution	<5
PU	≈ 5
luminosity	4
corrections	<10
HLT	≈ 0
total	30



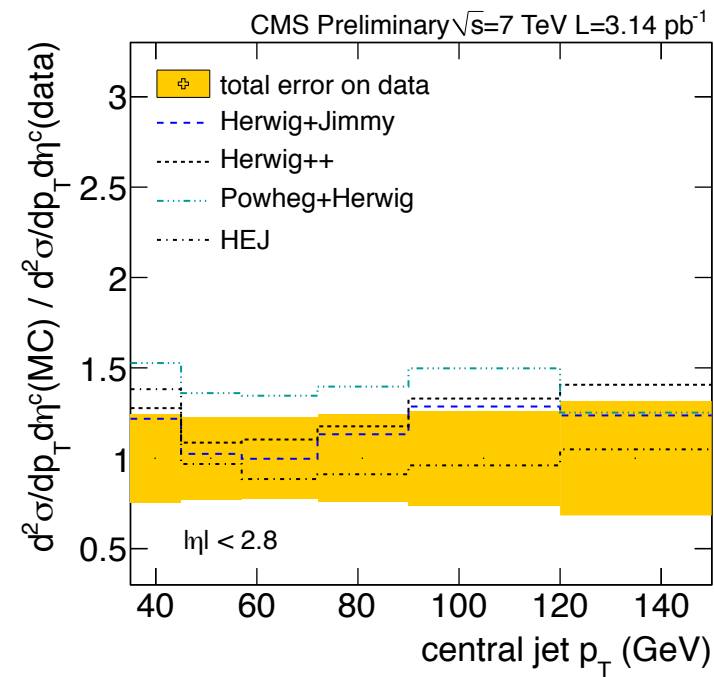
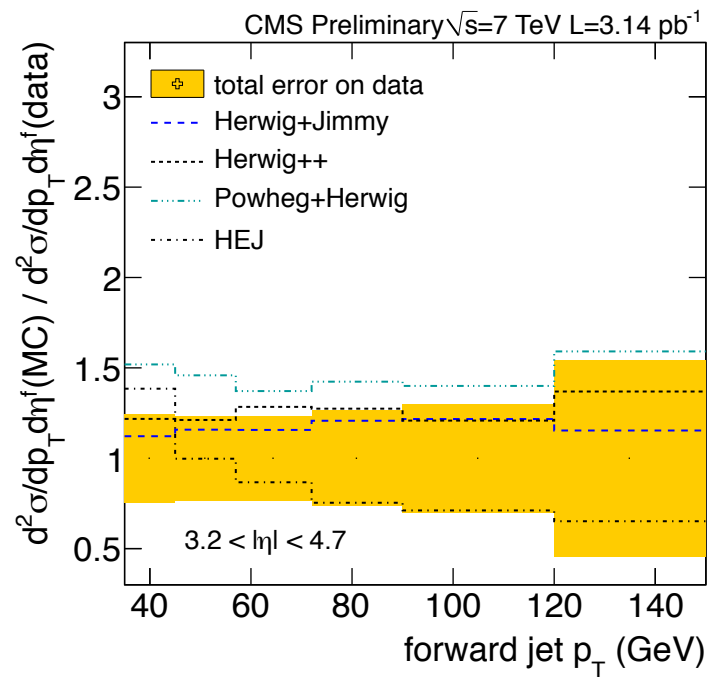
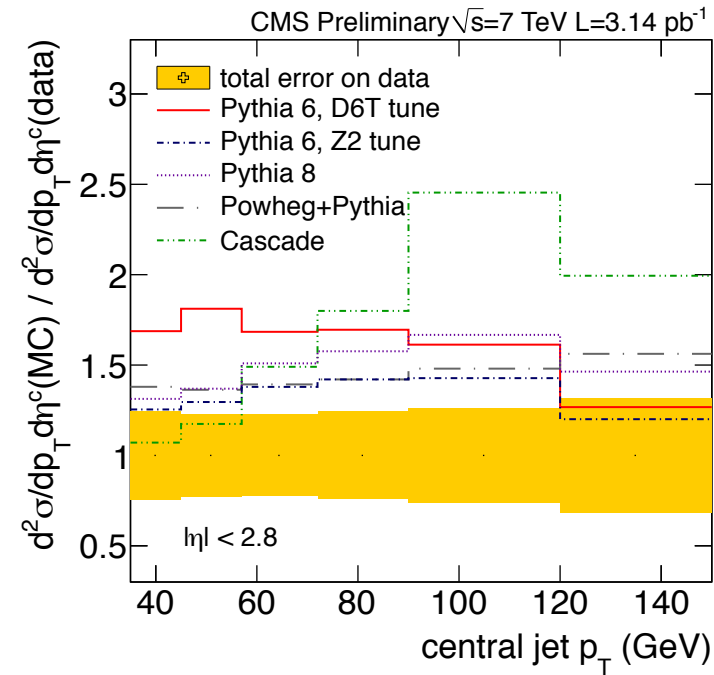
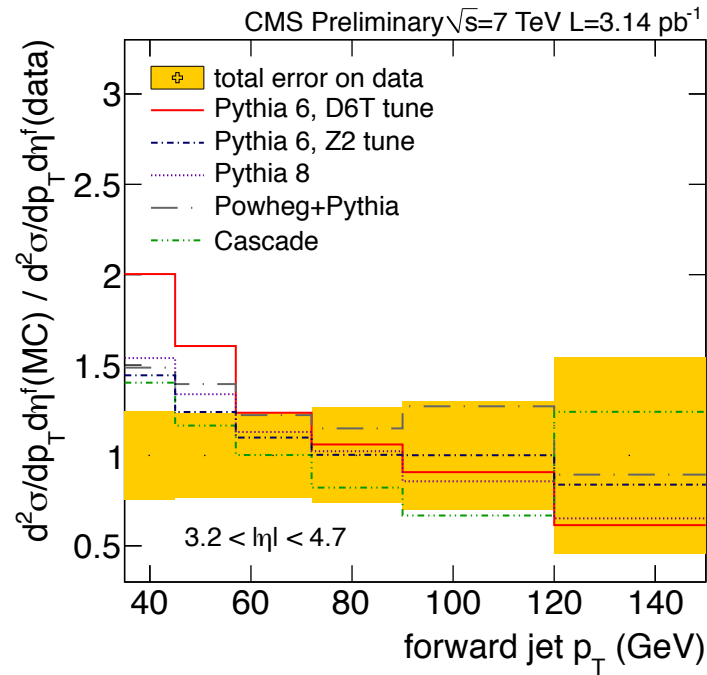
- propagating the initial uncertainties through the analysis
 - **jet energy scale (JES)**: coherently vary all the jets p_T of $\pm\sigma$ in the analysis and compare the results
 - **jet energy resolution**: assume a better (worse) resolution and propagate the effect
- comparing the effects of different initial choices (PU, corrections)
 - **PU**: perform the analysis with all the events, or the ones with a single vertex
 - **unfolding**: calculate the factors with several simulations and combine the results

the cross-section

- such final state can give informations on multi-parton interaction and multi-jet production
- study different types of parton radiation dynamics (DGLAP, BKKL, CCFM)
- compare the results to simulations that implement the different behaviours



the result



H > WW > lvlv

one plot for the Higgs boson

